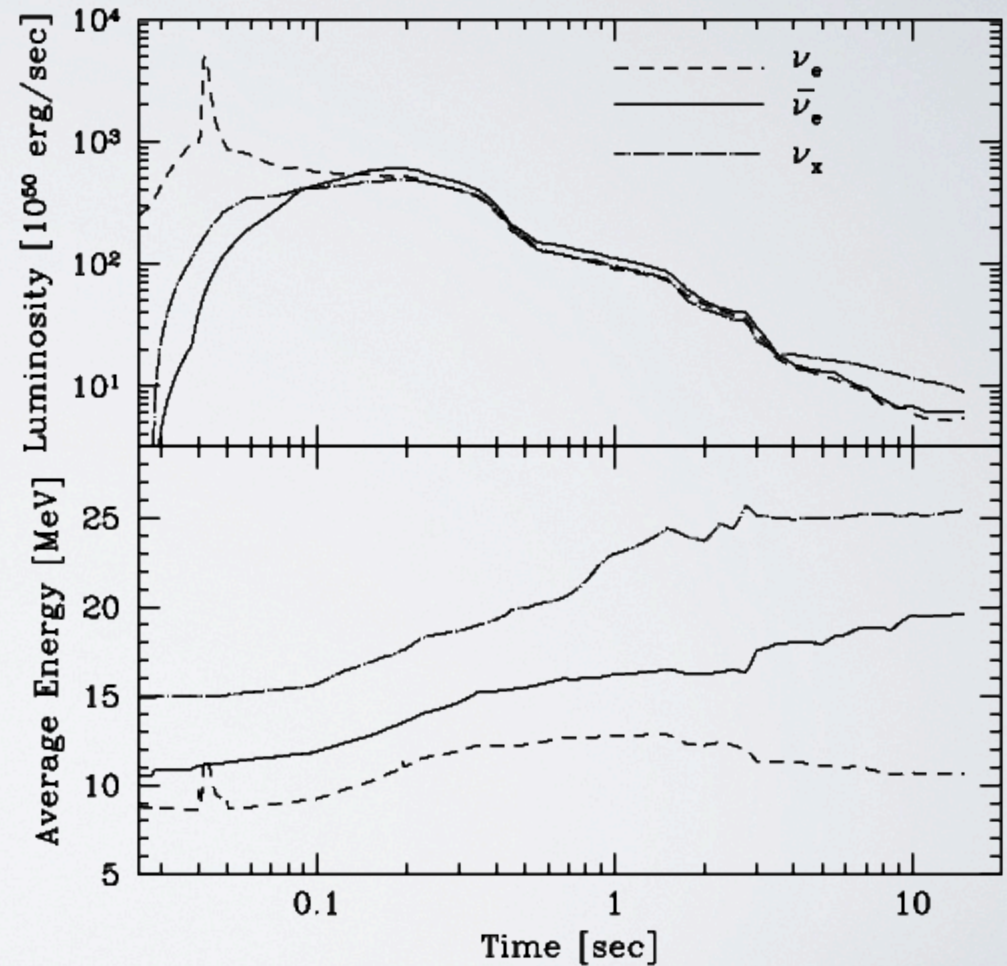


NEUTRINOS FROM SN COLLAPSE

Neutrino collapse and neutrino production
Detection of a galactic supernova: SN 1987A
Current detectors



(Hubble image)



What did we learn?

- Luminosity \approx total energy budget
 - Energy emitted is of *gravitational* nature:

$$L_\nu \approx G M_f^2/R_f - G M_i^2/R_i \sim 3 \cdot 10^{53} \text{ ergs } \checkmark \quad (R_f \sim 10 \text{ Km})$$
- Energy spectrum: \sim *Fermi Dirac (thermal)*
 - $E \approx 3.15 T \sim 15\text{-}20 \text{ MeV } \checkmark$
- Duration of neutrino burst \sim diffusion time
 - Time \approx (size²)/(mean free path) $\sim 10 \text{ s } \checkmark$

Limit on neutrino masses now well overcome:

$$v = \frac{p}{E} = \frac{\sqrt{E^2 - m^2}}{E} \approx 1 - \frac{1}{2} \frac{m_\nu^2}{E_\nu^2} \quad t_a - t_e = d \times \left(1 + \frac{1}{2} \frac{m_\nu^2}{E_\nu^2}\right) \quad |\Delta t_a - \Delta t_e| = \frac{1}{2} d m_\nu^2 \frac{|E_1^2 - E_2^2|}{E_1^2 E_2^2}$$

propagation time = arrival - emission time

Measure energy of 2 events and the time dispersion between the 2 events =>

$$m_{\nu_e} \lesssim 20 \text{ eV}$$

$$F(\bar{\nu}_e) = (1 - p) \cdot F_0(\bar{\nu}_e) + p \cdot F_0(\bar{\nu}_\mu).$$

The anti- ν_e spectrum is a mixture of the 2 original spectra for anti- ν_e , ν_μ permutation factor $< 35\%$ (99%CL)

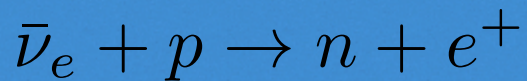
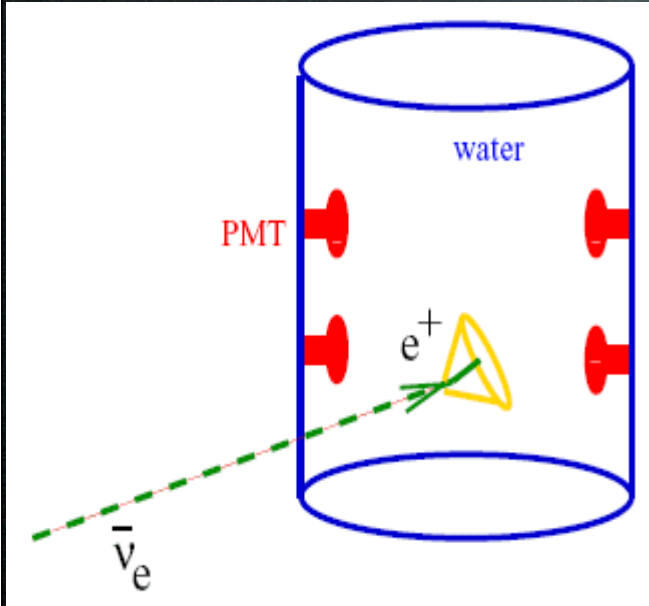
SN detectors

Detector	Type	Location	Mass (ton)	Events @ 8kpc	Status
Super-K	Water	Japan	32000	8000	Running (SK IV)
LVD	Scintillator	Italy	1000	300	Running
KamLAND	Scintillator	Japan	1000	300	Running
Borexino	Scintillator	Italy	300	100	Running
IceCube	Ice	South Pole	600000	1 million	Running
Baksan	Scintillator	Russia	330	50	Running
Mini-BOONE	Scintillator	USA	700	200	Running
HALO	Lead	Canada	76	85	Under construction
Icarus	Liquid argon	Italy	600	230	Almost ready
NOvA	Scintillator	USA	15	3000	Construction started
SNO+	Scintillator	Canada	1000	300	Funded
Lena	Scintillator	?	50000	$\bar{\nu}_e$: 15000 ν_e : 800 ν_x : 6000	Design study 08-10
Hyper-K	Water	DUSEL	300000	$\bar{\nu}_e$: 70000 ν_e : 3000 ν_x : 3000	Design study S4-funded 10 events Andromeda

Detection

SN ν interactions
inverse beta in
Water Čerenkov

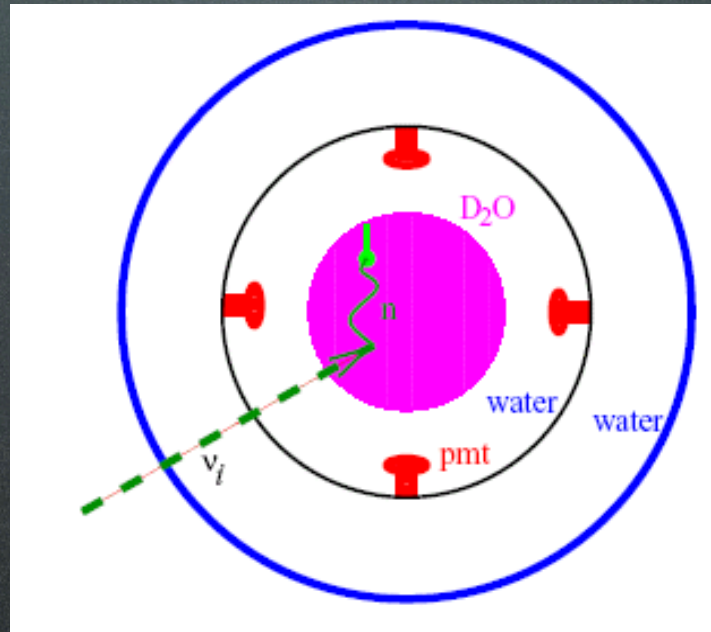
SK, IceCube



1. Deep underground location
2. Using low radioactivity materials
3. Anti-coincidence system
4. Using reactions with good signature
5. Coincidence of signals

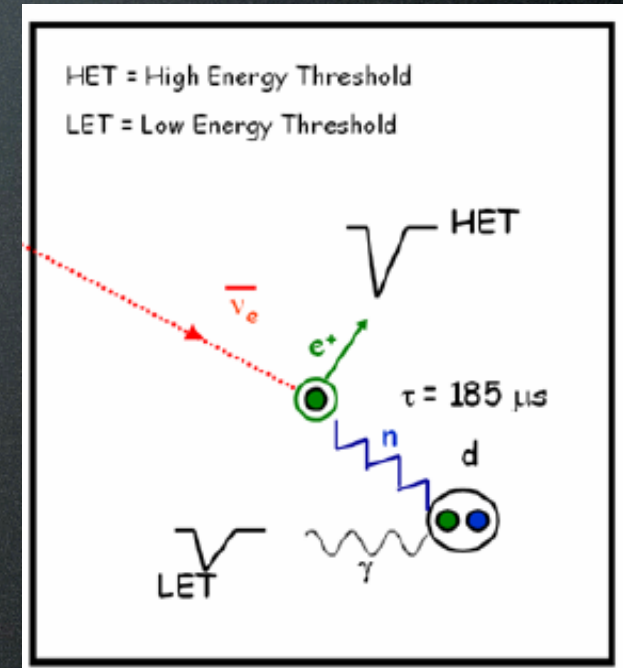
SN ν interactions in
heavy water
Čerenkov

SNO



SN ν interactions in
Liquid Scintillator

LVD, LENA, Borexino,
SNO+



The SN Early Warning System system

Distribution of alerts from experiments looking for SN burst
(currently LVD, SK, SNO, IceCube)

The combination of information allows to reduce the false alarm rate and possibly to make a triangulation between detectors. Triangulation is difficult because only

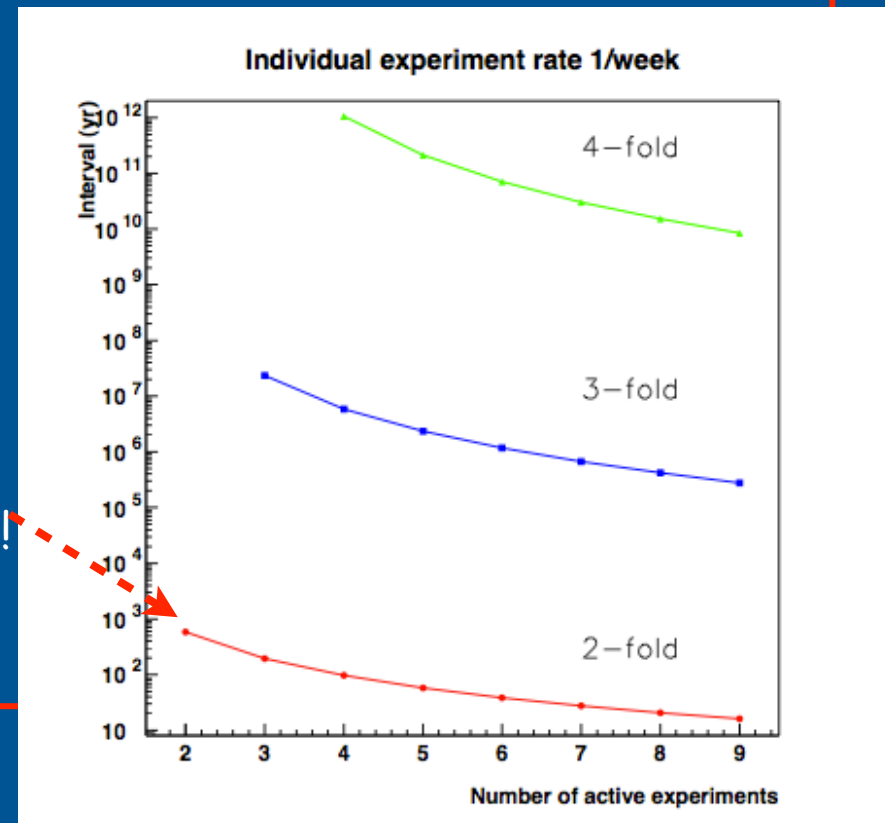
$\nu + e^- \rightarrow \nu + e^-$ has directional

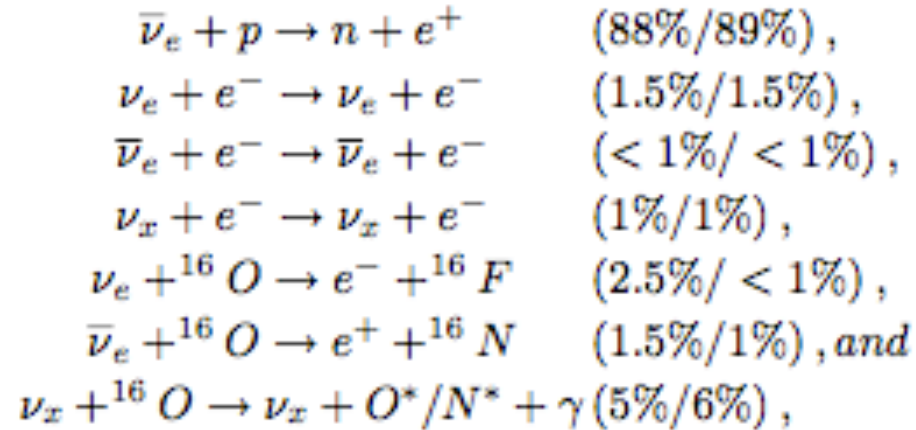
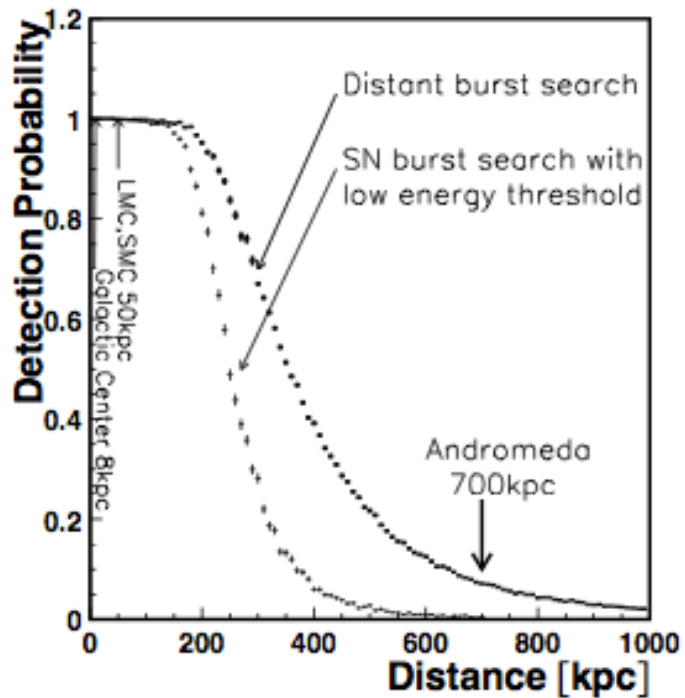
capabilities but it is not dominant in any detector (only in SK is 2% of signal).

Every experiment looks for SN burst and sends alarms at average rate of 1/week. The average interval of time between accidental alerts for a 2-fold coincidence in 10 s for 2 active experiments with this alarm rate is about 500 yrs!

Prompt information to astronomical community

<http://snews.bnl.gov/alert.html>





Signal from SN at 10 kpc: 10^4 events

100% efficient out to 100 kpc.

upper limit 0.32 SN yr^{-1} for 2589 d

IceCube: 2.6 yr^{-1} AMANDA 1997-2004

Baksan (best): 0.093 yr^{-1}

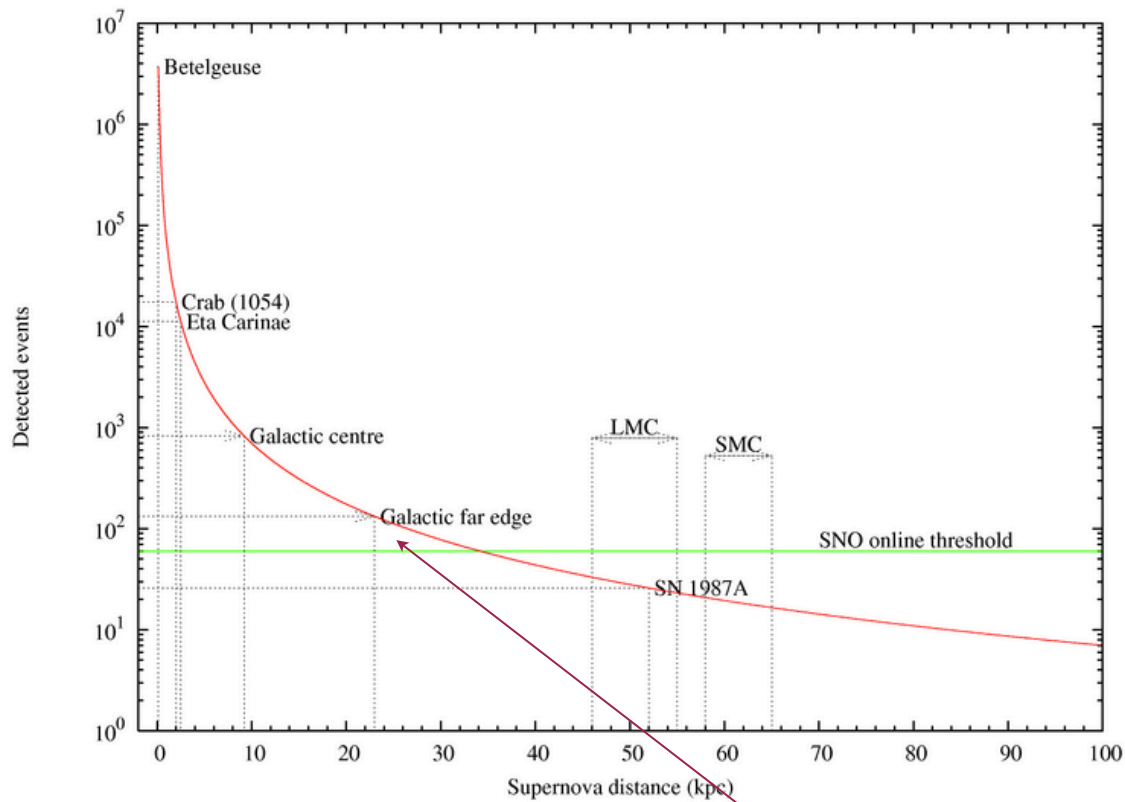
EGADS – Evaluating Gadolinium’s Action on Detector Systems: a dedicated Gd test facility of 200 ton built to explore feasibility of adding Gd to SK water for efficient anti-electron neutrino tagging.

[see Beacom and Vagins, *Phys. Rev. Lett.*, **93**:171101, 2004]

The cross section for thermal neutron capture on natural Gd is 49,000 barns, compared to 0.3 barns on free protons. With 0.2% GdCl_3 in water, 90% of n captures are on Gd with gamma decay of 8 MeV, 0.2% on Cl with gamma of 8.6 MeV and the rest on H (gamma 2.2 MeV not detectable in SK)

SNO

14 kt, 1 kt of D₂O, H₂O



Type	Reaction	Number of events	Number of counts (MC)		
			D ₂ O	Salt	NCD
H ₂ O $\bar{\nu}\bar{\nu}$	$\bar{\nu}_e + p \rightarrow n + e^+$	356 e ⁺	331		
D ₂ O CC	$\nu_e + d \rightarrow p + p + e^-$	83 e ⁻	72		
D ₂ O $\bar{C}\bar{C}$	$\bar{\nu}_e + d \rightarrow n + n + e^+$	106 n + 53 e ⁺	82	138	90
D ₂ O NC	$\nu'_{e\mu\tau} + d \rightarrow \nu'_{e\mu\tau} + p + n$	264 n	84	226	151
ES	$\nu'_{e\mu\tau} + e^- \rightarrow \nu'_{e\mu\tau} + e^-$	47 e ⁻	36		
Total:		909	605	803	680

IceCube Lab

50 m

IceTop

81 Stations, each with
2 IceTop Cherenkov detector tanks
2 optical sensors per tank
324 optical sensors

IceCube Array

86 strings including 8 DeepCore strings
60 optical sensors on each string
5160 optical sensors

December, 2010: Project completed, 86 strings

1450 m

DeepCore

8 strings-spacing optimized for lower energies
480 optical sensors



Eiffel Tower
324 m

86 strings

(8 Deep Core)

81 IceTop stations

5484 optical sensors

1 billion tons of ice

Bedrock



IceCube

Halzen, Jacobsen & Zas
astro-ph/9512080

5160 Optical Modules between 1450-2450 m -43 / -20 °C with dark count rates of about 540 Hz (approximately 16 Hz from muons, 10 Hz thermal emission from photocathode, 30 Hz afterpulses and the rest is a radioactive decays of ^{40}K , U, Th). The rate is reduced enforcing an artificial dead time after each count to 280 Hz at the cost of 13% dead time for signal.

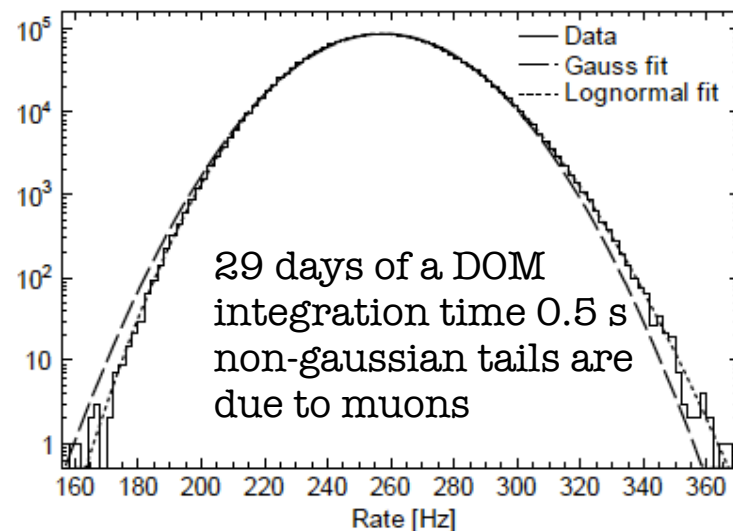
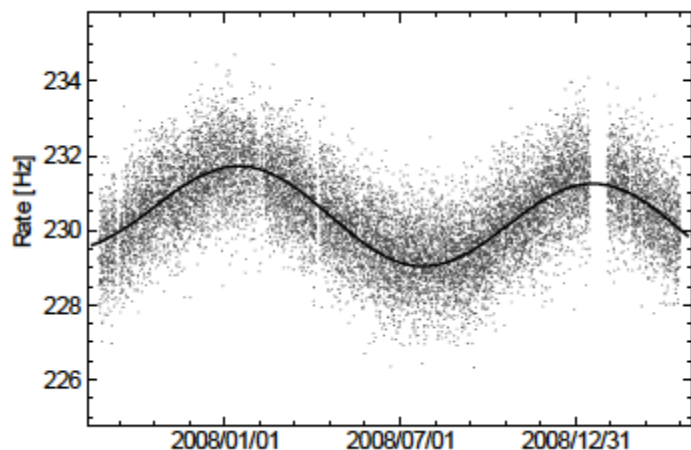
The contribution due to muons shows an annual modulation.

IceCube's advantages:

- Best detector for fine details in neutrino light curve for close supernovae
 - location far from other supernova detectors (triangulation, earth effect differences)

IceCube's disadvantages:

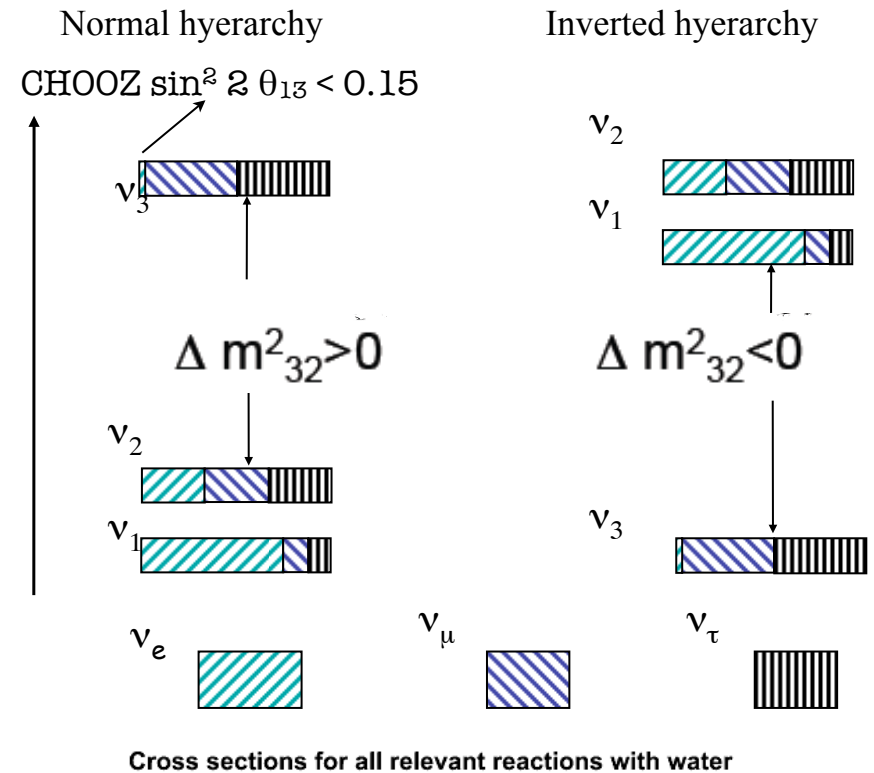
- no information on direction, neutrino type, energy, reach limited to 50 kpc
 - currently limited time resolution
 - limited sensitivity to ν_e



Signal rates and x-section for inverse beta and electron scattering

1 km³ of ice for Normal (NH) and Inverted (IH) Hierarchy for SN1987A at 10 kpc

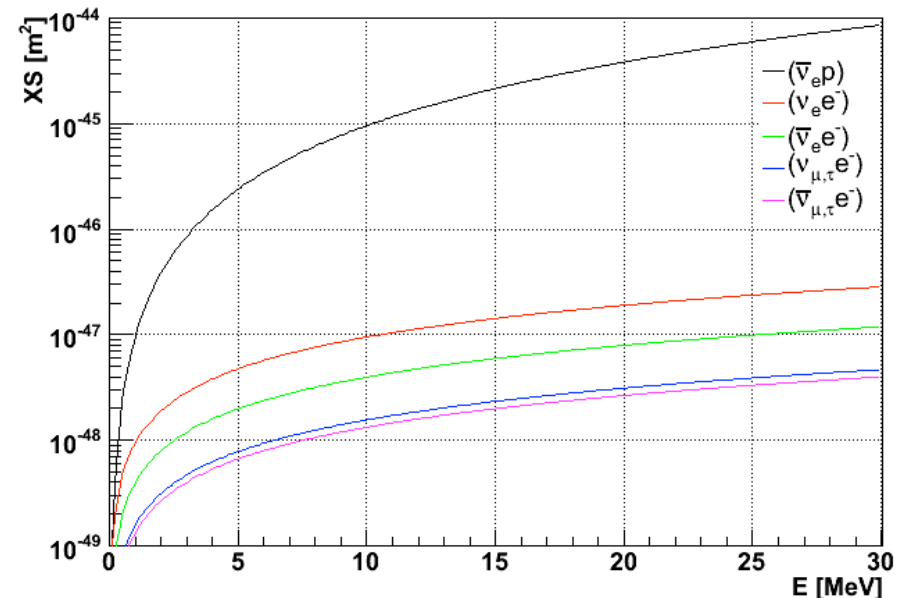
Reaction	# Targets	# Signal Hits	Signal Fraction
$\bar{\nu}_e + p \rightarrow e^+ + n$	$6 \cdot 10^{37}$	134 k (157 k)	93.8 % (94.4 %)
$\nu_e + e^- \rightarrow \nu_e + e^-$	$3 \cdot 10^{38}$	2.35 k (2.25 k)	1.7 % (1.4 %)
$\bar{\nu}_e + e^- \rightarrow \bar{\nu}_e + e^-$	$3 \cdot 10^{38}$	660 (720)	0.5 % (0.4 %)
$\nu_{\mu+\tau} + e^- \rightarrow \nu_{\mu+\tau} + e^-$	$3 \cdot 10^{38}$	700 (720)	0.5 % (0.4 %)
$\bar{\nu}_{\mu+\tau} + e^- \rightarrow \bar{\nu}_{\mu+\tau} + e^-$	$3 \cdot 10^{38}$	600 (570)	0.4 % (0.4 %)
$\nu_e + {}^{16}\text{O} \rightarrow e^- + \text{X}$	$3 \cdot 10^{37}$	2.15 k (1.50 k)	1.5 % (0.9 %)
$\bar{\nu}_e + {}^{16}\text{O} \rightarrow e^+ + \text{X}$	$3 \cdot 10^{37}$	1.90 k (2.80 k)	1.3 % (1.7 %)
$\nu_{\text{all}} + {}^{16}\text{O} \rightarrow \nu_{\text{all}} + \text{X}$	$3 \cdot 10^{37}$	430 (410)	0.3 % (0.3 %)
$\nu_e + {}^{17/18}\text{O}/{}^2_1\text{H} \rightarrow e^- + \text{X}$	$6 \cdot 10^{34}$	270 (245)	0.2 % (0.2 %)



Inverse beta on p dominates

$$E_{\text{th}} = (m_n - m_p + m_e) c^2 = 1.8 \text{ MeV}$$

Reactions on Oxygen have
 $E_{\text{th}} > 10 \text{ MeV}$.



Rate

$$R(t) = \epsilon_{\text{deadtime}} \frac{n_{\text{target}} L_{\text{SN}}^{\nu}(t)}{4\pi d^2 \bar{E}_{\nu}(t)} \int_0^{\infty} dE_e \int_0^{\infty} dE_{\nu} \\ \times \frac{d\sigma}{dE_e}(E_e, E_{\nu}) N_{\gamma}(E_e) V_{\gamma}^{\text{eff}} f(E_{\nu}, \bar{E}_{\nu}, \alpha_{\nu}, t),$$

number of radiated Cherenkov photons $\approx 188 E_e$ in 300-600 nm

Mean positron track length with $v > c/n$

$$\bar{x}_e(E_e) = (0.579 \pm 0.017) \text{cm} \cdot E_e / \text{MeV}$$

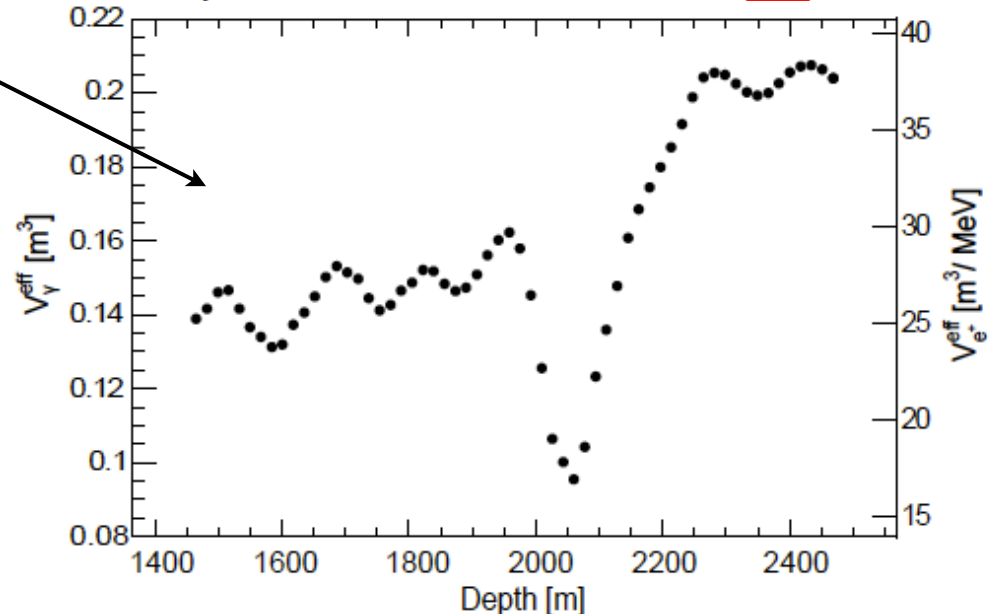
$$\frac{d^2}{d\lambda dx} N_{\gamma C} = \frac{2\pi\alpha z^2}{\lambda^2} \left(1 - \frac{1}{\beta^2 n_{\text{ice}}^2}\right)$$

$$\bar{N}_{\gamma C} = 2\pi\alpha \left(1 - \frac{1}{\beta^2 n_{\text{ice}}^2}\right) \bar{x}_e \int \frac{1}{\lambda^2} d\lambda = 325.4 \frac{\bar{x}_e}{\text{cm}} = (188.3 \pm 5.5) \cdot E_e / \text{MeV}$$

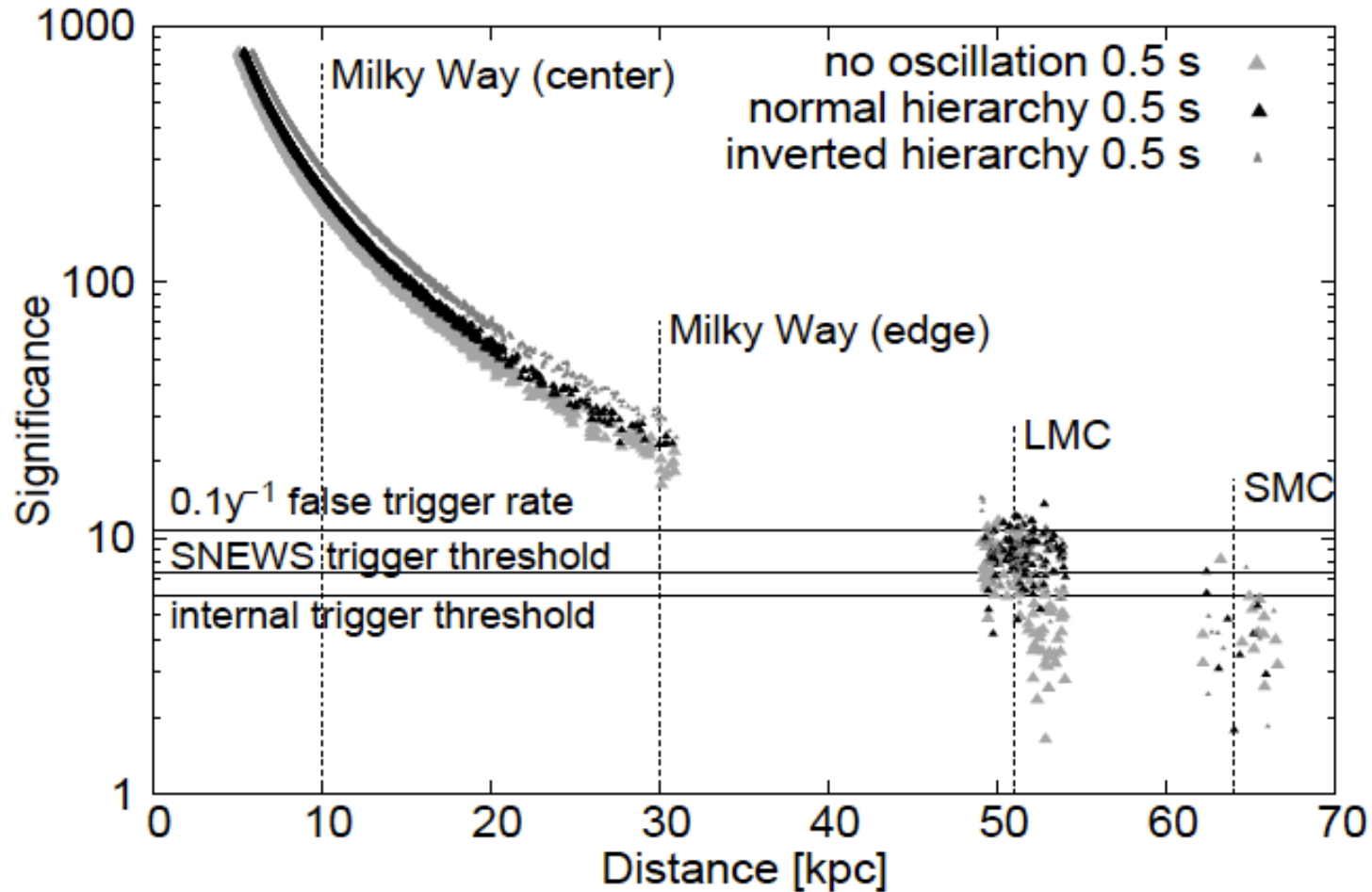
Effective Volume/PMT per single photon detection and for positrons

Average over all DOMs

$$V_{\gamma}^{\text{eff}} = 0.160 \pm 0.004 (\text{stat.}) \pm 0.020 (\text{syst.}) \text{m}^3$$



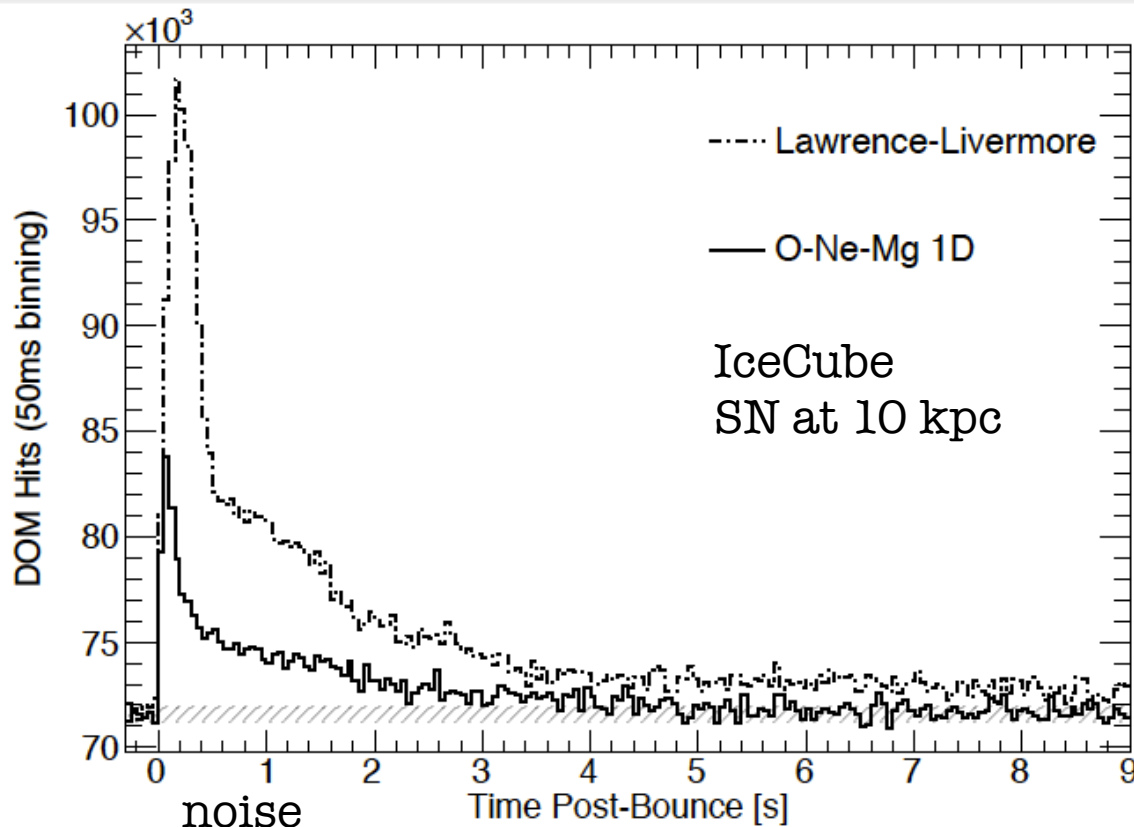
Significance



The significances are increased by neutrino oscillations in the star by typically 15% in case of a normal hierarchy and 40% in case of an inverted hierarchy

Astrophysics: SN core

Neutrinos can distinguish between a SN with iron core (Fe-core SN) and one with an oxygen-neon-magnesium core (ONeMg-core or low mass core SN), for which the density distribution and the shock propagation are completely different. Stars with 8-10 M_{Sun} are expected to develop ONeMg cores, which may undergo gravitational collapse before Ne ignition due to rapid electron captures on Mg and Ne (electron-capture supernovae)

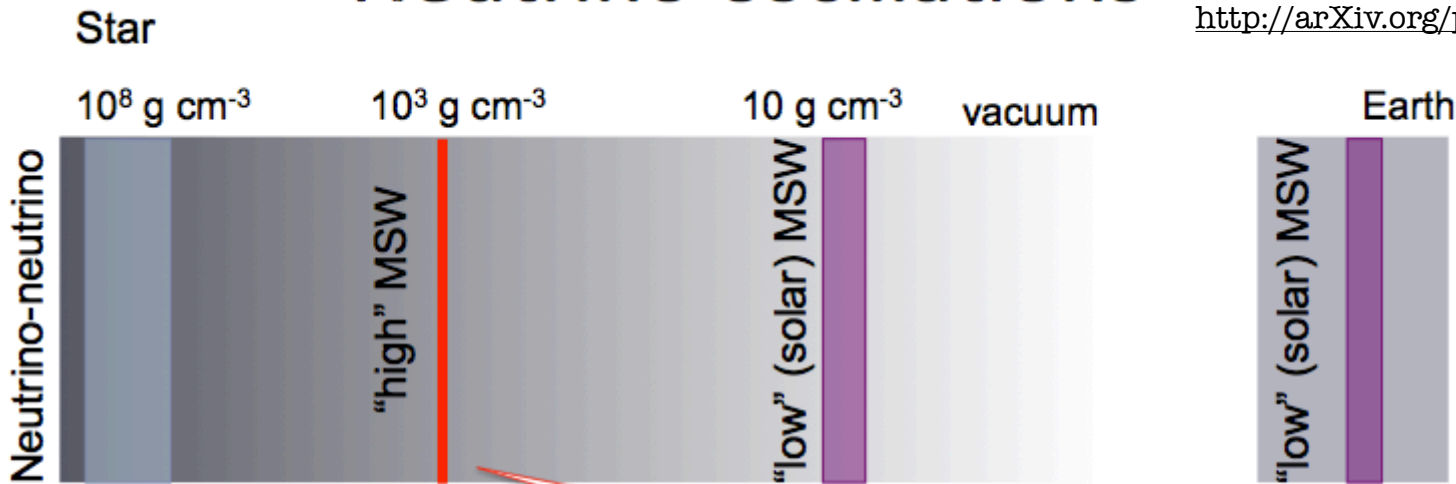


Models: Garching 8.8 solar masses
ONeMg core – arXiv:0912.0260 0.3M
neutrinos
Reference: Lawrence-Livermore
("classic" model with 20 solar masses)
1M neutrinos

IceCube
SN at 10 kpc

Neutrino oscillations

<http://arXiv.org/pdf/hep-ph/9907423>



- Matter effects:

*Unique of
supernovae!*

Given what we know of mixing parameters,

$$\sin^2 \theta_{12} = 0.32, \Delta m_{21}^2 = 8 \times 10^{-5} \text{ eV}^2, |\Delta m_{31}^2| = 3 \cdot 10^{-3} \text{ eV}^2, \sin \theta_{13} \leq 2 \cdot 10^{-2}$$

neutrinos encounter the MSW resonant density in the star that produce a distortion of spectra and reduction of the neutronization peak. The effect depends on the sign of

Δm_{31}^2 and the value of θ_{13} (we ignore δ_{CP})

$$F_e = pF_e^0 + (1 - p)F_x^0$$

$$F_{\bar{e}} = \bar{p}F_{\bar{e}}^0 + (1 - \bar{p})F_{\bar{x}}^0$$

\underline{p} = survival probability of neutrinos
 $\underline{\bar{p}}$ = surv. prob. of anti-neutrinos

RESONANCE

#electrons/nucleon

$$\Delta m^2 \cos 2\theta = 2\sqrt{2}G_F E N_e$$

$$N_e = \frac{\rho Y_e}{m_N}$$

Density of resonance layer nucleon mass

$$\rho_{res} \approx \frac{1}{2\sqrt{2}G_F} \frac{\Delta m^2}{E} \frac{m_N}{Y_e} \cos 2\theta$$

$$\rho_{res} \sim 1.4 \times 10^6 \text{ g/cc} \left(\frac{\Delta m^2}{1 \text{ eV}^2}\right) \left(\frac{10 \text{ MeV}}{E}\right) \left(\frac{0.5}{Y_e}\right) \cos 2\theta$$

H resonance

In a SN there are two resonance layers, one at higher densities corresponding to atmospheric neutrinos $\Delta m^2_{atm} \sim 2.5 \times 10^{-3} \text{ eV}^2$

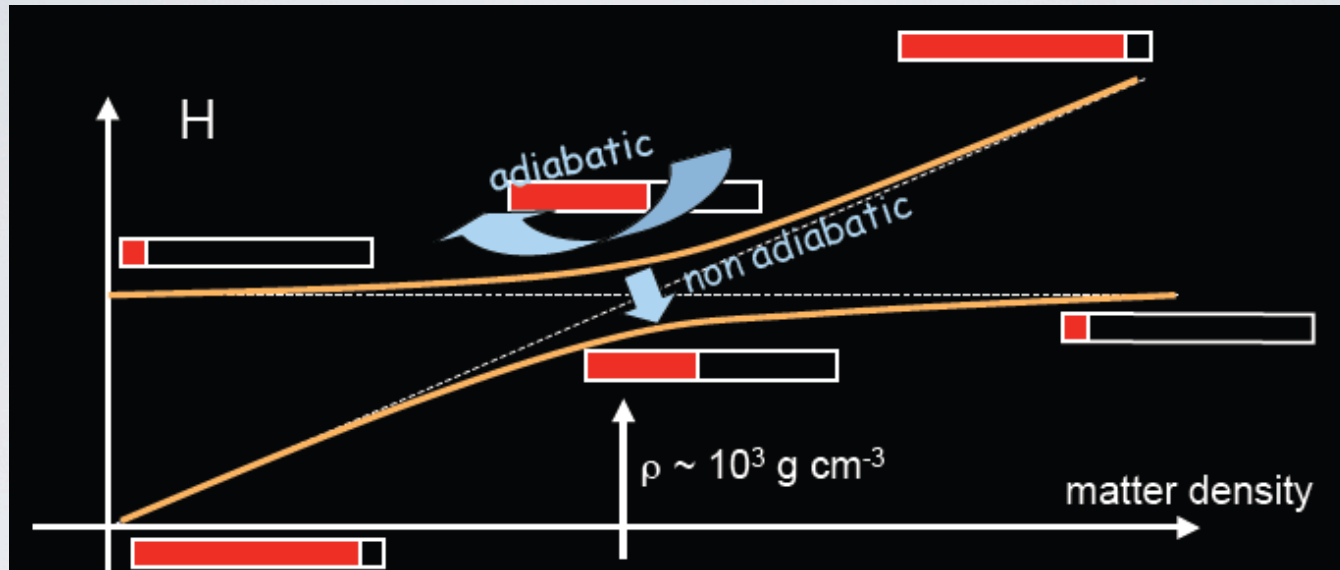
$$\rho_H \sim 10^3 - 10^4 \text{ g/cc}$$

and the other at lower densities (far from the core) corresponding to Δm^2_{sol}

$$\rho_L = \begin{cases} 5 - 15 & \text{g/cc} & \text{for SMA} \\ 10 - 30 & \text{g/cc} & \text{for LMA} \\ < 10^{-4} & \text{g/cc} & \text{for VO} \end{cases}$$

L resonance

ADIABATICITY



when the matter density changes the adiabaticity parameter governs the dynamics of the conversion

$$\gamma \equiv \frac{\Delta m^2 \sin^2 2\theta}{2E \cos 2\theta} \frac{1}{(1/N_e)(dN_e/dr)}$$

Jumping prob. that a neutrino in one matter eigenstate jumps to another.

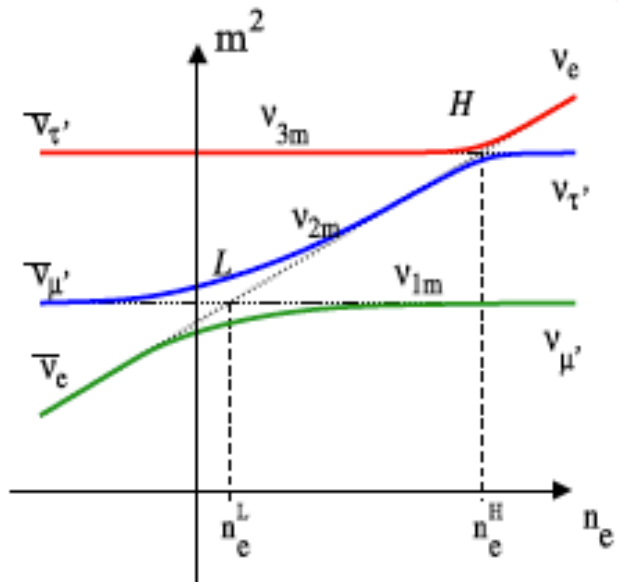
$$P_f = \exp\left(-\frac{\pi}{2}\gamma\right)$$

$\gamma \gg 1$ $P_f = 0$ adiabatic resonance: if the density changes slowly \Rightarrow strong flavor conversion (neutrino stays in mass state and flavor conversion occurs)

$\gamma \ll 1$ $P_f = 1$ non adiabatic, little flavor conversion

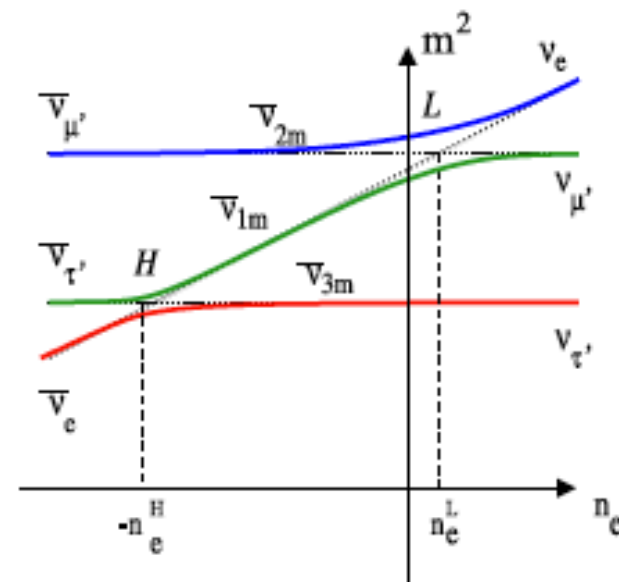
Inside a SN

Normal mass ordering



negative N_e is to show the case for anti-neutrinos

Inverted mass ordering



AD, A.Smirnov, PRD62, 033007 (2000)

H resonance: $(\Delta m_{\text{atm}}^2, \theta_{13}), \rho \sim 10^3 - 10^4 \text{ g/cc}$

- In $\nu(\bar{\nu})$ for normal (inverted) hierarchy
- Adiabatic (non-adiabatic) for $\sin^2 \theta_{13} \gtrsim 10^{-3} (\lesssim 10^{-5})$

L resonance: $(\Delta m_{\odot}^2, \theta_{\odot}), \rho \sim 10 - 100 \text{ g/cc}$

- Always adiabatic, always in ν

The H resonance ‘jumping probability’ is:

Inside a SN

$$P_H = \exp \left[- \left(\frac{E_{na}}{E} \right)^{2/3} \right],$$

$$E_{na} \simeq 1.08 \cdot 10^7 \text{ MeV} \left(\frac{|\Delta m_{32}^2|}{10^{-3} \text{ eV}^2} \right) C^{1/2} \sin^3 \theta_{13}$$

for the SN density profile $\rho \propto r^{-3}$:

$$\rho(r) = 10^{13} C \left(\frac{10 \text{ km}}{r} \right)^3 \text{ g} \cdot \text{cm}^{-3}$$

1-15 \rightarrow

- for θ_{13} small $\Rightarrow P_H \sim 1$ and for both mass hierarchy

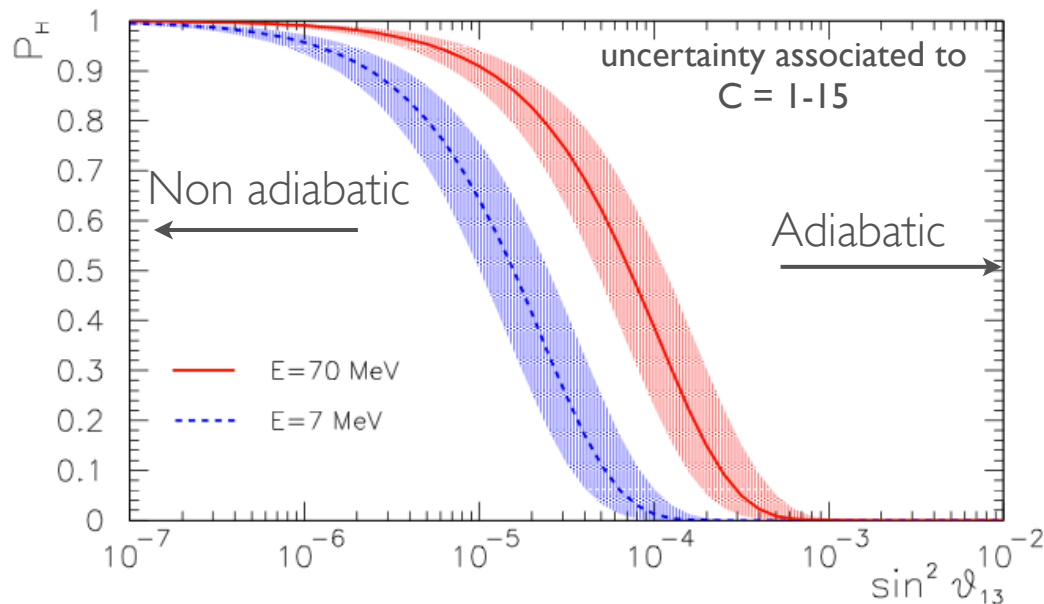
$$p \simeq \sin^2 \theta_{12}, \quad \bar{p} \simeq \cos^2 \theta_{12}$$

- for θ_{13} large $\Rightarrow P_H \sim 0$ (adiabatic conversion)

$$p \simeq \sin^2 \theta_{13}, \quad \bar{p} \simeq \cos^2 \theta_{12} (1 - \sin^2 \theta_{13}) \approx \cos^2 \theta_{12} \quad (\text{n.h.}),$$

and

$$p \simeq \sin^2 \theta_{12} (1 - \sin^2 \theta_{13}) \approx \sin^2 \theta_{12}, \quad \bar{p} \simeq \sin^2 \theta_{13} \quad (\text{i.h.})$$



sensitivity to θ_{13} to about 10^{-5} !

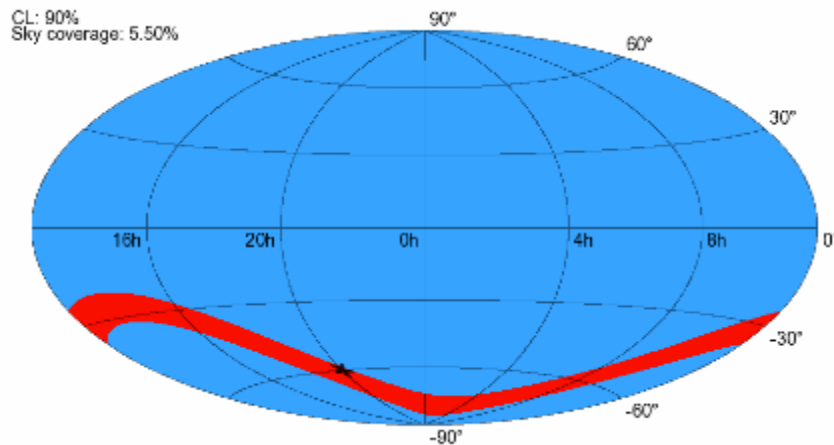
Matter oscillations in the Earth and triangulation

Scholberg et al., <http://arxiv.org/pdf/0910.3174v1>

Earth matter effects can be identified comparing signals in detectors, when the shadowing by the earth in different. Since the frequency of oscillations in energy depends on the pathlength traveled by the neutrinos in the Earth, an observed energy spectrum contains information about the SN direction.

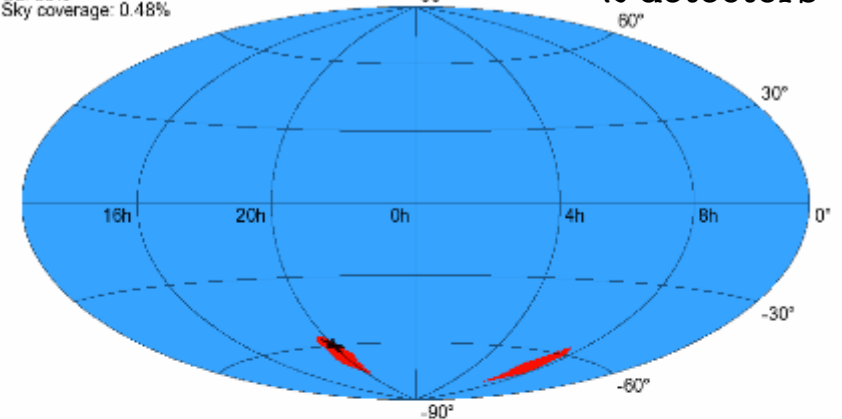
90%CL allowed region for SN position indicated by the black star

1 detector with good energy resolution

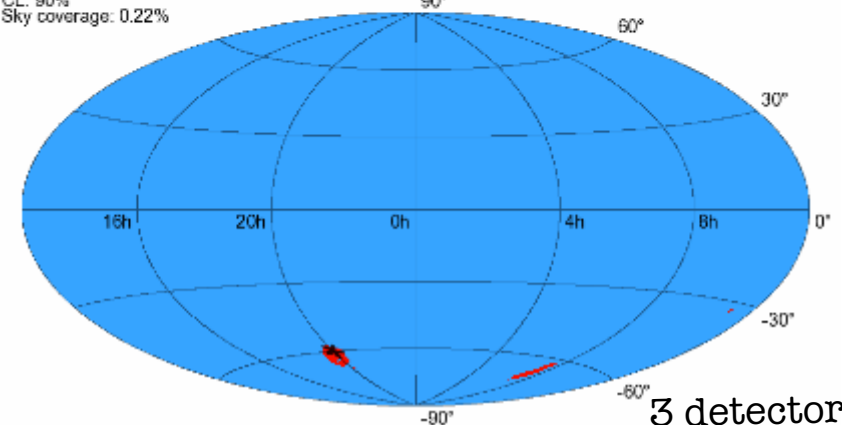


CL: 90%
Sky coverage: 0.48%

2 detectors



CL: 90%
Sky coverage: 0.22%



3 detectors

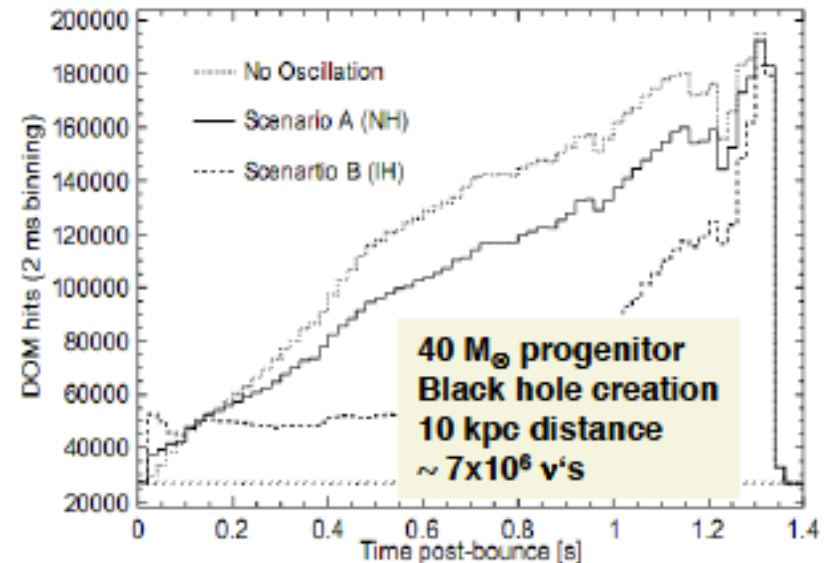
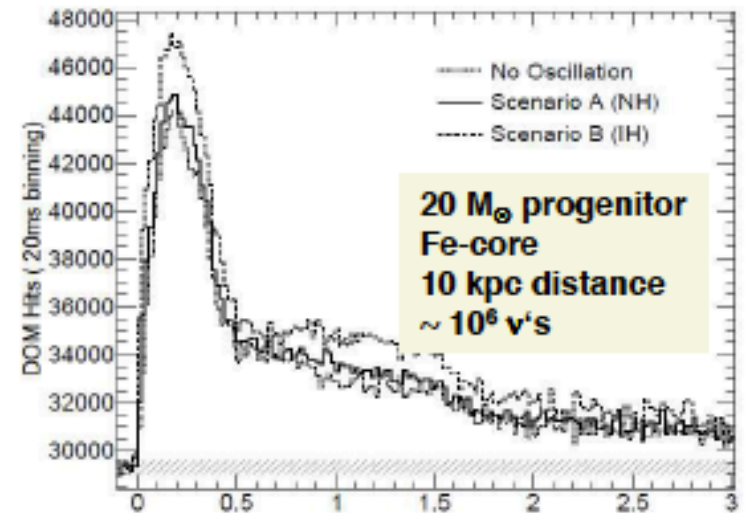
Neutrino properties

- learn about neutrino hierarchy and θ_{13} mixing angle from, various forms of ν -oscillations

MSW-oscillation in star
and Earth

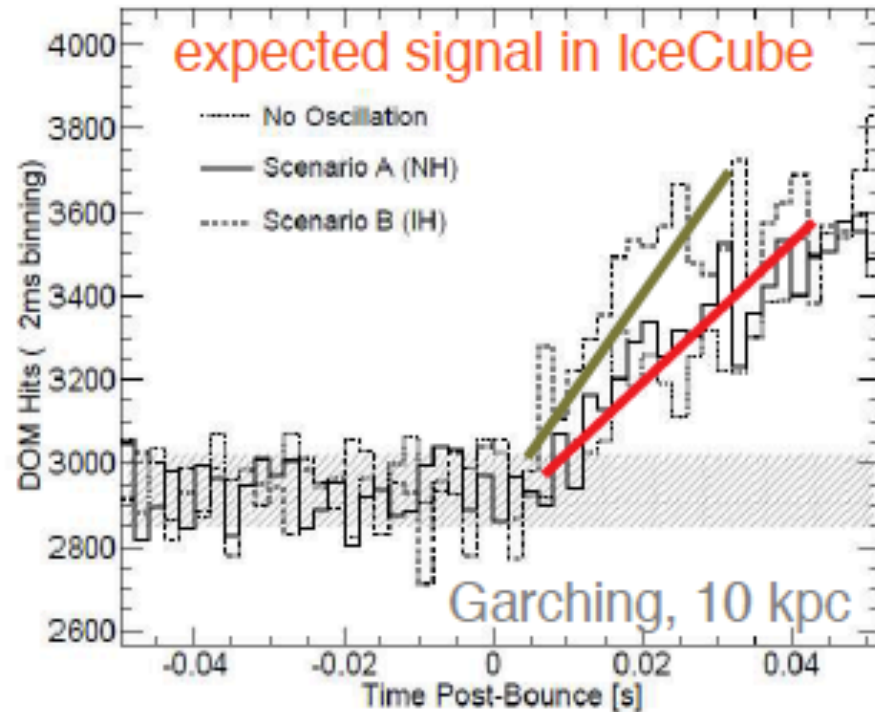
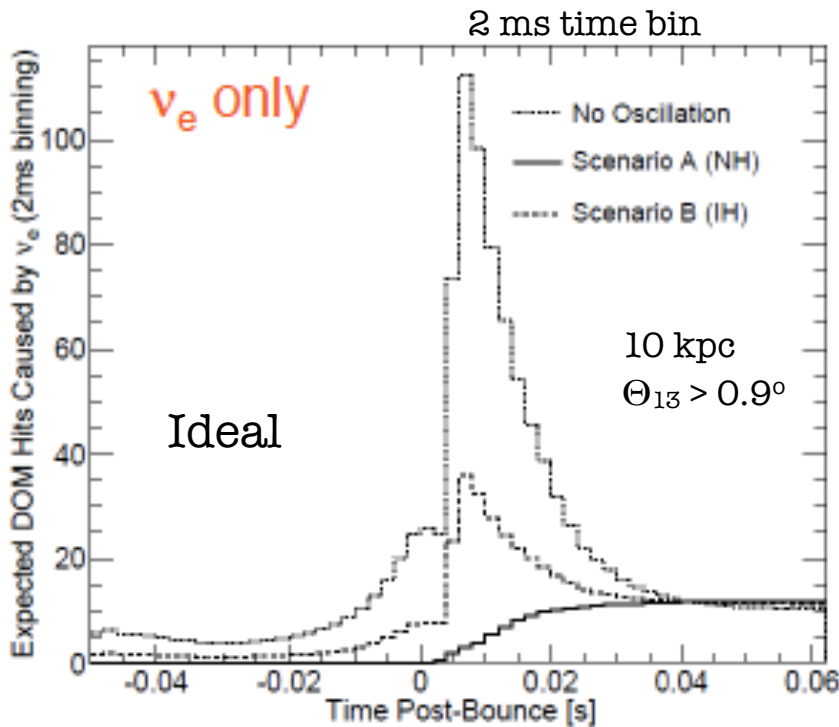
Large shape differences between
Normal and inverted hierarchies

IceCube simulation following Totani et al. (1997)



Deleptonization peak and oscillations (< 10 ms)

Detected by the elastic scattering reaction dominantly $\nu_e + e^- \rightarrow \nu_e + e^-$
smaller cross section x number of targets by 50 compared to inverse- β

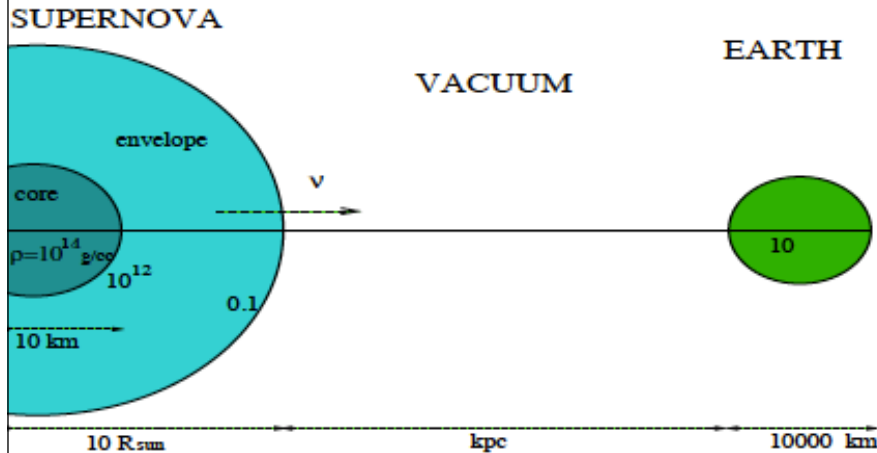


With noise: features are washed out but

signal visible up to a couple kpc distance ... however, rising slope bears information on ν hierarchy (3σ up to 6 kpc)

Collective effects and MSW resonance

Dighe et al, <http://arXiv.org/pdf/0904.3542>



Inside the SN: *flavour conversion*

Collective effects and MSW matter effects

Between the SN and Earth: *no flavour conversion*

Mass eigenstates travel independently

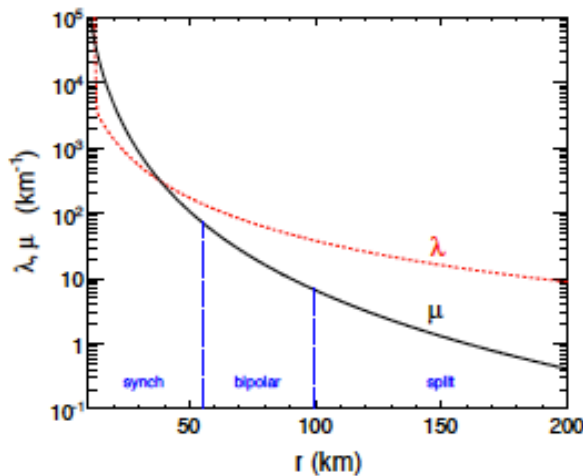
Inside the Earth: *flavour conversion*

MSW matter effects (*if detector is on the other side*)

Collective effects (non-linear ν oscillations + neutrinos-neutrinos interactions) produce exchanges of ν_e (anti- ν_e) spectra with ν_x (anti- ν_x) in some energy ranges and introduce spectral distortions for NH and IH

- $\mu \equiv \sqrt{2} G_F (N_\nu + N_{\bar{\nu}})$

- $\lambda \equiv \sqrt{2} G_F N_e(r)$ MSW



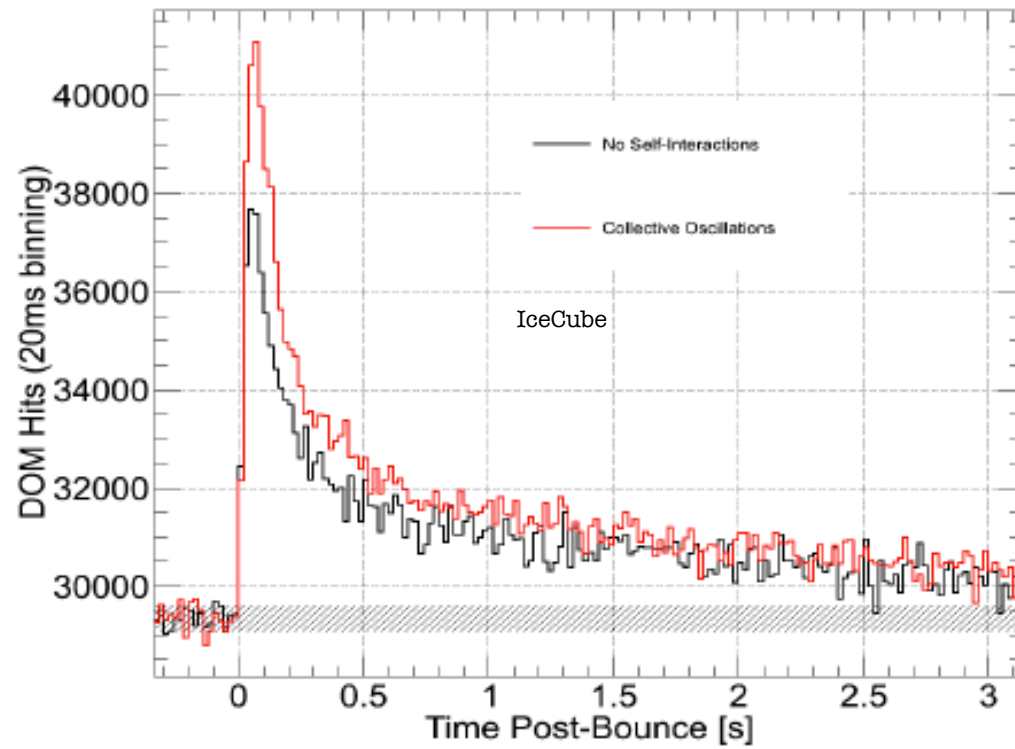
- $r \lesssim 200$ km: collective effects dominate

- $r \gtrsim 200$ km: standard MSW matter effects dominate

G.L.Fogli, E. Lisi, A. Marrone, A. Mirizzi, JCAP 0712, 010 (2007)

Neutrino self-interactions

Test model: Garching 8.8 solar masses ONeMg core



Future: LAGUNA

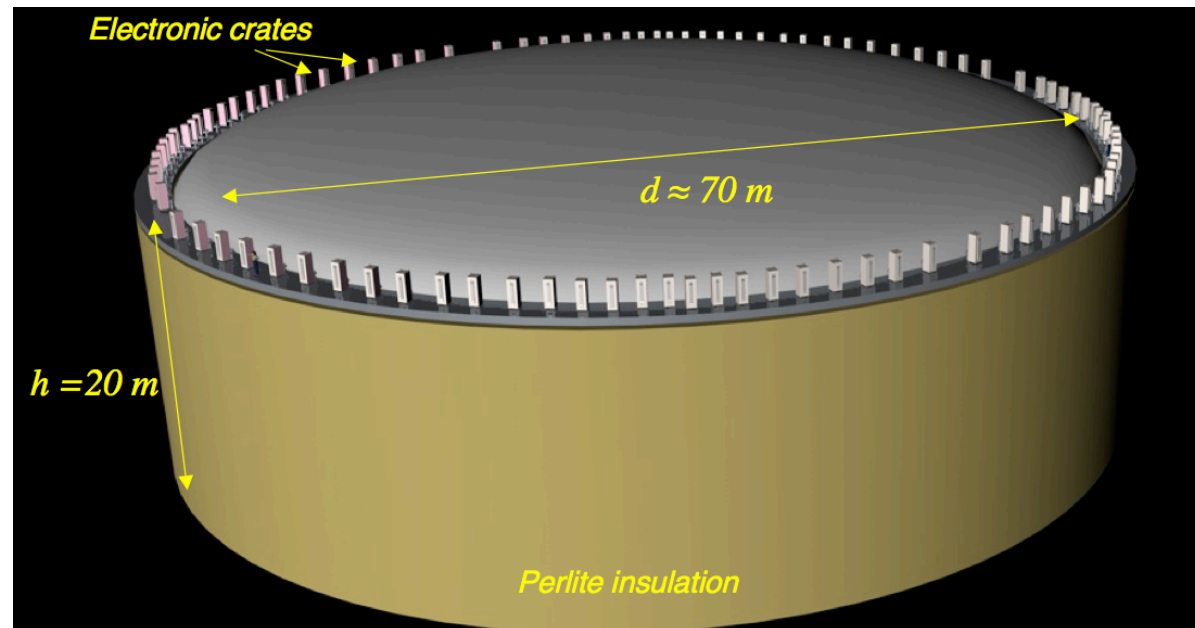
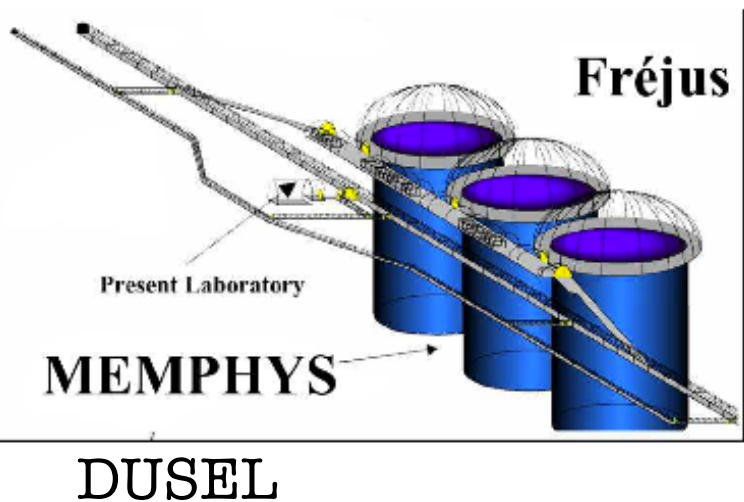
Order of 10,000 events from SN at 10 kpc

Glacier, LANND LAr TPC 100 kt

	CC
$\nu_e + {}^{40}\text{Ar} \rightarrow {}^{40}\text{K}^* + e^- + \text{gamma decay}$	2.5e4
$\bar{\nu}_e + {}^{40}\text{Ar} \rightarrow {}^{40}\text{Cl}^* + e^+$	540
$\nu_\alpha + {}^{40}\text{Ar} \rightarrow {}^{40}\text{Ar}^* + \nu_\alpha + \text{gamma decay}$	NC
	3e4
$\nu + e^- \rightarrow \nu + e^-$	ES
	1e3

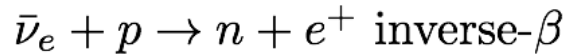
<http://arXiv.org/pdf/0705.0116v2>

WC 750 kt



Future: LENA or Hano Hano

Golden channel:



$n + p \rightarrow d + \gamma(2.2\text{MeV})$ after an average of $250 \mu\text{s}$



Reaction	Type	Events for $\langle E_\nu \rangle$ values		
		12 MeV	14 MeV	16 MeV
$\bar{\nu}_e p \rightarrow n e^+$	CC	1.1×10^4	1.3×10^4	1.5×10^4
$\nu p \rightarrow p \nu$	NC	1.3×10^3	2.6×10^3	4.4×10^3
$\nu e \rightarrow e \nu$	CC/NC	6.2×10^2	6.2×10^2	6.2×10^2
$\nu {}^{12}\text{C} \rightarrow {}^{12}\text{C}^* \nu$				
${}^{12}\text{C}^* \rightarrow {}^{12}\text{C} \gamma$	NC	6.0×10^2	1.0×10^3	1.5×10^3
$\bar{\nu}_e {}^{12}\text{C} \rightarrow {}^{12}\text{B} e^+$				
${}^{12}\text{B} \rightarrow {}^{12}\text{C} e^- \bar{\nu}_e$	CC	1.8×10^2	2.9×10^2	4.2×10^2
$\nu_e {}^{12}\text{C} \rightarrow {}^{12}\text{N} e^-$				
${}^{12}\text{N} \rightarrow {}^{12}\text{C} e^+ \nu_e$	CC	1.9×10^2	3.4×10^2	5.2×10^2

$E_{\text{tot}} = 3 \times 10^{53}$ erg for a SN at 10 kpc, equipartitioned among all neutrino species, and Maxwell-Boltzmann spectra with $\langle E_\nu \rangle = 12, 14$ and 16 MeV

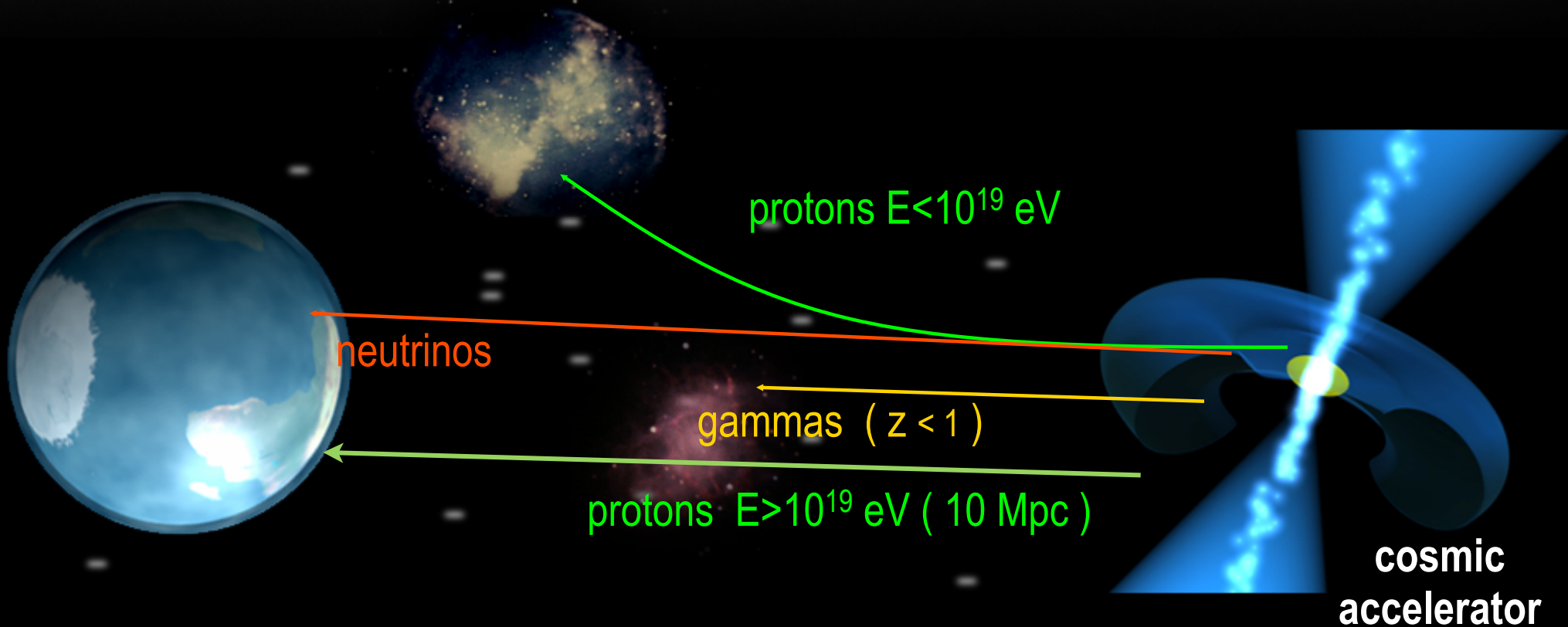
Three charged-current (CC) reactions measure ν_e and anti- ν_e fluxes and spectra while three neutral-current (NC) processes, sensitive to all flavors, give information on the total flux.

Selected topics on high energy neutrino astrophysics

- cosmic ray -gamma - neutrino connection
- GRBs as an instructive example
- UHE neutrinos
- Neutrino - gamma connection

MESSENGERS FROM THE UNIVERSE

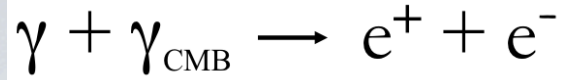
Discovery messengers: Neutrinos and Gravitational Waves



photons: absorbed on dust and radiation; reprocessed at source
protons/nuclei: deviated by magnetic fields, absorbed on radiation (GZK)

OBSERVABLE UNIVERSE

10^3 TeV photons barely reach us from the Galactic Centre

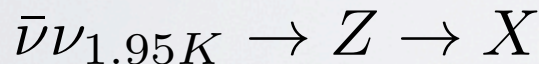


Proton horizon (GZK cut-off):



$$L_p = \frac{1}{\sigma_{p-\gamma_{\text{CMB}}} n_\gamma} \sim \frac{1}{10^{-28} \text{cm}^2 \times 400 \text{cm}^{-3}} \sim 10 \text{ Mpc}$$

The neutrino horizon is comparable to the universe!

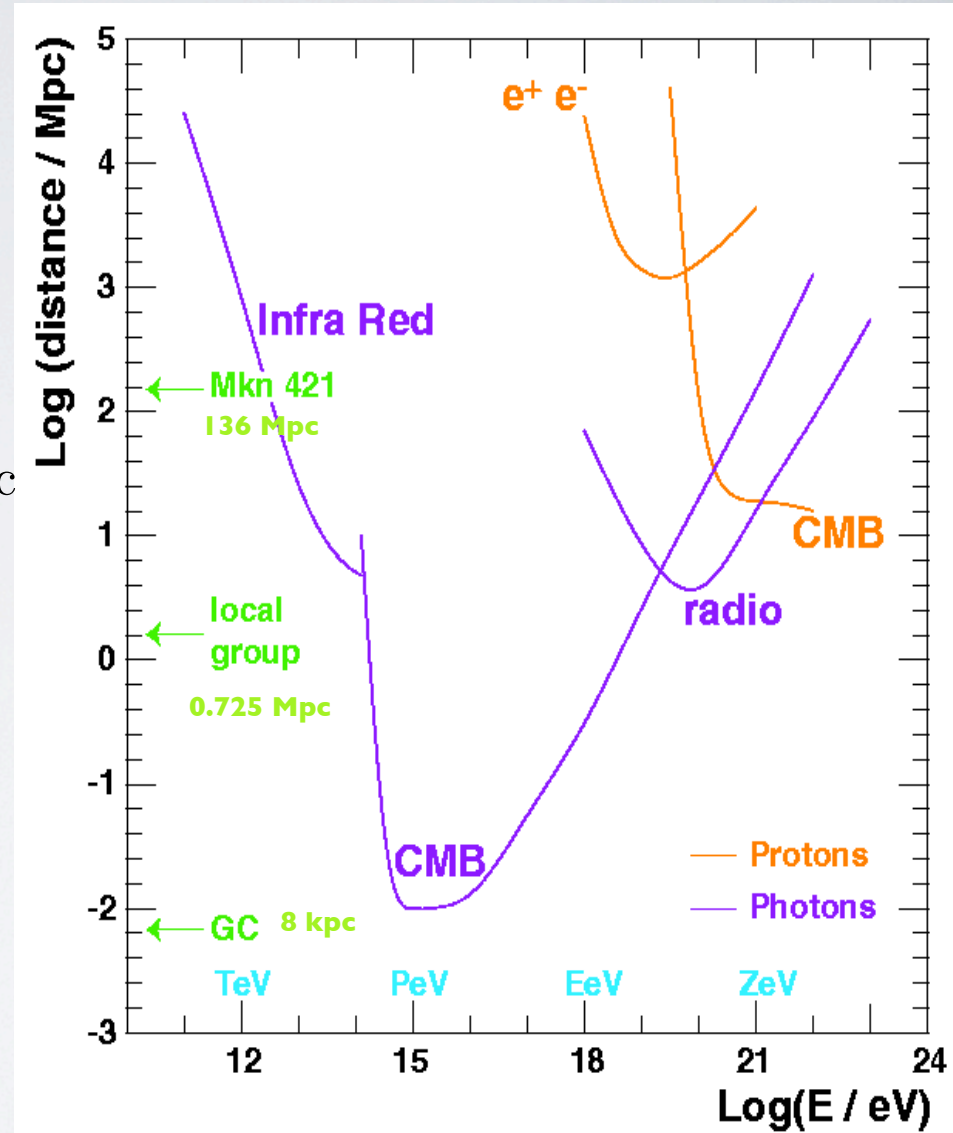


$$E_{\text{res}} = \frac{M_Z^2}{2m_\nu} \cong 4 \times 10^{21} \left(\frac{1\text{eV}}{m_\nu} \right) \text{eV}$$

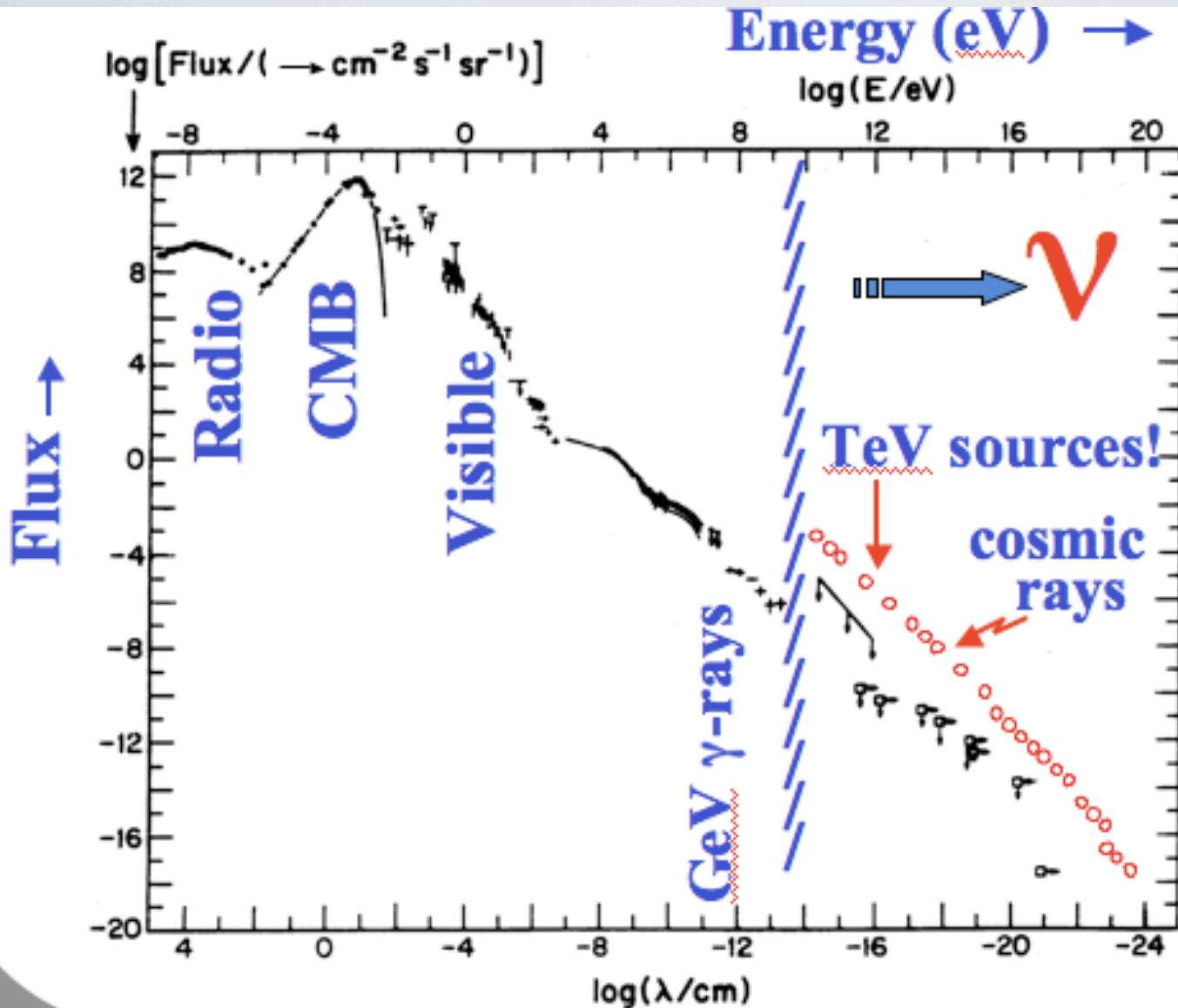
$$L_\nu = \frac{1}{\sigma_{\text{res}} \times n} = \frac{1}{5 \times 10^{31} \text{cm}^2 \times 112 \text{cm}^{-3}} \approx 6 \text{Gpc}$$

Particle horizon about 14 Gpc

1Mpc = 3.26 Mly = 3.1×10^{24} cm



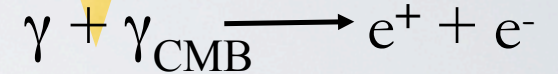
UNDERSTANDING ACCELERATION PROCESSES IN THE UNIVERSE



Gamma astronomy



$< 100 \text{ TeV}$

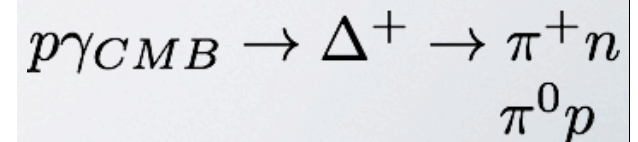


gamma are absorbed on backgrounds

Neutrino astronomy



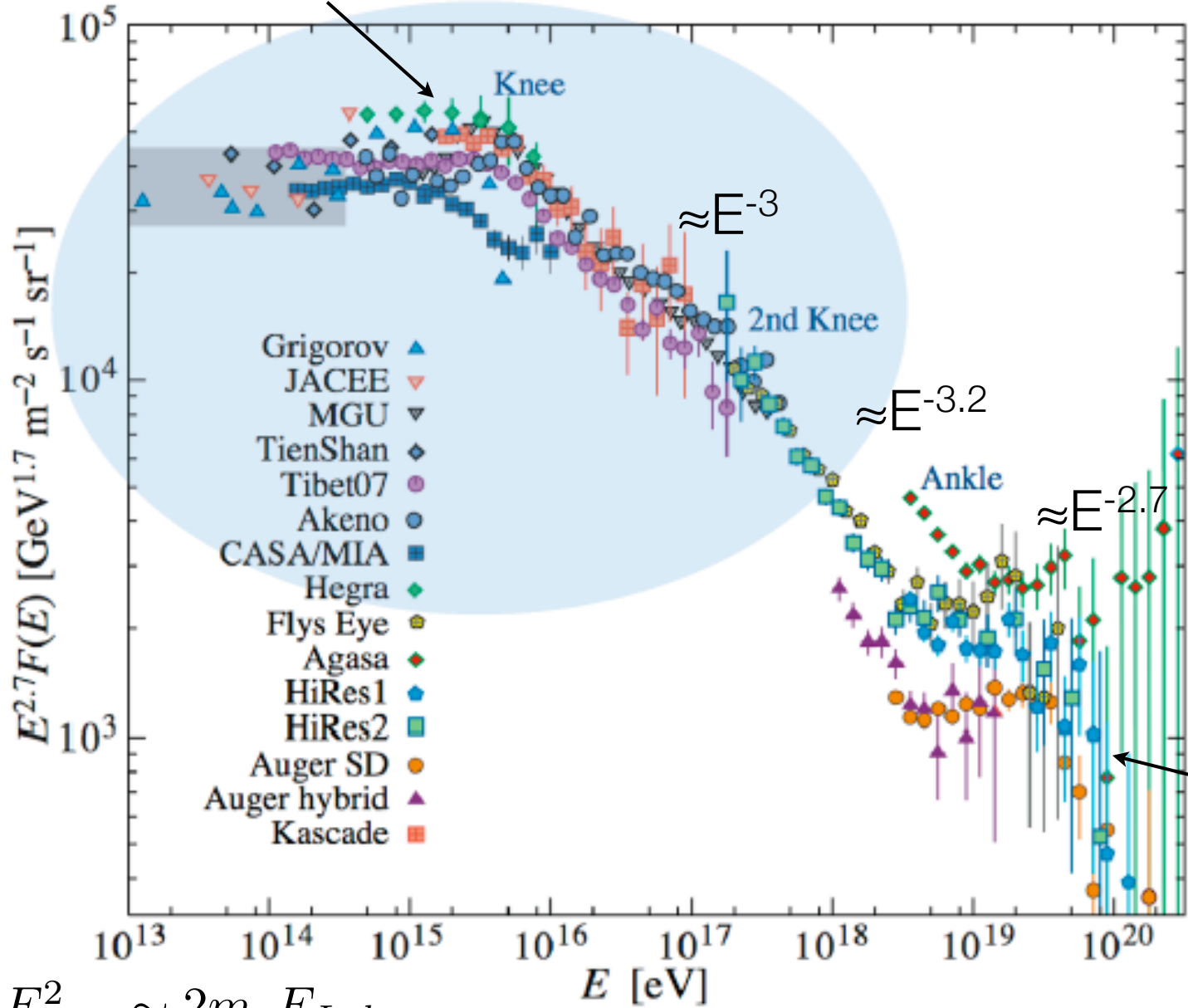
$> 10 \text{ EeV}$



Proton astronomy

Cosmic Ray Spectrum

Luminosity of the CR beam decreases steeply with energy



Galaxy containment ($B \sim 4 \mu\text{G}$)

$R \sim 400 \text{ pc @ } 10^{18} \text{ eV}$
 $R \sim 0.4 \text{ pc @ } 10^{15} \text{ eV}$

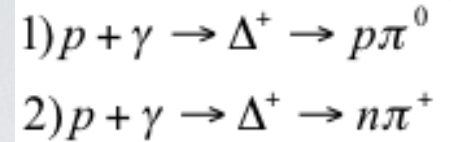
$$E_{CM}^2 \sim 2m_p E_{Lab}$$

For LHC: $E_{CM} = 14 \text{ TeV}$, $m_p \sim 1 \text{ GeV} \rightarrow E_{Lab} \sim 10^{17} \text{ eV}$

POWER OF SOURCES OF COSMIC RAYS

energy density flux = velocity x density

$$4\pi \int dE \left(E \frac{dN}{dE} \right) = c \rho_E$$



Waxman & Bahcall, PRD59, 1999 and PRD64, 2001)

Galactic

galactic CR: $\rho_E \sim 10^{-12} \text{ erg/cm}^3$
 Power needed: $\rho_E / \tau_{\text{esc}} \approx 10^{-26} \text{ erg/cm}^3 \text{ s}$
 $\tau_{\text{esc}} \approx 3 \times 10^6 \text{ yrs}$ escape time from Galaxy

10^{51} erg/SN every 30 years $\sim 10^{-25} \text{ erg/cm}^3 \text{ s}$
 for Galactic disk volume $\sim 10^{67} \text{ cm}^3$

10% of SN provides the environment and energy to explain the galactic CRs!

1934 Baade and Zwicky
Acc mechanism then proposed
by Fermi in 1949

Extragalactic

Above the ankle:

$$E \left\{ E \frac{dN_{\text{CR}}}{dE} \right\} = \frac{3 \times 10^{10} \text{ GeV}}{(10^{10} \text{ cm}^2)(3 \times 10^7 \text{ s}) \text{ sr}}$$

$$= 10^{-7} \text{ GeV cm}^{-2} \text{ s}^{-1} \text{ sr}^{-1}$$

Energy density in extra-galactic CRs:

$$\rho_E = \frac{4\pi}{c} \int_{E_{\text{min}}}^{E_{\text{max}}} \frac{10^{-7}}{E} dE \frac{\text{GeV}}{\text{cm}^3} \sim 3 \times 10^{-19} \frac{\text{erg}}{\text{cm}^3}$$

$$E_{\text{max}}/E_{\text{min}} \sim 10^3$$

Power needed by a population of sources of p with E^{-2} to generate ρ_E over the Hubble time = $10^{10} \text{ yrs} \approx 10^{44} \text{ erg Mpc}^{-3} \text{ yr}^{-1}$

- $3 \times 10^{39} \text{ erg/s}$ per galaxy
- $3 \times 10^{42} \text{ erg/s}$ per cluster of galaxies
- $2 \times 10^{44} \text{ erg/s}$ per AGN
- $2 \times 10^{52} \text{ erg}$ per cosmological GRB.

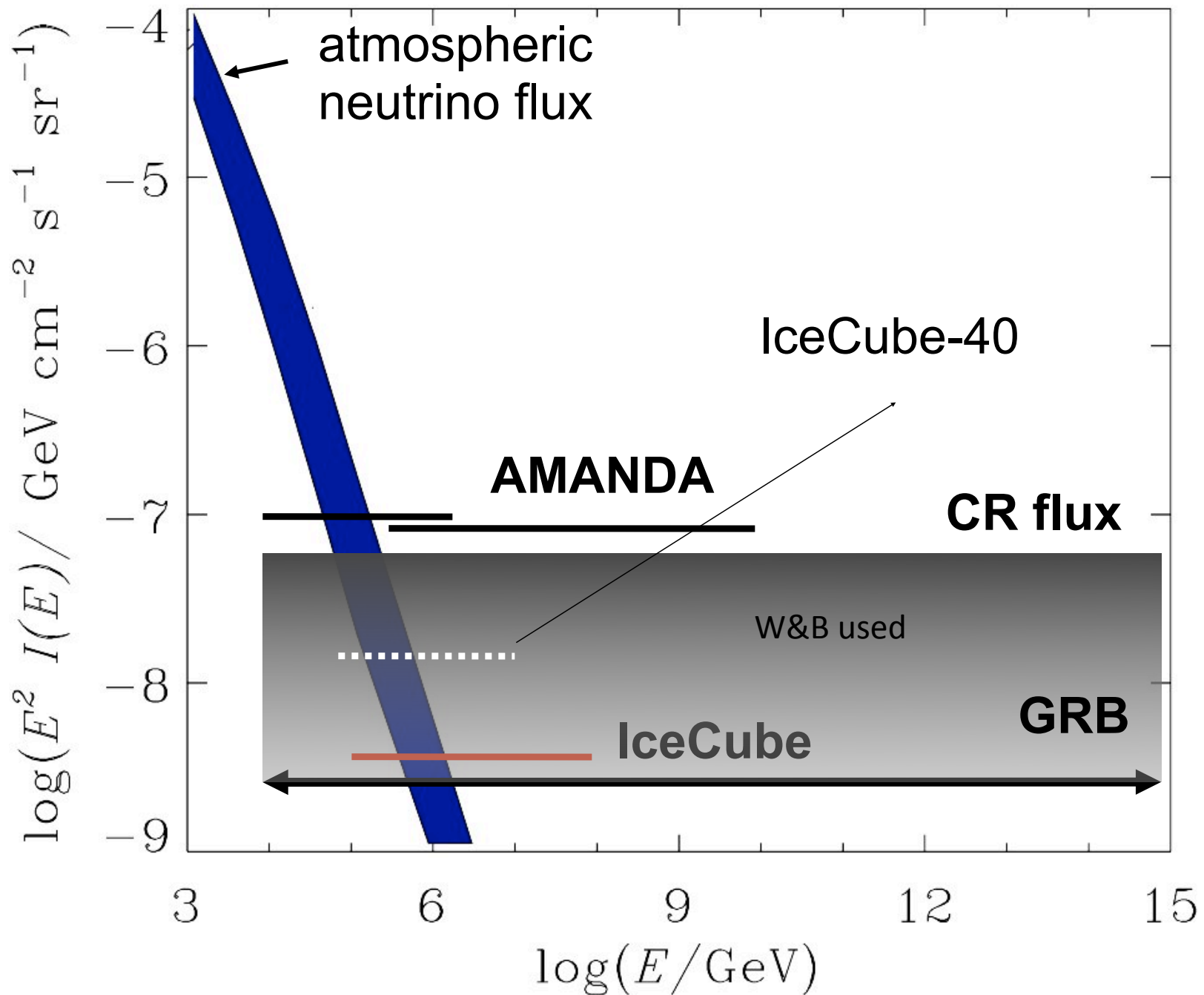
*300 GRB per Gigaparsec³ per year
for 10¹⁰ years (Hubble time)*

$$2 \times 10^{51} \text{ erg} \times \frac{300}{\text{Gpc}^3 \text{ yr}} \times 10^{10} \text{ yr} = 3 \times 10^{-19} \frac{\text{erg}}{\text{cm}^3}$$

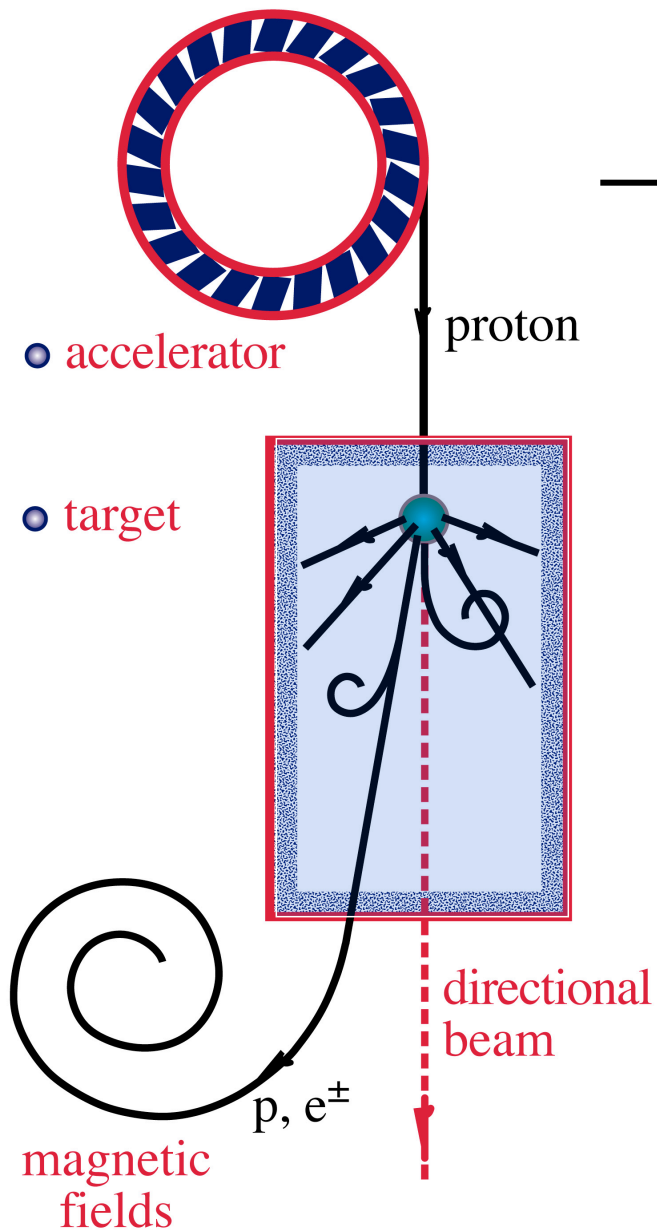
$$1 \text{ Gpc}^3 = 2.9 \times 10^{82} \text{ cm}^3 \quad \text{Hubble time} = 10^{10} \text{ years}$$

observed energy
density of extragalactic
CRs:
 $\sim 10^{-19} \text{ erg / cm}^3$

neutrinos associated with extragalactic cosmic rays



ν and γ beams : heaven and earth

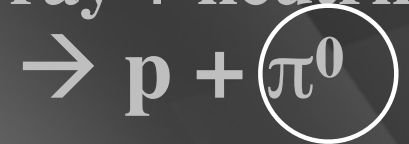


black hole

**radiation
and dust**



\sim cosmic ray + neutrino



\sim cosmic ray + gamma

Neutrino production: Threshold



IN CM

$$E'_p > \frac{m_\Delta^2 - m_p^2}{4E'_\gamma}$$

$$s = E_{CM}^2 = E_\gamma^2 + E_p^2 - 2 \times (E_p E_\gamma - p_p E_\gamma \cos \theta) = m_p^2 - 2E_p E_\gamma (1 - \cos \theta)$$

For opposite beams $\cos(\theta) = -1$. The threshold is calculated for final states at rest:
 $m_\Delta^2 = m_p^2 - 4E_p E_\gamma$

In the comoving proton frame
and in the observer frame

$$E_p > 1.4 \times 10^{16} \text{ eV} \left(\frac{\gamma}{300} \right)^2 \left(\frac{1 \text{ MeV}}{E_\gamma} \right) \quad \text{a gamma for each energy}$$

$$E_\nu = 1/4 \langle X_F \rangle E_p \simeq 1/20 E_p \simeq 700 \text{ TeV}$$

Fraction of proton energy in pions $\langle X_F \rangle \sim 0.2$ that is equally distributed in leptons. Flux in neutrinos from fraction of energy into a neutrino and total energy in CRs

$$\frac{dN_\nu}{dE} \sim n_{int} \times \frac{1}{x_\nu} \times \frac{dN_{CR}}{dE} \left(\frac{E_{CR}}{x_\nu} \right)$$

$$x_\nu \sim \frac{1}{4} x_F \sim \frac{1}{20}$$

$n_{int} \cong 1$ for GRB fireball

FRACTION OF PROTON ENERGY IN
NEUTRINOS, x_F p \rightarrow pion

The neutrino - proton connection

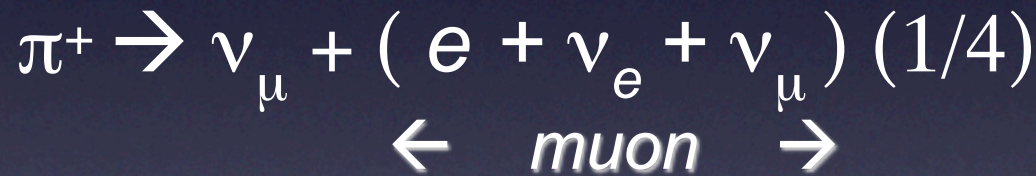
p-gamma interactions (require higher threshold than pp and ambient photons)

$$s = E_{CM}^2 = E_\gamma^2 + E_p^2 - 2 \times (E_p E_\gamma - p_p E_\gamma \cos \theta) = m_p^2 - 2E_p E_\gamma (1 - \cos \theta)$$

For opposite beams $\cos(\theta) = -1$. The threshold is calculated for final states at rest:

$$m_\Delta^2 = m_p^2 - 4E_p E_\gamma$$

Isospin conservation: Delta production produces pions that share 1/3 of p energy



$$\frac{dN_\nu}{dE_\nu} = 2 \times \frac{1}{3} \frac{1}{x_\nu} \frac{dN_p}{dE_p} (E_p)$$

$$\frac{dN_\nu}{dE_\gamma} = 2 \times \frac{2}{3} \frac{1}{x_\gamma} \frac{dN_p}{dE_p} (E_p)$$

	$p\pi^+$	$p\pi^0$	$p\pi^-$	$n\pi^+$	$n\pi^0$	$n\pi^-$
Δ^{++}	1					
Δ^+		2/3		1/3		
Δ^0			1/3		2/3	
Δ^-						1

$$x_\nu = \frac{E_\nu}{E_p} = \frac{1}{4} \langle x_F \rangle = \frac{1}{20}$$

$$x_\gamma = \frac{E_\gamma}{E_p} = \frac{1}{2} \langle x_F \rangle = \frac{1}{10}$$

$$dE_\nu = x_\nu dE_p$$

$$dE_{\nu,\gamma} = x_{\nu,\gamma} dE_p$$

The neutrino - gamma connection

Assuming an E^{-2} spectrum of protons we
have for the p spectrum

$$E_p = \frac{E_\nu}{x_\nu} \Rightarrow \frac{dN_p}{dE_p}(E_p) \propto E_\nu^{-2} x_\nu^2 \Rightarrow$$
$$\frac{dN_\nu}{dx_\nu} \propto 2 \times \frac{1}{3} \times \frac{1}{20} E_\nu^{-2}$$

$$\frac{dN_\gamma}{dx_\gamma} \propto 2 \times \frac{2}{3} \times \frac{1}{10} E_\gamma^{-2}$$

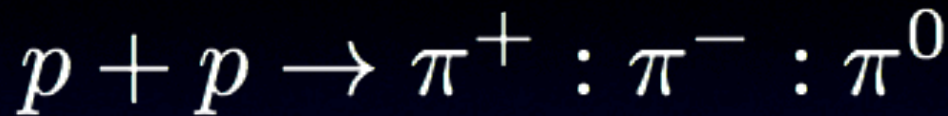
$$\frac{dN_\nu}{dE} \sim \frac{1}{4} \frac{dN_\gamma}{dE}$$

factor of 8 between gamma and nu fluxes (no oscillations)

Kelner and Aharonian for a full calculation [http://prd.aps.org/
abstract/PRD/v78/i3/e034013](http://prd.aps.org/abstract/PRD/v78/i3/e034013)

Galactic SNR and pp

In Galactic SN shocks CRs interact with the H in the Galactic disk (pp interactions, lower threshold than p-gamma)



$$E_{p,th} = \frac{(2m_p + m_\pi)^2 - 2m_p^2}{2m_p} \sim 1.23 \text{ GeV}$$

if all muons decay and for E^{-2} p spectrum:

$$\frac{dN_\nu}{dE} \sim 2 \times \frac{2}{3} \times \frac{1}{20}$$

← 2 pions x 1/3

$$\frac{dN_\gamma}{dE} \sim 2 \times \frac{1}{3} \times \frac{1}{10}$$

Assume always:

$$x_\nu = \frac{E_\nu}{E_p} = \frac{1}{4} \langle x_F \rangle = \frac{1}{20}$$

$$x_\gamma = \frac{E_\gamma}{E_p} = \frac{1}{2} \langle x_F \rangle = \frac{1}{10}$$

↙ 0.2 ↘

Ignoring oscillations there is a factor of about 1 between the gamma and neutrino flux.

For a full calculation see: <http://arxiv.org/pdf/astro-ph/0606058> for pp

**relation between pionic
gamma rays and neutrinos
robust (after oscillations)**

$$\frac{dN_{\nu}}{dE} = \frac{1}{2} \frac{dN_{\gamma}}{dE} \quad \text{for } p + p$$

$$\frac{dN_{\nu}}{dE} = \frac{1}{8} \frac{dN_{\gamma}}{dE} \quad \text{for } \gamma + p$$

E^{-2} assumed

NEUTRINO OSCILLATION REMINDER

Conversion probability (3 neutrino flavors and in vacuum):

$$P(\nu_\alpha \rightarrow \nu_\beta; L) = \delta_{\alpha\beta} - \sum_{j \neq k} U_{\alpha j}^* U_{\beta j} U_{\alpha k} U_{\beta k}^* (1 - e^{-i\Delta E_{jk}L})$$

Kayser's lectures

Where the mixing matrix is:

solar $U_{e1}, U_{e2} \leftrightarrow \theta_{12}$ CHOOZ $U_{e3} \leftrightarrow \theta_{13}$

$$U = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} \end{pmatrix} = \begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{13}s_{23}e^{i\delta} & c_{12}c_{23} - s_{12}s_{13}s_{23}e^{i\delta} & c_{13}s_{23} \\ s_{12}s_{23} - c_{12}s_{13}c_{23}e^{i\delta} & -c_{12}s_{23} - s_{12}s_{13}c_{23}e^{i\delta} & c_{13}c_{23} \end{pmatrix}$$

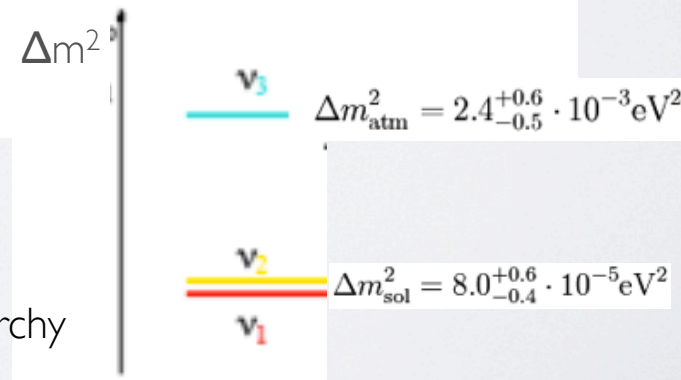
atmospheric $U_{e3} \leftrightarrow \theta_{13}$ $U_{\mu 3}, U_{\tau 3} \leftrightarrow \theta_{23}$

$$s_{ij} \equiv \sin \theta_{ij}, c_{ij} \equiv \cos \theta_{ij}.$$

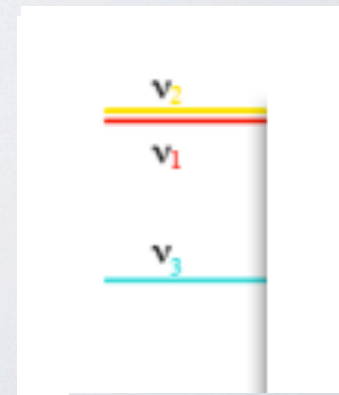
Observed values of oscillation parameters

- $\sin^2(\theta_{13}) < 0.032$ at 95% confidence level ($\theta_{13} < 10.3^\circ$) [4] CHOOZ
- $\tan^2(\theta_{12}) = 0.45_{-0.07}^{+0.09}$. This corresponds to $\theta_{12} \equiv \theta_{\text{sol}} = 33.9^\circ_{-2.2^\circ}^{+2.4^\circ}$ ("sol" stands for solar) [6]
- $\sin^2(2\theta_{23}) = 1_{-0.1}^{+0}$, corresponding to $\theta_{23} \equiv \theta_{\text{atm}} = 45 \pm 7^\circ$ ("atm" for atmospheric) [7]
- $\Delta m_{21}^2 \equiv \Delta m_{\text{sol}}^2 = 8.0_{-0.4}^{+0.6} \cdot 10^{-5} \text{eV}^2$ [6]
- $|\Delta m_{31}^2| \approx |\Delta m_{32}^2| \equiv \Delta m_{\text{atm}}^2 = 2.4_{-0.5}^{+0.6} \cdot 10^{-3} \text{eV}^2$ [7]

direct hierarchy



inverted hierarchy



ASTROPHYSICAL NEUTRINO OSCILLATIONS

In the limit $L \rightarrow \infty$, we have

$$P(\nu_\alpha \rightarrow \nu_\beta; L = \infty) = \delta_{\alpha\beta} - \sum_{j \neq k} U_{\alpha j}^* U_{\beta j} U_{\alpha k} U_{\beta k}^* = \sum_j |U_{\alpha j}|^2 |U_{\beta j}|^2$$

average over rapid oscillations

$$P(\nu_\alpha \rightarrow \nu_\beta) = \sum_i |U_{\alpha,i}|^2 |U_{\beta,i}|^2$$

$$P(\nu_e \rightarrow \nu_e) = \sum_i |U_{ei}|^2 |U_{ei}|^2 = |U_{e1}|^4 + |U_{e2}|^4 + |U_{e3}|^4 = 0.82^4 + 0.57^4 + 0 = 0.56$$

$$P(\nu_e \rightarrow \nu_\mu) = \sum_i |U_{ei}|^2 |U_{\mu i}|^2 = |U_{e1}|^2 |U_{\mu 1}|^2 + |U_{e2}|^2 |U_{\mu 2}|^2 + |U_{e3}|^2 |U_{\mu 1}|^2 = 0.82^2 \cdot 0.4^2 + 0.57^2 \cdot 0.58^2 + 0 = 0.22$$

$$P(\nu_e \rightarrow \nu_\tau) = \sum_i |U_{ei}|^2 |U_{\tau i}|^2 = |U_{e1}|^2 |U_{\tau 1}|^2 + |U_{e2}|^2 |U_{\tau 2}|^2 + |U_{e3}|^2 |U_{\tau 1}|^2 = 0.82^2 \cdot 0.4^2 + 0.57^2 \cdot 0.58^2 + 0 = 0.22$$

$\nu_\alpha \backslash \nu_\beta$	ν_e	ν_μ	ν_τ
ν_e	60%	20%	20%
ν_μ	20%	40%	40%
ν_τ	20%	40%	40%

At source:

$$\nu_e : \nu_\mu : \nu_\tau \sim 1 : 2 : 0$$

60% of ν_e survive and 2x20% come from

$2 \times \nu_\mu = 100\%$

$2 \times 40\% = 80\%$ of $2 \times \nu_\mu$ survive and 20% come

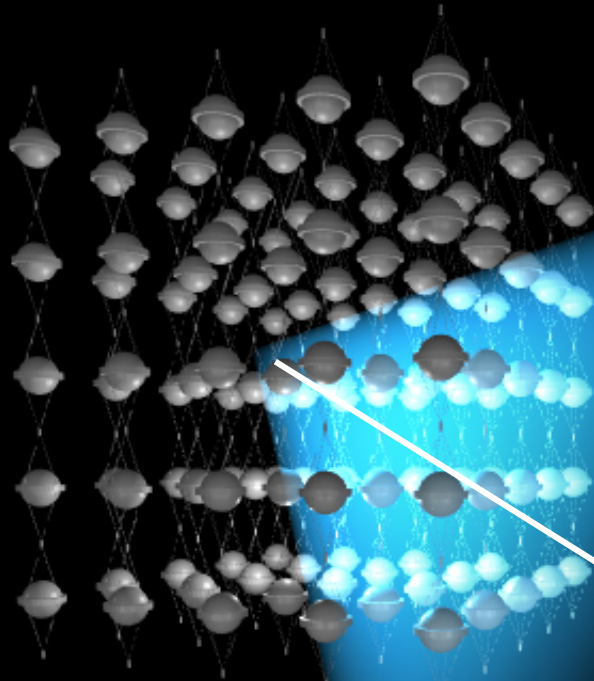
from $\nu_e = 100\%$

20% of ν_τ come from ν_e and $2 \times 40\%$

from $\nu_\mu = 100\%$

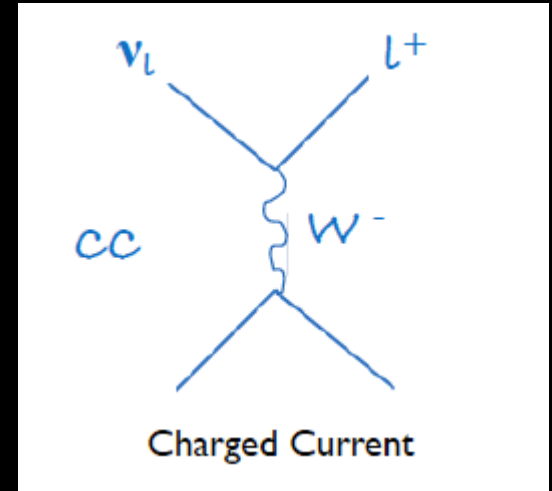
At Earth: 1: 1: 1

- shielded and optically transparent medium



$$P_{\mu \otimes \nu} = \frac{\lambda_{\mu}}{\lambda_{\nu}} = n R_{\mu} \sigma_{\nu}$$

μ



ν

- lattice of photomultipliers

Effective area

$$N_{\mu} = \int A_{\text{eff}}^{\nu} (E_{\nu}, \theta_{\nu}, \phi_{\nu}) \frac{d\Phi_{\nu}}{dE_{\nu} d\Omega_{\nu}} dE_{\nu} d\Omega_{\nu}$$

Neutrinos flux model

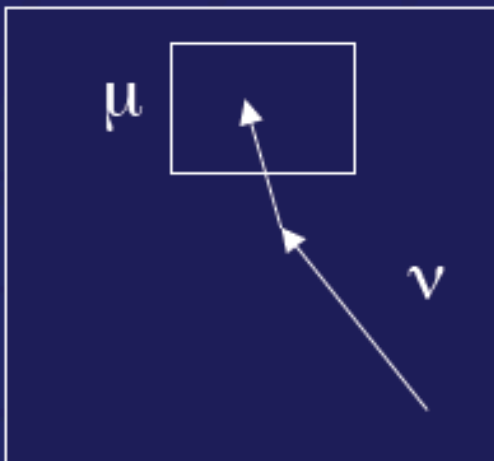
Target nucleon density

Neutrino-nucleon cross-section

$$A_{\text{eff}}^{\nu} = V_{\text{gen}} \times \frac{N_{\text{xxx}}(E_{\nu}, \theta_{\nu}, \phi_{\nu})}{N_{\text{gen}}(E_{\nu}, \theta_{\nu}, \phi_{\nu})} \times (\rho N_A) \sigma(E_{\nu}) \times P_{\text{earth}}(E_{\nu}, \theta_{\nu})$$

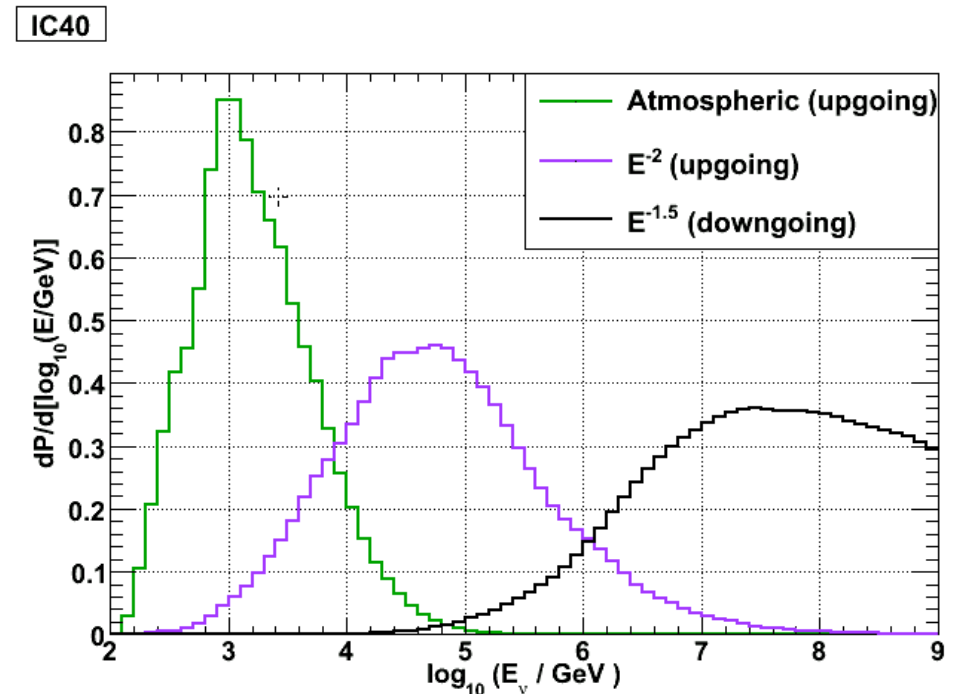
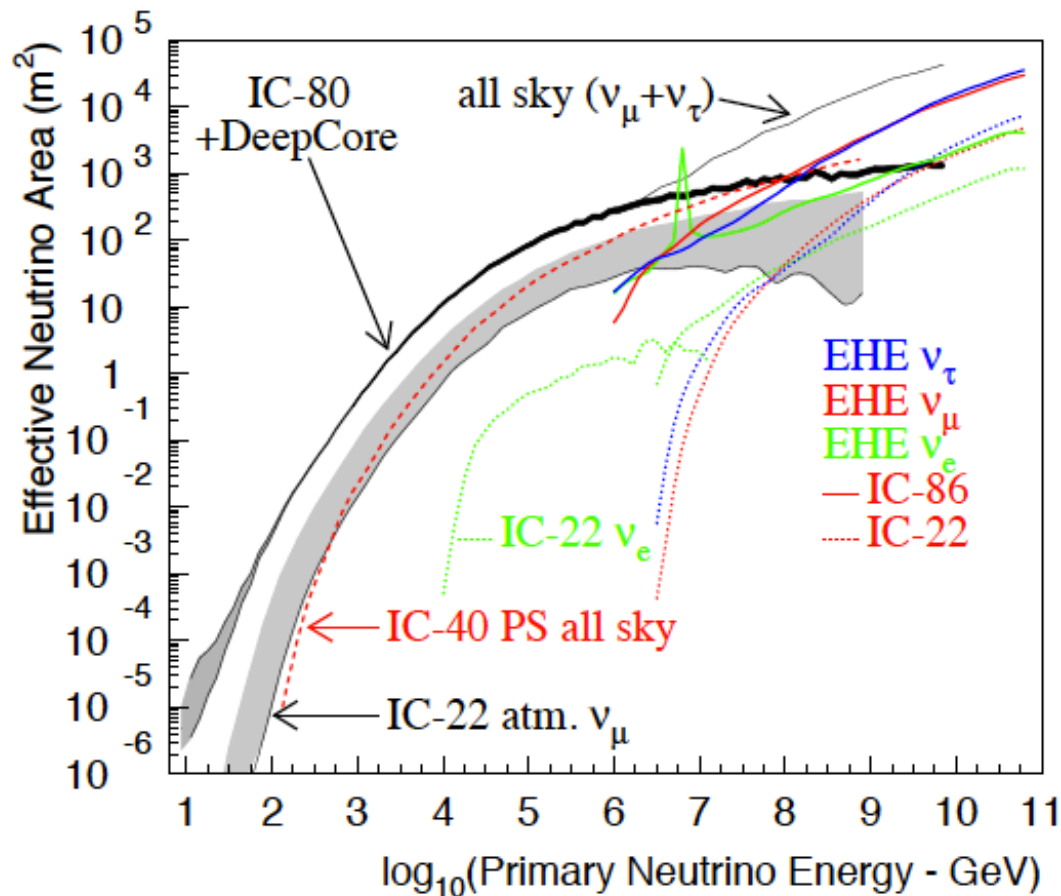
$$P_{\text{earth}}(E_{\nu}, \theta_{\nu}) = e^{-N_A \sigma(E_{\nu}) \int \rho dl}$$

Shadowing effect



V_{gen} = generation volume of all neutrino interaction that potentially would produce the signal in the detector

Effective area (useful for rates)



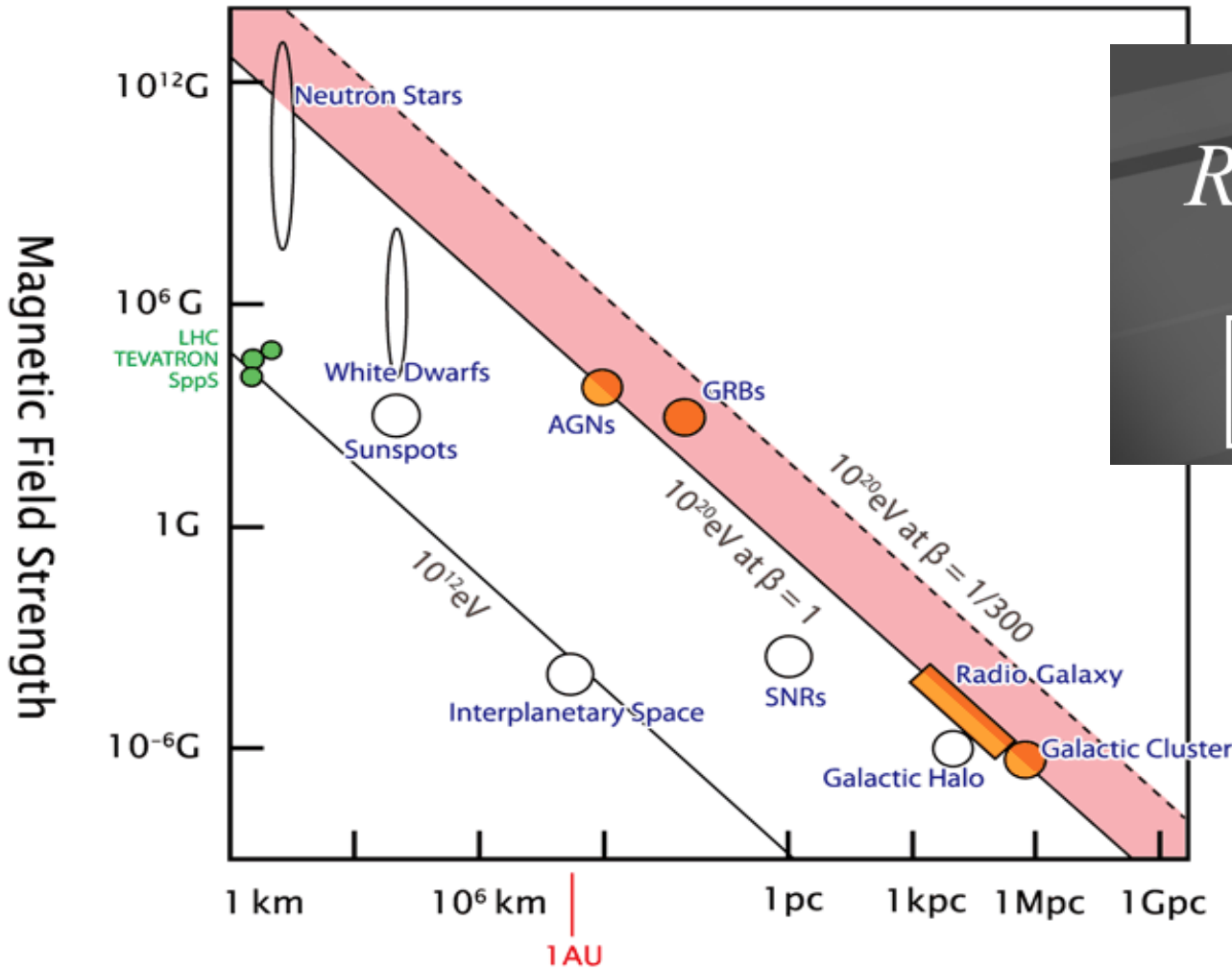
you get it from the generation volume and it is analysis dependent

$$A_{eff}^\nu(E_\nu, \theta_\nu) = V_{eff}(E_\nu, \theta_\nu) \times (\rho N_A) \times \sigma(E_\nu) \times P_{Earth}(E_\nu, \theta_\nu) \quad V_{eff}(E_\nu, \theta_\nu) = \frac{N_{selected}(E_\nu, \theta_\nu)}{N_{generated}(E_\nu, \theta_\nu)} \times V_{generation}$$

you use it to get the event rate

$$\int_{E_{\nu, min}}^{E_{\nu, max}} dE A_{eff}^\nu(E_\nu, \delta_{source}) \frac{dN}{dE} dE_\nu$$

ACCELERATION: HILLAS' PLOT



$$R_{gyro} \left(= \frac{E}{vqB} \right) \leq R$$

$$E \leq v q B R$$

$$r_L = \sqrt{\frac{1}{4\pi\alpha} \frac{E}{ZeB}}$$

$$= \frac{1.1}{Z} \left(\frac{E}{\text{EeV}} \right) \left(\frac{B}{\mu\text{G}} \right)^{-1} \text{ kpc.}$$

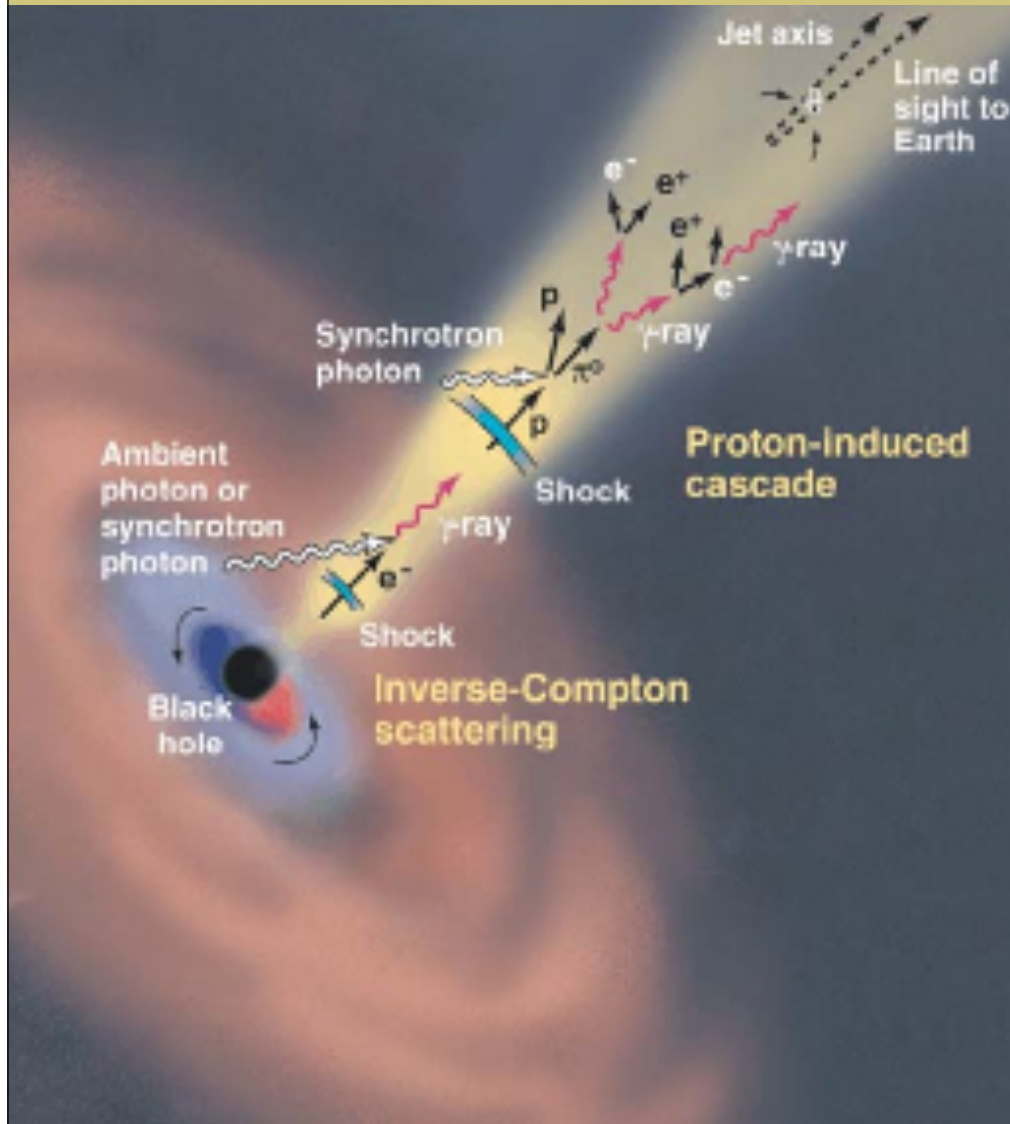
Size

$$E_{\text{max}} \simeq Z \left(\frac{B}{\mu\text{G}} \right) \left(\frac{R_{\text{source}}}{\text{kpc}} \right) \times 10^9 \text{ GeV.}$$

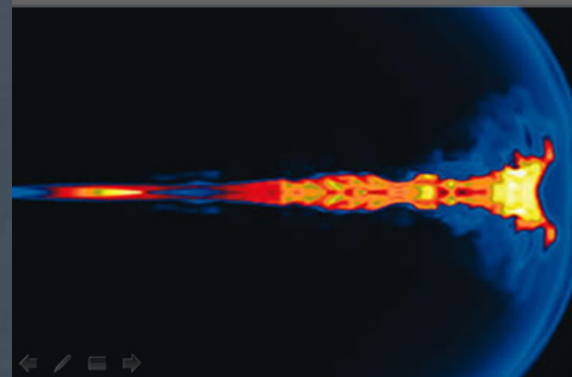
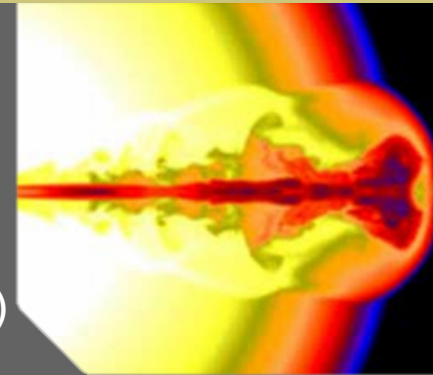
Extragalactic sources

AGN's

gamma ray bursts



collapse of massive star produces a spinning black hole (long bursts)

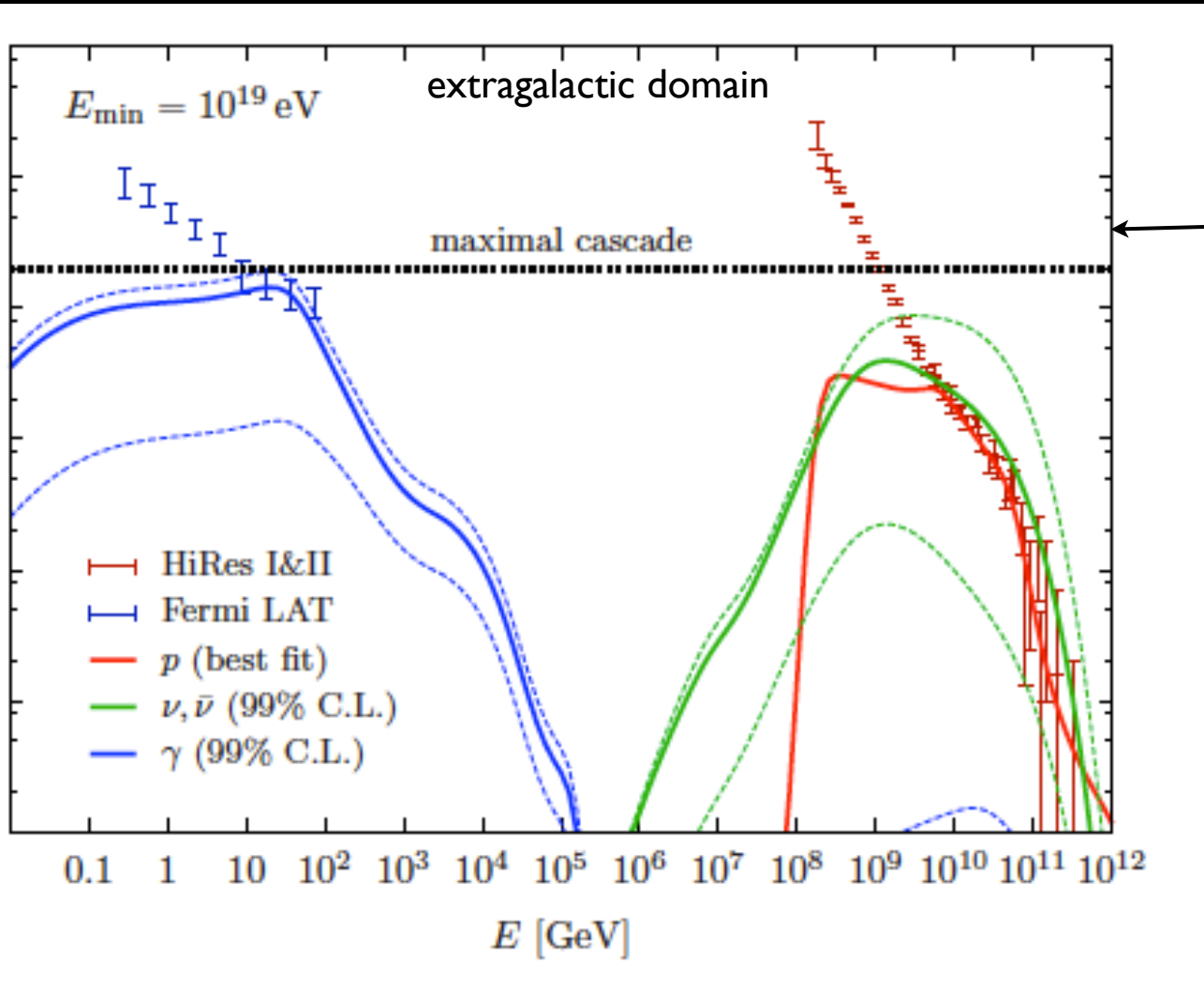


shocks produced in the outflow of the spinning black hole: electrons and protons ?

Associated to SN

compact binary mergers (short bursts)

Fermi-LAT extragalactic photon background - UHECR - neutrino connection



gamma flux in GeV-TeV saturating the
 total EM radiation energy density from
 proton energy losses

normalization to Fermi diffuse flux is
 tricky because gamma cascading and
 galactic point source flux subtraction.



~ cosmic ray + neutrino

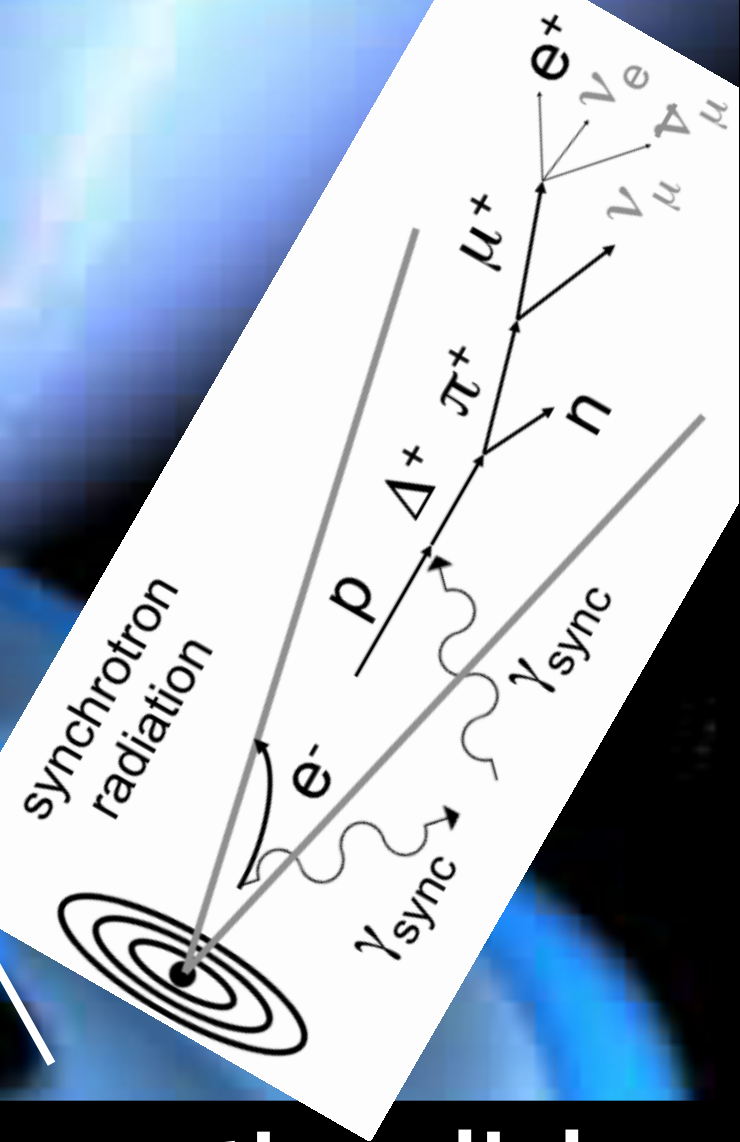


~ cosmic ray + gamma

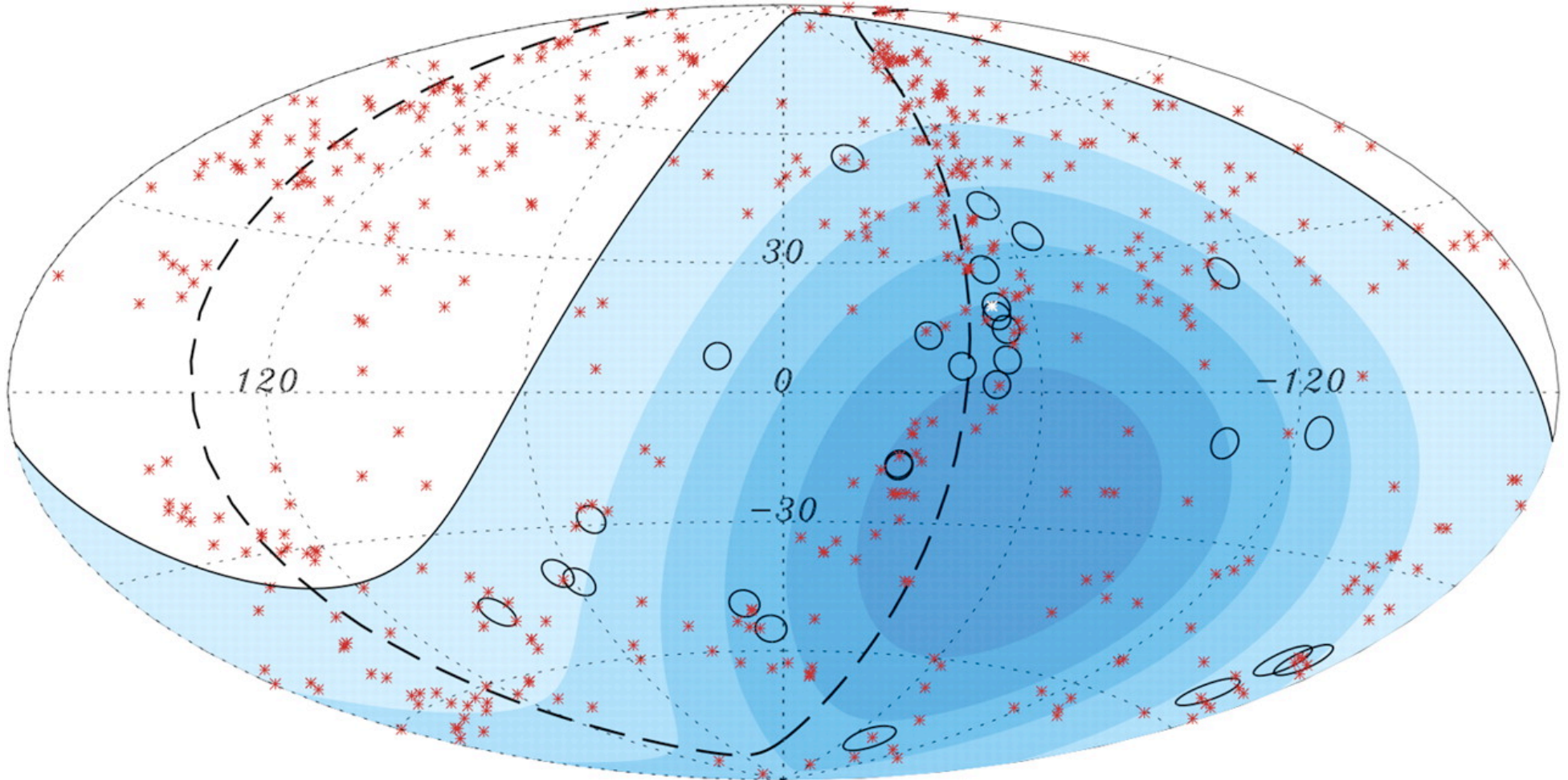
active galaxy

supermassive
black hole

- accretion disk
- jet with blobs



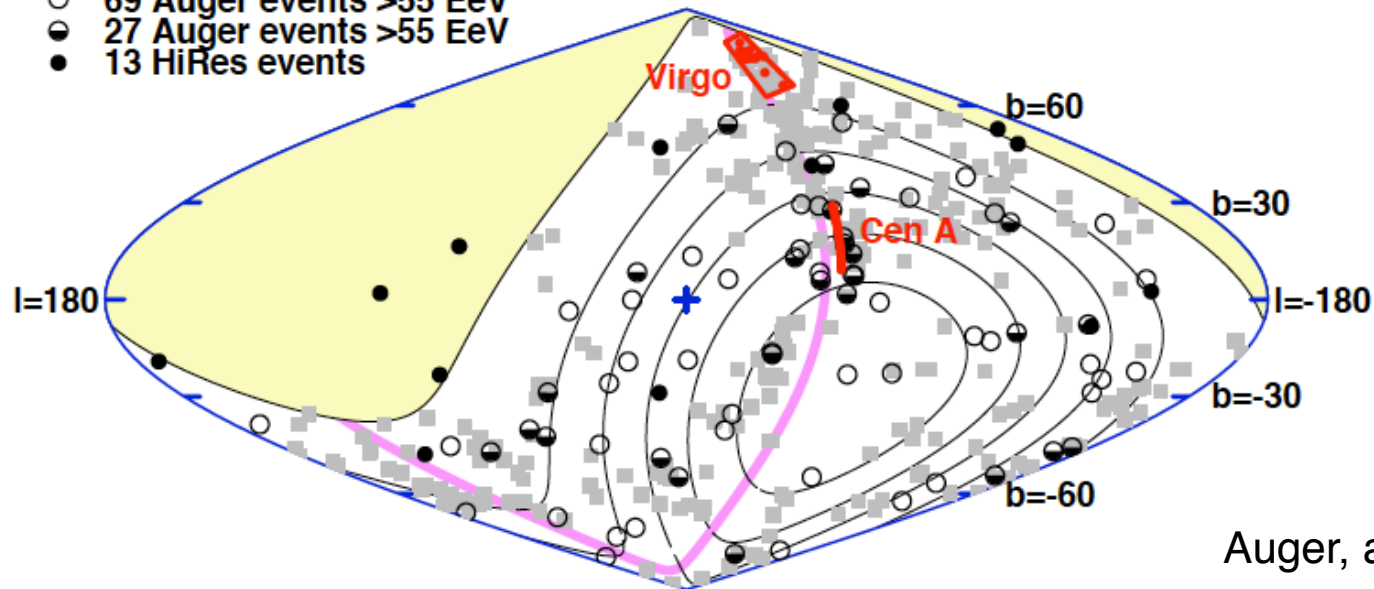
Proton astronomy?



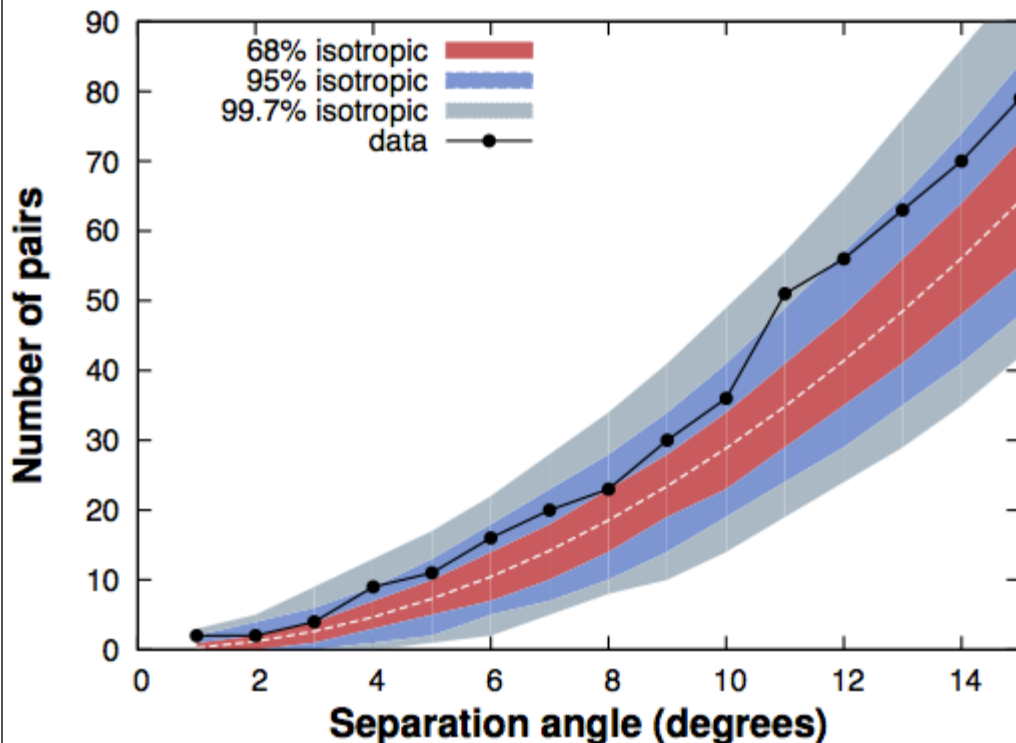
$$\theta \cong \frac{d}{R_{gyro}} = \frac{dB}{E} \quad \frac{\theta}{1^\circ} \sim \frac{(d/10\text{Mpc})(B/10^{-9}\text{Gauss})}{E/3 \times 10^{20}\text{eV}}$$

The sources revealed: AGNs?

- 69 Auger events >55 EeV
- ◐ 27 Auger events >55 EeV
- 13 HiRes events



Auger, arXiv:1009.1855



Auger: correlating fraction of $E > 55$ EeV events with AGN close-by catalogue decreased from 69% to 37% (21% expected to occur by chance for an isotropic flux).

18% of 69 events are inside 18° from Cen A but no event from M87.

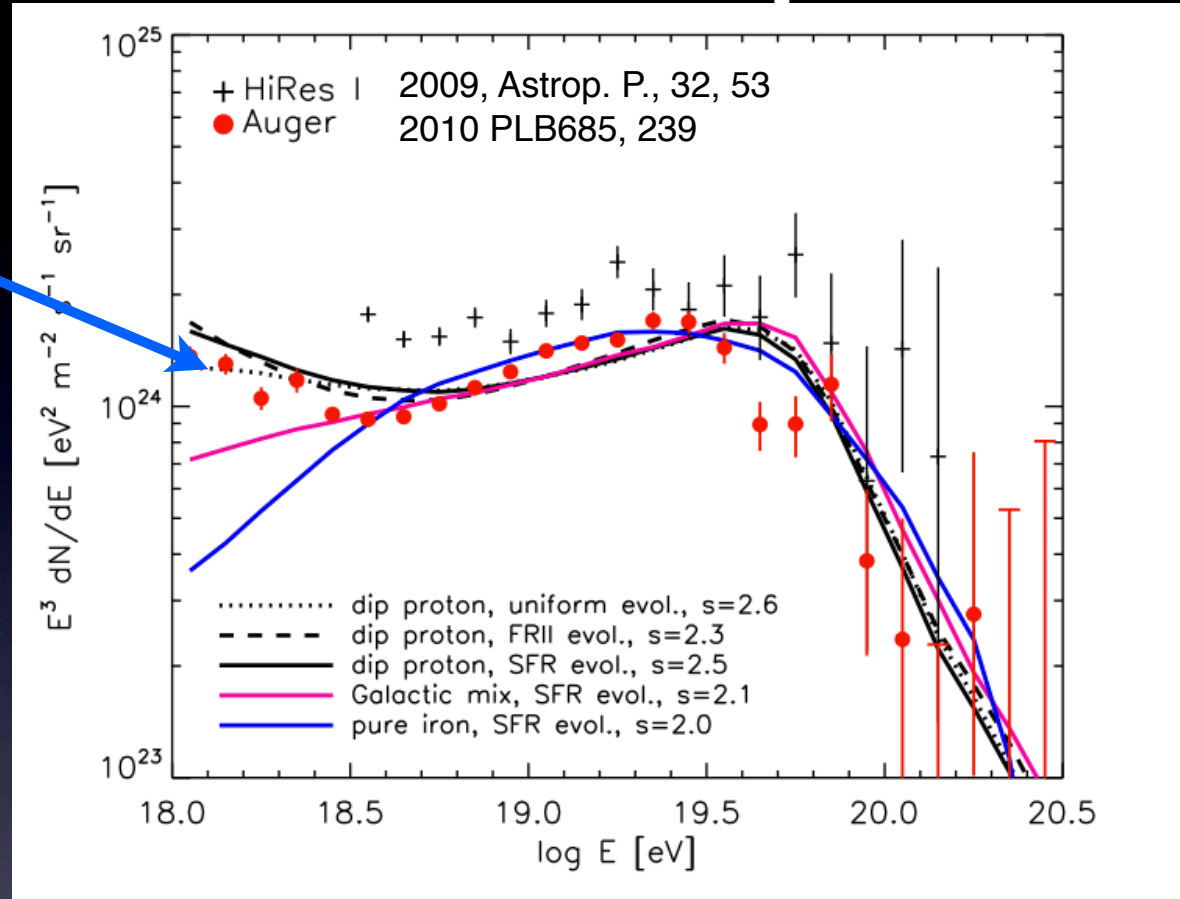
Autocorrelation function: largest deviation from isotropic distribution at 11° .

For an isotropic distribution 1.3% pairs of the 69 events have 51 or more pairs inside 11° .

The lack of knowledge of B-fields prevents to establish if UHECR astronomy is possible.

GZK-cut off proved

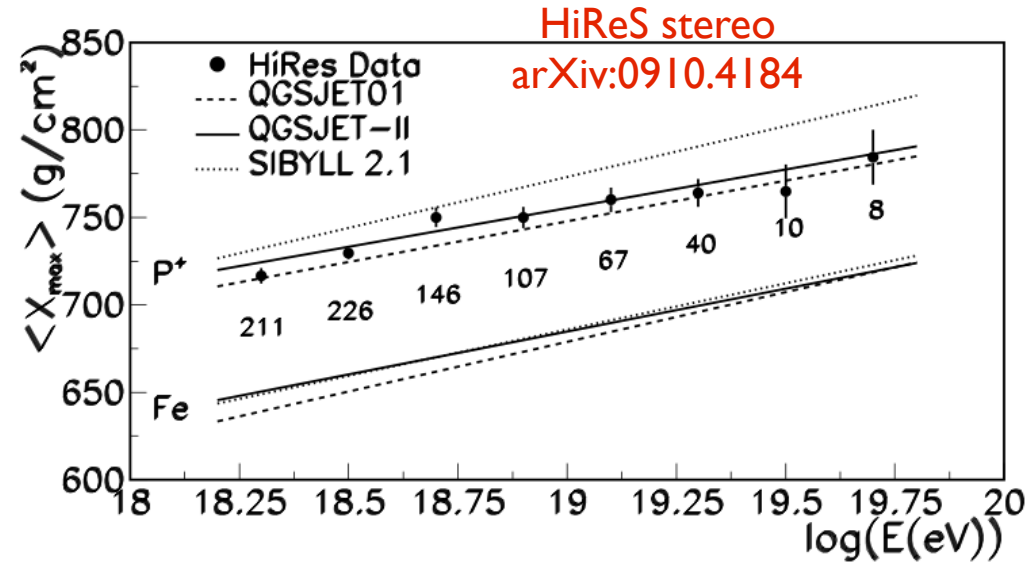
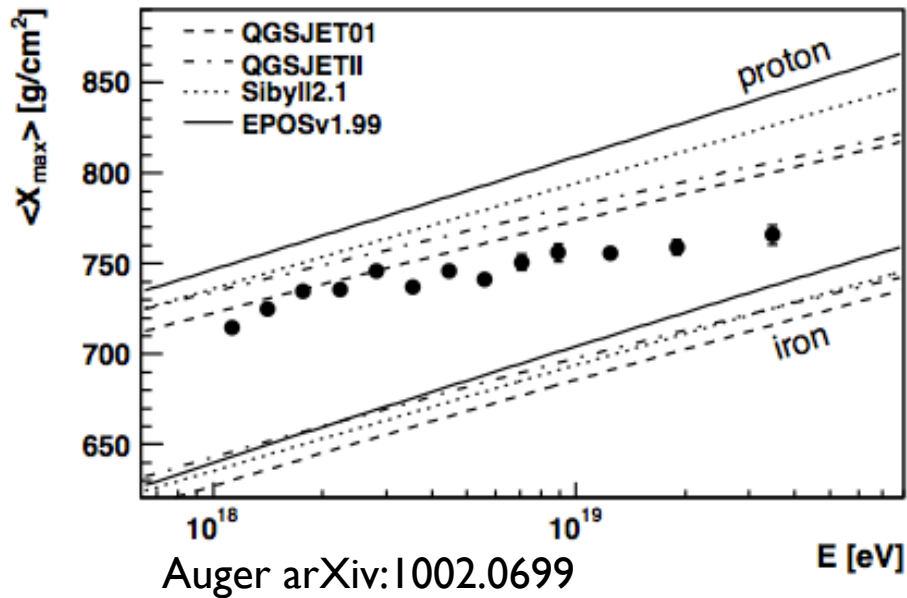
Where is the galactic - extra-galactic transition?



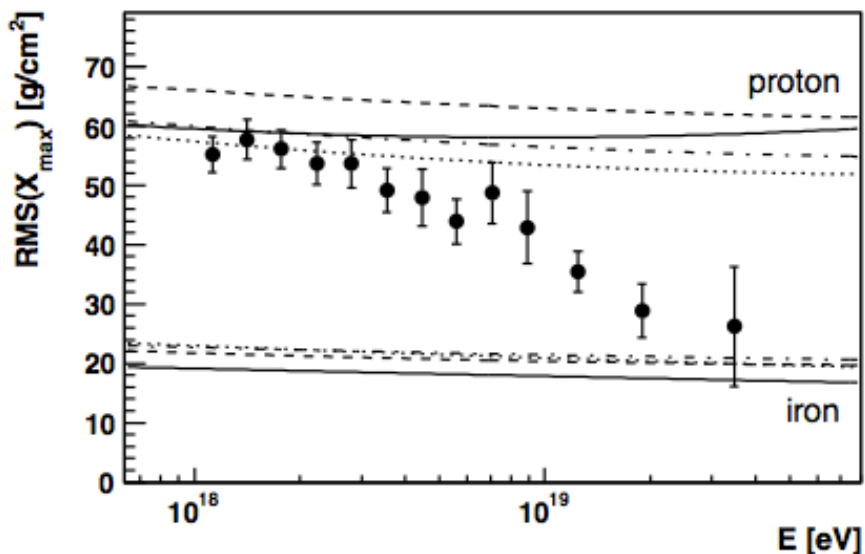
Nature trick: spectrum shows same feature for p and Fe. While jets can easily accelerate Fe ($E_{\max} \propto Z$), Fe would not survive photo-disintegrations when injected in extra-galactic accelerators such as a GRB fireball.

Heavy/light composition?

In a heavy composition scenario anisotropies would be washed out by B-fields.



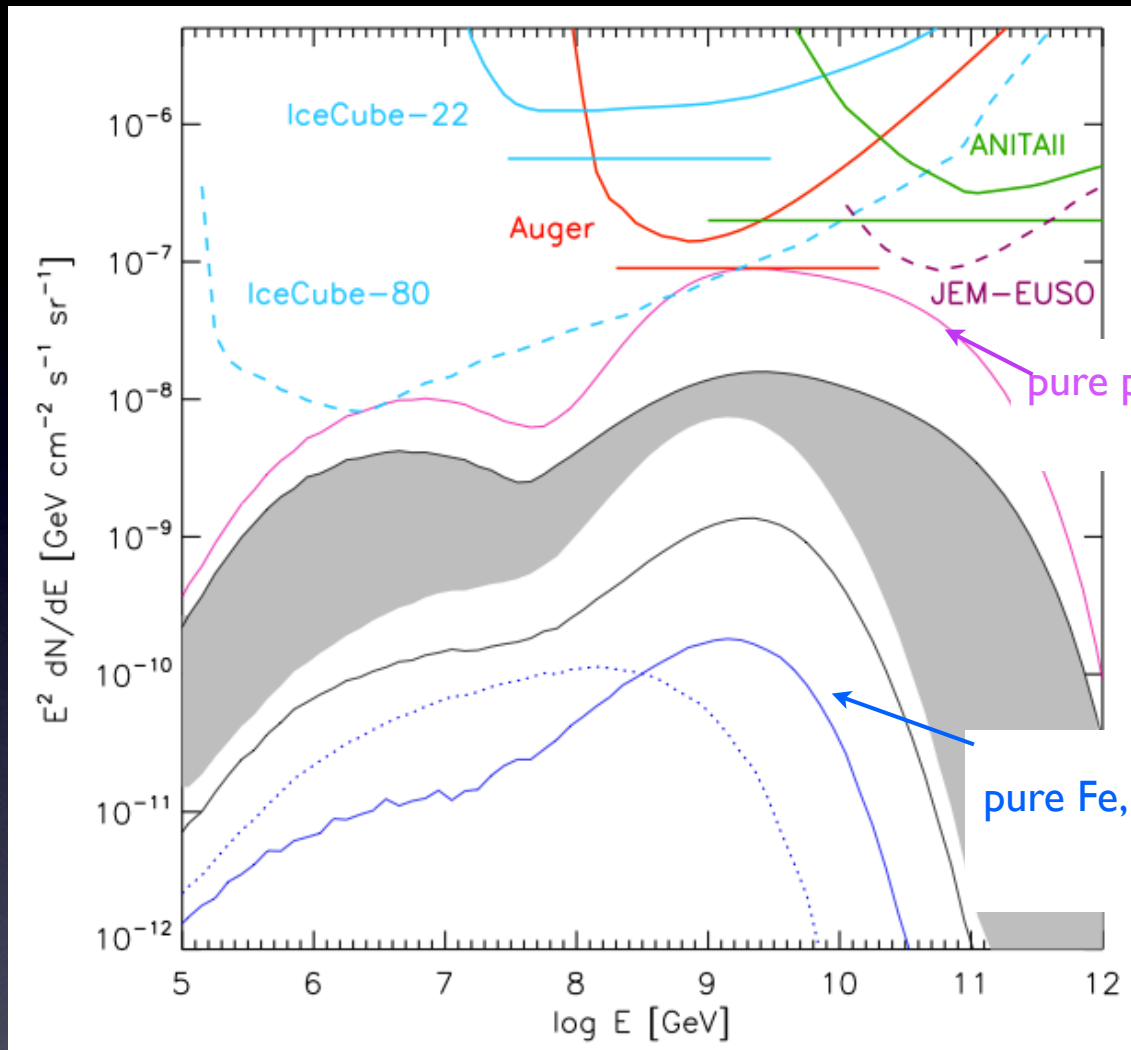
HiReS and Auger X_{\max} are not directly comparable due to fiducial cuts.



An evidence of neutrinos in coincidence with UHECR would indicate that the extra-galactic sources are not GRBs because p take much longer than neutrinos. It would also indicate a p-dominated composition.

IceCube: no correlation evidence in 40-strings, 40+59 strings to be presented at ICRC 2011

Fluxes of UHECR neutrinos



pure p, strong evolution, $E_{\max} = 3 \times 10^{21}$ eV

pure Fe, uniform evolution, $E_{\max} = Z \times 100$ EeV

The diffuse extragalactic background measured by Fermi can be used as a constraint accounting for photon cascading.

UHECR neutrino fluxes depend on:

- 1) the transition energy between galactic and extra-galactic CRs
- 2) E_{\max}
- 3) injection spectrum and source evolution

Kotera, Olinto's review arXiv: 1101.4256

Useful references

Halzen and Hooper, Rept.Prog.Phys.65:1025-1078,2002

<http://arxiv.org/abs/astro-ph/0204527v2>

Anchordoqui & TM, Annual Review of Nuclear and Particle Science (2009) 60

Anchordoqui, arXiv:1104.0509v1

J. Becker, High-energy neutrinos in the context of multimessenger astrophysics

Neutrino unbound

Textbooks:

T.K. Gaisser book on CRs

Longair, High Energy Astrophysics

Berezinski, Neutrino Astrophysics 1995