

QCD fits in diffractive DIS revisited

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Diffraction parton distributions

- Diffractive PDF's are essential ingredients for the description of hard diffractive processes
- DPDF's **evolution** is known (Trentadue, Veneziano '94)
- DPDF's **factorization** has been demonstrated in DDIS (Grazzini & al. '98, Collins '98)
- Once extracted from DDIS data they allow factorization tests in
 - diffractive dijet production in DIS
 - diffractive dijet photoproduction
 - single and double diffractive hard processes in hadron-hadron collisions
- A **more flexible** and **less model dependent** pQCD-approach has been designed to exploit the improved quality of forthcoming data
- Results of a benchmark study on H1 LRG06 data (EPJ 2006)

Data and Observable

- Data set : H1 LRG06 (diffractive selection by large rapidity gap method)
- Kinematic variables

$$\beta = \frac{Q^2}{Q^2 + M_X^2}; \quad x_P = \frac{x_B}{\beta}; \quad y = \frac{Q^2}{sx_B}$$

- Kinematic coverage @ $\sqrt{s}=318$ GeV :

$$3.5 \leq Q^2 \leq 1600 \text{ GeV}^2 \quad 3 \cdot 10^{-4} \leq x_P \leq 3 \cdot 10^{-2} \quad 10^{-3} \leq \beta \leq 0.8$$

- Observable:

$$\sigma_r^{D(3)}(\beta, Q^2, x_P) = F_2^{D(3)}(\beta, Q^2, x_P) - \frac{y^2}{1 + (1 - y)^2} F_L^{D(3)}(\beta, Q^2, x_P)$$

- Measurements integrated over $|t| \leq 1 \text{ GeV}^2$, $M_Y < 1.6 \text{ GeV} \rightarrow$ [DPDF's definition](#)

Theory

- Evolution and convolution performed with QCDNUM17 package (M.Botje 2011)
- Fixed Flavour Number Scheme with $n_f = 3$, evolution @NLL
- Heavy Quarks in G(eneral) M(assive) scheme: NLO massive coefficients from HQSTFUN
- Coupling @NLO : $\alpha_s^{(n_f=3)}(M_Z^2) = 0.107 \leftrightarrow \alpha_s^{(n_f=5)}(M_Z^2) = 0.118$
- $m_c = 1.4 \text{ GeV}$, $m_b = 4.5 \text{ GeV}$
- $\mu_F^2 = \mu_R^2 = Q^2$

The new method

- Widely used Regge-inspired approach + vertex factorisation:

$$F_i^D(\beta, Q^2, x_{\mathbb{P}}, t) = f_{\mathbb{P}/P}(x_{\mathbb{P}}, t) F_i^{\mathbb{P}}(\beta, Q^2) + f_{\mathbb{R}/P}(x_{\mathbb{P}}, t) F_i^{\mathbb{R}}(\beta, Q^2) + \dots$$

- Factorization theorem holds **at fixed** $x_{\mathbb{P}}$ and t
- DPDF's parametrize the parton content of the incoming proton under the condition that it is quasi elastically scattered with a given $x_{\mathbb{P}}$ and t
- Basic new assumption : for each $x_{\mathbb{P}}$ one has (in principle) **different** DPDF's
- Perform QCD-fits at **fixed** $x_{\mathbb{P}}$ with a common initial condition at Q_0^2

$$\Sigma(\beta, Q_0^2) = A_q \beta^{B_q} (1 - \beta)^{C_q} e^{\frac{-0.01}{1-\beta}}$$

$$g(\beta, Q_0^2) = A_g e^{\frac{-0.01}{1-\beta}}$$

- **4 free parameters**, inspired by H1 FitB parametrization

Fixed x_P -fits: results

- Same cuts as in H1 LRG '06 paper
- $M_X^2 \geq 4 \text{ GeV}^2$, $Q_{\min}^2 \geq 8.5 \text{ GeV}^2$, 190 points after cuts (out of 276)

x_P	χ^2	# Fitted points after cuts	# Total points
0.0003	-	3	12
0.0010	1.23	19	40
0.0030	1.01	39	59
0.0100	1.08	61	80
0.0300	0.63	68	85

- Almost good overall description despite only statistical errors are used

$x_{\mathbb{P}}$ -dependence of input distributions

- A generalisation of the initial condition to include the $x_{\mathbb{P}}$ -dependence is needed
- $x_{\mathbb{P}}$ -dependence fully contained in parameters

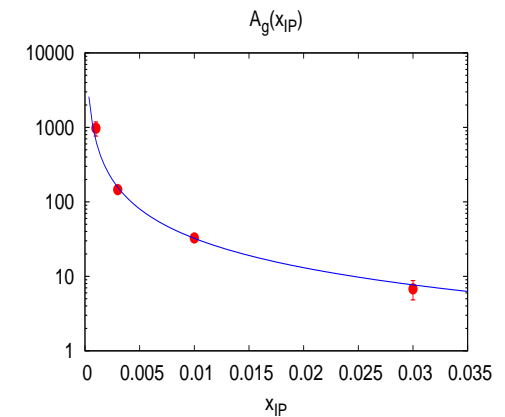
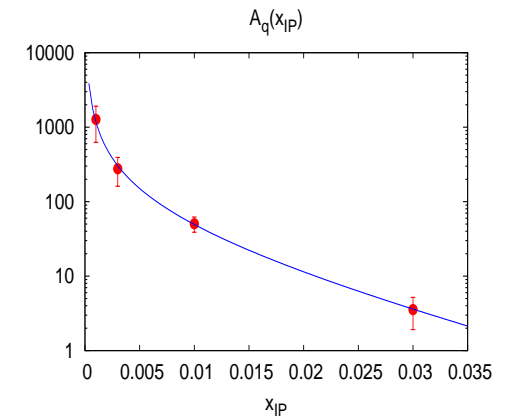
$$\Sigma(\beta, x_{\mathbb{P}}, Q_0^2) = A_q(x_{\mathbb{P}}) \beta^{B_q(x_{\mathbb{P}})} (1 - \beta)^{C_q(x_{\mathbb{P}})} e^{\frac{-0.01}{1-\beta}}$$

$$g(\beta, x_{\mathbb{P}}, Q_0^2) = A_g(x_{\mathbb{P}}) e^{\frac{-0.01}{1-\beta}}$$

- $x_{\mathbb{P}}$ -dependence is however controlled by **non-perturbative dynamics**
- ..get **inspiration** from fixed $x_{\mathbb{P}}$ -fits

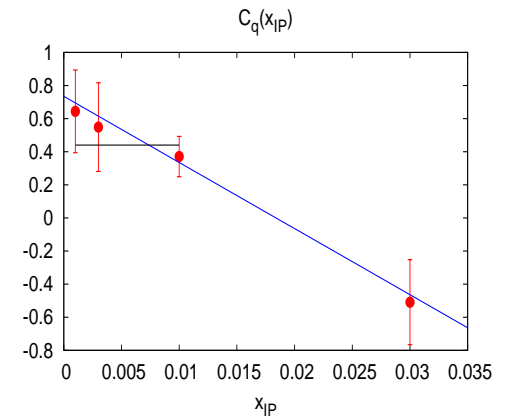
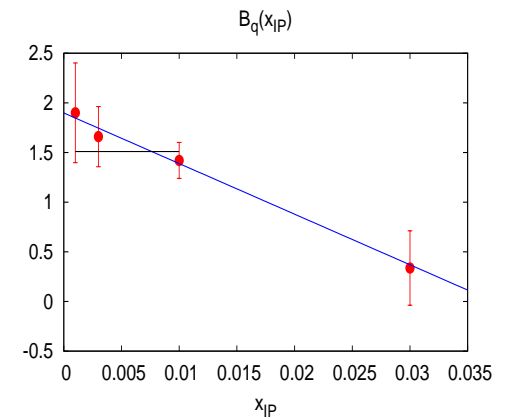
$x_{\mathbb{P}}$ -dependence of input distributions

- The NP dependence of DPDF parameters on $x_{\mathbb{P}}$ is inferred by fixed $x_{\mathbb{P}}$ -fits
- The singlet and gluon normalizations are parametrized as follows:
 - $A_q(x_{\mathbb{P}}) = A_{q,0} \cdot (x_{\mathbb{P}})^{A_{q,1}} \cdot (1 - x_{\mathbb{P}})^{A_{q,2}}$
 - $A_g(x_{\mathbb{P}}) = A_{g,0} \cdot (x_{\mathbb{P}})^{A_{g,1}}$
 - **Distinct** powers for singlet and gluons (different w.r.t. to fluxes)



$x_{\mathbb{P}}$ -dependence of input distributions

- The NP dependence of DPDF parameters on $x_{\mathbb{P}}$ is inferred by fixed $x_{\mathbb{P}}$ -fits
- The singlet shape-modulation parameters are parametrized as follows:
 - $B_q(x_{\mathbb{P}}) = B_{q,0} + B_{q,1} \cdot (x_{\mathbb{P}})$
 - $C_q(x_{\mathbb{P}}) = C_{q,0} + C_{q,1} \cdot (x_{\mathbb{P}})$
 - Black line : Regge inspired (pomeron only)
 - H1 FPS and Tevatron go up to 0.1 in $x_{\mathbb{P}}$...



Combined- $x_{\mathbb{P}}$ fit : results (1)

- $M_X^2 > 4 \text{ GeV}^2$, $Q_{\min}^2 \geq 8.5 \text{ GeV}^2$, 190 points
- In combined- $x_{\mathbb{P}}$ fit, systematics errors taken into account (Pascaud and Zomer LAL-95-05)
- Systematics parameters (21) are fitted along with theory parameters (9)

$x_{\mathbb{P}}$	partial χ^2 in the combined fit	χ^2 for fixed $x_{\mathbb{P}}$ fits	Points after cuts
0.0003	3.1	-	3
0.0010	26.2	18.4	19
0.0030	38.9	35.4	39
0.0100	58.6	61.8	61
0.0300	36.4	40.2	68

- **Good description** in most of the bins

Combined- $x_{\mathbb{P}}$ fit : results (2)

- Best $\chi^2 = 165.8$ for 9 free parameters, $\chi^2/d.o.f. = 0.92$
- $x_{\mathbb{P}}$ -combination costs $\Delta\chi^2 = +5$ units
- Comparable quality with respect to H1 published pQCD fits
- Same power controlling the $x_{\mathbb{P}}$ -dependence of singlet and gluon normalization:

$$A_{q,1} = A_{g,1} \rightarrow \chi^2=171.0$$

- $A_{q,1} = A_{g,1}$ & no shaping for the singlet:

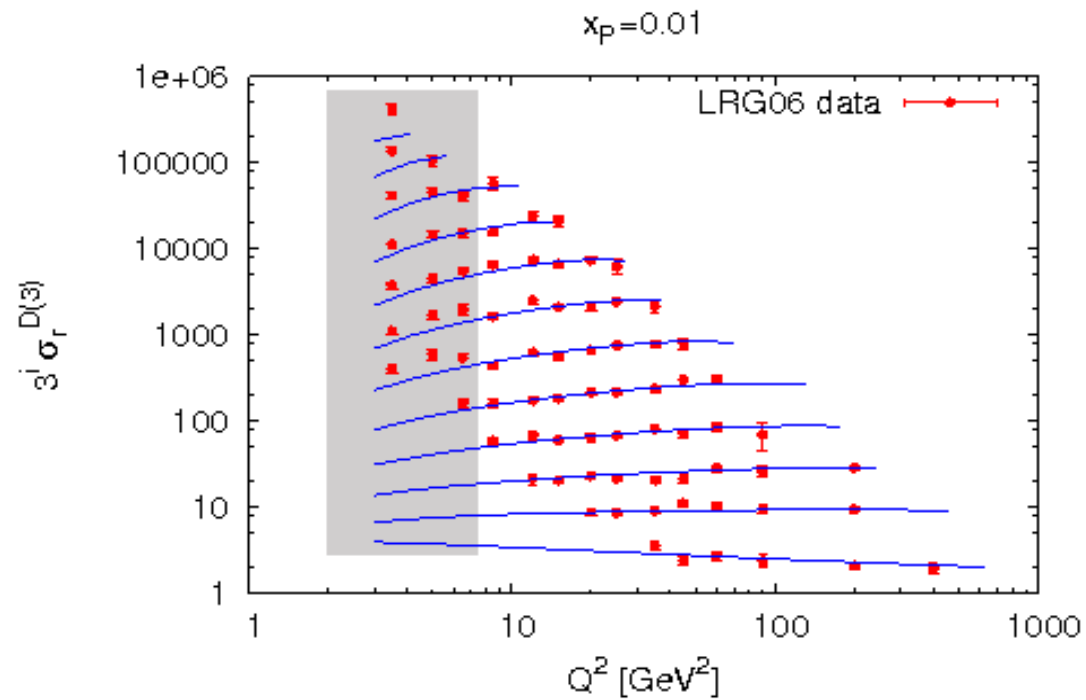
$$B_{q,1} = C_{q,1} = 0 \rightarrow \chi^2=187.6$$

Study of the stability of the fit

- Restrict the phase space to $Q_{\min}^2 \geq 8.5 \text{ GeV}^2$ and $y < 0.5$
- ▶ Both χ^2 's and parameters unchanged :
 - consistent. Same conclusion as the recent F_L^D analysis by H1 (EPJ 2012)
- Relax the cut on Q_{\min}^2 and keep $0 < y < 1$
- ▶ χ^2 's deteriorate and parameters are strongly affected
- Same results with no cut on Q_{\min}^2 and $y < 0.5$
- Situation does not improve even with more flexible inputs (!)
 - consistent with H1 analysis (EPJ 2006)

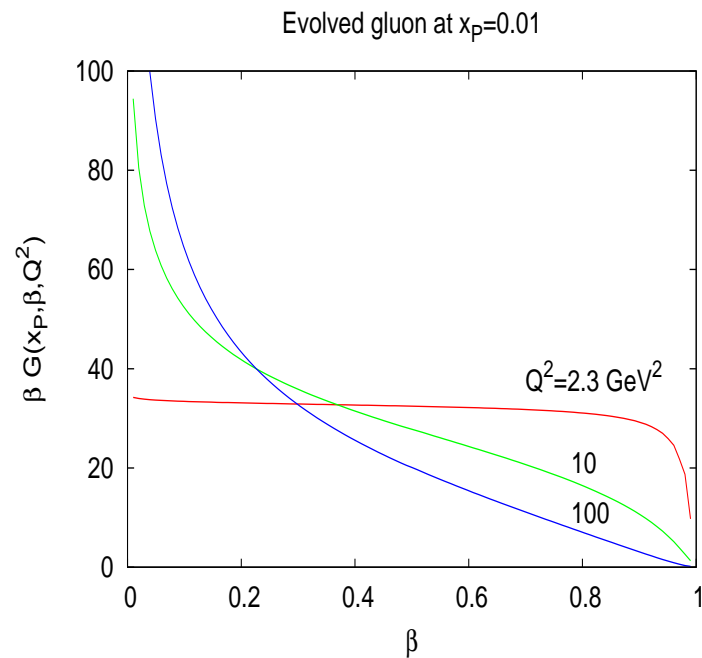
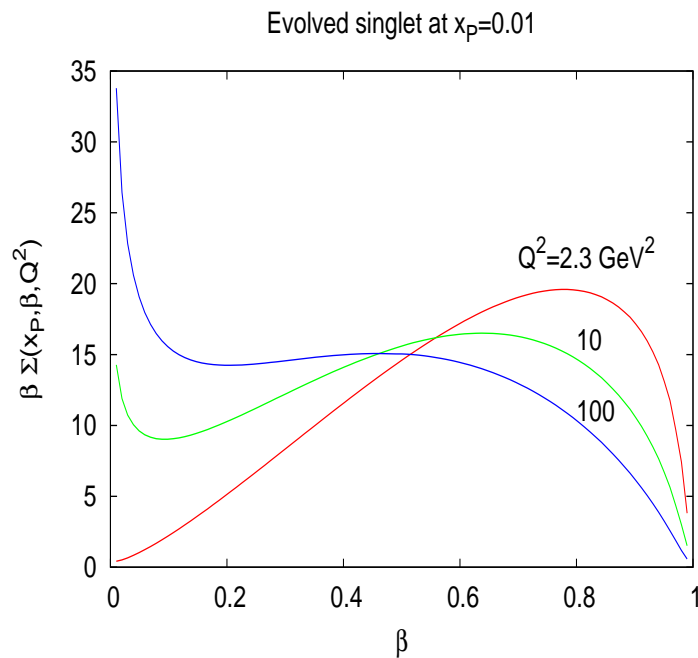
Reduced cross-sections at $x_P = 0.01$

- The grey area indicates the cut $Q_{\min}^2 \geq 8.5 \text{ GeV}^2$



Singlet and gluon Q^2 -evolution at $x_P = 0.01$

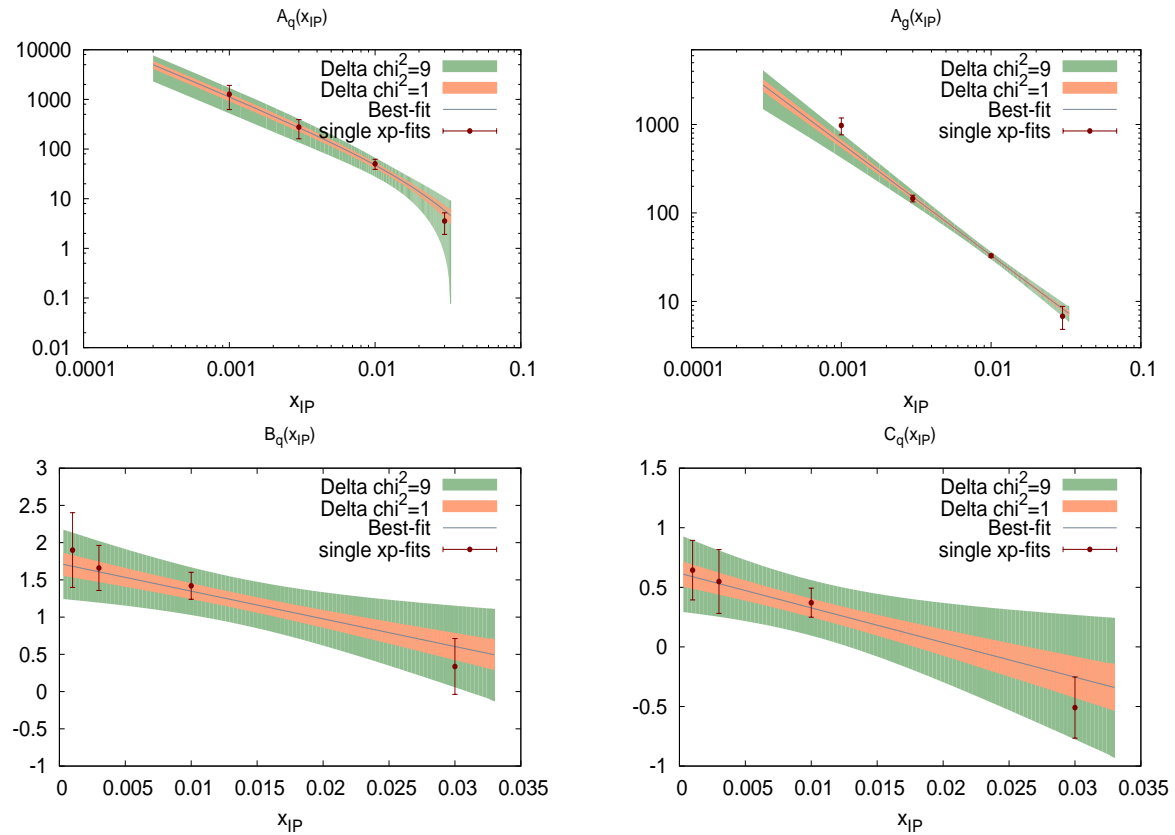
- Evolution at high Q^2 washes away any structure at high β



- However the initial condition at $Q_0^2 = 2.3 \text{ GeV}^2$ is far in the unmeasured region

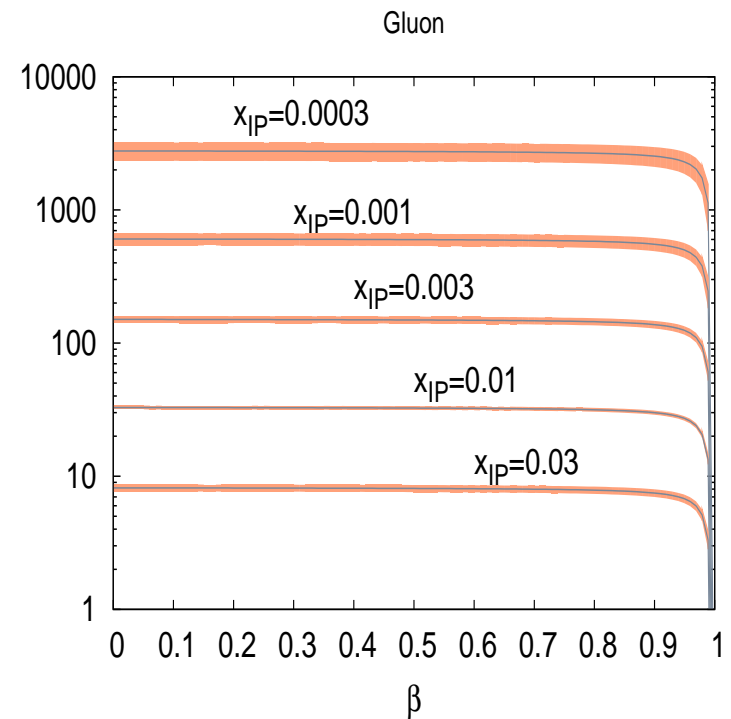
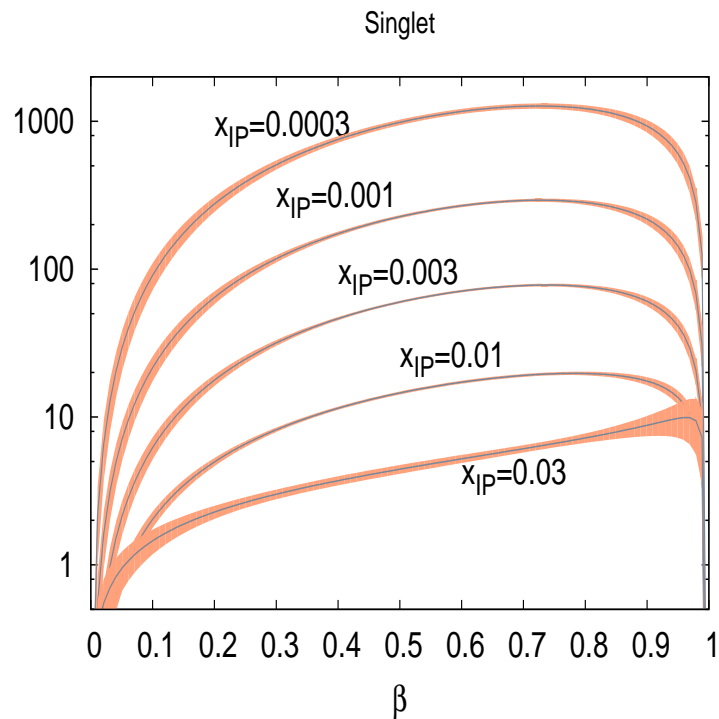
Parameters check

- Comparison of parameters from fixed- x_{IP} fits and combined-fit



DPDF's vs β at $Q_0^2 = 2.3 \text{ GeV}^2$ for different x_{IP}

- The singlet slightly evolves with x_{IP} (initial conjecture)



Conclusions

- The new method offers flexibility in the choice of the input distributions : **surface** in $(\beta, x_{\mathbb{P}})$ properly interfaced with pQCD
- Singlet β -shape **slightly evolves** in $x_{\mathbb{P}}$
- $x_{\mathbb{P}}$ -combined fit is possible (important for $x_{\mathbb{P}}$ -integrated cross-sections, *i.e.* jet CS, portability of the code..)
- Quality of the fit comparable with respect to pQCD+Regge-inspired fit by H1 on the same data
- The approach can be extended to treat the t -dependence of the cross-section
- The precision of forthcoming data will further tests these ideas