

Accuracy complements energy

Electroweak precision tests at Tera-Z

Based on VM, Stefanek, You, 2412.14241



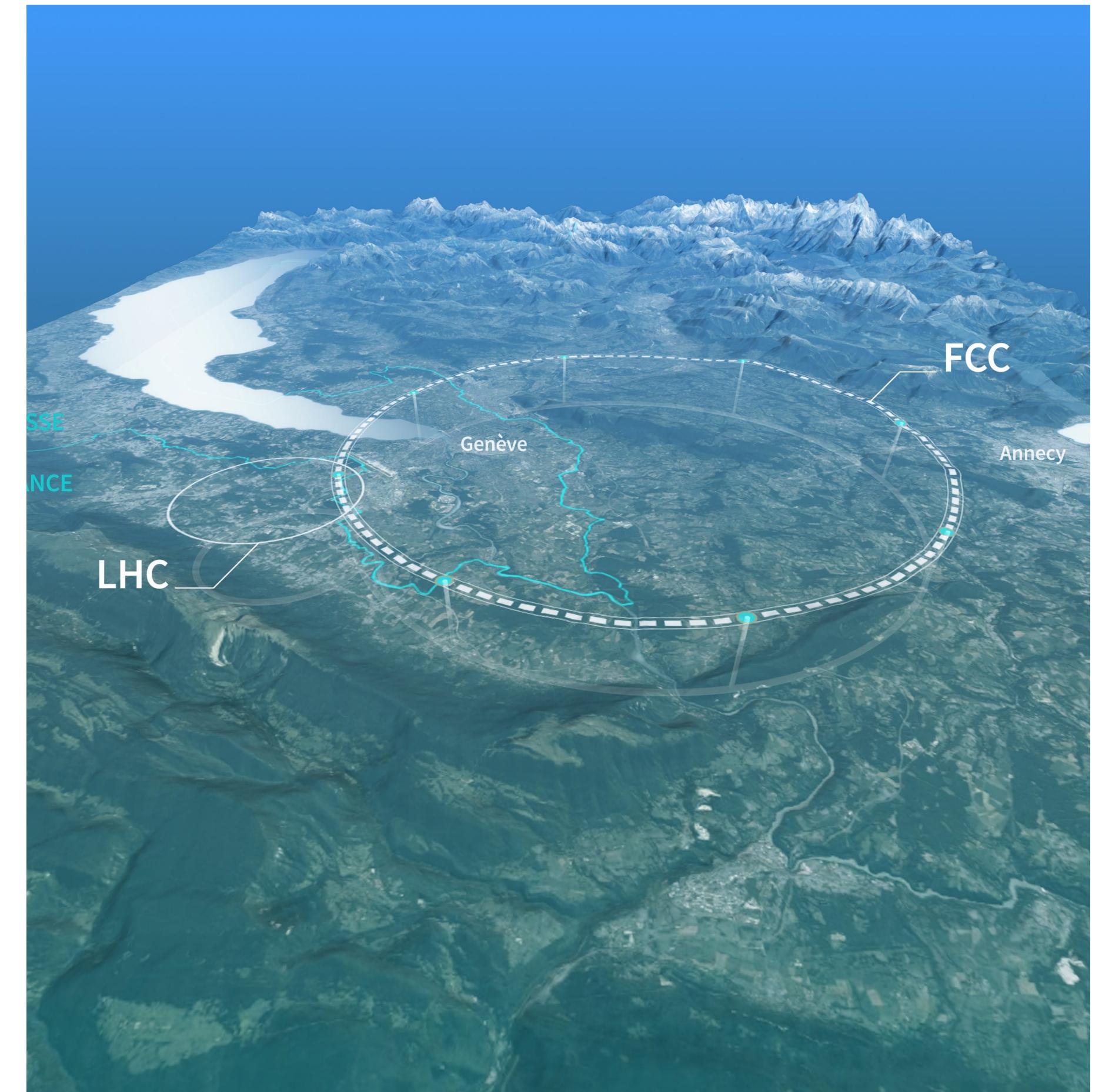
Victor Maura Breick, 19.02.2025



FCC Early-Career Forum, 19.02.2025

Overview

1. A theorists' view of FCCee
2. Standard Model Effective Field Theory (SMEFT)
3. Accuracy complements energy (ACE)
4. Implications



* We have 10+10 minutes, so if you have questions, please interrupt me and I will try to address them. Otherwise, save them for the end and I can go more in depth!

A theorist's view of FCCee

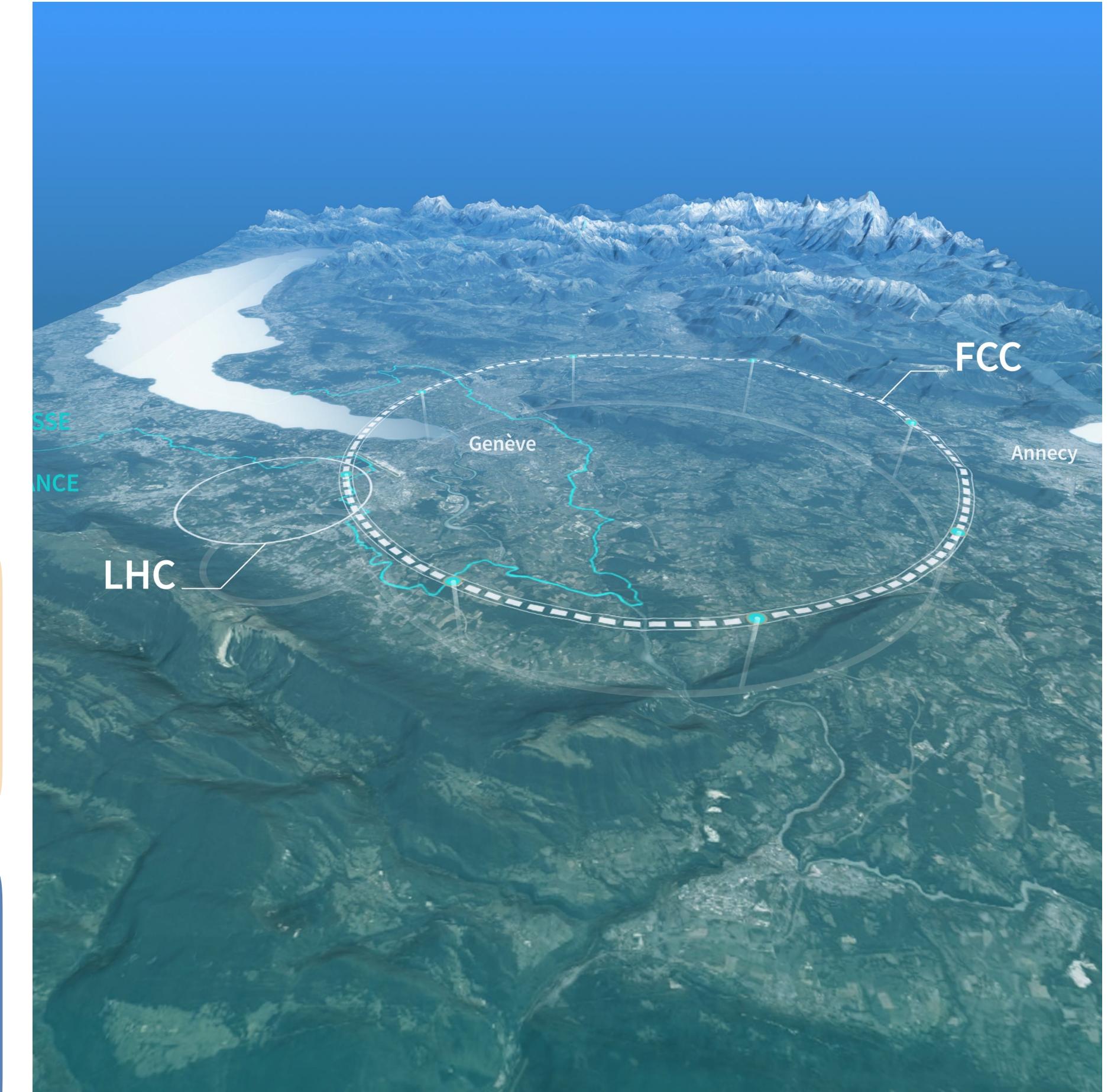
	Z – pole	WW	Zh	t <bar>t</bar>
Energy [GeV]	91.2	163	240	365
Accuracy	$6 \cdot 10^{12} Z$	$2.4 \cdot 10^8 WW$	$2.2 \cdot 10^6 h$	$2 \cdot 10^6 t\bar{t}$

On Pole/Z-pole:

$$O_{\text{on-pole}} = \left\{ \Gamma_Z, \sigma_{\text{had}}, R_l, A_{\text{FB}}^{0,l}, R_b, R_c, A_b^{\text{FB}}, A_c^{\text{FB}}, A_l, A_b, A_c, A_s, m_W, \Gamma_W \right\}$$

Above/off-pole:

$$O_{\text{off-pole}} = \left\{ \sigma(e^+e^- \rightarrow W^+W^-), \sigma(e^+e^- \rightarrow ZH), \sigma(e^+e^- \rightarrow f\bar{f}), \mu_{b\bar{b}}, \mu_{c\bar{c}}, \mu_{\tau\bar{\tau}}, \mu_{\mu\bar{\mu}} \right\}$$



Standard Model Effective Field Theory

Effective Field Theory:

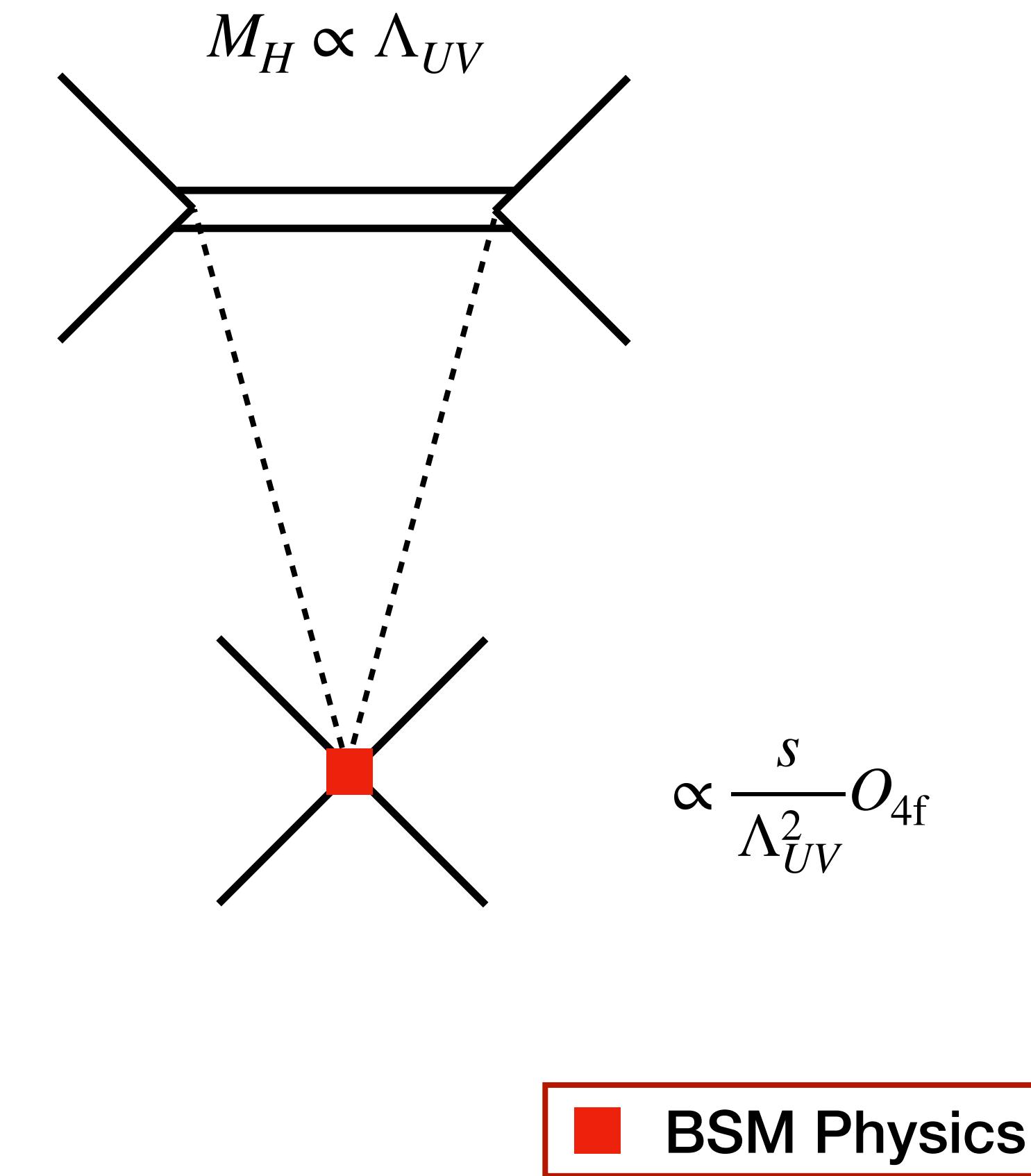
- Non-renormalisable QFT with **clear separation between UV and IR modes** and a **power counting** parametrised by $\alpha \ll 1$
- Allows to separate calculation of observables (LO,NLO...) in the IR and BSM in the UV
- Model independent global fits

SMEFT:

$$\mathcal{L} = \mathcal{L}_{SM} + \sum_{i \in S_2} \left(\frac{E_{CM}}{\Lambda_{UV}} \right)^2 C_i(\mu) O_i + \mathcal{O}(\Lambda^{-3})$$

Parametrise New Physics!
Scale Dependent Couplings! ⚡

IR physics, always the same



Increasing the Energy:

1. New Physics effects are **energy enhanced**

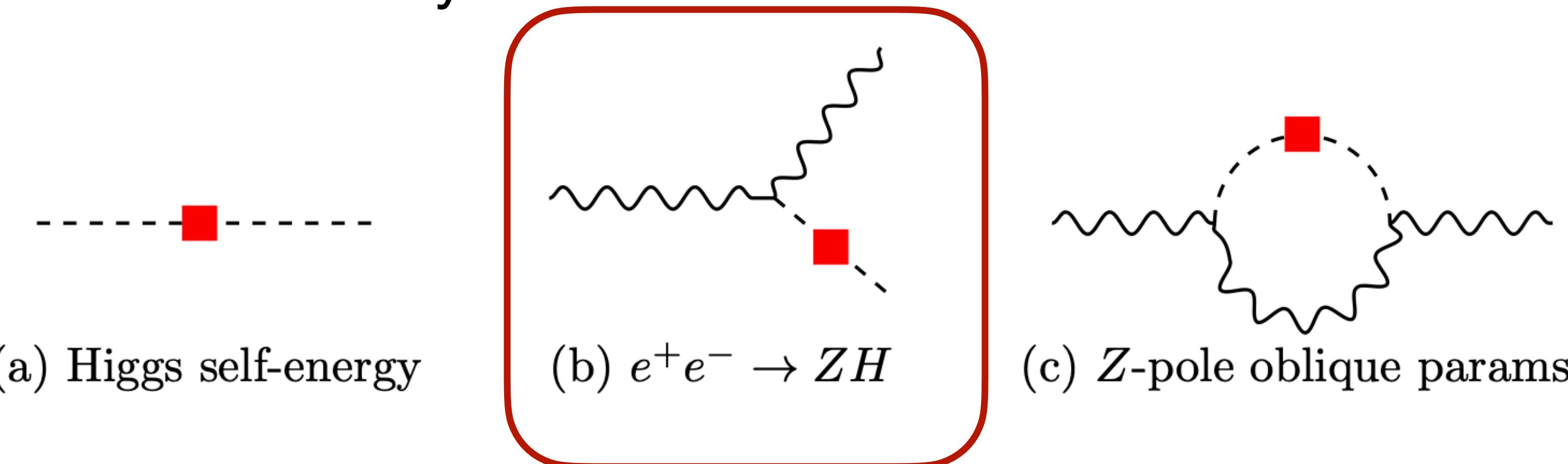
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2. **Direct** access to New Physics

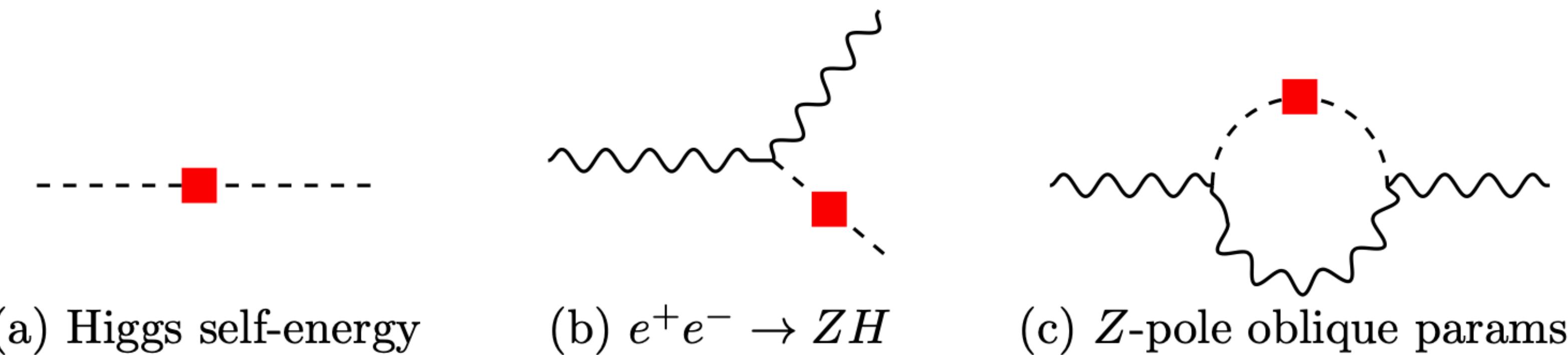


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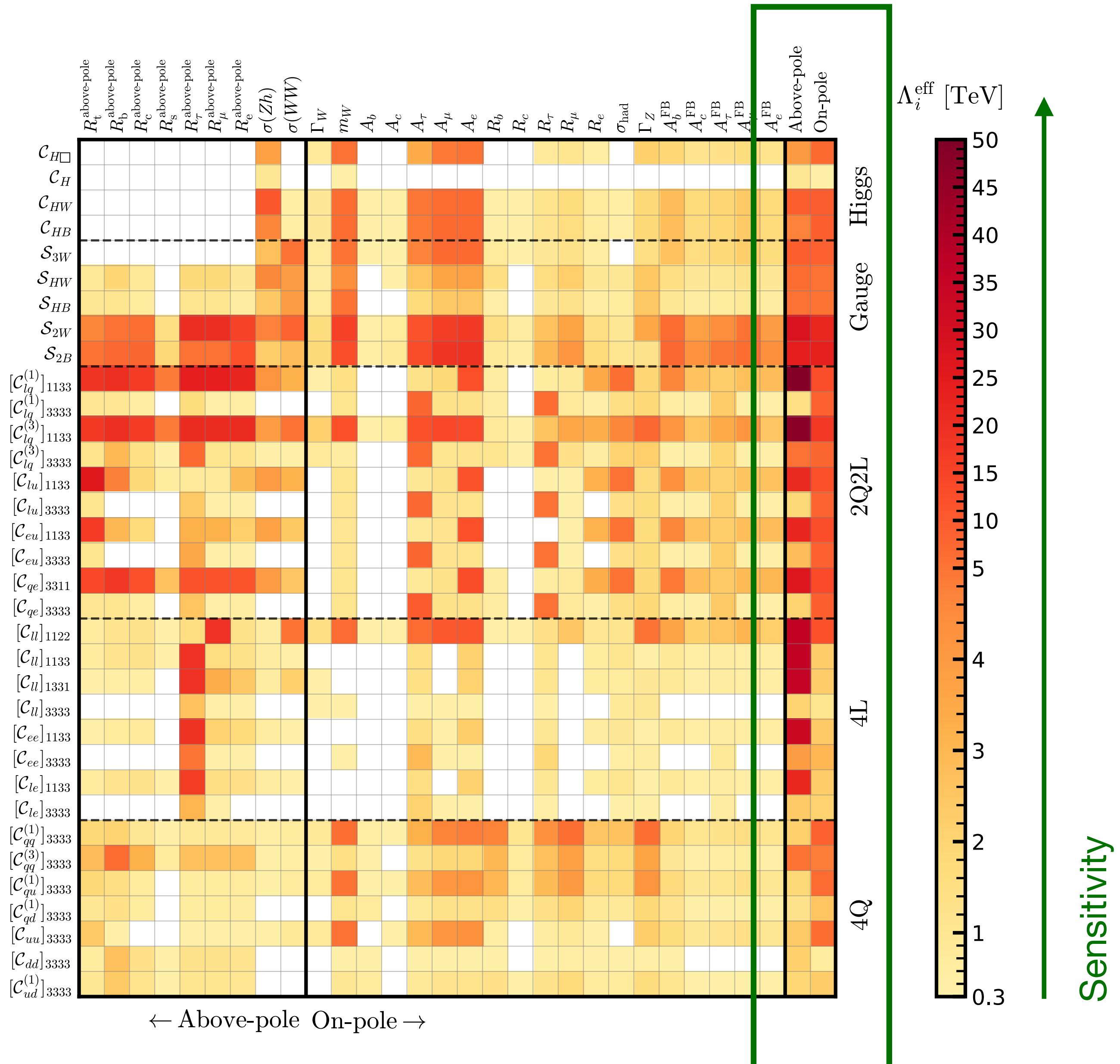
2. **Direct** access to New Physics



Can the increased precision of the FCCee Z-pole run compensate the relative suppression?

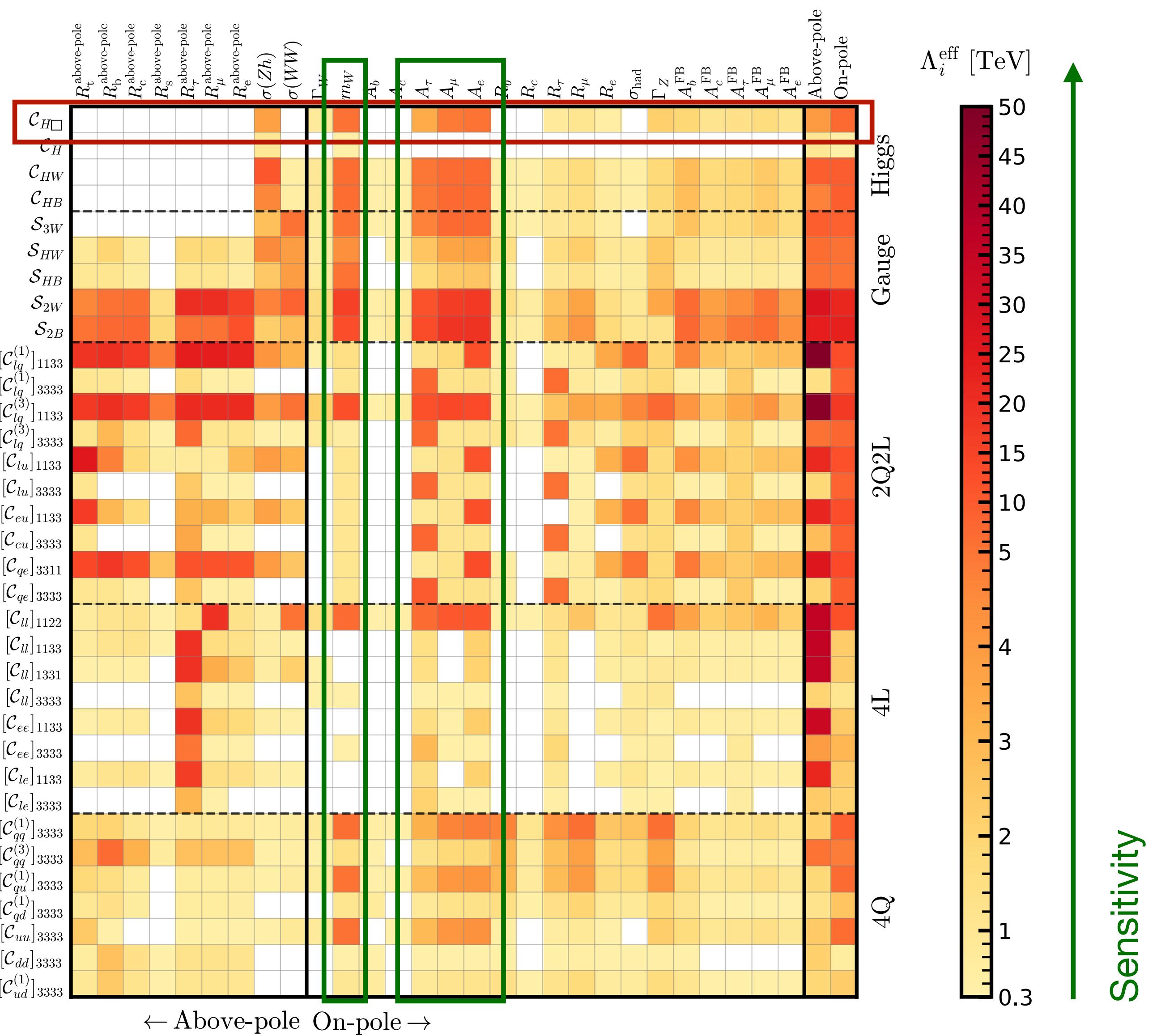
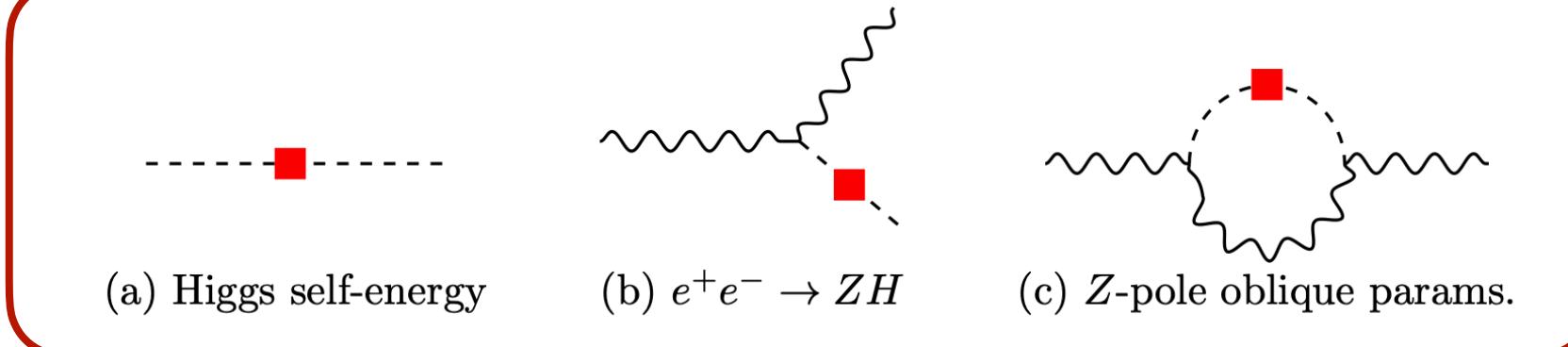
Accuracy Complements Energy

Yes!



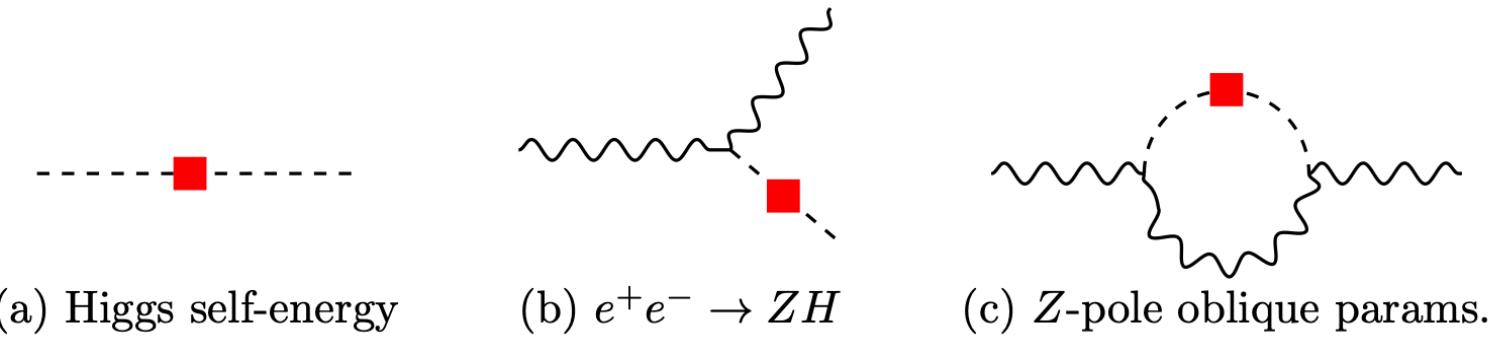
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Higgs 2 and 3 pt functions

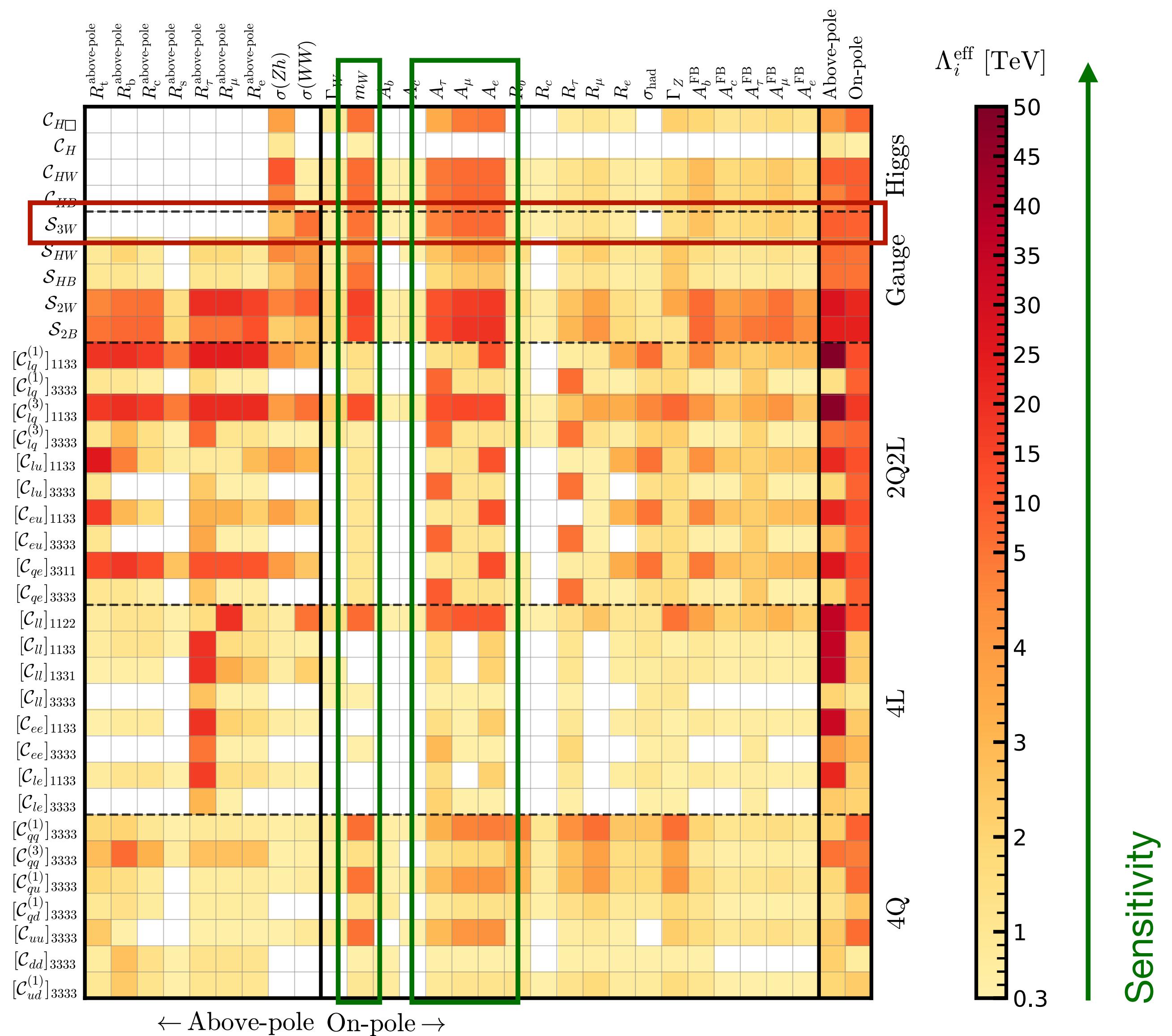
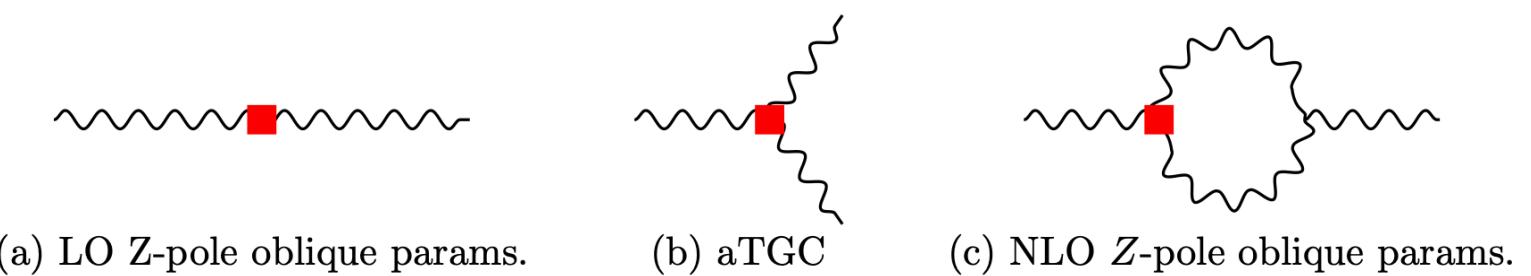


Accuracy Complements Energy

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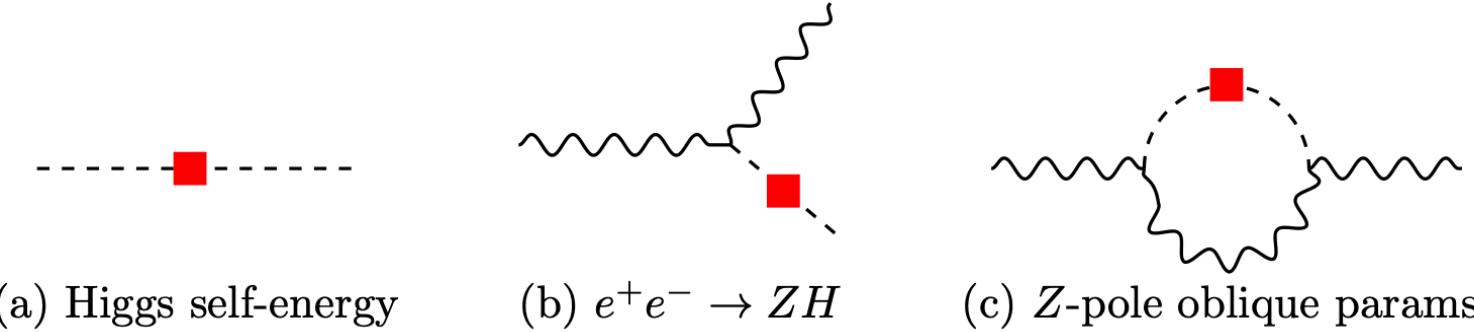


Gauge 2 and 3 pt functions

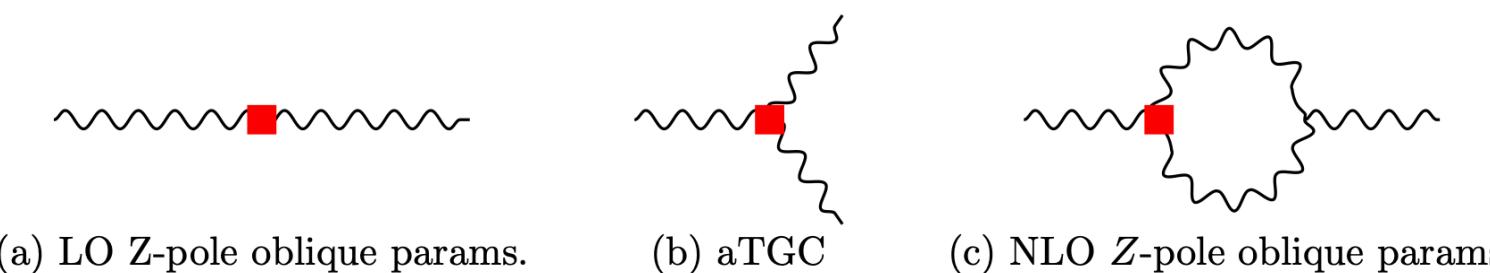


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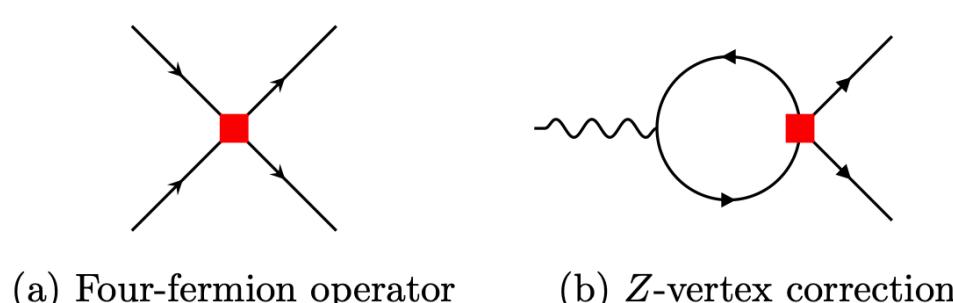
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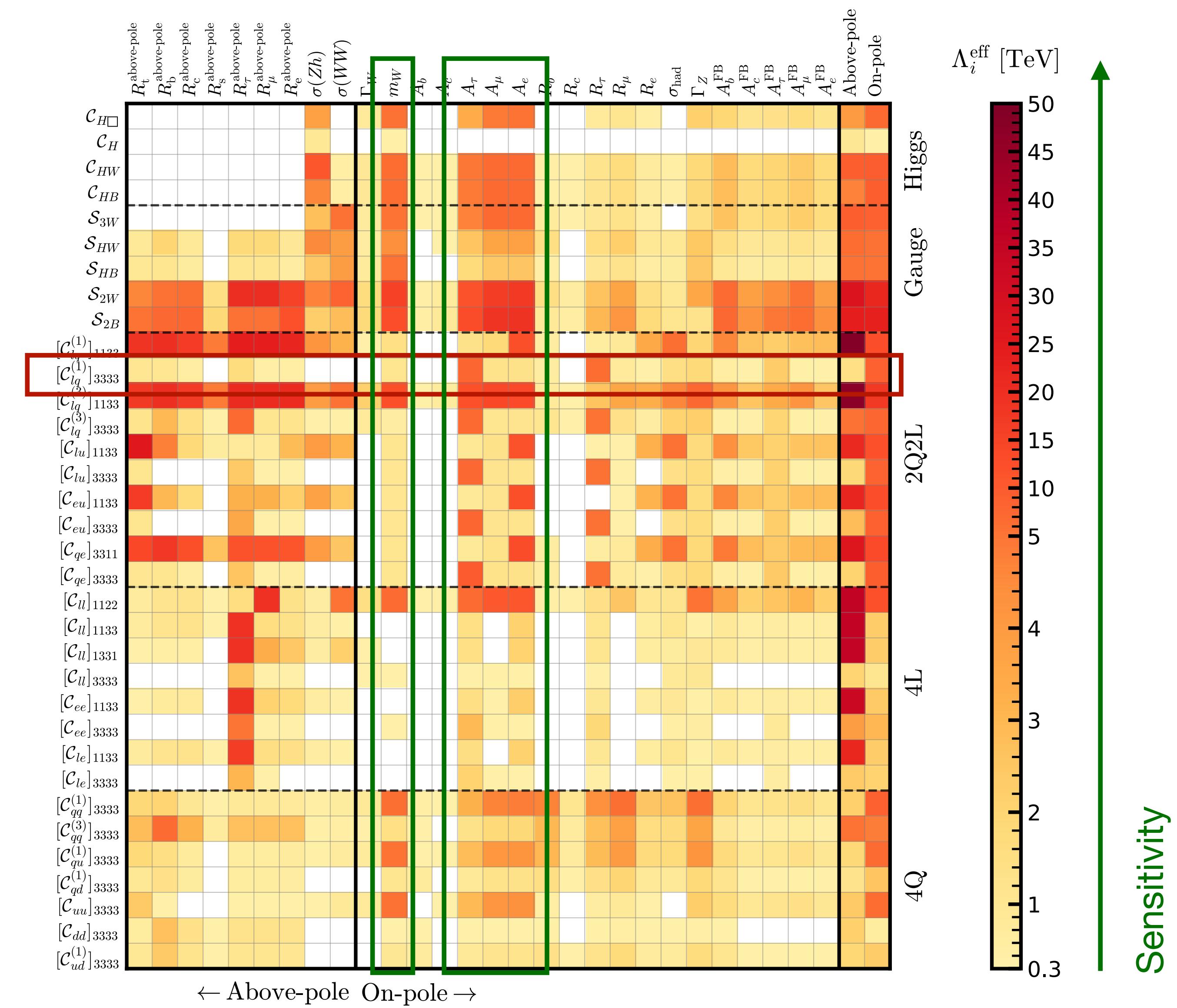
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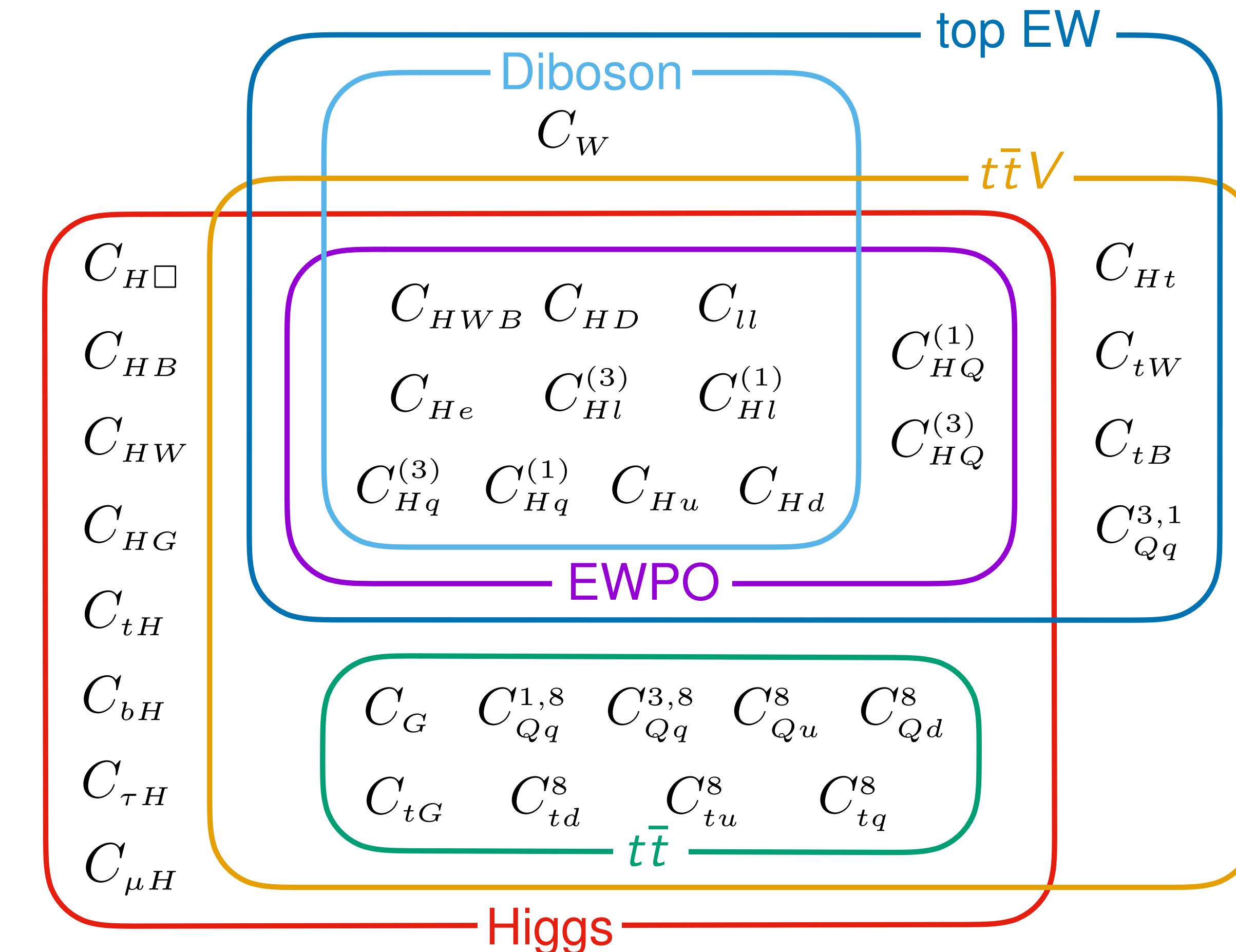
4 Fermion Operators



Best studied e.g. 2311.00020, 2407.09593,
2304.00029, 2411.02485...

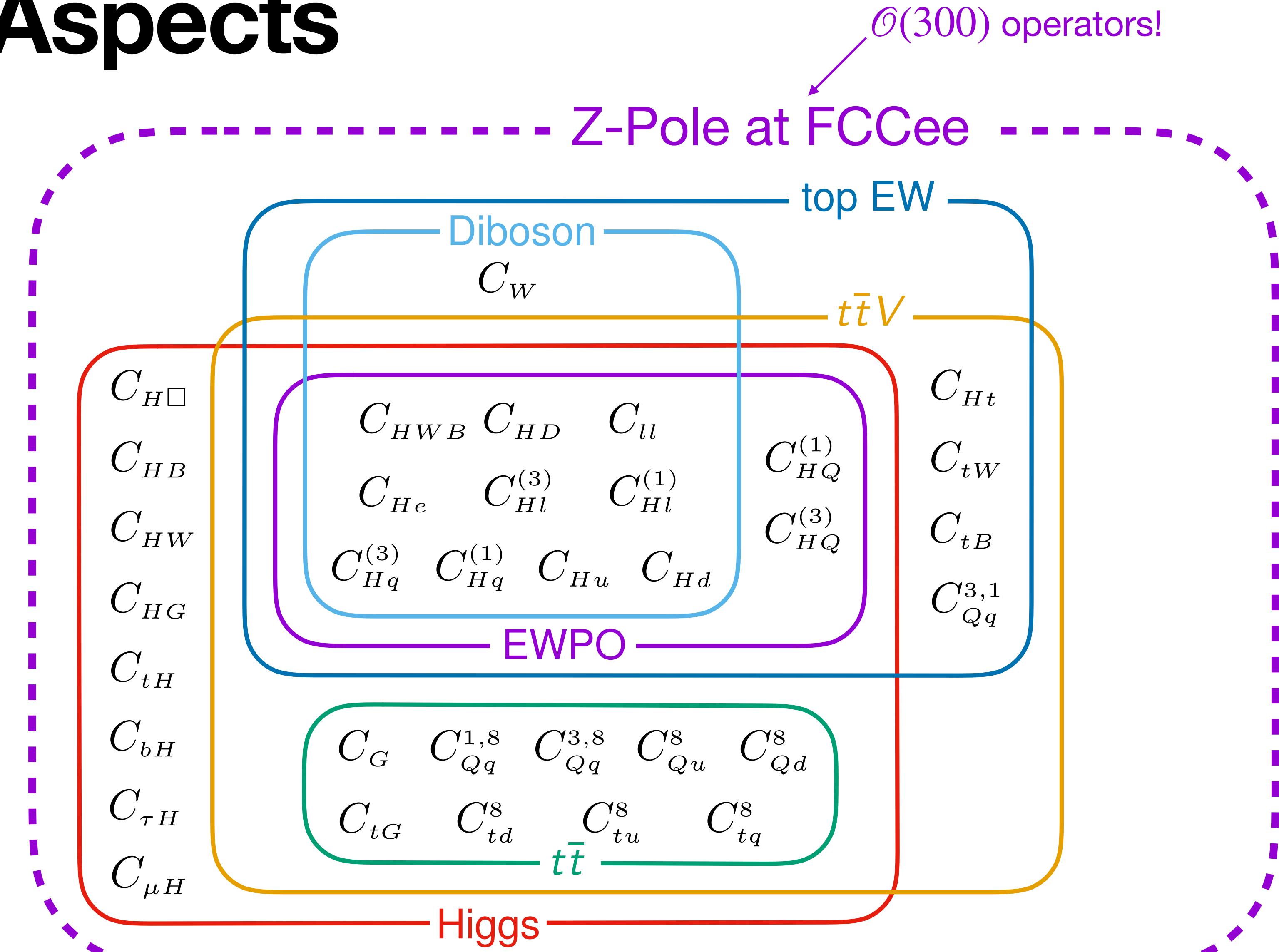


ACE: Negative Aspects



*Adapted from 2012.02779

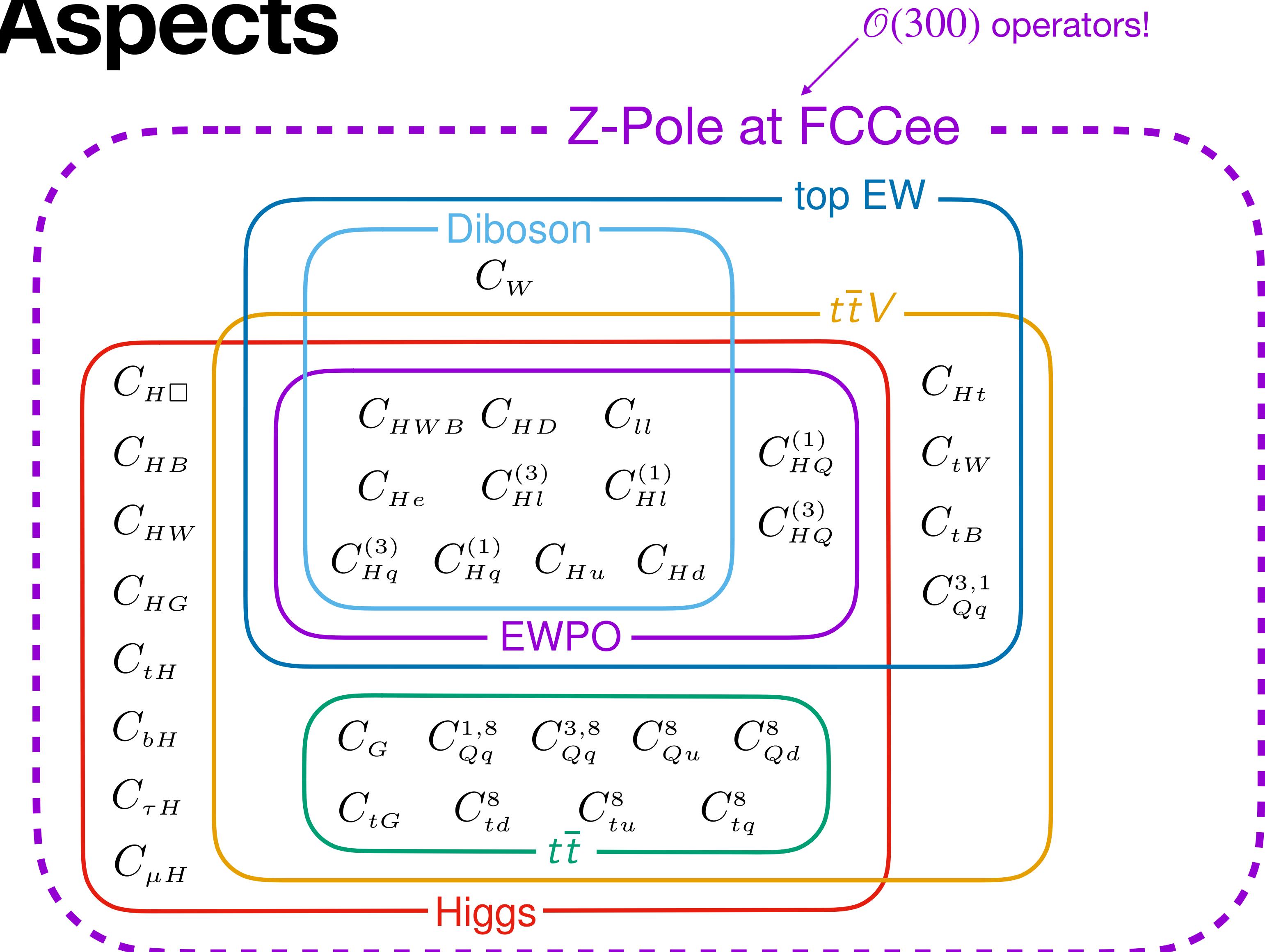
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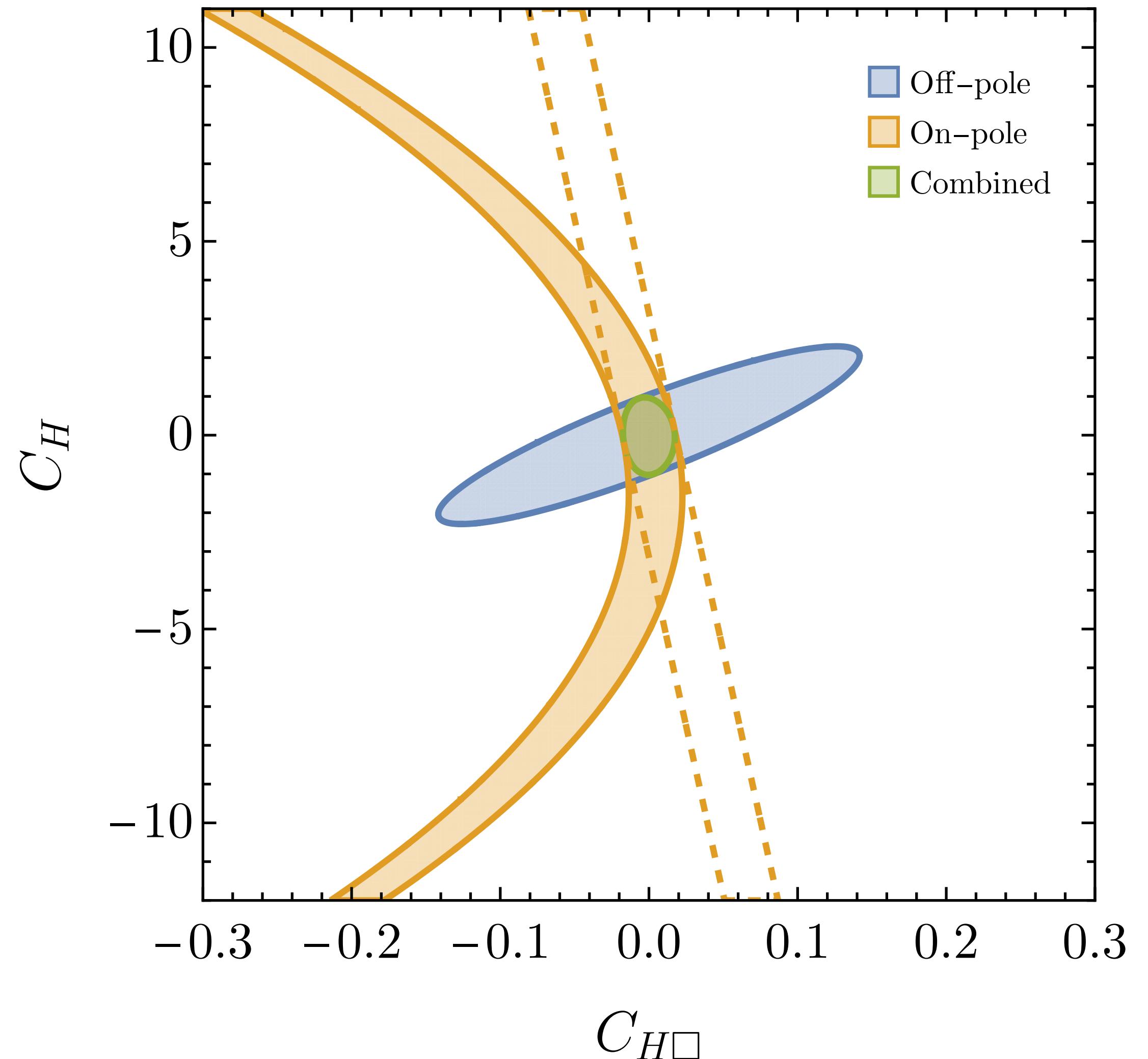
ACE: Negative Aspects

- Access to many operators at the Z pole!
- Running and NLO effects on Z-pole are not negligible!
- Flavour assumptions matter!
- Closed global fits very complicated!



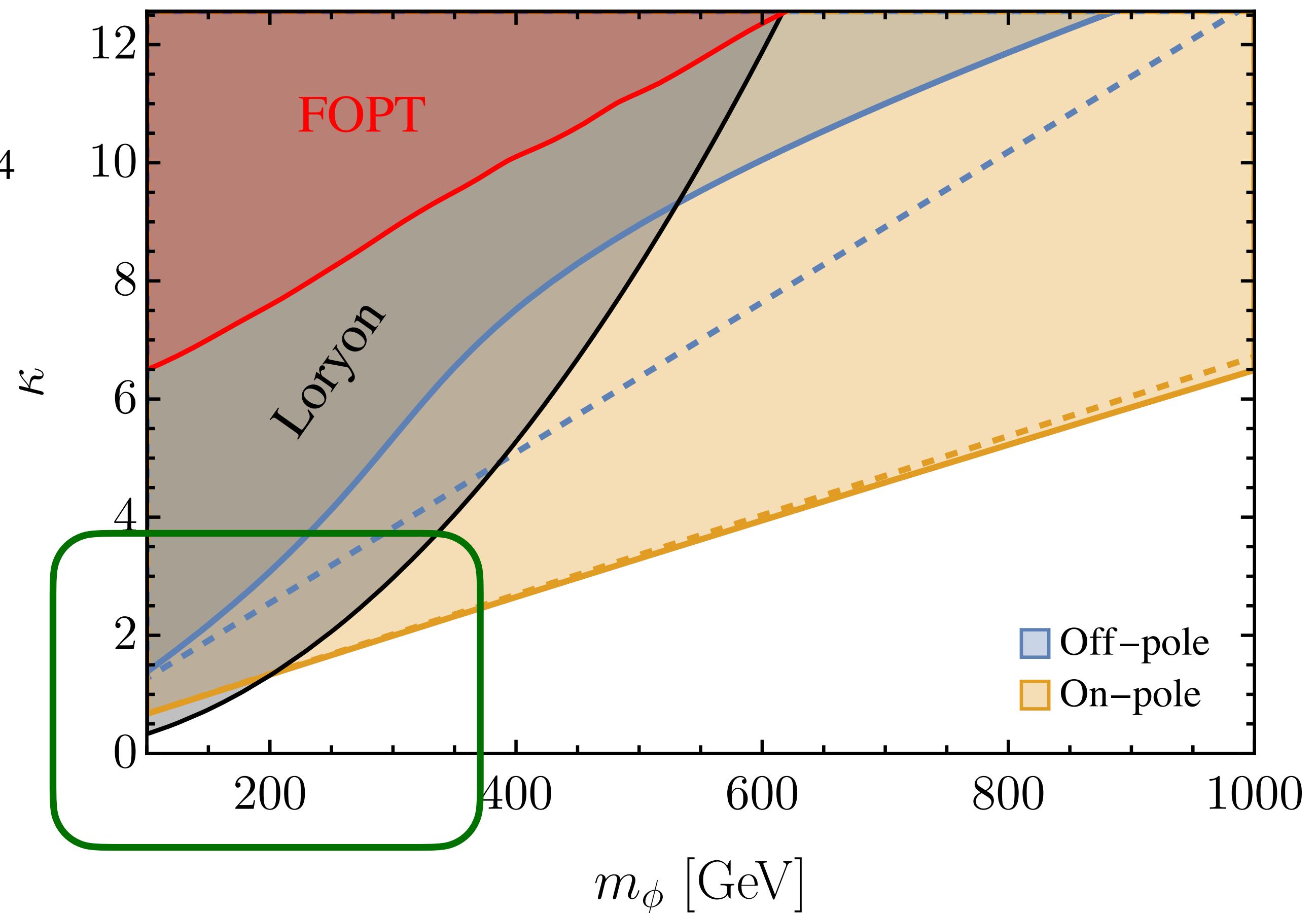
ACE: Bright Side

- **Access to many operators at the Z pole!**
- Global flat directions, but overall volume of parameter space allowed shrinks
- BSM models activate correlated operators controlled by few params
- For concrete BSM models, this is **undoubtedly good news!**



ACE in action: Real Singlet Scalar

- Real Singlet Scalar with \mathbb{Z}_2 -symmetry
- $\mathcal{L} \supset (\partial_\mu \phi)^2 - \frac{1}{2} m_\phi^2 \phi^2 - \frac{1}{2} \kappa \phi^2 |H|^2 - \frac{1}{4!} \lambda \phi^4$
- **Mass contribution from Higgs $\propto \kappa v^2$**
- **Loryon:** Most of the mass comes from the Higgs
- Only loryon still allowed after FCCee (2409.18177)
- **Z pole covers entire parameter space!**



**EFT breaks down Full two loop calculation: WIP (VM, Stefanek, You, 25xx.xxxxxx)

Conclusion

- Extreme precision at Z pole can **compensate for an O(100) suppression or enhancement** in the cross section
- Loop effects + RGE have **significant phenomenological implications**
- Any fit should be done **globally including full NLO Z-pole observables + RGE**
- Z-pole run is **extremely versatile and will likely herald any new physics** that can be further explored at higher energy runs
- Beautiful **complementarity between ALL FCCee runs**

Accuracy Complements Energy

Thank you for your attention!

Questions?

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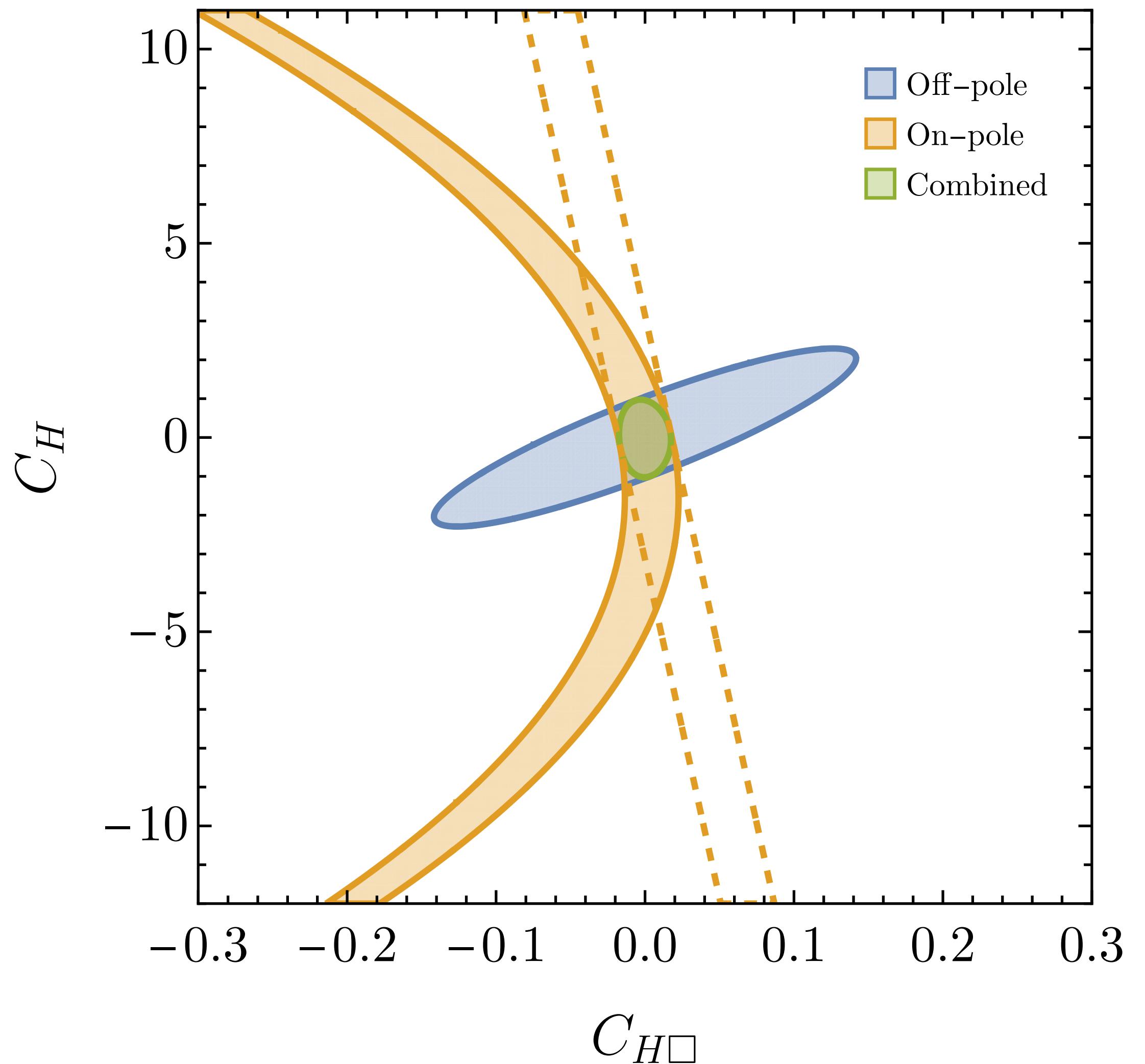
- Why are there two lines in the Singlet scalar plot?
- You have a constraint on C_H at the Z-pole, where is that coming from?
- What do you mean by “flavour assumptions matter”?
- Can you elaborate why you say that RGE effects are important?
- Are there other models for which this is relevant?
-

Backup Slides

ACE in action: Real Singlet Scalar

- Real Singlet Scalar with \mathbb{Z}_2 -symmetry

$$\mathcal{L} \supset (\partial_\mu \phi)^2 - \frac{1}{2} m_\phi^2 \phi^2 - \frac{1}{2} \kappa \phi^2 |H|^2 - \frac{1}{4!} \lambda \phi^4$$

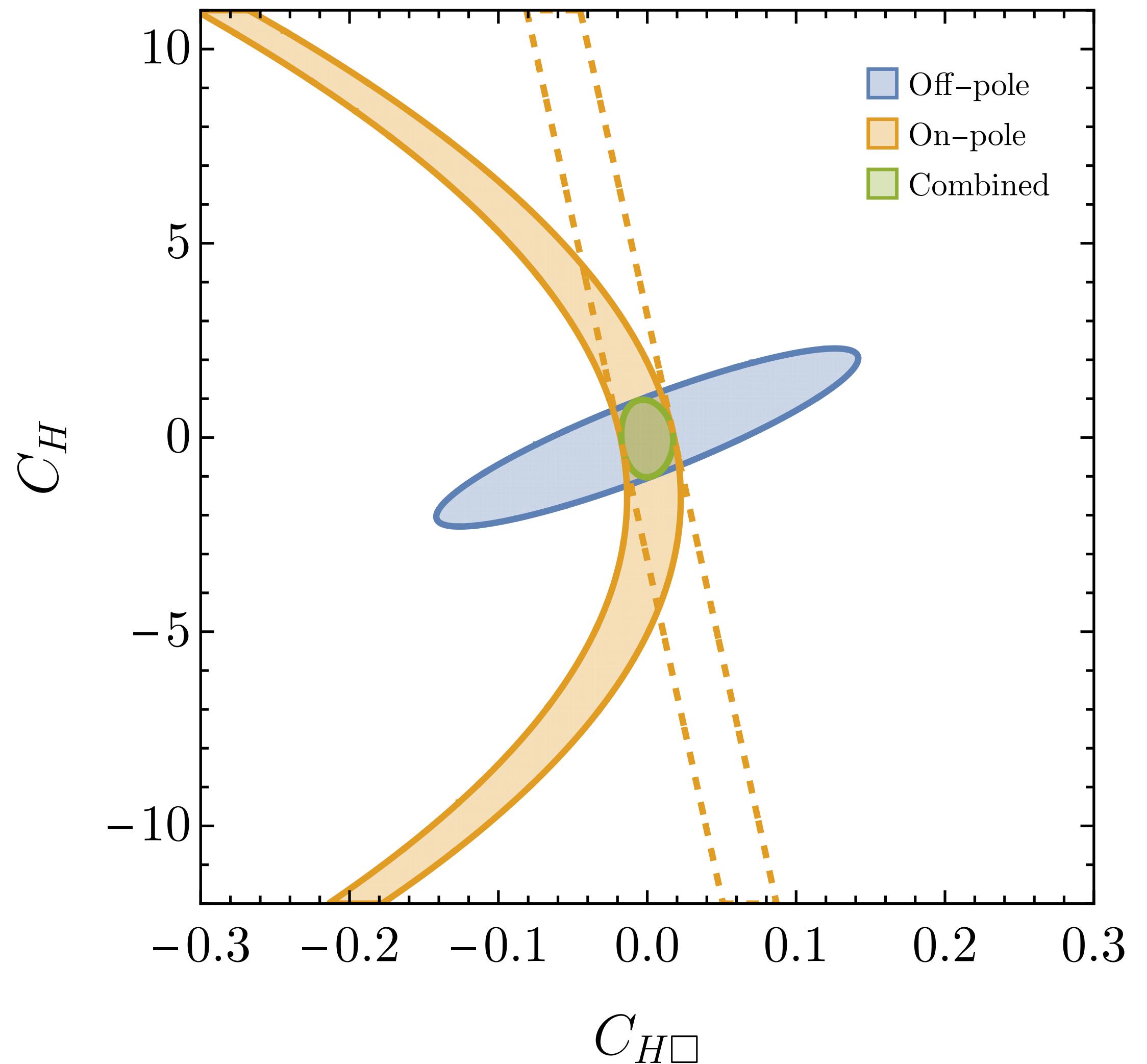


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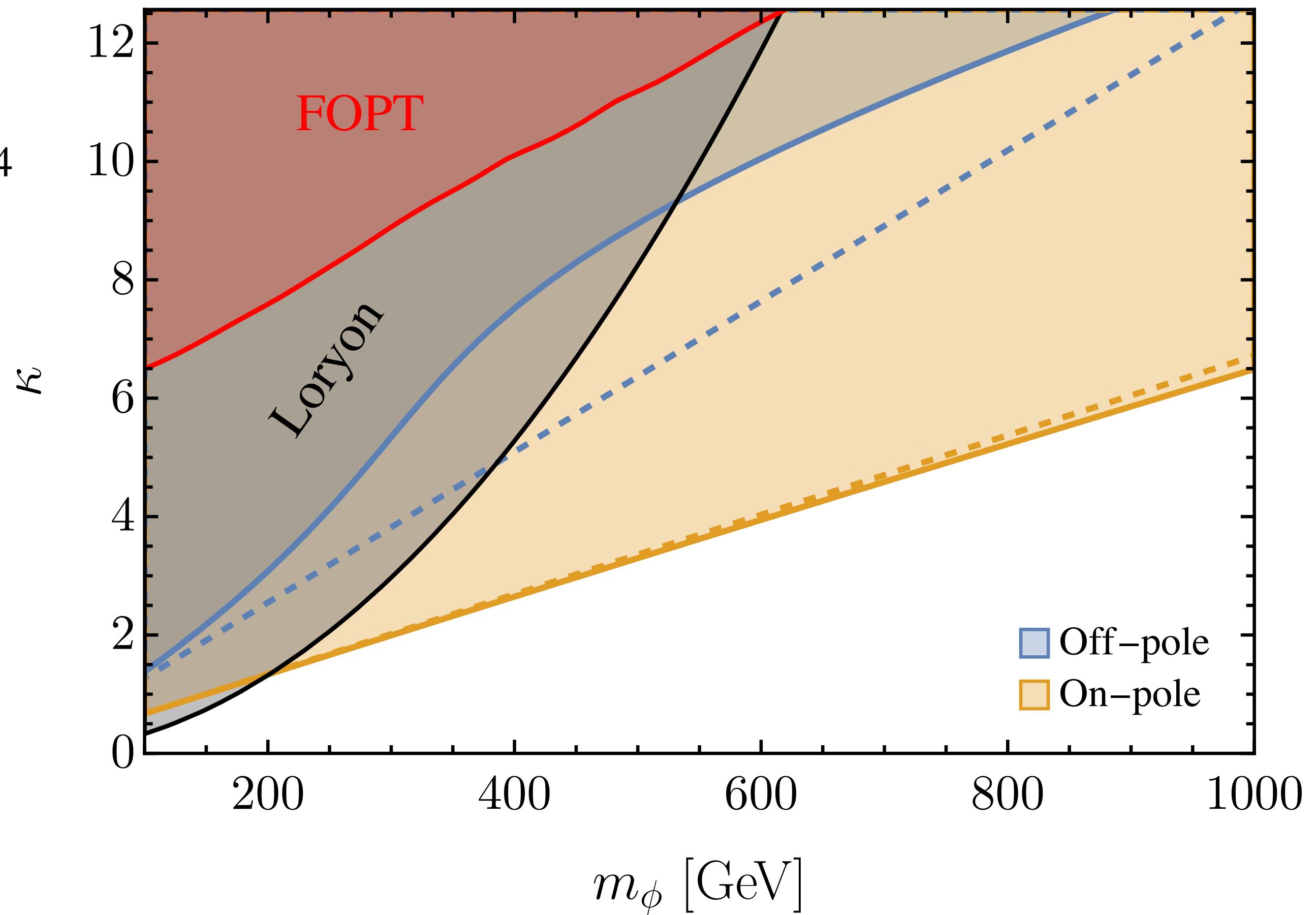
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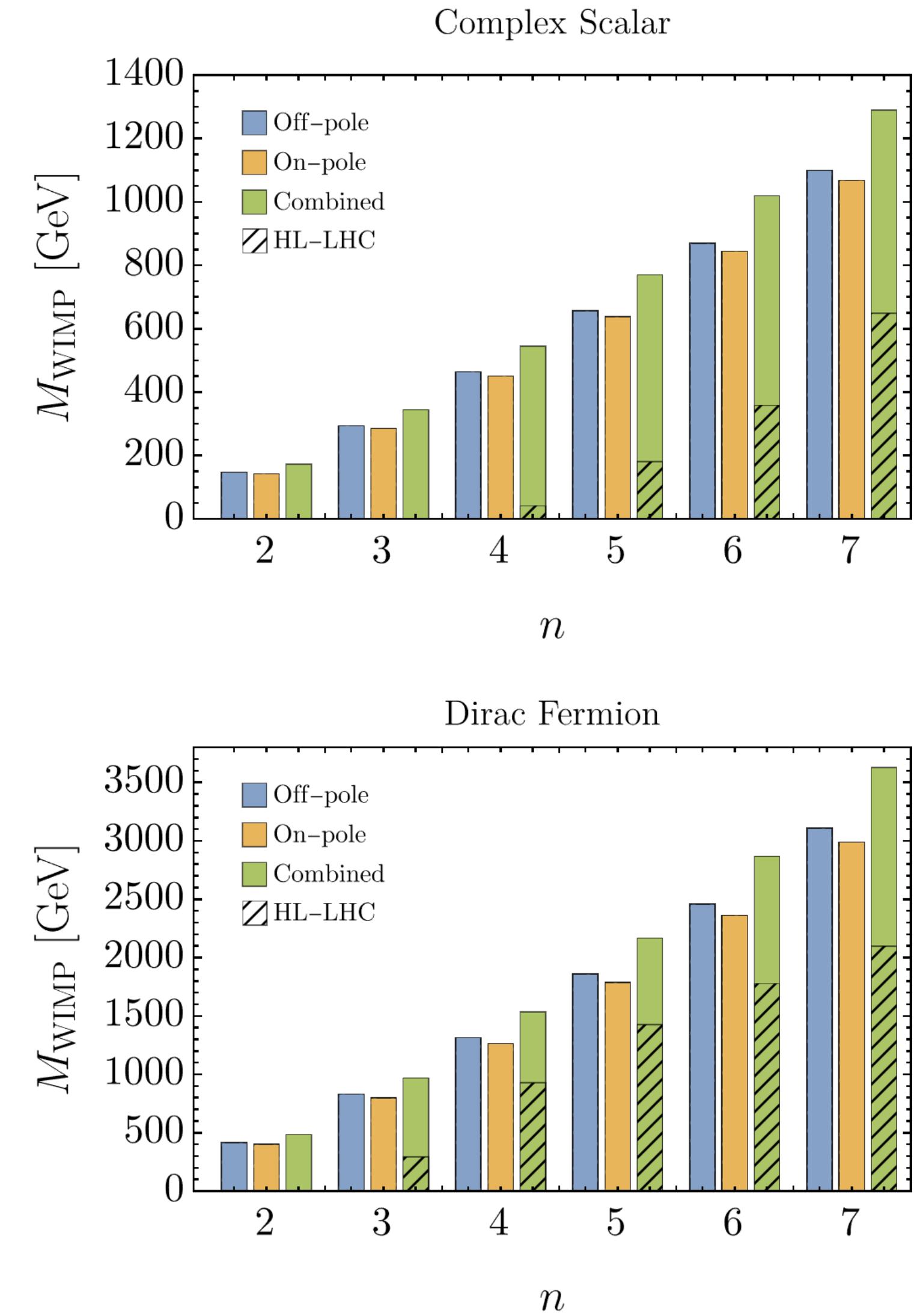
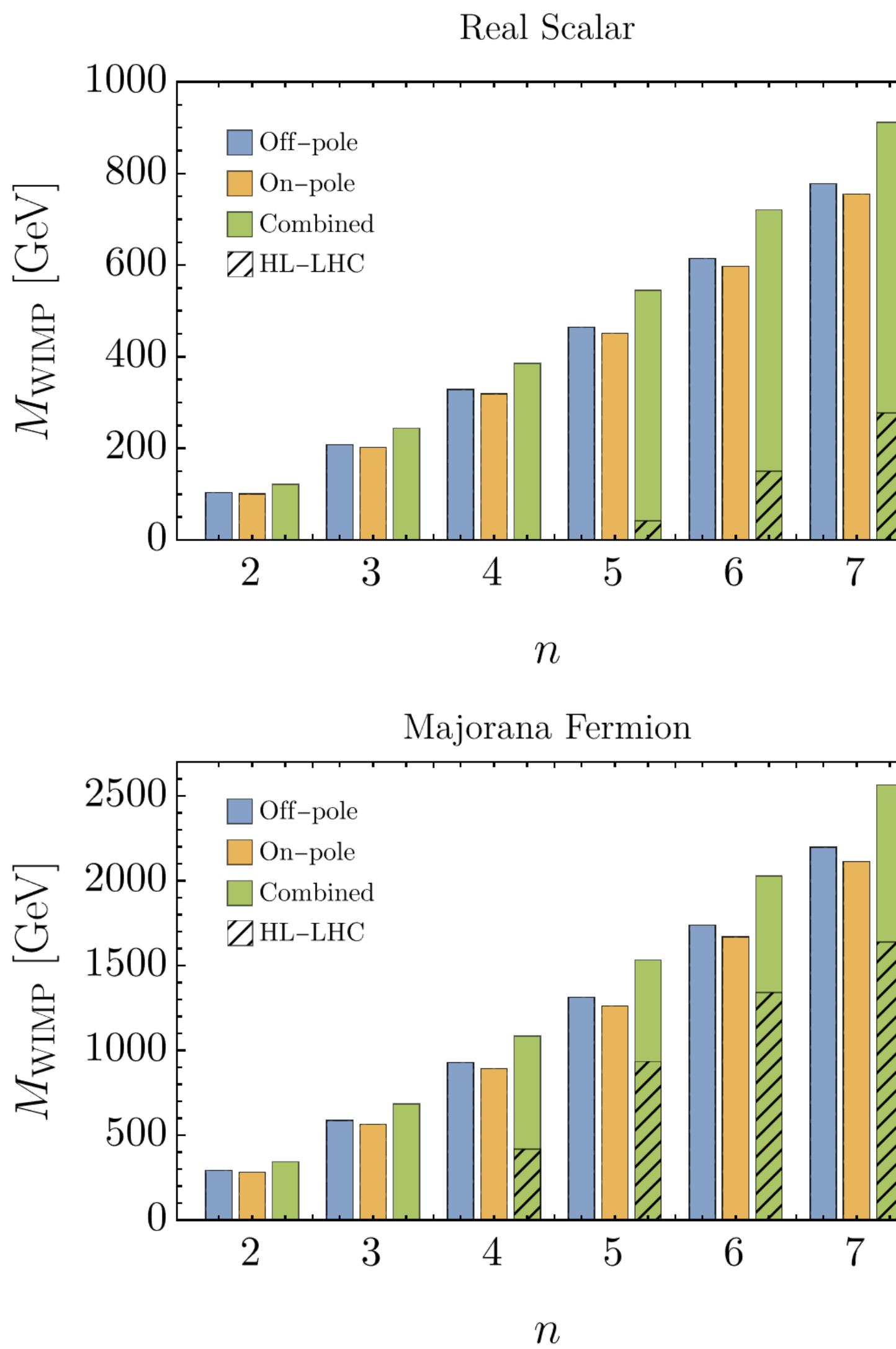
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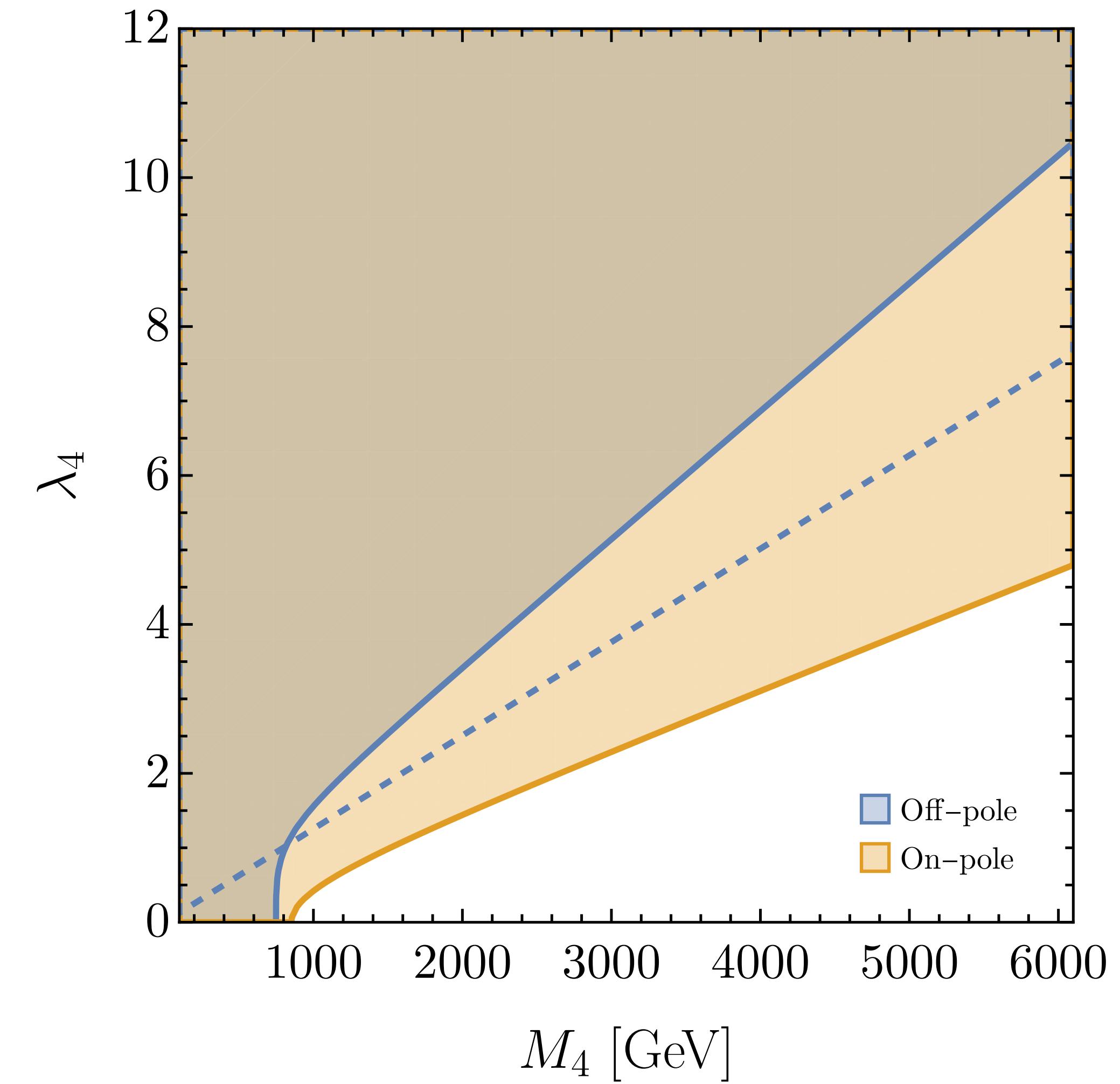
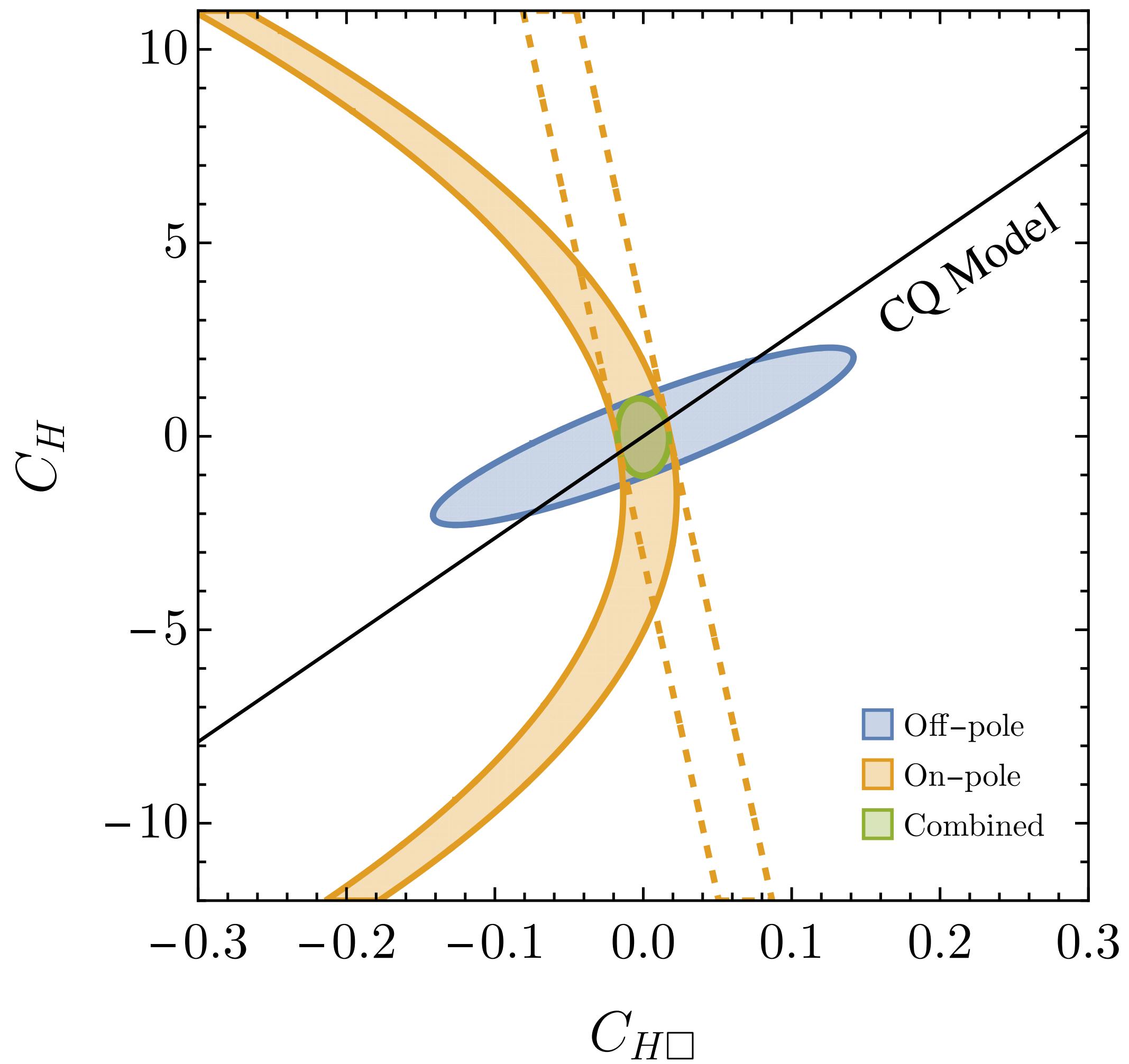


ACE in action: WIMPs

- Higher dimensional Representations of $SU(2)_L$
- Could be Dark Matter
- Can **significantly improve upon HL-LHC** constraints
- Competitive with direct production of low n multiplets!



Custodial Quadruplet



Warsaw Basis

X^3		H^6 and H^4D^2		ψ^2H^3	
\mathcal{O}_G	$f^{ABC}G_\mu^{A\nu}G_\nu^{B\rho}G_\rho^{C\mu}$	\mathcal{O}_H	$(H^\dagger H)^3$	\mathcal{O}_{eH}	$(H^\dagger H)(\bar{l}_p e_r H)$
$\mathcal{O}_{\tilde{G}}$	$f^{ABC}\tilde{G}_\mu^{A\nu}G_\nu^{B\rho}G_\rho^{C\mu}$	$\mathcal{O}_{H\square}$	$(H^\dagger H)\square(H^\dagger H)$	\mathcal{O}_{uH}	$(H^\dagger H)(\bar{q}_p u_r \tilde{H})$
\mathcal{O}_W	$\varepsilon^{IJK}W_\mu^{I\nu}W_\nu^{J\rho}W_\rho^{K\mu}$	\mathcal{O}_{HD}	$(H^\dagger D^\mu H)^*(H^\dagger D_\mu H)$	\mathcal{O}_{dH}	$(H^\dagger H)(\bar{q}_p d_r H)$
$\mathcal{O}_{\widetilde{W}}$	$\varepsilon^{IJK}\tilde{W}_\mu^{I\nu}W_\nu^{J\rho}W_\rho^{K\mu}$				
X^2H^2		ψ^2XH		ψ^2H^2D	
\mathcal{O}_{HG}	$H^\dagger H G_{\mu\nu}^A G^{A\mu\nu}$	\mathcal{O}_{eW}	$(\bar{l}_p \sigma^{\mu\nu} e_r) \tau^I H W_{\mu\nu}^I$	$\mathcal{O}_{Hl}^{(1)}$	$(H^\dagger i \overset{\leftrightarrow}{D}_\mu H)(\bar{l}_p \gamma^\mu l_r)$
$\mathcal{O}_{H\tilde{G}}$	$H^\dagger H \tilde{G}_{\mu\nu}^A G^{A\mu\nu}$	\mathcal{O}_{eB}	$(\bar{l}_p \sigma^{\mu\nu} e_r) H B_{\mu\nu}$	$\mathcal{O}_{Hl}^{(3)}$	$(H^\dagger i \overset{\leftrightarrow}{D}_\mu^I H)(\bar{l}_p \tau^I \gamma^\mu l_r)$
\mathcal{O}_{HW}	$H^\dagger H W_{\mu\nu}^I W^{I\mu\nu}$	\mathcal{O}_{uG}	$(\bar{q}_p \sigma^{\mu\nu} T^A u_r) \tilde{H} G_{\mu\nu}^A$	\mathcal{O}_{He}	$(H^\dagger i \overset{\leftrightarrow}{D}_\mu H)(\bar{e}_p \gamma^\mu e_r)$
$\mathcal{O}_{H\widetilde{W}}$	$H^\dagger H \tilde{W}_{\mu\nu}^I W^{I\mu\nu}$	\mathcal{O}_{uW}	$(\bar{q}_p \sigma^{\mu\nu} u_r) \tau^I \tilde{H} W_{\mu\nu}^I$	$\mathcal{O}_{Hq}^{(1)}$	$(H^\dagger i \overset{\leftrightarrow}{D}_\mu H)(\bar{q}_p \gamma^\mu q_r)$
\mathcal{O}_{HB}	$H^\dagger H B_{\mu\nu} B^{\mu\nu}$	\mathcal{O}_{uB}	$(\bar{q}_p \sigma^{\mu\nu} u_r) \tilde{H} B_{\mu\nu}$	$\mathcal{O}_{Hq}^{(3)}$	$(H^\dagger i \overset{\leftrightarrow}{D}_\mu^I H)(\bar{q}_p \tau^I \gamma^\mu q_r)$
$\mathcal{O}_{H\tilde{B}}$	$H^\dagger H \tilde{B}_{\mu\nu} B^{\mu\nu}$	\mathcal{O}_{dG}	$(\bar{q}_p \sigma^{\mu\nu} T^A d_r) H G_{\mu\nu}^A$	\mathcal{O}_{Hu}	$(H^\dagger i \overset{\leftrightarrow}{D}_\mu H)(\bar{u}_p \gamma^\mu u_r)$
\mathcal{O}_{HWB}	$H^\dagger \tau^I H W_{\mu\nu}^I B^{\mu\nu}$	\mathcal{O}_{dW}	$(\bar{q}_p \sigma^{\mu\nu} d_r) \tau^I H W_{\mu\nu}^I$	\mathcal{O}_{Hd}	$(H^\dagger i \overset{\leftrightarrow}{D}_\mu H)(\bar{d}_p \gamma^\mu d_r)$
$\mathcal{O}_{H\widetilde{W}B}$	$H^\dagger \tau^I H \tilde{W}_{\mu\nu}^I B^{\mu\nu}$	\mathcal{O}_{dB}	$(\bar{q}_p \sigma^{\mu\nu} d_r) H B_{\mu\nu}$	\mathcal{O}_{Hud}	$i(\tilde{H}^\dagger D_\mu H)(\bar{u}_p \gamma^\mu d_r)$
$(\bar{L}L)(\bar{L}L)$		$(\bar{R}R)(\bar{R}R)$		$(\bar{L}L)(\bar{R}R)$	
\mathcal{O}_{ll}	$(\bar{l}_p \gamma_\mu l_r)(\bar{l}_s \gamma^\mu l_t)$	\mathcal{O}_{ee}	$(\bar{e}_p \gamma_\mu e_r)(\bar{e}_s \gamma^\mu e_t)$	\mathcal{O}_{le}	$(\bar{l}_p \gamma_\mu l_r)(\bar{e}_s \gamma^\mu e_t)$
$\mathcal{O}_{qq}^{(1)}$	$(\bar{q}_p \gamma_\mu q_r)(\bar{q}_s \gamma^\mu q_t)$	\mathcal{O}_{uu}	$(\bar{u}_p \gamma_\mu u_r)(\bar{u}_s \gamma^\mu u_t)$	\mathcal{O}_{lu}	$(\bar{l}_p \gamma_\mu l_r)(\bar{u}_s \gamma^\mu u_t)$
$\mathcal{O}_{qq}^{(3)}$	$(\bar{q}_p \gamma_\mu \tau^I q_r)(\bar{q}_s \gamma^\mu \tau^I q_t)$	\mathcal{O}_{dd}	$(\bar{d}_p \gamma_\mu d_r)(\bar{d}_s \gamma^\mu d_t)$	\mathcal{O}_{ld}	$(\bar{l}_p \gamma_\mu l_r)(\bar{d}_s \gamma^\mu d_t)$
$\mathcal{O}_{lq}^{(1)}$	$(\bar{l}_p \gamma_\mu l_r)(\bar{q}_s \gamma^\mu q_t)$	\mathcal{O}_{eu}	$(\bar{e}_p \gamma_\mu e_r)(\bar{u}_s \gamma^\mu u_t)$	\mathcal{O}_{qe}	$(\bar{q}_p \gamma_\mu q_r)(\bar{e}_s \gamma^\mu e_t)$
$\mathcal{O}_{lq}^{(3)}$	$(\bar{l}_p \gamma_\mu \tau^I l_r)(\bar{q}_s \gamma^\mu \tau^I q_t)$	\mathcal{O}_{ed}	$(\bar{e}_p \gamma_\mu e_r)(\bar{d}_s \gamma^\mu d_t)$	$\mathcal{O}_{qu}^{(1)}$	$(\bar{q}_p \gamma_\mu q_r)(\bar{u}_s \gamma^\mu u_t)$
		$\mathcal{O}_{ud}^{(1)}$	$(\bar{u}_p \gamma_\mu u_r)(\bar{d}_s \gamma^\mu d_t)$	$\mathcal{O}_{qu}^{(8)}$	$(\bar{q}_p \gamma_\mu T^A q_r)(\bar{u}_s \gamma^\mu T^A u_t)$
		$\mathcal{O}_{ud}^{(8)}$	$(\bar{u}_p \gamma_\mu T^A u_r)(\bar{d}_s \gamma^\mu T^A d_t)$	$\mathcal{O}_{qd}^{(1)}$	$(\bar{q}_p \gamma_\mu q_r)(\bar{d}_s \gamma^\mu d_t)$
				$\mathcal{O}_{qd}^{(8)}$	$(\bar{q}_p \gamma_\mu T^A q_r)(\bar{d}_s \gamma^\mu T^A d_t)$
$(\bar{L}R)(\bar{R}L)$ and $(\bar{L}R)(\bar{L}R)$		B -violating			
\mathcal{O}_{ledq}	$(\bar{l}_p^j e_r)(\bar{d}_s q_t^j)$	\mathcal{O}_{duq}	$\varepsilon^{\alpha\beta\gamma} \varepsilon_{jk} [(d_p^\alpha)^T C u_r^\beta] [(q_s^\gamma)^T C l_t^k]$		
$\mathcal{O}_{quqd}^{(1)}$	$(\bar{q}_p^j u_r) \varepsilon_{jk} (\bar{q}_s^k d_t)$	\mathcal{O}_{qqu}	$\varepsilon^{\alpha\beta\gamma} \varepsilon_{jk} [(q_p^{\alpha j})^T C q_r^{\beta k}] [(u_s^\gamma)^T C e_t]$		
$\mathcal{O}_{quqd}^{(8)}$	$(\bar{q}_p^j T^A u_r) \varepsilon_{jk} (\bar{q}_s^k T^A d_t)$	\mathcal{O}_{qqq}	$\varepsilon^{\alpha\beta\gamma} \varepsilon_{jn} \varepsilon_{km} [(q_p^{\alpha j})^T C q_r^{\beta k}] [(q_s^{\gamma m})^T C l_t^n]$		
$\mathcal{O}_{lequ}^{(1)}$	$(\bar{l}_p^j e_r) \varepsilon_{jk} (\bar{q}_s^k u_t)$	\mathcal{O}_{duu}	$\varepsilon^{\alpha\beta\gamma} [(d_p^\alpha)^T C u_r^\beta] [(u_s^\gamma)^T C e_t]$		
$\mathcal{O}_{lequ}^{(3)}$	$(\bar{l}_p^j \sigma_{\mu\nu} e_r) \varepsilon_{jk} (\bar{q}_s^k \sigma^{\mu\nu} u_t)$				

Feynman Diagrams