

Hard Exclusive ρ^0 -Meson Production at COMPASS

Heiner Wollny for the COMPASS collaboration

CEA Saclay, 91191 Gif-sur-Yvette, France

DOI: <http://dx.doi.org/10.3204/DESY-PROC-2012-02/127>

New results for the transverse target spin azimuthal asymmetry $A_{UT}^{\sin(\phi-\phi_S)}$ for hard exclusive ρ^0 -meson production on a transversely polarised ${}^6\text{LiD}$ and NH_3 target will be presented. The measurement was performed with the COMPASS detector using the 160 GeV/c muon beam of the SPS at CERN. The asymmetry is sensitive to the nucleon helicity-flip generalised parton distribution E , which is related to the orbital angular momentum of quarks in the nucleon.

1 Introduction

Hard exclusive vector meson production on nucleons is an important tool to study strong interactions. Moreover, in Bjorken kinematics it provides access to generalised parton distributions (GPDs) [1, 2, 3]. The GPDs contain a wealth of information on the structure of the nucleon. In particular, they embody both nucleon electromagnetic form factors and parton distribution functions. Furthermore GPDs correlate longitudinal momenta and transverse spatial position of partons referred to as 3-dimensional nucleon tomography [4].

At leading twist, vector meson production is described by the GPDs H^f and E^f , where f denotes a quark of flavour f or a gluon. The GPDs are functions of t , x and ξ , where t is the squared four-momentum transfer to the nucleon, x is the average and ξ is half the difference of the longitudinal momenta carried by the struck parton in the initial and final state. The GPDs H^f describe the case where the nucleon-helicity is conserved in the scattering process, whereas the GPDs E^f describe the case of nucleon-helicity flip. Hence, in the latter case angular momentum must be involved in order to conserve total angular momentum. The Ji relation, connects the total angular momentum of a given parton species f to the second moment of the sum of GPDs H^f and E^f [3]:

$$J^f = \frac{1}{2} \lim_{t \rightarrow 0} \int_{-1}^1 dx x [H^f(x, \xi, t) + E^f(x, \xi, t)]. \quad (1)$$

This relation attracted much attention since it is the only known way to constrain the quark angular momentum to the nucleon spin budget. In vector meson production on unpolarised nucleons the GPDs E^f are suppressed in COMPASS kinematics and one is only sensitive to the GPDs H^f . However, in the cross-section for transversely polarised nucleons the GPDs E^f appear at leading twist in the azimuthal asymmetry $A_{UT}^{\sin(\phi-\phi_S)}$. Here, the indices U and T refer to the unpolarised beam and transversely polarised target, respectively. The superscript $\sin(\phi - \phi_S)$ indicates the type of azimuthal modulation of the cross-section, where ϕ is the azimuthal angle between lepton scattering plane and the plane defined by the virtual photon

and the produced meson, and ϕ_S is the azimuthal angle of the target spin vector relative to the lepton scattering plane. In this paper the asymmetry $A_{UT}^{\sin(\phi-\phi_S)}$ for exclusive ρ^0 -meson production on transversely polarised deuterons and protons is presented.

2 Data sample and event selection

The presented analysis is performed on data taken with the COMPASS spectrometer [5] by scattering positive muons of 160 GeV/c from the CERN SPS off transversely polarised solid state targets. The data taken in 2003 and 2004 with a ${}^6\text{LiD}$ target and the data taken in 2007 and 2010 with a NH_3 target were analysed. For the ${}^6\text{LiD}$ material the average target polarisation is about 0.5, while for NH_3 it is about 0.8. The target dilution factor for exclusive ρ^0 production is typically 0.45 and 0.25 for the ${}^6\text{LiD}$ and NH_3 target, respectively. Before 2006 the target consisted of two separated cells of equal length, which were oppositely polarised. Since 2006 three target cells were used, where neighbouring cells were oppositely polarised and the length of the two outer cells matches the length of the middle one. This ensures a better balanced acceptance for cells with opposite polarisation. Both target layouts allow a simultaneous measurement of both target spin directions compensating flux dependent systematic uncertainties. In addition, to reduce systematic effects of the acceptance the polarisation of the target cells were reversed about every week.

A new solenoid magnet installed during the shut down in 2005 increased the angular acceptance of the experiment from ± 70 mrad to the design value of ± 180 mrad.

Events in the DIS regime are selected by cuts on squared four-momentum transfer $Q^2 > 1 (\text{GeV}/c)^2$, on the fractional energy lost of the muon $0.1 < y < 0.9$ and on the invariant mass of the $\gamma^* - N$ system $W > 5 \text{ GeV}/c^2$. An upper cut on $Q^2 < 10 (\text{GeV}/c)^2$ is applied to remove the region where the fraction of non-exclusive background is large. The interaction vertex is required to be inside the polarised target material and the extrapolated track of the beam muon must traverse the full length of the target, in order to ensure equal flux in all target cells. The ρ^0 -meson is detected via its decay into a $\pi^+\pi^-$ pair, with invariant mass $0.5 \text{ GeV}/c^2 < M_{\pi\pi} < 1.1 \text{ GeV}/c^2$. Only events with an incident muon track, a scattered muon track and two additional tracks with opposite charge are considered. Suppression of semi-inclusive deep-inelastic (SIDIS) production is achieved by a cut on the energy of the ρ^0 -meson and on the missing energy $E_{\text{miss}} = (M_X^2 - M_P^2)/(2M_P)$, where M_X is the invariant mass of the undetected system and M_P is the proton mass: $E_{\rho^0} > 15 \text{ GeV}$ and $-2.5 \text{ GeV} < E_{\text{miss}} < 2.5 \text{ GeV}$. In order to suppress background of coherent production of exclusive ρ^0 production on nuclei of the target a lower cut on the transverse momentum p_T of the ρ^0 -meson is applied: $0.1 (\text{GeV}/c)^2 < p_T^2 < 0.5 (\text{GeV}/c)^2$ for ${}^6\text{LiD}$ and $0.05 (\text{GeV}/c)^2 < p_T^2 < 0.5 (\text{GeV}/c)^2$ for NH_3 , where the upper cut on p_T^2 is applied to remove the region, where non-exclusive production dominates.

After all cuts the final samples of incoherent exclusive ρ^0 production consist of about 97000 and 797000 events for the ${}^6\text{LiD}$ and NH_3 target, respectively. Still, about 20% of the events originate from SIDIS production. In order to correct for this contribution fits to the E_{miss} distribution are performed in each bin required for the asymmetry extraction. This means in bins in x_{Bj} , Q^2 , or p_T^2 per target cell, and also in $\phi - \phi_S$ and according to the target spin orientation. For the signal a Gaussian distribution is fitted. The shape for the background is obtained from a parameterisation of SIDIS Monte Carlo (MC) data generated with LEPTO binned in the same way as the real data, except in polarisation state and $\phi - \phi_S$, since no

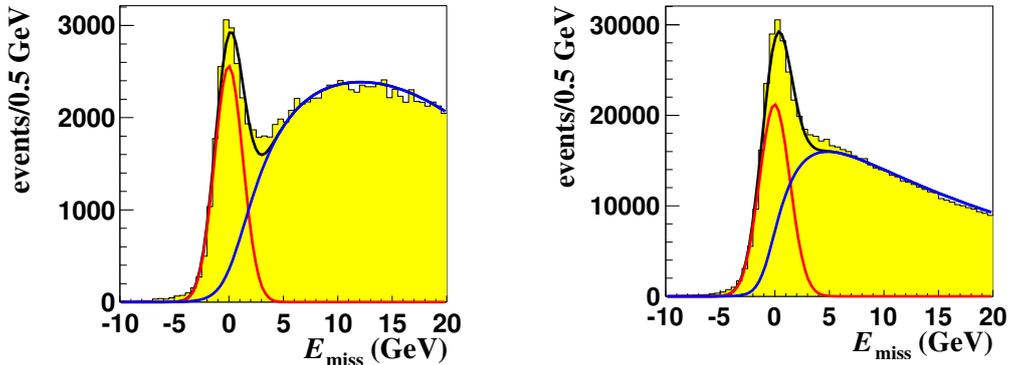


Figure 1: Signal plus background fit to the E_{miss} distribution for the ${}^6\text{LiD}$ data (left) and NH_3 data (right) in the region $2.4(\text{GeV}/c)^2 < Q^2 < 10(\text{GeV}/c)^2$ where the SIDIS background is the largest. The different shapes of the two distributions for large E_{miss} is explained by the increased angular acceptance of the COMPASS detector due to the new target magnet installed in 2005.

polarisation effects are simulated. The shape obtained from MC is fixed and only the normalisation is fitted to the data. This is exemplarily shown in Fig. 1 for both the ${}^6\text{LiD}$ and NH_3 data in the region $2.4(\text{GeV}/c)^2 < Q^2 < 10(\text{GeV}/c)^2$ integrated over the target cells, polarisation state and angle $\phi - \phi_S$. The resolution in E_{miss} is not sufficient to resolve ρ^0 production with diffractive dissociation of the target nucleon. It is visible in Fig. 1 right of the exclusive peak as a slight enhancement over the SIDIS background. The amount was estimated using Monte Carlo generated with HEPGEN [6] to be $\approx 14\%$. No attempt was made to remove this background. This is motivated by HERA results for unpolarised protons, which demonstrated that the angular distributions for proton-dissociative production are consistent with those of exclusive ρ^0 production [7, 8, 9]. The asymmetry is extracted from the numbers of exclusive events, after subtraction of SIDIS background, using a binned maximum likelihood fit in 12 bins of $\phi - \phi_S$.

3 Results

The results for the transverse target spin asymmetries $A_{UT}^{\sin(\phi - \phi_S)}$ measured on proton and deuteron as a function of x_{Bj} , Q^2 or p_T^2 , are shown in Fig. 2. For both targets the asymmetries are found to be small and compatible with zero within statistical uncertainties. For transversely polarised deuterons it is the first measurement. The proton results are compatible with the results measured by HERMES [10]. Note that the COMPASS proton results are more precise by a factor of about 3 and cover a larger kinematic domain. The results are compared with predictions of the GPD model by Goloskokov and Kroll [11], taking into account only contributions from valence quark GPDs E^u and E^d . A reasonable agreement is achieved. The small value of the asymmetry can be explained by an approximate cancellation of comparable contributions of opposite signs from GPDs E^u and E^d : $E^u \approx -E^d$. For proton the asymmetry is sensitive to $2/3E^u + 1/3E^d$, while for deuteron it depends on $E^u + E^d$.

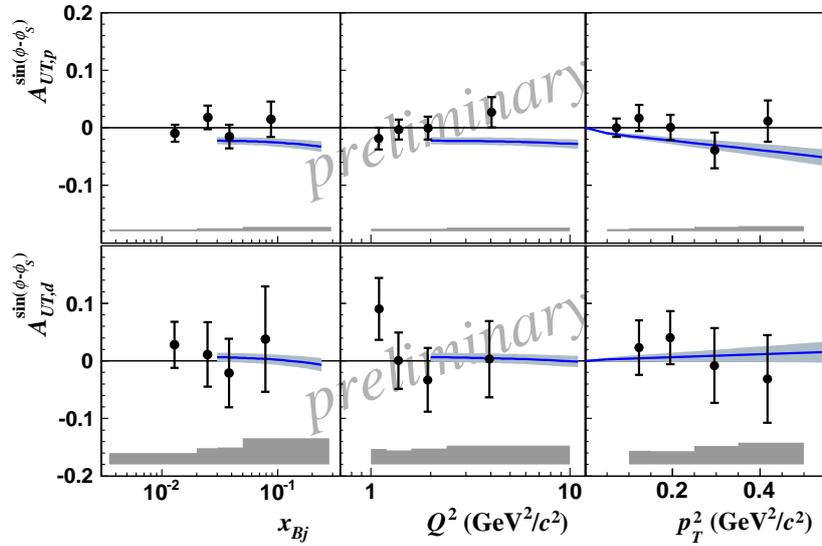


Figure 2: Results for $A_{UT}^{\sin(\phi-\phi_S)}$ measured on proton (upper) and deuteron (lower) as a function of x_{Bj} , Q^2 and p_T^2 . The systematic uncertainties are indicated by grey bands. The curves show the predictions of the GPD model [11] using the set of parameters called ‘variant 1’. The theoretical error bands take into account uncertainties of GPD parameterisations.

References

- [1] D. Mueller, D. Robaschik, B. Geyer, F. Dittes, and J. Horejsi. *Fortsch.Phys.* **42** (1994) 101, [arXiv:hep-ph/9812448](#) [hep-ph].
- [2] A. Radyushkin. *Phys.Lett.* **B385** (1996) 333–342, [arXiv:hep-ph/9605431](#) [hep-ph].
- [3] X.-D. Ji. *Phys.Rev.Lett.* **78** (1997) 610–613, [arXiv:hep-ph/9603249](#) [hep-ph].
- [4] M. Burkardt. *Phys.Rev.* **D62** (2000) 071503, erratum *ibid.* **D66** (2002) 119903, [arXiv:hep-ph/0005108](#) [hep-ph].
- [5] P. Abbon *et al.* *Nucl. Instrum. Meth.* **A577** (2007) 455–518, [arXiv:hep-ex/0703049](#).
- [6] A. Sandacz.
<http://project-gpd-full-chain-mc.web.cern.ch/project-gpd-full-chain-mc/hepgen>.
- [7] J. Breitweg *et al.* *Eur.Phys.J.* **C12** (2000) 393–410, [arXiv:hep-ex/9908026](#) [hep-ex].
- [8] S. Chekanov *et al.* *PMC Phys.* **A1** (2007) 6, [arXiv:0708.1478](#) [hep-ex].
- [9] F. Aaron *et al.* *JHEP* **1005** (2010) 032, [arXiv:0910.5831](#) [hep-ex].
- [10] A. Rostomyan and J. Dreschler. [arXiv:0707.2486](#) [hep-ex].
- [11] S. Goloskokov and P. Kroll. *Eur.Phys.J.* **C59** (2009) 809–819, [arXiv:0809.4126](#) [hep-ph].