

Probing New Bosons & Nuclear Structure With Ytterbium Isotope Shifts



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QuantumFrontiers
Cluster of Excellence



Leibniz
Universität
Hannover

based on work with

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Nuclear Shape and BSM Searches at Colliders, CERN, 13th January 2025

Outline

Isotope Shifts

King Plots

Ytterbium King Plot

Input For Nuclear Physics

Sensitivity To New Physics

Outline

Isotope Shifts

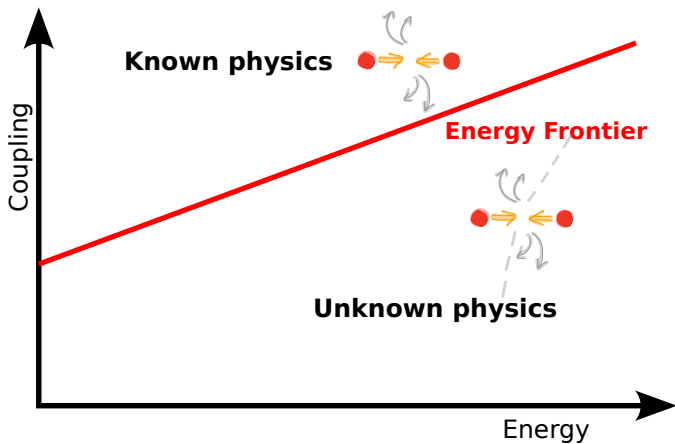
King Plots

Ytterbium King Plot

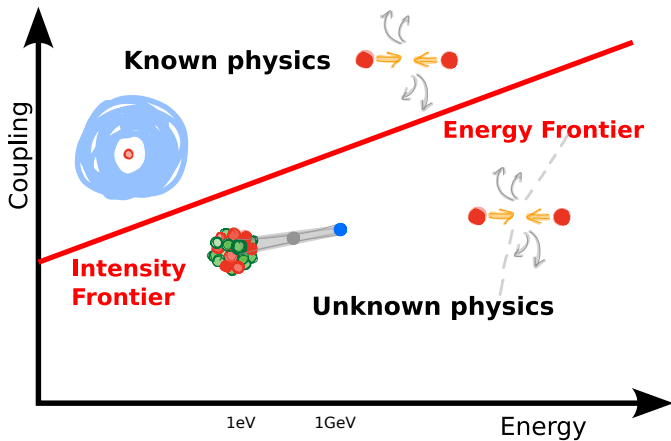
Input For Nuclear Physics

Sensitivity To New Physics

Where is the New Physics?



Where is the New Physics?



Why Isotope Shifts?

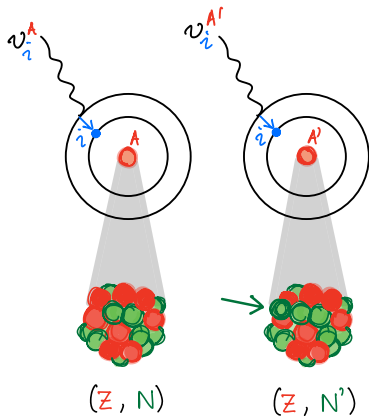
The most accurately measured numbers in physics are ratios of atomic clock transition frequencies:

- $\nu_{\text{Al}^+}/\nu_{\text{Hg}^+} = 1.052871833148990438(55)$ (NIST; $\sigma_\nu/\nu \sim 5.2 \times 10^{-17}$)
[Rosenband et al. *Science* 319, 1808 (2008)]
- $\nu_{\text{Yb}}/\nu_{\text{Sr}} = 1.207507039343337749(55)$ (RIKEN; $\sigma_\nu/\nu \sim 4.6 \times 10^{-17}$)
[Nemitz et al. *Nat. Photonics* 10, 258 (2016)]
- $\nu_{\text{E3}}/\nu_{\text{E2}} = 0.932829404530965376(32)$ (PTB; $\sigma_\nu/\nu \sim 3.4 \times 10^{-17}$)
[Lange et al. *PRL* 126 011102 (2021)]
- $\nu_{\text{In}^+}/\nu_{\text{Yb}^+} = 1.973773591557215789(9)$ (PTB; $\sigma_\nu/\nu \sim 4.4 \times 10^{-18}$)
[Hausser et al. *arXiv: 2402.16807* (2024)]

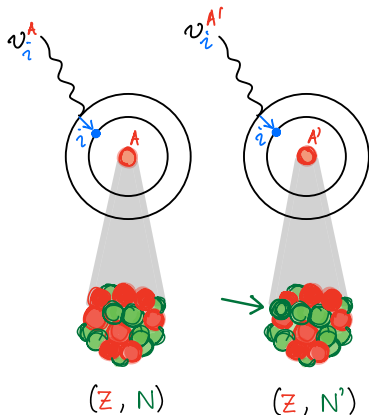
⇒ These are sensitive to “everything”, but we cannot calculate the spectrum below around 1% accuracy.

So what can we do with these?

Isotope Shifts



Isotope Shifts



Isotope shifts:

$$\begin{aligned} \nu_i^{AA'} &\equiv \nu_i^A - \nu_i^{A'} \\ &= K_i \mu^{AA'} + F_i \delta \langle r^2 \rangle^{AA'} + \dots \end{aligned}$$

i : transition index

AA' : isotope pair index

K_i, F_i, \dots : electronic coeffs.

$\mu^{AA'}, \delta \langle r^2 \rangle^{AA'}, \dots$: nuclear coeffs.

Z : number of protons

N, N' : number of neutrons in A, A'

Isotope Shifts: Mass Shift & Field Shift

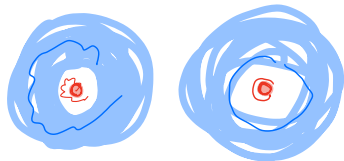
$$\nu_i^{AA'} = K_i \mu^{AA'} + F_i \delta \langle r^2 \rangle^{AA'} + \dots$$

Mass Shift

Nuclear motion in A vs. A'

⇒ Correction to e^- kin. energy

$$\propto \mu^{AA'} = \frac{1}{M^A} - \frac{1}{M^{A'}}$$



Isotope Shifts: Mass Shift & Field Shift

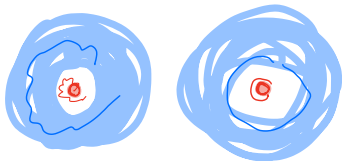
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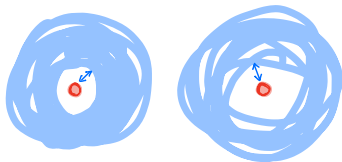
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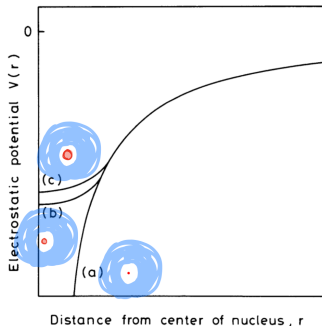


Field Shift

$$\propto \delta \langle r^2 \rangle^{AA'} = \langle r^2 \rangle^A - \langle r^2 \rangle^{A'}$$



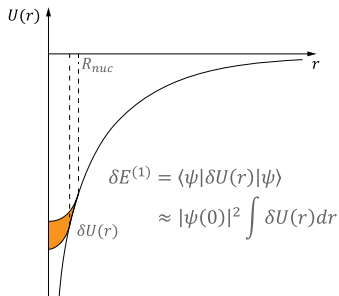
Field Shift: Nuclear Size Effect



- (a) Coulomb $V = -\frac{Ze}{4\pi r}$
(b) Finite size nucleus
(c) Larger nucleus

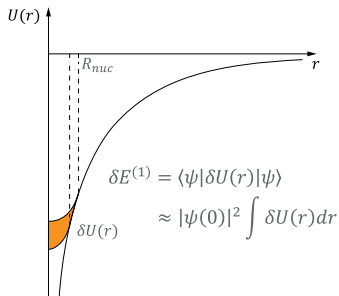
- **Inside the atom**, electron wavefct. affected by non-Coulombic nuclear potential, dep. on
 - Radial coordinate r
 - Nuclear **charge radius** $\langle r^2 \rangle = \frac{\int \rho_N(\mathbf{r})r^2 d\mathbf{r}^3}{\int \rho_N(\mathbf{r})d\mathbf{r}^3}$

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- Radial coordinate r

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⇒ Shift in $\langle r^2 \rangle \Rightarrow$ Energy shift

$$\delta E_i \equiv F_i \delta \langle r^2 \rangle^{AA'}$$

F_i : (Electronic) field shift constant

$\delta \langle r^2 \rangle^{AA'}$: Charge radius variance

$$\delta \langle r^2 \rangle^{AA'} = \langle r^2 \rangle^A - \langle r^2 \rangle^{A'}$$

Isotope Shifts: Mass Shift & Field Shift

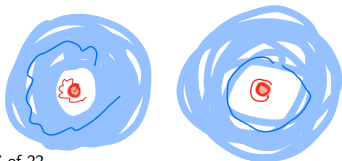
$$\nu_i^{AA'} = K_i \mu^{AA'} + F_i \delta \langle r^2 \rangle^{AA'} + \dots$$

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$$\propto \mu^{AA'} = \frac{1}{M^A} - \frac{1}{M^{A'}}$$

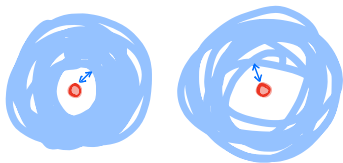


Field Shift

Nuclear charge distr. in A vs. A'

⇒ Difference in contact interactions between e^- & nuclei in A vs. A'

$$\propto \delta \langle r^2 \rangle^{AA'} = \langle r^2 \rangle^A - \langle r^2 \rangle^{A'}$$



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The King-Plot: Trade Data for Nuclear Physics

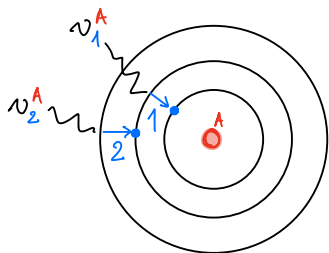
[W. King, J. Opt. Soc. Am. 53, 638 (1963)]

Issue: Large uncertainty on **charge radius**
variance $\delta\langle r^2 \rangle^{AA'}$

$$\nu_1^{AA'} = K_1 \mu^{AA'} + F_1 \delta\langle r^2 \rangle^{AA'}$$

The King-Plot: Trade Data for Nuclear Physics

[W. King, J. Opt. Soc. Am. 53, 638 (1963)]



Issue: Large uncertainty on **charge radius variance** $\delta\langle r^2\rangle^{AA'}$

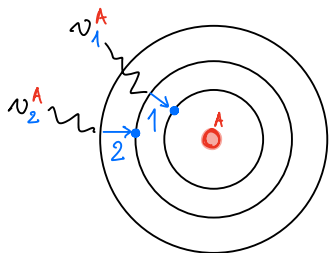
\Rightarrow Measure isotope shifts for 2 transitions

$$\nu_1^{AA'} = K_1 \mu^{AA'} + F_1 \delta\langle r^2\rangle^{AA'}$$

$$\nu_2^{AA'} = K_2 \mu^{AA'} + F_2 \delta\langle r^2\rangle^{AA'}$$

The King-Plot: Trade Data for Nuclear Physics

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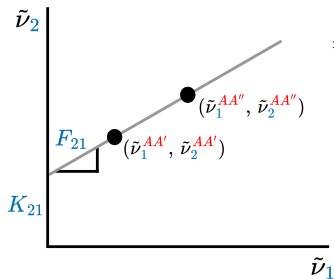
$$\nu_2^{AA'} = K_2 \mu^{AA'} + F_2 \delta\langle r^2\rangle^{AA'}$$

\Rightarrow Eliminate **charge radius variance** $\delta\langle r^2\rangle^{AA'}$

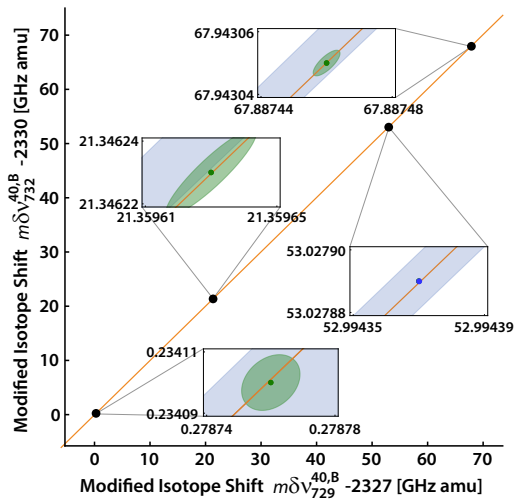
$$\tilde{\nu}_2^{AA'} = K_{21} + F_{21} \tilde{\nu}_1^{AA'}$$

$$\tilde{\nu}_i^{AA'} \equiv \nu_i^{AA'} / \mu^{AA'} \quad \Rightarrow \text{data}$$

$$F_{21} \equiv F_2 / F_1 \quad K_{21} \equiv K_2 - F_{21} K_1 \quad \Rightarrow \text{fit}$$



Example of a Linear King Plot: Ca^+ [arXiv:2311.17337]



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Isotope Shifts

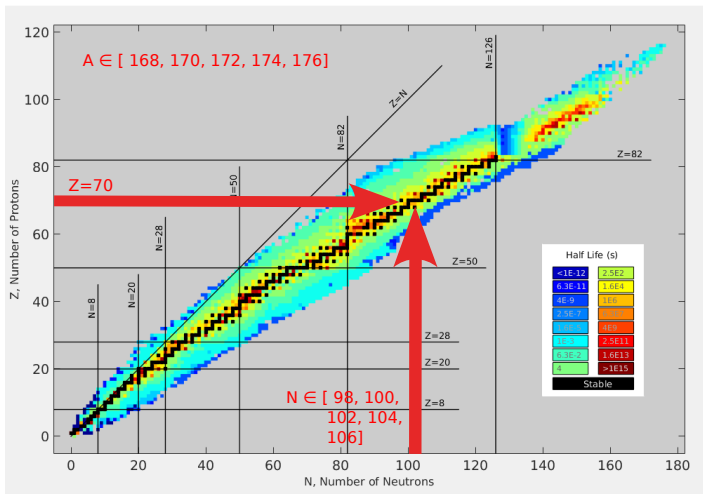
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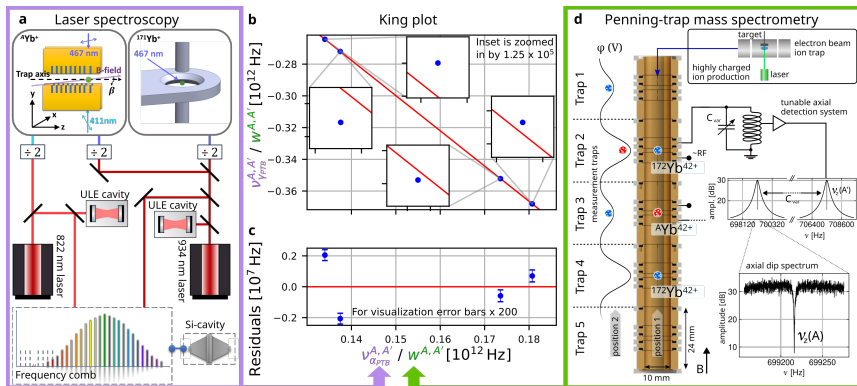
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Ytterbium and its Stable Isotopes



PTB + MPIK = New Yb King Plot [arXiv:2403.07792]

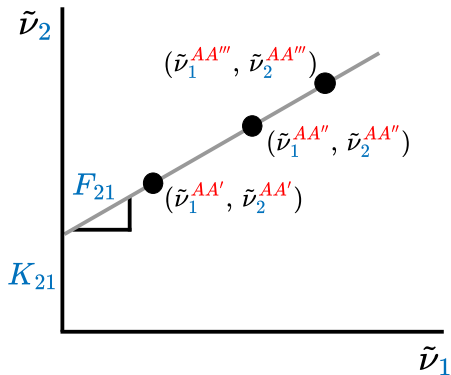


$$\text{Frequencies: } \frac{\delta\nu}{\nu} \sim \mathcal{O}(10^{-9}) \quad \text{Mass ratios } \eta_A = \frac{m_A}{m_{172}} : \frac{\delta\eta}{\eta} \sim \mathcal{O}(10^{-12})$$

⇒ Hz-precision for isotope shifts

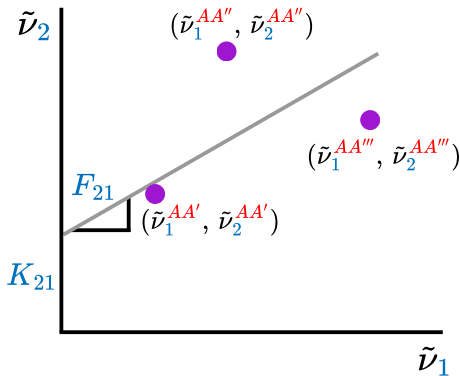
Observed King plot nonlinearity: $\sim 20.17(2)$ kHz

The King-Plot: Fit to Isotope Shift Data



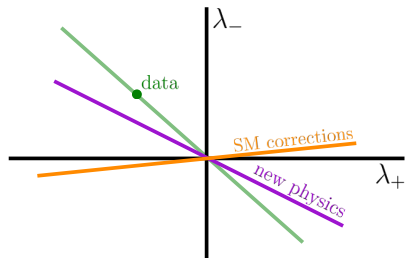
$$\begin{aligned}\tilde{\nu}_2^{AA'} &= K_{21} + F_{21}\tilde{\nu}_1^{AA'} \\ \tilde{\nu}_2^{AA''} &= K_{21} + F_{21}\tilde{\nu}_1^{AA''} \\ \tilde{\nu}_2^{AA'''} &= K_{21} + F_{21}\tilde{\nu}_1^{AA'''}\end{aligned}$$

The King-Plot: Fit to Isotope Shift Data



$$\begin{aligned}\tilde{\nu}_2^{AA'} &= K_{21} + F_{21}\tilde{\nu}_1^{AA'} + ? \\ \tilde{\nu}_2^{AA''} &= K_{21} + F_{21}\tilde{\nu}_1^{AA''} + ? \\ \tilde{\nu}_2^{AA'''} &= K_{21} + F_{21}\tilde{\nu}_1^{AA'''} + ?\end{aligned}$$

The Nonlinearity Decomposition Plot



- Plane of King linearity ($\mathbf{1} = (1, 1, 1, 1)$)

$$\tilde{\nu}_j \approx F_{j1} \tilde{\nu}_1 + K_{j1} \mathbf{1}, \quad j > 1.$$

- Project isotope-shift data onto $\tilde{\nu}_1, \mathbf{1}, \Lambda_+, \Lambda_-$ with $\Lambda_{\pm} \perp (\tilde{\nu}_1, \mathbf{1})$:

$$\tilde{\nu}_j = (\tilde{\nu}_1, \mathbf{1}, \Lambda_+, \Lambda_-) (F_{j1}, K_{j1}, \lambda_+, \lambda_-)^T$$

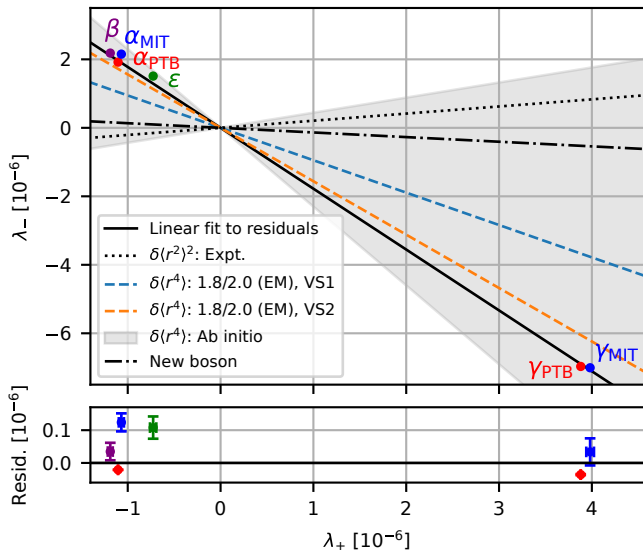
In presence of just one nonlinearity,

$$\tilde{\nu}_j \approx F_{j1} \tilde{\nu}_1 + K_{j1} \mathbf{1} + G_{j1}^{(4)} \delta \langle \tilde{r}^4 \rangle, \quad j > 1.$$

$$\text{slope: } \frac{\lambda_-}{\lambda_+} \equiv \frac{G_{j1}^{(4)} \delta \langle \tilde{r}^4 \rangle_-}{G_{j1}^{(4)} \delta \langle \tilde{r}^4 \rangle_+} = \frac{\delta \langle \tilde{r}^4 \rangle_-}{\delta \langle \tilde{r}^4 \rangle_+} \Rightarrow \text{transition-universal}$$

[arXiv:2004.11383, arXiv:2201.03578]

The Nonlinearity Decomposition Plot [arXiv:2403.07792]



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Extracting Nuclear Physics from Isotope-Shifts

- **Assuming $\delta\langle r^4 \rangle$ dominates**, what does the isotope-shift data tell us about the evolution of $\delta\langle r^4 \rangle$ along the isotope chain?

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⇒ “Put the King plot on it’s head.”: Experimental data (under control)

King Plot (in transition space): Eliminate charge radius variance $\delta\langle r^2 \rangle^{AA'}$

$$\nu_i^{AA'} = F_i \delta\langle r^2 \rangle^{AA'} + K_i \mu^{AA'} + G_i^{(4)} \delta\langle r^4 \rangle^{AA'}$$

$$\nu_1^{AA'} = F_1 \delta\langle r^2 \rangle^{AA'} + K_1 \mu^{AA'} + G_1^{(4)} \delta\langle r^4 \rangle^{AA'}$$

$$\Rightarrow \nu_i^{AA'} / \mu^{AA'} = F_{i1} \nu_1^{AA'} / \mu^{AA'} + K_{i1} + G_{i1}^{(4)} \delta\langle r^4 \rangle^{AA'} / \mu^{AA'}$$

⇒ Fit $F_{i1}, K_{i1}, G_{i1}^{(4)} = G_i^{(4)} - F_{i1} G_1^{(4)}$ (double-index electronic coefficient): AMBiT

Extracting Nuclear Physics from Isotope-Shifts

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$$\begin{array}{r} \nu_i^{AA'} = F_i \delta\langle r^2 \rangle^{AA'} + K_i \mu^{AA'} + G_i^{(4)} \delta\langle r^4 \rangle^{AA'} \\ \nu_1^{AA'} = F_1 \delta\langle r^2 \rangle^{AA'} + K_1 \mu^{AA'} + G_1^{(4)} \delta\langle r^4 \rangle^{AA'} \\ \hline \Rightarrow \nu_i^{AA'} / \mu^{AA'} = F_{i1} \nu_1^{AA'} / \mu^{AA'} + K_{i1} + G_{i1}^{(4)} \delta\langle r^4 \rangle^{AA'} / \mu^{AA'} \end{array}$$

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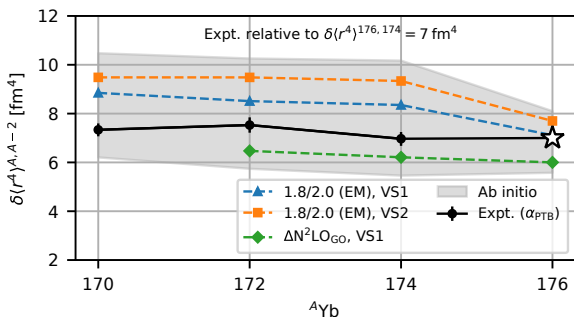
$\delta\langle r^4 \rangle$ -Extraction (in isotope pair space): Eliminate mass shift coefficient K_i

$$\begin{array}{r} \nu_i^{AA'} = F_i \delta\langle r^2 \rangle^{AA'} + K_i \mu^{AA'} + G_i^{(4)} \delta\langle r^4 \rangle^{AA'} \\ \nu_i^{RR'} = F_i \delta\langle r^2 \rangle^{RR'} + K_i \mu^{RR'} + G_i^{(4)} \delta\langle r^4 \rangle^{RR'} \\ \hline \nu_i^{AA'} - \frac{\mu^{AA'}}{\mu^{RR'}} \nu_1^{AA'} = F_i \left(\delta\langle r^2 \rangle^{AA'} - \frac{\mu^{AA'}}{\mu^{RR'}} \delta\langle r^2 \rangle^{AA'} \right) + G_i^{(4)} \underbrace{\left(\delta\langle r^4 \rangle^{AA'} - \frac{\mu^{AA'}}{\mu^{RR'}} \delta\langle r^4 \rangle^{RR'} \right)}_{Q^{AA'RR'}} \end{array}$$

⇒ Fit $F_i, \delta\langle r^2 \rangle^{AA'}: [\text{Angeli, Marinova}], G_i^{(4)}: \text{AMBiT}$

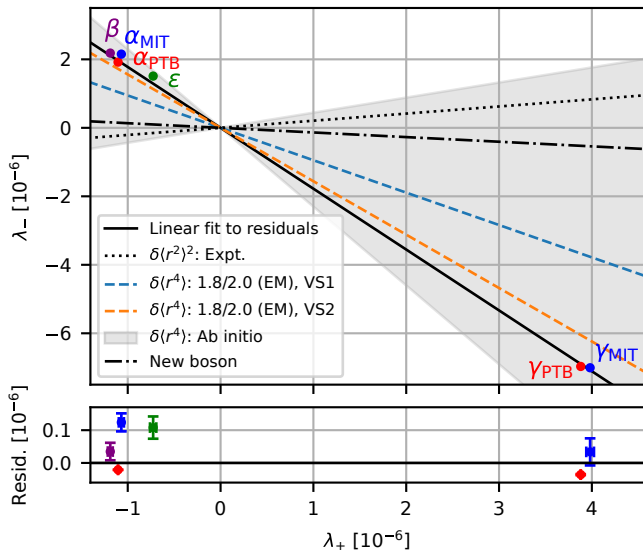
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blue, orange, green: Calculations by group of Prof. Achim Schwenk
black: new spectroscopic method, fixed at \star

The Nonlinearity Decomposition Plot [arXiv:2403.07792]



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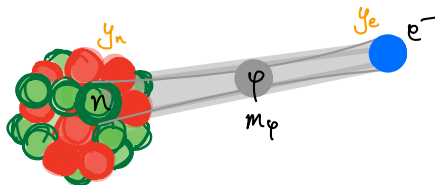
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Dark Portals

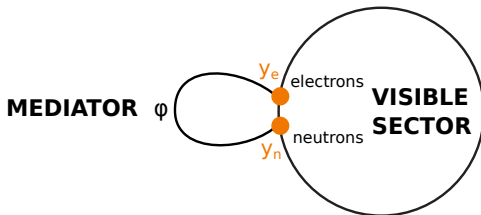


Question that can be addressed by King plot analyses:

Given a mediator mass m_φ in the **eV-GeV** range, ...

...how large can the coupling $y_e y_n$ be?

Dark Portals and Isotope Shift Measurements

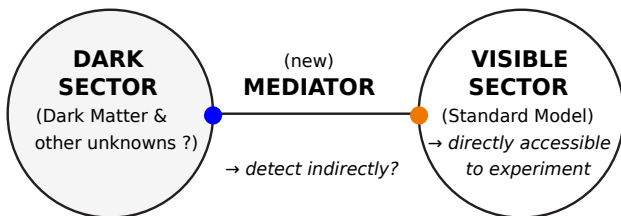


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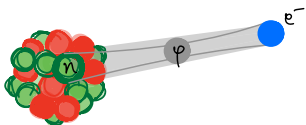


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King-Plot Bounds on New Bosons [arXiv:1704.05068,2005.06144]

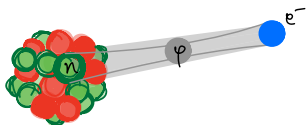


New effective Yukawa-potential

$$V_{\phi}(r) = -\alpha_{\text{NP}}(A - Z) \frac{e^{-m_{\phi}r}}{r}$$

with $\alpha_{\text{NP}} = (-1)^s \frac{y_e y_n}{4\pi}$, $s = 0, 1, 2$ (spin)

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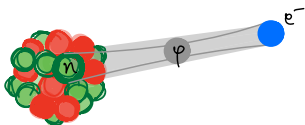
Induces new term in the King-relation:

$$\tilde{\nu}_2^{AA'} = K_{21} \tilde{\mu}^{AA'} + F_{21} \tilde{\nu}_1^{AA'} + G_{21}^{(4)} \delta \langle r^4 \rangle^{AA'} + \alpha_{\text{NP}} X_{21} \tilde{\gamma}^{AA'}$$

$X_{21} = X_2 - F_{21} X_1$: NP electronic coefficient

$\tilde{\gamma}^{AA'} \equiv (A - A') / \mu^{AA'}$: NP nucl. coeff.

King-Plot Bounds on New Bosons [arXiv:1704.05068,2005.06144]



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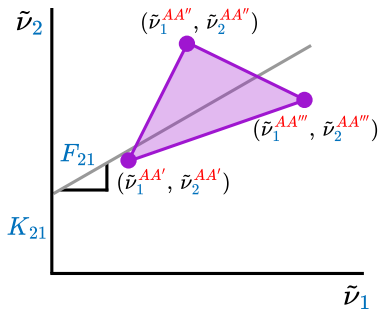
⇒ Extract α_{NP} from fraction of volumes spanned by frequency vectors:

$$\alpha_{\text{NP}} = \frac{\text{Vol.}}{\text{Vol.}|_{th, \alpha_{\text{NP}}=1}} = \frac{2 \det \left(\vec{\nu}_1, \vec{\nu}_2, \vec{\nu}_3, \vec{\mu} \right)}{\varepsilon_{ijk} \det \left(X_i \vec{\gamma}, \vec{\nu}_j, \vec{\nu}_k, \vec{\mu} \right)}$$

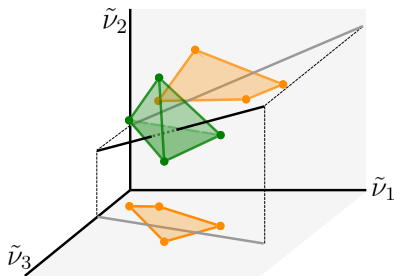
$\{\vec{\nu}_j\}$: data vect. in isotope-pair space, $\vec{\mu} \equiv (1, 1, 1, 1)$, $X_i, \vec{\gamma}$: theory input

King-Plot Method in Presence of Nuclear Effects: The Generalised King Plot

[arXiv:2005.06144]



⇒ Test King linearity

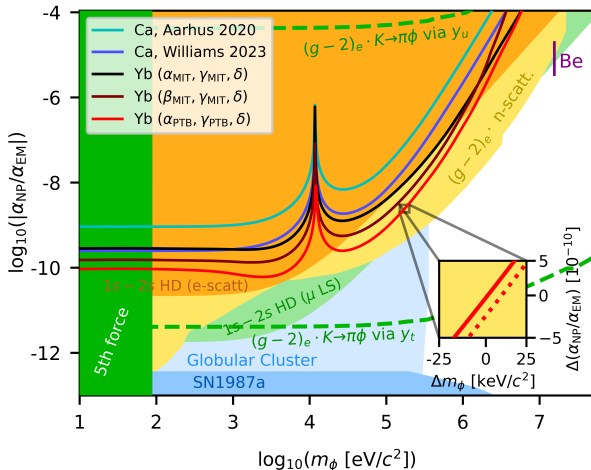


⇒ Account for one King nonlinearity

⇒ Put bound on 2^{nd}

⇒ **King-plot method also works in presence of nuclear effects.**

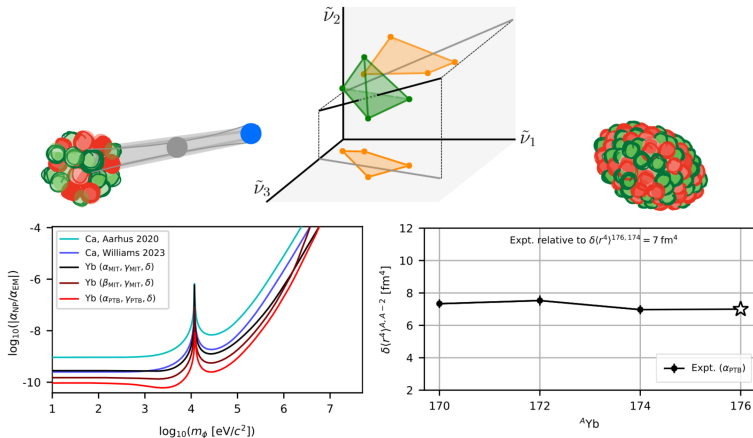
New Spectroscopy Bounds on New Physics



- $m_\phi \rightarrow 0$: $>$ size atom
- $m_\phi \rightarrow \infty$: not sensitive to contact interactions
- "Peaks" due to cancellations among electronic coefficients

Conclusions

Atomic clocks are sensitive probes for



New mediators between n & e^-

Nuclear structure

Check out our paper:

Yb King plot: [arXiv:2403.07792](https://arxiv.org/abs/2403.07792)

Stay tuned for:

- kifit: Global King-plot analysis

Thank you for your attention.

Backup slides

α_{NP} from Determinants

(No-Mass King-Plot:)

$$\vec{\nu}_1 = K_1 \vec{\mu} + F_1 \overrightarrow{\delta\langle r^2 \rangle} + \alpha_{\text{NP}} X_1 \vec{\gamma}$$

$$\vec{\nu}_2 = K_2 \vec{\mu} + F_2 \overrightarrow{\delta\langle r^2 \rangle} + \alpha_{\text{NP}} X_2 \vec{\gamma}$$

$$\vec{\nu}_3 = K_3 \vec{\mu} + F_3 \overrightarrow{\delta\langle r^2 \rangle} + \alpha_{\text{NP}} X_3 \vec{\gamma}$$

$$\Rightarrow \det(\vec{\nu}_1, \vec{\nu}_2, \vec{\nu}_3) = \alpha_{\text{NP}} \det(\vec{K}, \vec{F}, \vec{X}) \det(\vec{\mu}, \overrightarrow{\delta\langle r^2 \rangle}, \vec{\gamma})$$

$$\begin{aligned} \Rightarrow \alpha_{\text{NP}} &= \frac{\text{Vol}}{\text{Vol}|_{th, \alpha_{\text{NP}}=1}} = \frac{\det(\vec{\nu}_1, \vec{\nu}_2, \vec{\nu}_3)}{\det(\vec{K}, \vec{F}, \vec{X}) \det(\vec{\mu}, \overrightarrow{\delta\langle r^2 \rangle}, \vec{\gamma})} \\ &= \frac{\det(\vec{\nu}_1, \vec{\nu}_2, \vec{\nu}_3)}{\frac{1}{2} \varepsilon_{ijk} \det(X_i \vec{\gamma}, \vec{\nu}_j, \vec{\nu}_k)} \end{aligned}$$

Choose your King-Plot

Extraction of α_{NP} using the “determinant method” requires

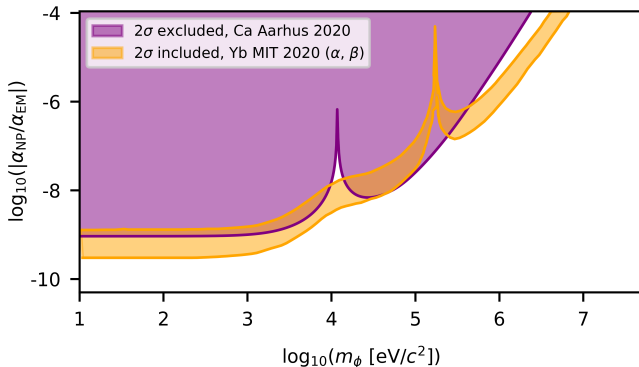
Type of King-Plot	Isotope-Pairs	Transitions	
Generalised King-Plot:	n	$n - 1$	[PRR 2, 043444 (2020)]
No-Mass King-Plot:	n	n	[PRR 2, 043444 (2020)]

$n \geq 3$ (else cannot search for nonlinearities)

$$\alpha_{\text{NP}} = \frac{V}{V|_{\text{th}, \alpha_{\text{NP}}=1}} = \frac{(n-2)! \det(\vec{v}_1, \dots, \vec{v}_{n-1}, \vec{\mu})}{\varepsilon_{i_1, \dots, i_{n-1}} \det(X_{i_1} \vec{\gamma}, \vec{v}_{i_2}, \dots, \vec{v}_{i_{n-1}}, \vec{\mu}_{i_n})}$$

$$\alpha_{\text{NP}} = \frac{v}{v|_{\text{th}, \alpha_{\text{NP}}=1}} = \frac{(n-1)! \det(\vec{v}_1, \vec{v}_2, \dots, \vec{v}_n)}{\varepsilon_{i_1, i_2, \dots, i_n} \det(X_{i_1} \vec{\gamma}, \vec{v}_{i_2}, \dots, \vec{v}_{i_n})}$$

Upper Bounds on $|\alpha_{NP}|$ vs. New Mediator Mass m_ϕ



Nonlinear King plot relation:

$$\tilde{\nu}_2^{AA'} = K_{21}\tilde{\mu}^{AA'} + F_{21}\tilde{\nu}_1^{AA'} + G_{21}^{(2)}\delta\langle r^2 \rangle^2 + G_{21}^{(4)}\delta\langle r^4 \rangle + \dots?$$

X Coefficients

Overlap of new physics potential and electronic wavefunction

$$X_i = \int d^3r \frac{e^{-m_\phi r}}{r} [|\psi_b(r)|^2 - |\psi_a(r)|^2]$$

$|\psi(r)|^2$: electron density in absence of new physics,
 a, b initial, final states

Requirement for searches for new light bosons:

- At least one of ψ_a or ψ_b should have good overlap with new potential.
- For tight bounds on α_{NP} , one X_i needs to be large.

Recipe for the Nonlinearity Decomposition Plot

[PRL 125, 123002 (2020), PRL 128, 163201 (2022)]

1. Arrange the isotope-shift data for all transitions $\tau \in \{\alpha, \beta, \gamma, \delta, \epsilon\}$ in n -vectors $\tilde{\nu}_\tau$, where n is the number of isotope pairs (here 4):

$$\tilde{\nu}_\tau = (\tilde{\nu}_\tau^{168,170}, \tilde{\nu}_\tau^{170,172}, \tilde{\nu}_\tau^{172,174}, \tilde{\nu}_\tau^{174,176})$$

2. Choose a reference transition, say δ .
3. Plane of King linearity is defined by the relations ($\mathbf{1} = (1, 1, 1, 1)$)

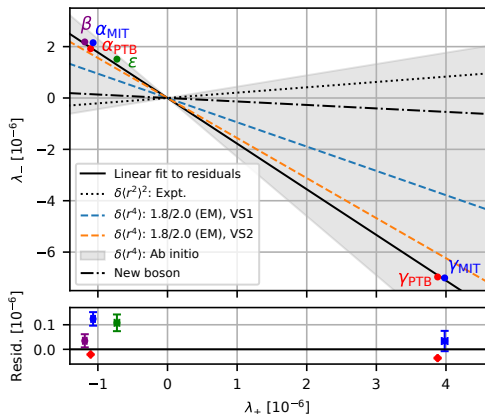
$$\tilde{\nu}_\tau \approx F_{\tau\delta} \tilde{\nu}_\delta + K_{\tau\delta} \mathbf{1}.$$

4. Define two ($n = 4$)–vectors Λ_\pm that are orthogonal to $\tilde{\nu}_\delta$, $\mathbf{1}$.
5. Project all isotope-shift data onto the four vectors $\tilde{\nu}_\delta$, $\mathbf{1}$, Λ_+ , Λ_- :

$$\tilde{\nu}_\tau = (\tilde{\nu}_\delta \quad \mathbf{1} \quad \Lambda_+ \quad \Lambda_-) \begin{pmatrix} F_{\tau\delta} & K_{\tau\delta} & \lambda_+^{(\tau)} & \lambda_-^{(\tau)} \end{pmatrix}^T$$

6. Plot all points $(\lambda_+^{(\tau)}, \lambda_-^{(\tau)})$ in the same plane.

The Nonlinearity Decomposition Plot



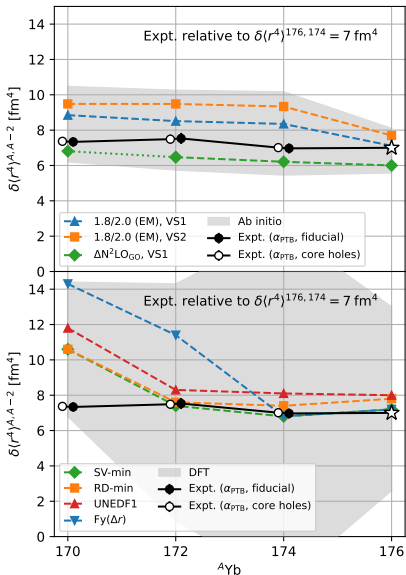
Notation	Transition	Refs.
$\alpha_{MIT,PTB}$	$2S_{1/2} \rightarrow 2D_{5/2}$ E2 in Yb ⁺	MIT, t.w.
β	$2S_{1/2} \rightarrow 2D_{3/2}$ E2 in Yb ⁺	MIT
$\gamma_{MIT,PTB}$	$2S_{1/2} \rightarrow 2F_{7/2}$ E3 in Yb ⁺	MIT, t.w.
δ	$1S_0 \rightarrow 3P_0$ in Yb	Kyoto
ϵ	$1S_0 \rightarrow 1D_2$ in Yb	Mainz

- $\delta\langle r^2 \rangle^2$ estimated using Angeli & Marinova Tables of experimental nuclear ground state charge radii
- $\delta\langle r^4 \rangle$: Calculations by group of Prof. Achim Schwenk, TU Darmstadt

In presence of just one nonlinearity, e.g. $G^{(4)}\delta\langle r^4 \rangle$,

$$\text{slope: } \frac{\lambda_-^{(\tau)}}{\lambda_+^{(\tau)}} = \frac{G_\tau^{(4)}\delta\langle r^4 \rangle_-}{G_\tau^{(4)}\delta\langle r^4 \rangle_+} = \frac{\delta\langle r^4 \rangle_-}{\delta\langle r^4 \rangle_+} \equiv \frac{\lambda_-}{\lambda_+} \Rightarrow \text{transition-universal}$$

$\delta\langle r^4 \rangle$ Calculations: Ab initio vs. DFT



- Experimental $\delta\langle r^4 \rangle^{AA'}$ values relative to $\delta\langle r^4 \rangle^{176,174} = 7 \text{ fm}^4$ extracted from isotope shifts from the α transition using atomic theory (fiducial, core holes)
- Above: ab initio calculations (t.w.)
- Below: density functional theory calculations (PRL.128.163201)
- Gray bands: estimated theory uncertainties