# [CMS] dE/dX measurements for long-lived particle search

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#### Intro to HSCPs





#### Ionization observables

 $\frac{2j-2}{N}$  $\frac{N}{J} + \sum_{j=1}^N \left[ P_j \left( P_j - \right. \right.$  $\frac{1}{12N}$  $\frac{1}{\infty}$ 



 $|I_h - C|$ 

 $\overline{K}$ 

# Event preselection



$$
I_{\text{PF}}^{\text{rel}} = \frac{(\sum_{\Delta R < \Delta R_{\text{mini-iso}}} p_{\text{T}}^{\text{PF}}) - p_{\text{T}}^{\text{HSCP}}}{p_{\text{T}}^{\text{HSCP}}}
$$

$$
I_{\text{trk}} = (\sum_{\Delta R < \Delta R_{\text{mini-iso}}} p_{\text{T}}^{\text{trk}}) -
$$

 $\int$  0.2,  $\Delta R_{\text{mini-iso}} = \begin{cases} 0.27 & \text{pT} = 0.0 \text{ GeV} \\ 10 \text{ GeV} / p_{\text{T}}^{\text{HSCP}}, & 50 \text{ GeV} < p_{\text{T}}^{\text{HSCP}} < 200 \text{ GeV} \\ 0.05, & p_{\text{T}}^{\text{HSCP}} \geq 200 \text{ GeV}. \end{cases}$  $\mathbf{H}\cap\mathbf{S}$ 

$$
\Delta R < \Delta R_{\text{mini-iso}}
$$
  

$$
p_{\text{T}}^{\text{HSCP}} \leq 50 \text{ GeV}
$$

 $p_{\rm T}^{\rm HSCP}$  ,

$$
\frac{J/Psi}{\sqrt{2\pi\left(\frac{1}{2}\right)^2}}\left(\frac{1}{2}\right)^{\frac{1}{2}}
$$

# Background prediction – two methods

We use two independent methods for the background prediction \* Both of them are data driven \* Reuse the trigger + preselection + signal systematics

#### Given the ATLAS excess this gives us an extra handle

Ionization method: new approach, relying on independence of the pixels and strips detectors, uses a transfer function and invokes a shape based analysis

Mass method: improved historical approach, assuming the independence of  $I_h$  and p as well as the  $p_T$  and  $G_i^{String}$  as a counting experiment in dedicated mass windows





#### Results

Ionization method: Signal region is defined by  $p_T > 200$  GeV,  $F_i^{\text{pixels}} > 0.9$  and the full shape of Gistrips is used



Signal region is defined by  $p_T > 70$  GeV and  $G_i^{String} > 0.22$ , counting experiment in mass window





### Interpretations summary

This model [\(2205.04473\)](https://arxiv.org/abs/2205.04473) was specifically created as an explanation to the ATLAS excess: provides a highly ionizing track with beta  $\sim$  1



For the cross section limits the ionization method provides better limits at low signal masses while the mass method is more efficient at large masses





### Conclusions





# Specific model for ATLAS excess

#### Pheno proposal to explain ATLAS excess

arXiv:2205.04473v3 [hep-ph] 20 Jul 2022 dE/dx from boosted long-lived particles

Gian F. Giudice, Matthew McCullough and Daniele Teresi CERN, Theoretical Physics Department, Geneva, Switzerland





generation Q=2 leptons

**Standard particles** 

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Tau' could be multicharged, q=1e and q=2e with two production modes



# Ionization observables (3-4)

Ionization estimator using the squared harmonic mean in order to suppress the tails of the Landau distribution  $dE/dx$ 

$$
I_{\rm h} = \left(\frac{1}{N} \sum_j^N (dE/dx_j)^{-2}\right)^{-1/}
$$

 $I_{\rm h} = K \frac{m^2}{p^2} + C$ 

Mass of the HSCP can be interpreted through approximating the Bethe-Bloch formula with  $I_h$ (m,p,K,C), where the empirical parameters K and C are determined using a sample of lowmomentum particles composed of protons, kaons and pions.







#### Pixel and strips detectors





Table 2: Summary of number of APV25 chips per module and strip pitch (strip pitch range) for barrel (endcap) sensor geometries in the strips detector, and their labels.



#### Detailed silicon simulations with PixelAV

The current track reconstruction uses a special reconstruction method (CPE) that was developed to simulate irradiated sensors using an external software called PixelAV.

Charge deposition based on Bichsel pion-Si cross-sections

Delta-ray range using continuously slowing-down approach with NIST ESTAR dEdx data

Multiple scattering and magnetic curvature of delta-rays

Carrier transport based on Runge-Kutta integration of saturated drift

- E-field is coming from ISE TCAD simulation of a pixel cell
- Includes charge trapping, diffusion, induction on implants

Electronics simulation: noise, linearity, thresholds, miscalibration





# Trigger choice

Trigger choice: Single muon trigger (HLT\_Mu50)

\* Motivated by the ATLAS excess being compatible with muons.

- \* Less QCD and more EW processes as background.
- \* Using MET triggers require a deeper understanding of R-hadrons trigger efficiency, and will be included in a next paper.



#### Background prediction – ionization method

Fully data driven method relying in the independence of the pixels and strips detectors

PASS and FAIL regions defined by the  $F_i^{Pixels}$  variable, then doing joint fit in bins of the  $G_i^{String}$  distribution using the full shape





A control regions defined by  $50 < p_T < 55$  GeV and  $55 < p_T < 200$  GeV show very good performance

No signal curve is shows as there is no signal in this CR

# Background prediction – mass method

Fully data driven method assuming the independence of  $I_h$  and p as well as the  $p_T$  and  $G_i^{String}$  as a counting experiment in dedicated mass windows

The shape of the background mass spectrum in D comes from fitted distributions of  $I_h$  in the control region B and p in the control region C

This is done in bins of pseudorapidity to take the correlation between  $p_T$  and p into account

 $\Rightarrow$   $m = p$ 

$$
p_{T}\n\begin{array}{c}\n\begin{array}{c}\n\hline\n\end{array} & \text{Mass spectrum} \\
\begin{array}{c}\n\hline\n\end{array} & \begin{array}{c}\n\hline\n\end{array} &
$$



 $I_h - C$ 

 $\overline{K}$ 

#### Results – ionization method

Signal region is defined by  $p_T > 200$  GeV,  $F_i^{\text{pixels}} > 0.9$  and the full shape of G $_i^{\text{String}}$  is used

No significant excess beyond the SM is observed (expected 0.4, observed 1)



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### Results – mass method

Signal region is defined by  $p_T > 70$  GeV and G $_i^{String} > 0.22$ , counting experiment in mass window







# Signal systematic uncertainties (1)

Dominant uncertainty coming from trigger efficiency: HSCPs that are moving too slowly could be associated with the next bunch crossing collision

Efficiency curves are derived as a function of the generator-level beta for different bins in eta, and variations on these curves have been conservatively estimated assuming a delay of 1.5 ns in the muon chambers, that is equivalent to the time resolution of the chambers.



#### Interpretations

This model ([2205.04473\)](https://arxiv.org/abs/2205.04473) was specifically created as an explanation to the ATLAS excess: it provides a highly ionizing track with beta  $\sim$  1



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Pair produced stau obs (exp) mass limits for ion method: 0.52-0.69 (0.53-0.71) TeV GMSB produced stau obs (exp) mass limits for ion method: 0.85 (0.86) TeV





Pair produced tau'-1e obs (exp) mass limits for ion method: 1.02 (1.05) TeV Pair produced tau'-2e obs (exp) mass limits for ion method: 1.32 (1.35) TeV



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## Conclusions

A signature based, model independent search for HSCPs has been presented

The background was studied with MC samples, and the preselection was optimized

Two methods of data-driven background predictions were developed: \* a novel approach relying on the independence of the ionization in the tracking detectors \* an improved version of the historical mass method

No significant excess over the SM expectation was found

The results were interpreted in 10 different models, one of them is a direct response to an excess seen by ATLAS

All figures are part of the results submitted to JHEP! ArXiv:2410.09164, SUPP material EXO-18-002.

HEPData record containing the (re)interpretation material: doi:10.17182/hepdata.153850

