

Plasma Diagnostics of Ion Sources

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Diagnostics – The Window to the Knowledge



- ▶ **Langmuir probes:** ϕ_{pl} , n_e , T_e
- ▶ **Absorption techniques:** $n_s \rightarrow Cs, H^-$
- ▶ **Emission spectroscopy:** $n_s, T_s \rightarrow e, H, H_2, H^-$

Monitoring and Quantification – Spatial and Temporal Resolution

CERN Accelerator School, Senec, Slovakia

29th May – 8th June 2012

What do I want to know?

- ▶ identify the quantity (and the reason for it)
- ▶ define the required precision
- ▶ temporal behaviour and required time resolution
- ▶ necessity for spatial resolution
- ▶ ...



Adequate diagnostic technique?

- ▶ extensive or simple setup
- ▶ data acquisition and evaluation
- ▶ reliability
- ▶ costs and time (manpower) needed
- ▶ ...



Accessibility of the ion source?

- ▶ diagnostic ports
- ▶ test stand or continuous operation
- ▶ risks and feasibility
- ▶ reliability
- ▶ ...

invasive – non-invasive ; active – passive ; basic – specific parameters

Method	Standard	Sophisticated	Extras
Langmuir probe	single probe (cylindrical or planar)	double or triple probe emissive probe	special method: Boyd-Twiddy
emission spectroscopy	optical wavelength range with fibre optics & survey spectrometer	extended wavelength range VUV, UV or IF	sophisticated system spectral resolution, type of detector
absorption spectroscopy	white light absorption technique	tunable laser absorption	cavity-ringdown spectroscopy
Laser methods	laser induced fluorescence	TDLAS	
mass spectrometry	residual gas analyser	energy resolved mass spectrometry	

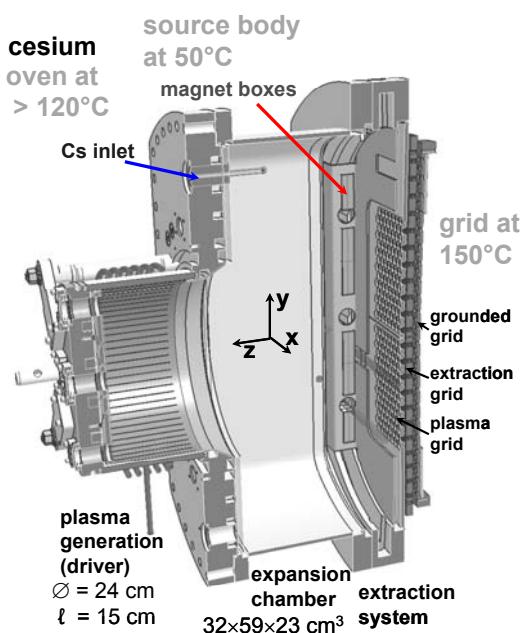
Introduction into techniques and applications for example in [1] – [7]

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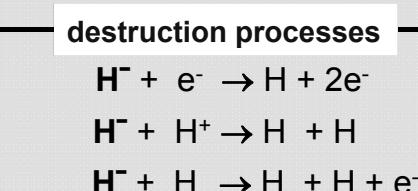
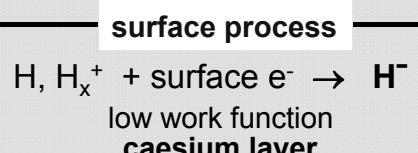
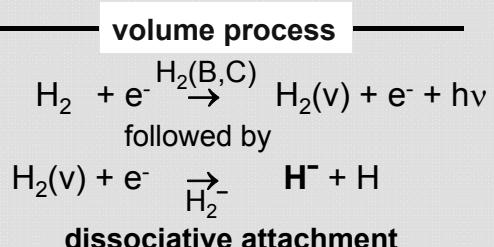
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ICP: f = 1 MHz, P = 70 kW, p = 0.3 Pa

6s plasma (4s beam), every 3 min: BATMAN
cw, up to 1 hour, every 3 min: MANITU

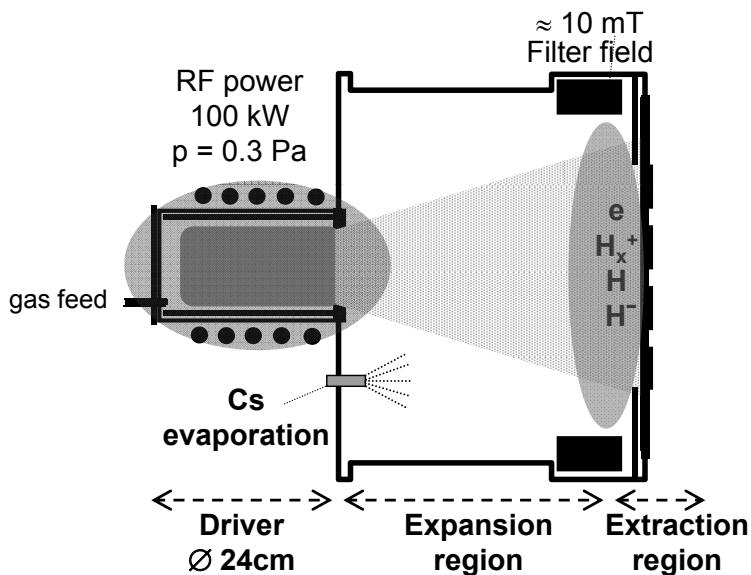


H⁻ formation and losses ...



... determine source optimisation

Ion sources for negative hydrogen ions: ionising – recombining plasma



$H, H_x^+ + \text{surface } e^- \rightarrow H^-$
Cs evaporation → low work function

Plasma generation ionising plasma

- ▶ ionisation: $\alpha \approx 0.1$
- ▶ dissociation: $\delta \approx 0.3$
- ▶ $T_e \approx 10 \text{ eV}, n_e \approx 5 \times 10^{18} \text{ m}^{-3}$

H^- generation recombining plasma

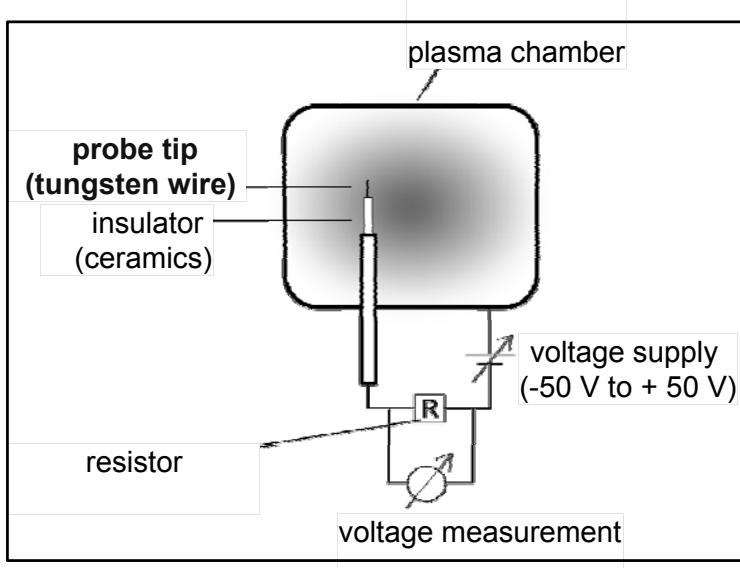
- ▶ $T_e \approx 1 \text{ eV}, n_e \approx 5 \times 10^{17} \text{ m}^{-3}$
- ▶ $H^-/n_e \approx 0.1 - 5$
- ▶ $Cs^+/n_e \approx 0.01 - 0.1$

Main issues

- ▶ Production and destruction of negative ions
- ▶ Extraction of negative ions
- ▶ Reduction of co-extracted electrons

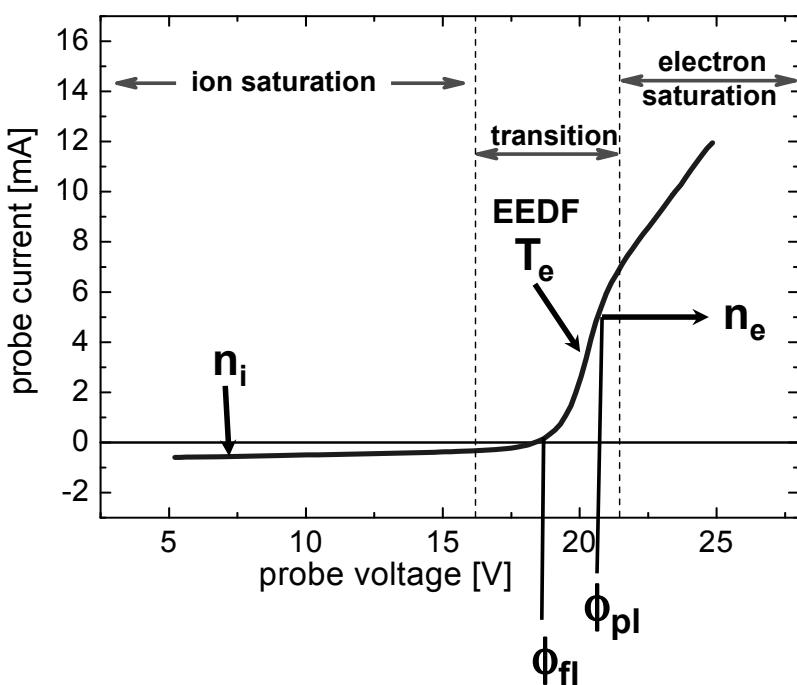
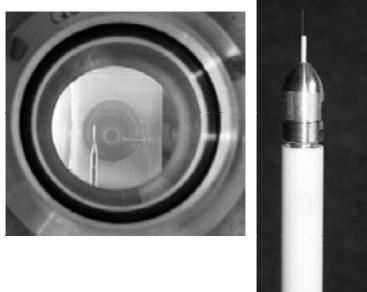
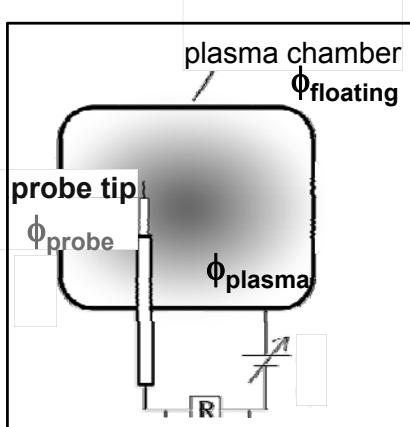
focus of diagnostics
(and modelling)
on plasma close to the grid

Main principle

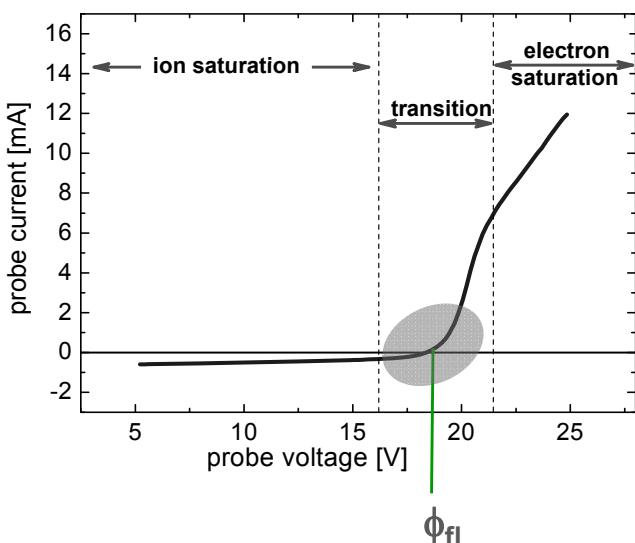


- ▶ stick a wire into the plasma tungsten, $\varnothing \approx 100 \mu\text{m}$, $l = 1 \text{ cm}$
- ▶ choose a reference electrode potential of plasma chamber
- ▶ apply a variable voltage typ. from - 50 V to + 50 V

$I - V$
characteristics

Main principle

Ursel Fantz, p. 7

CAS on Ion Sources, 6th June 2012**Floating potential**

- same fluxes for ions and electrons

$$\Gamma_{\text{ions}} = \Gamma_{\text{electrons}}$$

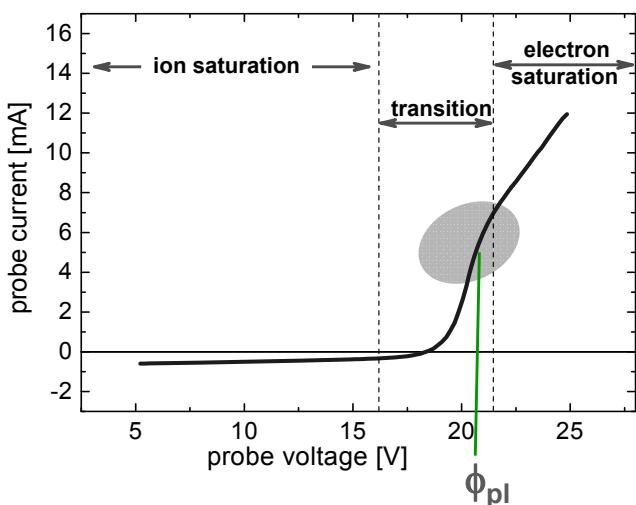
same currents

- no probe current



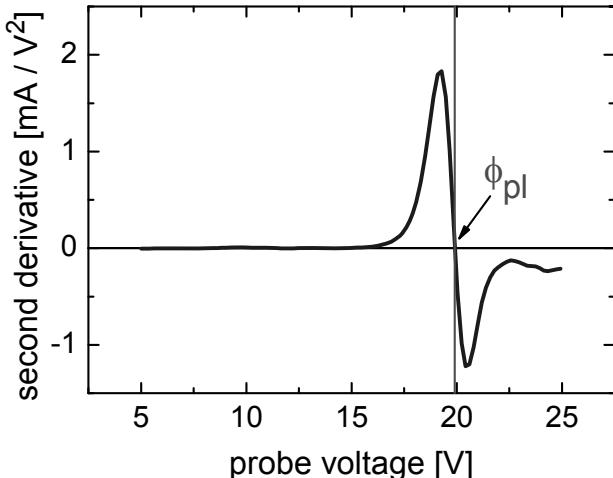
$$I_{\text{probe}} = 0 \rightarrow \phi_{\text{fl}}$$

Plasma potential

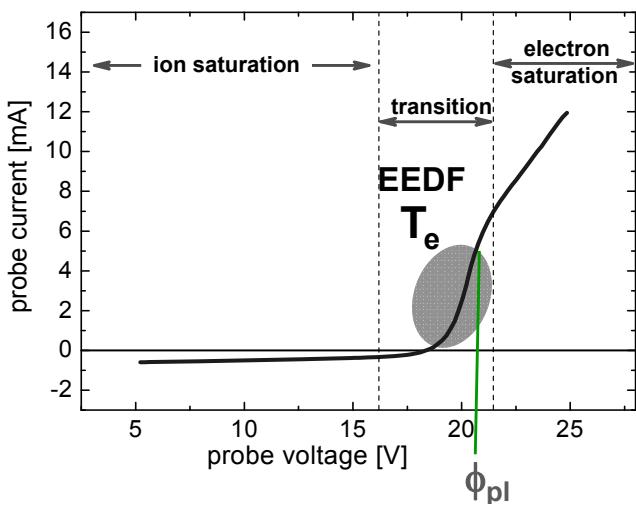


- for curves with high noise level
crossing of linear fits to
electrons saturation current and
transition close to turning point

- determined by ambipolar diffusion
 - turning point of I-V characteristics
 - zero-crossing of second derivative
- $d^2 I / d \phi_{pr}^2 \rightarrow \phi_{pl}$

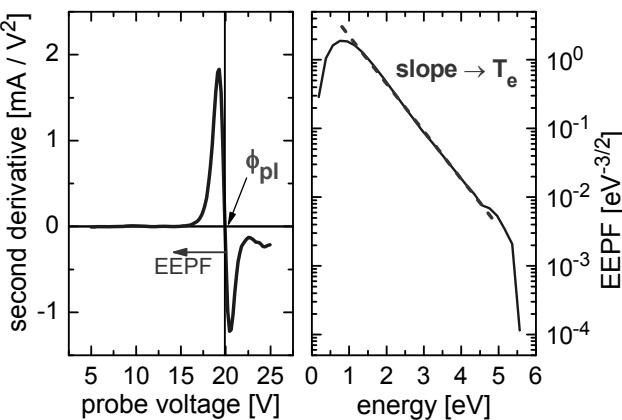


Electron energy distribution function (EEDF) and electron temperature

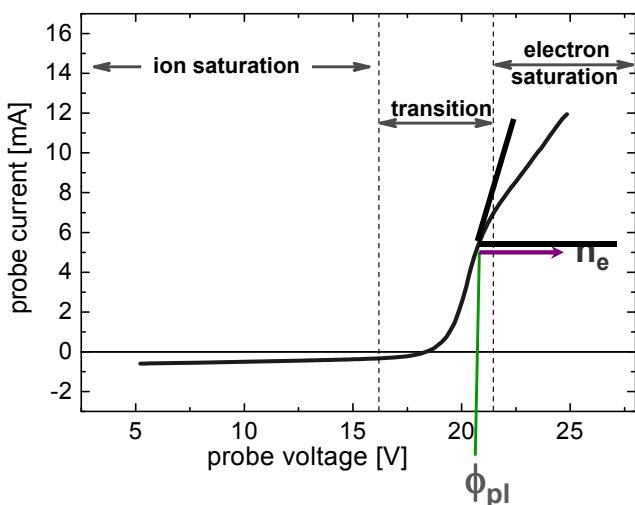


- T_e also from potential difference
 $\phi_{pl} - \phi_{fl} = k_b T_e \times \ln(\sqrt{m_i/(2\pi m_e)})$
 $\approx 2-3 \times T_e$

- distinguish: **EEDF = $\sqrt{E} \times EEPF$**
e.g.: Maxwell function probability function
for both: normalisation
- plot $\log(I_{pr})$ versus $E = \phi_{pl} - \phi_{pr}$
- slope yields T_e for Maxwell EEDF
 $d(\ln I) / dE = e / (k_b T_e)$



The electron density



- electron saturation current

$$I_{e,sat} = \frac{1}{4} n_e e v_e A_{eff}$$

- problem: effective probe area due to increase of plasma sheath



- take current at plasma potential
→ $A_{eff} = A_{probe}$

$$n_e = \frac{I(\phi_{pl})}{r_{pr} l_{pr} e} \sqrt{\frac{m_e}{2\pi k_B T_e}}$$

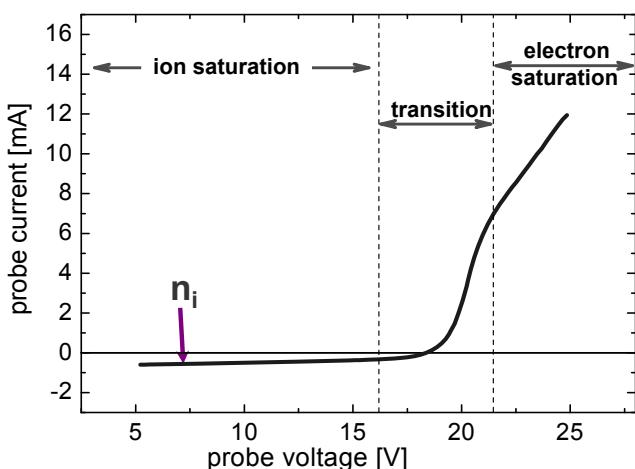
needs T_e

Probe geometry

influences shape of electron saturation current

- cylindrical probe (standard case)
- planar probe → $A_{pr} \gg$ sheath
- spherical probe

Ion density (positive ions)



- ion saturation current

- basically three theories available

OML: Orbital-Motion-Limited

ABR: Allen-Boyds-Reynolds

BRL: Bernstein-Rabinowitz-Langmuir

all of them assuming collision-less plasma sheath, i.e. $\lambda(\text{ions}) > r(\text{sheath})$

- simplest case: OML ($r_{pr} / r_{\text{sheath}} < 3$)

$$I_i = n_i e A_{pr} \sqrt{\frac{k_B T_e}{m_i}}$$

needs T_e

needs m_i

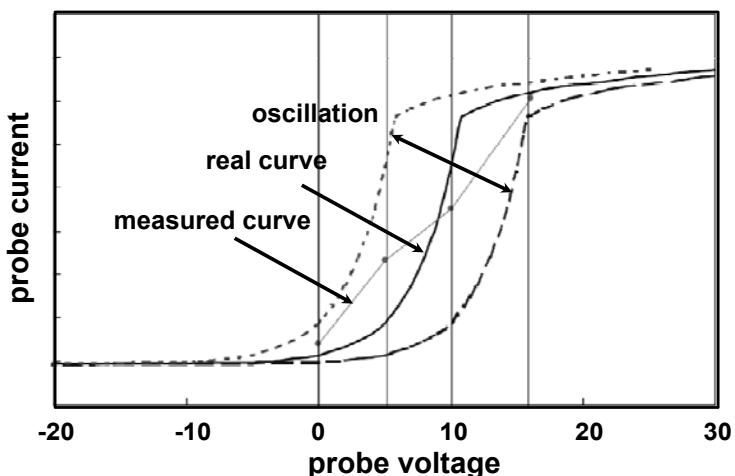
often unclear, e.g. hydrogen

→ H^+ , H_2^+ , H_3^+

- choose proper ϕ_{pr}
guide line: $\phi = \phi_{pl} - 10 \times kT_e$
- check if $n_e = n_i$ is fulfilled
- $I_{sat,i}$ at fixed ϕ_{pr} : useful as monitor signal

Specific features – to keep in mind

- ▶ invasive method → probe size versus plasma volume
 - ▶ level of noise → EEPF for typically three orders of magnitude
average, smoothing, filtering
 - ▶ RF field → oscillating ϕ_{pl}
measured curve \neq real curve
- ↓
- RF compensation
active or passive
- ▶ magnetic field → gyro motion
use I_{ion} instead of I_e
 - ▶ negative ions → I_- instead of I_e
for same mass $I_{sat,ion} = I_{sat}$, symmetric curves

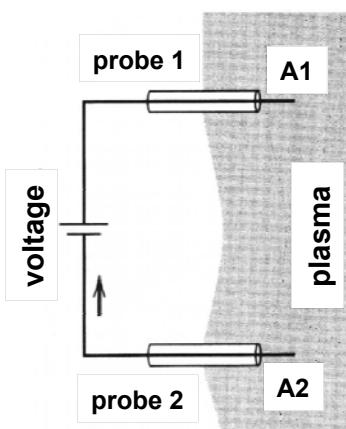


For monitoring or quantification with spatial and temporal resolution !

Ursel Fantz, p. 13

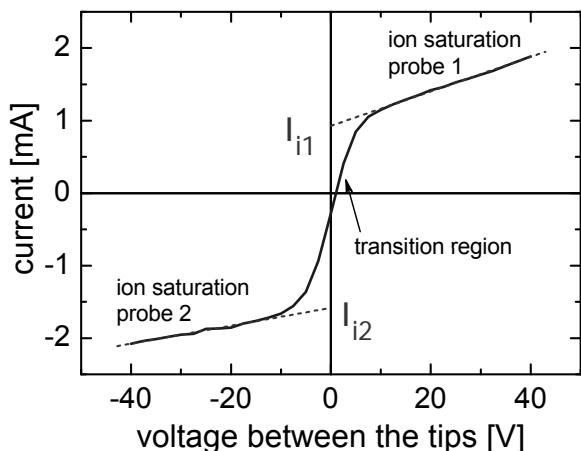
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Double probe



- ▶ two probe tips of same size
distance > 2 Debye length
- ▶ voltage between the probes
both probes are floating
no reference potential needed

compatible with
quartz or ceramic
chamber
and RF field



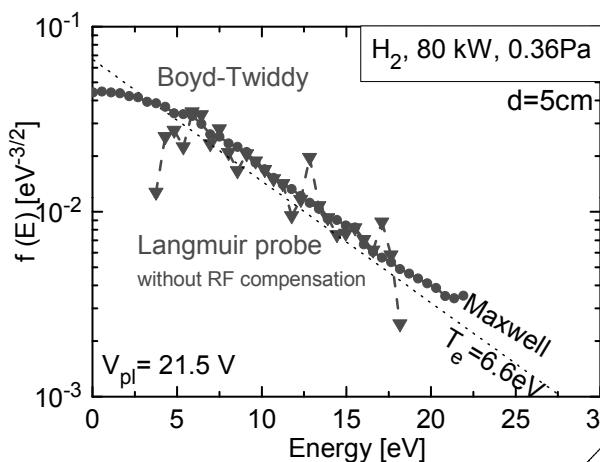
- ▶ symmetric curve: ion saturation current

$$I_{i,sat} = \frac{I_{isat,+} + I_{isat,-}}{2} \quad n_i = \frac{I_{i,sat}}{eA \sqrt{\frac{k_B T_e}{m_i}}}$$

- ▶ T_e from transition region

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Comparison of EEDF**Boyd-Twiddy method:****direct measurement of EEDF**

- voltage ramp is superimposed by AC modulated signal
- measure frequency spectrum of probe current

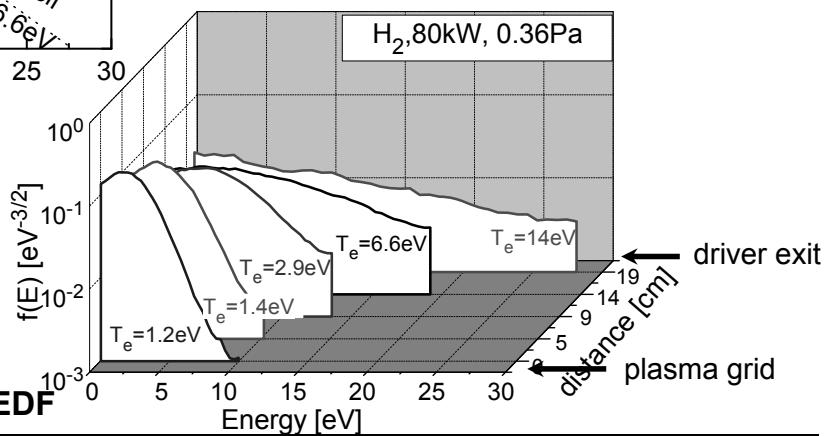
B. Crowley, S. Dietrich PSST 18 (2009) 014010

**Sophisticated
Langmuir probe system**

for RF ion sources

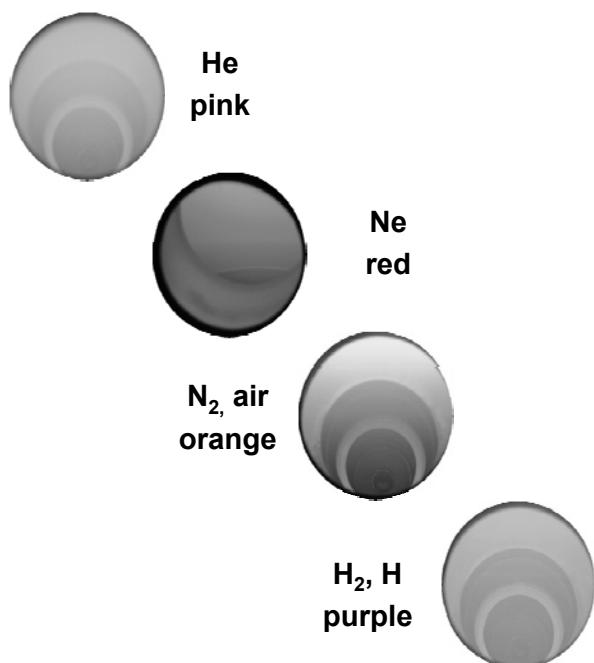
McNeely et al.

PSST 18 (2009) 014011

Maxwellian EEDF

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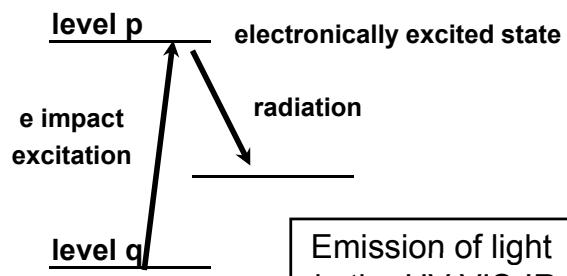
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Radiation of low temperature plasmas**Colourful plasmas !**

- **Neutrals** atoms and molecules
- **Ions** single charged
- **Electrons** $n_e \ll n_n$

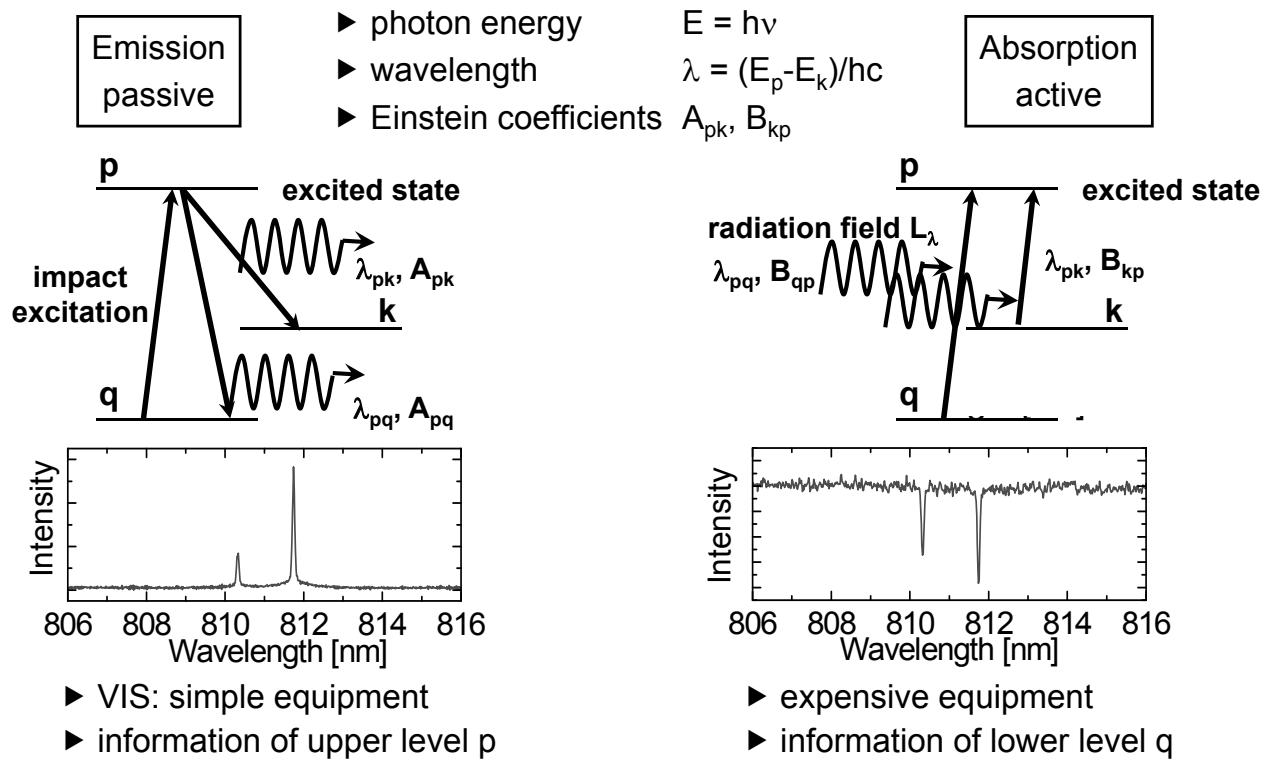


collisions and spontaneous emission



Emission of light
in the UV-VIS-IR

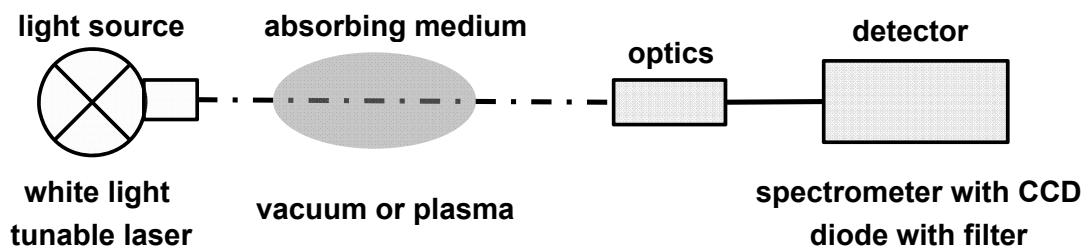
Emission versus absorption spectroscopy



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Main principle of line absorption



Non-invasive and line of sight integrated method !

Absorption in a medium
with path length l

$$I(\lambda, l) = I(\lambda, 0) \exp[-\kappa(\lambda)l]$$

absorption coefficient $\kappa(\lambda)$
statistical weights g_i, g_k

$$k(\lambda) = \frac{1}{l} \ln \left(\frac{I(\lambda, 0)}{I(\lambda, l)} \right)$$

with Ladenburg relation $\int_{\text{line}} \kappa_{ki}(\lambda) d\lambda = n_k \frac{g_i}{g_k} \frac{\lambda_0^4}{c} \frac{A_{ik}}{8\pi}$

Density of lower state
ground state, metastables

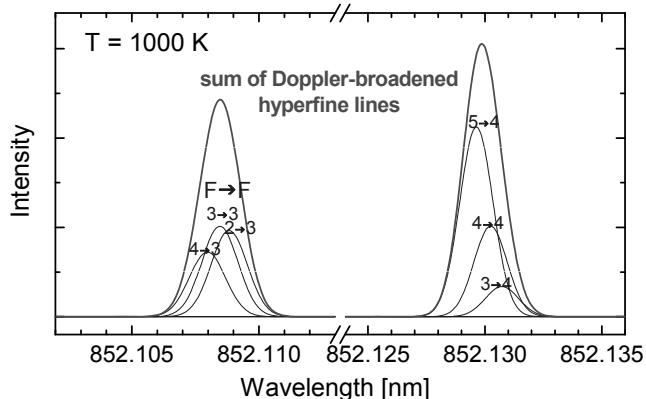
$$n_k = \frac{8\pi c}{\lambda_0^4} \frac{g_k}{g_i} \frac{1}{A_{ik} l} \int_{\text{line}} \ln \left(\frac{I(\lambda, l)}{I(\lambda, 0)} \right) d\lambda$$

Example: Cs line at 852.1 nm**resonance line $6^2\text{S}_{1/2} - 6^2\text{P}_{3/2}$ with hyperfine structure**

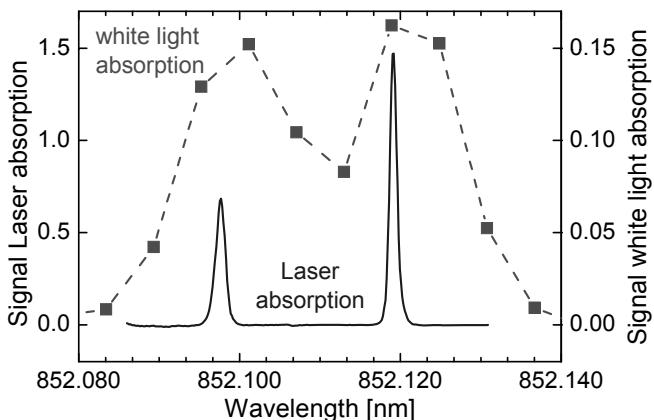
take line broadening into account

Doppler profile

$$\Delta\lambda_{D,FWHM} = \frac{\lambda}{c} \sqrt{\frac{8k_B T \ln 2}{m}}$$



appearance depends on technique

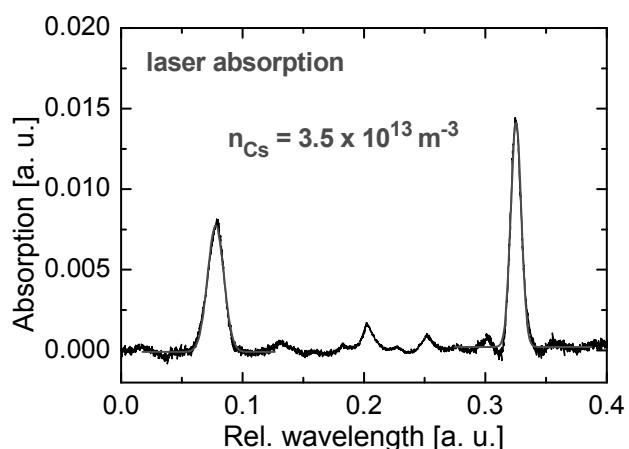
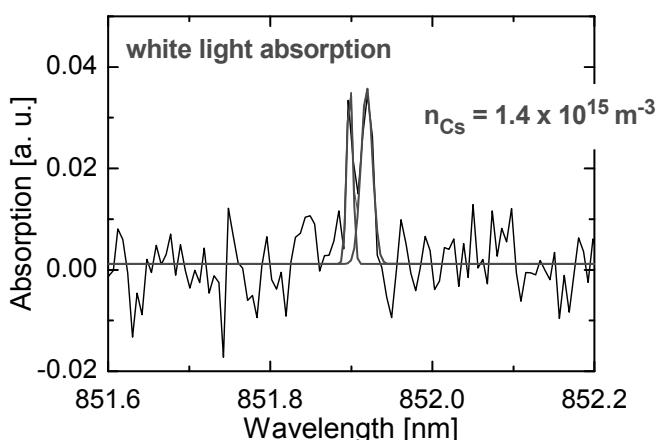


white light absorption ↔ laser absorption

apparatus profile (spectrometer) Doppler profile

six hyperfine lines overlap to two peaks: $\Delta\lambda \approx 21 \text{ pm}$ **Example: Cs line at 852.1 nm**U. Fantz, C. Wimmer
J. Phys. D 44 (2011) 335202

white light absorption versus laser absorption

**Improved detection limit for laser absorption: factor > 40 !**sensitivity range: $3 \times 10^{13} \text{ m}^{-3} - 10^{17} \text{ m}^{-3}$ (path length = 15 cm)
being perfectly in the range required for the ion sources

Straightforward analysis

- ▶ in vacuum
- ▶ with plasma → subtract emission

$$n_k = \frac{8\pi c}{\lambda_0^4} \frac{g_k}{g_i} \frac{1}{A_{ik} l_{\text{line}}} \int \ln \left(\frac{I(\lambda, l)}{I(\lambda, 0)} \right) d\lambda$$

... but

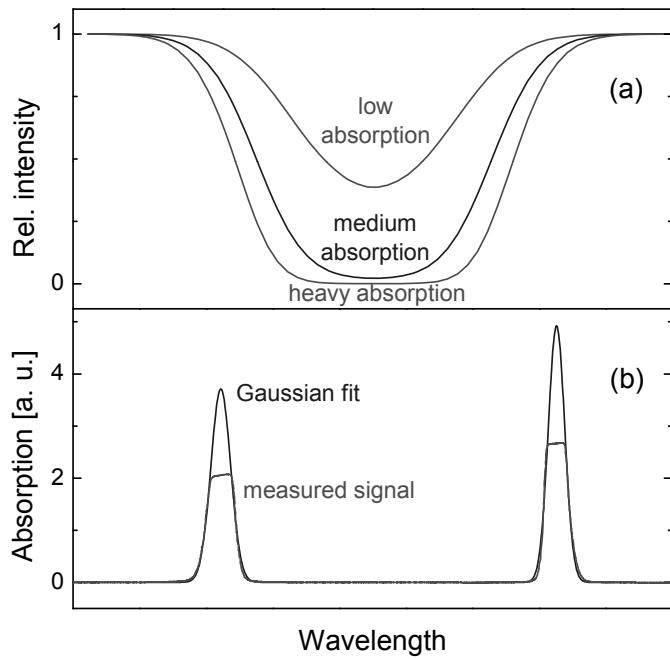
Line saturation

- ▶ strong absorption: $n_k \times l$
- ▶ correction factors by profile fitting

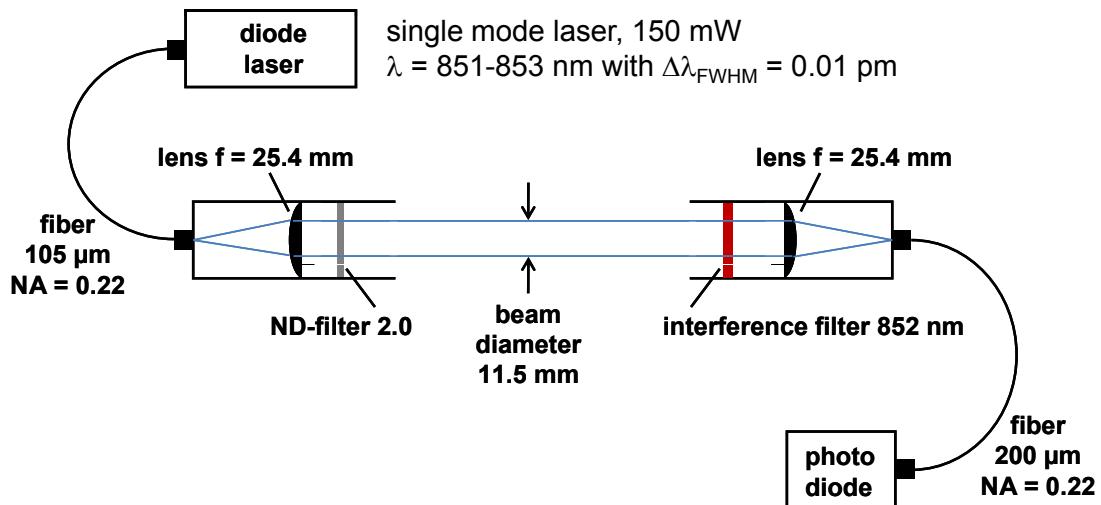
... and

Depopulation effect

- ▶ strong intensity
- ▶ attenuation of laser to $\approx 1\%$ → trade-off with temporal resolution



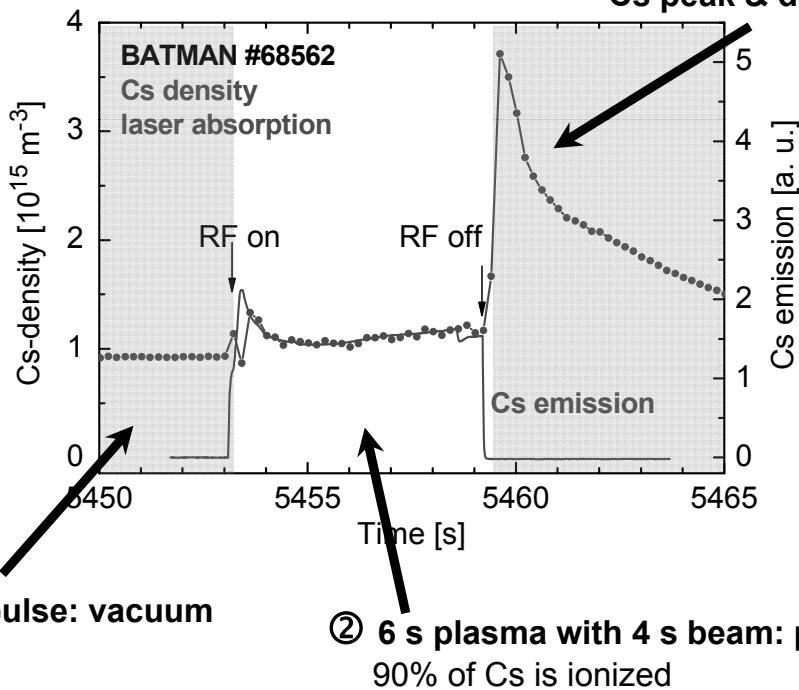
Tunable diode laser – Fibre optics – Photo diode with interference filter



Simple and robust setup for application to ion sources !

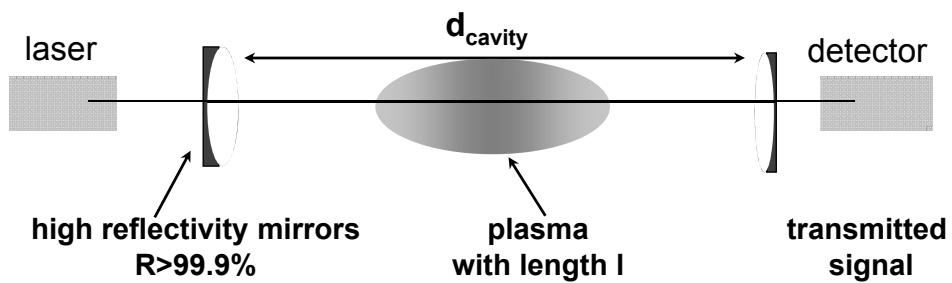
On-line monitoring: vacuum – plasma (6s with 4s beam) - vacuum

③ 30 s after pulse:
Cs peak & decay time



Cavity – Ringdown – Absorption – Spectroscopy → CRDS

pumping of an optical cavity by a (tunable) laser source
 measurement of signal decay after laser source is switched off



Measurement of laser light attenuation trapped in a high-finesse optical cavity

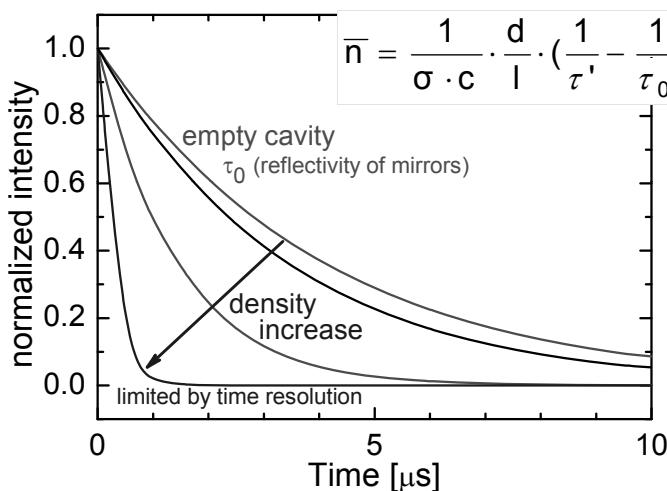
transfers absorption signal from wavelength into time dependence

empty cavity with decay time τ_0

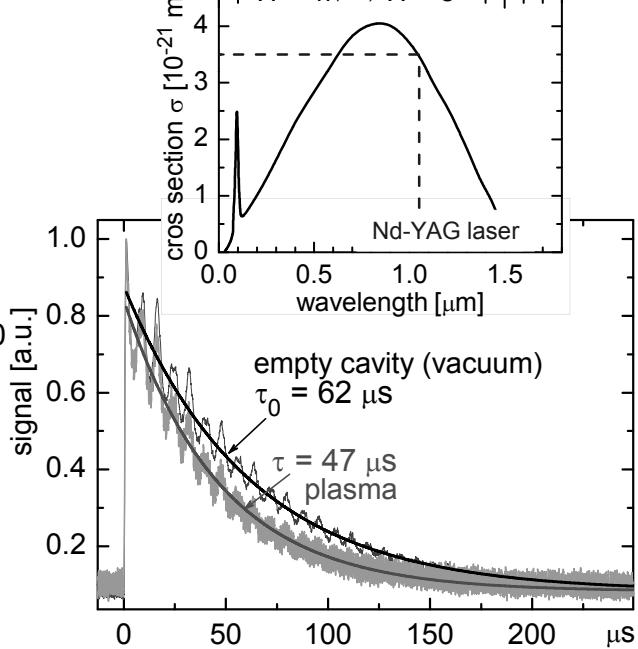
$$I(t) = I_0 \cdot e^{-\frac{t}{\tau_0}} ; \quad \tau_0 = \frac{d}{c(1-R)}$$

additional absorption $\tau_0 \rightarrow \tau'$

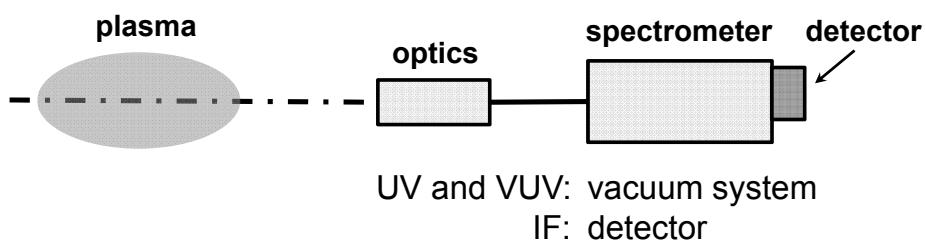
$$\bar{n} = \frac{1}{\sigma \cdot c} \cdot \frac{d}{l} \cdot \left(\frac{1}{\tau'} - \frac{1}{\tau_0} \right)$$

Cavity – Ringdown – Absorption – Spectroscopy → CRDS

Example H^-
cross section: photodetachment
 $\text{H}^- + h\nu \rightarrow \text{H} + e$

 **H^- density: line of sight averaged**

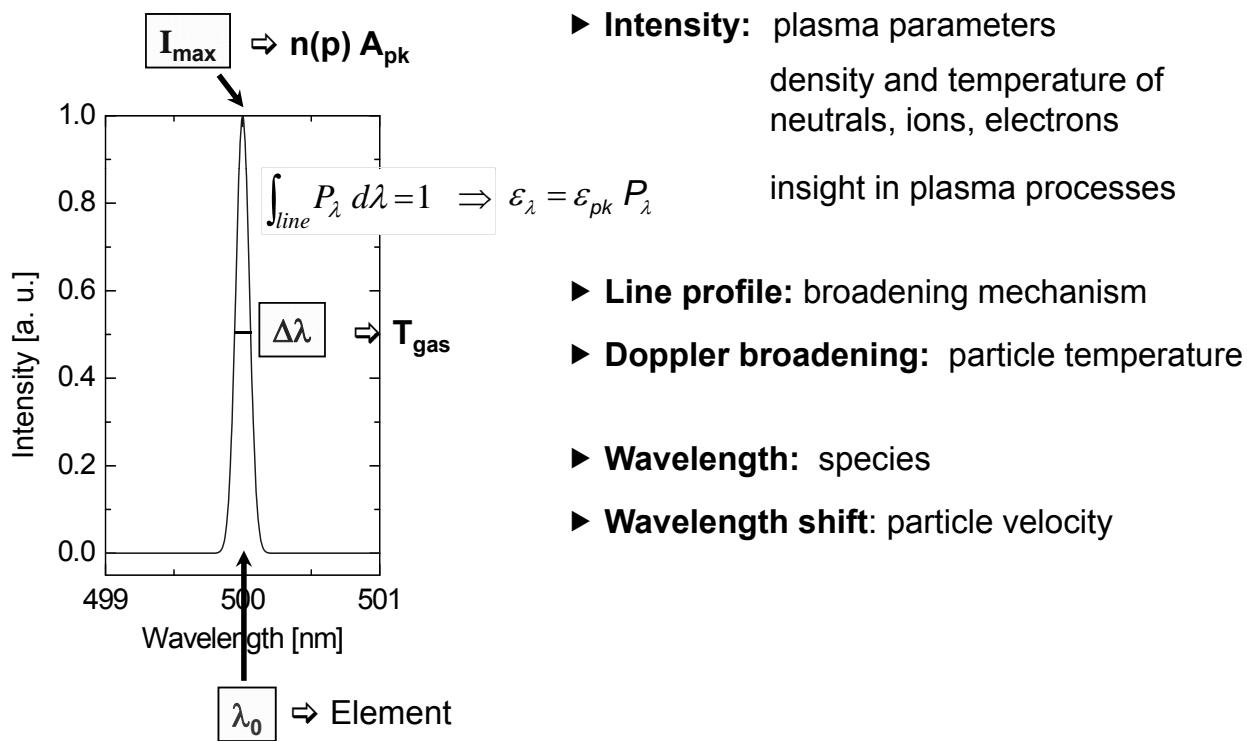
- ▶ detection limit $\approx 10^{15} \text{ m}^{-3}$
- ▶ $\text{H}^- = 5 \times 10^{17} \text{ m}^{-3} \rightarrow t = 8 \mu\text{s}$

The main principle**Non-invasive and line of sight integrated method !****Measures density of excited state ...**

$$\varepsilon_{pk} = n(p) A_{pk} \frac{hc/\lambda}{4\pi}$$

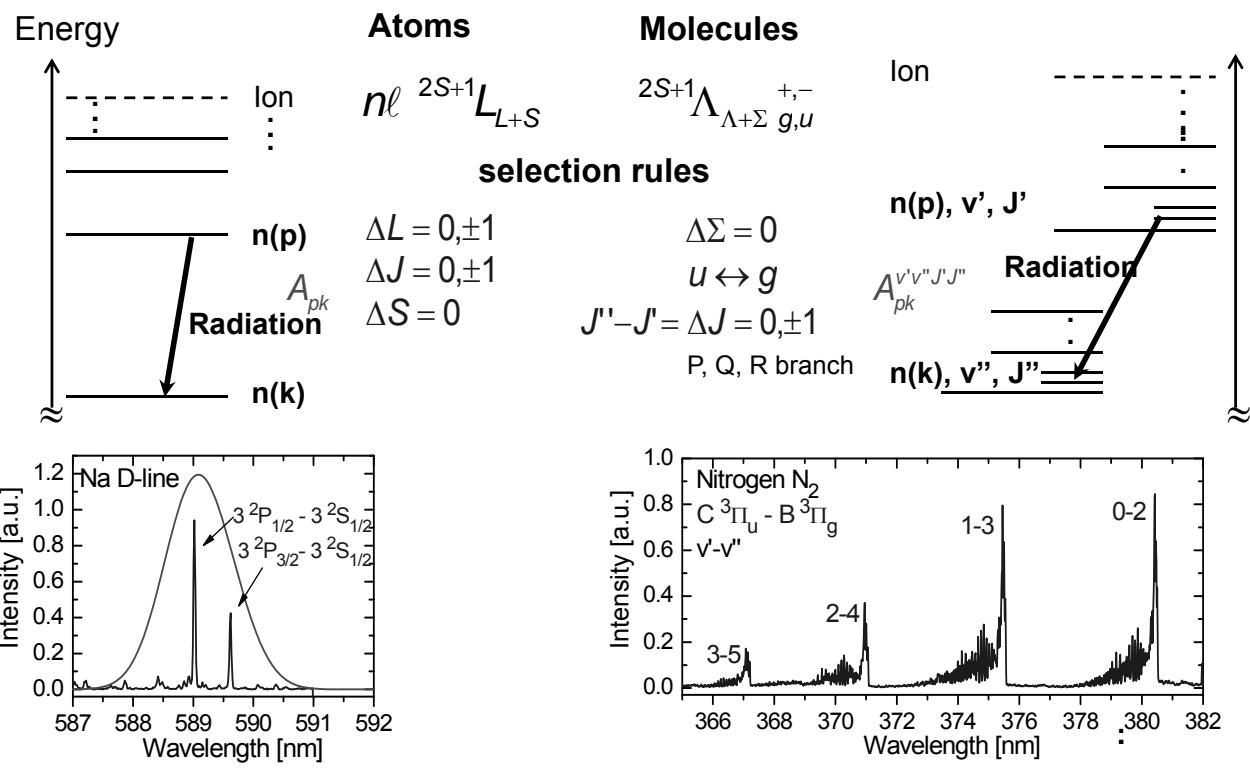
... which depends on plasma parameters !

What information can be obtained from the line emission ?



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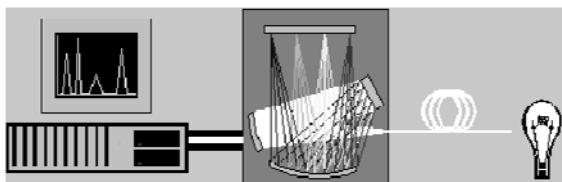


Appearance depends on spectral resolution !

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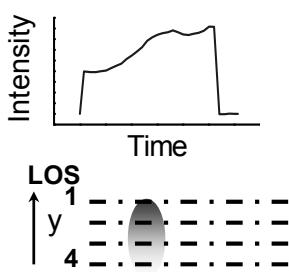
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Spectroscopic system

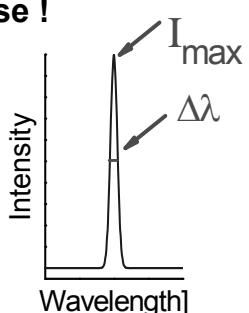


Detector	Spectrometer	Optics
► photomultiplier λ scan $\Delta\lambda, \Delta t$	► focus length spectral resolution $\Delta\lambda$	► fibre very flexible VIS: glass, quartz, UV enhanced
► diode array λ range	► grating spectral resolution Blaze - intensity	► lens and aperture Imaging optics solid angle
► CCD, ICID pixel size - $\Delta\lambda$ intensity	► slits entrance slit $\Delta\lambda$ exit slit - detector	

The spectroscopic system is determined by the purpose !



- time resolution detector
- spatial resolution detector, lines of sight
- intensity detector, spec., optics
- spectral resolution detector, spec., optics



survey spectrometer	pocket size	$\Delta\lambda \approx 1-2 \text{ nm}$
1m spectrometer	good optics	$\Delta\lambda \approx 20 \text{ pm}$
Echelle spectrometer	high resolution	$\Delta\lambda \approx 1-2 \text{ pm}$

→ line profile
line shift

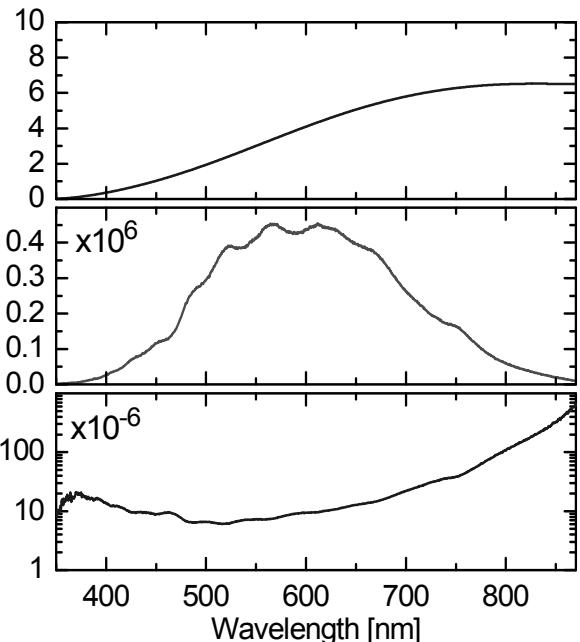
line monitoring
very simple
 Δt , poor $\Delta\lambda$
less information

relative intensities
common technique
poor Δt , $\Delta\lambda$, Δx , flexible
moderate information

absolute intensities
expensive technique
poor Δt , $\Delta\lambda$, Δx , flexible
powerful tool

Calibration of the spectroscopic system

Wavelength: pixel → nm
spectral lamps, plasma, λ tables



Radiance - Intensity
counts → $W/m^2/sr, ph/m^2/s$

Ulbricht sphere

Calibrated spectrum
spectral radiance
[$W/m^2/sr/nm$]



Measurement
intensity [counts/s]

Conversion factor
spectral sensitivity

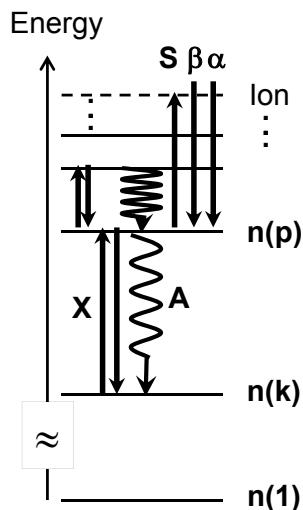
$$\left[\frac{W}{m^2 sr nm (counts/s)} \right] \times \frac{4\pi\lambda}{hc} = \left[\frac{\text{photons}}{m^2 snm (counts/s)} \right]$$

↑ exposure time

From intensity $I_{pk} = n(p) A_{pk}$ to plasma parameter

Population models

high n_e	Boltzmann distribution (TE, LTE)	↑ relevance of photons
↓	Collisional radiative model	
low n_e	Corona equilibrium	



Rate equations for excitation and de-excitation processes

- electron impact excitation and de-excitation
- absorption and emission, heavy particle collisions,

Corona model

$$\frac{dn(p)}{dt} = n_1 n_e X_{1p}(T_e) - n(p) \sum_k A_{pk} = 0$$

rate coefficient

CR model

$$I_{pk} = n_0 n_e X_{pk}(T_e) \quad \text{with} \quad X_{pk} = X_{1p}(T_e) A_{pk} / \sum_k A_{pk}$$

emission rate coefficient

$$I_{pk} = n_0 n_e X_{pk}^{eff}(T_e, n_e, \dots)$$

Electron temperature from absolute line emission

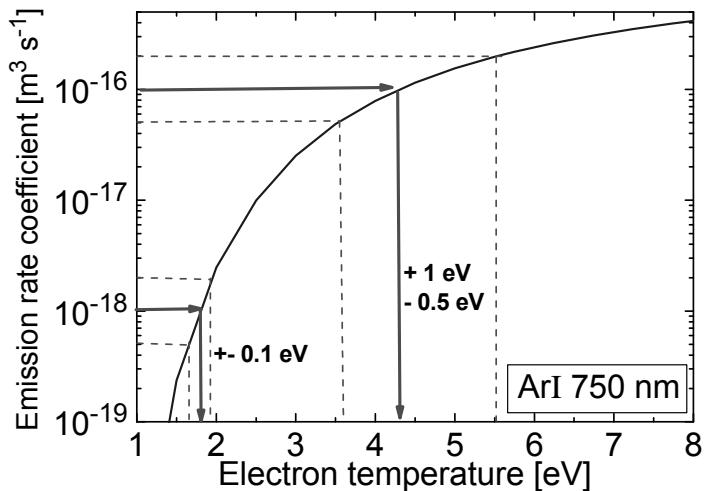
$$I_{pk} = n_0 n_e X_{pk}^{eff}(T_e, n_e, \dots)$$

$$n_0, n_e \text{ known} \Rightarrow X_{pk}^{eff}(T_e, n_e, \dots) = \frac{I_{pk}}{n_0 n_e}$$

Find suitable gases and diagnostic lines

- ▶ admixture of small amount of diagnostic gas
- ▶ prominent example: Ar

Very sensitive for low T_e !

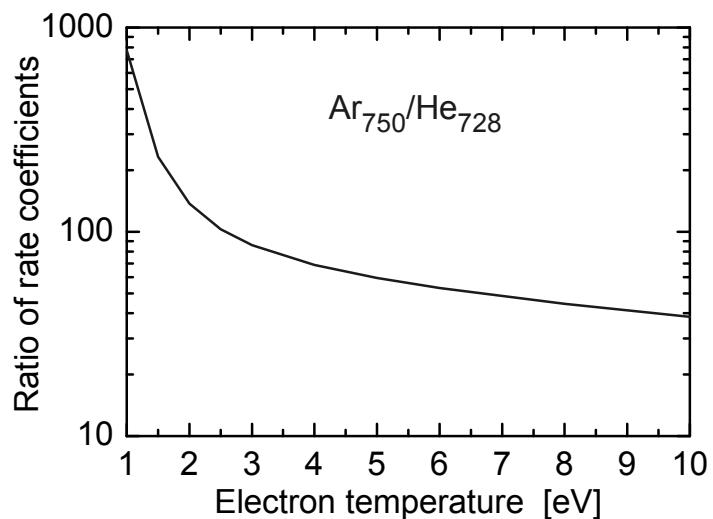


Electron temperature from line ratio (relative calibration)

$$\frac{I_{pk}^1}{I_{pk}^2} = \frac{n_1 \gamma_e^1 X_{pk}^1(T_e)}{n_2 \gamma_e^2 X_{pk}^2(T_e)} \rightarrow \text{ratio of rate coefficients for known densities}$$

Find suitable gases and diagnostic lines

- ▶ n_1, n_2 inert gases or $n_1 = n_2$
- ▶ I_{pk} undisturbed lines
- ▶ ground state excitation
- ▶ X_{pk} ratio depends on T_e



U. Fantz et al., Nucl. Fusion 49 (2009) 125007

Actinometry: density ratio from line ratio (relative calibration)

$$\frac{I_{pk}^1}{I_{pk}^2} = \frac{n_1 \gamma_e X_{pk}^1(T_e)}{n_2 \gamma_e X_{pk}^2(T_e)} \rightarrow \text{ratio of densities for known rate coefficients}$$

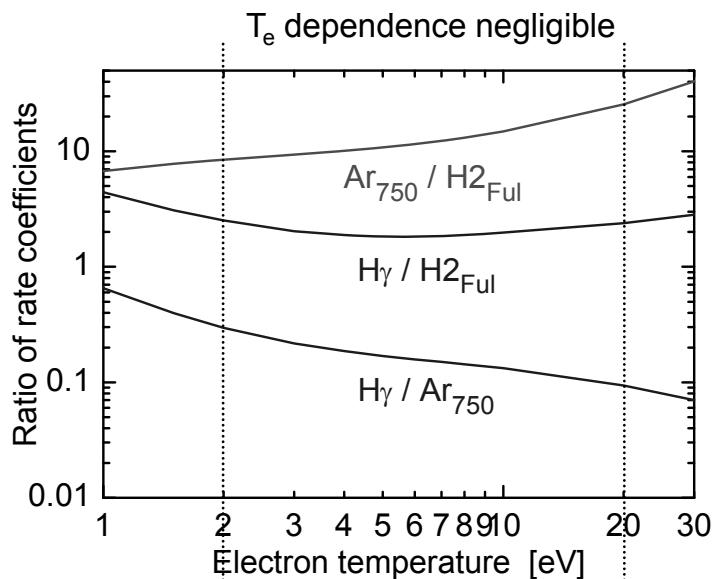
- density ratio (n_H/n_{H_2})

$$\frac{I_{434}^H}{I_{H_2}^{Ful}} = \frac{n_H \gamma_e X_{434}^H(T_e)}{n_{H_2} \gamma_e X_{Ful}^{H_2}(T_e)} \text{ independent on } n_e$$

- density ratio (n_{Ar}/n_{H_2})

$$\frac{I_{750}^{Ar}}{I_{H_2}^{Ful}} = \frac{n_{Ar} \gamma_e X_{750}^{Ar}(T_e)}{n_{H_2} \gamma_e X_{Ful}^{H_2}(T_e)}$$

- density ratio (n_{He}/n_{H_2})
dependence on T_e (factor of 10)
→ needs iteration



Ursel Fantz, p. 35

CAS on Ion Sources, 6th June 2012

Particle density from absolute line emission

$$I_{pk} = n_0 n_e X_{pk}^{eff}(T_e, n_e, \dots)$$

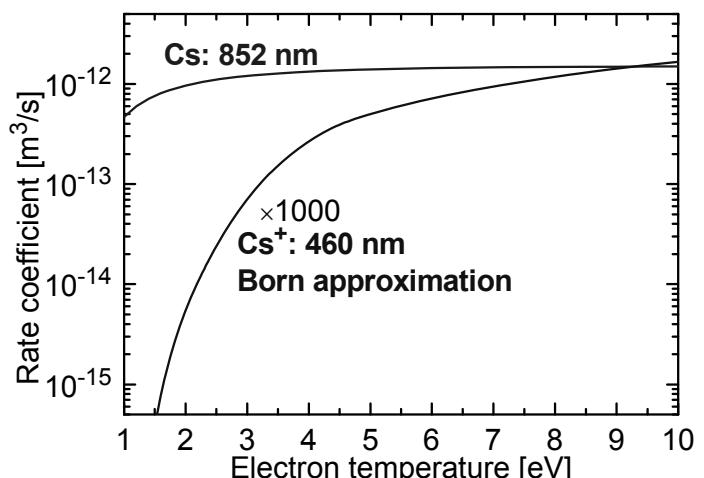
$$n_e, T_e \text{ known} \Rightarrow n_0 = \frac{I_{pk}}{n_e X_{pk}^{eff}(T_e, n_e, \dots)}$$

Knowledge of dominant excitation mechanism is essential !**Example: Cs and Cs⁺ lines**

$$\text{Cs: } I_{852}^{Cs} = n_{Cs} n_e X_{852}^{Cs}(T_e)$$

needs n_e , almost independent of T_e

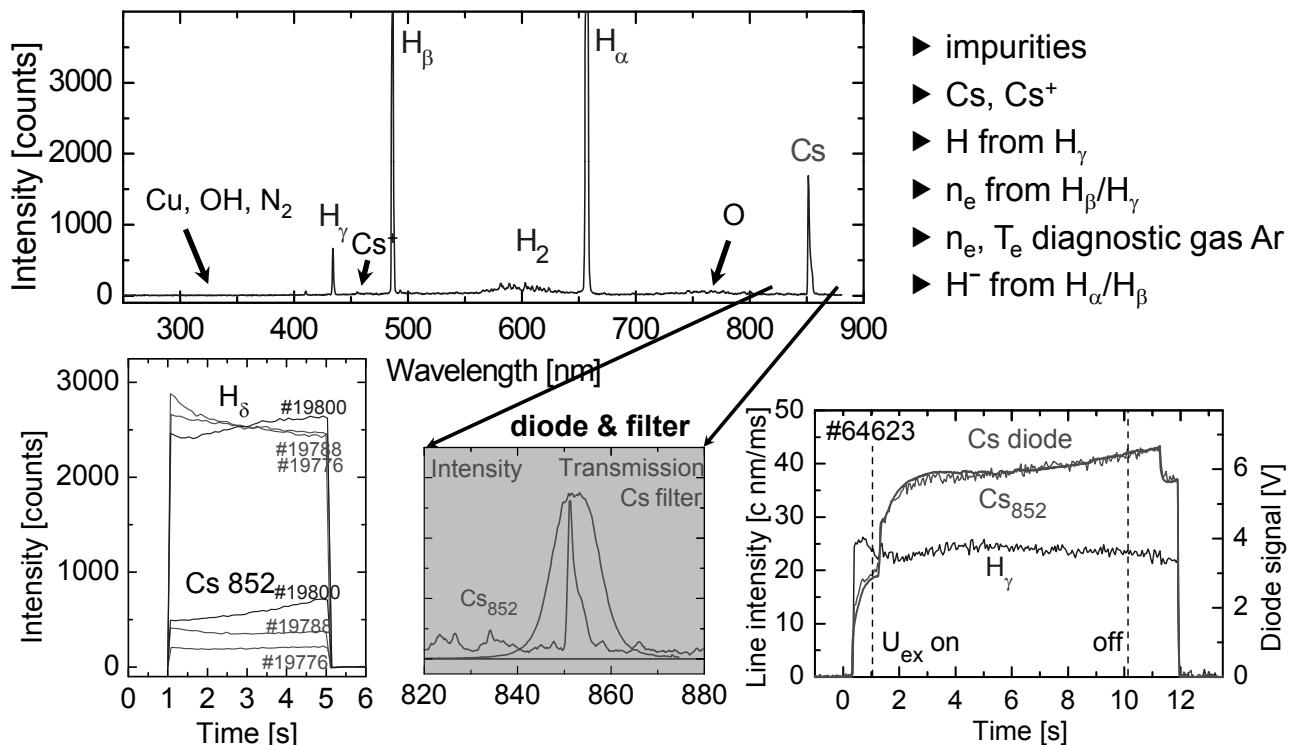
$$\text{Cs}^+: I_{460}^{Cs^+} = n_{Cs^+} n_e X_{460}^{Cs^+}(T_e)$$

needs n_e , strong dependence on T_e 

Ursel Fantz, p. 36

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Survey spectrometer and on-line monitoring



Ursel Fantz, p. 37

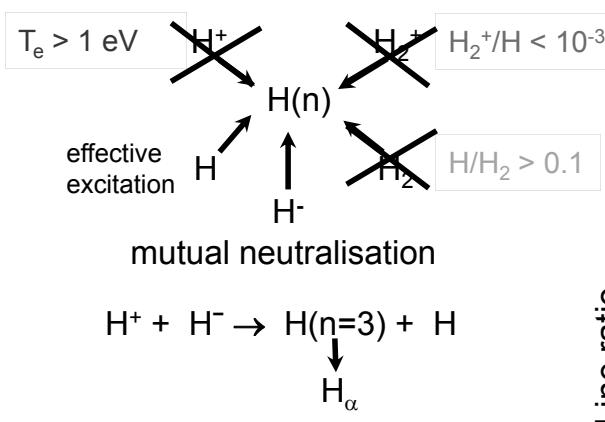
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U. Fantz, D. Wunderlich

A novel diagnostic technique for H⁻ volume density

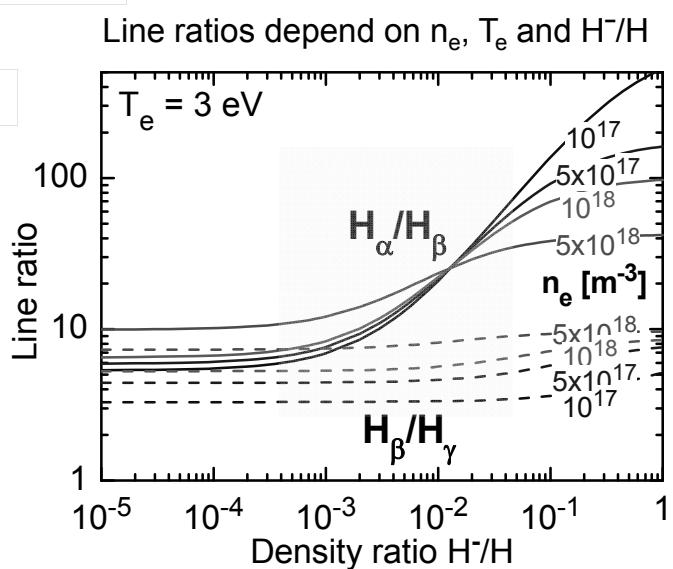
NJP 8 (2006) 301

Population mechanisms for H



Collisional radiative model

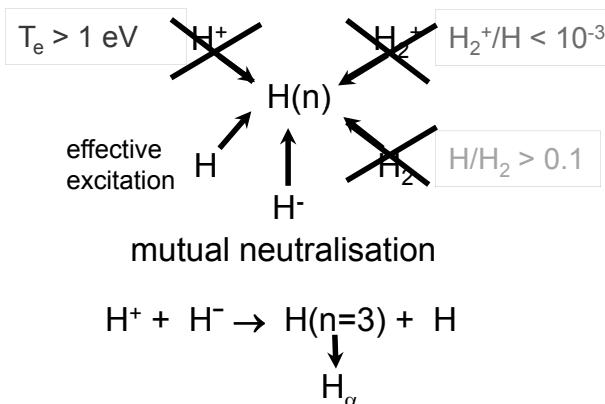
- Measurement of Balmer line ratios
- H_α/H_β depends on H⁻/H
 - H_β/H_γ reflects n_e and T_e



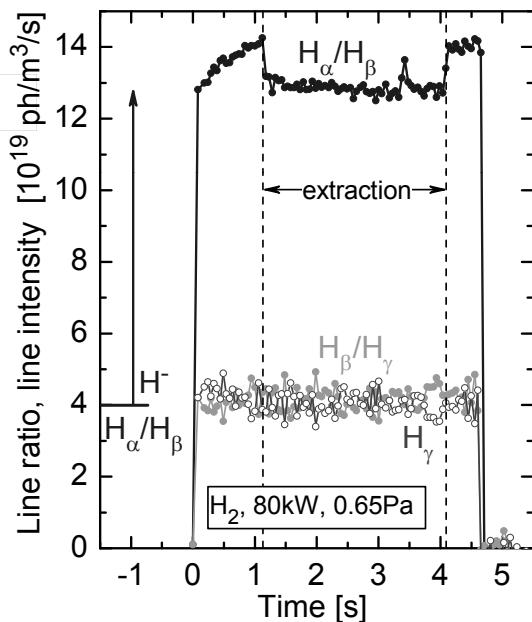
U. Fantz, D. Wunderlich

A novel diagnostic technique for H^- volume density

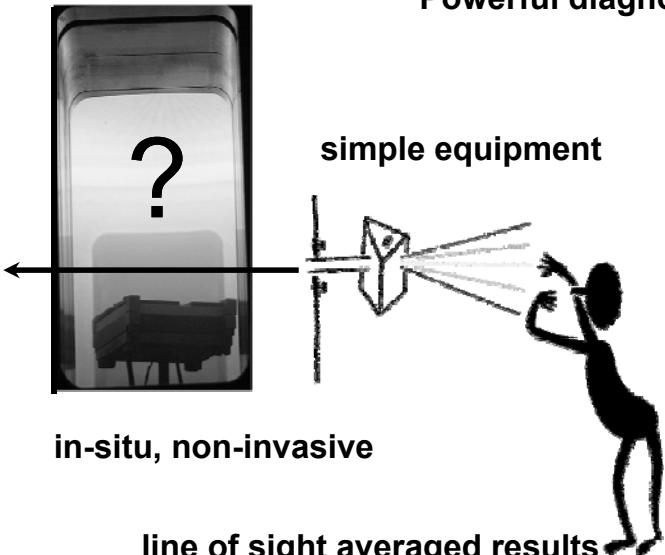
NJP 8 (2006) 301

Population mechanisms for H**Collisional radiative model****Measurement of Balmer line ratios**

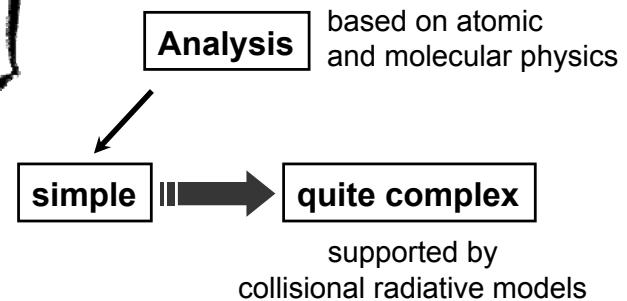
- H_α/H_β depends on H^-/H
- H_β/H_γ reflects n_e and T_e



Ursel Fantz, p. 39

CAS on Ion Sources, 6th June 2012**Powerful diagnostic tool**

- identification of species
- particle densities
- particle temperatures
- on-line monitoring
- insight in plasma processes
- spatial resolution by several lines of sight



Ursel Fantz, p. 40

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Plasma Diagnostics of Ion Sources

The three W's

- ▶ What do I want to know ? → and why?
- ▶ What is the adequate technique ? → effort versus gain!
- ▶ What is the accessibility of the source ? → feasibility !

The three examples

- ▶ Langmuir probes → ϕ_{pl} , n_e , T_e , (EEDF)
- ▶ Absorption techniques → $n_{species}$ → Cs, H^-
- ▶ Emission spectroscopy → n_s , $T_s \rightarrow e, H, H_2, H^-$

The three “keep-in-mind’s”

- ▶ Monitoring versus quantification → trends or full information
- ▶ Spatial resolution → averaged or x-resolved (step width!))
- ▶ Temporal resolution → averaged or t-resolved (time scale!)

Diagnostics – The Window to the Knowledge !

- [1] F. F. Chen, J.P Chang, *Lecture Notes on Principles of Plasma Processing* (Kluwer/Plenum, 2003)
- [2] M. Lieberman, A. Lichtenberg, *Principles of Plasma Discharges and Materials Processing* (Wiley, 1994)
- [3] B. Chapman, *Glow Discharge Processes* (Wiley, 1986)
- [4] R. Hippler, S. Pfau, *Low Temperature Plasma Physics* (Wiley, 2001)
- [5] D. Flamm, O. Auciello: *Plasma Diagnostics*, Volume 1 (Academic Press, Inc., 1989)
- [6] I.H. Hutchinson: *Principles of Plasma Diagnostics* (Cambridge University Press, 1987)
- [7] A. Thorne, *Spectrophysics: Principles and Applications* (Springer 1999)
- [8] <http://www.ee.ucla.edu/~ffchen/>
- [9] U. Fantz: *Basics of plasma spectroscopy*
Plasma Sources Sci. Technol. 15 (2006), 137
- [10] U. Fantz: *Emission Spectroscopy of Molecular Low Pressure Plasmas*
Contrib. Plasma Phys. 44 (2004) 508