

MILAN CHRISTMAS MEETING 2024



STUDY OF NON-PERTURBATIVE POWER CORRECTIONS TO EVENT SHAPES USING PanScales SHOWERS

SILVIA ZANOLI
University of Oxford

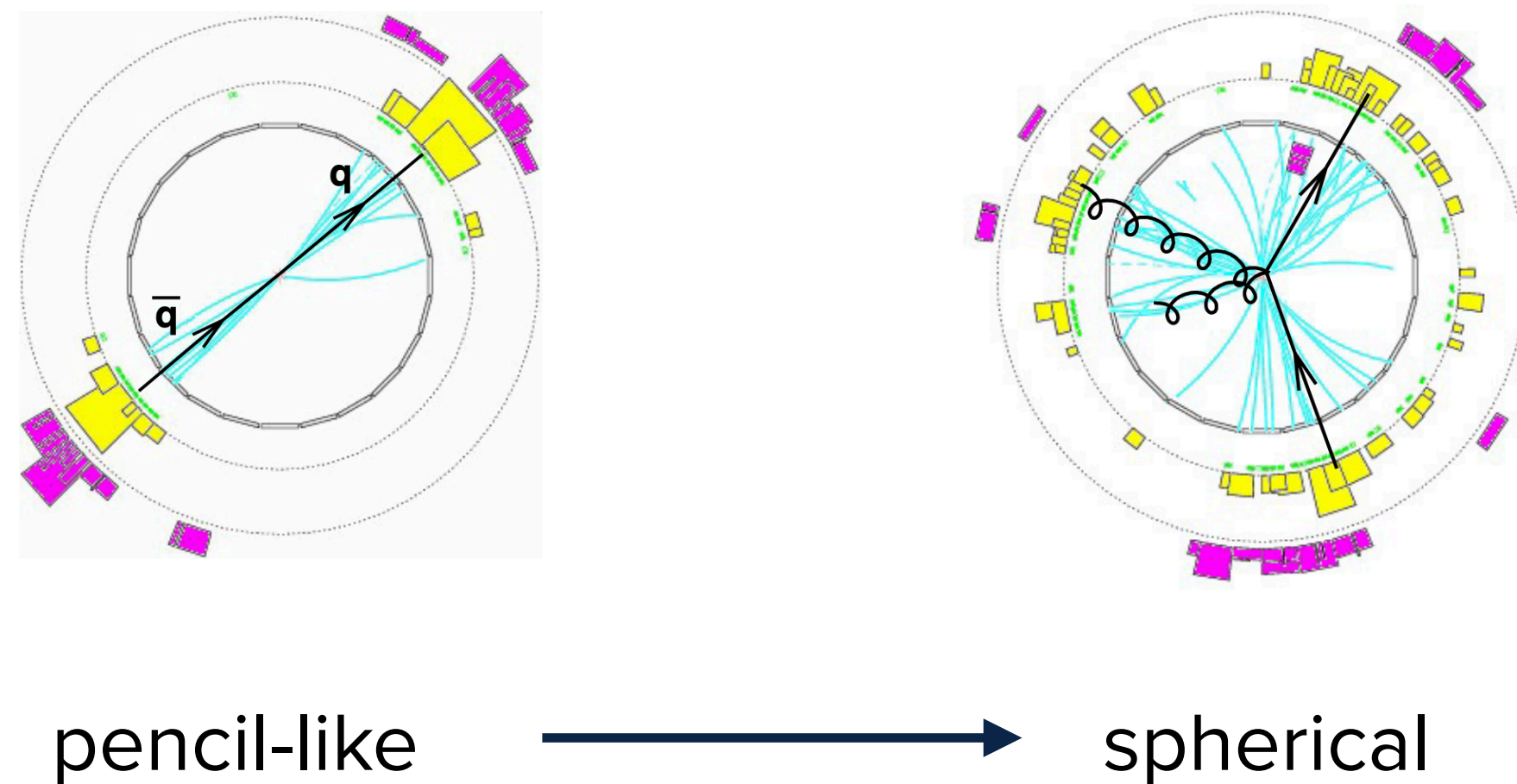


UNIVERSITY OF
OXFORD

Ongoing work with F. Caola, C. Farren-Colloty, J. Helliwell, R. Patel, G.P. Salam

EVENT SHAPES

Used mostly in the context of e^+e^- collisions, they provide information on the geometry of an event.



NICE PROPERTIES

- Very clean environment for experimental measurements.
- IR safe, so they can be computed in perturbative QCD.
- Known at high orders in QCD (including resummation).

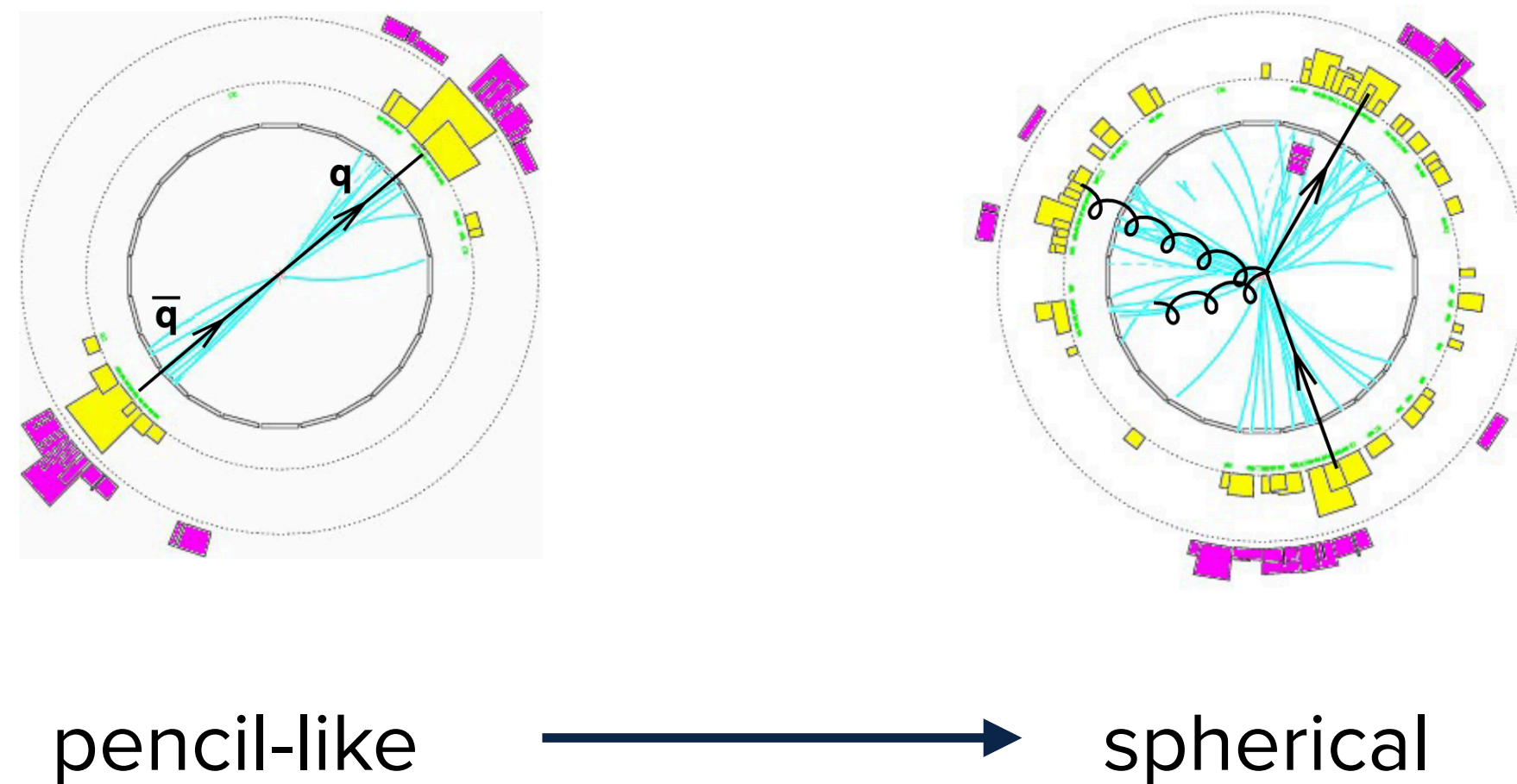


USED TO TEST QCD AND ITS DYNAMICS

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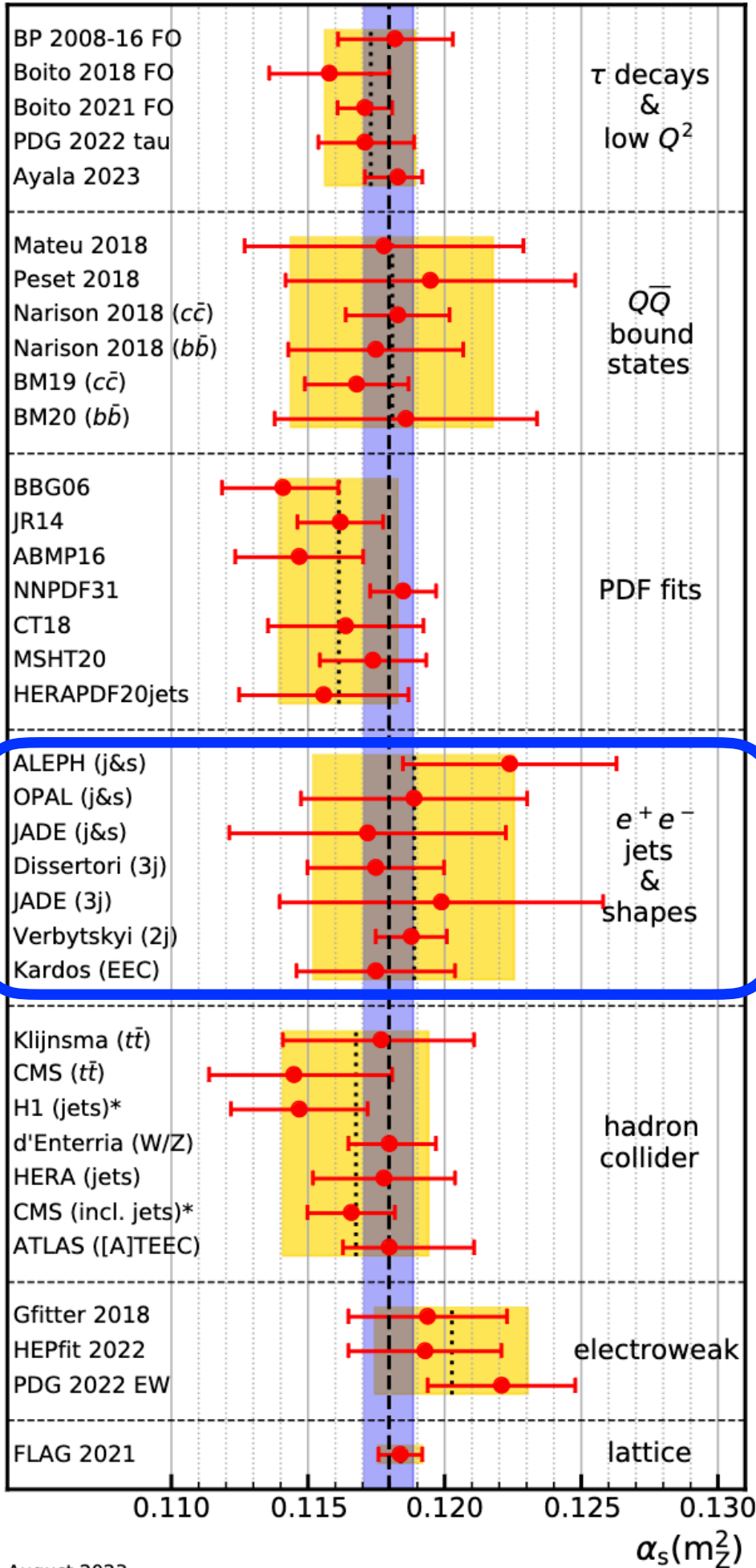


AFFECTED BY SIGNIFICANT NON-PERTURBATIVE CORRECTIONS (Λ/Q)

DETERMINATION OF α_s

Hadronic final state
of e^+e^- collisions:

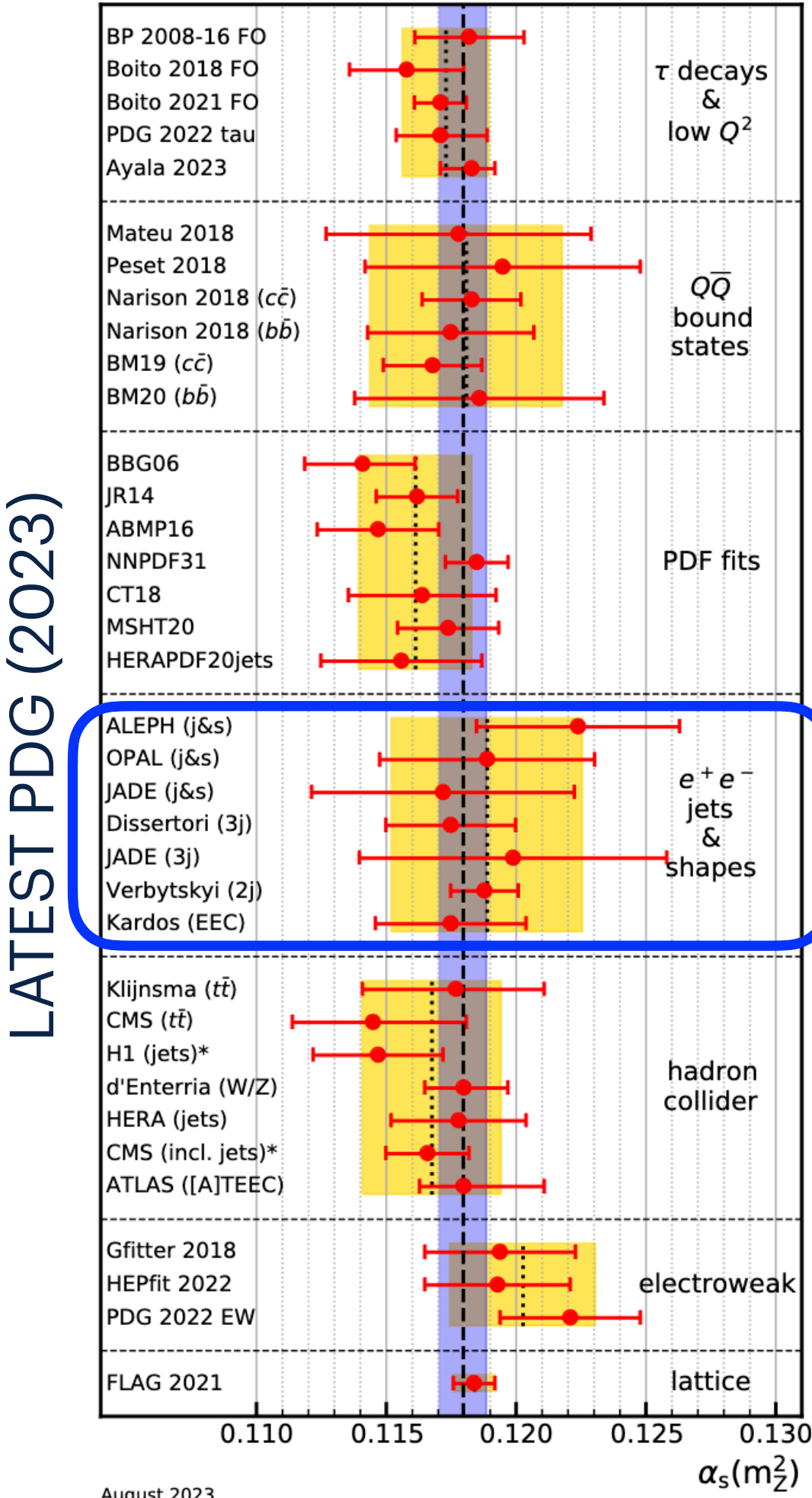
LATEST PDG (2023)



August 2023

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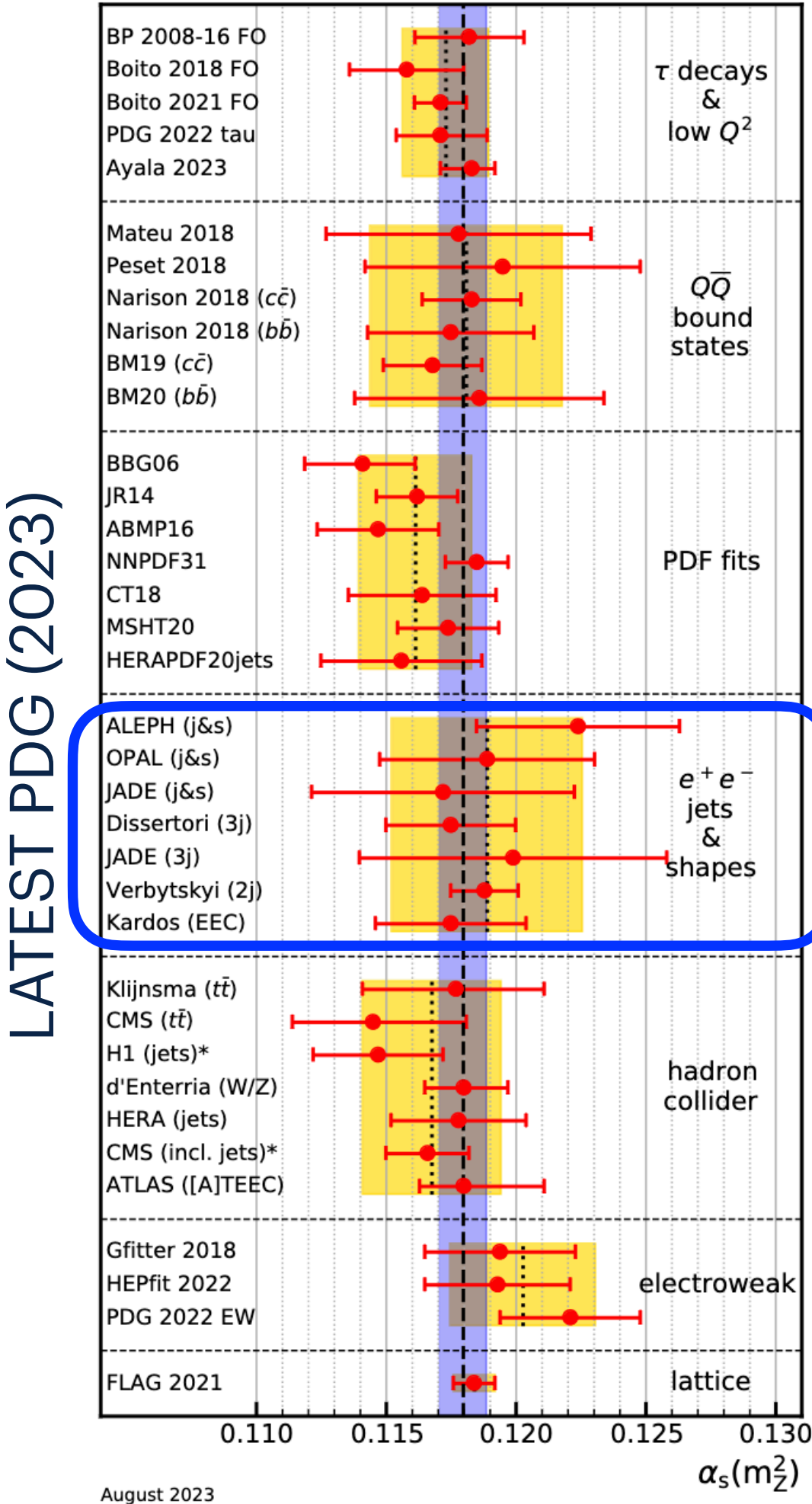
NON-PERTURBATIVE (=HADRONIZATION) CORRECTIONS CAN BE OBTAINED IN TWO DIFFERENT WAYS:

1 - FROM A MC GENERATOR

- ✓ Very practical: construct a migration matrix describing the parton to hadron transition, and apply it in the data/theory comparison.
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2 - FROM ANALYTIC MODELS

Fit of thrust data: $\alpha_s(m_Z^2) = 0.1135 \pm 0.0011$ [1]

$\alpha_s(m_Z^2) = 0.1134^{+0.0031}_{-0.0025}$ [2]

Fit of C parameter data: $\alpha_s(m_Z^2) = 0.1123 \pm 0.0015$ [3]

World average: $\alpha_s(m_Z^2) = 0.1180 \pm 0.0008$



[1: Abbate et al. 1006.3080] [2: Gehrmann, Luisoni, Monni 1210.6945][3: Hoang et al. 1501.04111]

TIMELINE OF ANALYTIC MODELS

90s

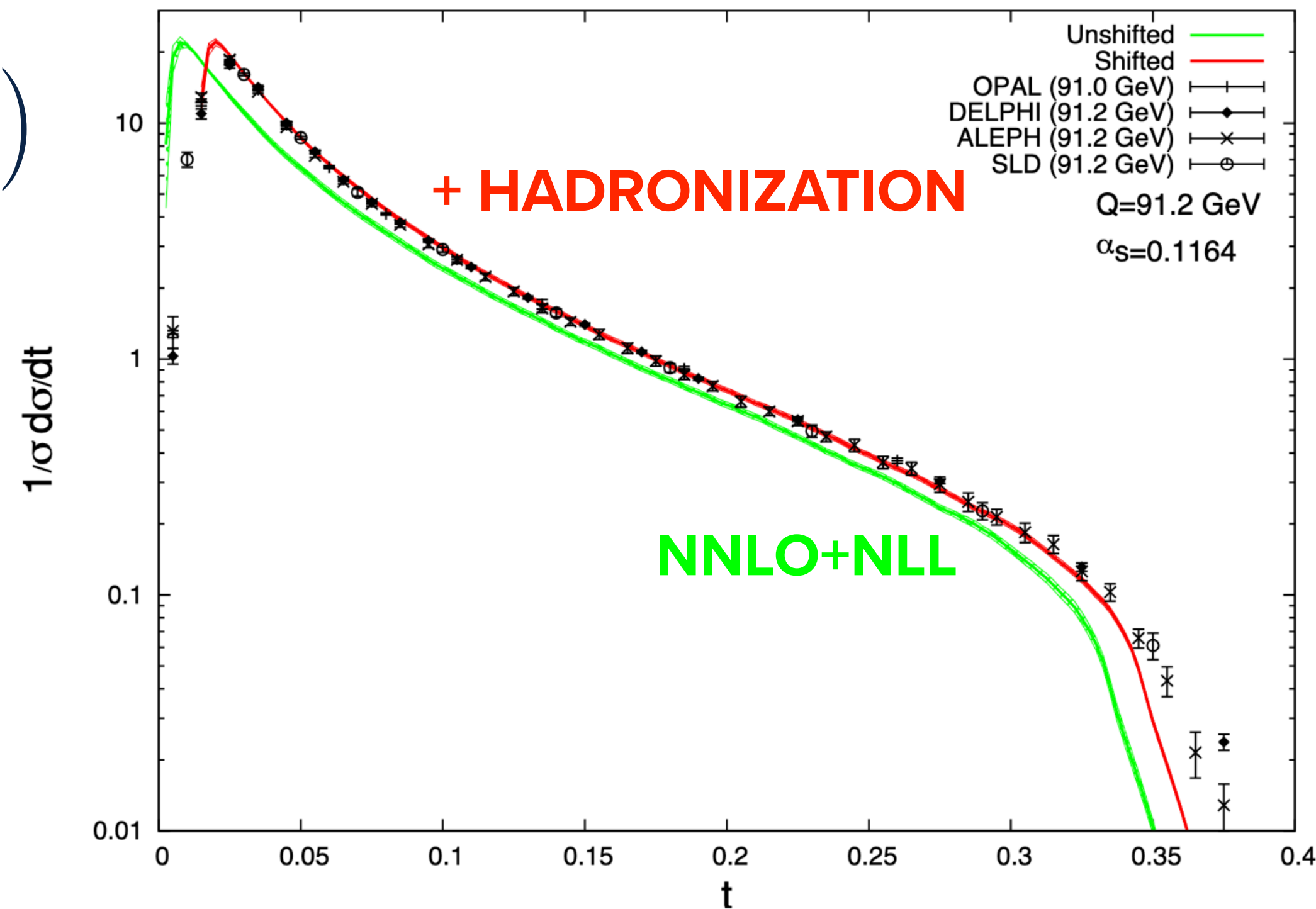
Analytic models are based on a 2-jet limit analysis, assuming that non perturbative effects are constant in the entire kinematic spectrum. This leads to α_s much smaller than the world average.

2021

2022

2023

$$t = 1 - T = \max_{\vec{n}_T} \left(\frac{\sum_i |\vec{p}_i \cdot \vec{n}_T|}{\sum_i |\vec{p}_i|} \right)$$



$$\left. \frac{1}{\sigma} \frac{d\sigma}{dt} \right|_t = \left(\left. \frac{1}{\sigma} \frac{d\sigma}{dt} \right|_{t+\delta t} \right)^{\text{pert}}$$

$\delta t = \text{constant } 1/Q \text{ shift}$

[Davison, Webber 0809.3326]

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The previous assumption is proven to be wrong calculating the leading non-perturbative correction to the C parameter in C=0 and C=3/4.

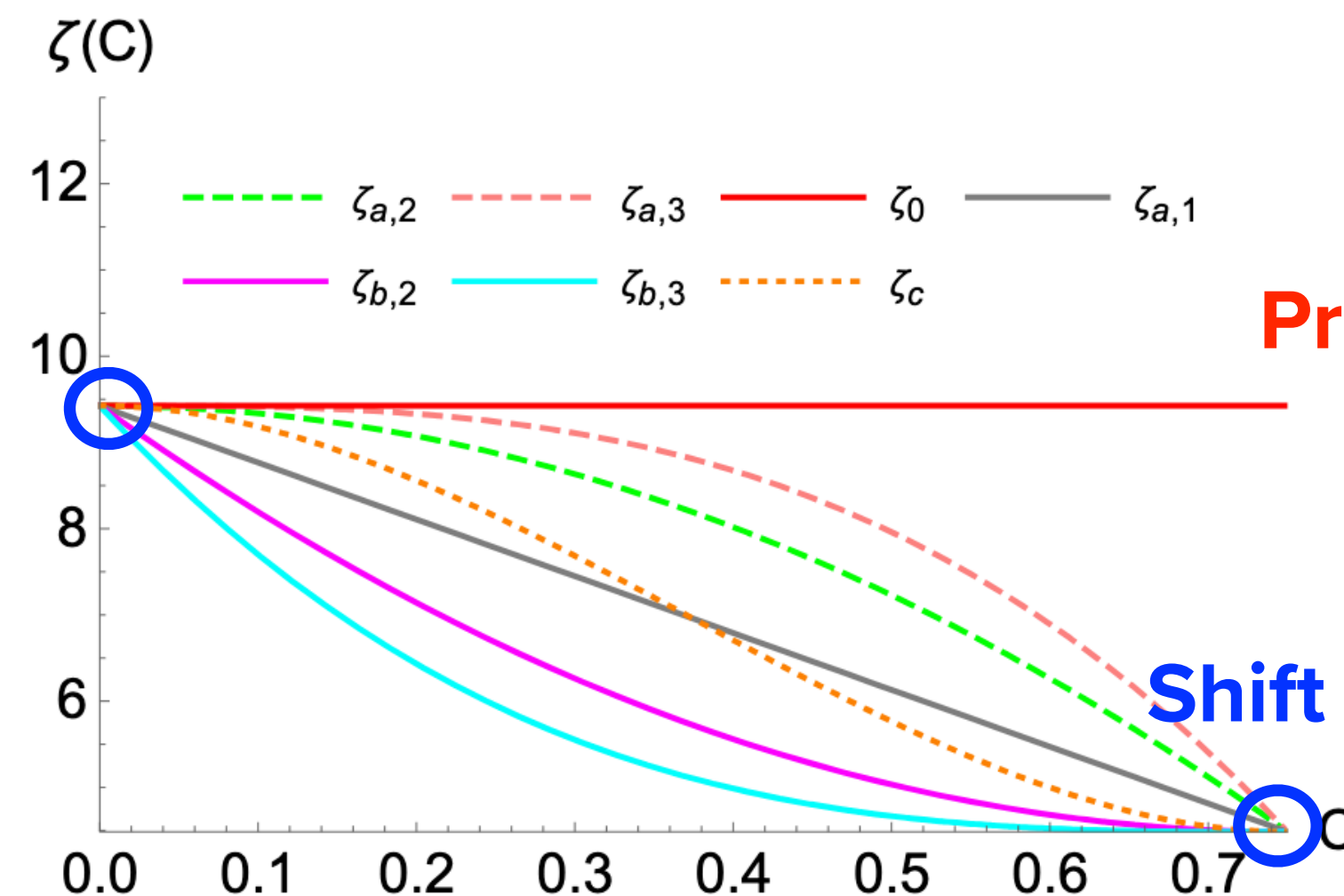
2022

2023

$$\Theta^{\alpha\beta} \frac{1}{\sum_i |p_i|} \sum_i \frac{p_i^\alpha p_i^\beta}{|\vec{p}_i|},$$

$$C = 3 \cdot (\lambda_1 \lambda_2 + \lambda_1 \lambda_3 + \lambda_2 \lambda_3)$$

Shift in the 2-jet limit



Previous assumption of constant shift

Shift in the symmetric 3-jet limit

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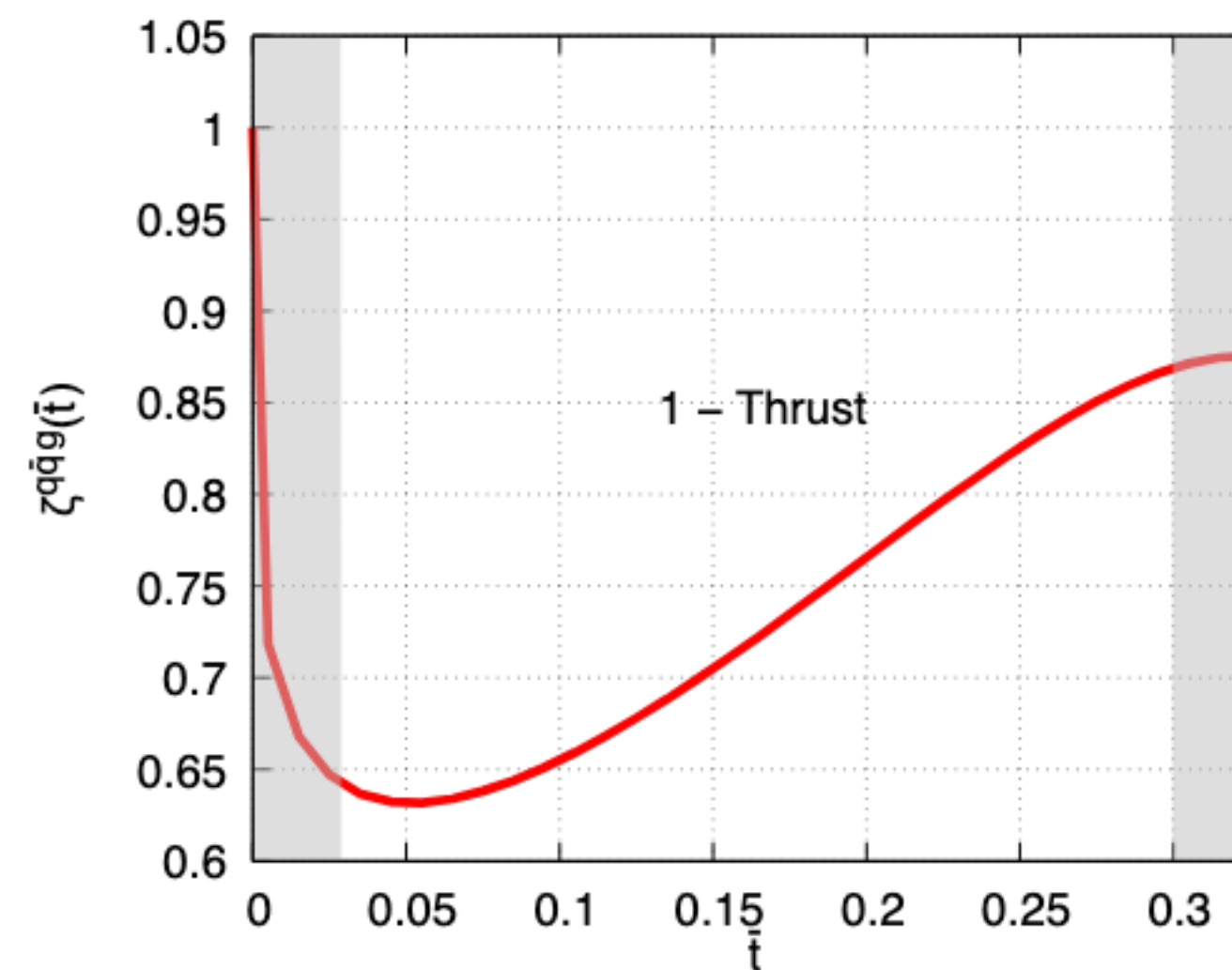
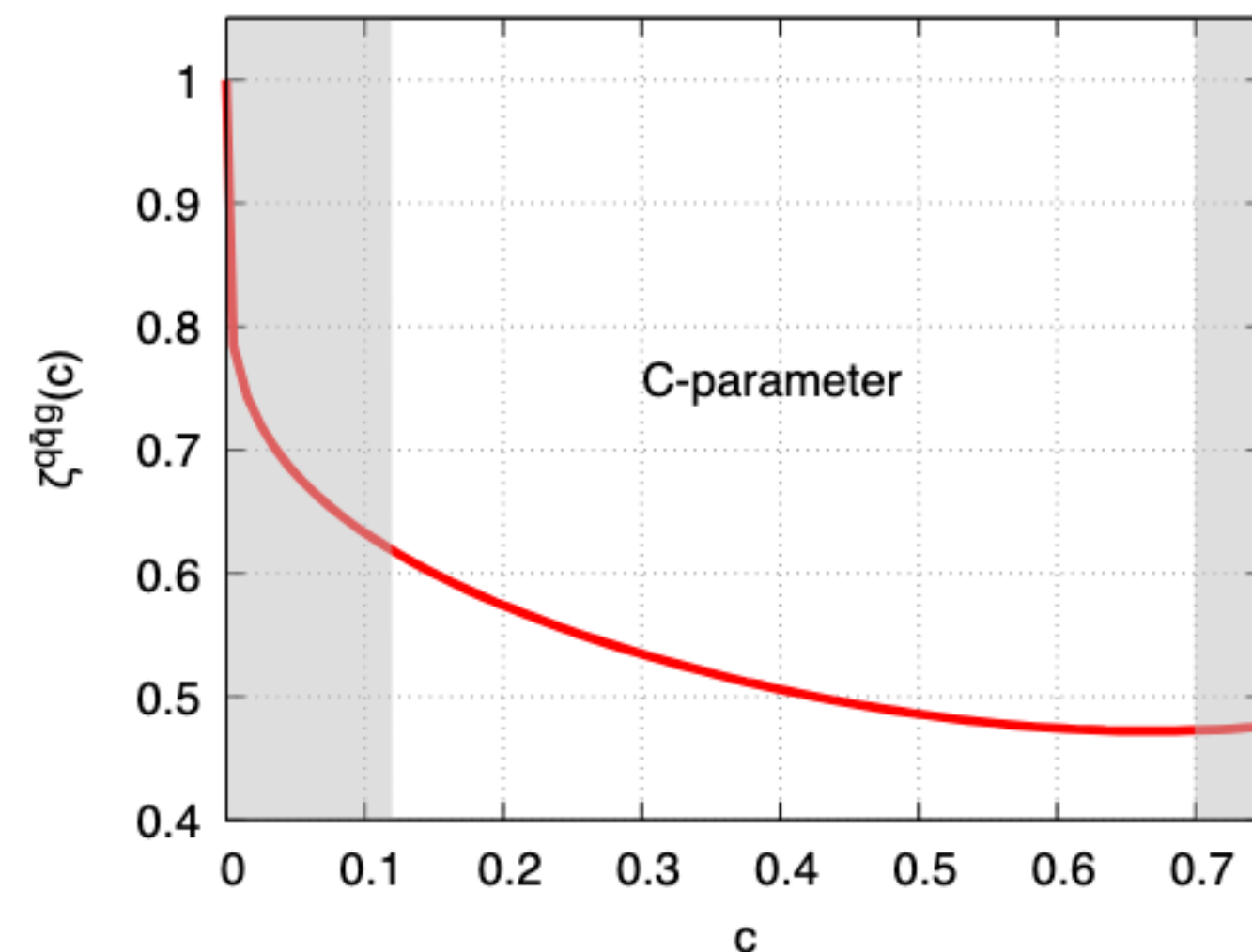
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Full calculation in the 3-jet limit for the thrust and the C parameter.

2023



NON TRIVIAL SHAPE!

[Caola, Ferrario Ravasio, Limatola, Melnikov, Nason 2108.08897; +Ozcelik 2204.02247]

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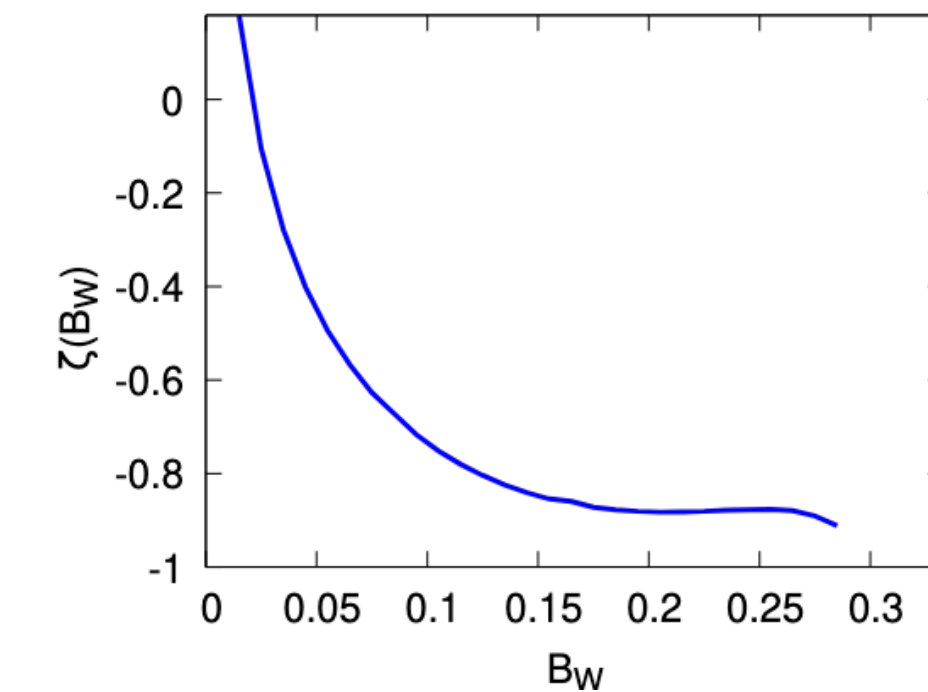
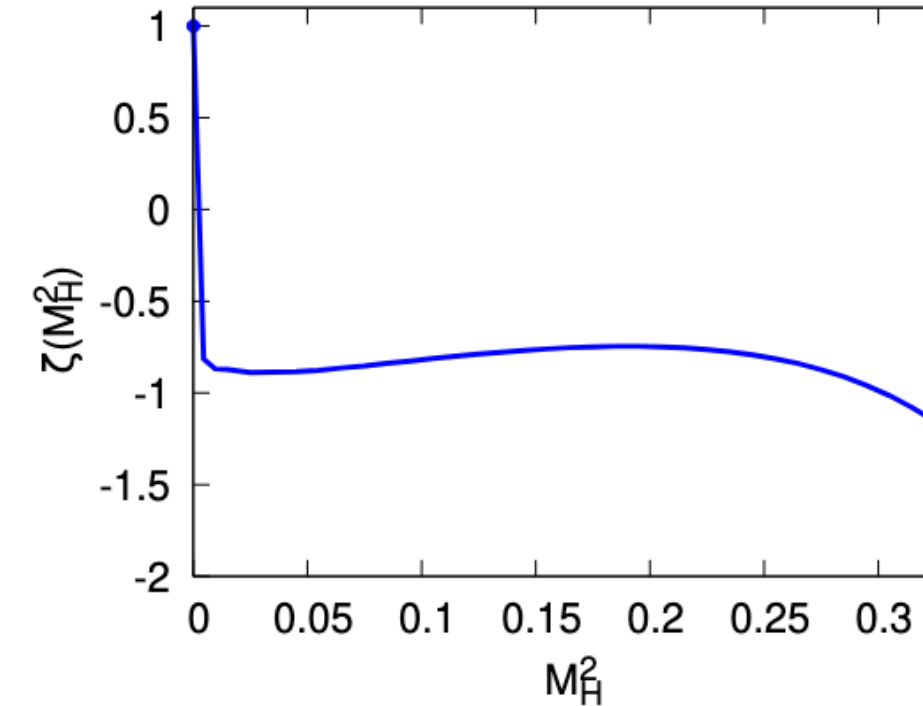
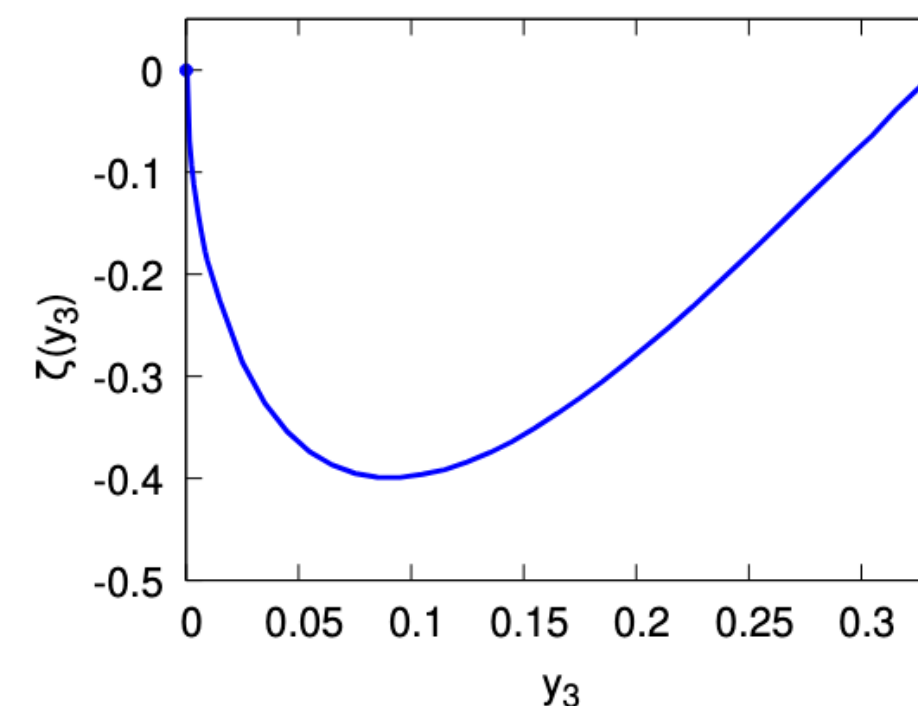
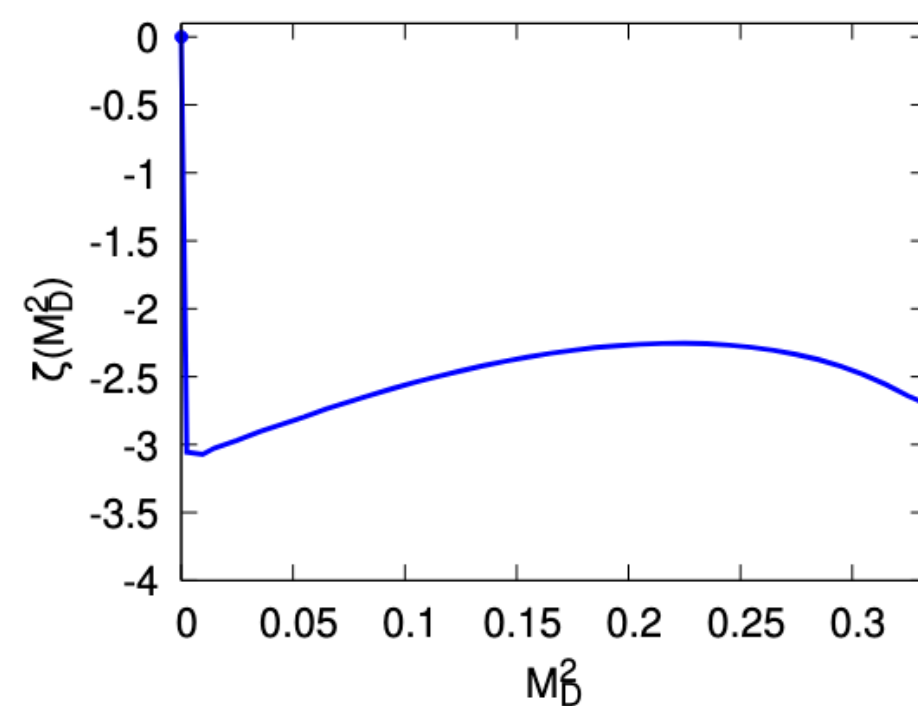
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Extension of the calculation to other event shapes, confirming the previous results.



[Nason, Zanderighi 2301.03607]

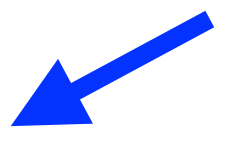
THE DOKSHITZER-WEBBER MODEL

[Dokshitzer, Webber hep-ph/9504219]

HADRONIZATION \equiv emission of soft $k_T \sim \Lambda$ non-perturbative gluon (= “gluer”)

The divergent behaviour of the running coupling at low scales is cured by an effective coupling that is finite:

Intrinsic ambiguity of pQCD
“renormalons picture”


$$\int_0^Q dk \alpha_s(k) = \int_0^{\mu_I} dk \alpha_s(k) + \int_{\mu_I}^Q dk \alpha_s(k) \longrightarrow \mu_I \bar{\alpha}_0(\mu_I) + \int_{\mu_I}^Q dk \alpha_s(k)$$

Matching scale
 $\mathcal{O}(\text{GeV})$

IR finite and universal
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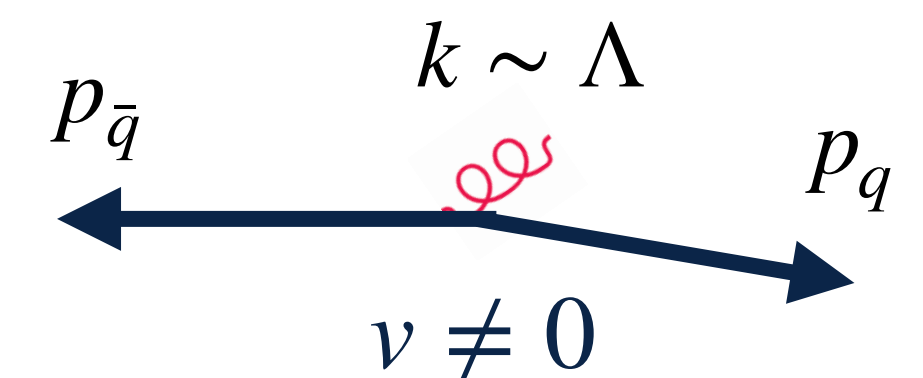
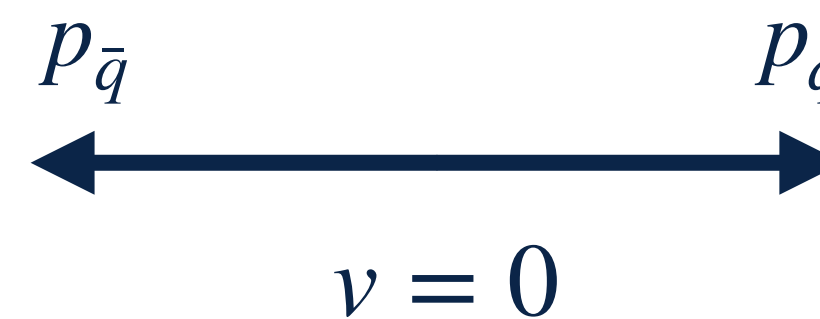
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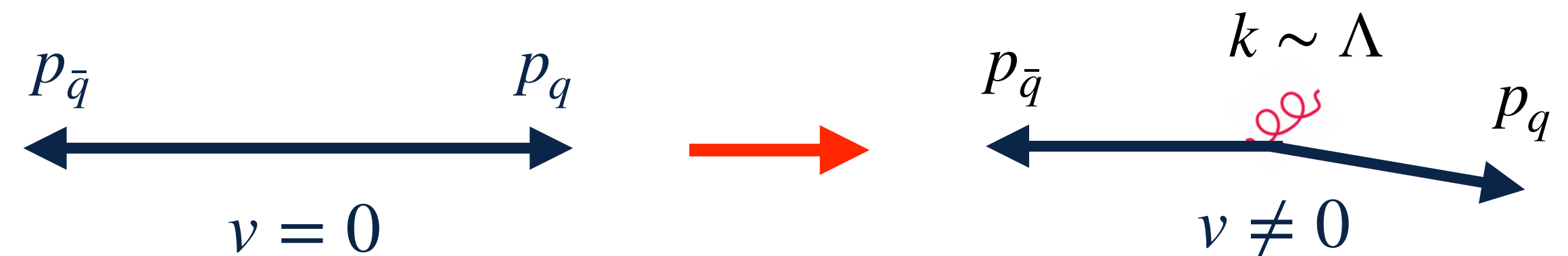
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2 Calculate the shift in the event shape:

$$\delta v = v(p_q, p_{\bar{q}}, k) - v(p_q, p_{\bar{q}}) = v(p_q, p_{\bar{q}}, k)$$

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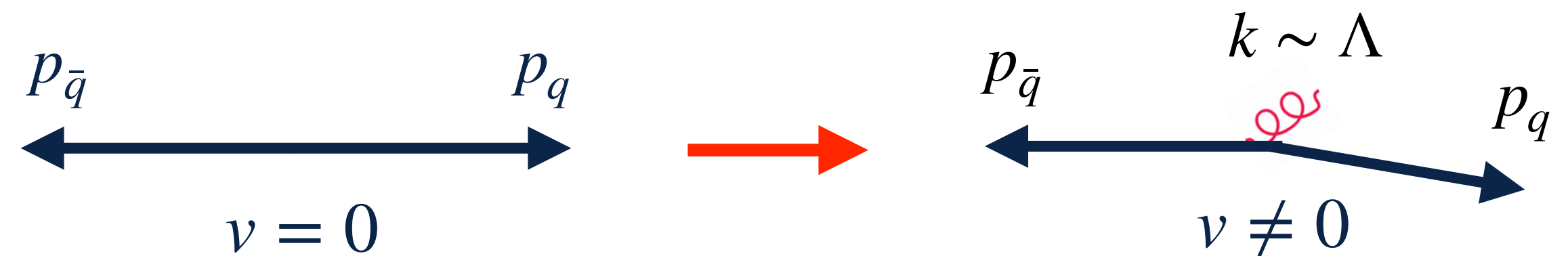
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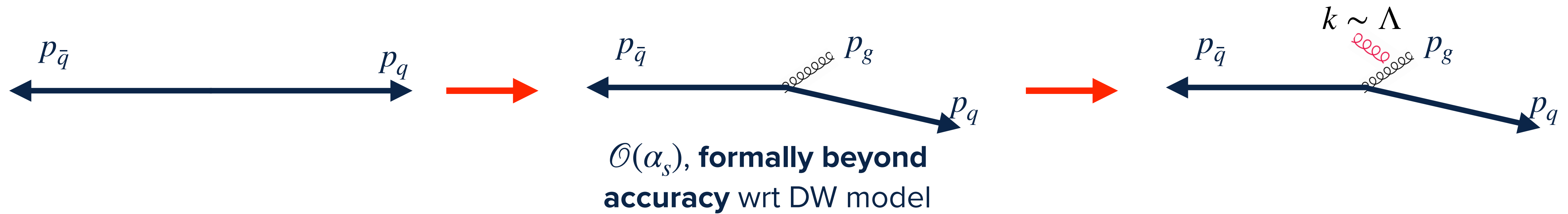
2 Calculate the shift in the event shape: $\delta v = v(p_q, p_{\bar{q}}, k) - v(p_q, p_{\bar{q}}) = v(p_q, p_{\bar{q}}, k)$

3 Average this shift over the gluer emission probability: $\langle \delta v \rangle^{NP} = \frac{1}{\sigma} \int d\Gamma \mathcal{M}^2 \delta v$

PERTURBATIVE EVOLUTION IN Q

[Dasgupta, Hounat 2411.16867]

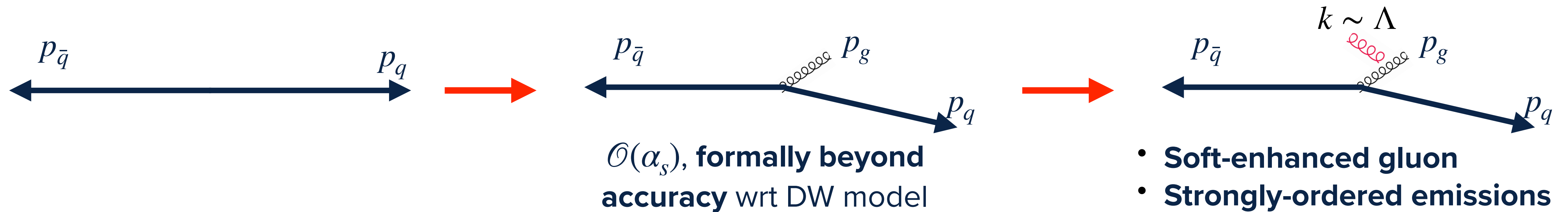
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Interplay between the soft emission and the gluon leading to **large logarithmic corrections**:

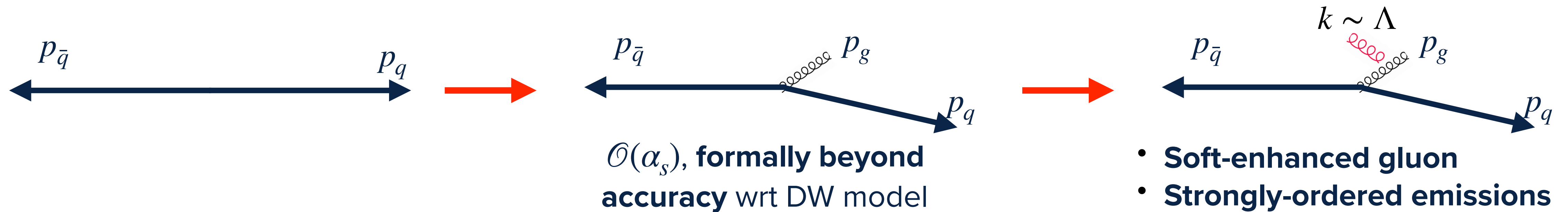
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ANOMALOUS DIMENSION

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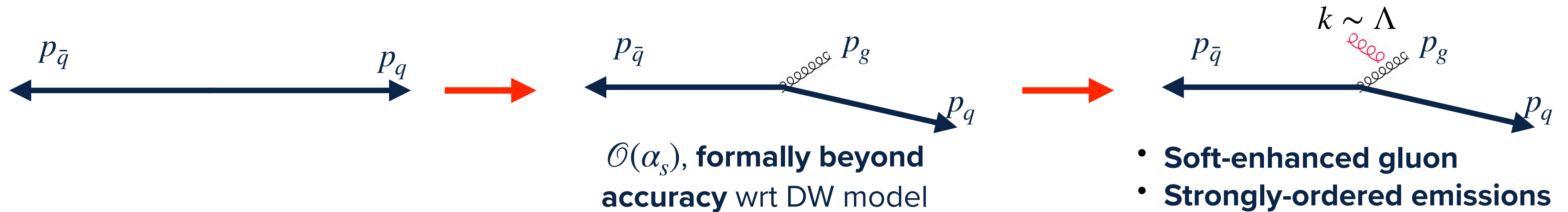
$$\langle \delta\tau \rangle^{\text{NP},1} = -C_F C_A \frac{\alpha_s}{2\pi} \frac{1}{Q} \int_0^{\mu_I} d\kappa_T \frac{\alpha_s(\kappa_T)}{2\pi} \ln \frac{Q}{\kappa_T} \times 19.64048824,$$

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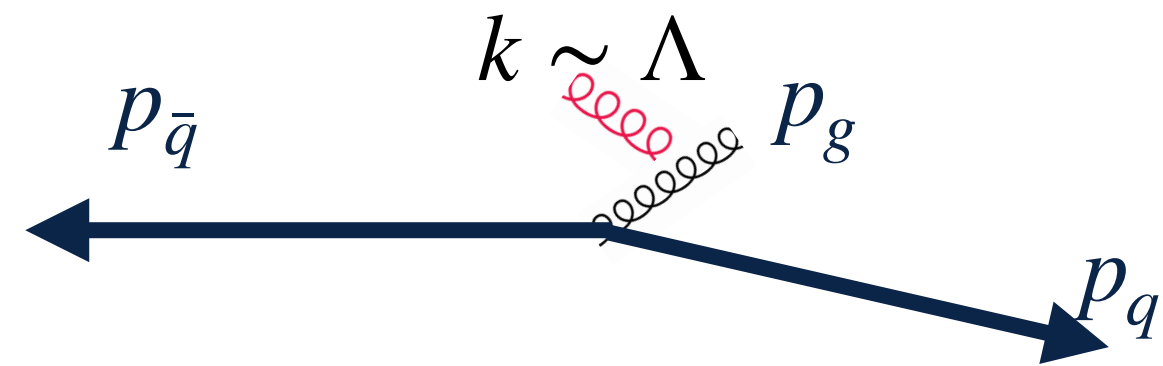
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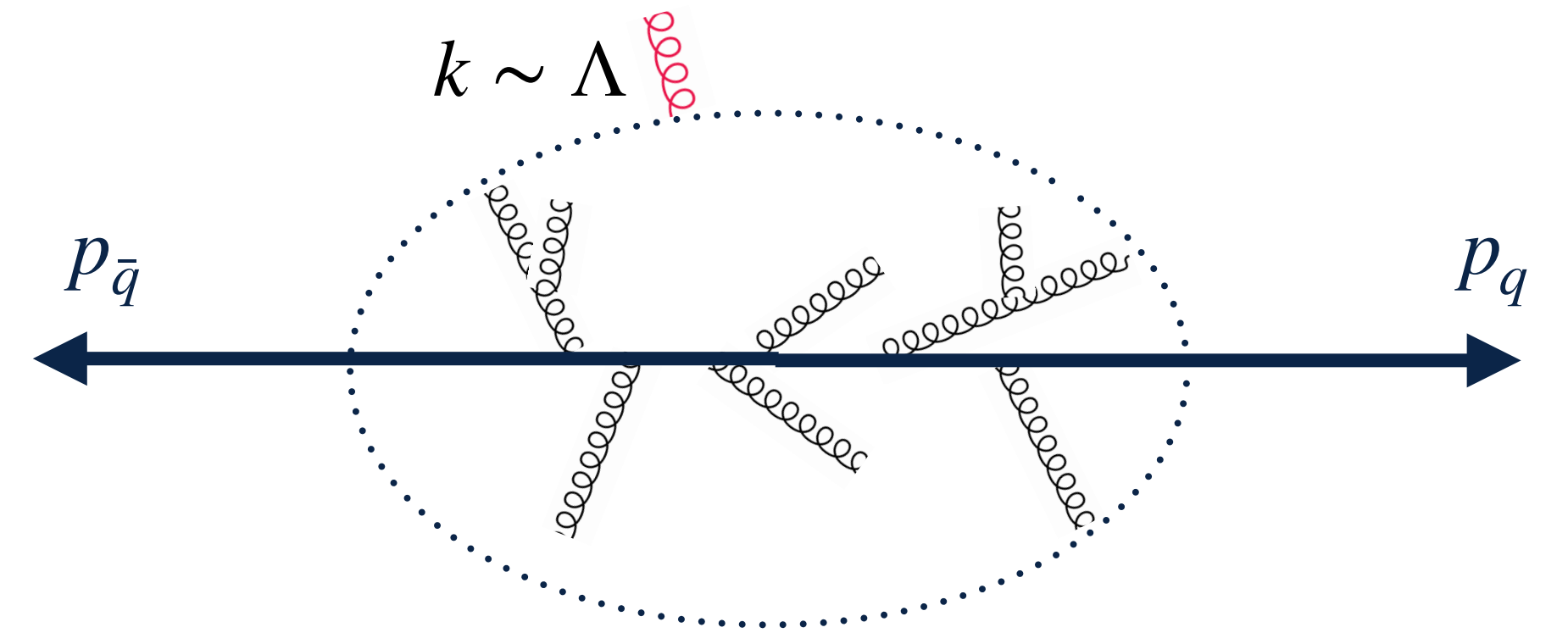
WHAT IF WE ADD MORE EMISSIONS?

RESUMMATION OF THE ANOMALOUS DIMENSION

$$Q \gg k_{T,1} \gg k_{T,2} \gg \dots \gg \Lambda$$



$$\frac{\Lambda}{Q} \cdot \alpha_s \ln \frac{Q}{\Lambda}$$

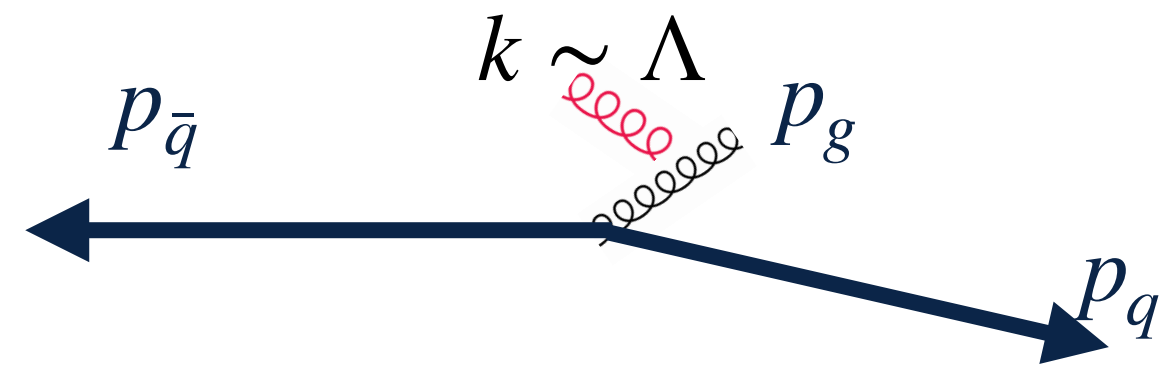


$$\frac{\Lambda}{Q} \cdot \left(\alpha_s \ln \frac{Q}{\Lambda} \right)^n$$

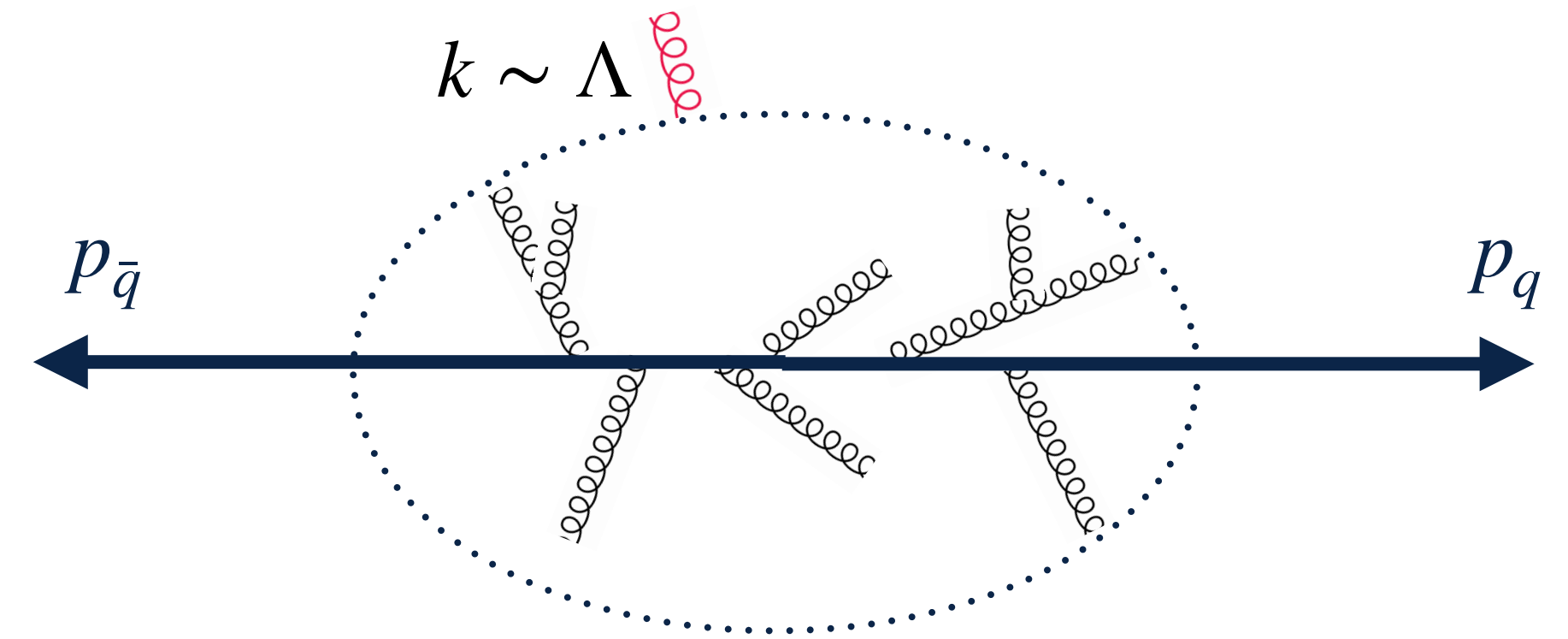
NON-GLOBAL RESUMMATION

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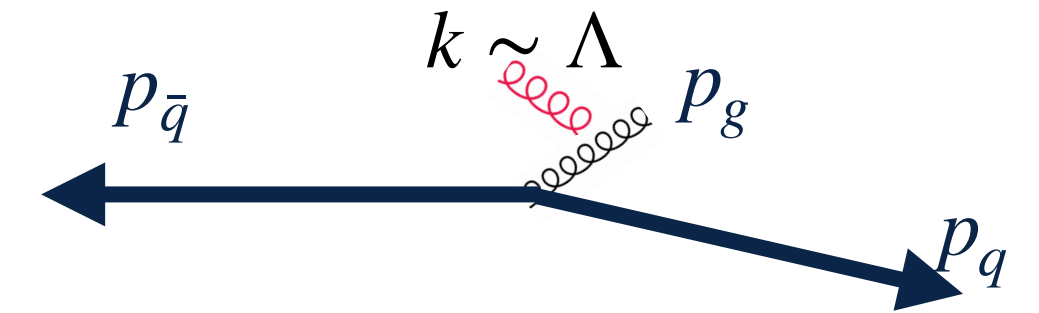
OUR PLAN: resum this contribution to all orders using PanScales showers.

Note: PanScales mappings are consistent with the smoothness criteria pointed out in [2108.08897; 2204.02247]

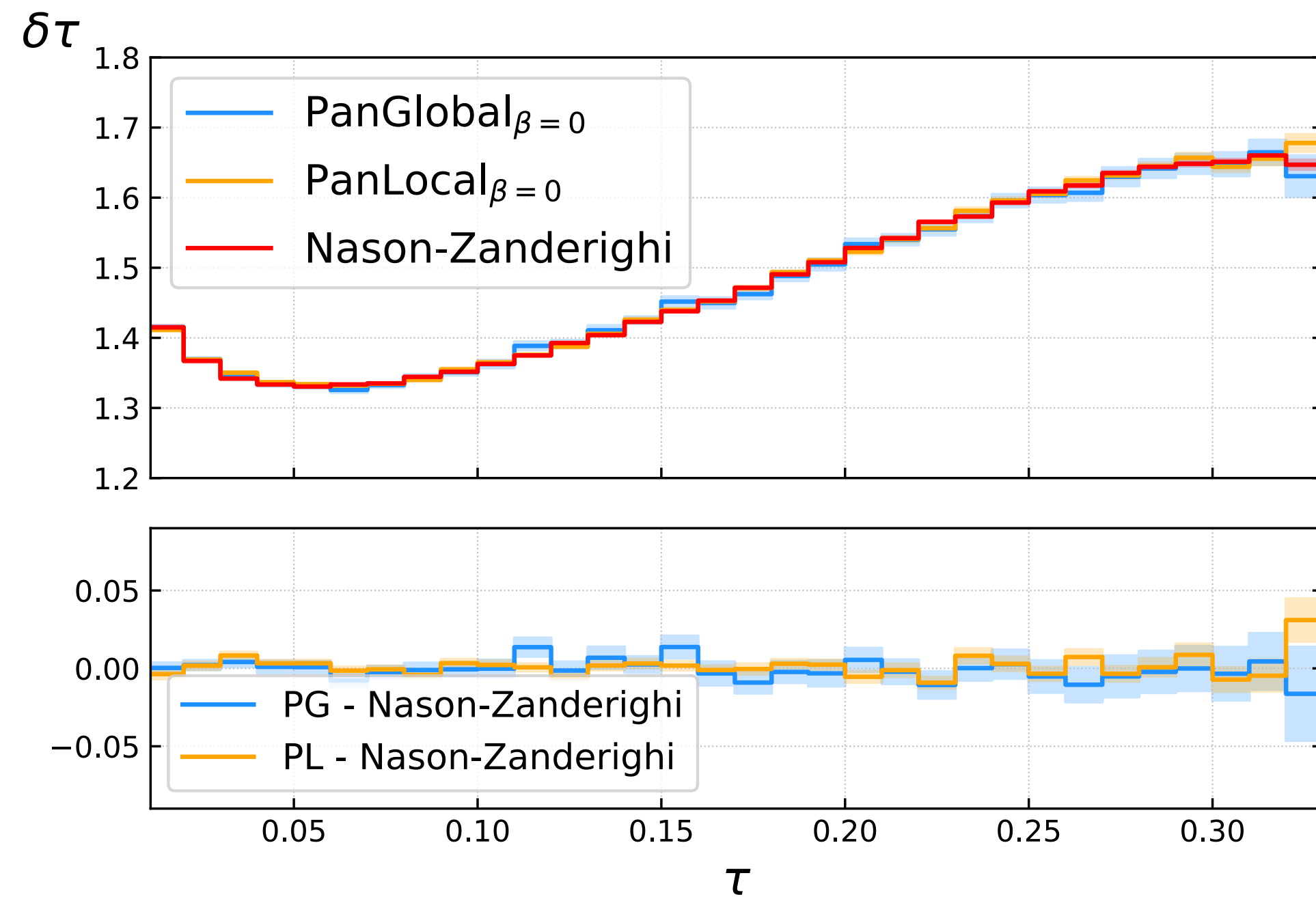
VALIDATION OF OUR METHOD (1)

COMPARISON TO NASON-ZANDERIGHI, 3-jet

We reproduce the non-perturbative shift in event shapes generating a 3-jet configuration and adding subsequently a gluer.

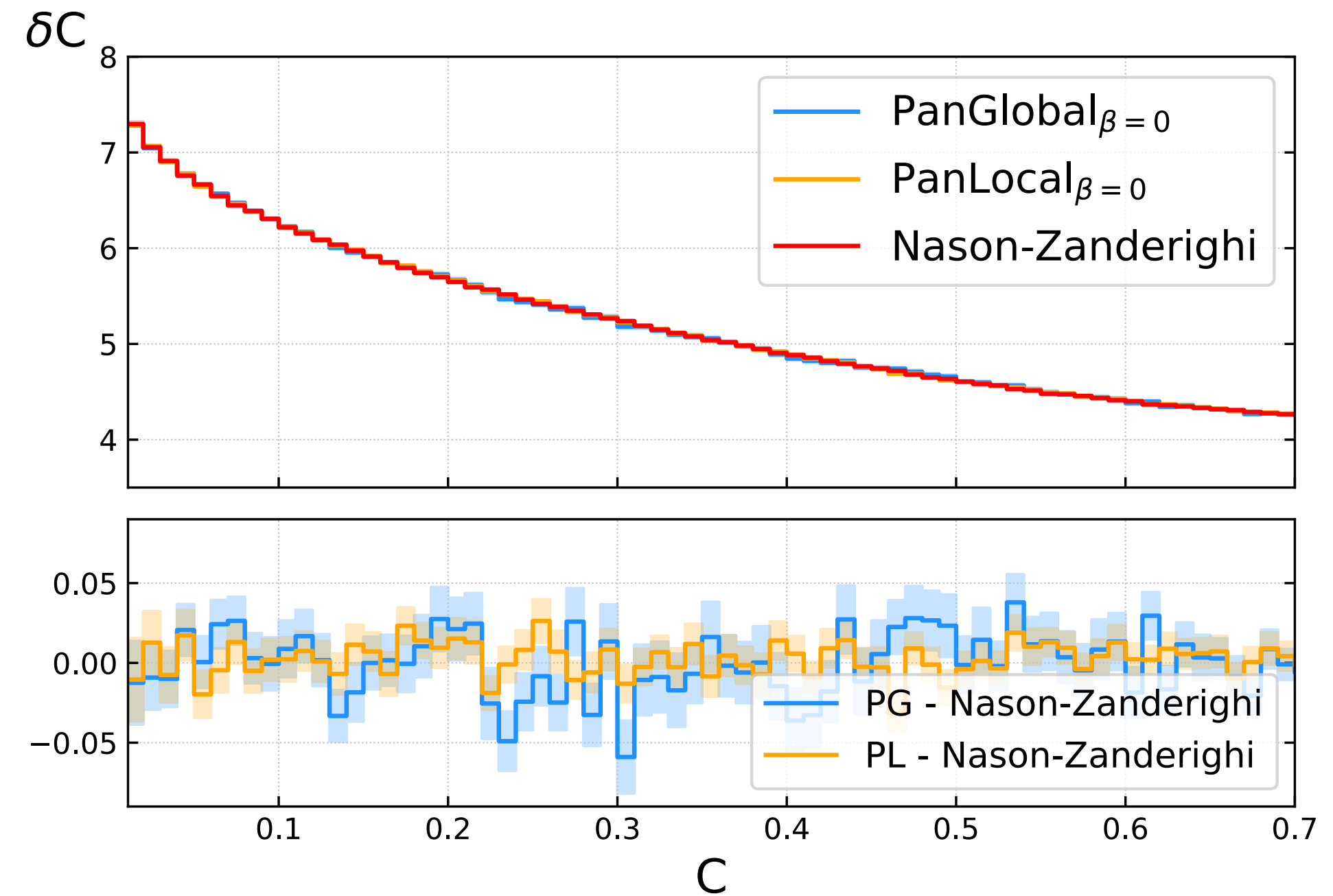


Event shape: τ



$C_F = C_A/2, \ln v_{min} = -12, \ln v_{gluer} = -16$

Event shape: C



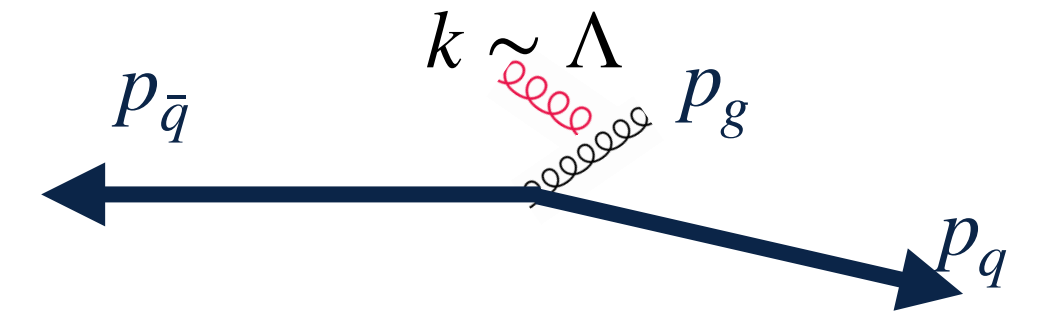
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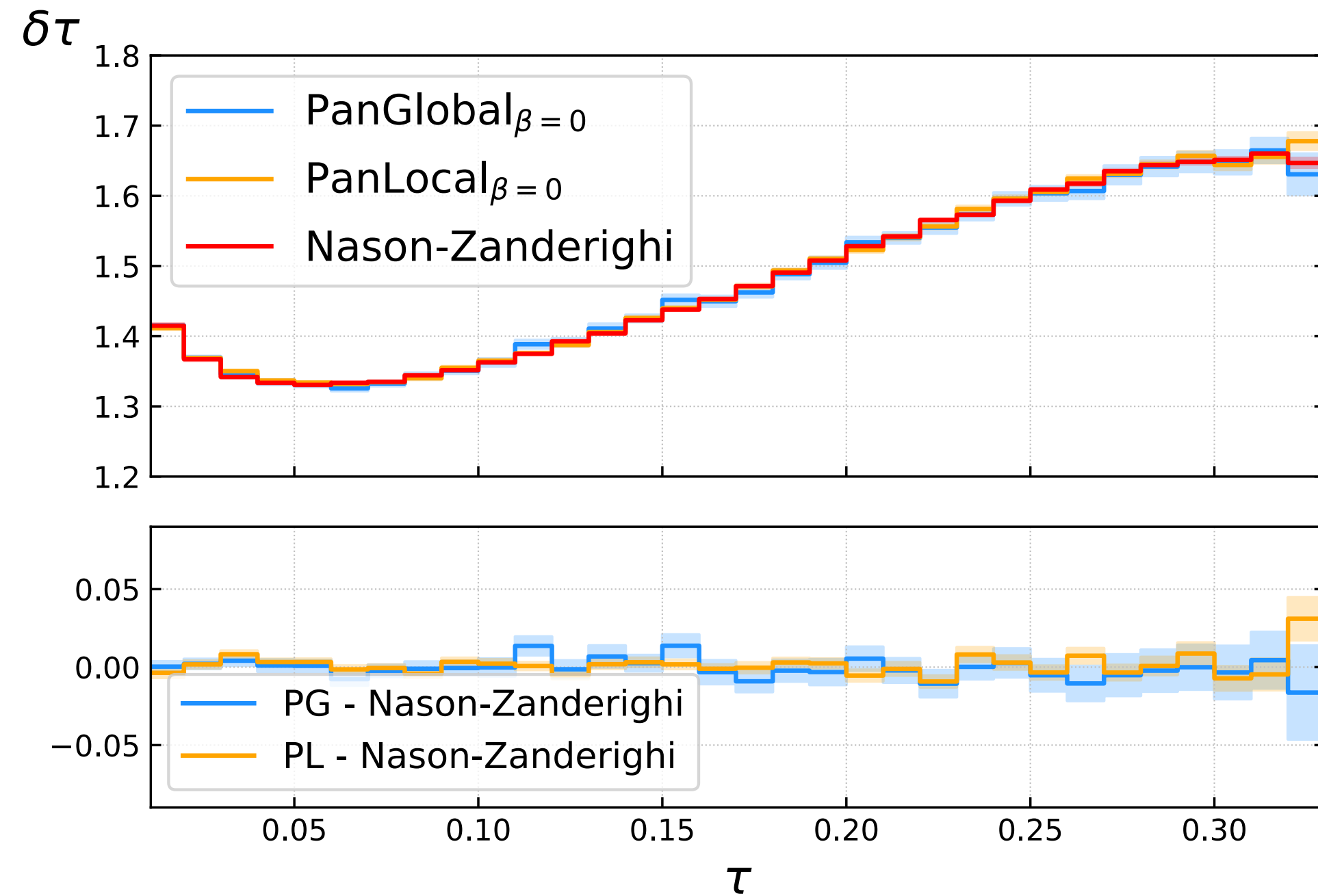
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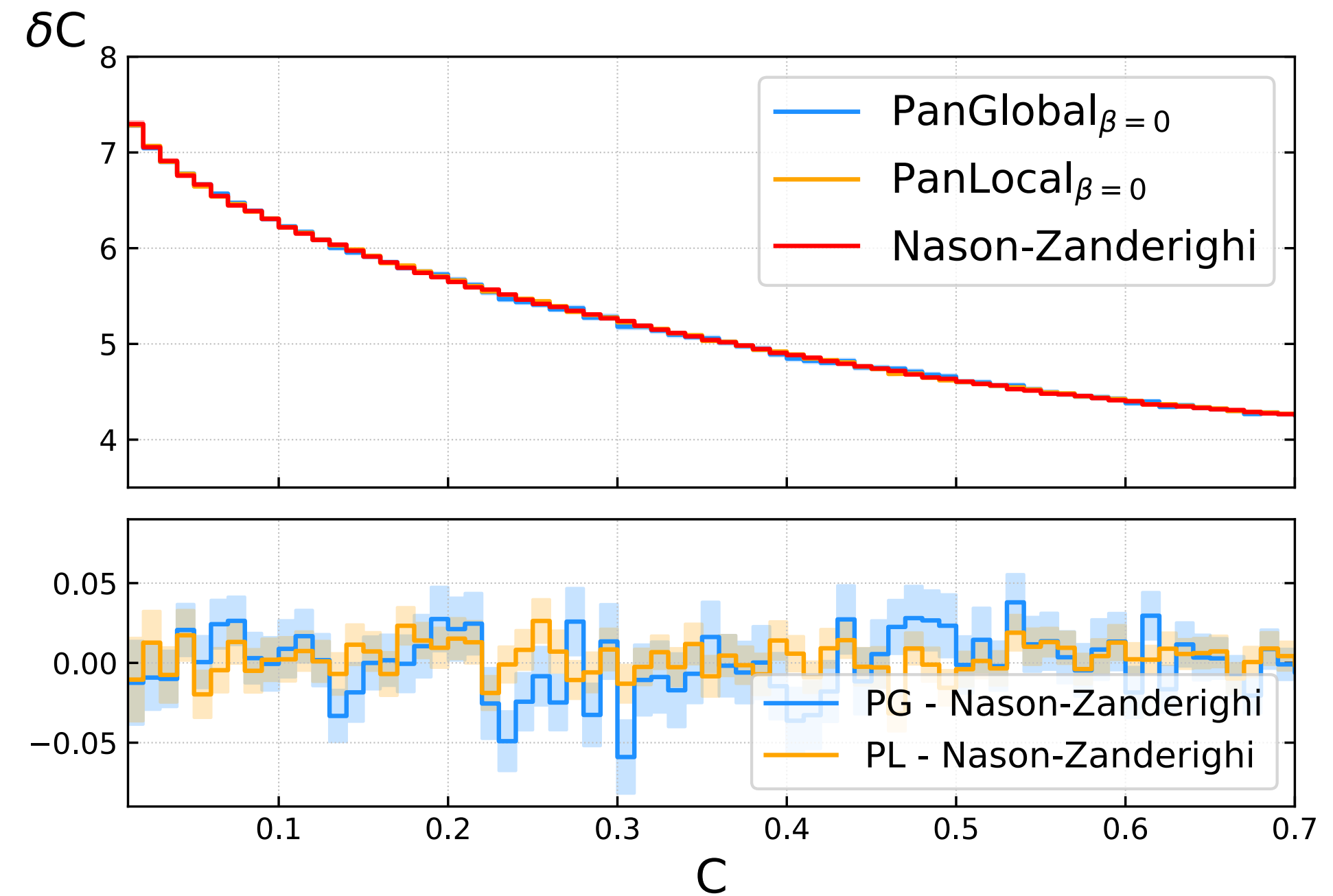


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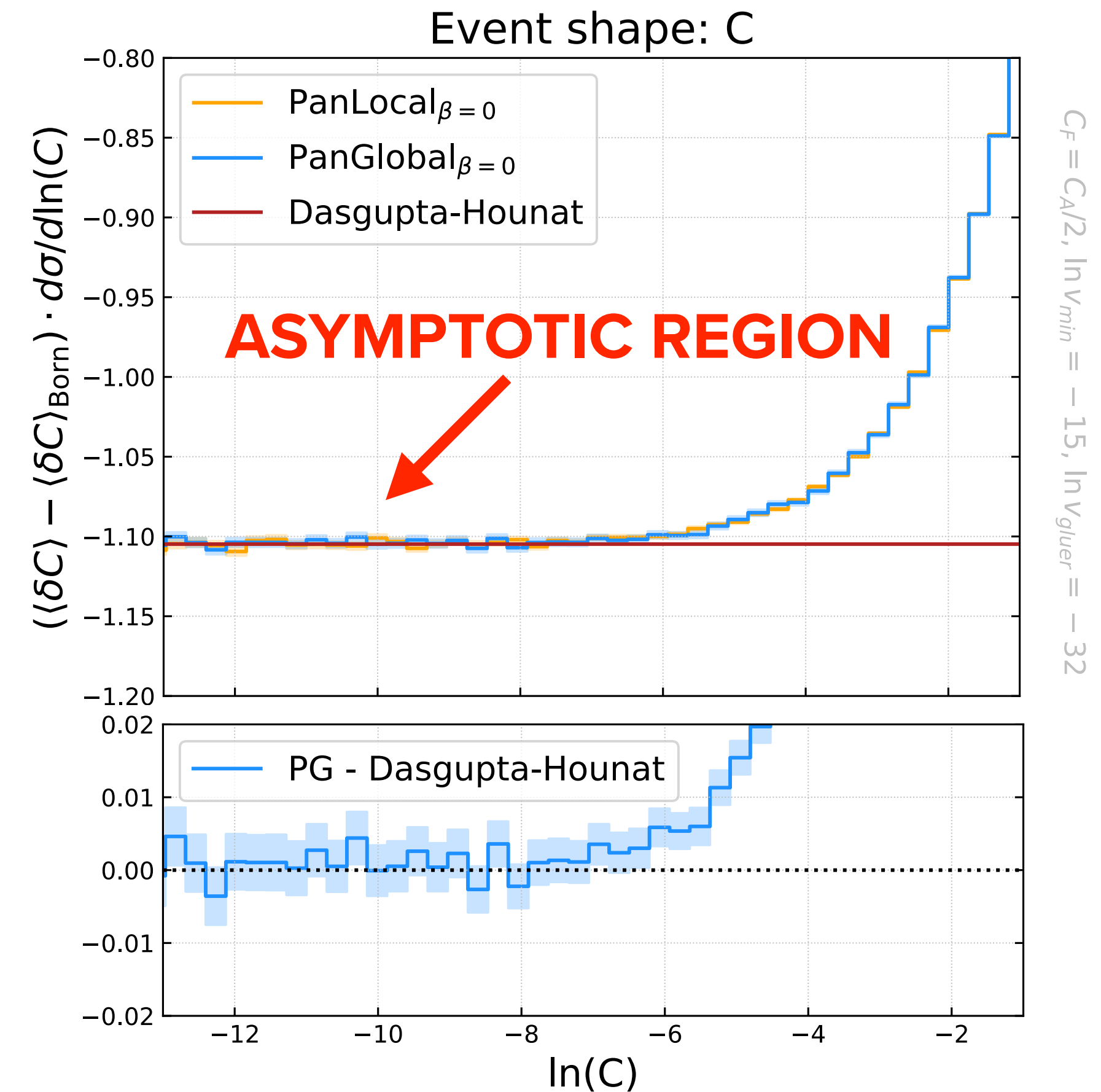
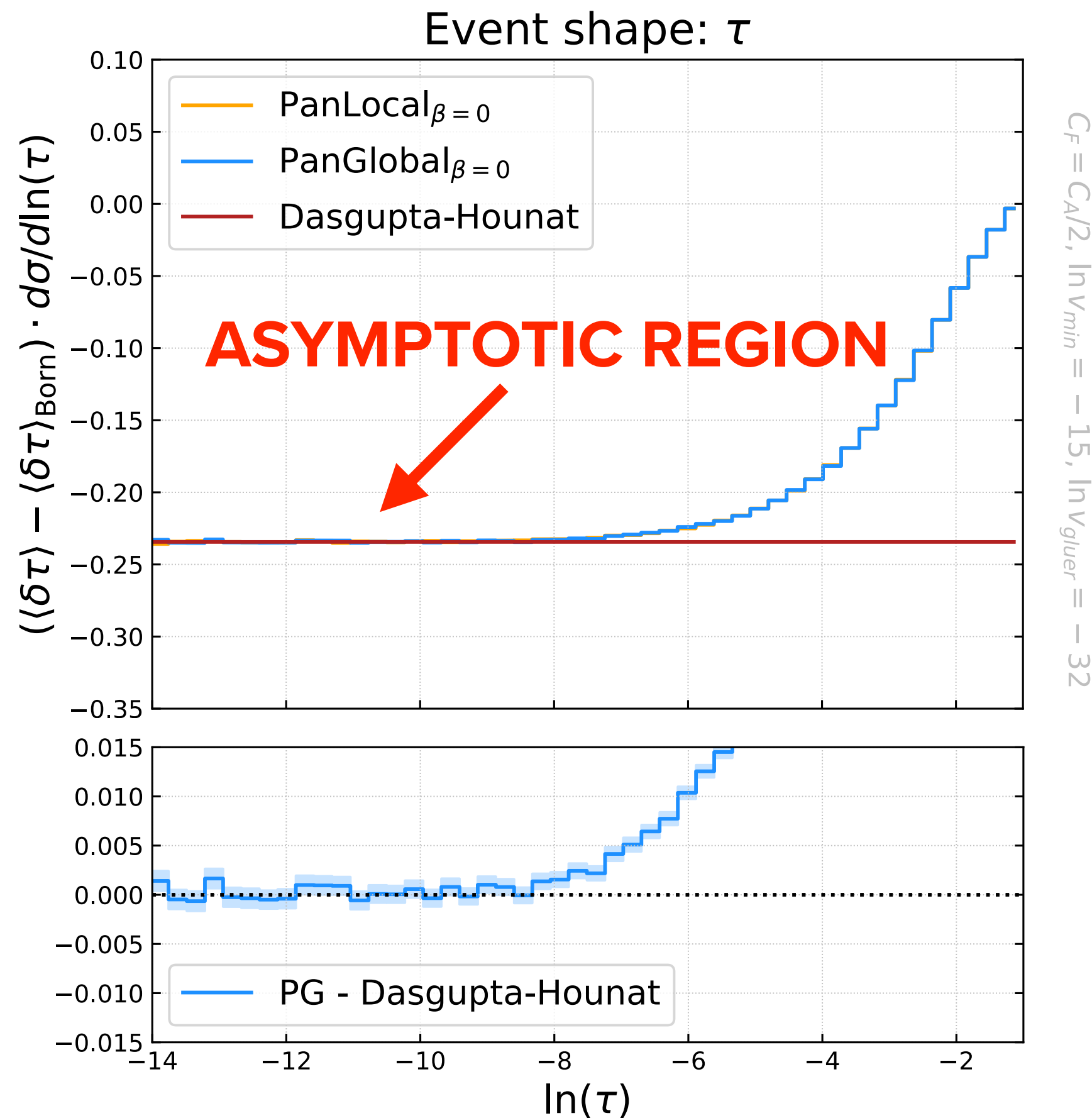
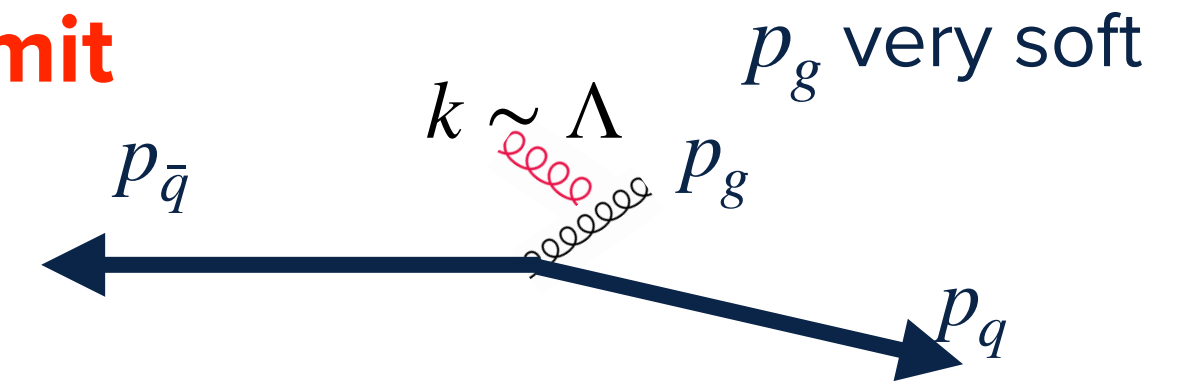


Good agreement - same conclusions hold for y23, broadening and heavy-jet mass.

VALIDATION OF OUR METHOD (2)

COMPARISON TO DASGUPTA-HOUNAT, 3-jet in asymptotic limit

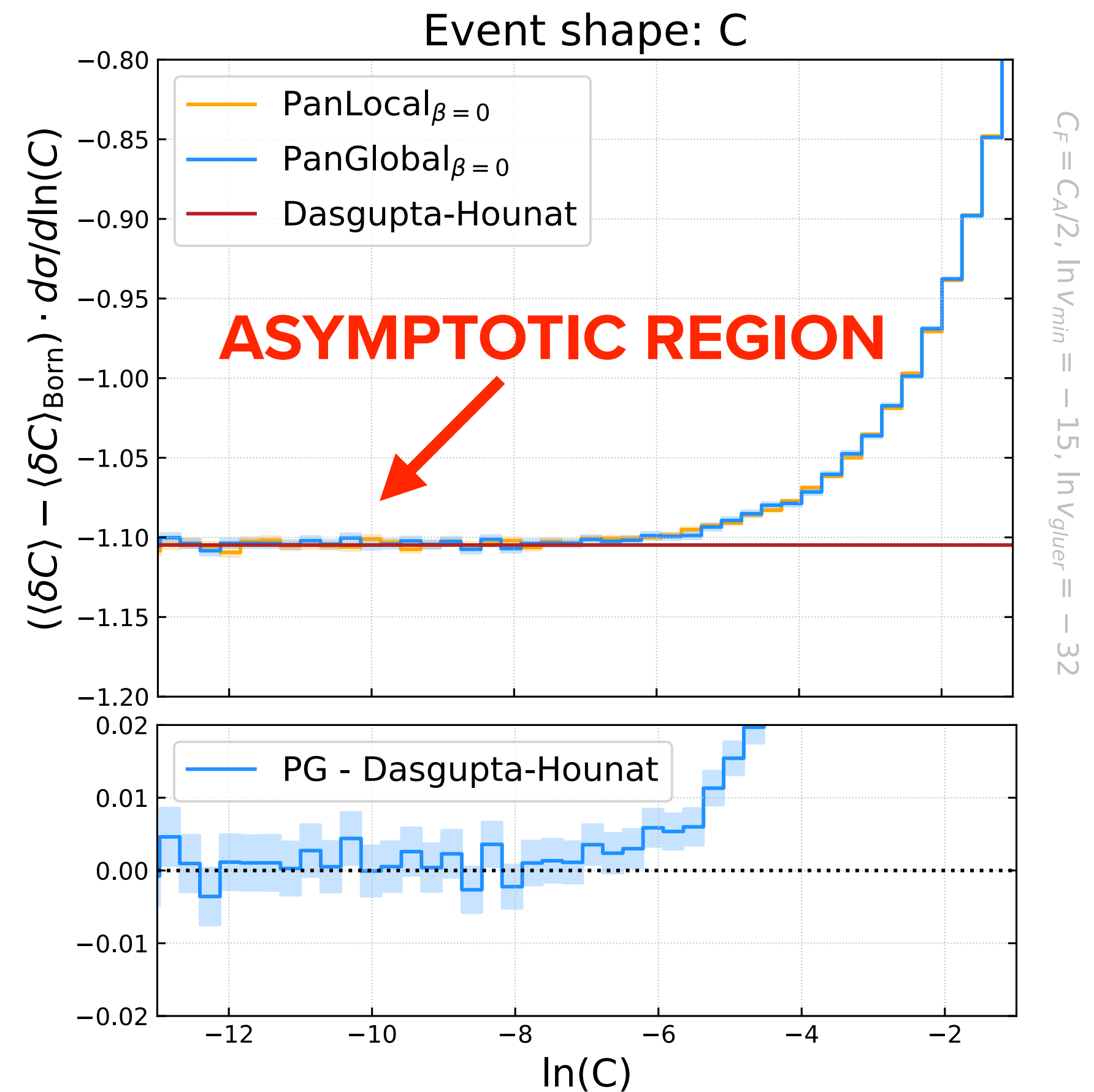
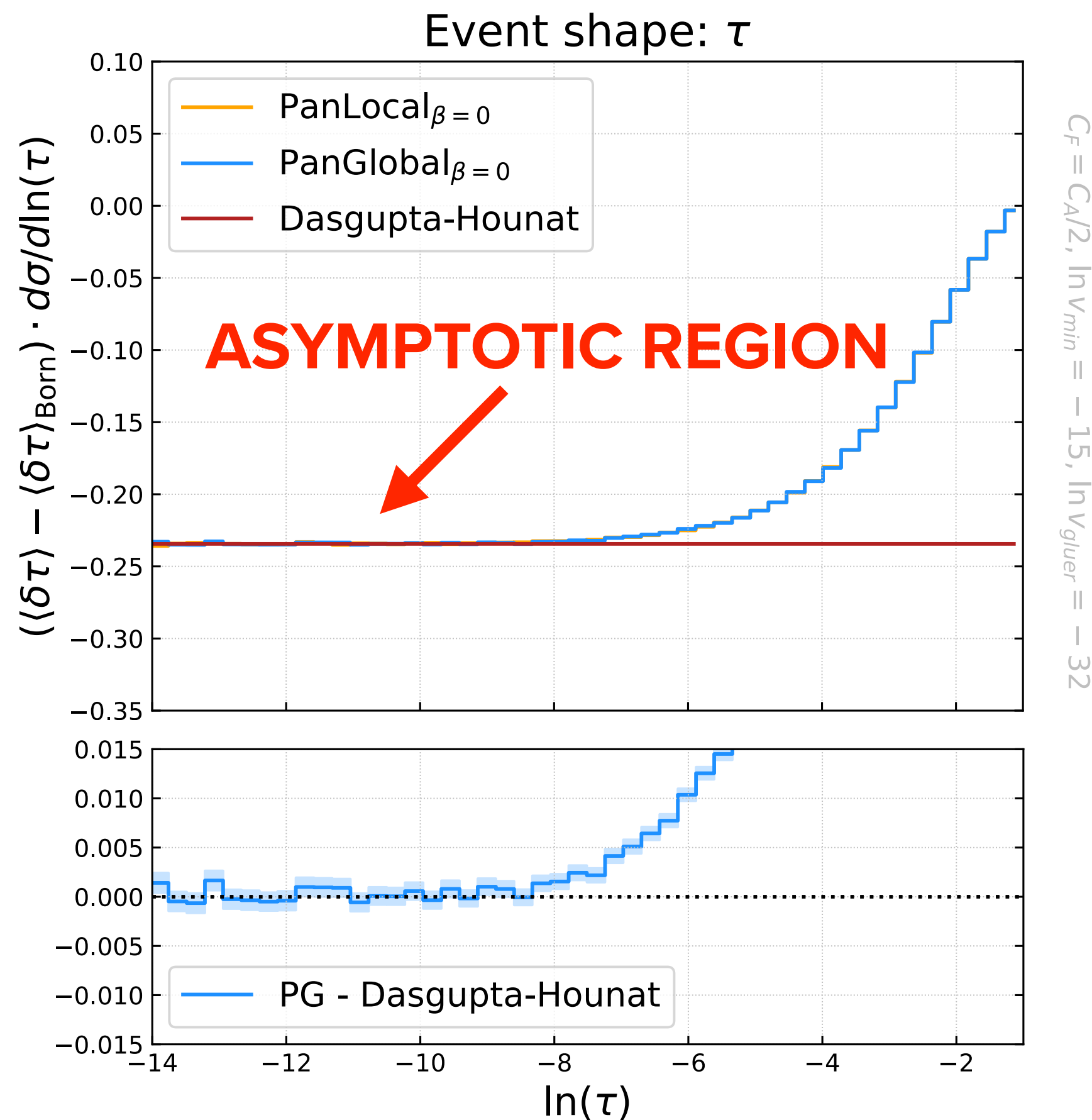
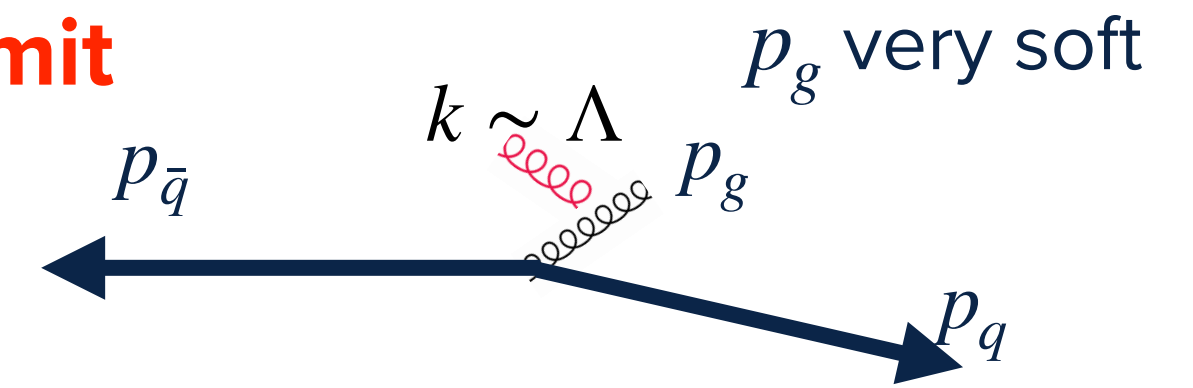
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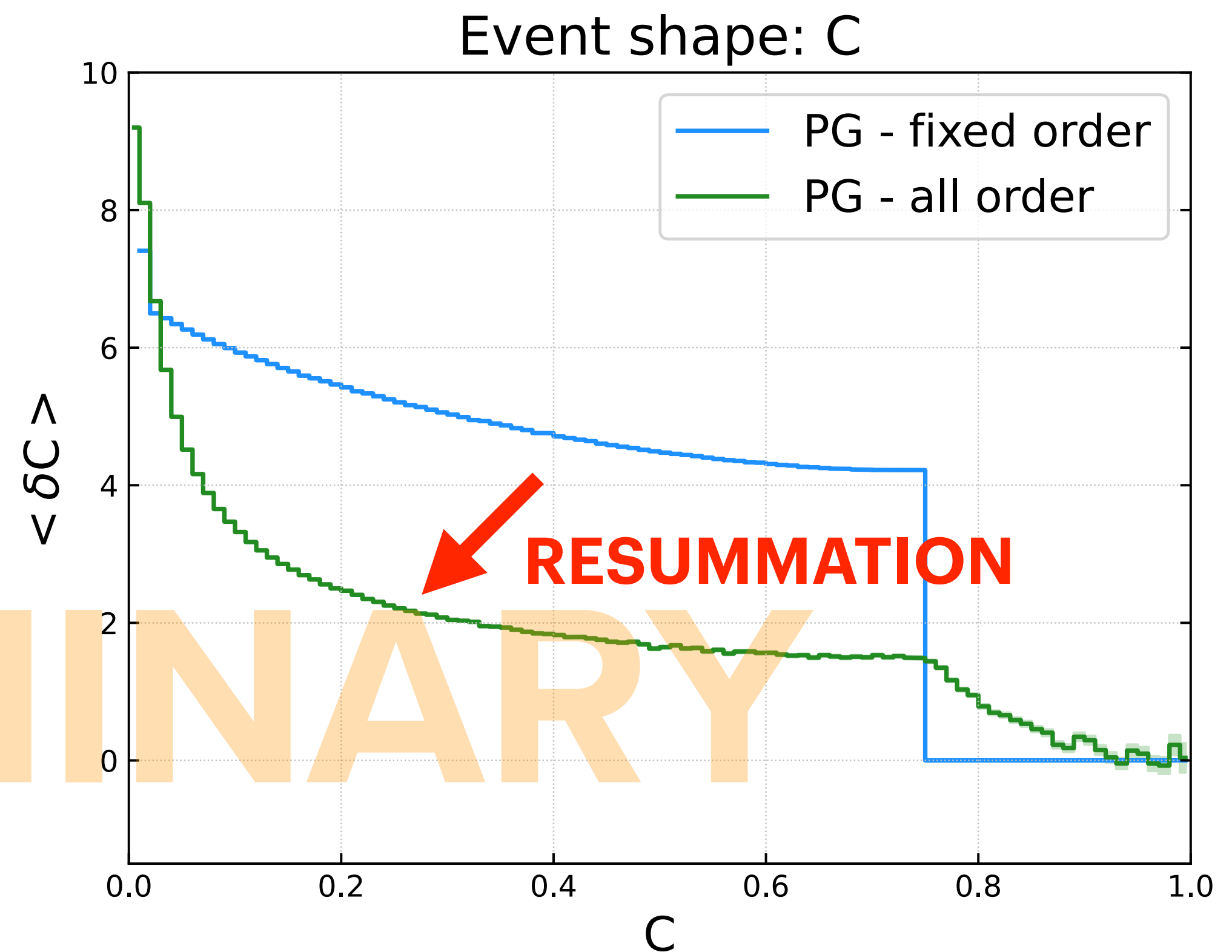
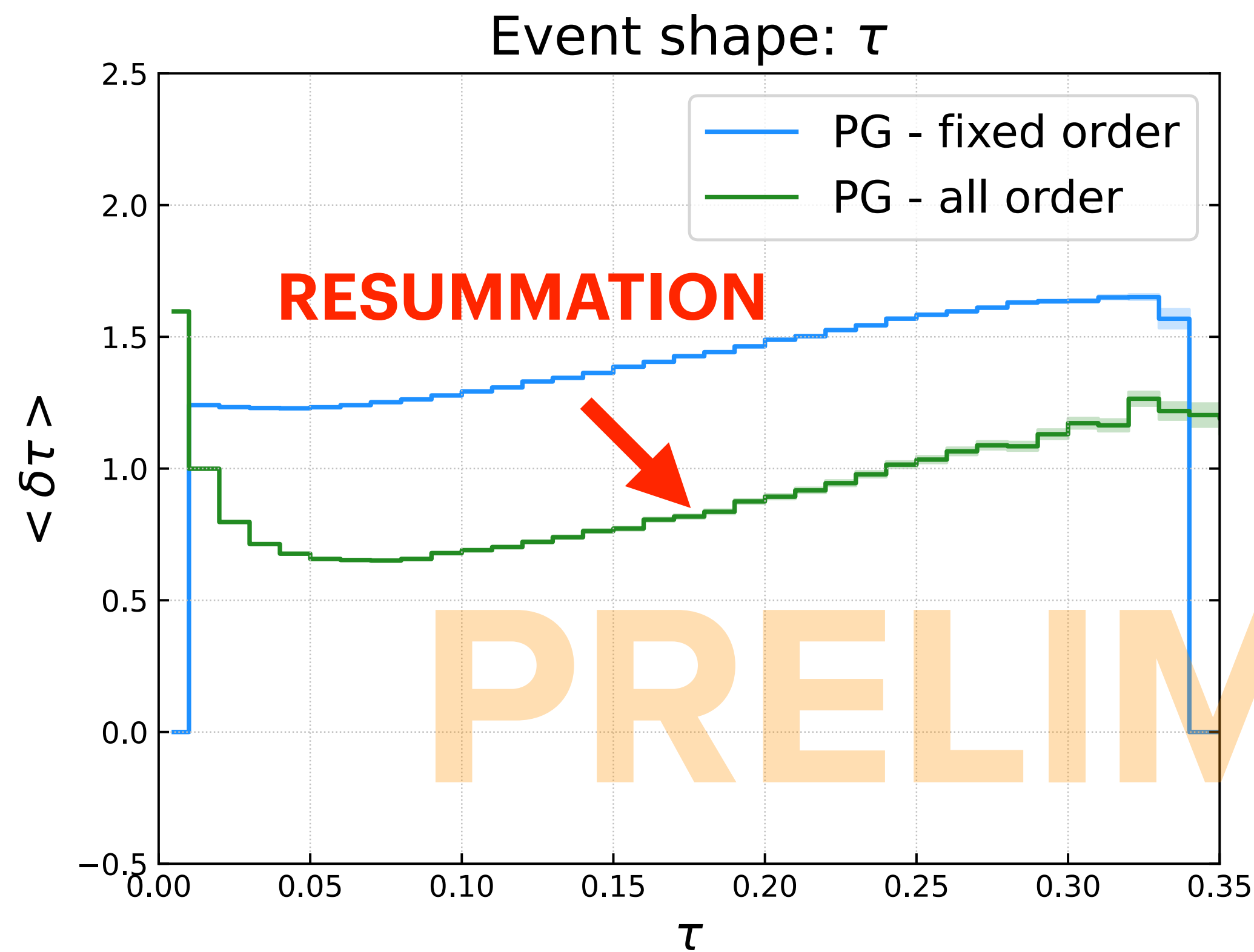
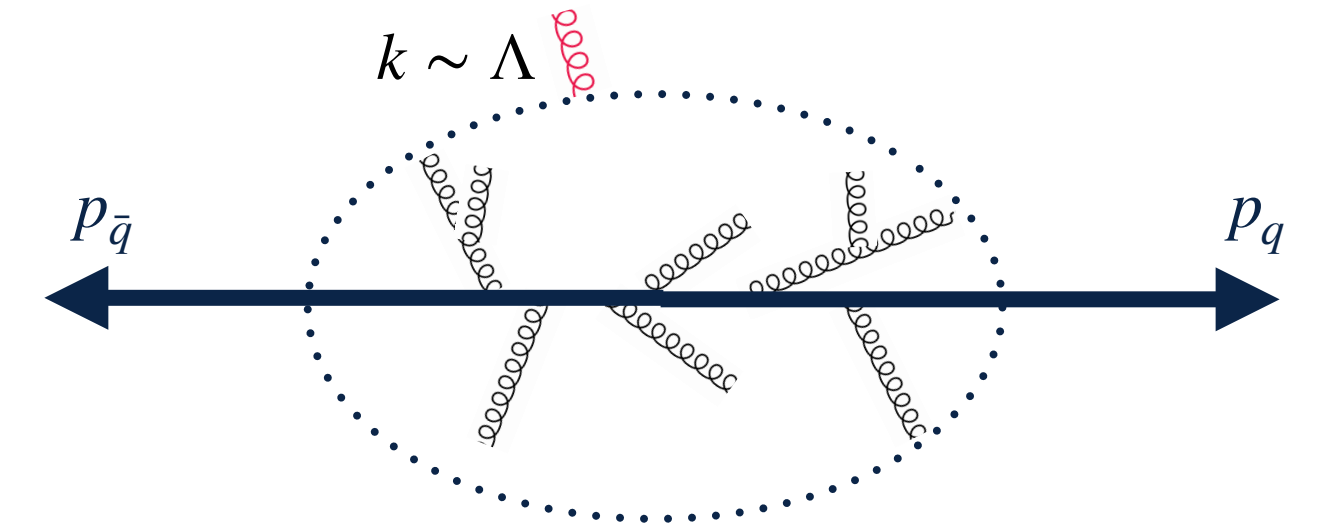
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Good agreement. Possibility to predict the same correction for other event shapes.

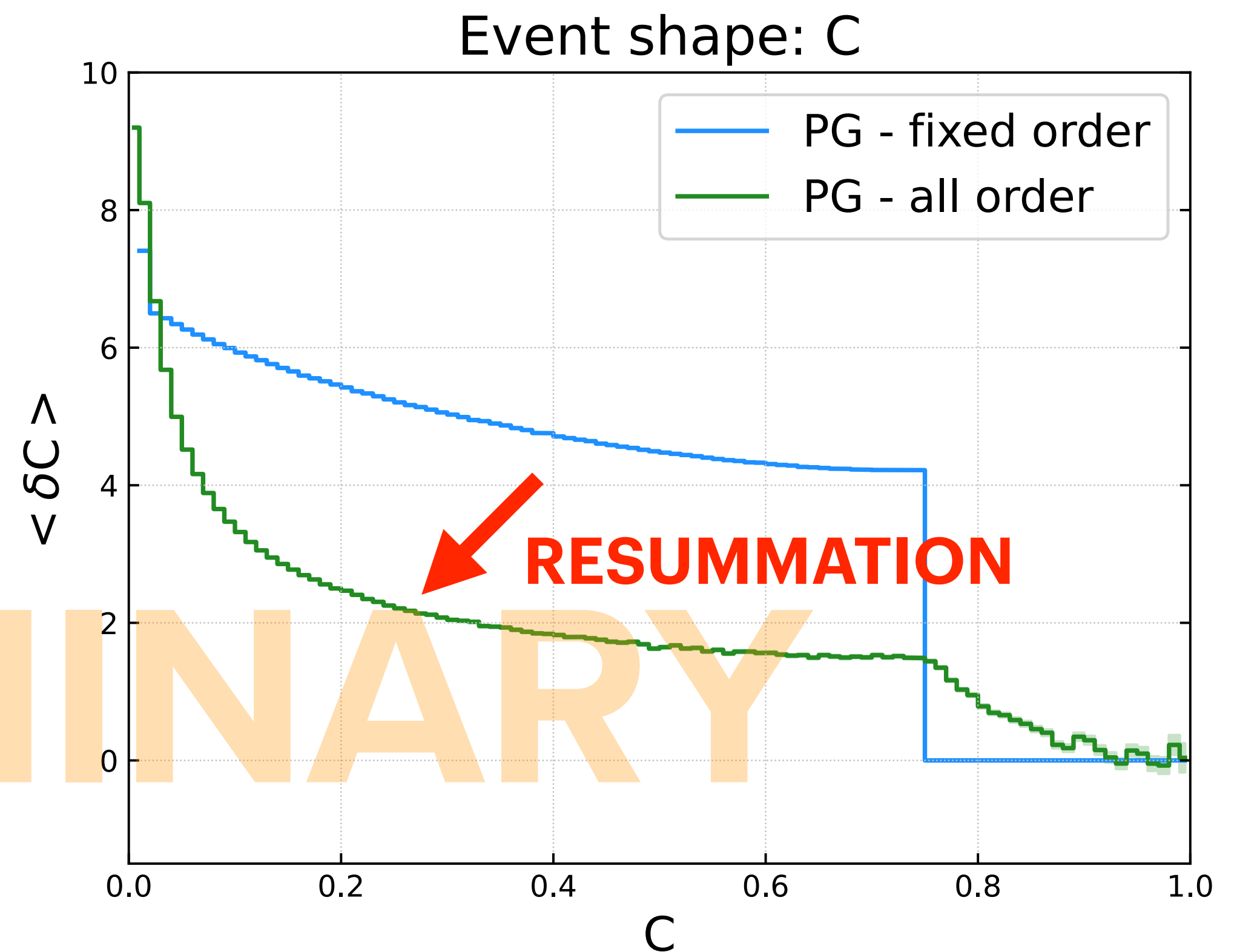
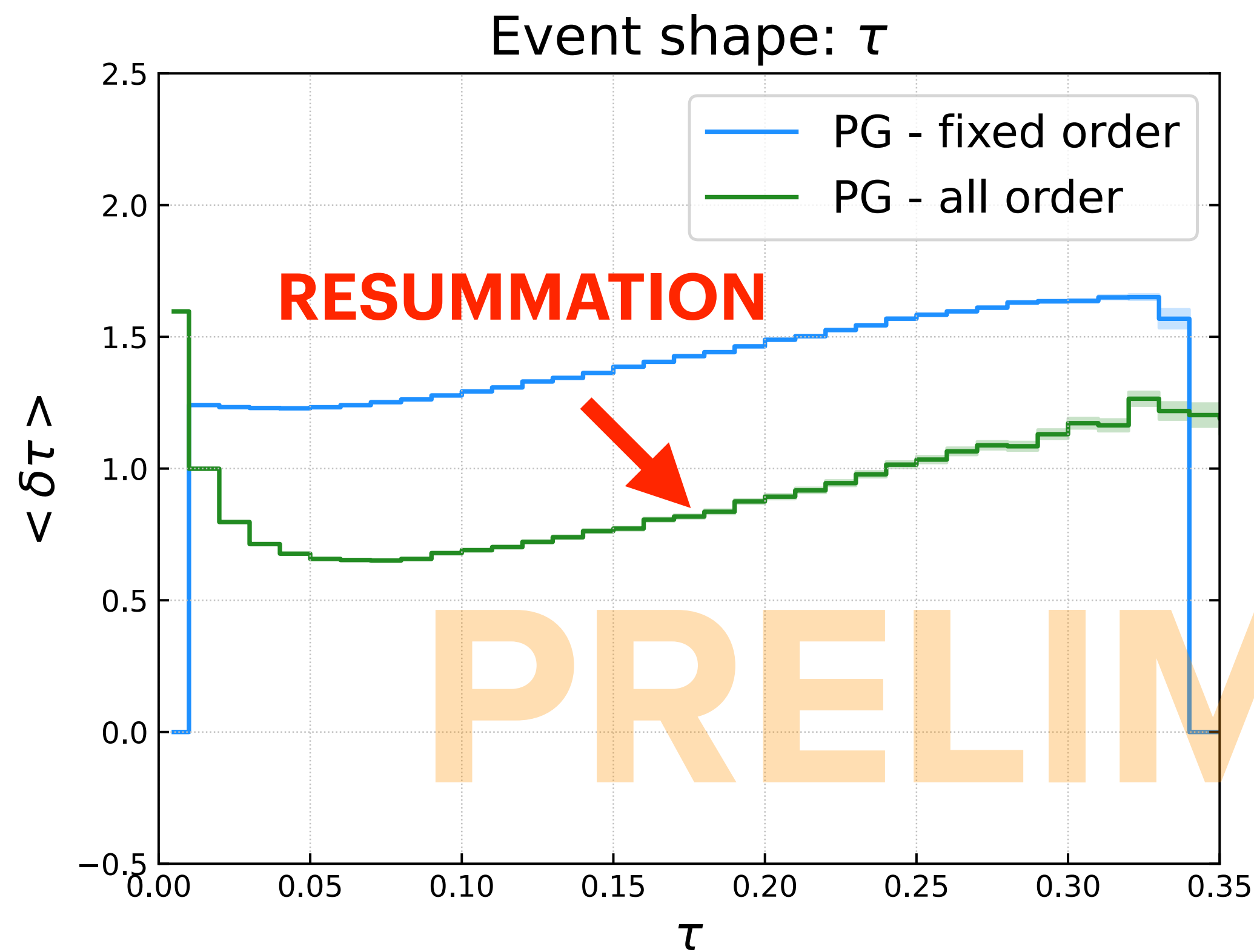
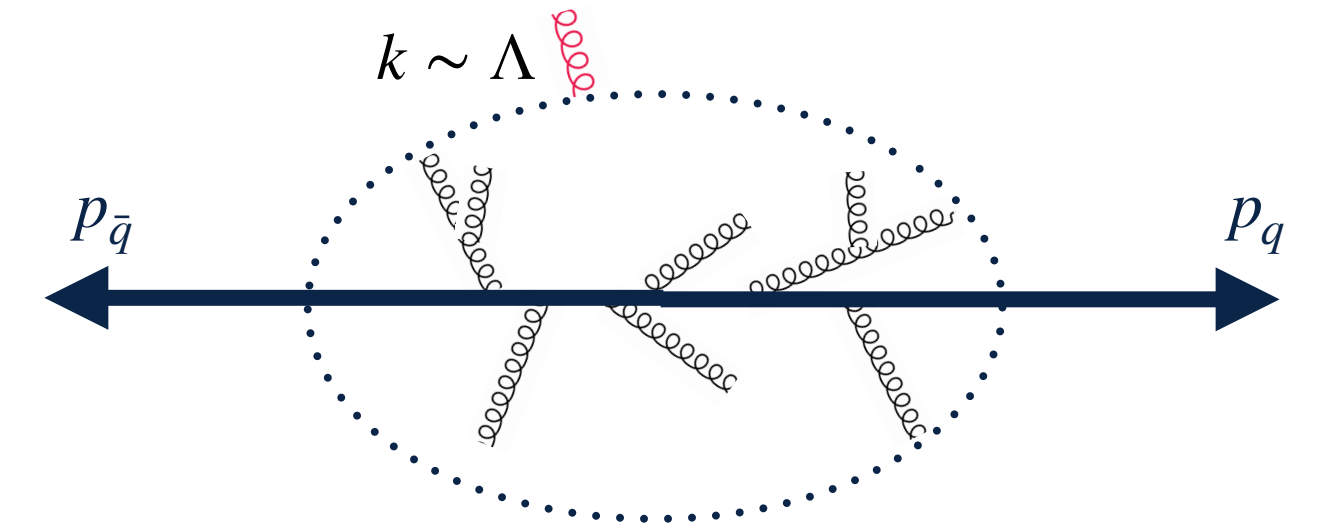
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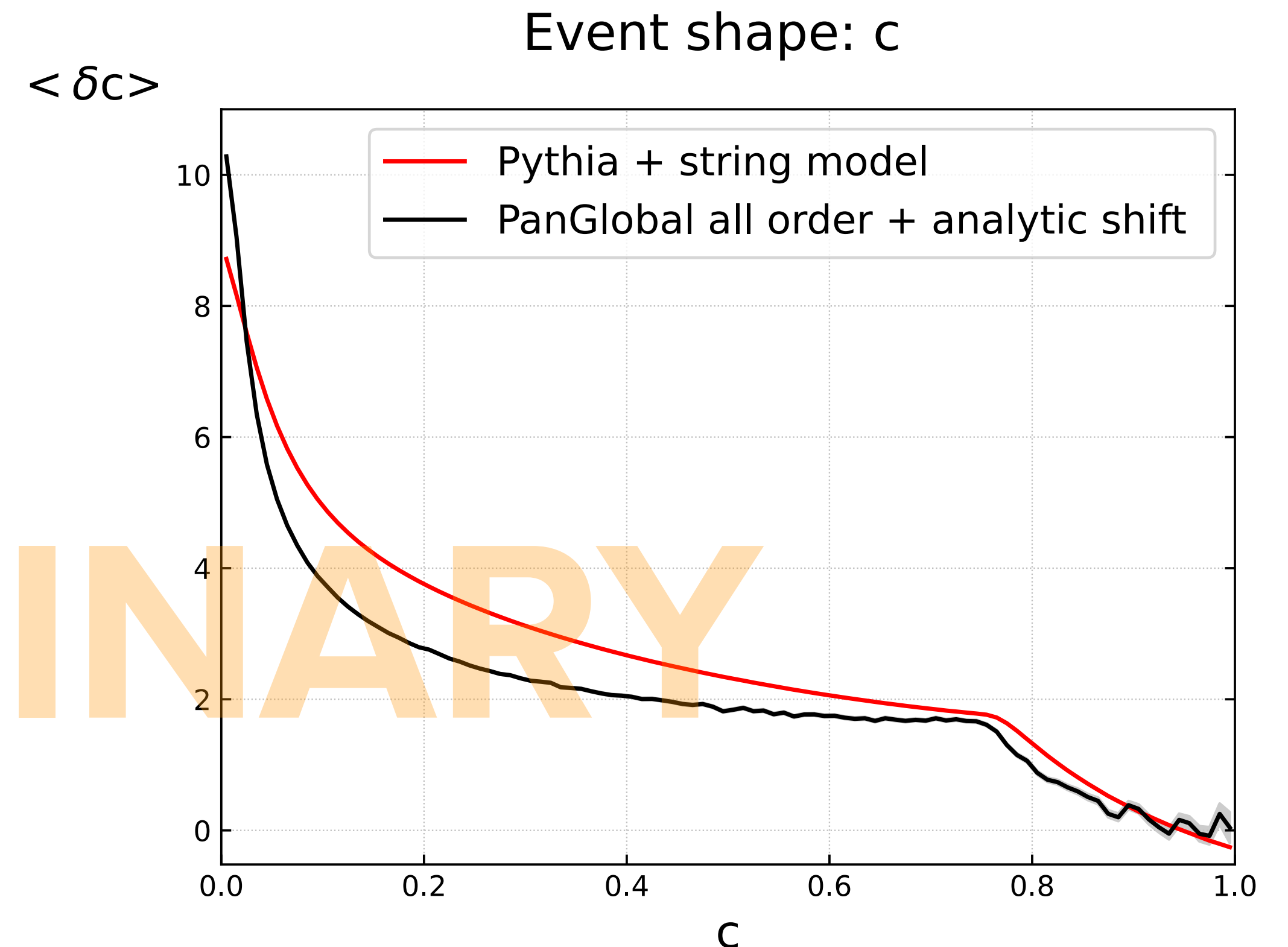
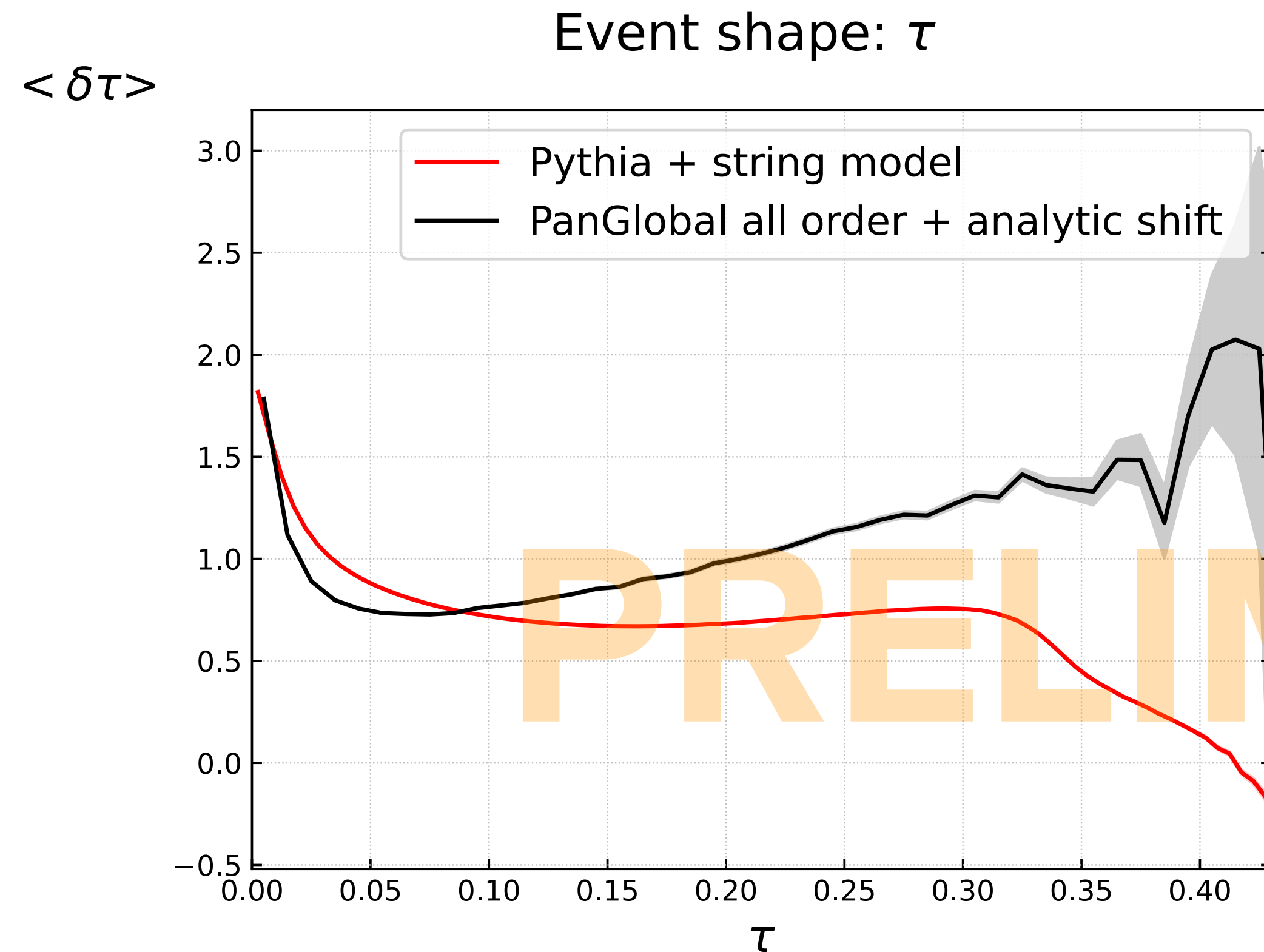
Very large effect that needs to be taken into account for e.g. α_s extrapolation.

CAN WE LEARN SOMETHING ON HADRONIZATION?

Q: Do usual hadronization models (e.g. string model in PYTHIA) capture this all order effect?

Pythia $\rightarrow \delta v = v^{\text{hadron}} - v^{\text{parton}}$
PanGlobal $\rightarrow \delta v = v^{\text{shower+gluer}} - v^{\text{shower}}$

- Curves normalized to the average shift in the thrust.
- Ambiguities in the treatment of masses (see backup)

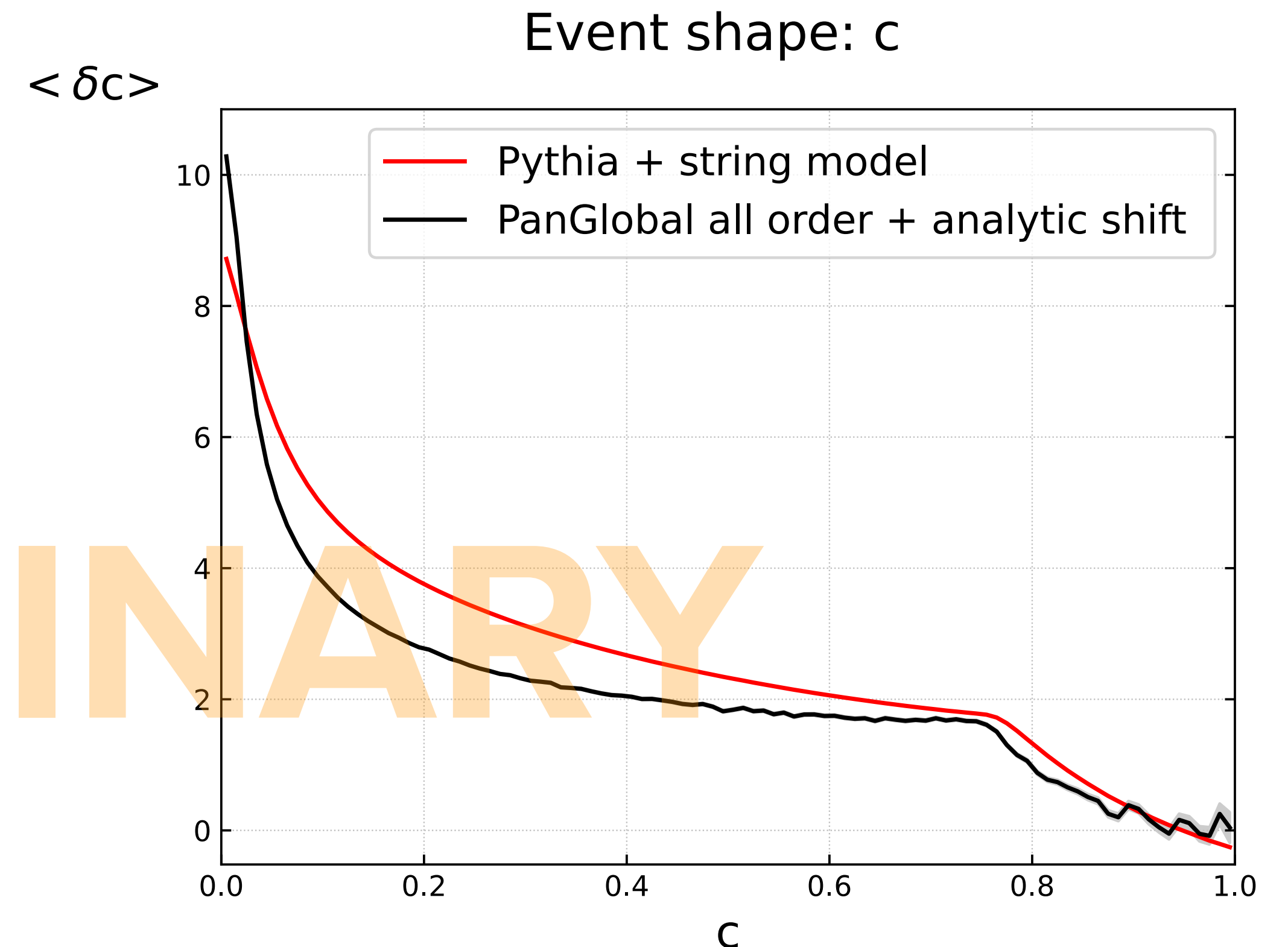
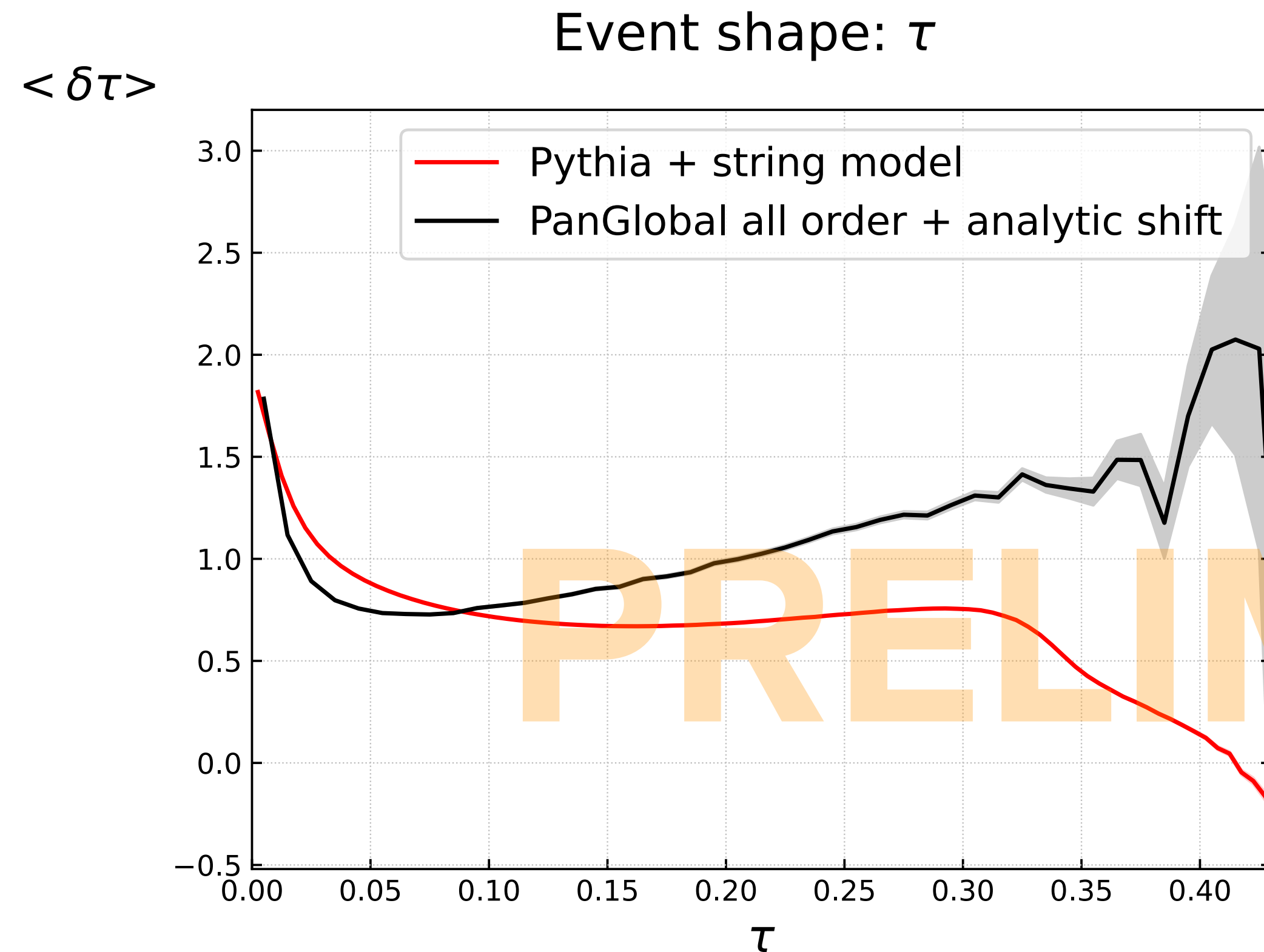


CAN WE LEARN SOMETHING ON HADRONIZATION?

Q: Do usual hadronization models (e.g. string model in PYTHIA) capture this all order effect?

Pythia $\rightarrow \delta v = v^{\text{hadron}} - v^{\text{parton}}$
PanGlobal $\rightarrow \delta v = v^{\text{shower+gluer}} - v^{\text{shower}}$

- Curves normalized to the average shift in the thrust.
- Ambiguities in the treatment of masses (see backup)



The shape looks consistent for C-parameter, not so clear for the thrust.

CONCLUSIONS

Our understanding of analytic models for the description of hadronization effects has evolved significantly in the past few years, but there is still room for improvements.

- **Analytic models are strongly affected by higher-orders effects.** We are currently working on the inclusion of all order effects from the **resummation of the anomalous dimension** in the DW model.
- We are exploring how we could gain **new directions of insight into the possible behaviour of hadronization.** *
- We plan to analyze the Q dependence of the analytic model presented before, compared to e.g. Pythia.

* Hadronization is fundamental also for the precision physics programme of the LHC: gaining insight into this beyond Herwig/Pythia/Sherpa differences is crucial.

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We ultimately plan to explore the impact of our findings on the fit of α_s .

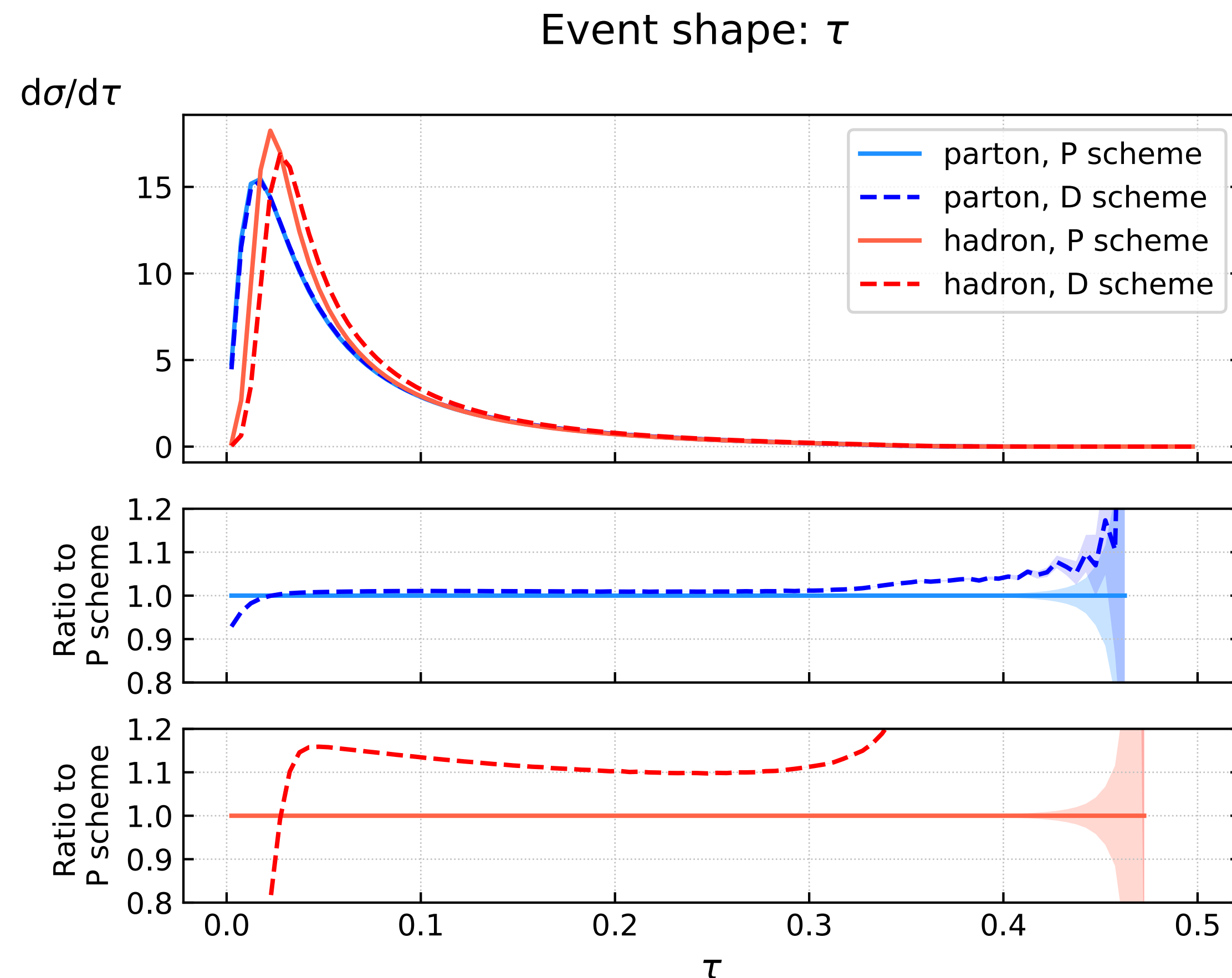


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BACKUP

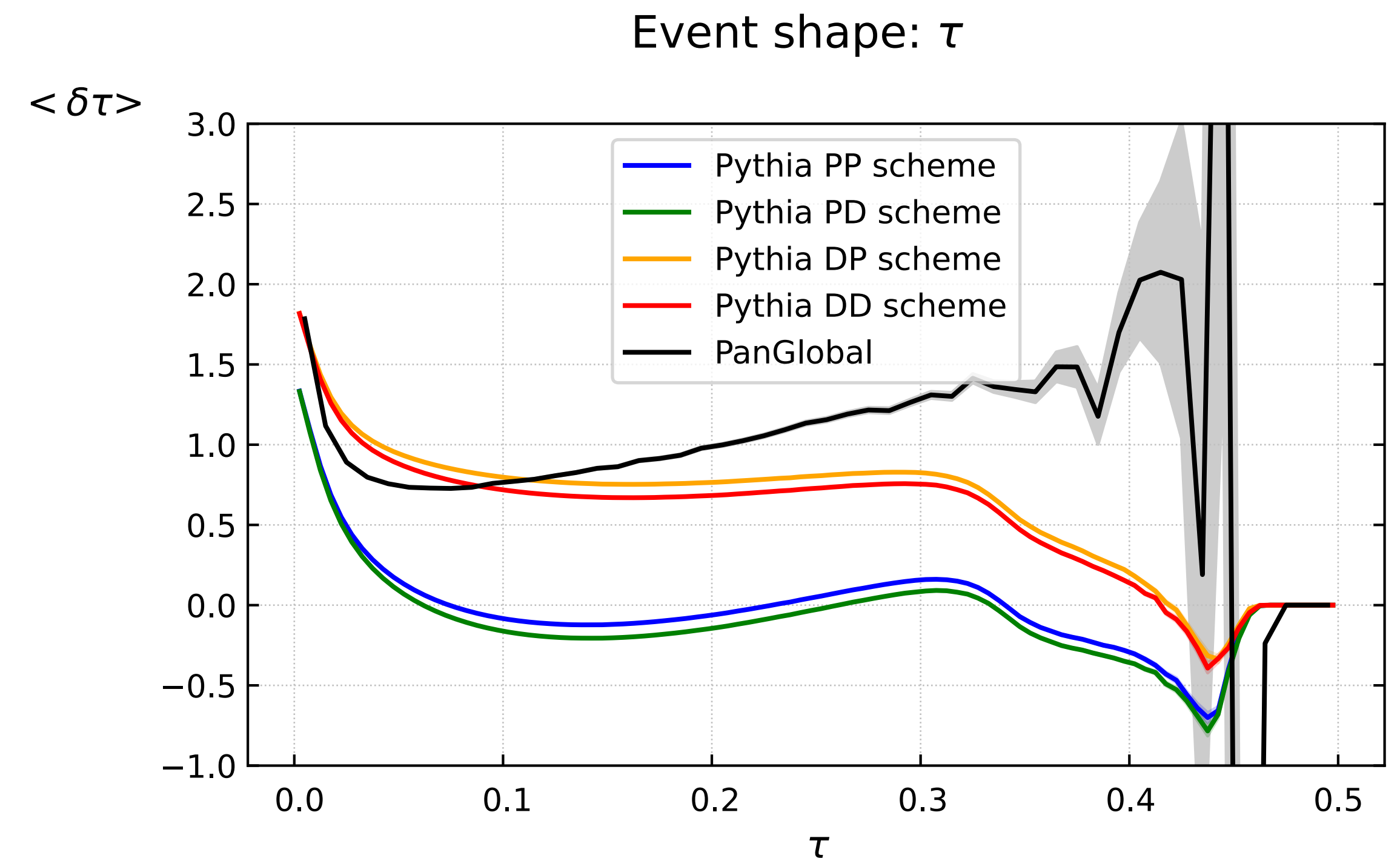
TREATMENT OF MASSES

Perturbative calculations assume massless partons, while experimental measurements deal with massive hadrons. In general, the definition of an event shape is different in the two cases.



P scheme: $\vec{p}_i \rightarrow \vec{p}_i, \quad E_i \rightarrow |\vec{p}_i|$

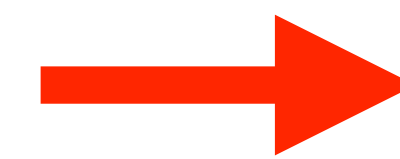
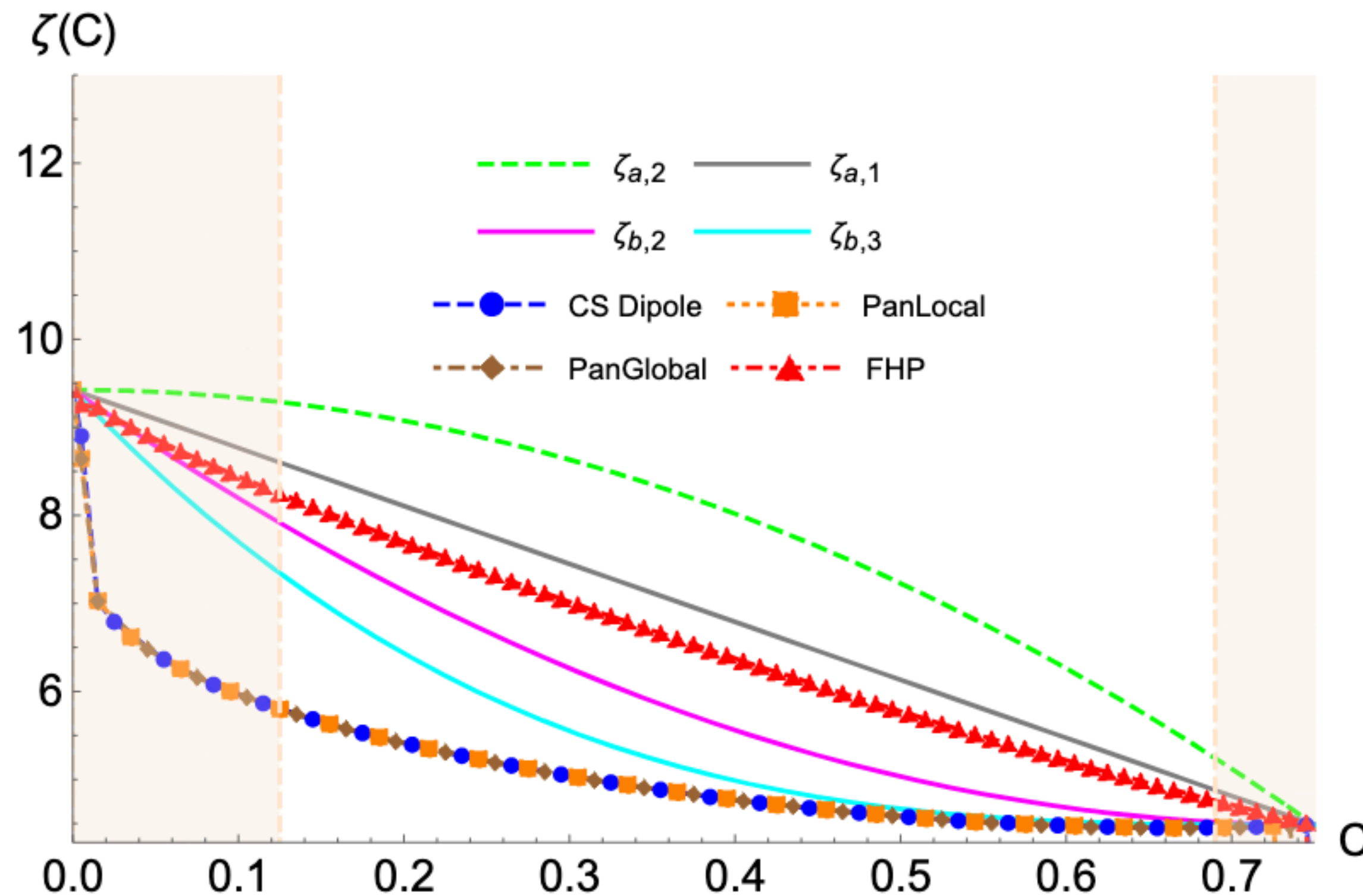
D scheme: massive particles are decayed isotropically into 2 massless partons



[Salam, Wicke hep-ph/0102343]

IMPACT OF MAPPINGS

The calculation of the non perturbative shift associated to a gluon emission requires a recoil scheme to enforce energy-momentum conservation. What is the impact of this choice on our result?



Same results for mappings in which the longitudinal recoil is kept local within the dipole

The mappings adopted in this work (PG and PL) satisfy the smoothness requirement in the soft limit that is needed in order for the recoil effects not to contribute to linear power corrections.

[Luisoni, Monni, Salam 2012.00622]

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