

# Non-Linearity of Deflecting Field in LHC crab-cavities

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# Acknowledgements

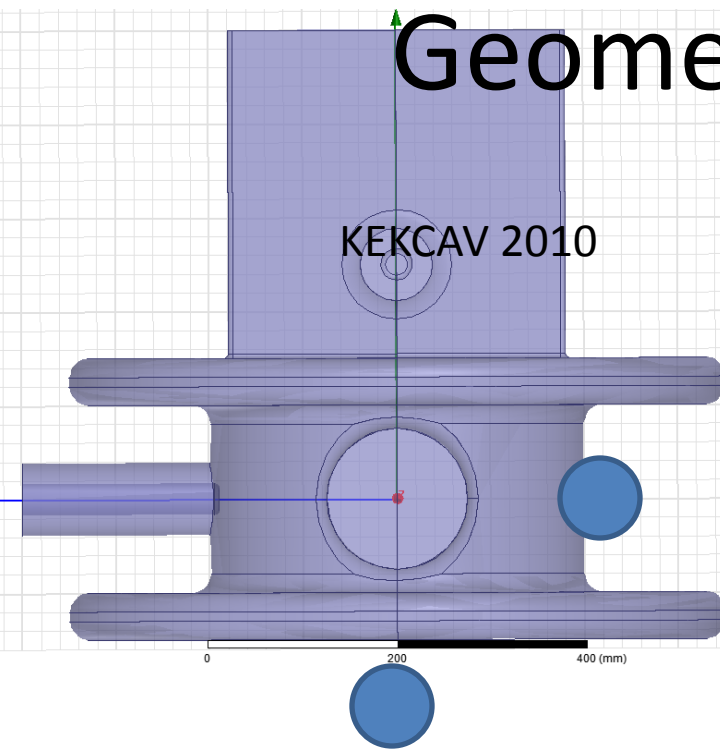
- Many thanks to Rama Calaga for stimulating discussions and introduction to the LHC crab cavity. Also for providing 3D models of the ODUCAV, KEKCAV, UKCAV, SRHW, QWAVE
- And to Riccardo de Maria for introducing to LHC bunch-crossing schemes
- 3D HFSS simulations were performed on the CLIC study workstation for HFSS simulations (128GB, 24CPU)

# Outline

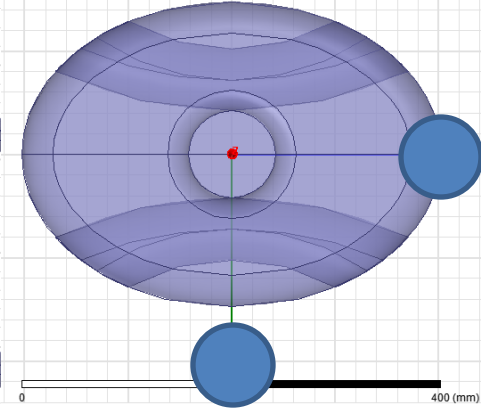
- Cavity geometry and mesh
- Non-linearity of the deflecting field treatment in terms of multipoles
- An example (KEKCAV)
- Comparison of RF multipoles to typical magnetic multipoles
- Conclusion
- If time permit, yet another compact crab cavity design

# Geometry of the cavities

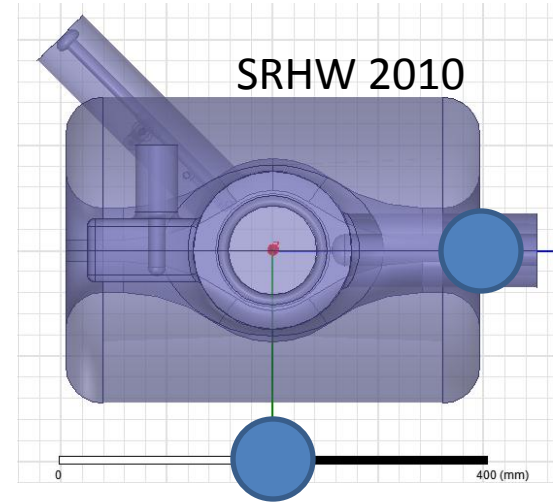
KEKCAV 2010



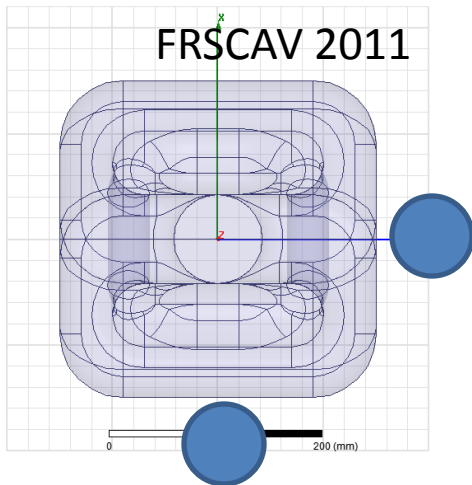
ODUCAV 2010



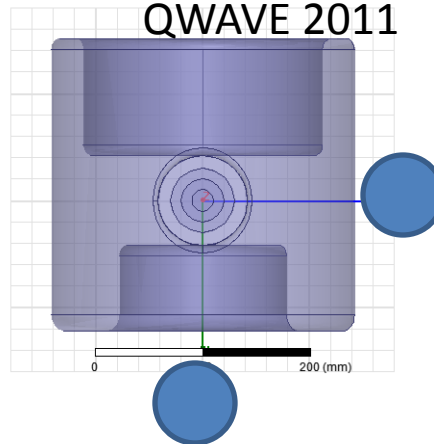
SRHW 2010



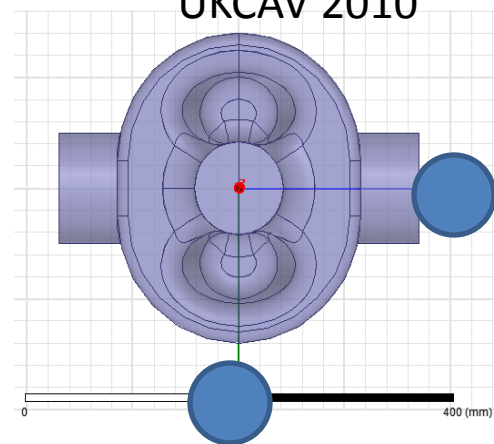
FRSCAV 2011



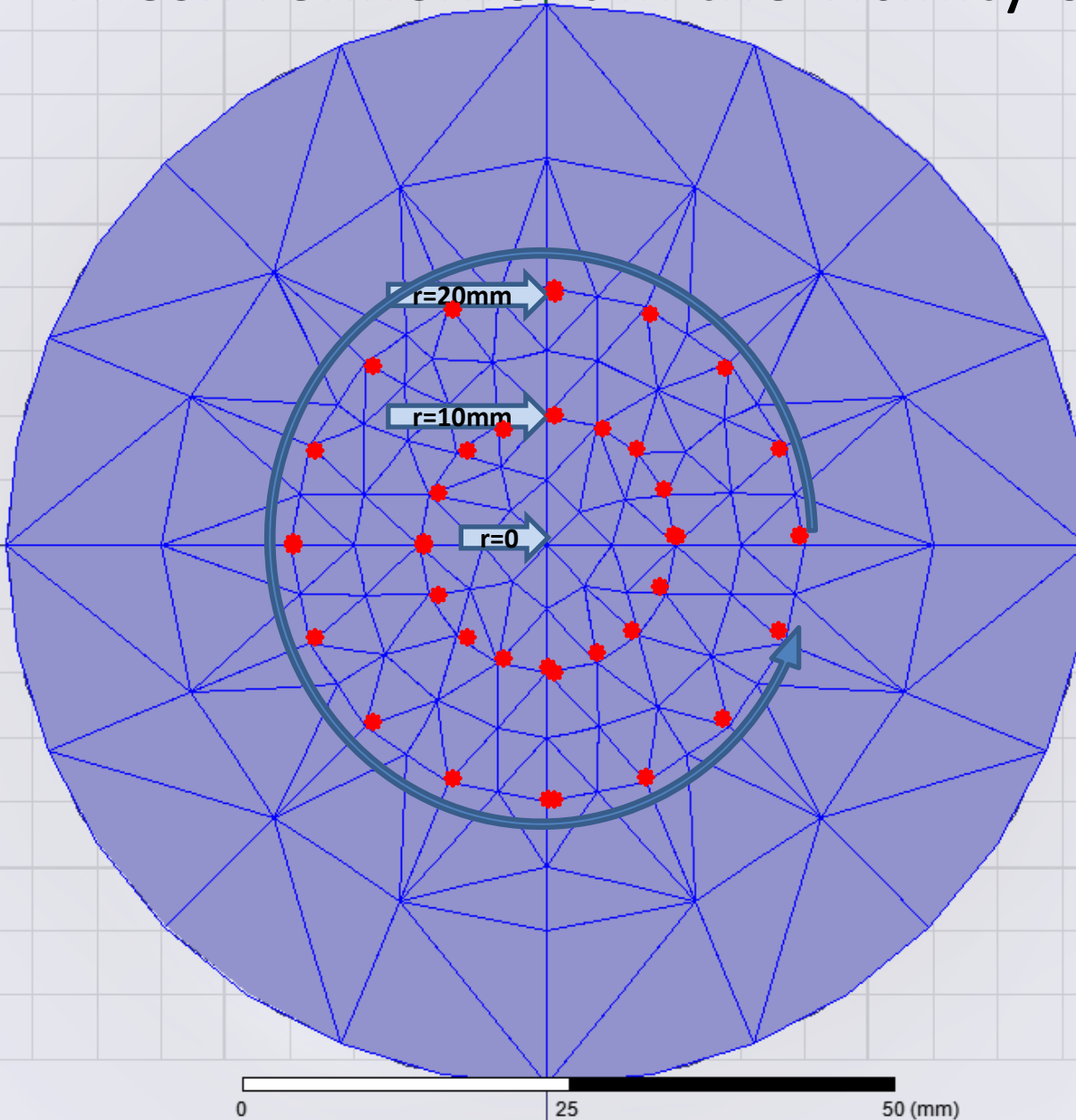
QWAVE 2011



UKCAV 2010



# Mesh refinement in the vicinity of the beam axis



- Surface approximation of 1mm is used
- 3 coaxial polyhedrones (n=16) of 10, 20, 30 mm radii were used for mesh symmetrisation near the axis

- Length based mesh refinement of 7 mm inside  $r=20\text{mm}$  and of 15 mm inside  $r=30\text{mm}$

# Multipole expansion of $E_z$

Accelerating gradient:

Multipole expansion:

$$E_{acc}(r, \varphi, z) = E_z(r, \varphi, z) \cdot e^{j\frac{\omega}{c}z}$$

$$E_{acc}(r, \varphi, z) = \sum_n E_{acc}^{(n)} r^n e^{jn\varphi} = \sum_n E_{acc}^{(n)} r^n \cos(n\varphi)$$

Skew components = 0  
due to the symmetry

**Panofsky-Wenzel (PW) theorem:**

$$p_{\perp}(r, \varphi) = \frac{+je}{\omega} \int_0^L dz [\nabla_{\perp} E_{acc}(r, \varphi, z)]; \quad \text{for } E \sim e^{+j\omega t}$$

$$\text{where: } \nabla_{\perp} = \vec{u}_r \frac{\partial}{\partial r} + \vec{u}_{\varphi} \frac{1}{r} \frac{\partial}{\partial \varphi}$$

Gives an expression for multipolar RF kicks:

$$p_{\perp}^{(n)}(r, \varphi) = \frac{je}{\omega} nr^{n-1} [\vec{u}_r \cos(n\varphi) + \vec{u}_{\varphi} \sin(n\varphi)] \int_0^L E_{acc}^{(n)}(z) dz$$

**Lorenz Force (LF):**

Gives an expression for kick directly from the RF EM fields:

$$F_{\perp} = e[E_{\perp} + v_z \times B_{\perp}] = \Big|_{v_z=c} e[E_{\perp} + Z_0 \vec{u}_z \times H_{\perp}]$$

$$p_{\perp}(r, \varphi) = \int_0^L \frac{F_{\perp}}{v_z} dz = \Big|_{v_z=c} \frac{e}{c} \int_0^L [E_{\perp} + Z_0 \vec{u}_z \times H_{\perp}] dz$$

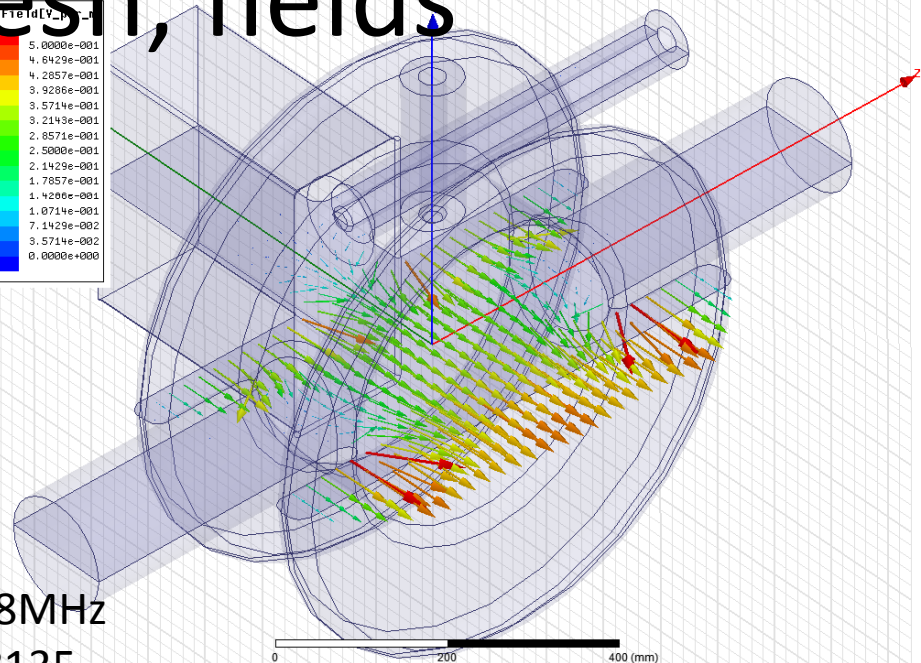
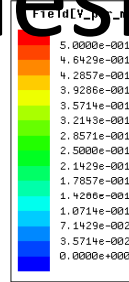
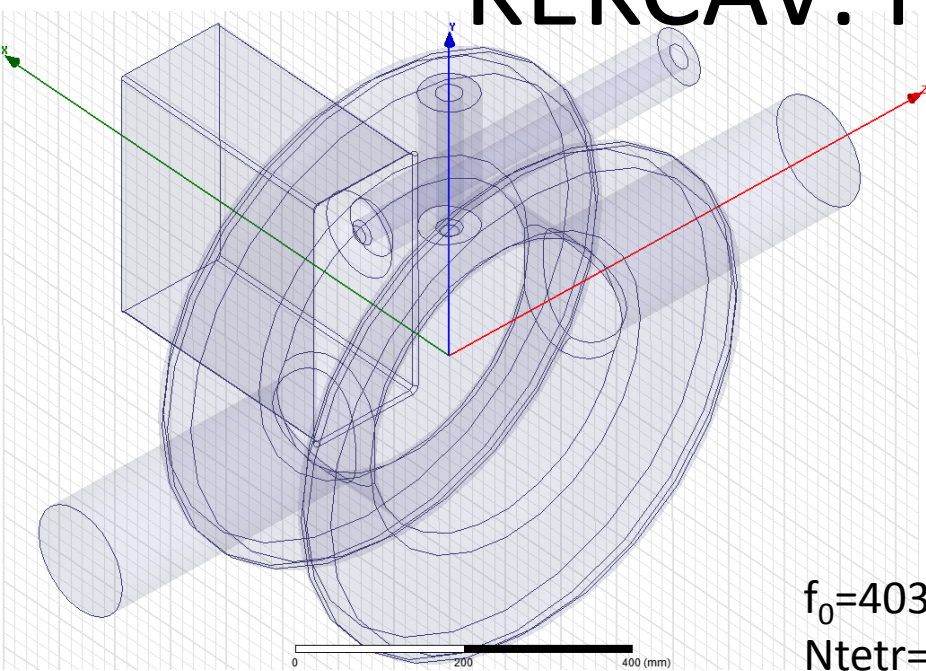
Which can be decomposed into multipoles:

$$p_{\perp}^{(n)}(r, \varphi) = \frac{1}{c} r^{n-1} [\vec{u}_r \cos(n\varphi) + \vec{u}_{\varphi} \sin(n\varphi)] \int_0^L F_{\perp}^{(n)} dz$$

Equating the RF and magnetic kicks, RF kick strength can be expressed in magnetic units:

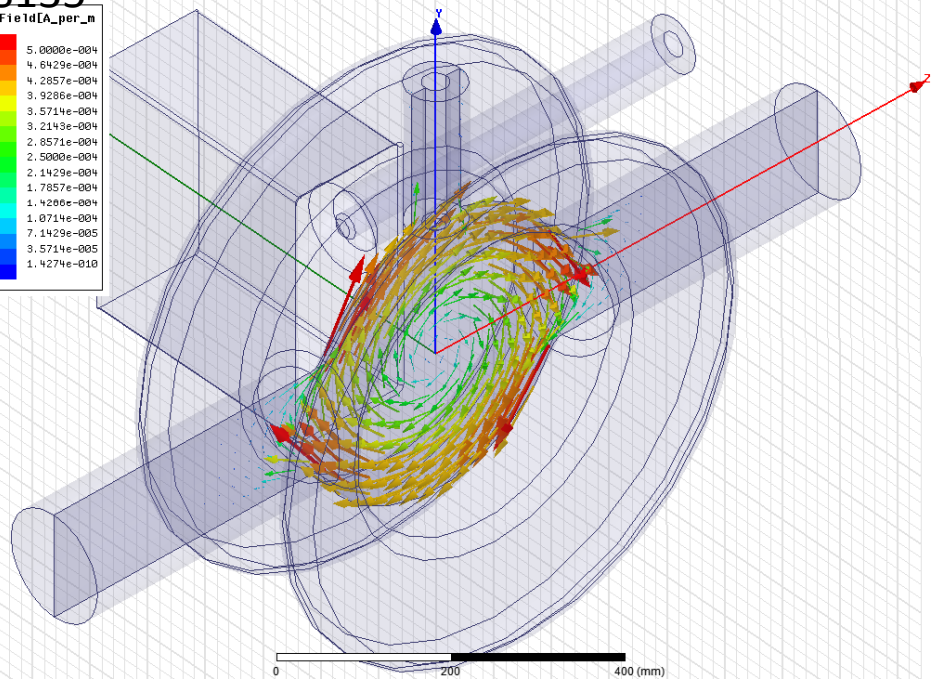
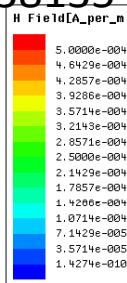
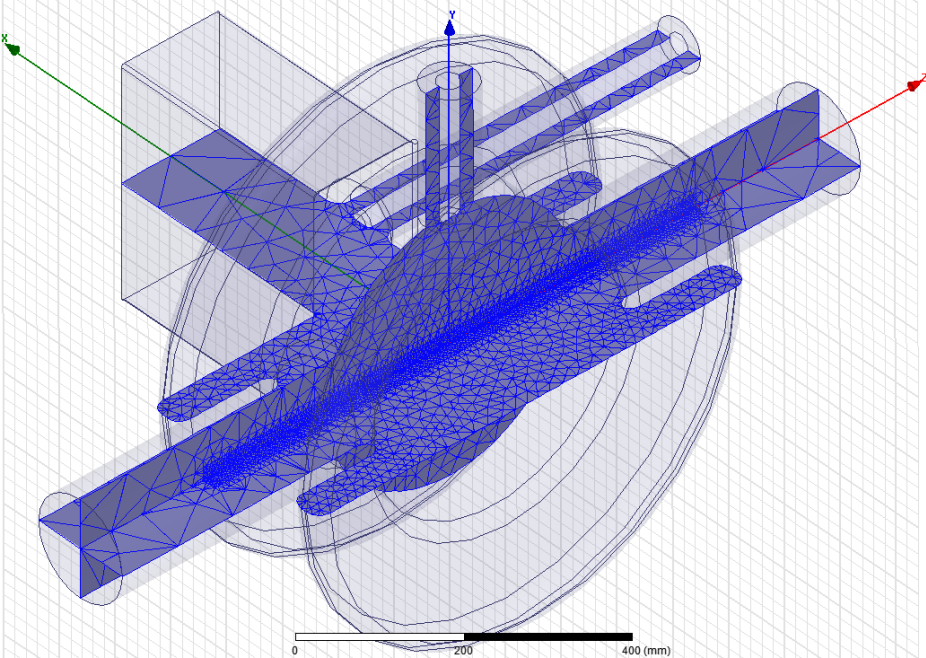
$$B_{\perp}^{(n)} = \frac{1}{ec} F_{\perp}^{(n)} = \frac{nj}{\omega} E_{acc}^{(n)} \quad [Tm / m^n]$$

# KEKCAV: mesh, fields

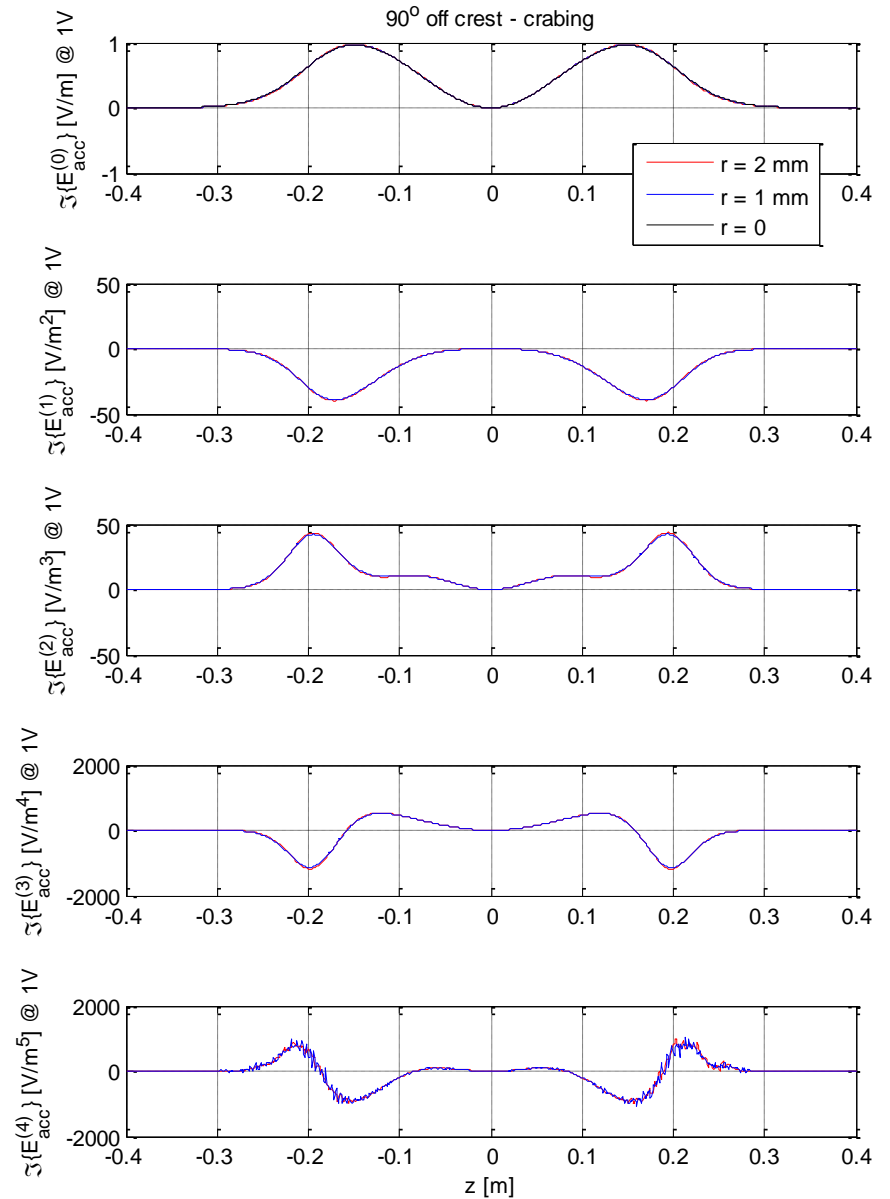
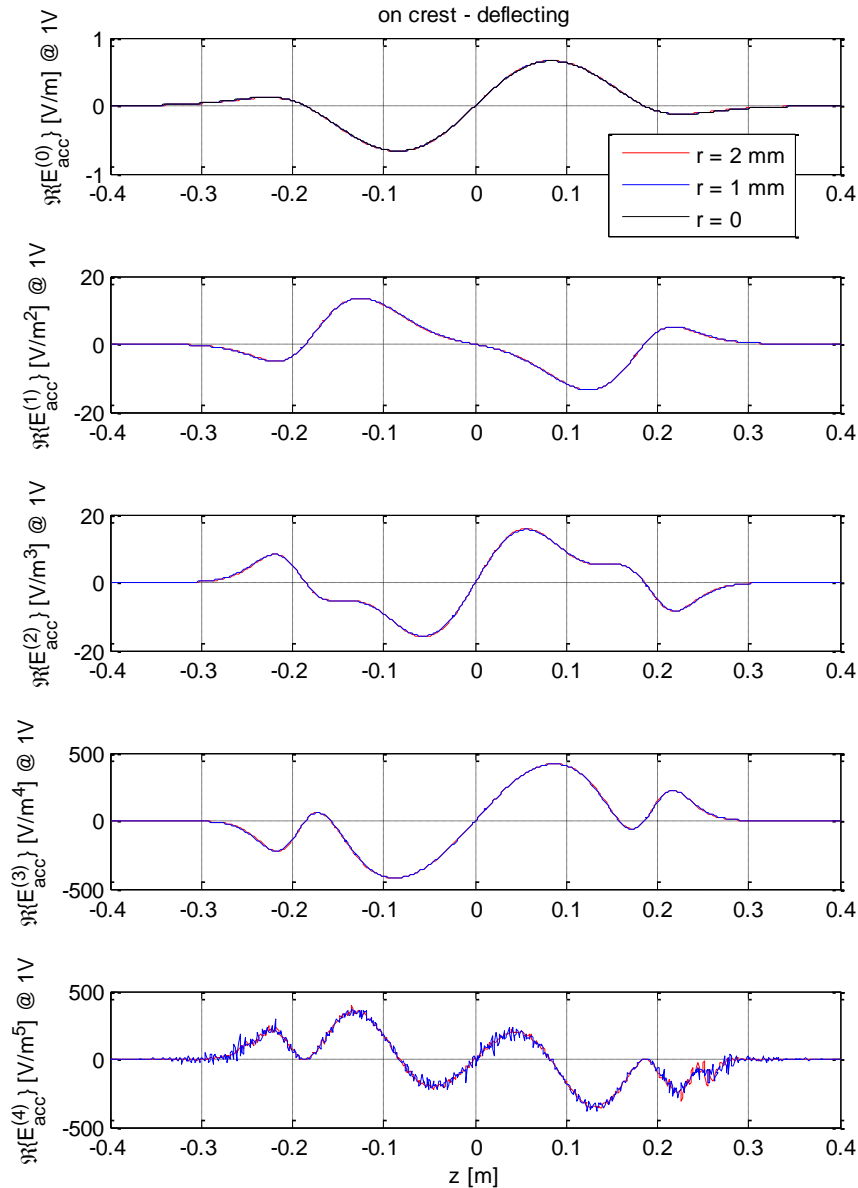


$f_0=403.358\text{MHz}$

Ntetr=338135



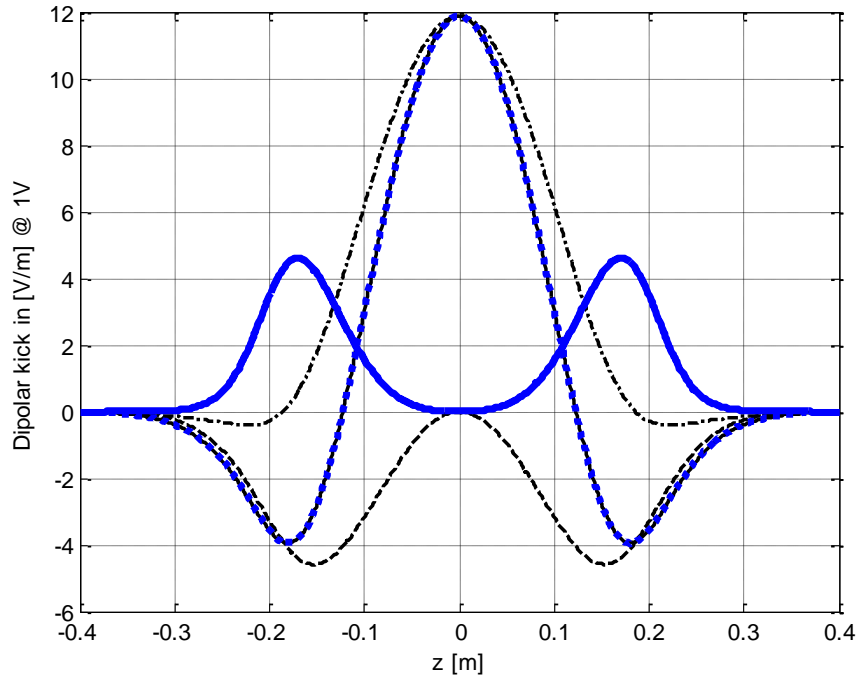
# KEKCAV: multipoles of $E_{\text{acc}}$ at $V_x=1\text{V}$



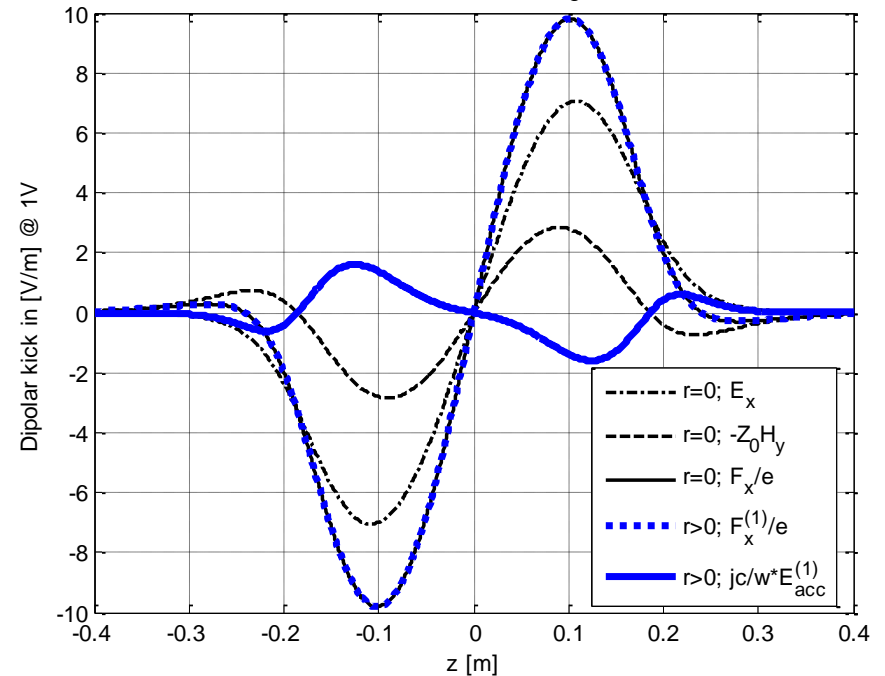


# KEKCAV: dipolar kick, PW versus LF

on crest - deflecting



90° off crest - crabbing



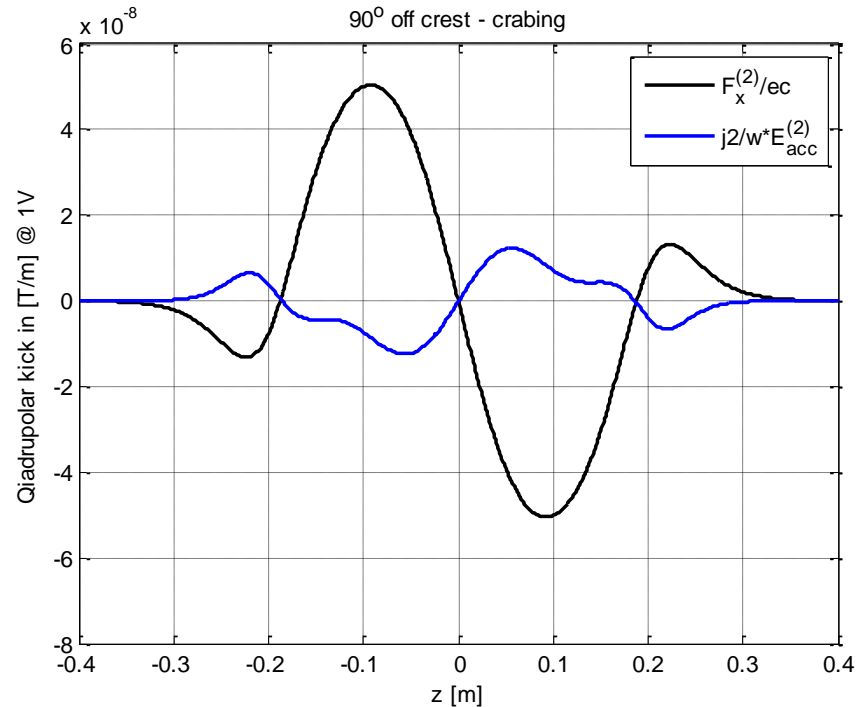
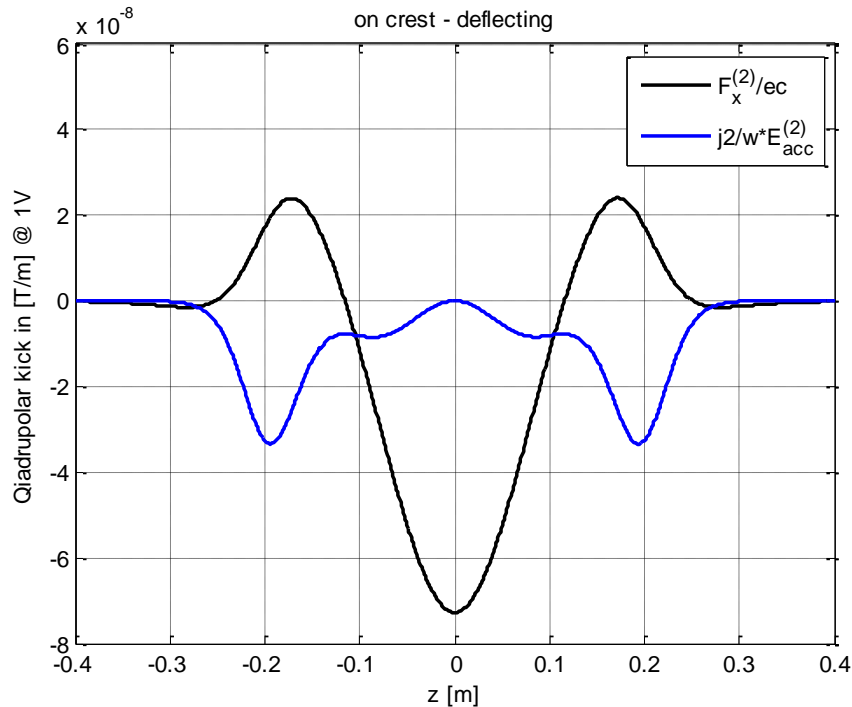
- Dipolar kick strength  $F_x$  and corresponding multipole of  $E_{acc}^{(1)}$  have very different dependence along the beam axis but the integrals are equal.
- Dipolar kick is mainly electric, but is reduced significantly by the magnetic part

Comparison  $V_x$

LF:  $1.0000 + 0.0000i$  [V]

PW:  $1.0018 + 0.0000i$  [V]

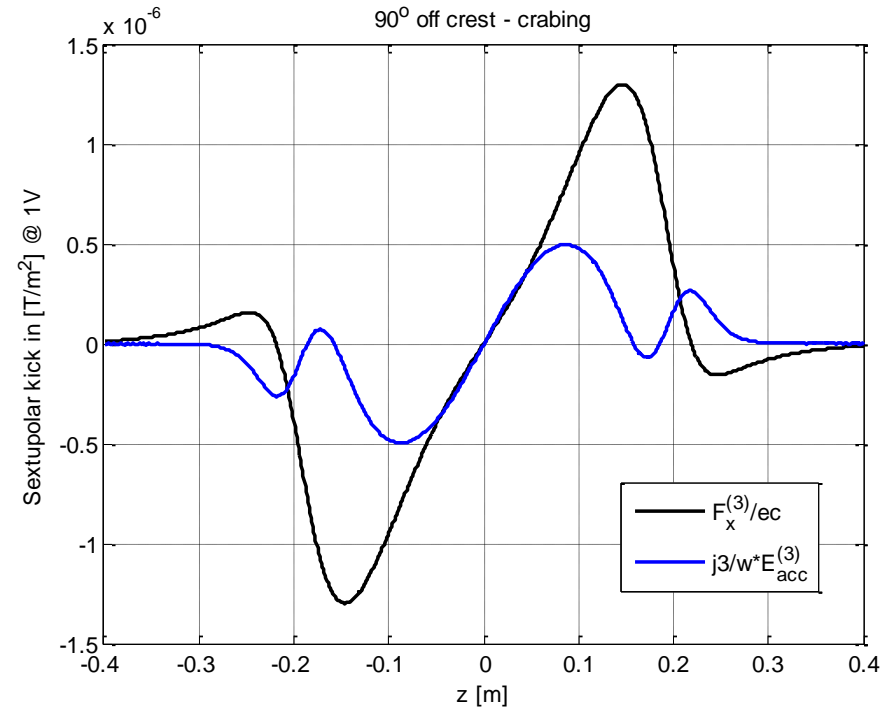
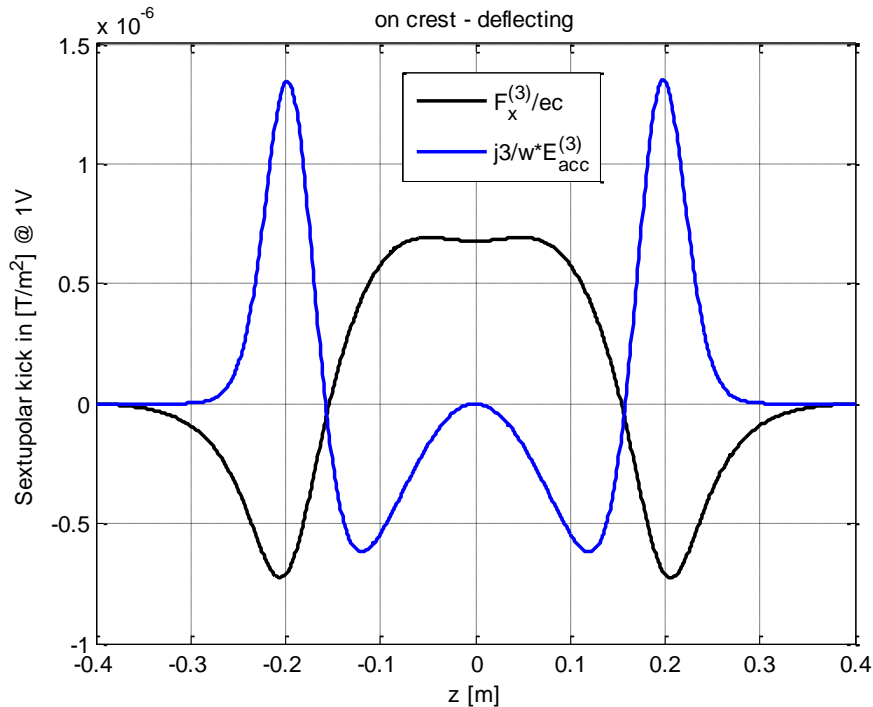
# KEKCAV: quadrupolar kick, PW versus LF



Quadrupolar kick strength  $F_x^{(2)}$  and corresponding multipole of  $E_{acc}^{(2)}$  have very different dependence along the beam axis but the integrals are equal.

Comparison  $\int B^{(2)} dz$  @  $V_x=1V$   
 LF:  $-6.53 -0.02i$  [nTm/m]  
 PW:  $-6.54 -0.02i$  [nTm/m]

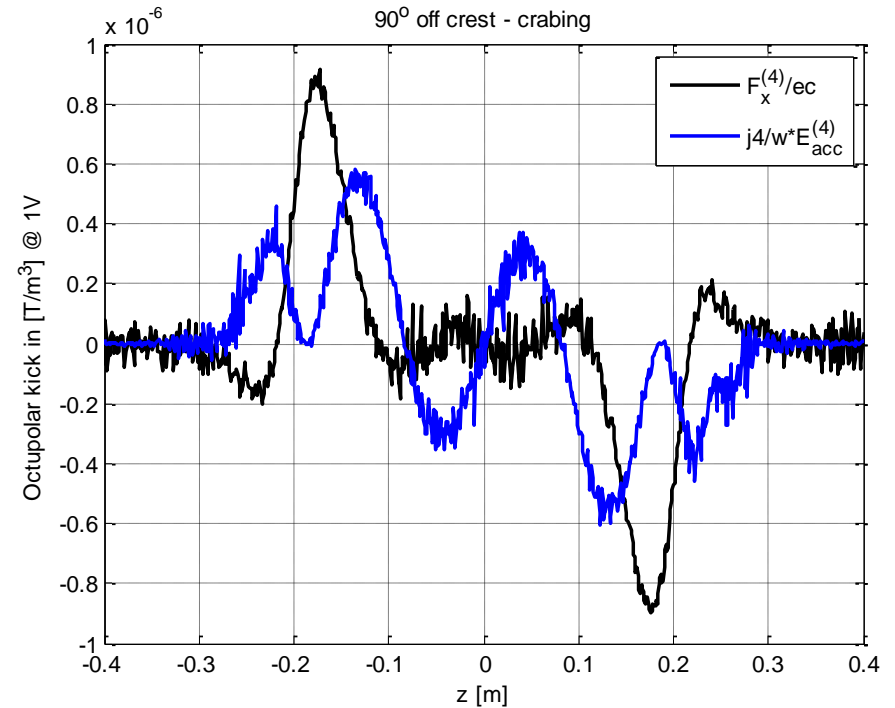
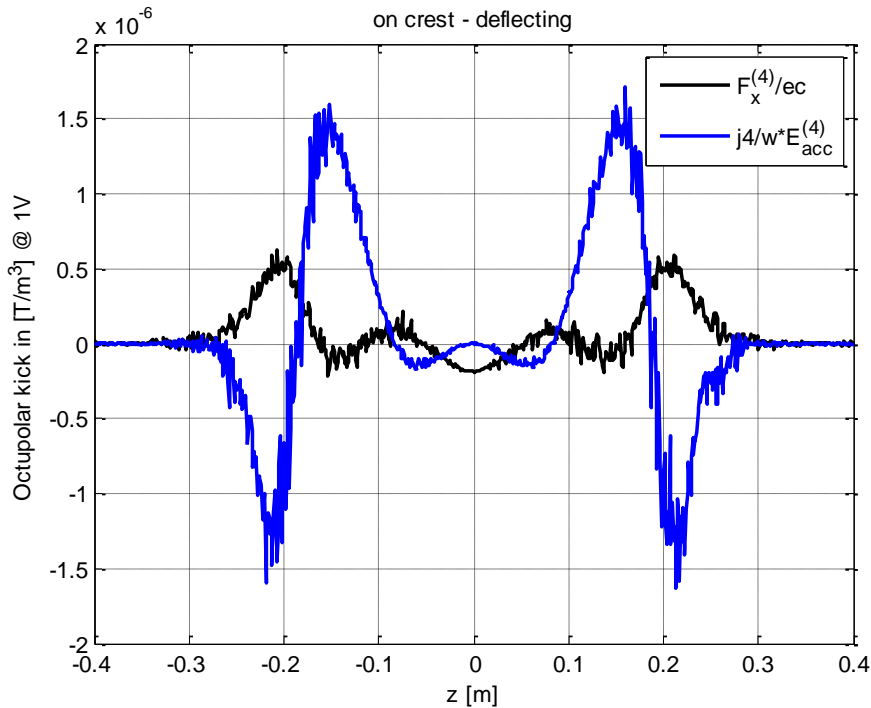
# KEKCAV: sextupolar kick, PW versus LF



Sextupolar kick strength  $F_x^{(3)}$  and corresponding multipole of  $E_{acc}^{(3)}$  have very different dependence along the beam axis but the integrals are equal.

Comparison  $\int B^{(3)} dz$  @  $V_x=1V$   
 LF: 49.9 -0.08i  $[nTm/m^2]$   
 PW: 50.2 -0.10i  $[nTm/m^2]$

# KEKCAV: octupolar kick, PW versus LF



Octupolar kick strength  $F_x^{(4)}$  and corresponding multipole of  $E_{acc}^{(4)}$  have very different dependence along the beam axis but the integrals are equal.

Comparison  $\int B^{(4)} dz$  @  $V_x=1\text{V}$   
 LF: 53.1 -0.9i  $[\text{nTm}/\text{m}^3]$   
 PW: 56.2 -1.4i  $[\text{nTm}/\text{m}^3]$

# Summary table for $V_x = 5\text{MV}$

Values are calculated from Lorenz force

	ODUCAV	SRHW	KEKCAV	UKCAV	QWAVER	FRSCAV
$V_z(x=0)$ [kV]	0.0	-2.1 - 2.5i	-4 + <b>1378i</b>	0.0	0 + 85.7i	-0.1 - 0.2i
$V_x$ [MV]	5	5	5	5	5	5
$B^{(2)}$ [mTm/m]	0	0 - 0.04i	-32.7 - 0.1i	0.02 + 0i	25 + 0i	0 + <b>108i</b>
$B^{(3)}$ [mTm/m <sup>2</sup> ]	1250 + 0i	229 + 0i	250 - 0i	<b>2452</b> - 0.5i	464 + 0i	-233 + 1i
$B^{(4)}$ [mTm/m <sup>3</sup> ]	0	0	266 - 5i	0	540 + 0i	-189 - <b>14209i</b>



NB:  $B_y^{(n)}(y=0, x=a) = B^{(n)}a^{n-1}$ . This is not MAD convention for multipolar strength.

There is the following dependences of the multipolar kick on the RF phase, where  $\delta\phi_{crab}$  is the deviation of the (macro)particle RF phase from the crabbing phase

$$p_{\perp}^{(n)}(r, \varphi, \delta\phi_{crab}) = er^{n-1} \left[ \vec{u}_r \cos(n\varphi) + \vec{u}_{\varphi} \sin(n\varphi) \right] \mathfrak{I} \left\{ B^{(n)} e^{j\delta\phi_{crab}} \right\}$$

# Multipole errors in 1e-4 units at R=17mm

	ODUCAV		SRHW		KEKCAV		UKCAV		QWAVER		FRSCAV	
	real	imag	real	imag	real	imag	real	imag	real	imag	real	imag
Vx [MV]	5	0	5	0	5	0	5	0	5	0	5	0
B <sup>(2)</sup> [mTm/m]	0	0	0	-0.04	-32.7	-0.1	0.02	0	25	0	0	<b>108</b>
B <sup>(3)</sup> [mTm/m <sup>2</sup> ]	1250	0	229	0	250	0	<b>2452</b>	-0.5	464	0	-233	1
B <sup>(4)</sup> [mTm/m <sup>3</sup> ]	0	0	0	0	266	-5	0	0	540	0	-189	<b>-14209</b>

## multipole error relative to dipole kick of the crab cavity itself

Vx [MV] => B1*L [mTm]	16.68	0	16.68	0	16.68	0	16.68	0	16.68	0	16.68	0
b2[1e-4@R17mm]	0	0	0	-0.41	-333.3	-1.02	0.204	0	254.8	0	0	<b>1100.8</b>
b3[1e-4@R17mm]	216.6	0	39.68	0	43.32	0	<b>424.9</b>	-0.09	80.40	0	-40.37	0.173
b4[1e-4@R17mm]	0	0	0	0	0.784	-0.01	0	0	1.591	0	-0.557	<b>-41.86</b>

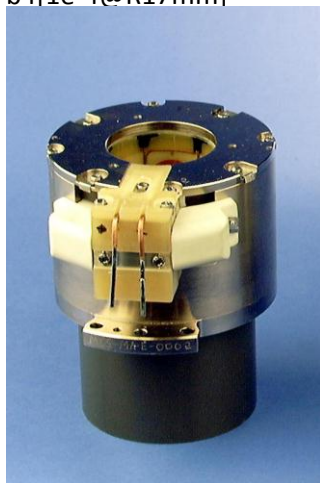
## multipole error relative to the LHC main dipole at injection 0.45 TeV

B1*L [Tm]	7.728	7.728	7.728	7.728	7.728	7.728	7.728	7.728	7.728	7.728	7.728	7.728
b2[1e-4@R17mm]	0	0	0	-0	-0.719	-0	4E-04	0	0.550	0	0	<b>2.376</b>
b3[1e-4@R17mm]	0.467	0	0.086	0	0.093	0	<b>0.917</b>	-0	0.174	0	-0.0871	0.0004
b4[1e-4@R17mm]	0	0	0	0	0.002	-0	0	0	0.003	0	-0.0012	<b>-0.090</b>

## multipole error relative to the LHC main dipole at 7TeV

B1*L [Tm]	119.2	119.2	119.2	119.2	119.2	119.2	119.2	119.2	119.2	119.2	119.2	119.2
b2[1e-4@R17mm]	0	0	0	-0	-0.047	-0	3E-05	0	0.036	0	0	<b>0.154</b>
b3[1e-4@R17mm]	0.03	0	0.006	0	0.006	0	<b>0.059</b>	-0	0.011	0	-0.006	2E-05
b4[1e-4@R17mm]	0	0	0	0	1E-04	-0	0	0	0.0002	0	-8E-05	<b>-0.006</b>

Multipole errors relative to the LHC main dipole at 7TeV is very small



For example, LHC sextupole corrector (MCS) parameters from the Blue Book: field strength: 1630 T/m<sup>2</sup>, magnetic length: 110 mm, integrated kick: 179300 mTm/m<sup>2</sup>

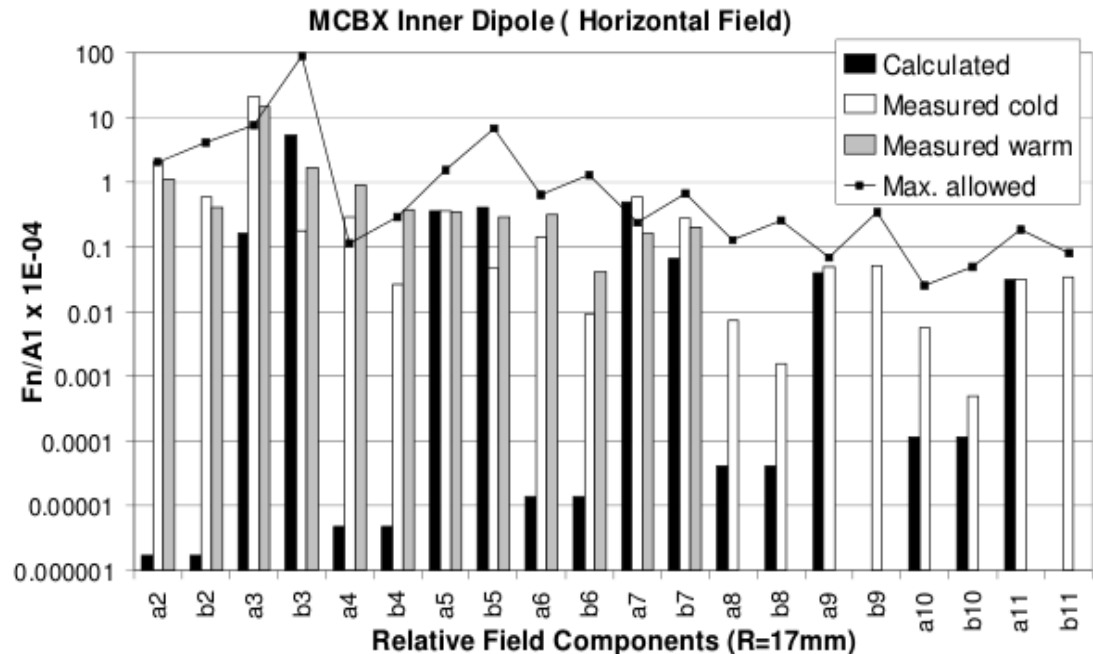
# Multipole errors in 1e-4 units at R=17mm

	ODUCAV		SRHW		KEKCAV		UKCAV		QWAVER		FRSCAV	
	real	imag	real	imag	real	imag	real	imag	real	imag	real	imag
$V_x$ [MV]	5	0	5	0	5	0	5	0	5	0	5	0
$B^{(2)}$ [mTm/m]	0	0	0	-0.04	-32.7	-0.1	0.02	0	25	0	0	<b>108</b>
$B^{(3)}$ [mTm/m <sup>2</sup> ]	1250	0	229	0	250	0	2452	-0.5	464	0	-233	1
$B^{(4)}$ [mTm/m <sup>3</sup> ]	0	0	0	0	266	-5	0	0	540	0	-189	-14209

multipole error relative to the LHC dipole corrector MCBXH

$B1 * L$ [Tm]	1.508	1.508	1.508	1.508	1.508	1.508	1.508	1.508	1.508	1.508	1.508	1.508
$b2[1e-4@R17mm]$	0	0	0	-0	-3.688	-0.01	0.002	0	2.8192	0	0	<b>12.179</b>
$b3[1e-4@R17mm]$	2.396	0	0.439	0	0.479	0	<b>4.701</b>	-0	0.8895	0	-0.4467	0.00192
$b4[1e-4@R17mm]$	0	0	0	0	0.009	-0	0	0	0.018	0	-0.0062	<b>-0.4631</b>

Multipole errors relative to the LHC dipole correctors MCBXH are comparable



# Conclusions

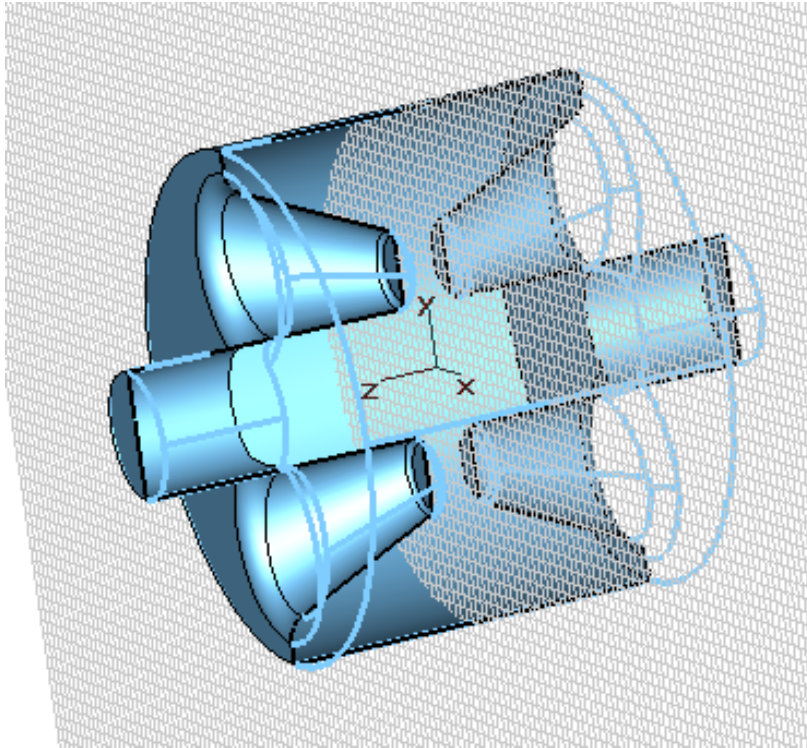
- Strength of the multipolar RF kicks associated with the EM field of deflecting mode has been calculated for ultra relativistic particles using HFSS code
- The strength has been expressed in terms of equivalent magnetic multipole strength giving the same transverse kick
- Whereas magnetic multipoles are constant in time the RF multipoles have harmonic dependence on the RF phase which has been introduced
- The calculated strength has been compared to the magnetic field strength/quality of different magnetic elements of LHC. The effect seems to be small. BUT those are static magnetic multipoles and not harmonic RF multipoles



Flat Rod Spoke cavity

# Idea

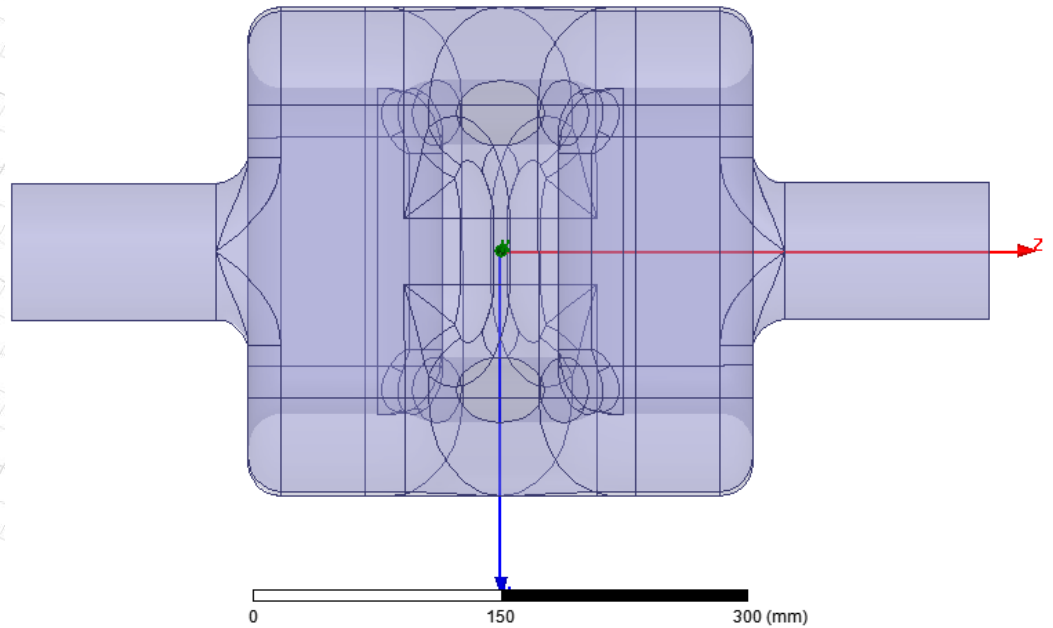
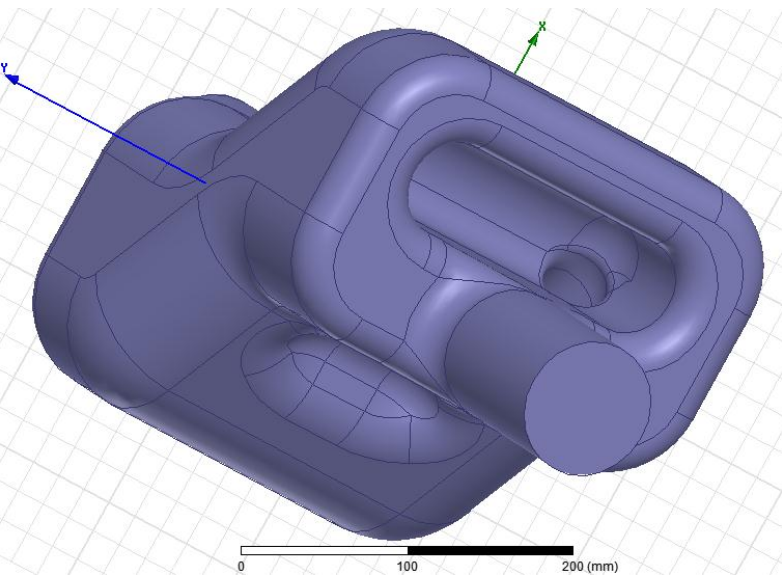
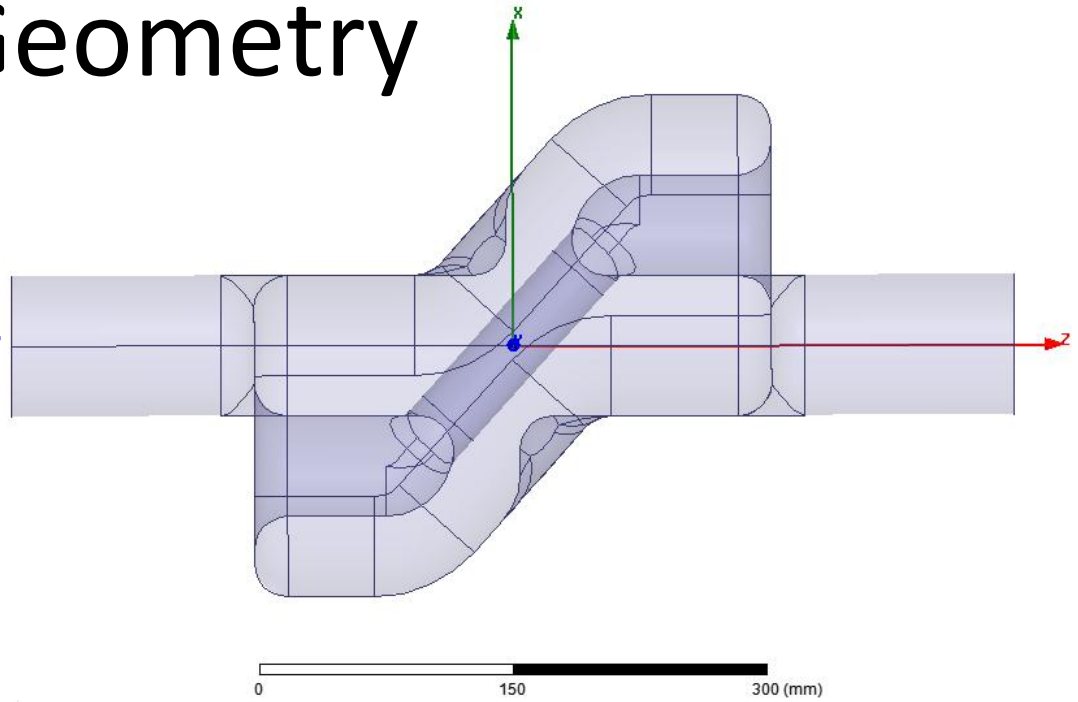
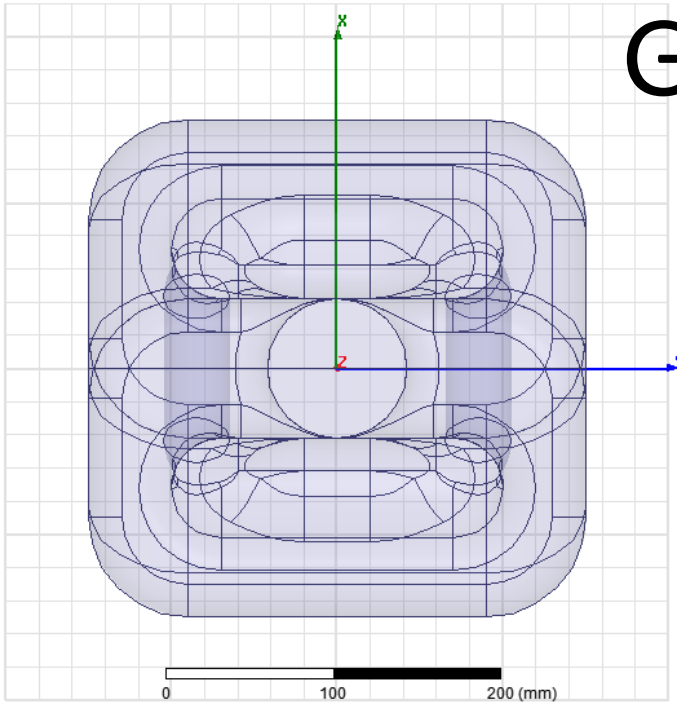
4-rod cavity = 4 x QW-resonators



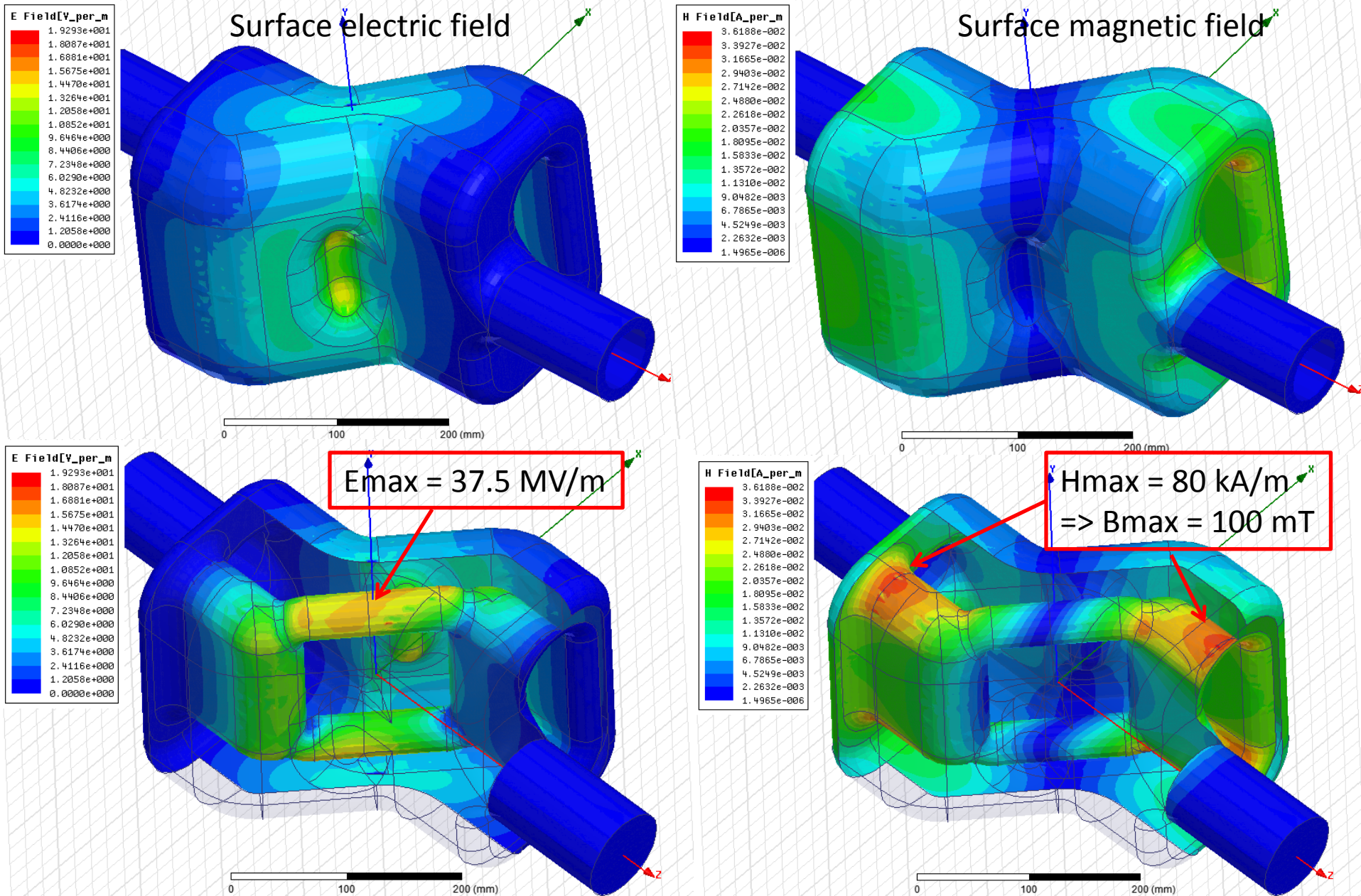
spoke cavity



# Geometry



# Surface EM fields for dipole kick: $V_x=2.5\text{MV}$



# R upon Q of the main mode and HOMs

