



# Study of nucleon spin structure by the Drell-Yan process in the COMPASS experiment

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LIP - Lisbon



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## Outline



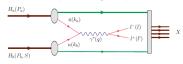
- Introduction to the Drell-Yan (DY) process
- Goals of the DY measurement @ COMPASS
- COMPASS Experiment Spectrometer description
- Feasibility of the measurement
- Acceptances, event rates and statistical errors





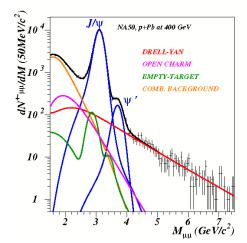


Annihilation of a  $q\bar{q}$  pair from a 2 hadrons collision, producing a lepton pair



- Strong  $\sigma_{DY}$  decreases with the dimuon mass  $(\sigma_{DY} \propto M_{\mu\mu}^{-4})$ .
- The Drell-Yan signal is very clean above 4  $GeV/c^2$  dimuon mass. It is the region where we are interested in, the High Mass Region (HMR).

Dimuon mass distribution for p @ 400 GeV/cin a Pb target (NA50 Collaboration)





## Drell-Yan angular distribution



The angular distribution of the DY events can be written as:

$$\frac{1}{\sigma}\frac{d\sigma}{d\Omega} = \frac{3}{4\pi}\frac{1}{\lambda+3}[1+\lambda\cos^2\theta + \mu\sin2\theta\cos\phi + \frac{\nu}{2}\sin^2\theta\cos2\phi]$$

The Lam-Tung sum rule  $|1 - \lambda - 2\nu = 0|$  relates the modulation amplitudes of the DY process.

The collinear hypothesis implies:  $\lambda=1,\,\mu=0,\,\nu=0$  .

It was shown by the NA10 (CERN) and E615 (Fermilab) experiments that this hypothesis is violated.

### They measured a modulation of $\cos 2\phi$ up to 30%

This means we cannot neglect the intrinsic transverse momentum  $k_T$  of quarks inside hadrons.

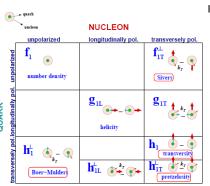


## Polarized DY @ COMPASS



The nucleon structure in first order QCD is described by:

- 3 PDFs in the collinear approximation
- ullet 8 PDFs taking into account the quark intrinsic transverse momentum,  $k_T$



In the DY COMPASS program we will access 4 of them:

- Transversity h<sub>1</sub>: describes the difference between the number density of quarks with parallel and anti-parallel spin w.r.t. the transversely polarised, relatively to the beam direction, father hadron
- Boer-Mulders h<sub>1</sub><sup>⊥</sup>: describes the correlation between the transverse spin and the transverse momentum of a quark in an unpolarised nucleon
- Sivers f<sub>1</sub><sup>±</sup>: describes the influence of the nucleon transverse spin in the transverse momentum distribution of the quark
- Pretzelosity  $h_{1T}^{\perp}$ : describes the transverse polarization of a quark, along its intrinsic transverse momentum direction



### Polarized DY @ COMPASS



Arnold *et al.*<sup>1</sup> derived the full expression of the  $\sigma_{DY}$ , for arbitrarily polarized beam and target.

Having an unpolarized beam and a transversely polarized target the  $\sigma_{DY}$  in LO can be written as:

$$\begin{split} \frac{d\sigma}{d^{4}qd\Omega} &= \frac{\alpha_{em}^{2}}{Fq^{2}} \hat{\sigma}_{U} \{ (1 + D_{[\sin^{2}\theta]} \middle| A_{U}^{\cos 2\phi} \middle| \cos 2\phi) + |\overrightarrow{S}_{T}| [A_{T}^{\sin \phi_{S}} \middle| \sin \phi_{S} \\ &+ D_{[\sin^{2}\theta]} (A_{T}^{\sin(2\phi + \phi_{S})} \middle| \sin(2\phi + \phi_{S}) + A_{T}^{\sin(2\phi - \phi_{S})} \middle| \sin(2\phi - \phi_{S}))] \} \end{split}$$

## where:

- ullet  $\theta$  and  $\phi$  are the polar and azimuthal angles of  $\mu^+$  in the Collins-Soper reference frame
- ullet  $\phi_{ extsf{S}}$  is the angle between the transverse spin of the target nucleon and the transverse momentum of the  $\gamma^*$
- F is given by  $F = 4\sqrt{(P_{\pi} \cdot P_{p})^{2} M_{\pi}^{2} M_{p}^{2}}$
- q is the  $\gamma^*$  four-momentum
- $\bullet$   $\hat{\sigma}_U$  is the part of the cross-section surviving the integration over the angles  $\phi$  and  $\phi_S$
- $|\overrightarrow{S}_T|$  is the target polarization value
- ullet  $D_{[\sin^2 heta]}$  is the virtual photon depolarization factor

<sup>&</sup>lt;sup>1</sup>S. Arnold et al. Phys.Rev. D79 (2009)034005



## Polarized DY @ COMPASS



$$\begin{split} \frac{d\sigma}{d^{4}qd\Omega} &= \frac{\alpha_{em}^{2}}{Fq^{2}} \hat{\sigma}_{U} \{ (1 + D_{[\sin^{2}\theta]} \middle| A_{U}^{\cos2\phi} \middle| \cos2\phi) + |\overrightarrow{S}_{T}| [ \middle| A_{T}^{\sin\phi_{S}} \middle| \sin\phi_{S} \\ &+ D_{[\sin^{2}\theta]} ( \middle| A_{T}^{\sin(2\phi + \phi_{S})} \middle| \sin(2\phi + \phi_{S}) + | A_{T}^{\sin(2\phi - \phi_{S})} \middle| \sin(2\phi - \phi_{S}))] \} \end{split}$$

The azimuthal asymmetries A contain a convolution of 2 PDFs of the target and beam hadrons:

- ullet  $A_U^{\cos2\phi}$  gives access to the Boer-Mulders functions of both hadrons.
- $A_T^{\sin\phi_S}$  gives access to the unpolarised PDF of beam hadron and the Sivers function of the target nucleon.
- $A_T^{sin(2\phi+\phi_S)}$  gives access to the Boer-Mulders function of the beam hadron and to the pretzelosity function of the target nucleon.
- $A_T^{sin(2\phi-\phi_S)}$  gives access to the Boer-Mulders function of the beam hadron and to the transversity function of the target nucleon.

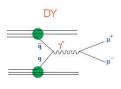
We need to disentangle the PDFs in each of these asymmetries  $\Rightarrow$  it requires some input.

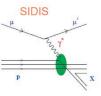


## DY vs SIDIS



The Sivers  $(f_{1T}^{\perp})$  and the Boer-Mulders  $(h_{1}^{\perp})$  functions are time-reversal odd functions. This leads to the prediction that they must change sign when accessed from DY or SIDIS<sup>2</sup>.





$$f_{1TDY}^{\perp} = -f_{1TSIDIS}^{\perp}$$

$$h_{1\ DY}^{\perp} = -h_{1\ SIDIS}^{\perp}$$

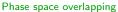
The experimental confirmation of this sign change is considered a crucial test of non-perturbative QCD.

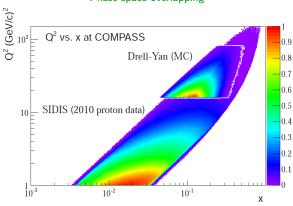
<sup>&</sup>lt;sup>2</sup>J.C. Collins, Phys. Lett. B536 (2002) 43, J.C. Collins, talk at LIGHT CONE 2008



# DY vs SIDIS









# $J/\psi$ duality



The  $J/\psi$  duality corresponds to consider a possible analogy in the production mechanisms of both  $J/\psi$   $(q\bar{q} \to J/\psi + X)$  and  $\gamma^*$ .

Studying the charmonium mass region in the dilepton decay channel:

- $lacktriangledisplays \begin{picture}(10,0) \put(0,0){\line(0,0){100}} \put(0,0){\line(0,0){100}}$
- ② Access PDFs from  $J/\psi$  events larger statistics available



# COMPASS @ CERN



## COmmon Muon Proton Apparatus for Structure and Spectroscopy

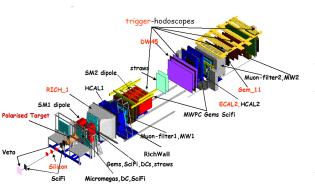




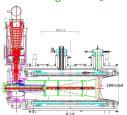
# Experimental setup



Beam:  $\pi^-$  @ 190 GeV/c



### Polarized target – $NH_3$



- Beam telescope:
  - SciFis
- Large Angle Spectrometer (LAS):
  - SM1 magnet
  - Tracking detectors
  - RICH
  - ECAL1
  - HCAL1
  - Muons Filter 1
- Small Angle Spectrometer (SAS):
  - SM2 magnet
  - Tracking detectors
  - ECAL2
  - HCAL2
  - Muons Filter 2



## Feasibility of the measurement



In 2007, 2008 and 2009 short Drell-Yan beam tests were performed, to check the feasibility of the measurement.

#### 2007 test:

- open spectrometer (without hadron absorber)
- $\pi^-$  beam @ 160 GeV/c ( $I_{beam}=4\times10^6~\pi/s$ ) and an unpolarized NH<sub>3</sub> target
- we collected  $\simeq 90000$  dimuons in  $\lesssim 12$  hours
- validation of the  $J/\psi$  yields expected

#### 2008 test:

- open spectrometer (without hadron absorber)
- $\bullet$  increase of the beam intensity up to  $6.5\times 10^6~\pi/s$
- high occupancy of detectors closer to the target region ⇒ confirmed that is mandatory the use of a hadron absorber
- verification of the spectrometer response and radiation doses



# Feasibility of the measurement



#### 2009 test was done using:

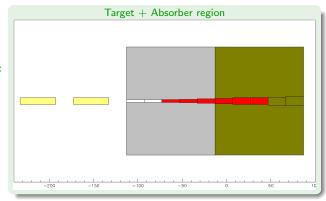
- $\bullet$   $\pi^-$  beam @ 190 GeV/c
- Two target cells (polyethylene) with 40 cm length and 5 cm diameter, spaced by 20 cm
- A prototype of an absorber:
  - $\bullet$  Two blocks, of concrete and stainless steel, 100 cm length each, and 80  $\times$  80  $cm^2$  in transverse dimensions each
  - A beam plug, inside the central part of absorber, made of W and steel disks.

Number of radiation lengths (multiple scattering for muons):

$$x/X_0 = 66.17$$

Number of interaction lengths (stopping power for pions):

$$x/\lambda_{int} = 6.69$$



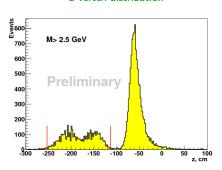


## 2009 DY beam test - results

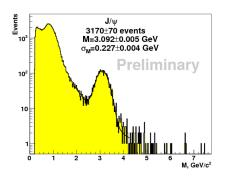


### 3 days test

z vertex distribution



### Dimuon mass distribution



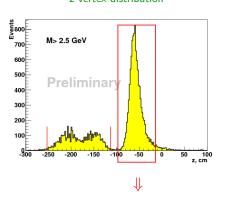


### 2009 DY beam test - results

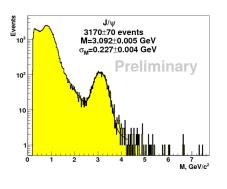


#### 3 days test

#### z vertex distribution



#### Dimuon mass distribution



Huge number of events due to the fact we didn't have a dimuon trigger.

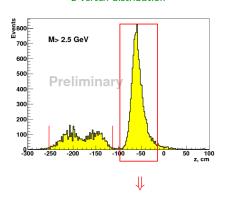


## 2009 DY beam test - results



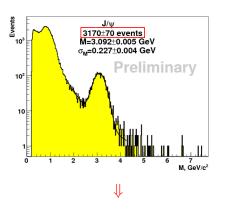
#### 3 days test

#### z vertex distribution



Huge number of events due to the fact we didn't have a dimuon trigger.

#### Dimuon mass distribution



The number of expected  $J/\psi$  was reached.



## Current DY setup status



The proposal setup was in the meanwhile optimized. The newer version is:

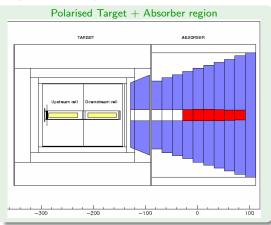
- Two target cells (NH<sub>3</sub>) inside the dipole with 55 cm length and 4 cm diameter, spaced by 20 cm;
- The absorber is 236 cm long, made of Al<sub>2</sub>O<sub>3</sub>;
- The plug inside the absorber is made of 6 disks of W, 20 cm long each and 20 cm of Alumina in the most downstream part (total of 140 cm).

Number of radiation lengths (multiple scattering for muons):

$$x/X_0 = 33.53$$
 (66.17 for 2009 test absorber)

Number of interaction lengths (stopping power for pions):

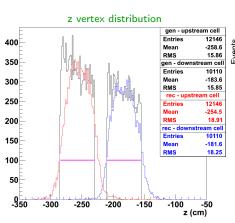
$$x/\lambda_{int} = 7.25$$
 (6.69 for 2009 test absorber)

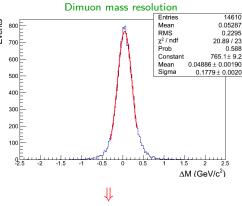




# Current DY setup status - MC results







 $\Delta z = 6 \ cm$  in HMR  $(M_{\mu\mu} > 4 \ GeV/c)$ 

 $\Delta M = 178~MeV/c^2$  in HMR  $(M_{\mu\mu} > 4~GeV/c)$ 

#### Contamination:

Upstream cell — 0.29 % Downstream cell — 1.16 %

⇒ A very low level of contamination.



## Acceptances, event rates and statistical errors



The dimuons acceptance in the HMR  $(M_{\mu\mu}>4~GeV/c)$  is 39%.

The accepted dimuons are:

$$\bullet$$
  $\mu_1$  (1<sup>st</sup> spectrometer) &  $\mu_2$  (1<sup>st</sup> spectrometer) - 22%

$$\bullet$$
  $\mu_1$  (2<sup>nd</sup> spectrometer) &  $\mu_2$  (2<sup>nd</sup> spectrometer) - 2%

$$\bullet$$
  $\mu_1$  (1<sup>st</sup> spectrometer) &  $\mu_2$  (2<sup>nd</sup> spectrometer) - 18%

We expect an DY event rate of 900 events/day in the HMR assuming:

$$\bullet$$
  $\pi^-$  beam with 190  $GeV/c$ 

• 
$$I_{beam} = 6 \times 10^7 \ particles/s$$

• 
$$L = 1.2 \times 10^{32} \text{ cm}^{-2} \text{s}^{-1}$$

Assuming two years of data taking we expect the following statistical errors in azimuthal asymmetries:

Asymmetry	Uncertainty in HMR (%)
$\delta A_U^{\cos 2\phi}$	0.43
$\delta A_T^{\sin \phi_S}$	1.34
$\delta A_T^{\sin(2\phi+\phi_S)}$	2.70
$\delta A_T^{\sin(2\phi-\phi_S)}$	2.70



# Azimuthal asymmetries and statistical errors (2010 status)

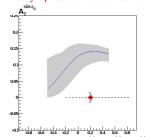


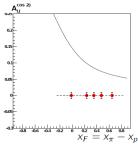
Theory update is needed to take into account the  $Q^2$  evolution.



Sivers

M. Anselmino *et al*, Phys. Rev. D 79 (2009) 054010





Boer-Mulders

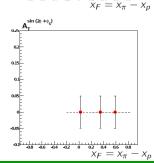
 $\otimes$ 

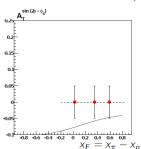
Boer-Mulders

B. Zhang *et al*, Phys. Rev. D 77 (2008)

054011







Boer-Mulders

Transversity

A. N. Sissakian et al, Phys. Part. Nucl. 41:

64-100, 2010



# Summary



- The opportunity to study, in the same experiment, the TMD PDFs from both SIDIS and the DY processes is unique.
- The sign change in Sivers and Boer-Mulders functions when accessed by DY and SIDIS will be checked.
- The feasibility of the measurement was proven after three past beam tests.
- The COMPASS II Proposal was approved by CERN for a first period of 3 years including 1 year for Drell-Yan.
- Drell-Yan data taking will start in 2014 and the second year of data taking is expected in 2017.