

WIMP and FIMP Dark Matter in Singlet-Triplet Fermionic Model

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Based on [JHEP11(2022)133] (arXiv:2208.00849) , in collaboration with
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LHC Higgs Working Group WG3 (BSM) -- Extended Higgs Sector subgroup meeting

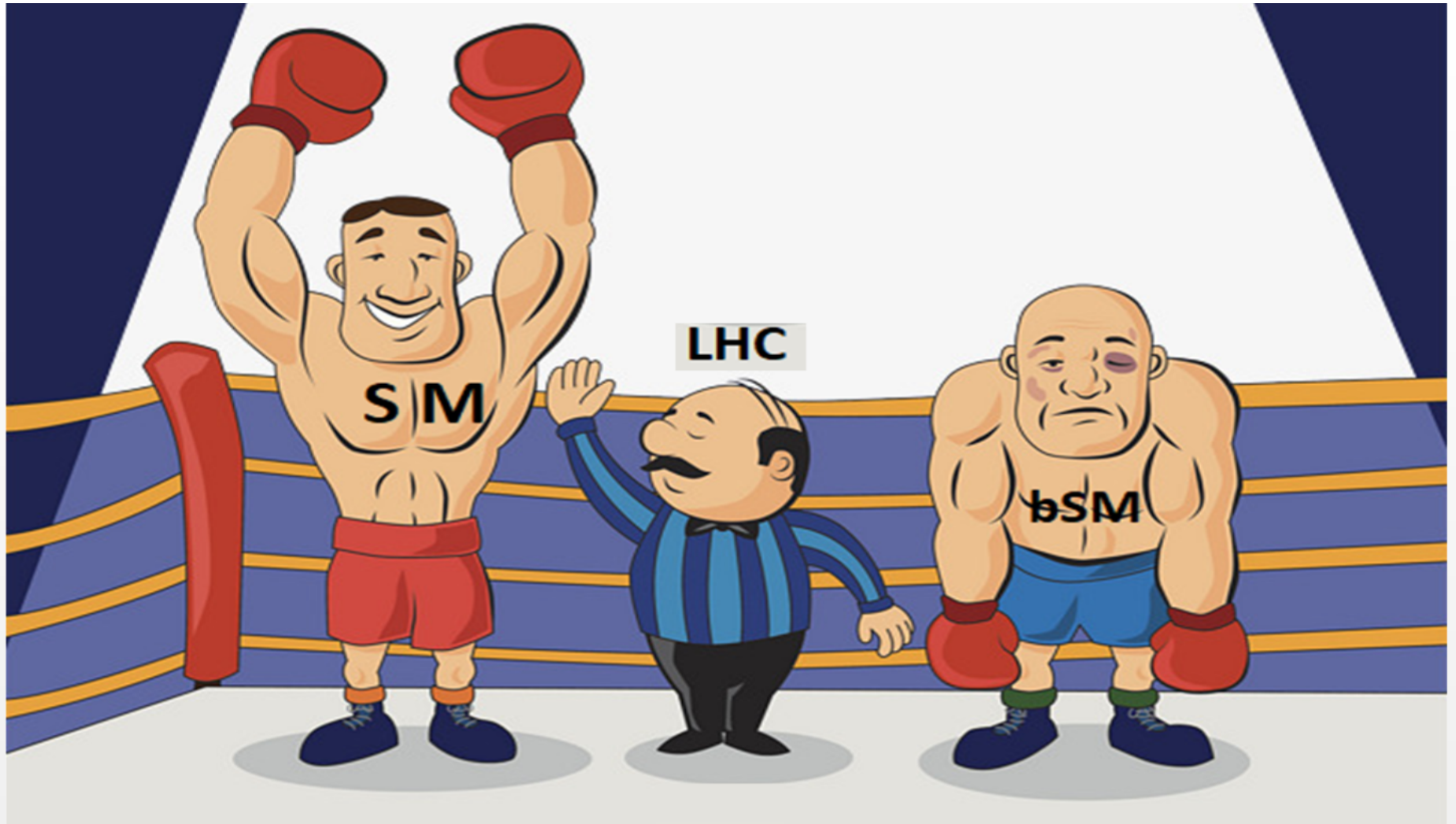
November 12th, 2024



Outline

- **Introduction**
- **Model Introduction and Motivation**
- **Constraints**
- **Dark Matter Production Mechanism**
- **Results based on STFM**
- **Summary**





LHC results able to confirm the validity of the SM, with no signatures of new physics.

Who can be a DM ?

- Should be massive
- Should be electrically neutral
- Should be present in early universe
- Should be stable or at least with half life greater than the age of the universe **Need a symmetry**

Singlet Scalar

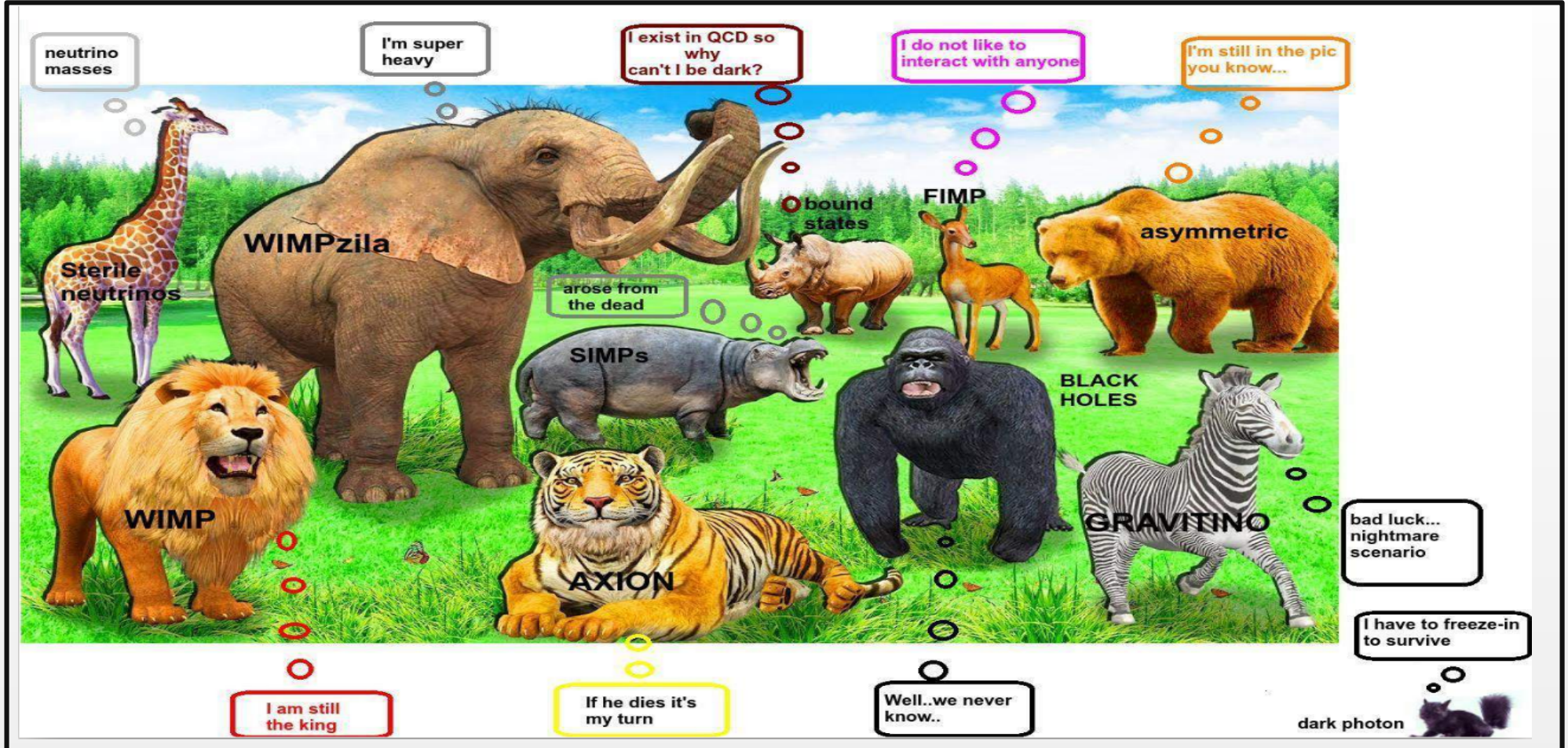
Singlet Fermion

Scalar in triplet repn

Fermion in triplet repn

...and many more

Zoo of Dark Matter Candidates



STFM to explain DM and neutrino mass

New Particles

Symmetry Group	Baryon Fields			Fermion Fields						Scalar Fields	
	Q_L^i	u_R^i	d_R^i	L_L^i	e_R^i	N'	ρ_1	ρ_2	ρ_3	ϕ_h	Δ
$SU(3)_c$	3	3	3	1	1	1	1	1	1	1	1
$SU(2)_L$	2	1	1	2	1	1	3	3	3	2	3
$U(1)_Y$	1/6	2/3	-1/3	-1/2	-1	0	0	0	0	1/2	0
\mathbb{Z}_2	+	+	+	+	+	-	+	+	-	+	+

Particle content and their corresponding charges under various symmetry groups.

The complete Lagrangian for the model:-

$$\begin{aligned}
 \mathcal{L} = & \mathcal{L}_{SM} + \sum_{i=1}^3 \text{Tr} [\bar{\rho}_i i \gamma^\mu D_\mu \rho_i] + \bar{N}' i \gamma^\mu D_\mu N' + \text{Tr}[(D_\mu \Delta)^\dagger (D^\mu \Delta)] - V(\phi_h, \Delta) \\
 & - \sum_{(i,j)=(1,1)}^{(3,2)} \lambda_{ij} \bar{L}_i \phi_h \rho_j^c - Y_{\rho\Delta} (\text{Tr}[\bar{\rho}_3 \Delta] N' + h.c.) - \sum_{i=1}^3 M_{\rho_i} \text{Tr}[\bar{\rho}_i^c \rho_i] - M_{N'} \bar{N}'^c N'
 \end{aligned}$$

Yukawa interaction term for neutrino mass generation in the Type III seesaw mechanism.

Higgs triplet is necessary

$$\begin{aligned}
 V(\phi_h, \Delta) = & -\mu_h^2 \phi_h^\dagger \phi_h + \frac{\lambda_h}{4} (\phi_h^\dagger \phi_h)^2 + \mu_\Delta^2 \text{Tr}[\Delta^\dagger \Delta] + \lambda_\Delta (\Delta^\dagger \Delta)^2 + \lambda_1 (\phi_h^\dagger \phi_h) \text{Tr}[\Delta^\dagger \Delta] \\
 & + \lambda_2 \left(\text{Tr}[\Delta^\dagger \Delta] \right)^2 + \lambda_3 \text{Tr}[(\Delta^\dagger \Delta)^2] + \lambda_4 \phi_h^\dagger \Delta \Delta^\dagger \phi_h + (\mu \phi_h^\dagger \Delta \phi_h + h.c.)
 \end{aligned}$$

- ϕ_h acquires vev and EWSB takes place.
- Δ_0 acquires an vev and takes the following form,

$$\phi_h = \begin{pmatrix} \phi^+ \\ \frac{v + H + i\xi}{\sqrt{2}} \end{pmatrix} \quad \Delta = \begin{pmatrix} \frac{\Delta^0 + v_\Delta}{2} & \frac{\Delta^+}{\sqrt{2}} \\ \frac{\Delta^-}{\sqrt{2}} & -\frac{\Delta^0 + v_\Delta}{2} \end{pmatrix}.$$

$$\langle \Delta^0 \rangle = v_\Delta = \frac{\mu v^2}{2 \left(\mu_\Delta^2 + (\lambda_4 + 2\lambda_1) \frac{v^2}{4} + (\lambda_3 + 2\lambda_2) \frac{v_\Delta^2}{2} \right)}$$

After symmetry breaking, CP even neutral Higgs mixes with each other.

$$\begin{aligned} H_1 &= \cos \alpha H + \sin \alpha \Delta^0 \\ H_2 &= -\sin \alpha H + \cos \alpha \Delta^0 \end{aligned}$$

SM like Higgs

CP even neutral Higgs mixing

The charged scalar also mixes with each other after EWSB takes the following form,

$$\begin{aligned} G^\pm &= \cos \delta \phi^\pm + \sin \delta \Delta^\pm \\ H^\pm &= -\sin \delta \phi^\pm + \cos \delta \Delta^\pm \end{aligned}$$

Charged Higgs

$$\tan \delta = \frac{2v_\Delta}{v}$$

Charged Higgs mixing

Dark Sector

Two neutral fermion states ρ_3^0 and N' also mixes.

$$\begin{aligned}\rho &= \cos \beta \rho_3^0 + \sin \beta N'^c \\ N &= -\sin \beta \rho_3^0 + \cos \beta N'^c.\end{aligned}$$

Mixing between singlet and triplet states due to

$$Y_{\rho\Delta} \text{Tr}[\bar{\rho}_3 \Delta] N'$$

where the mixing angle is,

$$\tan 2\beta = \frac{Y_{\rho\Delta} v_\Delta}{M_{\rho_3} - M_{N'}}.$$

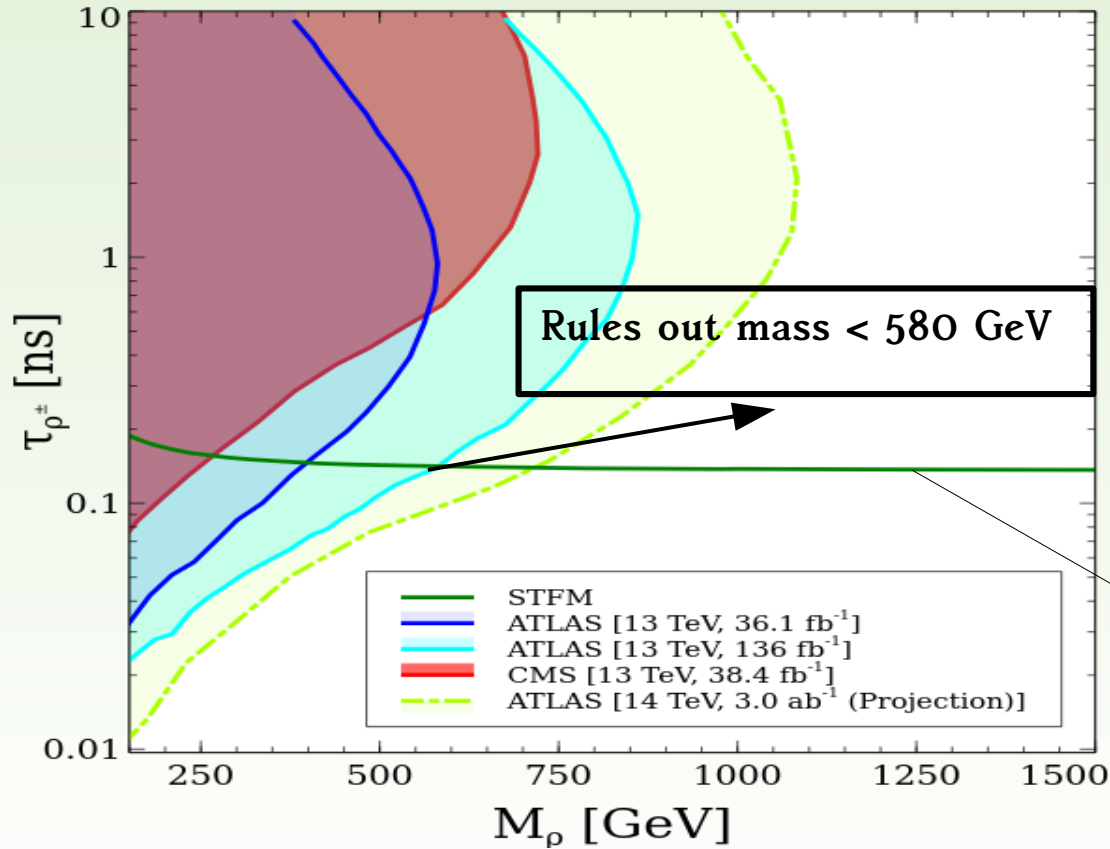
In the limit $Y_{\rho\Delta} \sim \mathcal{O}(10^{-10})$,

$$M_N \sim M_{N'}, \quad M_\rho \sim M_{\rho_3}$$

N is primarily singlet and ρ is primarily triplet states

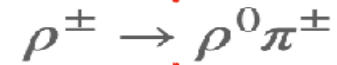
DM Constraints:-

Collider constraints on ρ :-



ρ contains ρ^0 and ρ^\pm

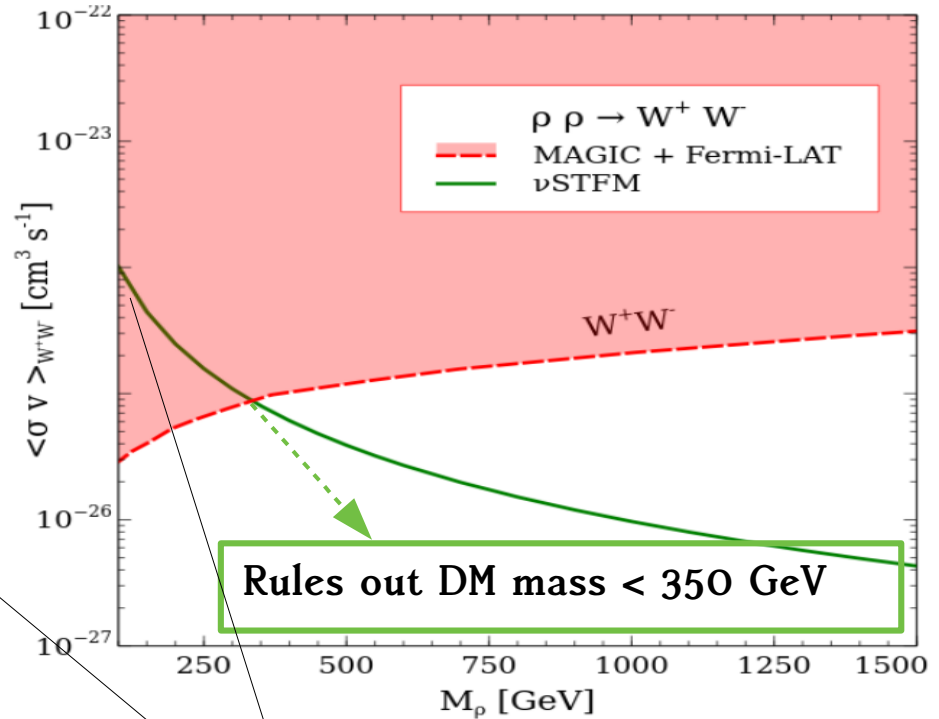
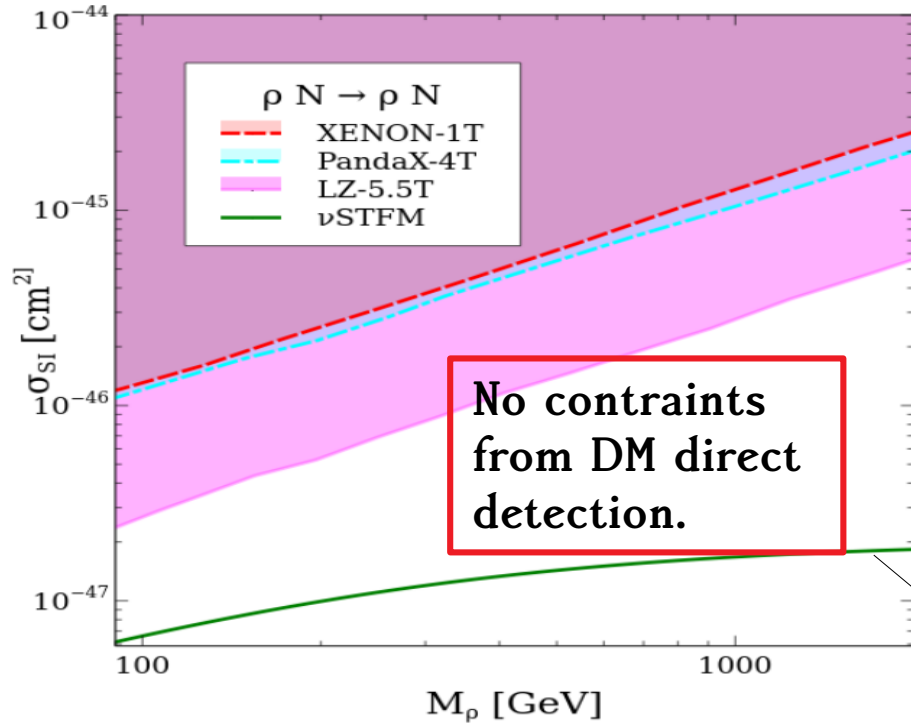
$$\Delta M = M_{\rho^\pm} - M_{\rho^0} = 167 \text{ MeV}$$



Decay of ρ^\pm manifest itself as disappearing track signal

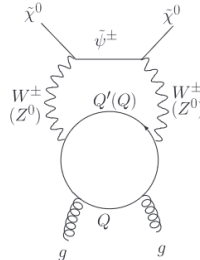
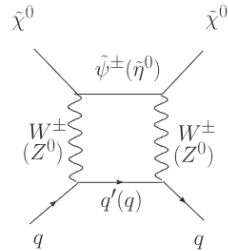
Theory Prediction

DM direct and Indirect Detection

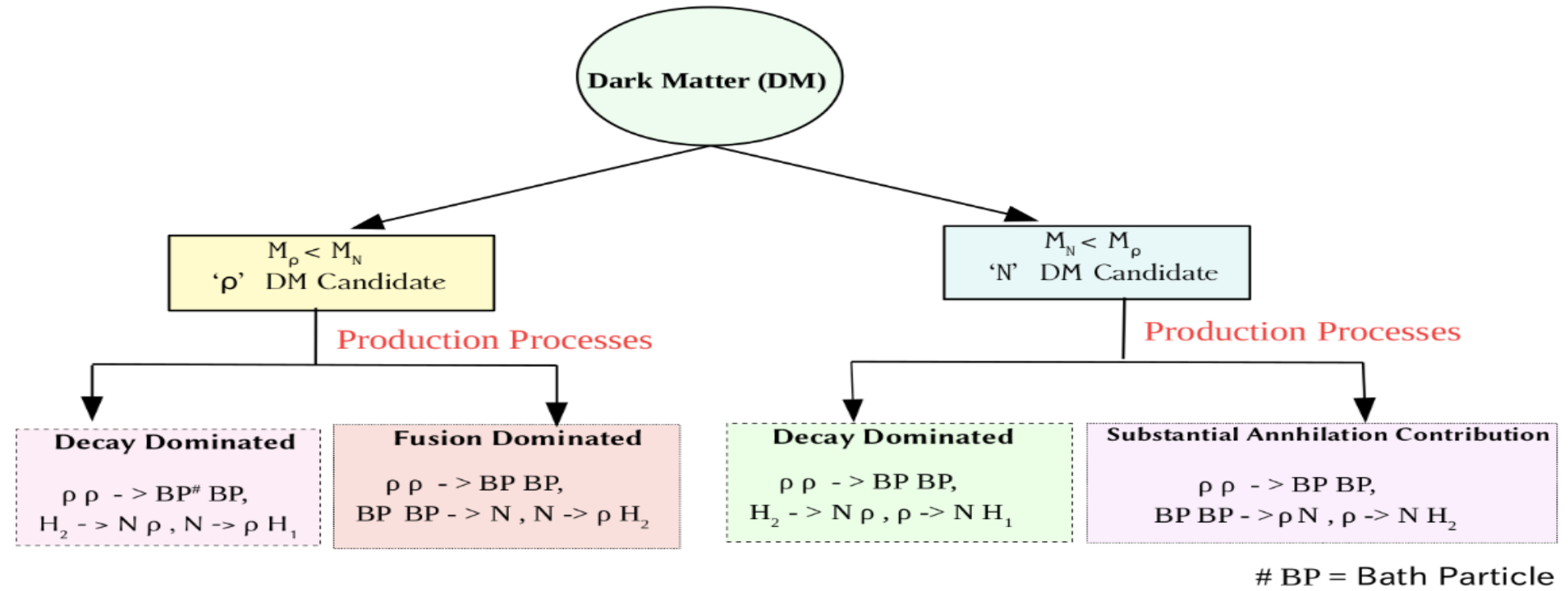


- No tree level diagram
- One loop mediated by W and two loop gluons

[JHEP07(2011)005]



Theory Prediction



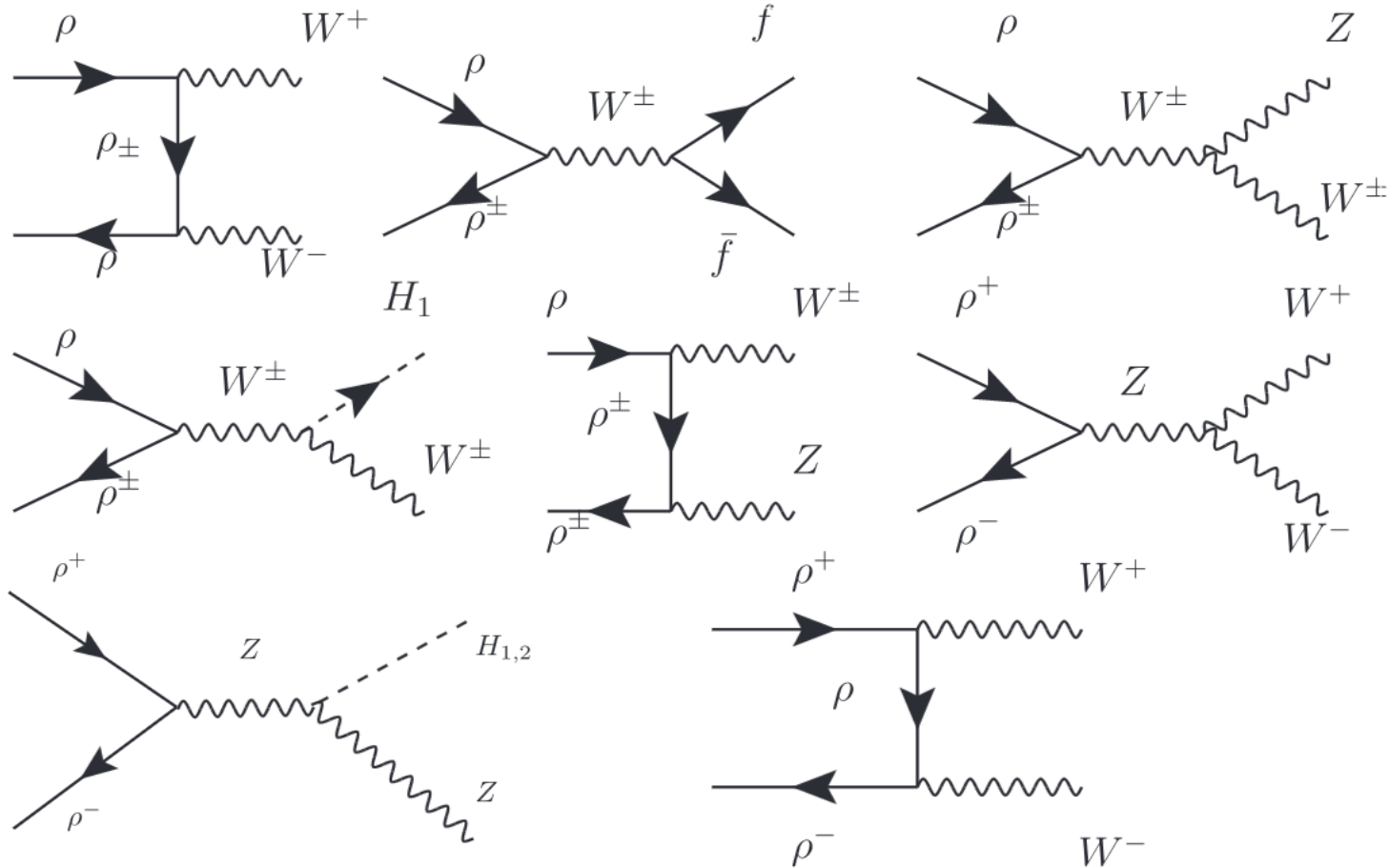
- WIMP Dark Matter
- DM is in thermal bath
- Annihilation of bath particles, decay of H_2 and late decay of N plays important role.

- FIMP Dark Matter
- N is non thermal
- Freeze-in production through decay and annihilation of bath particles.
- Late Decay of ρ also contributes substantially.

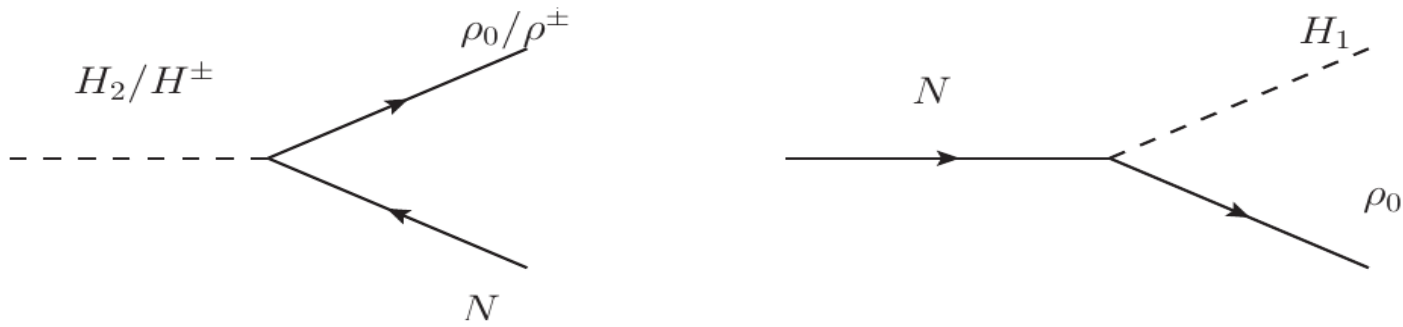
Scenario I

$$M_N > M_\rho$$

ρ is DM candidate.



Annihilation and co-annihilation channels of ρ in the early Universe .



Feynmann diag. for the dominant production of N as well its late decay to DM.

Boltzmann Equation for DM and NLOP:

$$\hat{L}f_N = \sum_{i=1,2} \mathcal{C}^{H_i \rightarrow N\rho} + \mathcal{C}^{AB \rightarrow N\rho} + \mathcal{C}^{N \rightarrow all},$$

B.eqn to determine the distribution function of 'N'

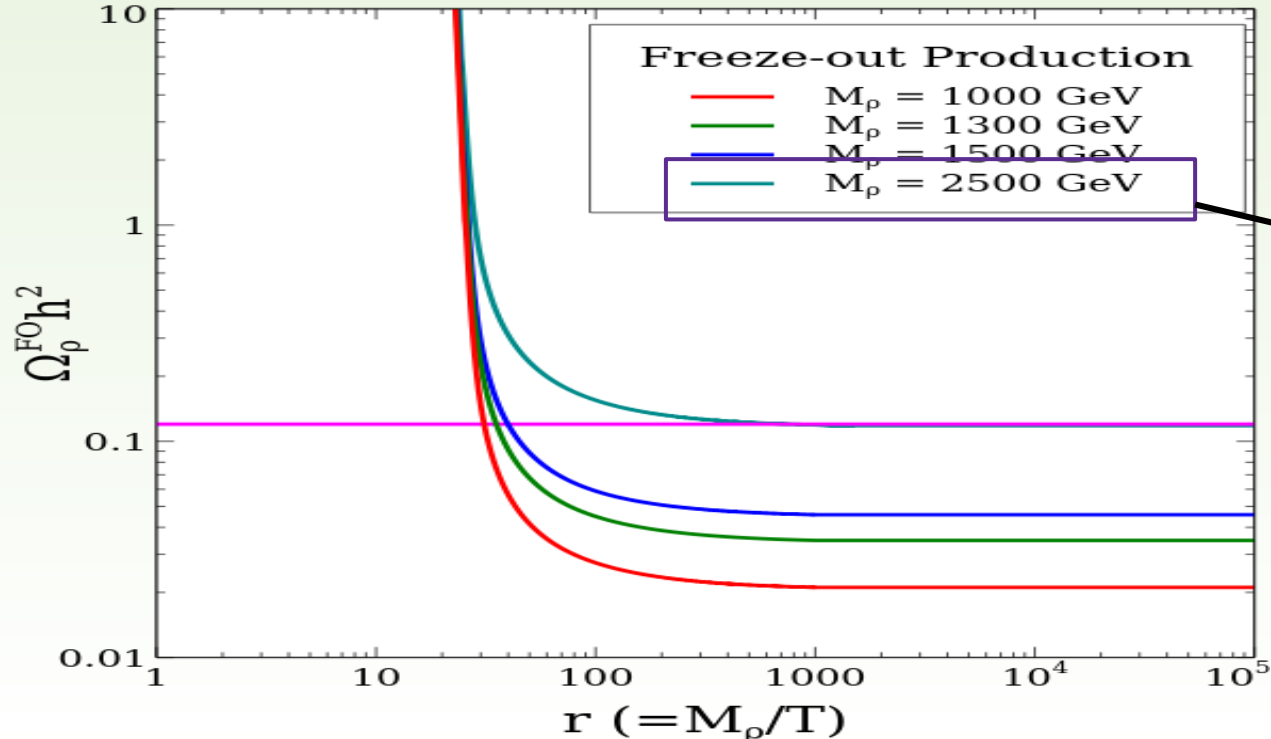
where,

$$\hat{L} = r H \left(1 + \frac{Tg'_s}{3g_s} \right)^{-1} \frac{\partial}{\partial r}$$

$$r = \frac{M_{sc}}{T}, \quad \xi_p = \left(\frac{g_s(T_0)}{g_s(T)} \right)^{1/3} \frac{p}{T}$$

B.eqn for the evolution of DM:

$$\frac{dY_\rho}{dr} = -\sqrt{\frac{\pi}{45G}} \frac{M_{Pl} \sqrt{g_*(r)}}{r^2} \langle \sigma_{eff} |v| \rangle (Y_\rho^2 - (Y_\rho^{eq})^2) + \frac{M_{Pl} r \sqrt{g_*(r)}}{1.66 M_{sc}^2 g_s(r)} [\langle \Gamma_{H_2 \rightarrow N\rho} \rangle (Y_{H_2} - Y_N Y_\rho) + \langle \Gamma_{N \rightarrow \rho A} \rangle_{NTH} (Y_N - Y_\rho Y_A)]$$

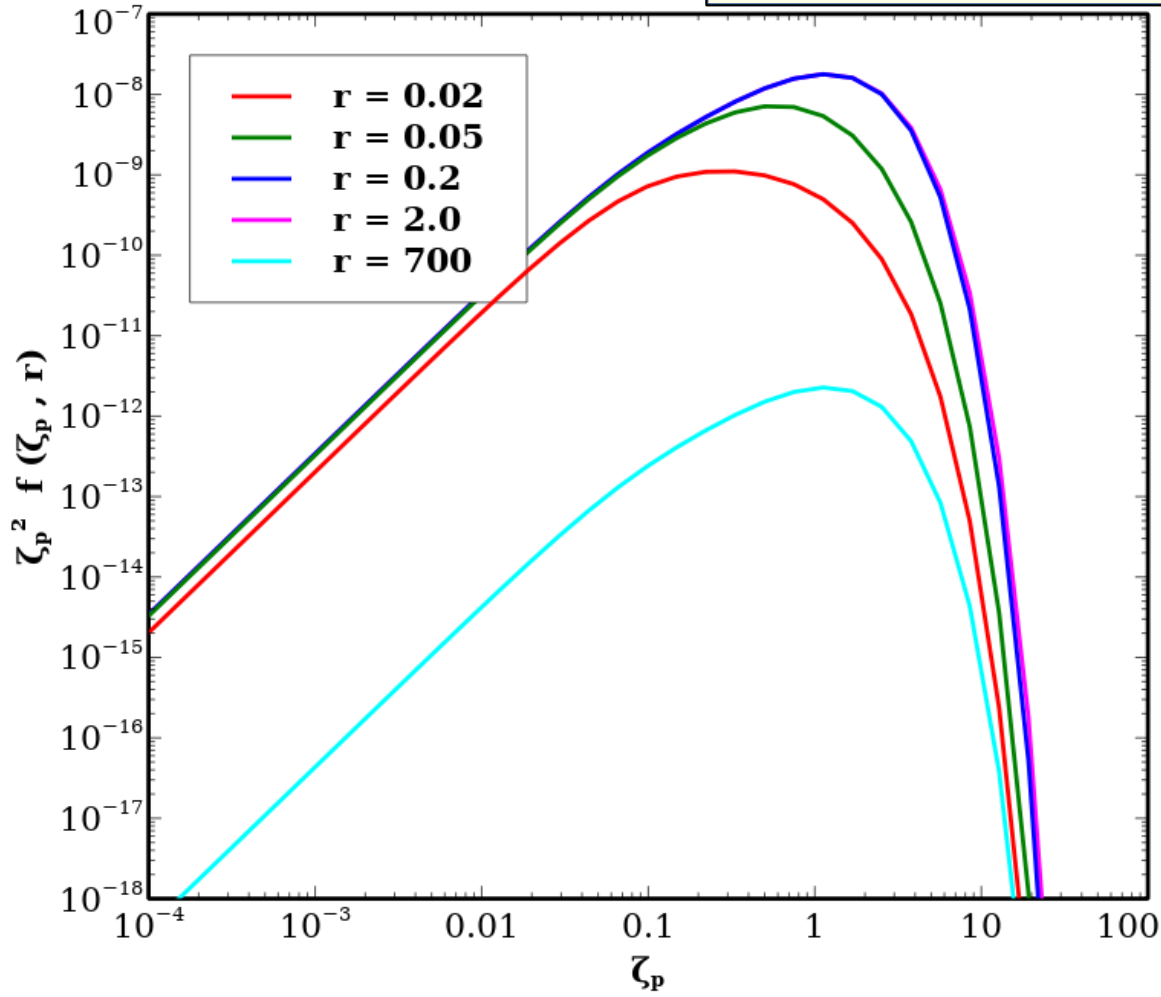


Satisfy the observed DM relic density

Evolution of distribution function for 'N'

Benchmark Point

$$M_N = 2000 \text{ GeV}, M_\rho = 1300 \text{ GeV}, Y_{\rho\Delta} = 2.5 \times 10^{-12}.$$



$$n_N(r) = \frac{gT^3}{2\pi^2} \mathcal{B}(r)^3 \int d\xi_p \xi_p^2 f_N(\xi_p, r)$$

This gives number density of 'N' at values of r .

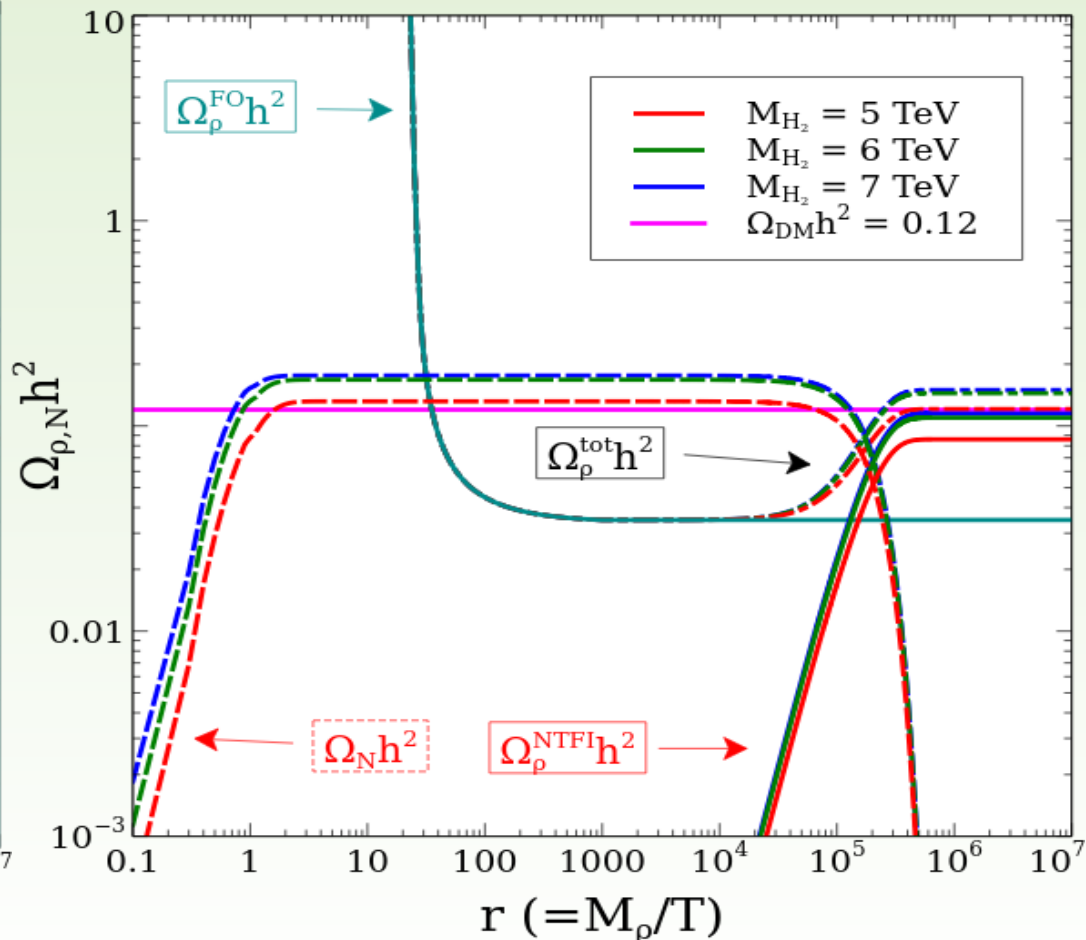
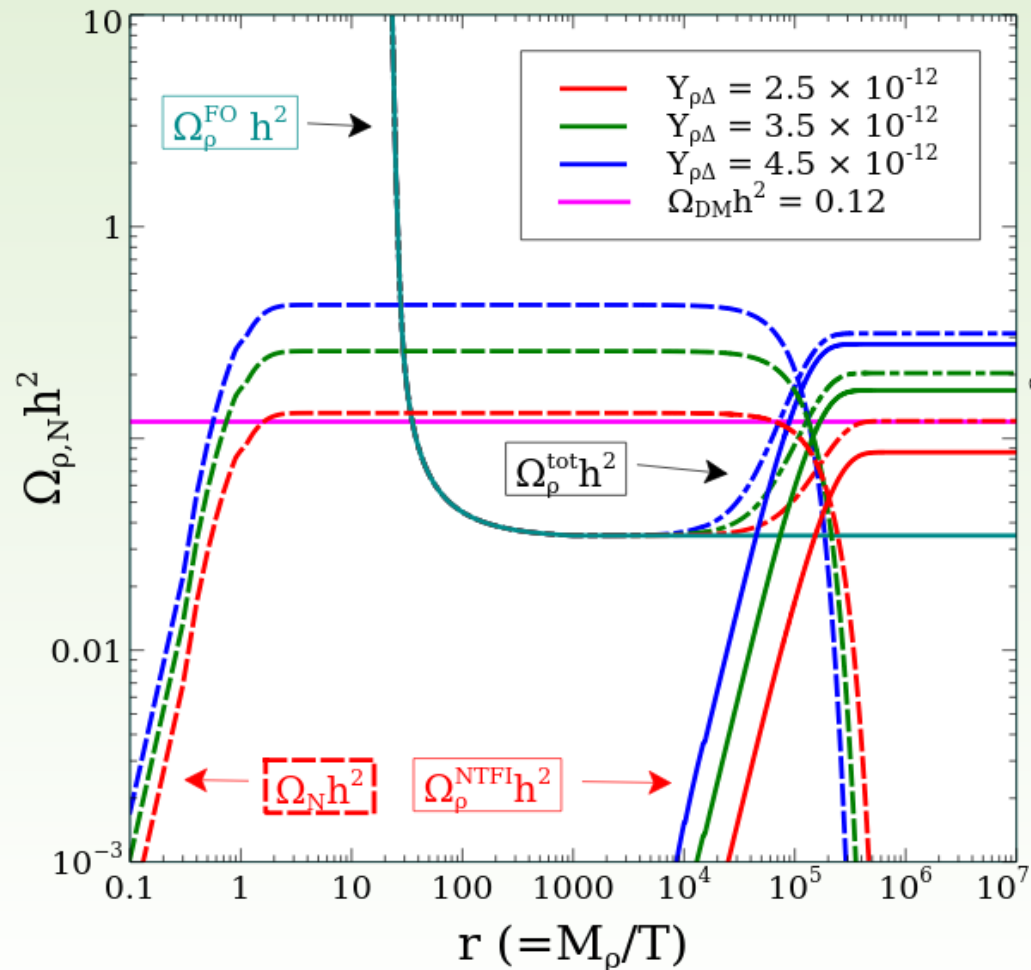
where,

$$\mathcal{B}(r) = \left(\frac{g_s(T_0)}{g_s(T)} \right)^{1/3} = \left(\frac{g_s(M_{sc}/r)}{g_s(M_{sc}/r_0)} \right)^{1/3}.$$

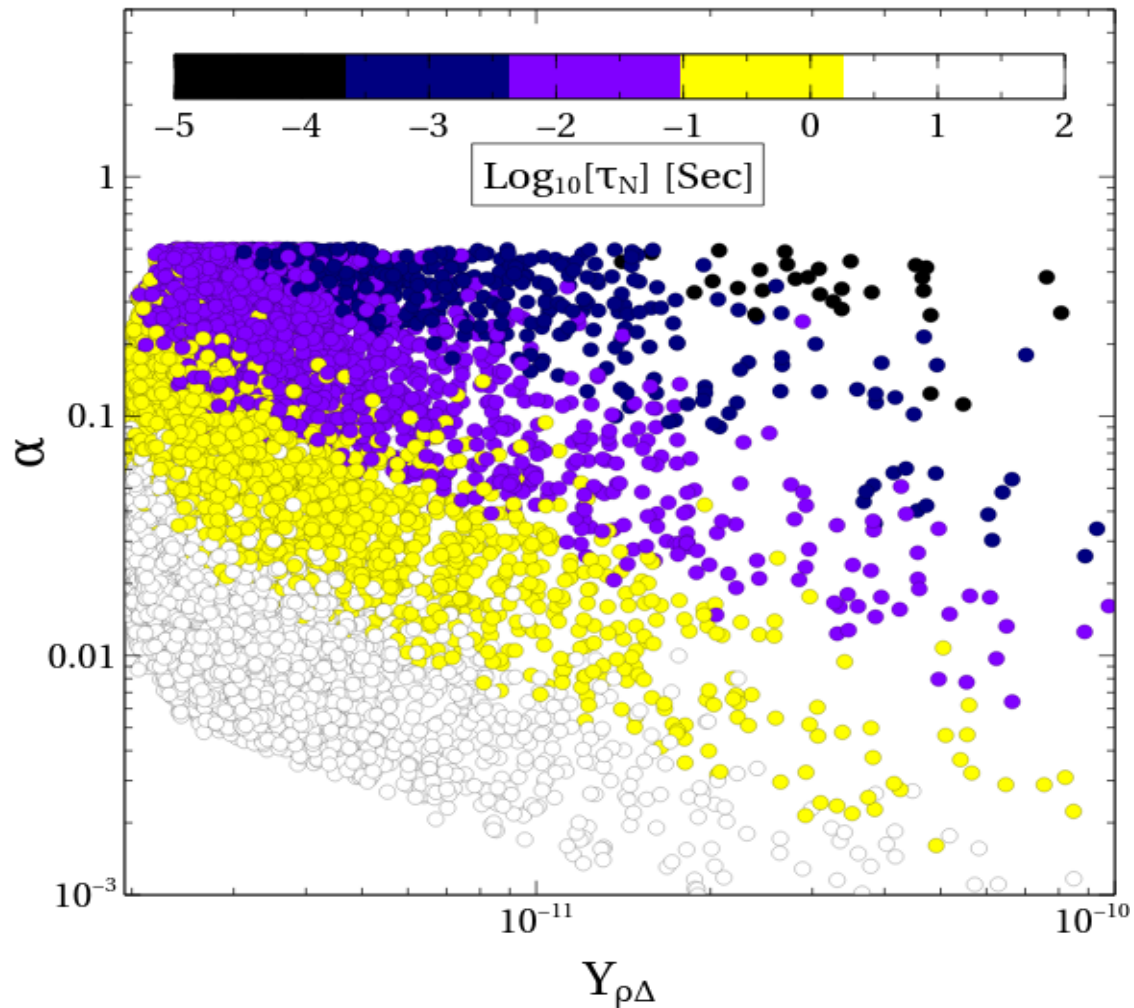
Results:-

$$M_N = 2000 \text{ GeV}, M_\rho = 1300 \text{ GeV}, Y_{\rho\Delta} = 2.5 \times 10^{-12}.$$

**Parameter
chosen
unless varied**



Results:-



Parameters Varied

$$700 \text{ GeV} < M_\rho < 1500 \text{ GeV},$$

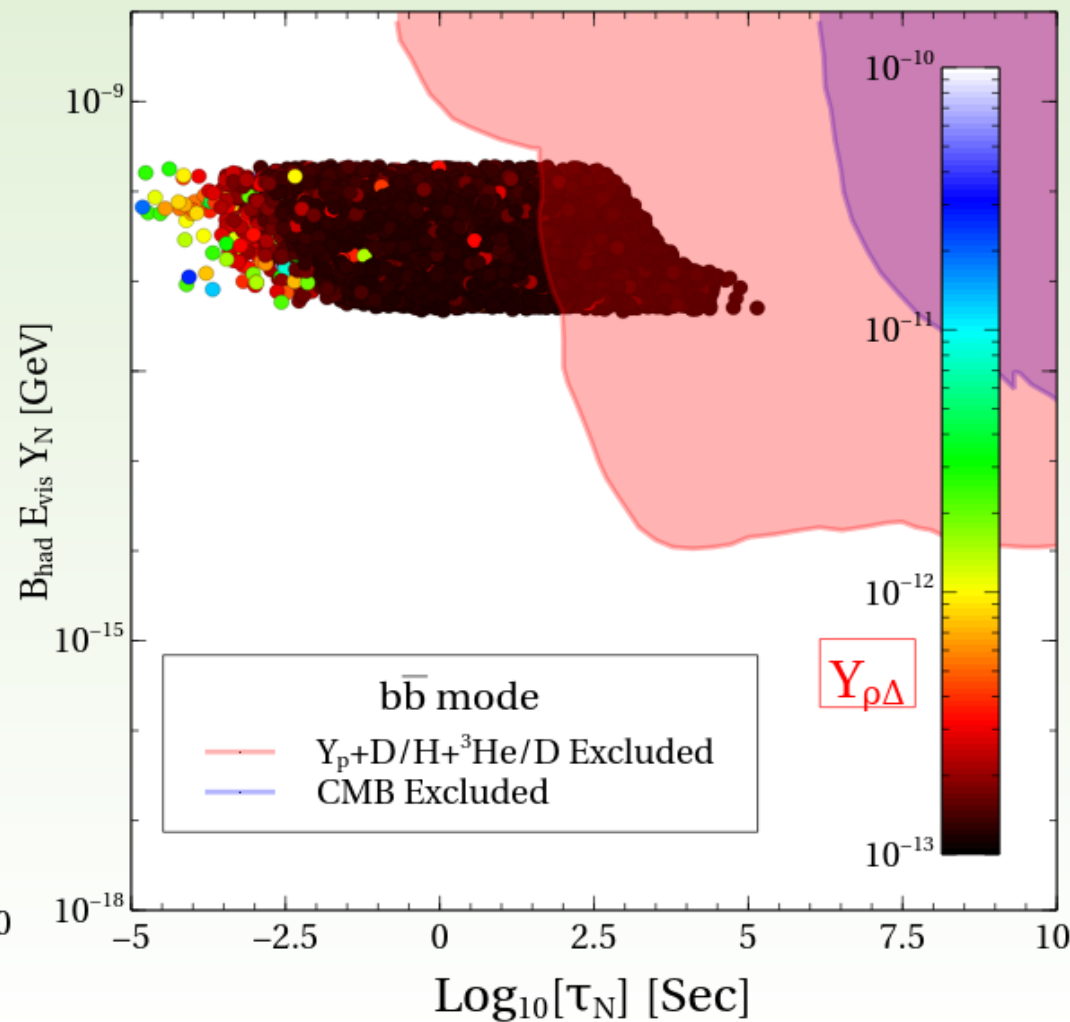
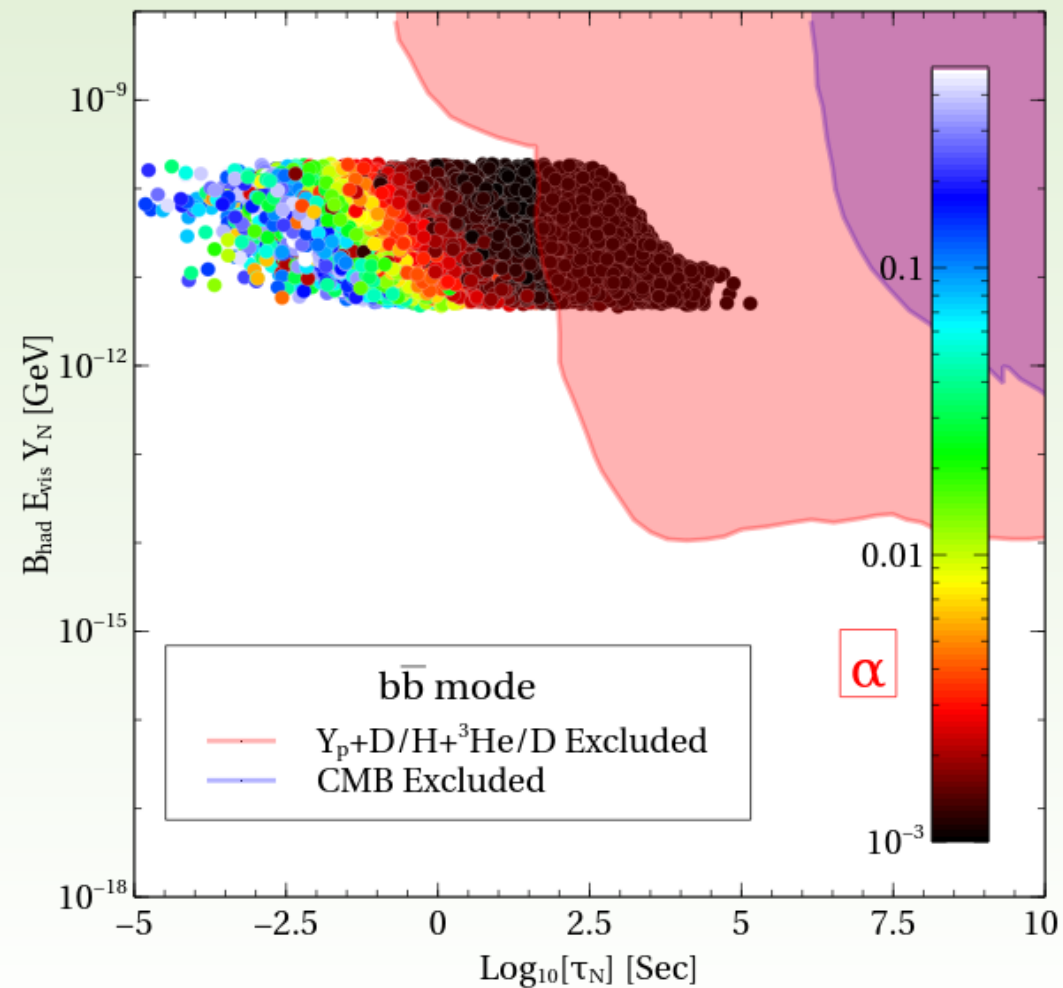
$$125 \text{ GeV} < M_N - M_\rho < 3000 \text{ GeV},$$

$$1500 \text{ GeV} < M_{H_2} < 20000 \text{ GeV},$$

$$10^{-13} < Y_{\rho\Delta} < 10^{-10}$$

$$10^{-3} < \alpha < 0.5.$$

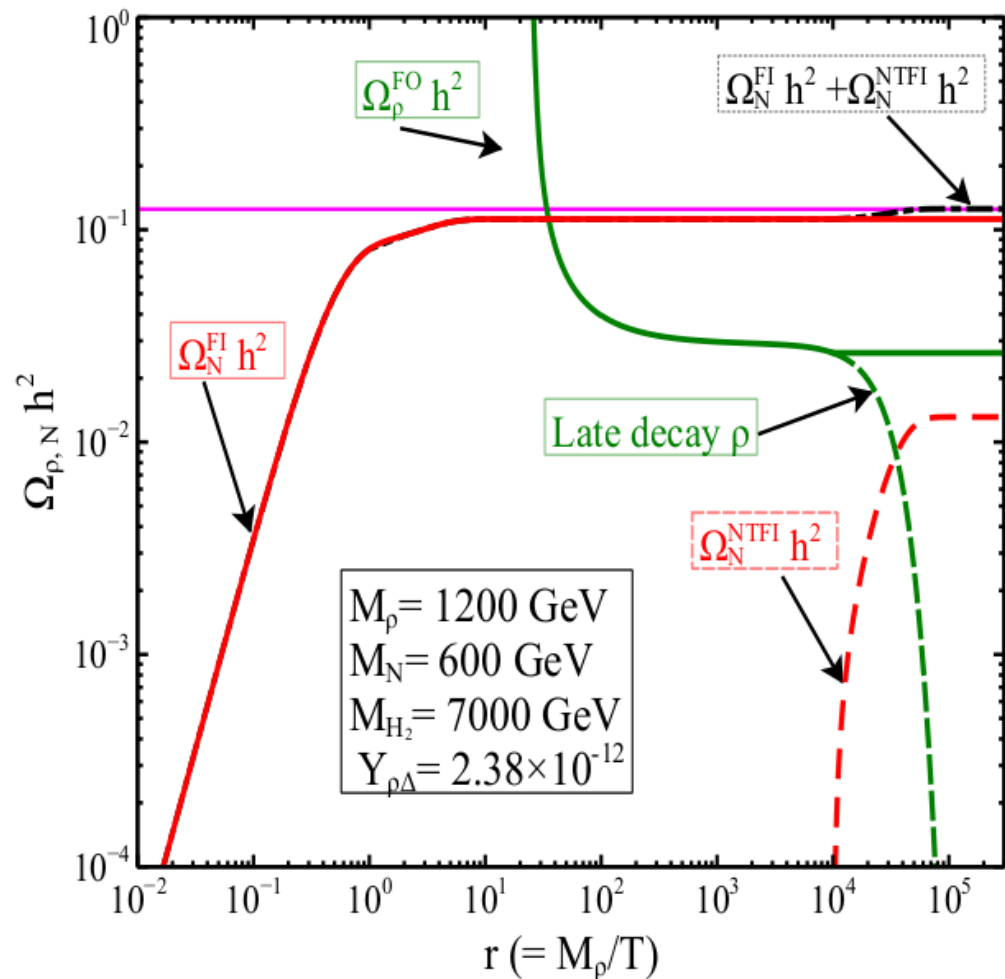
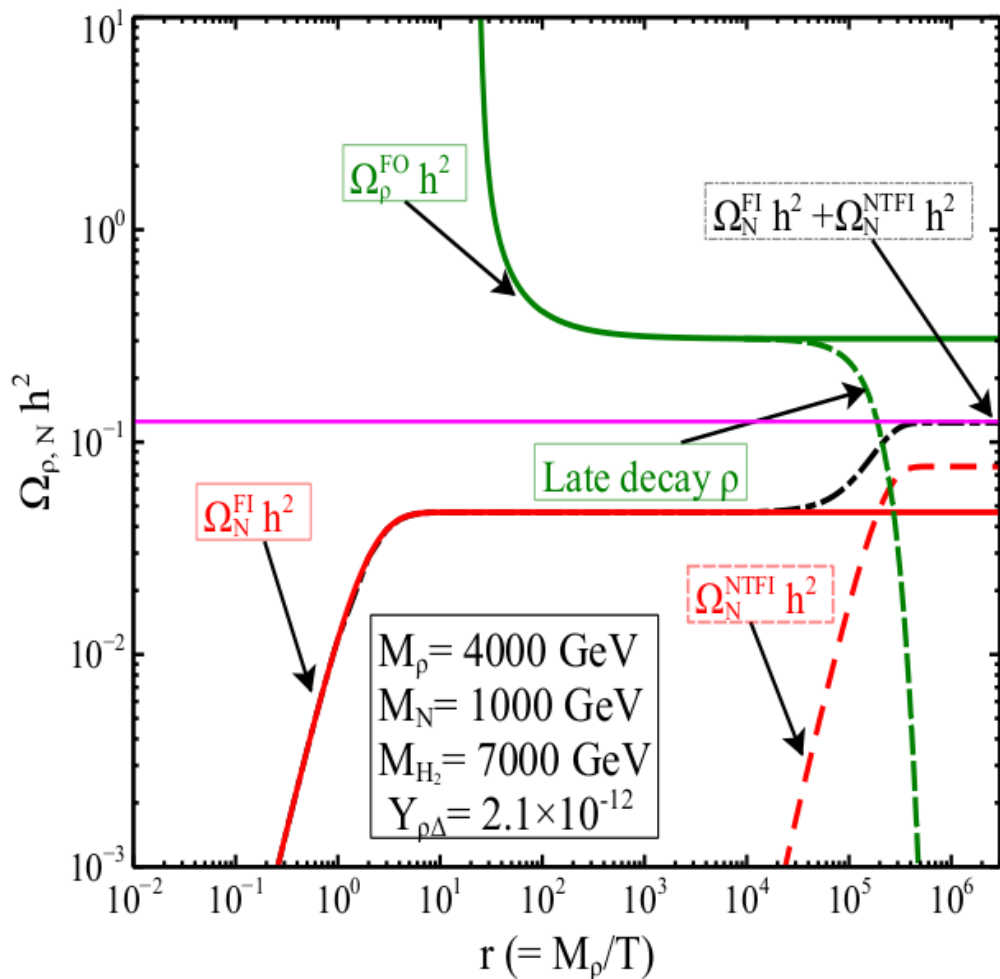
BBN Constraint



Scenario II

$$M_\rho > M_N$$

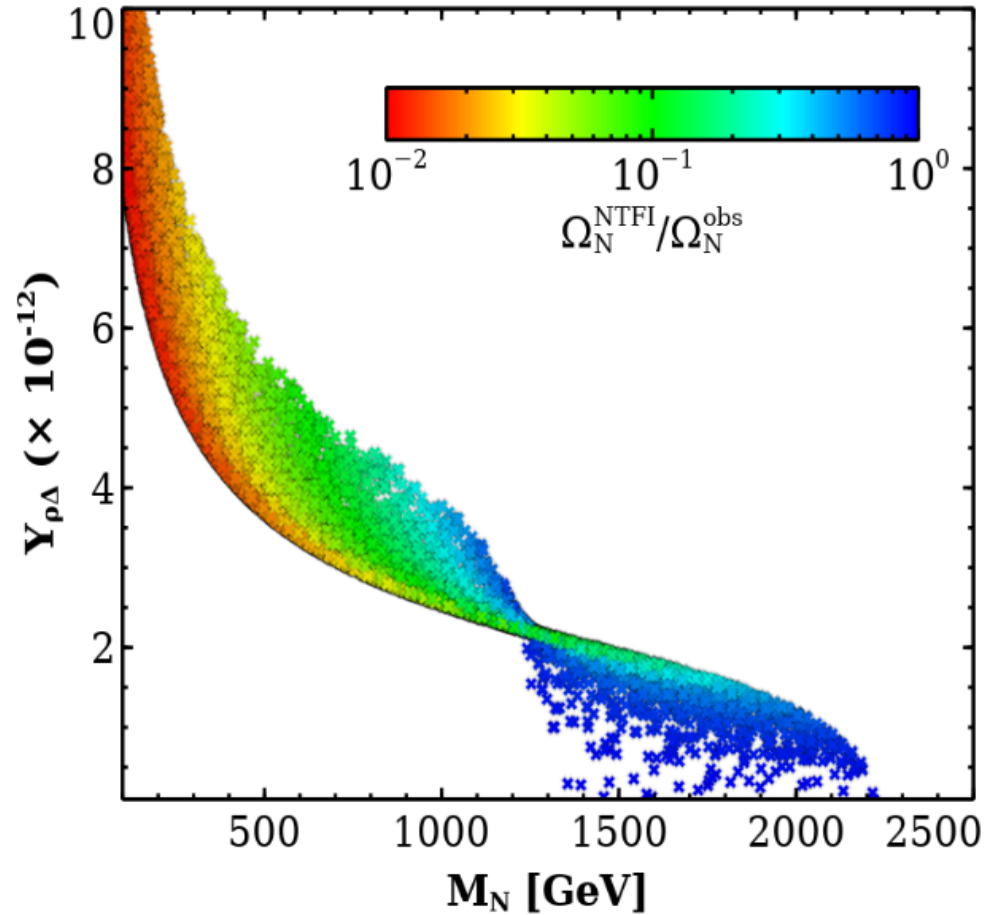
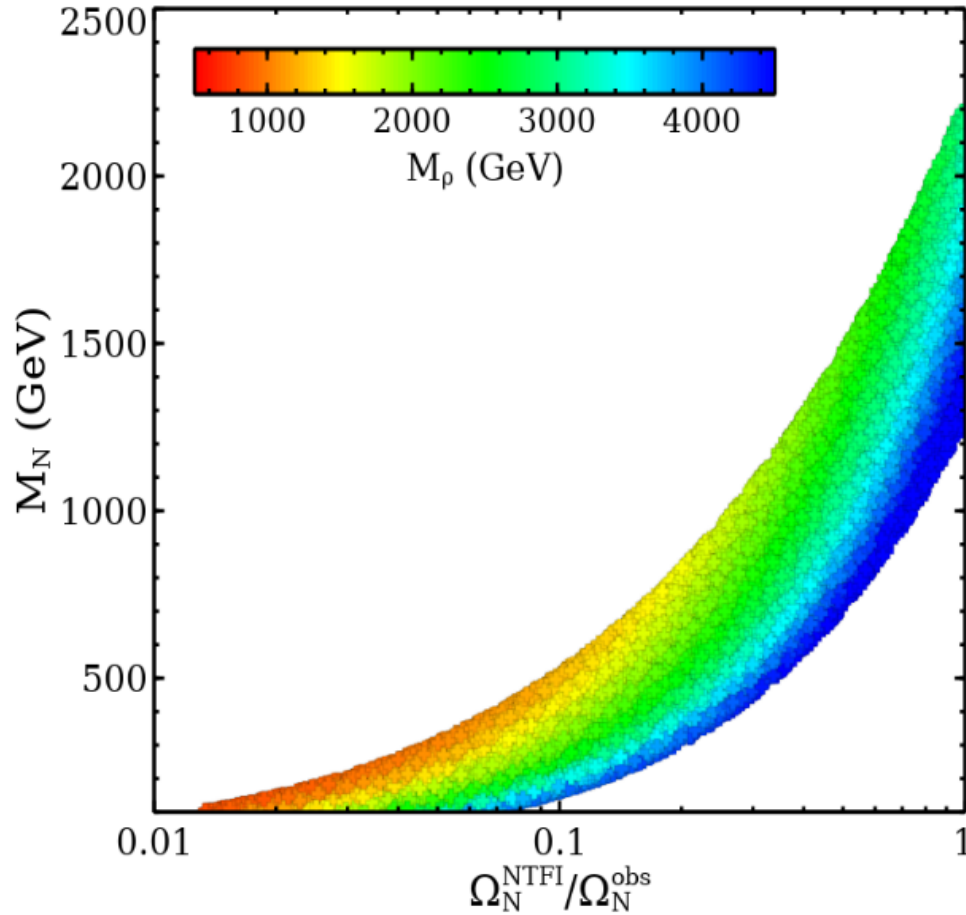
'N' is dark matter candidate.



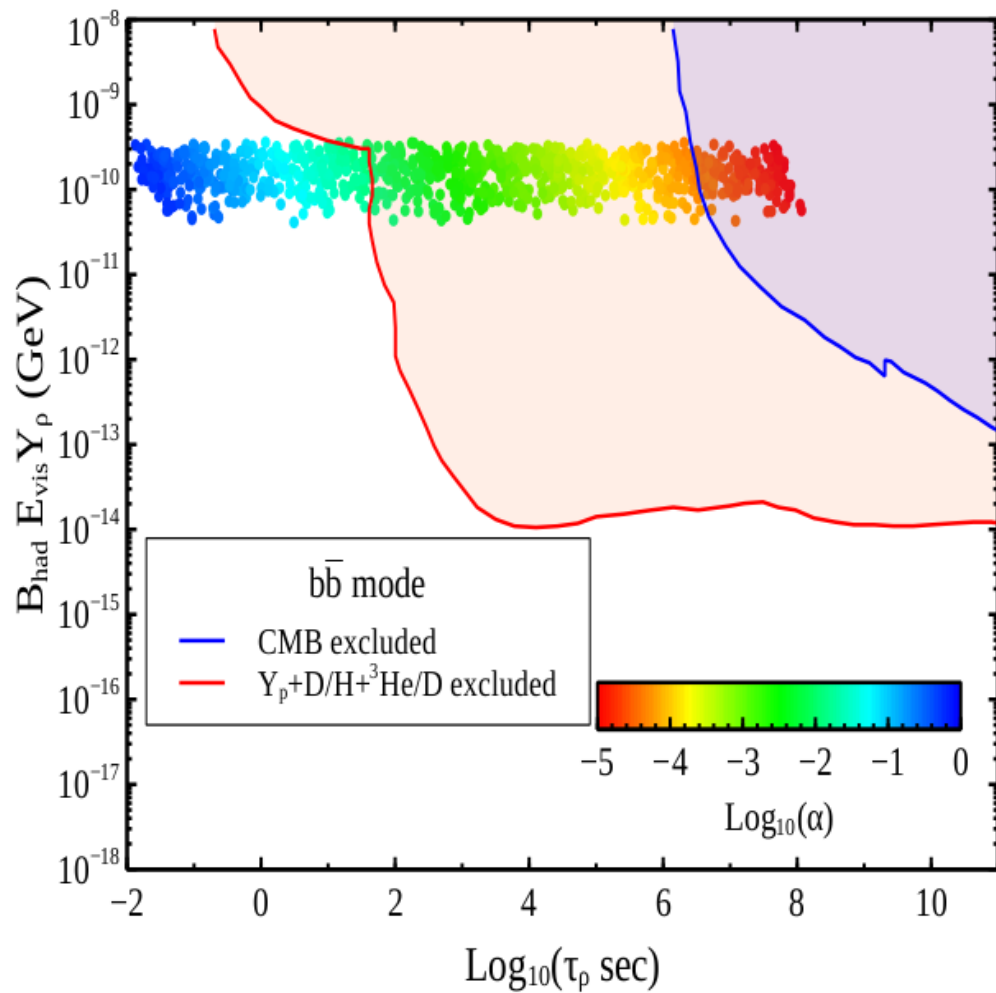
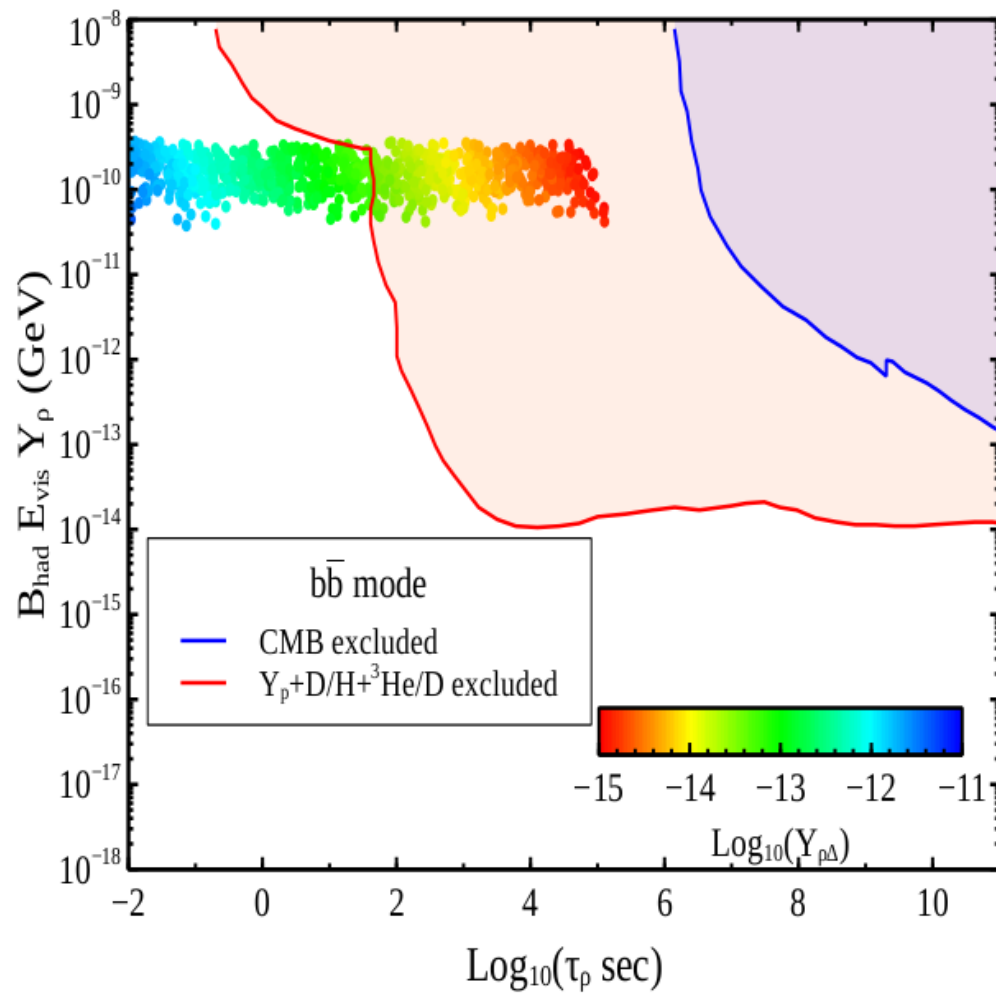
Results:-

Parameters Varied

$$10^{-11} < Y_{\rho\Delta} < 10^{-15}, 100 \text{ GeV} \leq M_N \leq 1800 \text{ GeV} \text{ and } 600 \text{ GeV} \leq M_\rho \leq 4500 \text{ GeV}$$

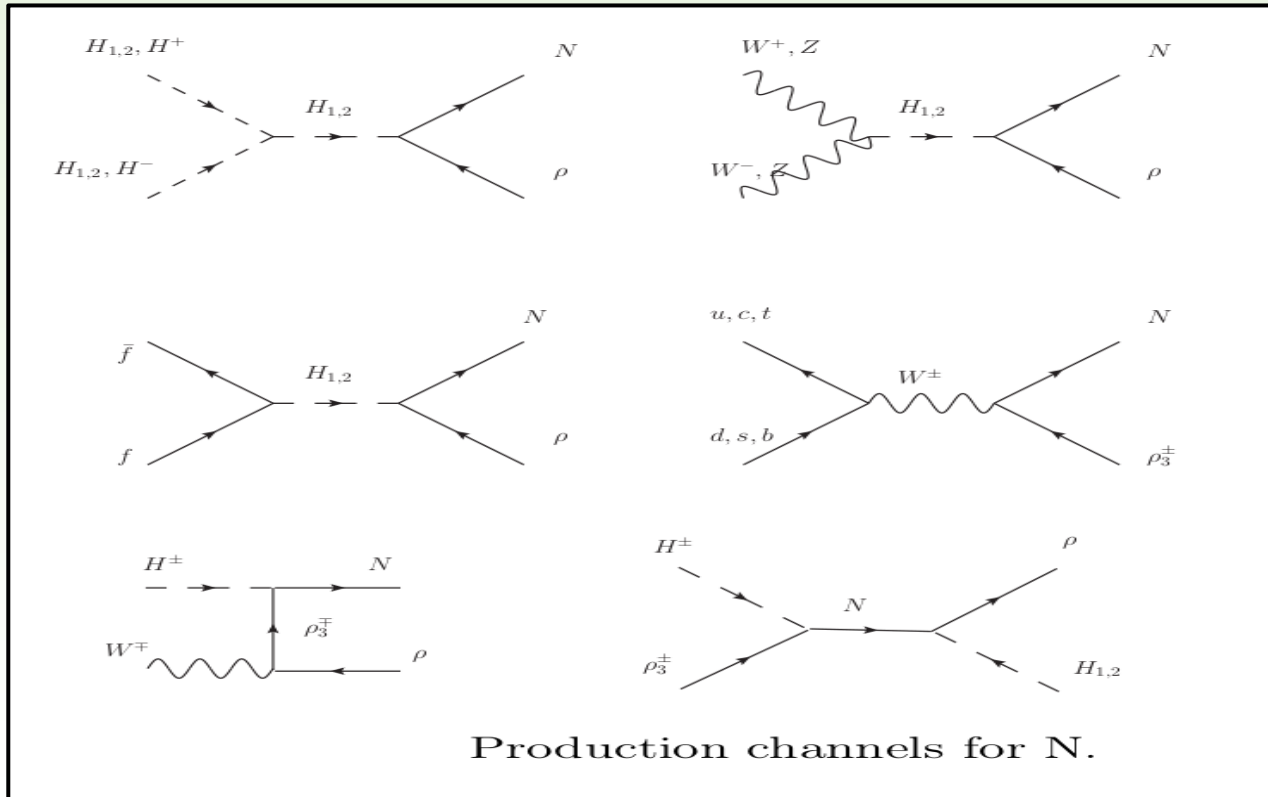


BBN Constraint:-

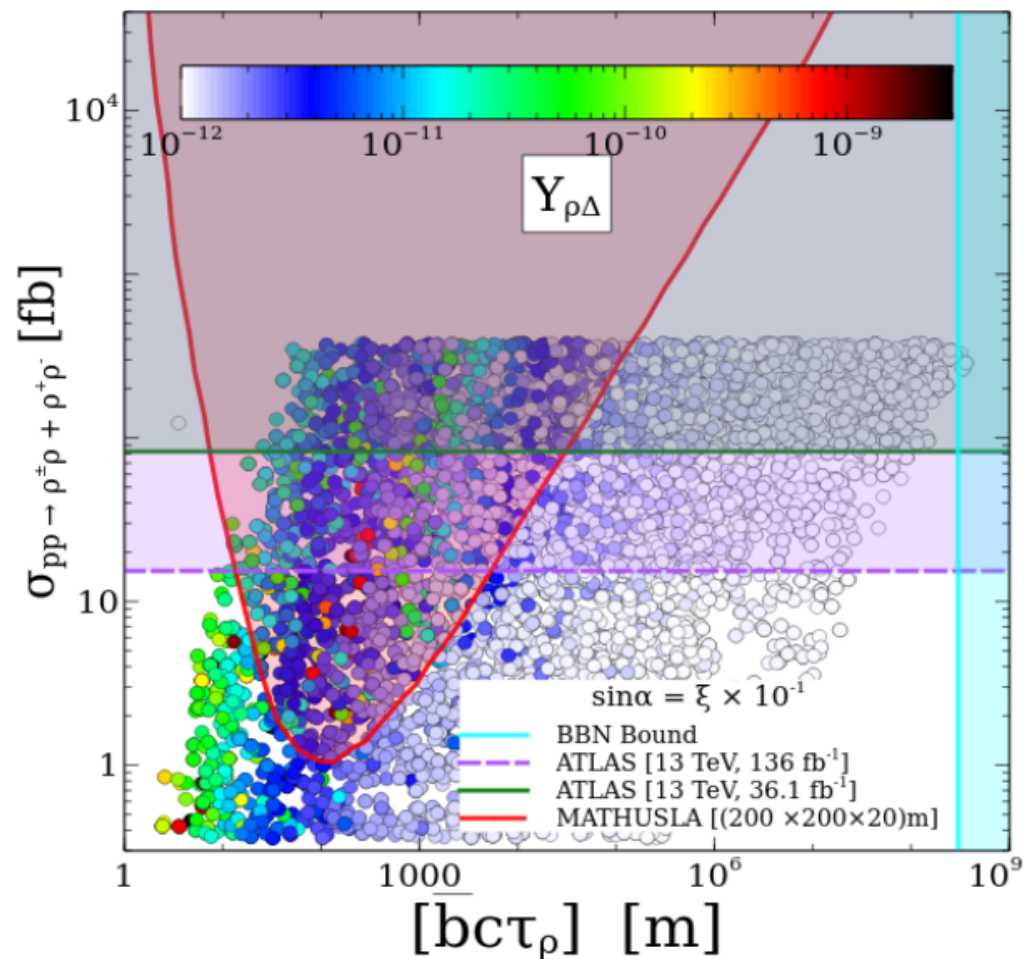
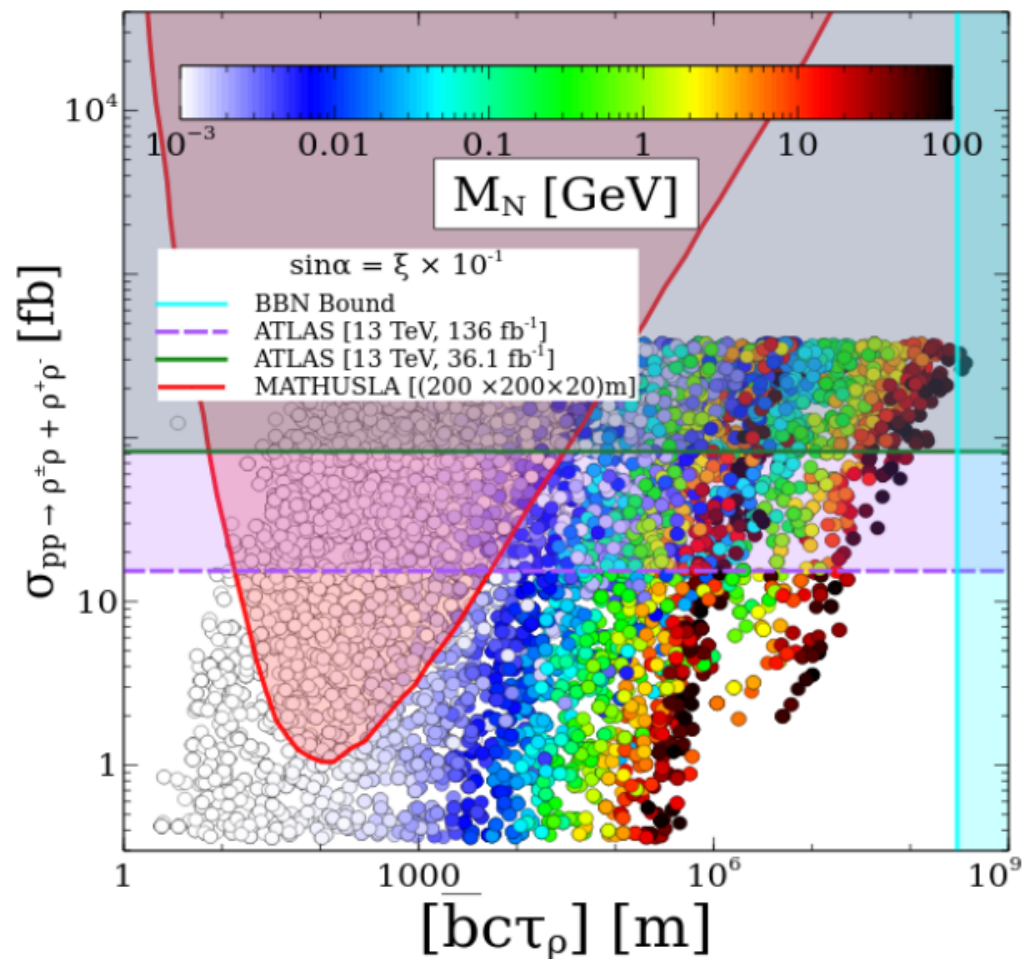


Results: allowing for a light scalar sector.

- In previous scenarios, 'N' is dominantly produced through decay at high temperature.
- Now, we assume 'N' is produced through annihilation of bath particles and production through decay is kinematically forbidden.



Substantial Annihilation Contribution: $M_N < M_\rho$ and $M_{H_2} < M_\rho + M_N$



Summary

- We extended SM with three fermion triplet, one singlet and one real triplet to explain DM and neutrino mass.
- We investigated different production mechanism for the production of DM.
- We also constrained our model parameters through BBN and found the model to be viable in large areas of parameter space.
- If the channel $H_2 \rightarrow \rho N$ is kinematically forbidden, observed dark matter relic density can also be realised with a few hundred GeV BSM Higgs and such scenario can be probed in MATHUSLA.



THANK YOU
for your
ATTENTION!