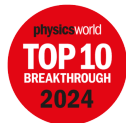


Modification of nPDFs by correlated nucleon pairs

Michael Klasen

ITP, University of Münster



CERN TH Heavy Ion Coffee, January 27, 2025

A. Denniston, T. Jezo, A. Kusina et al., Phys. Rev. Lett. 133 (2024) 15

Introduction

Short-range nucleon dynamics is fundamental:

- Deficiencies of mean-field model
- Change of nucleon structure in nuclei (EMC effect)
- Related to quark and gluon d.o.f. in nuclei (?)
- Properties of matter under extreme conditions

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- Complex nuclear many-body problem
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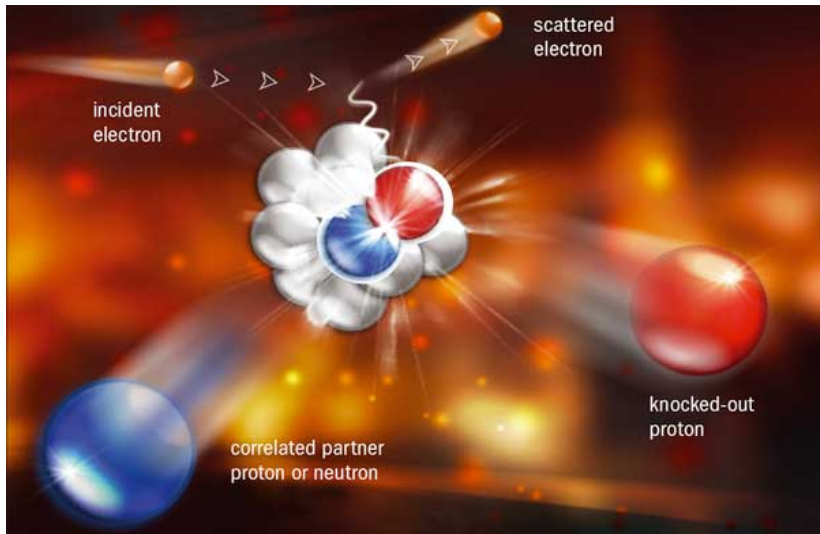
Short-range nucleon dynamics is challenging:

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Goal of this work:

- New nPDF ansatz: Single N in mean field + SRC NN pairs
- nCTEQ analysis of DIS, DY, W/Z-boson data at high energy
- Extract SRC fractions and properties of SRC pairs

Nuclei and short-range correlations



The nuclear many-body problem (1)

C. Ciofi degli Atti, Phys. Rep. 590 (2015) 1

Non-relativistic Schrödinger equation (even though $\frac{v_N}{c} \sim 0.1$):

$$\left[\sum_i \frac{p_i^2}{2m_N} + \sum_{i<j} v_2(x_i, x_j) + \sum_{i<j<k} v_3(x_i, x_j, x_k) \right] \psi_A^f(\{x_A\}) = E_A^f \psi_A^f(\{x_A\})$$

- V_2 dominates, $V_n \simeq \left(\frac{v_N}{c}\right)^{n-2} V_2$ [H. Primakoff, T. Holstein, PR 55 (1939) 1218]
- V_3 contributes $\leq 20\%$ to BE [S.C. Pieper, R.B. Wiringa, ARNPS 51 (2001) 53]

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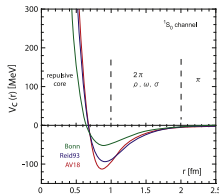
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Nuclear NN potentials:

[R.J. Furnstahl, K. Ebeler, Rep. Progr. Phys. 76 (2013) 126301]

- Depend on L, S, T ($L + S + T$ odd, Pauli)
- Attractive (1S_0 , spin-aligned) for $r \geq 0.7$ fm
- Tensor character for $S = 1$
- Repulsive at small $r \rightarrow$ pert. theory fails
- Off-shell effects, important at small r

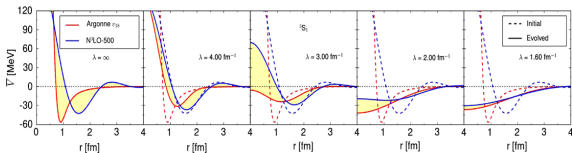


The nuclear many-body problem (2)

C. Ciofi degli Atti, Phys. Rep. 590 (2015) 1

Chiral perturbation theory: [H.W. Hammer, A. Nogga, A. Schwenk, RMP 85 (2013) 197]

- EFT based on approx., spont. broken chiral symmetry of QCD
- Systematic expansion in local operators and powers of $\Lambda \simeq m_\rho$
- Short range contact int., long range pion exchange ($p_N \simeq m_\pi$)
- Also repulsive at small r , solvable up to $A \leq 12$ with RGEs



- Off-shell NN interactions \leftrightarrow On-shell 3N forces

	pionless	chiral	chiral+ Δ
LO			
NLO			
N ² LO			

The nuclear many-body problem (3)

C. Ciofi degli Atti, Phys. Rep. 590 (2015) 1

Nuclear shell model:

- Independent Particle Model: [Jensen et al. PR 75 (1949) 1766; Mayer idib. 1969]

$$\left[-\frac{\hbar^2}{2m_N} \nabla_i^2 + U(x_i) \right] \phi_\alpha(x_i) = \epsilon_\alpha \phi_\alpha(x_i)$$

with isotropic mean field $U(r_i) = \frac{1}{2} \hbar \omega r_i^2 + DL_i^2 + \text{CLS}$.

Generated by $V_2 \rightarrow$ Reproduces magic numbers in Z and N

- Motivates "bound" nucleon PDFs
- V_3 modifies monopole, explains e.g. $N = 28$ in ^{20}Ca

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Modern approaches:

- No-Core Shell Model: [B.R. Barrett, P. Navratil J.P. Vary, PPNP 69 (2013) 131]
Variational many-body ansatz ($A \leq 12$), also used with ChPT
- Quantum Monte Carlo: [R. Cruz-Torres et al., Nature Phys. 17 (2021) 306]

Scale separation of short and long distances \rightarrow

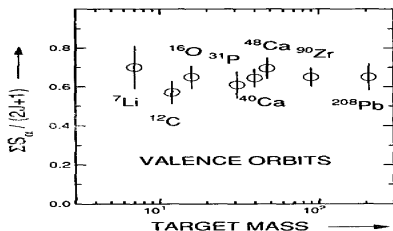
A-dependent SRC fractions, A-independent contact terms

Short-range correlations

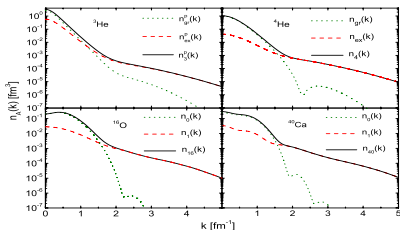
O. Hen et al., Rev. Mod. Phys. 89 (2017) 045002 [1611.09768]

Experimental and theoretical evidence:

- QE data show only 65% of single- N strength predicted by IPM



L. Lapikas, Nuclear Physics A 553 (1993) 297c



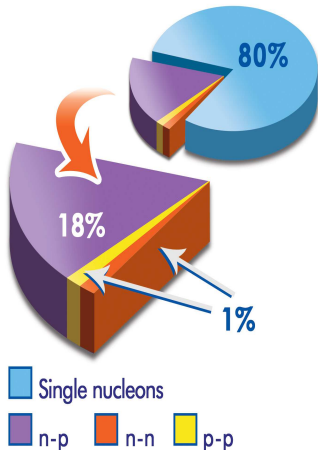
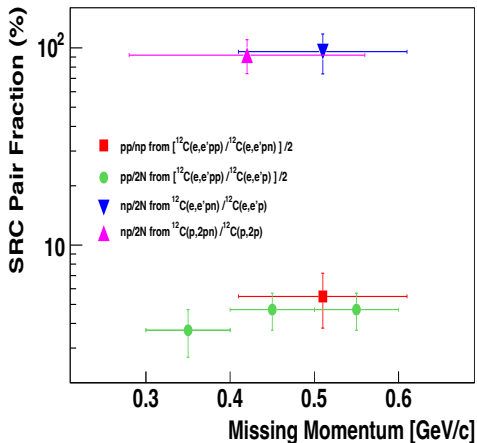
M. Alvioli et al., PRC 87 (2013) 034603 [1211.0134]

- Remainder in SRC pairs (small p^* , but $p_{\text{rel.}} > p_F \sim 250$ MeV)
- Dominated by nodeless rel. S -wave with $S = 1$, $T = 0$ ($\sim d$)
- Typical distance ~ 1 fm $<$ average 1.7 fm \rightarrow Higher density

Exclusive quasi-elastic scattering

R. Subedi et al. (Hall A), Science 320 (2008) 1476 [0908.1514]

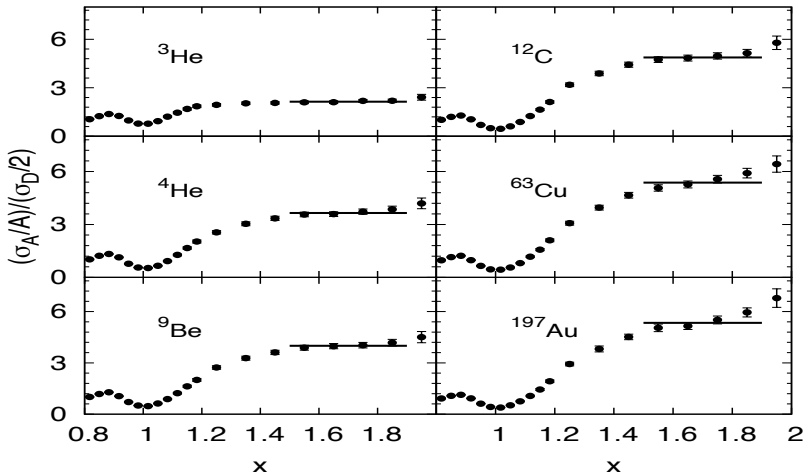
$^{12}\text{C}(e, e'p, pN)$ ($E_e = 4.627$ GeV, $\theta_e = 19.5^\circ$):



Inclusive quasi-elastic scattering

N. Fomin et al. (Hall C), Phys. Rev. Lett. 108 (2012) 092502 [1107.3583]

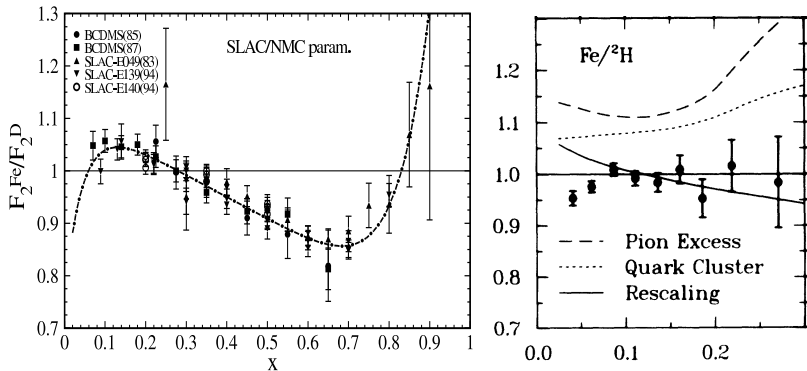
Ratio $a_2 = \text{const.}$ for $x > 1.5$ ($E_e = 5.766$ GeV, $\theta_e = 18^\circ$):



Inclusive deep-inelastic scattering

I. Schienbein et al. [nCTEQ], PRD 80 (2009) 094004 [0907.2357]; D.M. Alde et al., PRL 64 (1990) 2479

Nuclear modification of $F_2^A(x, Q^2)$ ($E_e = 4.5 \dots 280$ GeV):

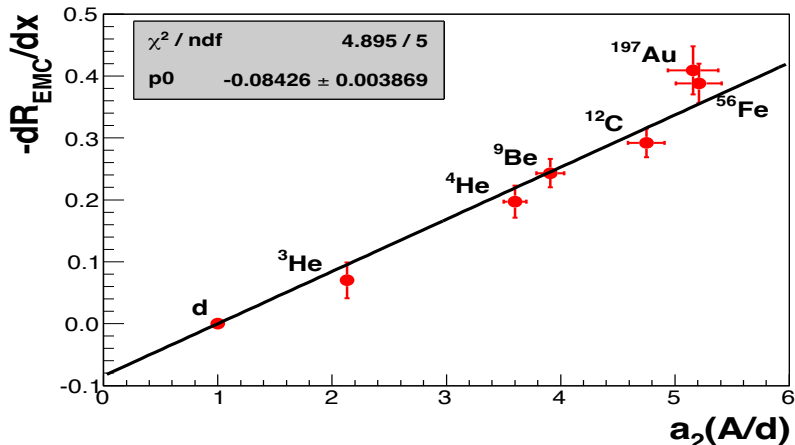


Explanation with N 's only violates baryon/momentum sum rules \rightarrow
 Need also π 's, but no \bar{q} enhancement seen in nuclear DY.
 Partonic? Fewer high-momentum quarks \rightarrow larger size (e.g. NN^*).

SRCs and the EMC effect

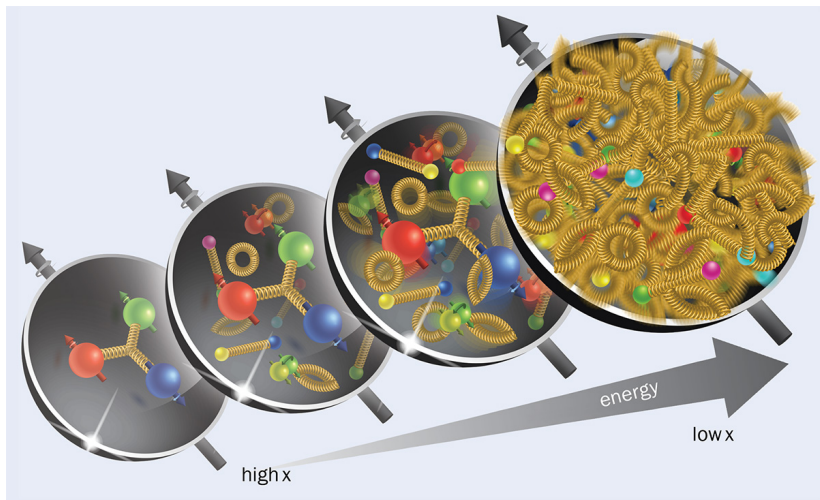
O. Hen, E. Piasetzky, L.B. Weinstein et al., PRL 106 (2011) 052301 and PRC 85 (2012) 047301 [1202.3452]

EMC slope in $0.35 \leq x \leq 0.7$ vs. ratio $a_2 = \sigma_A/\sigma_d$ for $x > 1.5$:



Correlation suggests that EMC effect is of short-distance nature.

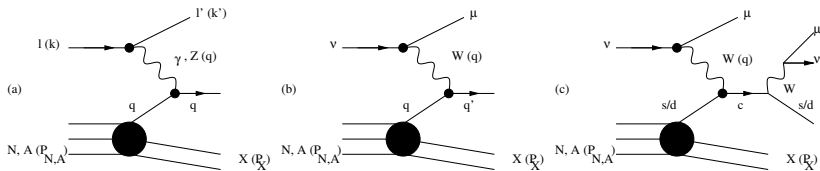
Deep-inelastic scattering and nuclear PDFs



Theoretical foundations

MK, H. Paukkunen, *Ann. Rev. Nucl. Part. Sci.* 74 (2024) 49 [2311.00450]

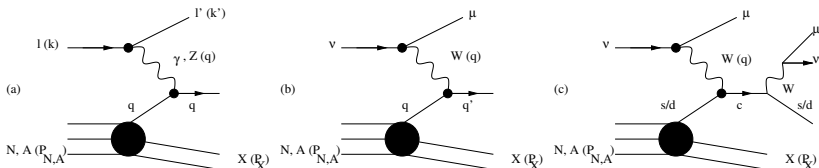
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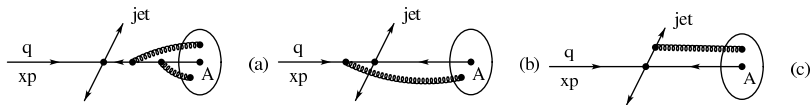
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Deep-inelastic scattering (NC, CC, dimuon production):



Hadronic collisions: Leading twist, higher-twist

[J.w. Qiu, 0305161]



- Transv. size, jet mass, rescattering: $\mathcal{O}\left(r_T^2 \sim \frac{1}{p_T^2}, \frac{m_J^2}{p_T^2}, \frac{\alpha_s(Q^2)\Lambda^2}{Q^2}\right)$
- Enhanced in nuclear collisions by $A^{1/3}$ due to many soft partons

(Perturbative) Quantum Chromodynamics

MK, H. Paukkunen, *Ann. Rev. Nucl. Part. Sci.* 74 (2024) 49 [2311.00450]

Nuclear structure function(s) in deep-inelastic scattering (DIS):

$$F_2^A(x, Q^2) = \sum_i f_i^{(A,Z)}(x, Q^2) \otimes C_{2,i}(x, Q^2)$$

QCD factorization theorem, Wilson coefficients $C_{2,i}$ at (N)NLO

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DGLAP evolution equations:

$$\frac{\partial f_i(x, Q^2)}{\partial \log Q^2} = \int_x^1 \frac{dz}{z} P_{ij} \left(\frac{x}{z}, \alpha_s(Q^2) \right) f_j(z, Q^2)$$

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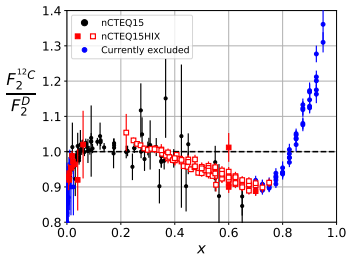
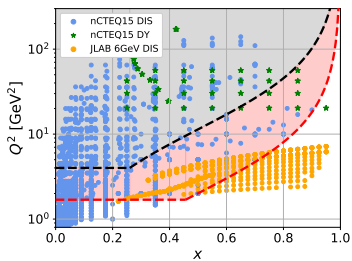
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Number and momentum sum rules, but also isospin symmetry:

$$\int_0^1 dx [f_{\{u,d\}}^{p/A} - f_{\{\bar{u},\bar{d}\}}^{p/A}(x)] = \{2, 1\} \quad ; \quad f_{d,u}^{n/A}(x, Q^2) = f_{u,d}^{p/A}(x, Q^2)$$

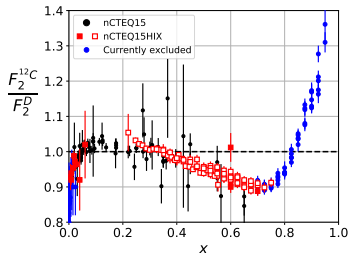
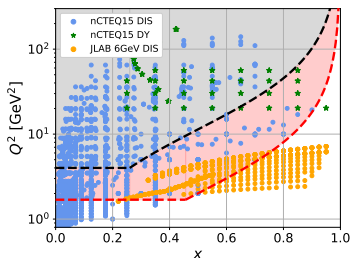
High- x JLab data: Deuteron, TMCs and HT

A. Accardi et al., Phys. Rev. D 93 (2016) 114017 [1602.03154]; E.P. Segarra et al., Phys. Rev. D 103 (2021) 114015 [2012.11566]



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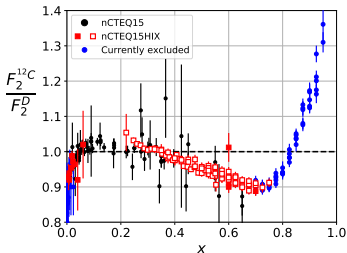
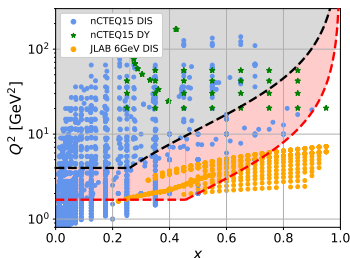


Deuteron:

- Loosely bound \rightarrow often isoscalar (pn) assumed, fitted with p
- Fermi motion, nucl. binding, off-shell effects (few %) [CJ15,CJ22]

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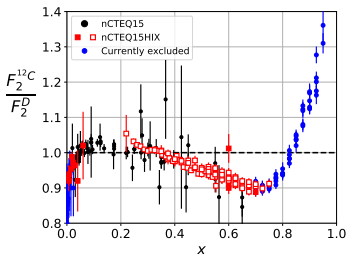
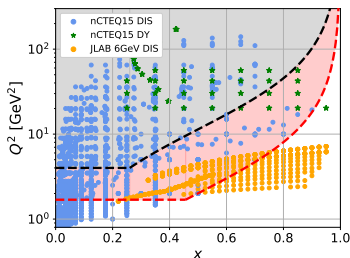
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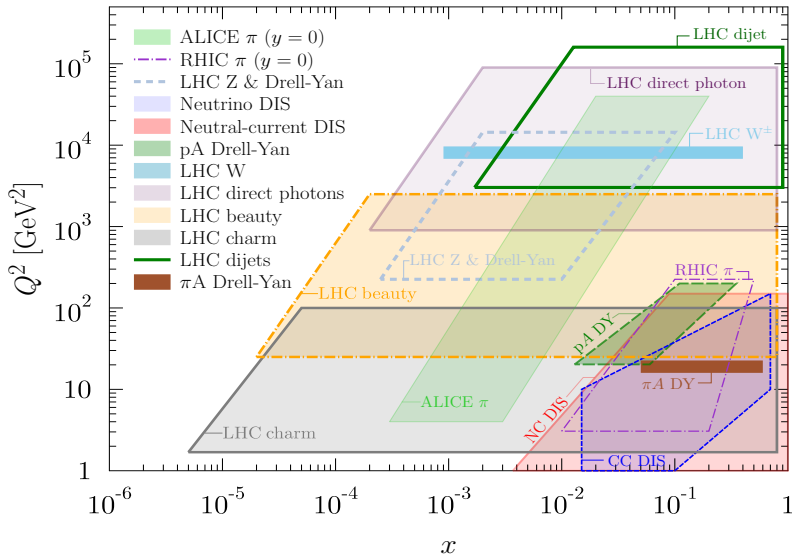
Higher twist (HT) corrections:

[CJ15,CJ22]

- $F_2^A(x, Q) \rightarrow F_2^A(x, Q) \left[1 + \frac{A^{1/3} h_0 x^{h_1} (1 + h_2 x)}{Q^2} \right]$

Experimental kinematic coverage in x and Q^2

MK, H. Paukkunen, Ann. Rev. Nucl. Part. Sci. 74 (2024) 49 [2311.00450]



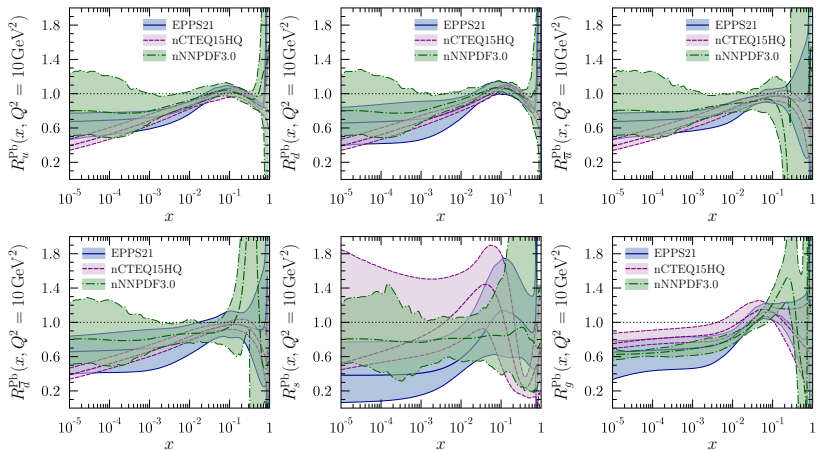
Global analyses of nuclear PDFs

MK, H. Paukkunen, Ann. Rev. Nucl. Part. Sci. 74 (2024) 49 [2311.00450]

ANALYSIS	nCTEQ15HQ	EPPS21	nNNPDF3.0	TUJU21	KSASG20
THEORETICAL INPUT:					
Perturbative order	NLO	NLO	NLO	NNLO	NNLO
Heavy-quark scheme	SACOT- χ	SACOT- χ	FONLL	FONLL	FONLL
Data points	1484	2077	2188	2410	4353
Independent flavors	5	6	6	4	3
Free parameters	19	24	256	16	18
Error analysis	Hessian	Hessian	Monte Carlo	Hessian	Hessian
Tolerance	$\Delta\chi^2 = 35$	$\Delta\chi^2 = 33$	N/A	$\Delta\chi^2 = 50$	$\Delta\chi^2 = 20$
Proton PDF	\sim CTEQ6.1	CT18A	\sim NNPDF4.0	\sim HERAPDF2.0	CT18
Deuteron corrections	$(\checkmark)^{a,b}$	\checkmark^c	\checkmark	\checkmark	\checkmark
FIXED-TARGET DATA:					
SLAC/EMC/NMC NC DIS	\checkmark	\checkmark	\checkmark	\checkmark	\checkmark
- Cut on Q^2	4 GeV ²	1.69 GeV ²	3.5 GeV ²	3.5 GeV ²	1.2 GeV ²
- Cut on W^2	12.25 GeV ²	3.24 GeV ²	12.5 GeV ²	12.0 GeV ²	
JLab NC DIS	$(\checkmark)^a$	\checkmark			\checkmark
CHORUS/CDHSW CC DIS	$(\checkmark/-)^b$	$\checkmark/-$	$\checkmark/-$	\checkmark/\checkmark	\checkmark/\checkmark
NuTeV/CCFR 2 μ CC DIS	$(\checkmark/\checkmark)^b$		$\checkmark/-$		
ρA DY	\checkmark	\checkmark	\checkmark		\checkmark
COLLIDER DATA:					
Z bosons	\checkmark	\checkmark	\checkmark	\checkmark	
W^\pm bosons	\checkmark	\checkmark	\checkmark	\checkmark	
Light hadrons	\checkmark	\checkmark^d			
Jets		\checkmark	\checkmark		
Prompt photons			\checkmark		
Prompt D ⁰	\checkmark	\checkmark	\checkmark^e		
Quarkonia (J/ψ , ψ' , Υ)	\checkmark				

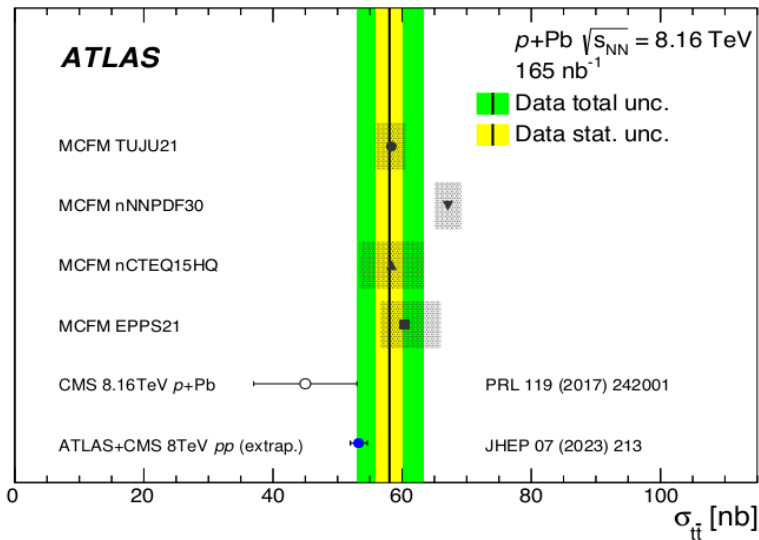
Nuclear PDFs after 10 years of LHC data

MK, H. Paukkunen, Ann. Rev. Nucl. Part. Sci. 74 (2024) 49 [2311.00450]

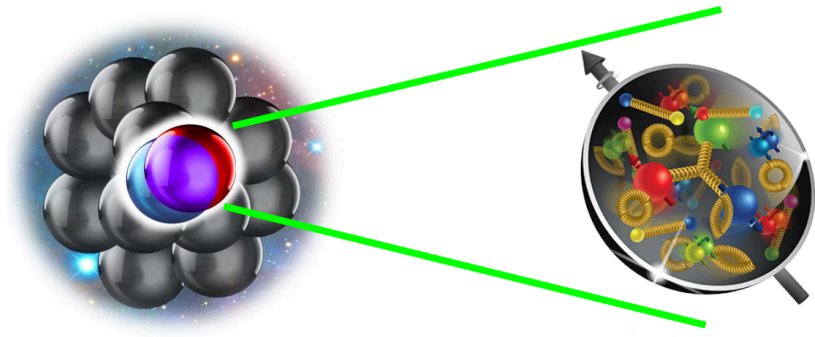


Top pair production in pPb with ATLAS (and CMS)

ATLAS Coll., JHEP 11 (2024) 101 [2405.05078]



Nuclear PDFs from short-range correlations



Experimental data

A. Denniston, T. Jezo, A. Kusina, MK et al., Phys. Rev. Lett. 133 (2024) 15 [2312.16293]

Inclusive FT (CERN, FNAL, SLAC) and collider (RHIC, LHC) exp.:

- Deep-inelastic scattering:

Kinematics: $Q^2 = -(p_e - p_{e'})^2$, $y = 1 - E_{e'}/E_e$, $x = Q^2/(sy)$

$$d\sigma(eA \rightarrow e'X) = \frac{4\pi\alpha^2}{Q^4} \left[F_2^A(x, Q^2) \left(\frac{y^2}{2} + 1 - y \right) - xy^2 F_L^A \right]$$

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- Vector-boson (Drell-Yan, W/Z) production:

$$d\sigma(pA \rightarrow VX) = \sum_{i,j} f_i^p \otimes f_j^A(x, Q^2) \otimes d\hat{\sigma}(ij \rightarrow VX)$$

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$$d\sigma(eA \rightarrow e'X) = \frac{4\pi\alpha^2}{Q^4} \left[F_2^A(x, Q^2) \left(\frac{y^2}{2} + 1 - y \right) - xy^2 F_L^A \right]$$

where $F_2^A(x, Q^2) = \sum_i f_i^A(x, Q^2) \otimes C_{2,i}(x, Q^2)$.

- Vector-boson (Drell-Yan, W/Z) production:

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- Other pA processes (jets, photons, light/hadrons) \rightarrow low x

Experimental data

A. Denniston, T. Jezo, A. Kusina, MK et al., Phys. Rev. Lett. 133 (2024) 15 [2312.16293]

Inclusive FT (CERN, FNAL, SLAC) and collider (RHIC, LHC) exp.:

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- nCTEQ15HIX/SRC, nCTEQ25: $Q > 1.3$ GeV, $W > 1.7$ GeV

Traditional theoretical approach

E.P. Segarra, T. Jezo, MK et al., Phys. Rev. D 103 (2021) 114015 [2012.11566]

Q^2 -dependence calculable to N³LO QCD [S. Moch et al., PLB 860 (2025) 139194]

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x -dependence at $Q_0 = 1.3 \text{ GeV}$ must be fitted to experiments.

Our traditional parameterisation (e.g. nCTEQ15HIX):

$$xf_i^{p/A}(x, Q_0) = c_0 x^{c_1} (1-x)^{c_2} e^{c_3 x} (1 + e^{c_4 x})^{c_5}$$

where $c_k(A) = p_k + a_k(1 - A^{-b_k})$, $k = \{1, \dots, 5\}$,

and $i = \{u_v, d_v, (\bar{u} + \bar{d}), (\bar{d}/\bar{u}), s, g\}$, $s = \bar{s} = \kappa(\bar{u} + \bar{d})/2$, $\kappa = 0.5$.

Reproduces free proton for $A \rightarrow 1$, open 19 free parameters.

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χ^2 test function for data set D (3 norm. par. for W/Z production):

$$\chi_D^2 = \sum_{i,j}^N \left(D_i - \frac{T_i}{N_{norm}} \right) (C^{-1})_{ij} \left(D_j - \frac{T_j}{N_{norm}} \right) + \left(\frac{1 - N_{norm}}{\sigma_{norm}} \right)^2$$

Covariance matrix: $C_{ij} = \sigma_i^2 \delta_{ij} + \sum_{\alpha}^S \bar{\sigma}_{i\alpha} \bar{\sigma}_{j\alpha}$

Hessian (or Markov Chain Monte Carlo) error analysis

SRC-motivated nuclear PDFs

A. Denniston, T. Jezo, A. Kusina, MK et al., Phys. Rev. Lett. 133 (2024) 15 [2312.16293]

Nuclear spectral function:

[C. Ciofi degli Atti, Phys. Rep. 590 (2015) 1]

$$S_A(k, E) = S_A^{\text{MF}}(\text{small } k, E) + S_A^{\text{SRC}}(\text{large } k, E)$$

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[R. Weiss et al., Phys. Lett. B 791 (2019) 242]

$$S_A^{\text{SRC}}(k, E) = \frac{Z}{A} \left[2C_{p/(pp)}^A \times S_A^{pp} + C_{p/(pn)}^A \times S_A^{pn} \right] \\ + \frac{N}{A} \left[2C_{n/(nn)}^A \times S_A^{nn} + C_{n/(pn)}^A \times S_A^{pn} \right]$$

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$$f_i^A(x, Q_0) = \frac{Z}{A} \left[(1 - C_p^A) \times f_i^p(x, Q_0) + C_p^A \times f_i^{p,\text{SRC}}(x, Q_0) \right] \\ + \frac{N}{A} \left[(1 - C_n^A) \times f_i^n(x, Q_0) + C_n^A \times f_i^{n,\text{SRC}}(x, Q_0) \right]$$

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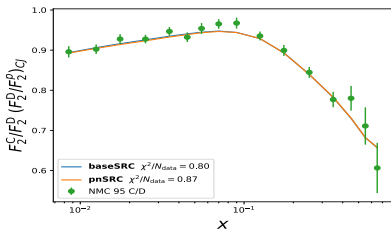
where we assume $C_p^A = 2C_{p/(pp)}^A + C_{p/(pn)}^A$ and similarly for $p \leftrightarrow n$.

Open 21 parameters (\sim HIX + s) + 30/19 for $C_{p,n}^A$ in base/pn.

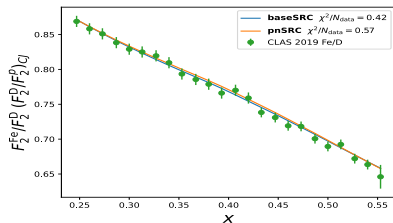
Fitted nuclei and selected comparisons

A. Denniston, T. Jezo, A. Kusina, MK et al., Phys. Rev. Lett. 133 (2024) 15 [2312.16293]

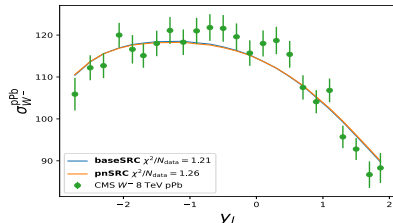
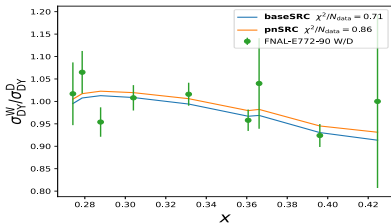
Nuclear A	2	3	4	6	9	12	14	27	40	56	64	84	108	119	131	184	197	208
# data	275	125	66	15	49	196	101	73	92	134	61	84	7	152	4	37	50	163



(a)



(b)



Quality of our fits

A. Denniston, T. Jezo, A. Kusina, MK et al., Phys. Rev. Lett. 133 (2024) 15 [2312.16293]

χ^2/N_{data}	DIS	DY	W/Z	JLab	χ_{tot}^2	$\frac{\chi_{\text{tot}}^2}{N_{\text{DOF}}}$
Reference	0.85	0.97	0.88	0.72	1408	0.85
SRC baseSRC	0.84	0.75	1.11	0.41	1300	0.80
SRC pnSRC	0.85	0.84	1.14	0.49	1350	0.82

- Better overall quality of the SRC fit

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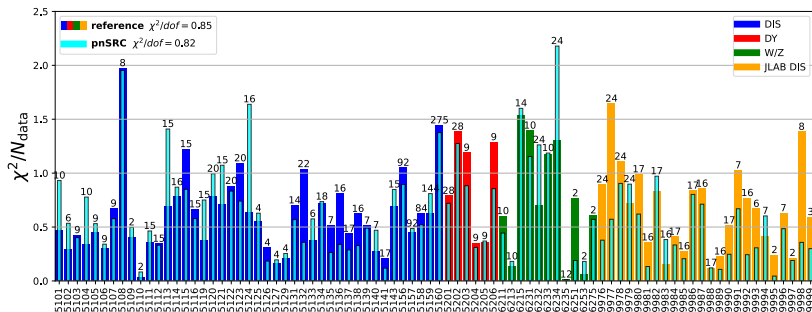
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- Slightly worse for W/Z bosons, which probe lower x

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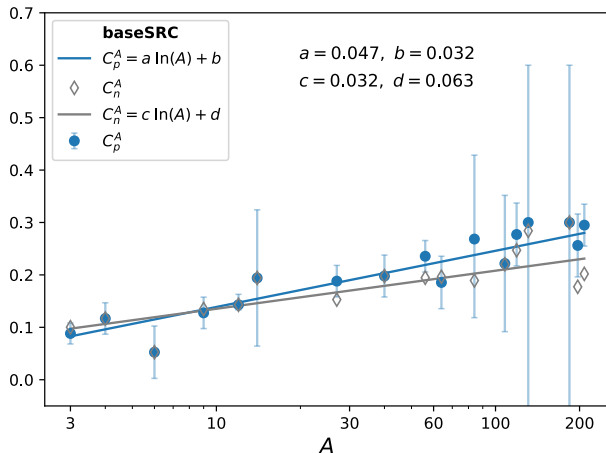
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A-dependence of C_p^A and C_n^A

A. Denniston, T. Jezo, A. Kusina, MK et al., Phys. Rev. Lett. 133 (2024) 15 [2312.16293]



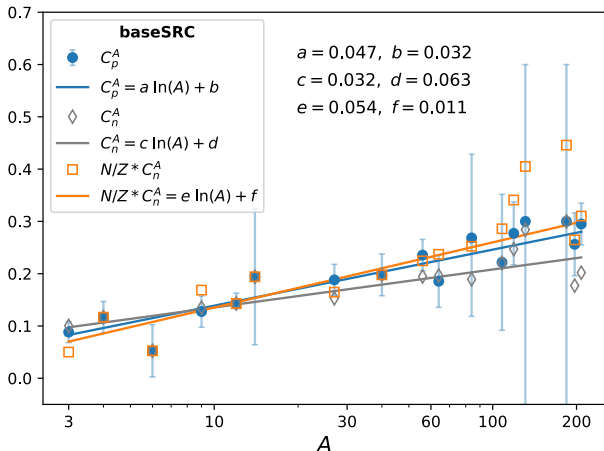
Almost equal number of protons and neutrons in SRC pairs:

- $^{197}_{79}\text{Au}$ ($C_p^A = 0.256, C_n^A = 0.178$): $79 C_p^A \simeq 20.2, 118 C_n^A \simeq 21.0$

- $^{208}_{82}\text{Pb}$ ($C_p^A = 0.295, C_n^A = 0.202$): $82 C_p^A \simeq 24.2, 126 C_n^A \simeq 25.5$

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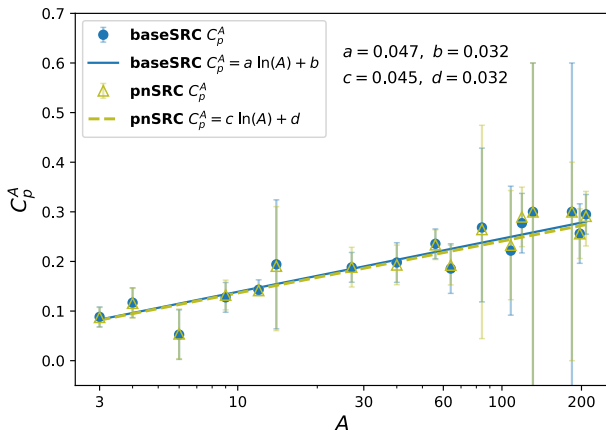
Even better agreement after correcting for neutron excess \rightarrow

Allows to reduce the number of free parameters: $C_n^A = (Z/N)C_p^A$.

Agrees with exclusive quasi-elastic scattering ($pn/NN \simeq 90 \pm 10\%$).

Comparison of the two fits

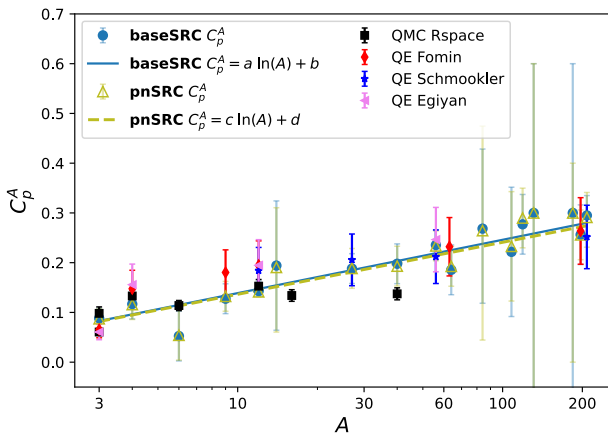
A. Denniston, T. Jezo, A. Kusina, MK et al., Phys. Rev. Lett. 133 (2024) 15 [2312.16293]



pnSRC fit very similar to baseSRC fit: $\chi^2 = 0.82/N_{\text{dof}}$ instead of 0.80.

Comparison with QE scattering and QMC calculations

A. Denniston, T. Jezo, A. Kusina, MK et al., Phys. Rev. Lett. 133 (2024) 15 [2312.16293]



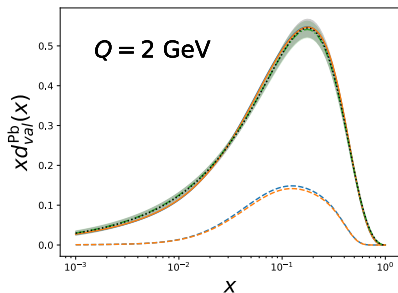
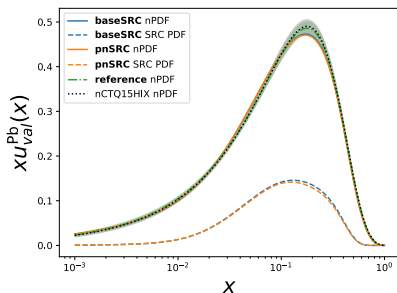
Excellent agreement with

- QE data [Fomin, PRL 108 (2012) 092502; Schmookler, Nature 566 (2019) 354; Egiyan, PRL 96 (2006)]
- QMC theory [R. Cruz-Torres et al., Nature Physics 17 (2021) 306]

Valence quark distributions

A. Denniston, T. Jezo, A. Kusina, MK et al., Phys. Rev. Lett. 133 (2024) 15 [2312.16293]

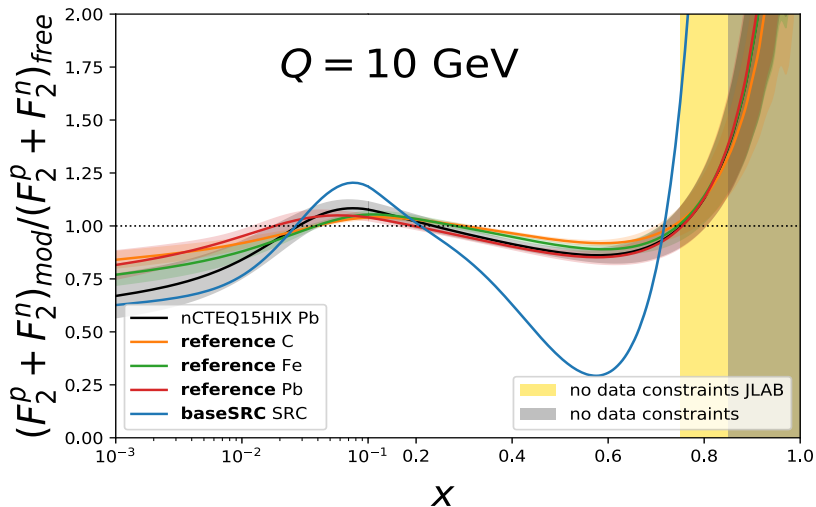
Sanity check:



SRC u, d valence quarks in Pb within previous error bands.
SRC components 20-30% of full nPDFs in agreement with $C_{p,n}^A$.

Ratio of bound to free pn -pair structure functions

A. Denniston, T. Jezo, A. Kusina, MK et al., Phys. Rev. Lett. 133 (2024) 15 [2312.16293]



Similar, but SRC more pronounced in EMC-region and universal!

Conclusion and outlook

Nuclear binding at low energy:

- Off-shell 2-body \leftrightarrow 3-body interactions \rightarrow Shell model
- Quasi-free nucleons in Mean Field motivate bound N PDFs
- But: Strong experimental and theoretical evidence for SRCs
- Striking correlation of QES and DIS in EMC region

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nCTEQ15HIX:

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nCTEQ15SRC:

- nCTEQ15HIX data + W/Z bosons \rightarrow 2 more parameters (s)
- Factorisation of energy scales \rightarrow **New SRC ansatz** for nPDFs
- **Very** good fit, in particular to JLab data, and consistent PDFs
- **p and n fractions** in SRC pairs **agree with LE data and theory**
- Dominated by **pn pairs**, again in agreement with LE data
- Partonic structure of SRC pairs **universal**, large EMC effect

Outlook

Future directions:

[O. Hen et al., Rev. Mod. Phys. 89 (2017) 045002]

- Test isospin breaking ($C_n^A \neq \frac{Z}{N} C_p^A$, $f_n \neq f_p$) with QED/PV
- Separate $C_{p,n/(pn)}^A$ and $C_{p,n/(pp,nn)}^A$ with tagging, mirror nuclei
- Separate spin 1 and spin 0 with polarised eA scattering

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Neutrino CC scattering:

[C. Ciofi degli Atti, Phys. Rep. 590 (2015) 1]

- π^0 production > baryon resonances [MiniBooNE, PRD 83 (2011) 052009]
- p production > QE prediction \rightarrow pp? [Minerva, PRL 111 (2013) 022502]

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Collective effects in QGP:

[C. Ciofi degli Atti, Phys. Rep. 590 (2015) 1]

- Glauber model of individual wounded nucleons (participants)
- SRCs reduce eccentricities [G. Denicol et al., 1406.7792]