

# Boost-invariant spin hydrodynamics with spin feedback effect



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This work was supported in part by the Polish National Science Centre (NCN) Grants No. 2022/47/B/ST2/01372 and No.2018/30/E/ST2/00432. We study one - dimensional and boost invariant expansion with second - order corrections in the spin polarization tensor to the energy-momentum tensor and baryon current without dissipative components. Consequently, we obtain feedback from spin dynamics on the hydrodynamic background which constrains possible spin polarization configurations. However, for a small magnitude of the spin polarization tensor (below unity in natural units), the permitted spin dynamics differ very little from that found in the case without the second-order corrections (Ref. [2]).

⇒ **Motivation:** theoretical description of the RHIC spin polarization data (Ref. [1]).

We use the form of polarization tensor with four - vectors decomposed in the orthonormal basis

$$\omega_{\mu\nu} = k_\mu U_\nu - k_\nu U_\mu + t_{\mu\nu}, \quad t_{\mu\nu} = \epsilon_{\mu\nu\alpha\beta} U^\alpha \omega^\beta, \quad (1)$$

$$k^\mu = C_{kx} X^\mu + C_{ky} Y^\mu + C_{kz} Z^\mu, \quad \omega^\mu = C_{\omega z} X^\mu + C_{\omega y} Y^\mu + C_{\omega x} Z^\mu, \quad t^\mu = V_x X^\mu + V_y Y^\mu + V_z Z^\mu. \quad (2)$$

With the tensor forms calculated in Ref. [3] and Ref. [4], we consider the conservation laws for baryon number current, energy - momentum tensor, and spin tensor respectively and as a result, we obtain equations

$$\partial_\mu N^\mu = \frac{d\bar{n}}{d\tau} + \frac{\bar{n}}{\tau} = 0, \quad (3) \quad \partial_\mu T^{\mu\nu} = \left[ \dot{\bar{\epsilon}} + \frac{\bar{\epsilon} + \bar{P}}{\tau} + \frac{P}{\tau} (C_{kz}^2 + C_{\omega z}^2) \right] U^\nu + \left[ \left( \dot{P}_t + \frac{P_t}{\tau} \right) V_x + P_t \dot{V}_x \right] X^\nu + \left[ \left( \dot{P}_t + \frac{P_t}{\tau} \right) V_y + P_t \dot{V}_y \right] Y^\nu + \left[ \left( \dot{P}_t + \frac{P_t}{\tau} \right) V_z + P_t \left( \dot{V}_z + \frac{V_z}{\tau} \right) \right] Z^\nu = 0, \quad (4)$$

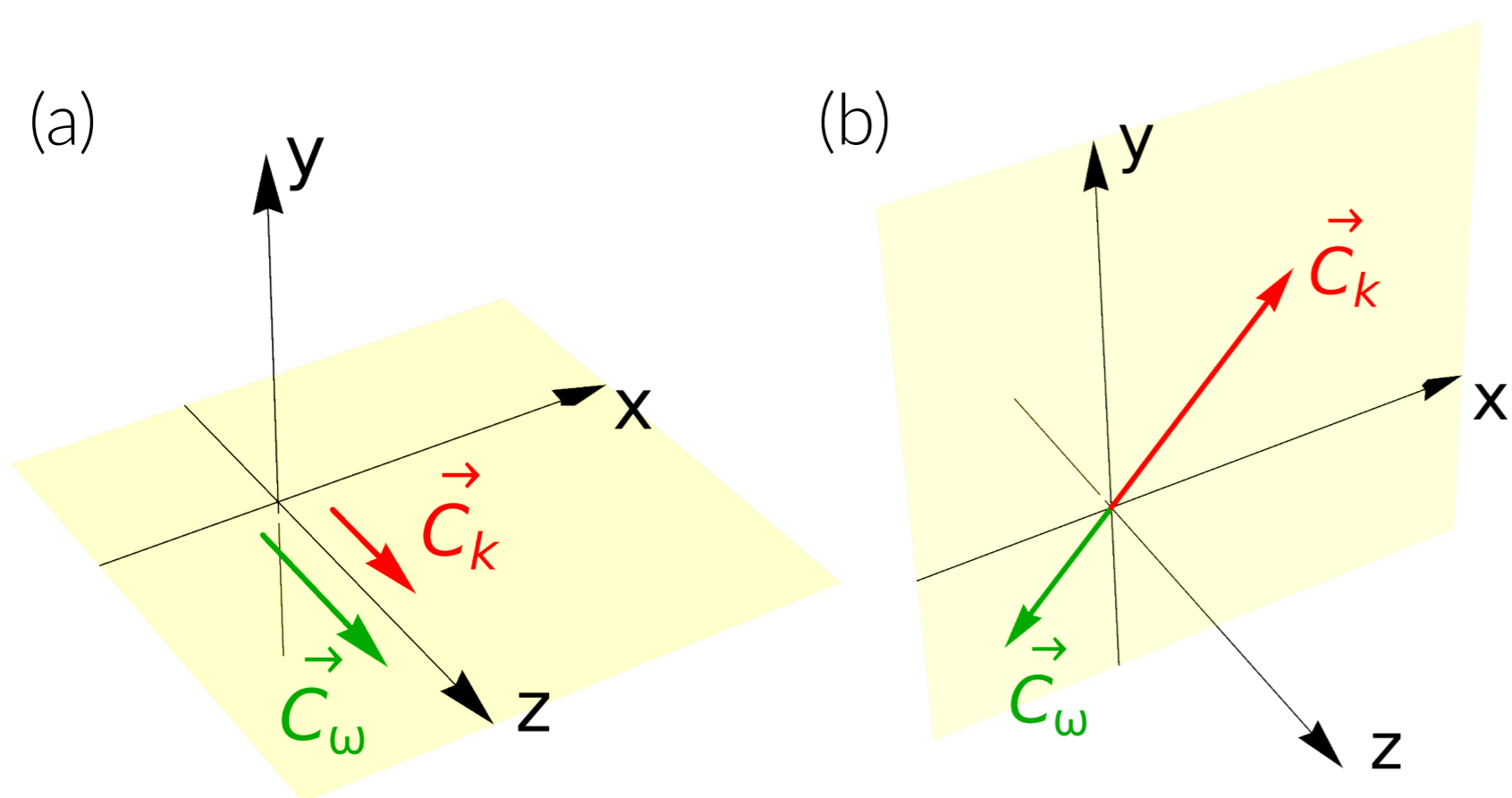


Figure 1: Schematic view for longitudinal (a) and transverse (b) configuration.

$$\partial_\lambda S^{\lambda,\mu\nu} = \frac{1}{\tau} [A (k^\mu U^\nu - k^\nu U^\mu) + A_1 t^{\mu\nu}] + (A k^\mu) U^\nu - (A k^\nu) U^\mu + \dot{A}_1 t^{\mu\nu} + A_1 \dot{t}^{\mu\nu} + \frac{1}{2} \partial_\lambda A_3 (t^{\lambda\mu} U^\nu - t^{\lambda\nu} U^\mu) + \frac{A_3}{2} [(\partial_\lambda t^{\lambda\mu}) U^\nu + t^{\lambda\mu} \partial_\lambda U^\nu - (\partial_\lambda t^{\lambda\nu}) U^\mu - t^{\lambda\nu} \partial_\lambda U^\mu - \frac{U^\mu}{\tau} k^\nu + \Delta^{\lambda\mu} \partial_\lambda k^\nu + \frac{U^\nu}{\tau} k^\mu - \Delta^{\lambda\nu} \partial_\lambda k^\mu] = 0. \quad (5)$$

## LONGITUDINAL CONFIGURATION

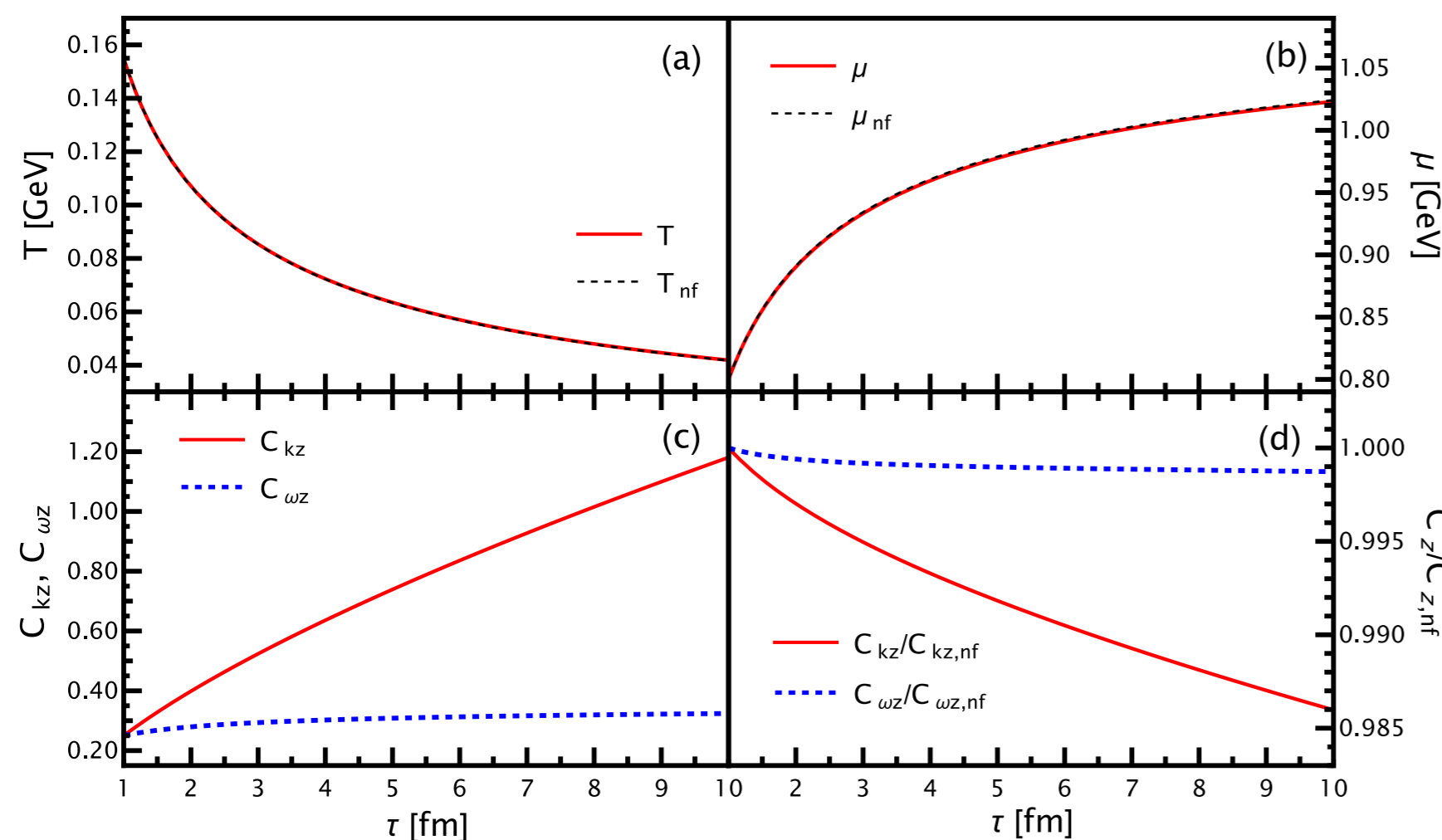


Figure 2: Longitudinal configuration with the initial values  $C_{kz}^0 = C_{\omega z}^0 = 0.25$ .

## TRANSVERSE CONFIGURATION

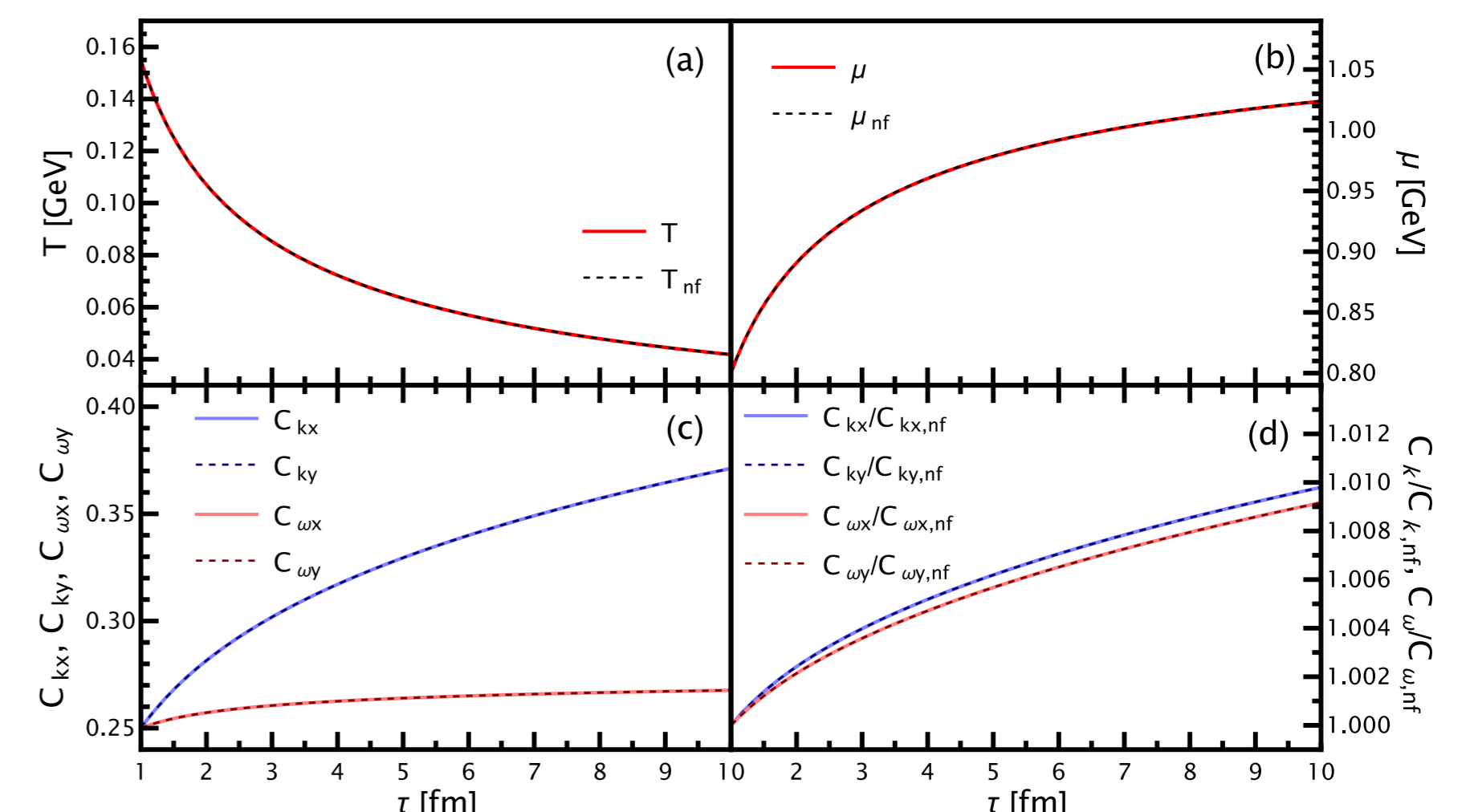


Figure 4: Transverse configuration with the initial values  $C_{kx}^0 = C_{ky}^0 = C_{\omega x}^0 = C_{\omega y}^0 = 0.25$ .

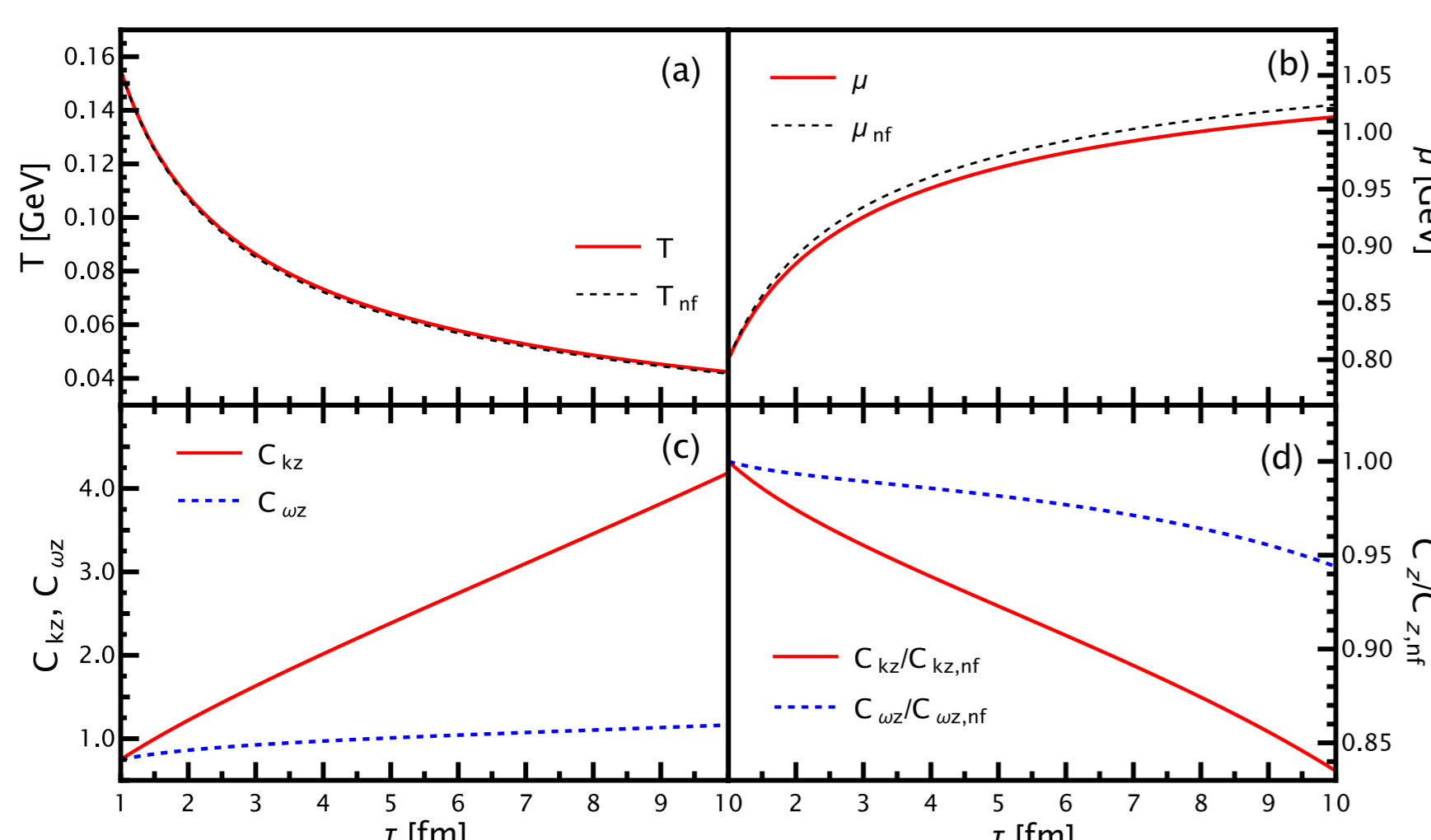


Figure 3: Longitudinal configuration with the initial values  $C_{kz}^0 = C_{\omega z}^0 = 0.75$ .

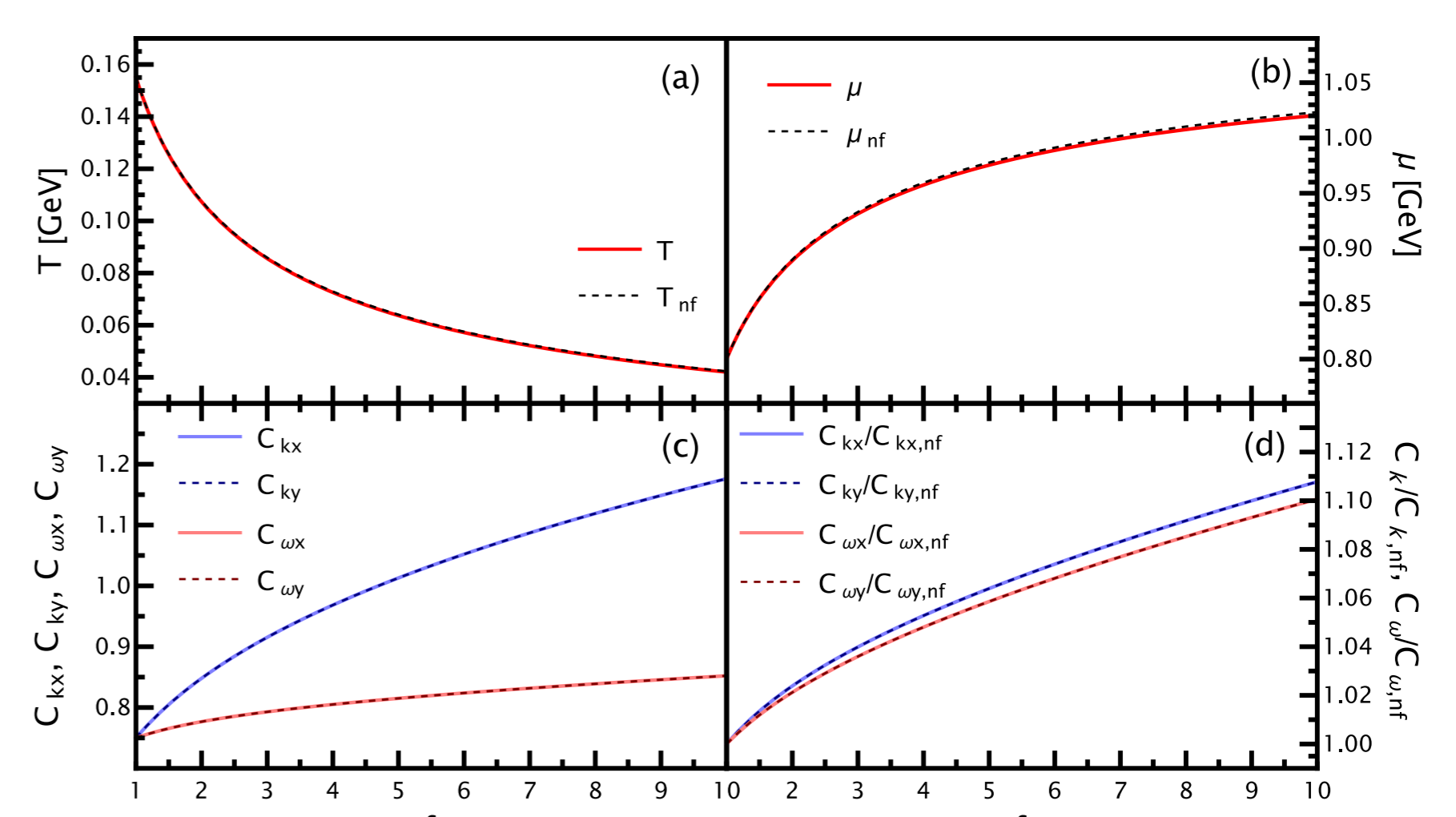


Figure 5: Transverse configuration with the initial values  $C_{kx}^0 = C_{ky}^0 = C_{\omega x}^0 = C_{\omega y}^0 = 0.75$ .

Panels describe proper-time dependence of (a) temperature  $T$ , (b) baryon chemical potential  $\mu$ , (c)  $C$  coefficients and (d) their ratios to no-feedback results.

## References

- [1] Francesco Becattini, Michael Lisa, *Polarization and Vorticity in the Quark Gluon Plasma*, Ann. Rev. Nucl. Part. Sci. 70 (2020) 395-423, arXiv:2003.03640.
- [2] W. Florkowski, A. Kumar, R. Ryblewski, and R. Singh, *Spin polarization evolution in a boost invariant hydrodynamical background*, Phys. Rev. C 99, 044910 (2019), arXiv:1901.09655.
- [3] W. Florkowski and M. Hontarenko, *Generalized thermodynamic relations for perfect spin hydrodynamics*, (2024), arXiv:2405.03263.
- [4] Z. Drogosz, W. Florkowski, and M. Hontarenko, *Hybrid approach to perfect and dissipative spin hydrodynamics*, (2024), arXiv:2408.03106.