Searching for resonant flavor-changing charged Higgs production at the LHC

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Based on W.-S. Hou and M. Krab, arXiv:2409.18474

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General two Higgs doublet model

In the Higgs basis, the general *CP*-conserving 2HDM scalar potential is given by Davidson and Haber, PRD'05; Hou and Kikuchi, EPL'18

$$V(\Phi, \Phi') = \mu_{11}^{2} |\Phi|^{2} + \mu_{22}^{2} |\Phi'|^{2} - (\mu_{12}^{2} \Phi^{\dagger} \Phi' + \text{H.c.}) + \frac{\eta_{1}}{2} |\Phi|^{4} + \frac{\eta_{2}}{2} |\Phi'|^{4} + \eta_{3} |\Phi|^{2} |\Phi'|^{2} + \eta_{4} |\Phi^{\dagger} \Phi'|^{2} + \left[\frac{\eta_{5}}{2} (\Phi^{\dagger} \Phi')^{2} + (\eta_{6} |\Phi|^{2} + \eta_{7} |\Phi'|^{2}) \Phi^{\dagger} \Phi' + \text{H.c.} \right], \quad (1)$$

with

$$\Phi = \begin{pmatrix} G^+ \\ (v+h_1+iG^0)/\sqrt{2} \end{pmatrix}, \qquad \Phi' = \begin{pmatrix} H^+ \\ (h_2+iA)/\sqrt{2} \end{pmatrix}.$$
 (2)

- \triangleright The usual Z_2 symmetry is dropped \implies FCNC at tree-level
- Many parameters and extra processes arise
- \triangleright EWBG, Absence of FCNC (e.g. $t \rightarrow ch_{125}$), ... could be explained
- ▷ Sub-TeV H, A, H^{\pm} bosons may still exist

General Yukawa interaction

Higgs-fermion interactions can be described by

Davidson and Haber, PRD'05

$$\mathcal{L}_{Y} = -\frac{1}{\sqrt{2}} \sum_{f=u,d,\ell} \bar{f}_{i} \left[\left(\lambda_{ij}^{f} s_{\gamma} + \rho_{ij}^{f} c_{\gamma} \right) h + \left(\lambda_{ij}^{f} c_{\gamma} - \rho_{ij}^{f} s_{\gamma} \right) H - i \operatorname{sgn}(Q_{f}) \rho_{ij}^{f} A \right] P_{R} f_{j} - \bar{u}_{i} \left[(V \rho^{d})_{ij} P_{R} - (\rho^{u\dagger} V)_{ij} P_{L} \right] d_{j} H^{+} - \bar{\nu}_{i} \rho_{ij}^{\ell} P_{R} \ell_{j} H^{+} + \operatorname{H.c.}$$
(3)

- $\triangleright \ \lambda^f$ matrices: diagonal, fixed by fermion mass
- $\triangleright
 ho^f$ matrices: non-diagonal (and in general complex) lead to FCNC
- ▷ Alignment ($c_{\gamma} \approx 0$) suppresses FCNC for h but allows FCNC for H and A
- $\triangleright \
 ho_{ij}$ are severely constrained by flavor physics
- ▷ Extra top couplings ρ_{tc} and ρ_{tt} could be $\mathcal{O}(1)$ and can each drive EWBG Fuyuto, Hou, Seneha, PLB'18

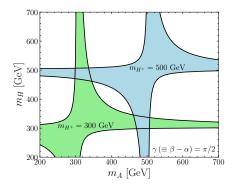
 \triangleright For simplicity, we set all $\rho_{ij} = 0$ except ρ_{tc} and ρ_{tt}

Constraints on G2DHM

G2HDM parameter space is subject to the following constraints:

- > Unitarity, perturbativity and vacuum stability
- \triangleright EWPD through oblique parameters S, T and U using the following fit result:

 $S = 0.05 \pm 0.08$, $T = 0.09 \pm 0.07$, $\rho_{ST} = 0.92$, [PDG]

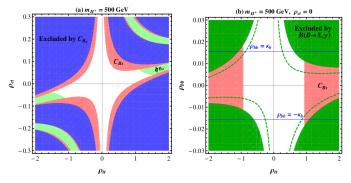


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Flavor physics and direct searches (next slides)

Flavor constraints

Because of a $|V_{cq}/V_{tq}|$ (q = d, s) (m_t/m_b) enhancement factor, ρ_{ct} (ρ_{bb}) is severely constrained by $B_q - \bar{B}_q$ $(b \rightarrow s\gamma)$.



B. Altunkaynak et al., PLB'15

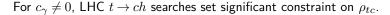
 \triangleright Constraints on ρ_{tc} are weak due to small m_c . An upper bound on ρ_{tc} was found to be $|\rho_{tc}| \lesssim 1.3 \ (1.7)$ for $m_{H^+} = 300 \ (500)$ GeV.

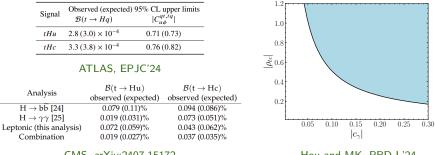
A. Crivellin et al., PRD'13

 $\triangleright \rho_{tc}$ and ρ_{tt} can still be sizable ($\lesssim O(1)$) under current data

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Limit from $t \rightarrow ch$ searches





CMS, arXiv:2407.15172

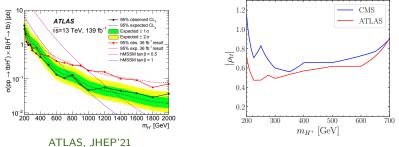
Hou and MK, PRD-L'24

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- $|\rho_{tc}| \gtrsim 0.5$ is excluded at 95% CL for $c_{\gamma} = 0.1$
- \triangleright The limit diminishes for $c_{\gamma} < 0.1$ and vanishes for $c_{\gamma} = 0$ (alignment)

Limit from $H^+ \rightarrow t\bar{b}$ searches

LHC searches for $pp \rightarrow \bar{t}bH^+ \rightarrow \bar{t}bt\bar{b}$ strongly constrain ρ_{tt} .

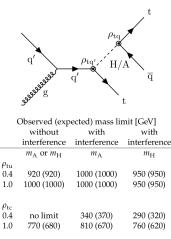


Hou and MK, PRD-L'24

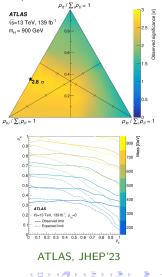
- \triangleright Limits are interpreted assuming $\mathcal{B}(H^+ \to t \bar{b}) = 100\%$
- ▷ Constraints from LHC searches for $pp \rightarrow H/A \rightarrow t\bar{t}$ and $pp \rightarrow t\bar{t}H/A \rightarrow t\bar{t}t\bar{t}$ are relatively weaker
- $\triangleright \rho_{tt}$ is safe from constraints from SM Higgs properties $(s_{\gamma} = 1)$

Search for G2HDM neutral Higgs bosons

With $t \to ch$ alignment-suppressed, it is natural to pursue $cg \to tH/tA \to tt\bar{c}/tt\bar{t}$ (same-sign top/triple top), which is controlled by $s_{\gamma} \simeq 1$.



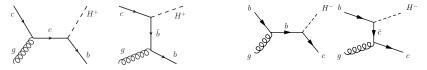
CMS, PLB'24



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Searching for H^+ with FC couplings

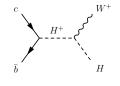
In G2HDM, where $\bar{c}bH^+$ couples with strength $\rho_{tc}V_{tb}$, $cg \rightarrow bH^+$ and $bg \rightarrow cH^-$ are not CKM-suppressed, compared to 2HDM-II.



Ghosh, Hou, Modak, PRL'20

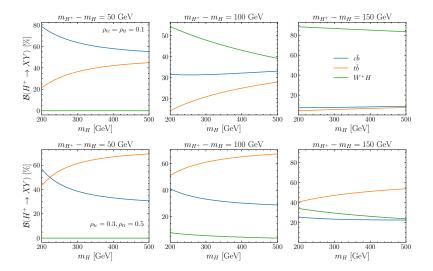
Hou and MK, PRD-L'24

 $c\bar{b} \rightarrow H^+ \rightarrow W^+H$, which goes through the same $\bar{c}bH^+$ coupling of $\rho_{tc}V_{tb}$, is suggested as a new avenue for discovering H^+ at the LHC.



Hou and MK, arXiv:2409.18474

Charged Higgs decay



Signal vs. Background

Signal: $c\bar{b} \rightarrow H^+ \rightarrow W^+H \rightarrow \ell^+\nu t (\rightarrow \ell^+\nu b)\bar{c} \rightarrow \ell^+\ell^+ + \nu\nu + b\bar{c}$ BKG: $t\bar{t}V$ (V = W,Z), tZj, t $\bar{t}h$, 4t, tW, WZ, ZZ

BP	η_2	η_3	η_4	η_5	η_7	m_H	m_A	m_{H^+}	μ_{22}^2/v^2
1	1.40	2.00	-0.82	-0.82	-0.55	200	300	300	0.49
2	2.88	4.75	-2.64	-2.64	0.16	300	500	500	1.75

Table: For BP1, $\rho_{tc} = \rho_{tt} = 0.1$, while for BP2, $\rho_{tc} = 0.3$, $\rho_{tt} = 0.5$.

Simulation: MadGraph5_aMC@NLO ($\sqrt{s} = 14 \text{ TeV}$) + Pythia + Delphes

$\triangleright N_j \ge 2$ with $P_T^j \ge 20$ GeV		
$V I V_j \ge 2$ with $I T \ge 20$ GeV	Background	Cross section
\triangleright At least one <i>b</i> -tagged ($N_b \ge 1$)	tW	1.61
\sim $ne \ o \ tagged (1) \geq 1)$	$tar{t}W$	1.09
$p_{1}(2) = p_{1}^{\ell}(2) + p_{2}^{\ell}(2) + p_{3}^{\ell}(2) + p_{3}^{\ell$	WZ	0.54
▷ SS2ℓ ($N_{\ell} = 2$), $P_T^{\ell_{1(2)}} \ge 25(20)$ GeV	tZj	0.40
	$t\bar{t}Z$	0.10
$\triangleright \Delta R_{\ell\ell}, \Delta R_{\ell j} > 0.4, E_T^{\text{miss}} > 35 \text{ GeV}$	tth	0.05
	ZZ	0.02
p _T sum of all jets and two SS	4t	0.0004
leptons $H_T < 400$ GeV	Q-flip	0.0018
$T_T < 400$ GeV	Fake	0.0002

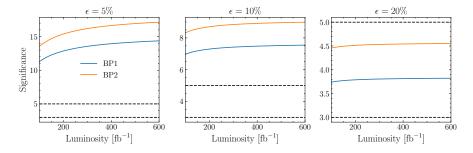
Significance

We estimate our signal sensitivity using

Kumar and Martin, PRD'15

$$\mathcal{Z} = \sqrt{2\left[(S+B)\ln\left(\frac{(S+B)\left(B+\Delta_B^2\right)}{B^2+(S+B)\Delta_B^2}\right) - \frac{B^2}{\Delta_B^2} + \ln\left(1 + \frac{\Delta_B^2 S}{B(B+\Delta_B^2)}\right) \right]},$$

with $\Delta_B = \epsilon B$, where S(B) is number of signal (background) events, and ϵ refers to systematic uncertainty in background estimation.



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Conclusion

- ▷ Charged Higgs bosons are actively searched for at the LHC.
- \triangleright However, it might be difficult to detect at the LHC via $bg \rightarrow tH^- \rightarrow t\bar{t}b$.
- ▷ In G2HDM, resonant $c\bar{b} \rightarrow H^+$ is induced by the FC top-charm coupling ρ_{tc} without CKM-suppression, and has a large cross section.
- \triangleright Our proposed $c\bar{b}\to H^+\to W^+H(\to t\bar{c})$ signal, with its same sign dilepton signature, could be promising.

Thank you!

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Thank you!