

Testing NRQCD factorization with J/ψ yield and polarization

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In collaboration with Mathias Butenschön

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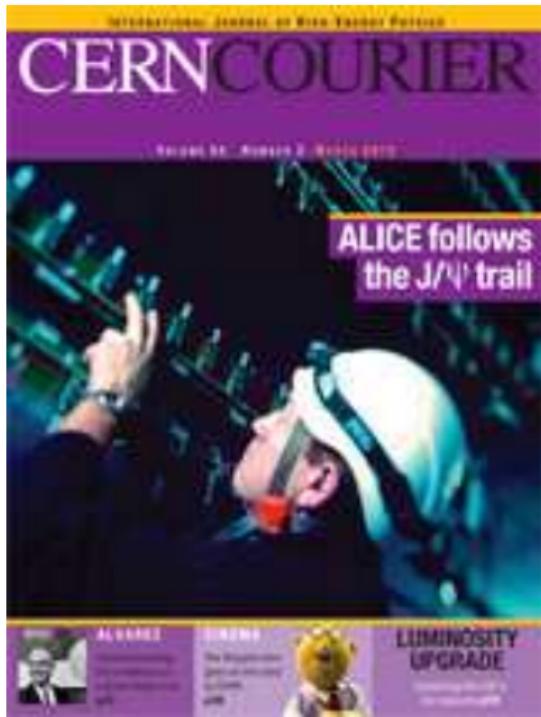
Phys. Rev. Lett. **106**, 022003 (2011)

Phys. Rev. D **84**, 051501(R) (2011)

Phys. Rev. Lett. **107**, 232001 (2011)

arXiv:1201.1872 [hep-ph], Phys. Rev. Lett. (in press)

CERN Courier, Volume 52, Issues 1 and 2



Outline

- 1 **Introduction:** Heavy quarkonia, NRQCD, J/ψ production
 - 2 **Technology:** Treatment of divergences, automation
 - 3 **Global fit:** Unpolarized J/ψ production
 - 4 **Polarized photoproduction:** HERA
 - 5 **Polarized hadroproduction:** Tevatron vs. LHC
 - 6 **Summary:** NRQCD at the crossroads

Heavy Quarkonia

Heavy quarkonia: Bound states of heavy quark and its antiquark.

- Charmonia ($c\bar{c}$) and Bottomonia ($b\bar{b}$)
 - Top decays too fast for bound state.

Charmonium spectrum ($c\bar{c}$):

$n^{2S+1}L_J$	Name	Mass
1^1S_0	η_c	2980 MeV
1^3S_1	J/ψ	3097 MeV
1^3P_0	χ_{c0}	3415 MeV
1^3P_1	χ_{c1}	3511 MeV
1^1P_1	h_c	3526 MeV
1^3P_2	χ_{c2}	3556 MeV
2^1S_0	η'_c	3637 MeV
2^3S_1	ψ'	3686 MeV

- 1974: Discovery of J/ψ :
First observation of heavy quarks
 - Long lifetime of $c\bar{c}$: Spectrum
and radiative transitions seen
 \implies Potential models
 - Calculation of energy spectrum:
Challenge for lattice QCD.
 - Production and decay rates:
One of first applications for
perturbative QCD.

Color-singlet model vs. NRQCD factorization

Classic approach: Color-singlet model

- $c\bar{c}$ pair in physical color-singlet state, e.g. $c\bar{c}[{}^3S_1^{[1]}]$ for J/ψ .
 - Nonperturbative information in J/ψ wave function at origin.
 - Leftover IR divergences for P-wave quarkonia \rightsquigarrow inconsistent!
 - Predicted cross section factor 10^1 – 10^2 below Tevatron data.

NRQCD factorization:

- Rigorous effective field theory [Bodwin, Braaten, Lepage]
 - Based on **factorization of soft and hard scales**
(Scale hierarchy: $Mv^2, Mv \ll \Lambda_{\text{QCD}} \ll M$)
 - Theoretically consistent: no leftover singularities.
 - NNLO proof of factorization in fragmentation [Nayak, Qiu, Sterman]
 - Can explain hadroproduction at Tevatron.

NRQCD factorization in a nutshell

Factorization theorem: $\sigma_{J/\psi} = \sum_n \sigma_{c\bar{c}[n]} \cdot \langle O^{J/\psi}[n] \rangle$

- n : every possible Fock state, including color-octet states.
- $\sigma_{c\bar{c}[n]}$: production rate of $c\bar{c}[n]$, calculated in perturbative QCD.
- $\langle O^{J/\psi}[n] \rangle$: long-distance matrix elements (LDMEs),
nonperturbative, extracted from experiment, universal?

Scaling rules: LDMEs scale with relative velocity v ($v^2 \approx 0.2$).

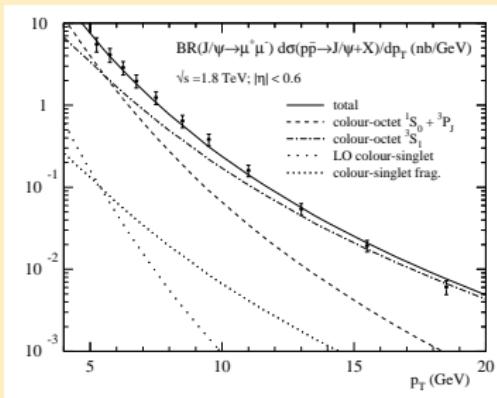
For J/ψ :

scaling	v^3	v^7	v^{11}
n	${}^3S_1^{[1]}$	${}^1S_0^{[8]}, {}^3S_1^{[8]}, {}^3P_{0/1/2}^{[8]}$	\dots

- Double expansion in v and α_s .
- Leading term in v ($n = {}^3S_1^{[1]}$) corresponds to color-singlet model.

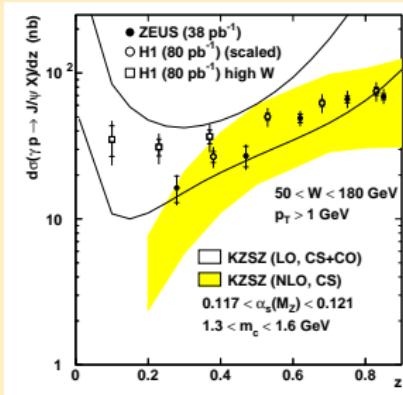
J/ ψ Production: NRQCD vs. Experiment

Hadroproduction at Tevatron:



- Color octet states important
⇒ Great success for NRQCD

Photoproduction at HERA:

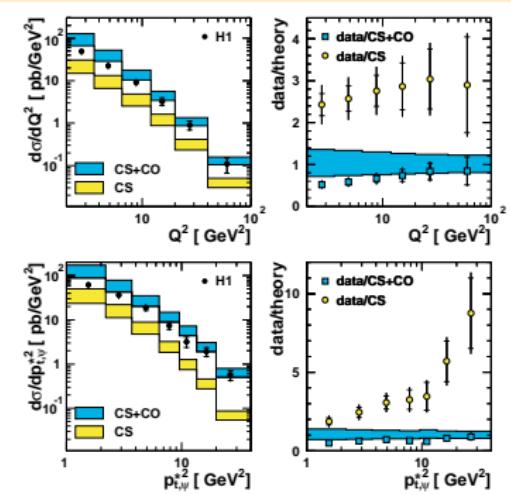


- MEs from fits to Tevatron data.
 - Importance of color octet unclear

Our work: NRQCD calculation for photo- and hadroproduction and e^+e^- annihilation at NLO
⇒ Establish universality of long distance matrix elements.

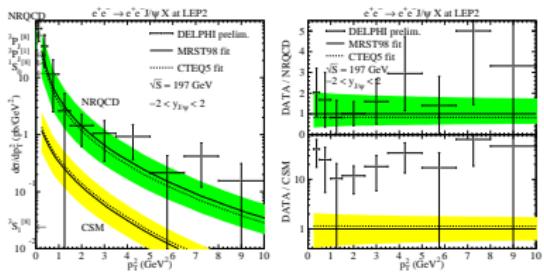
Production of J/ψ : NRQCD vs. Experiment (cont.)

Electroproduction at HERA:



BK, Zwirner, NPB(2002)

Two-Photon Collisions at LEP II:

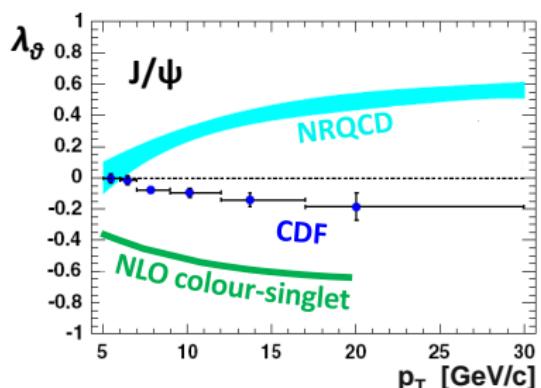


Klasen, BK, Mihaila, Steinhauser,
PRL(2002)

Evidence of color-octet mechanism at LO

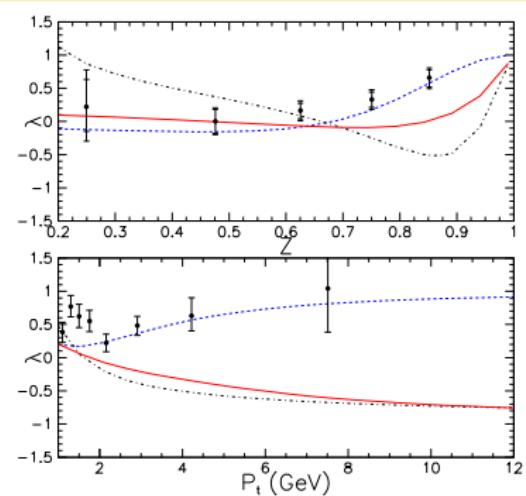
Production of J/ψ : NRQCD vs. Experiment (cont.)

Polarization at Tevatron:



Braaten, BK, Lee, PRD(2000);
Gong, Wang, PRL(2008)

Polarization at HERA:



Chang, Li, Wang, PRD(2009)

$$\frac{d\sigma}{d\cos\vartheta} \propto 1 + \lambda_\vartheta \cos^2\vartheta$$

Puzzling situation: LO NRQCD and NLO CSM fail
⇒ How about NLO NRQCD?

Organisation of our Calculation

FeynArts: Generate Feynman diagrams

Mathematica script: Apply color projectors. Evaluate color factors with FeynCalc.

FORM: Apply spin projectors. Treat squared amplitudes.

Two methods for virtual corrections:

FORM: Perform our tensor reduction
⇒ Scalar integrals

FORM: Use Integration by parts relations
⇒ Master integrals (uses AIR)

FORM: Cancel scalar products by denominators. Neg. propagator powers.
Directly apply IBP relations (uses AIR).
⇒ Master integrals
(Not for 1S_0 spin singlet state)

Mathematica script: Simplify results due to kinematics (Size dramatically reduced)

Analytical check: Results of two methods equal. All divergences cancel.

Numerical evaluation:

FORTRAN: Phase space integrations with VEGAS. Phase space slicing method.

Global fit at NLO in NRQCD

Fit	CO	LDMEs	to all available	world data	on J/ψ	inclusive production:
type	\sqrt{s}	collider	collaboration	reference		
pp	200 GeV	RHIC	PHENIX	PRD82(2010)012001		
$p\bar{p}$	1.8 TeV	Tevatron I	CDF	PRL97(1997)572; 578		
$p\bar{p}$	1.96 TeV	Tevatron II	CDF	PRD71(2005)032001		
pp	7 TeV	LHC	ALICE	NPB(PS)214(2011)56		
			ATLAS	PoS(ICHEP 2010)013		
			CMS	EPJC71(2011)1575		
			LHCb	EPJC71(2011)1645		
γp	300 GeV	HERA I	H1, ZEUS	EPJ25(2002)25; 27(2003)173		
γp	319 GeV	HERA II	H1	EPJ68(2010)401		
$\gamma\gamma$	197 GeV	LEP II	DELPHI	PLB565(2003)76		
e^+e^-	10.6 GeV	KEKB	BELLE	PRD79(2009)071101		

- Fit values:
 $10^{-2} \text{ GeV}^{3+2L}$

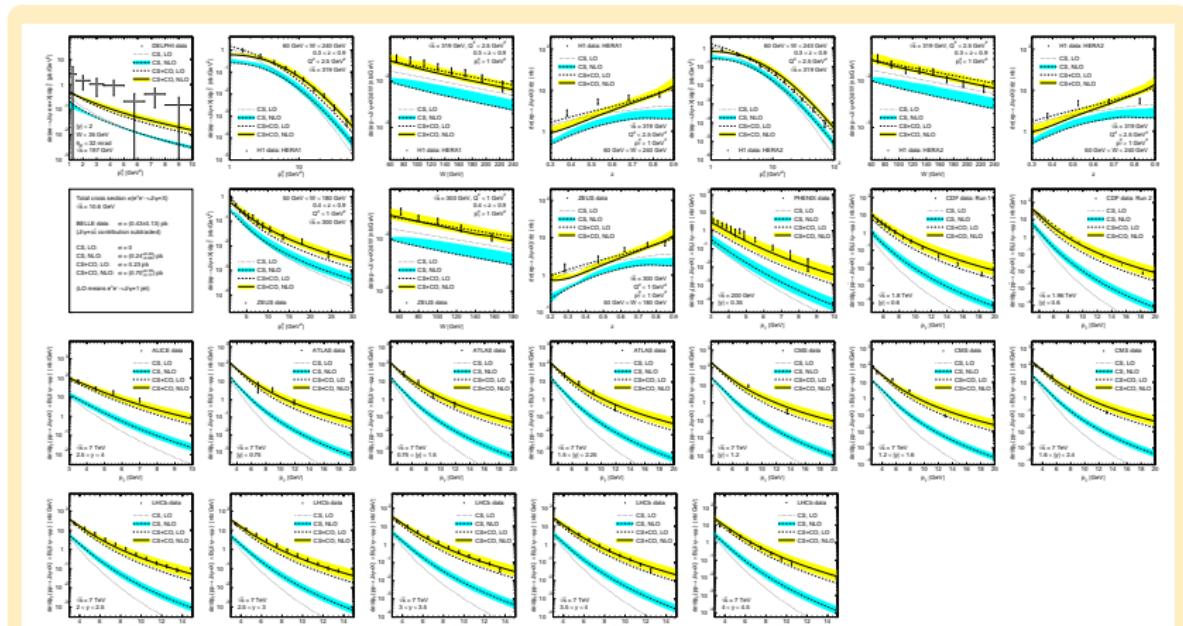
$$\langle \mathcal{O}(^1S_0^{[8]}) \rangle = 4.97 \pm 0.44$$

$$\langle \mathcal{O}({}^3S_1^{[8]}) \rangle = 0.224 \pm 0.059$$

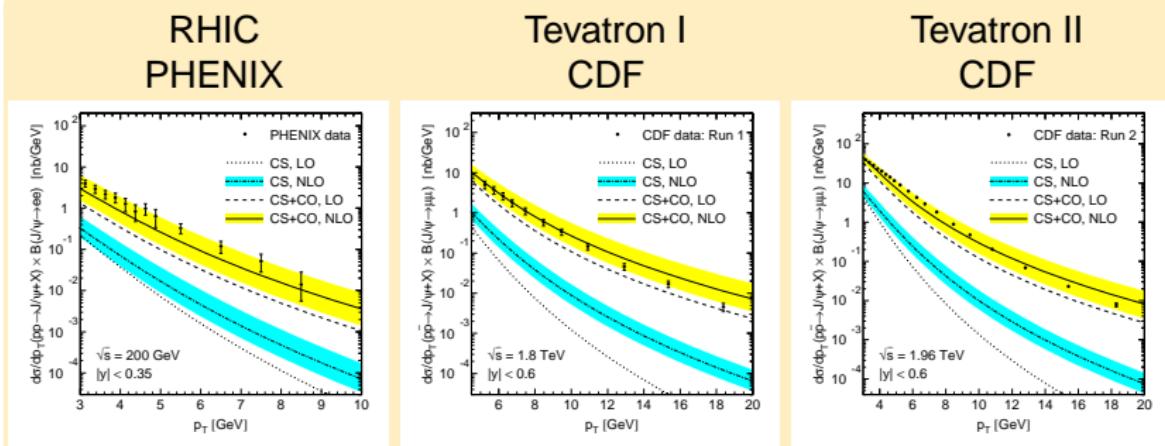
$$\langle \mathcal{O}(^3P_0^{[8]}) \rangle = -1.61 \pm 0.20$$

- $\chi^2/\text{d.o.f.} = 857/194 = 4.42$ for default prediction
 - $\propto v^4 \langle O_1(^3S_1) \rangle \sim \text{NRQCD velocity scaling rules}$ ✓

Comparison with world data



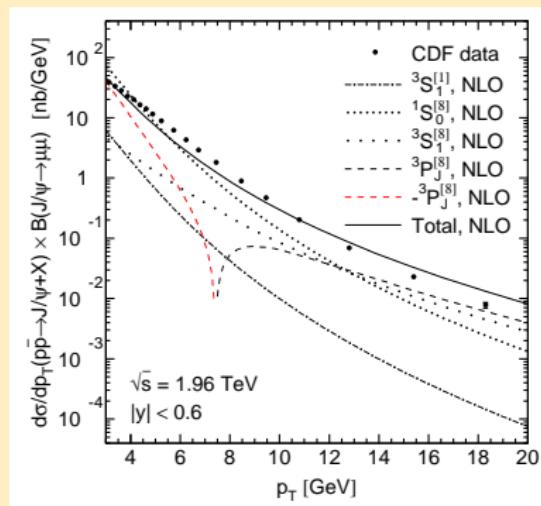
Comparison with RHIC and Tevatron



- Data well described by CS+CO at NLO.
 - CS orders of magnitudes below data.

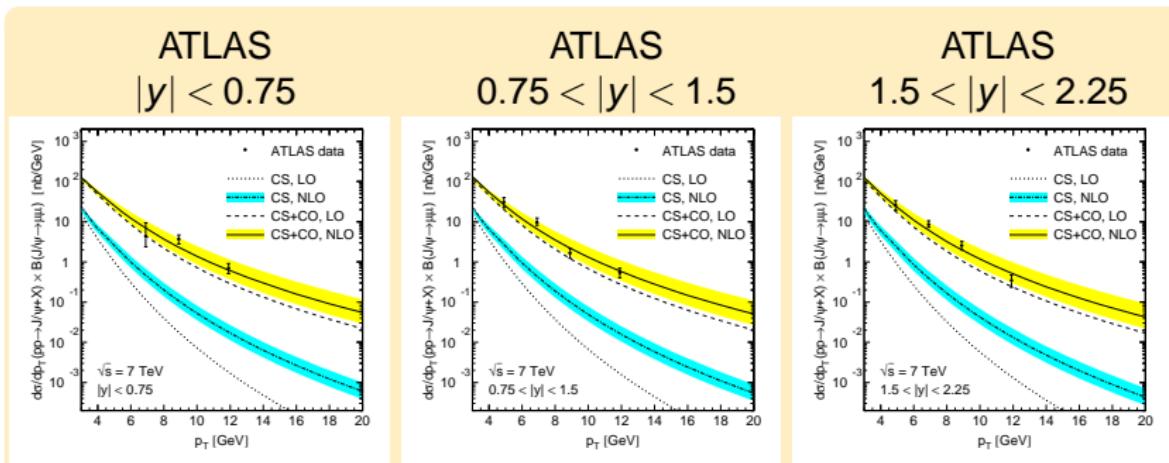
Comparison with Tevatron (cont.)

Relative importance of CO processes:



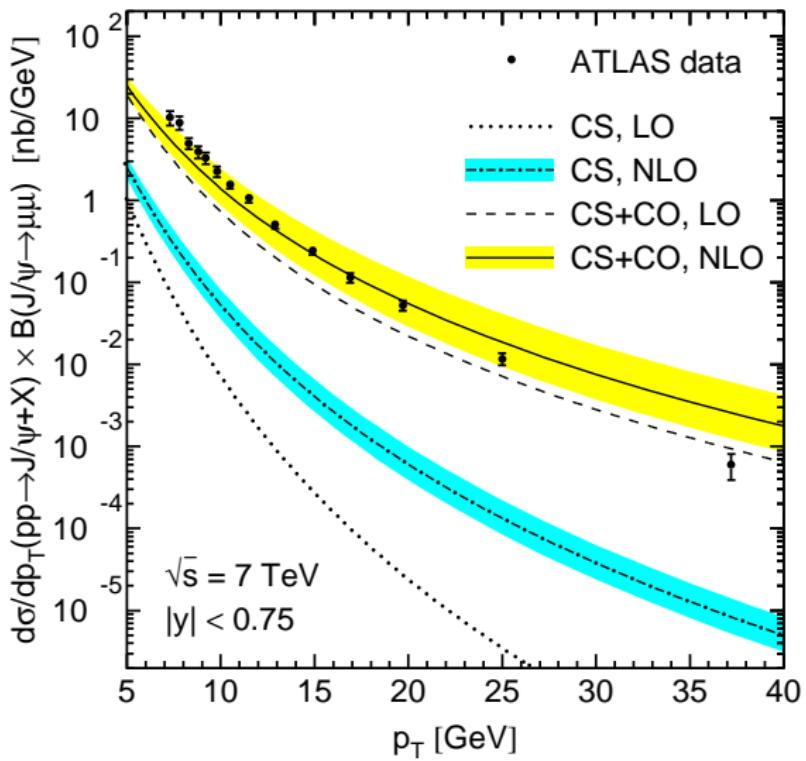
- Short-distance $\sigma(c\bar{c}[{}^3P_J^{[8]}]) < 0$ for $p_T \gtrsim 7$ GeV.
- But: Short-distance cross sections and LDMEs unphysical (NRQCD scale and scheme dependence) \leadsto No problem!

Comparison with ATLAS at LHC

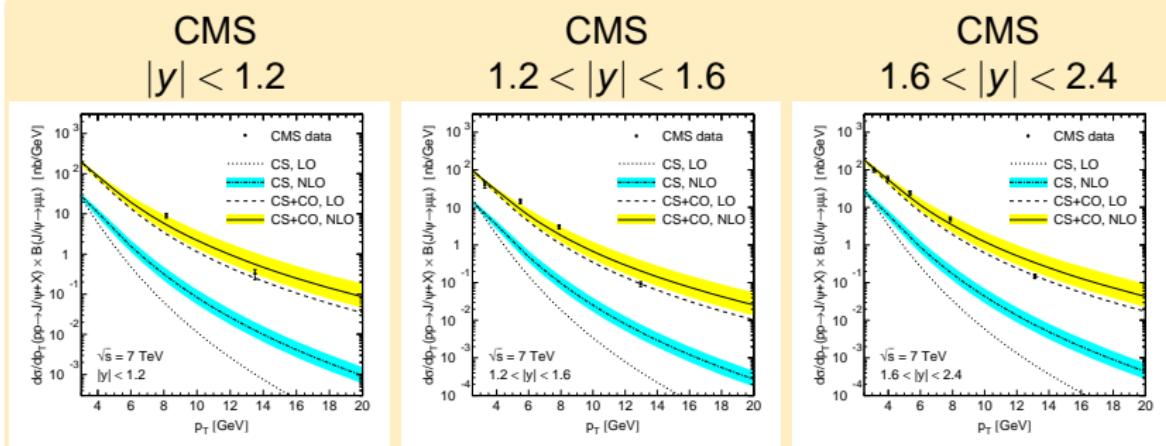


- Data well described by CS+CO at NLO.
 - CS orders of magnitudes below data.

Comparison with ATLAS (after fit) NPB850(2011)387

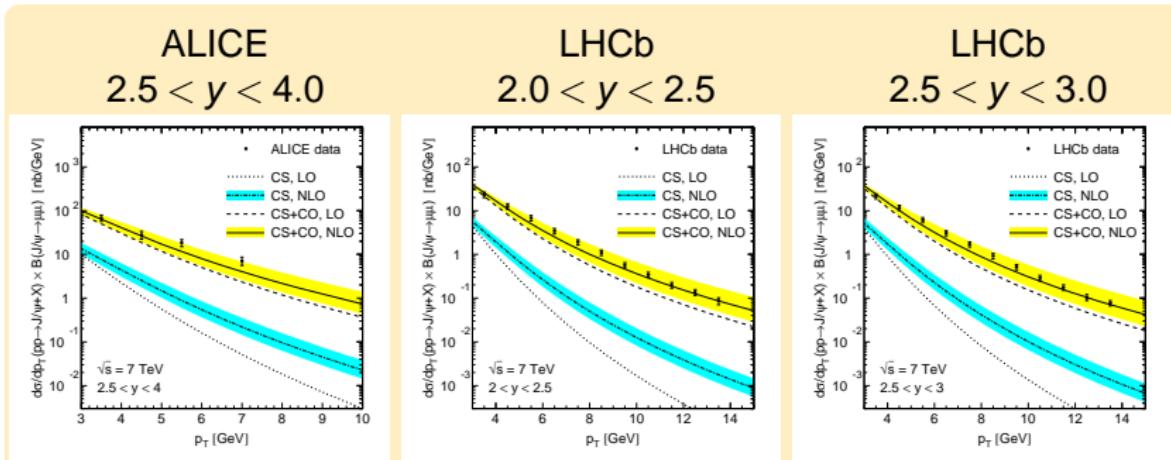


Comparison with CMS at LHC



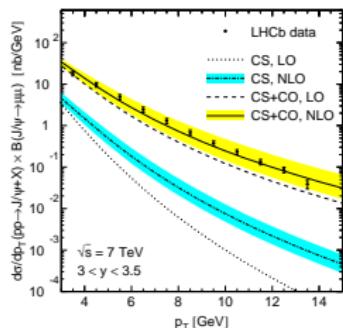
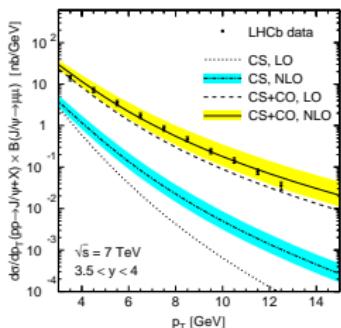
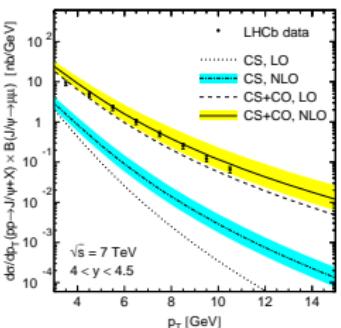
- Data well described by CS+CO at NLO.
 - CS orders of magnitudes below data.

Comparison with ALICE and LHBb at LHC



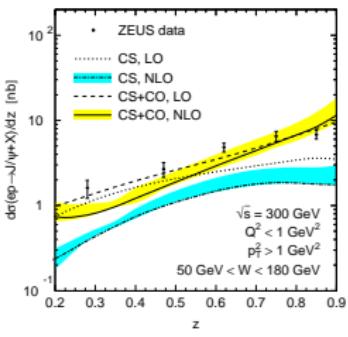
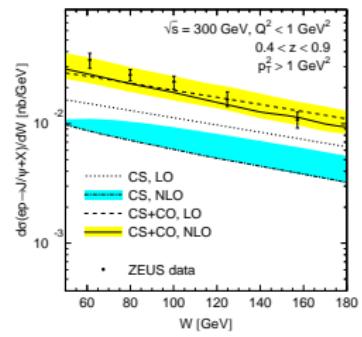
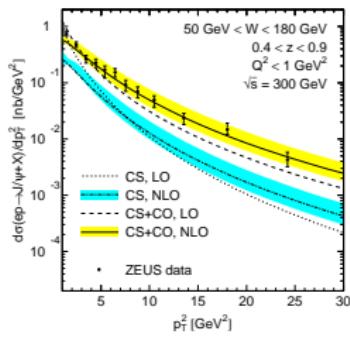
- Data well described by CS+CO at NLO.
 - CS orders of magnitudes below data.

Comparison with LHBb at LHC (cont.)

LHCb $3.0 < y < 3.5$ **LHCb** $3.5 < y < 4.0$ **LHCb** $4.0 < y < 4.5$ 

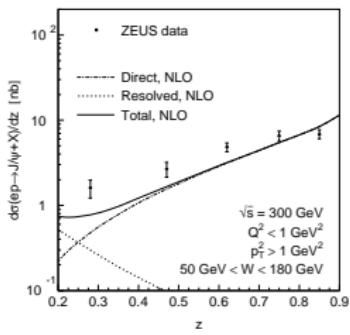
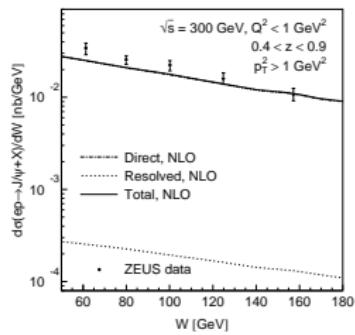
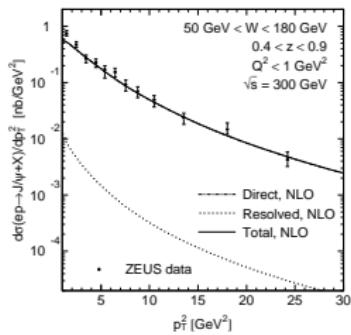
- Data **well described** by CS+CO at NLO.
- CS orders of magnitudes **below** data.

Comparison with ZEUS at HERA I (1)



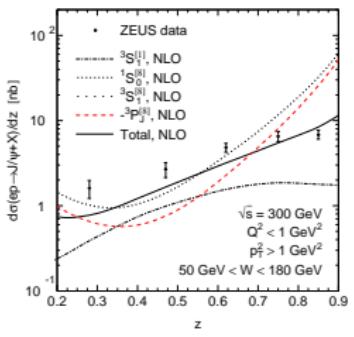
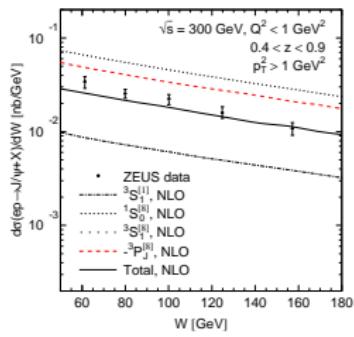
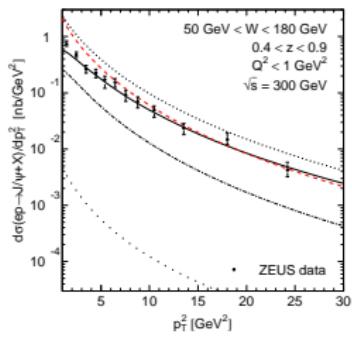
- $W = \gamma p$ CM energy.
 - $z =$ fraction of γ energy going to J/ψ in p rest frame.
 - Singularity for $z \rightarrow 1$ eliminated by shape function in SCET.
 - Data well described by CS+CO at NLO.
 - CS factor of 3–5 below the data.

Comparison with ZEUS at HERA I (2)



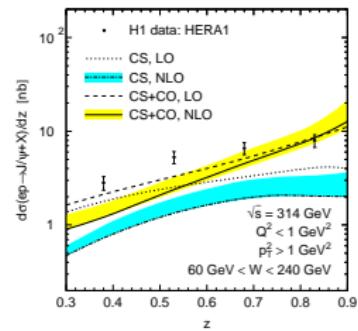
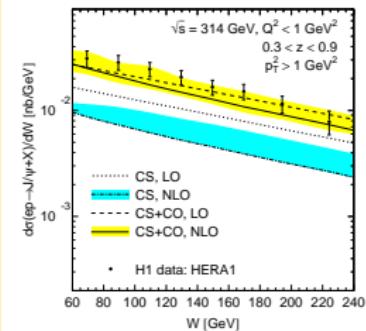
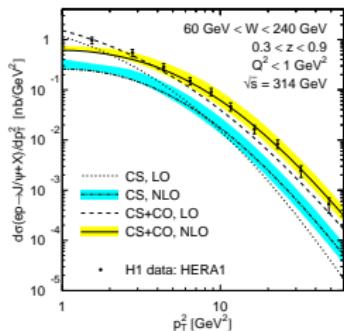
- Data for $0.4 < z < 0.9$ exhausted by direct photoproduction.
 - Resolved photoproduction only relevant for $z \lesssim 0.4$.

Comparison with ZEUS at HERA I (3)



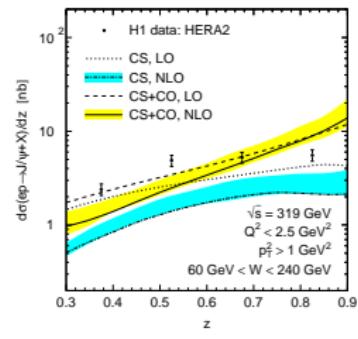
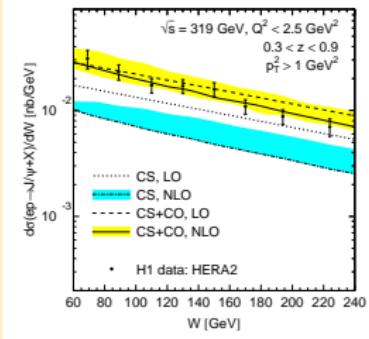
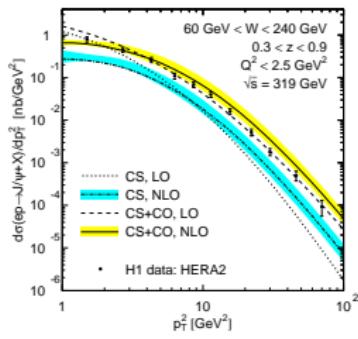
- $\langle \mathcal{O}(^3P_0^{[8]}) \rangle < 0 \rightsquigarrow ^3P_0^{[8]} \text{ contribution negative.}$
 - Negative interference with ${}^1S_0^{[8]}$ contribution beneficial.
 - ${}^3S_1^{[8]}$ contribution negligible here.

Comparison with H1 at HERA I



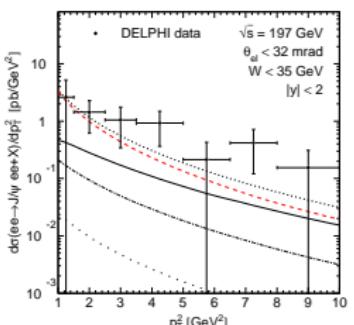
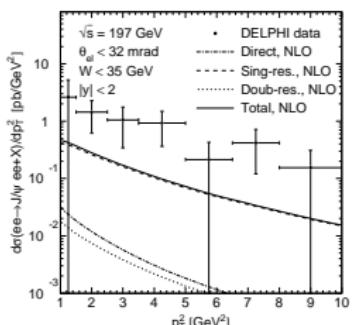
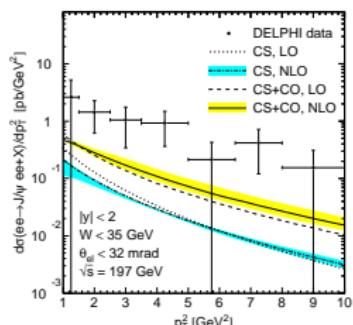
- Data well described by CS+CO at NLO.
 - CS factor of 3–5 below data.

Comparison with H1 at HERA II



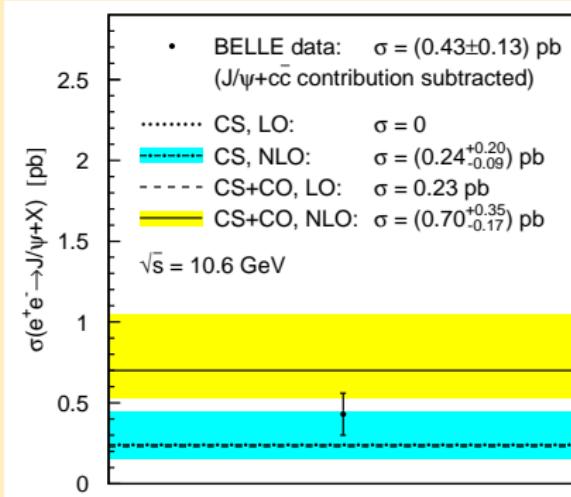
- Data well described by CS+CO at NLO.
 - CS factor of 3–5 below the data.

Comparison with DELPHI at LEP II



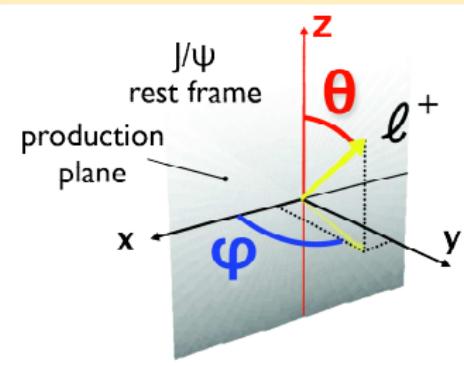
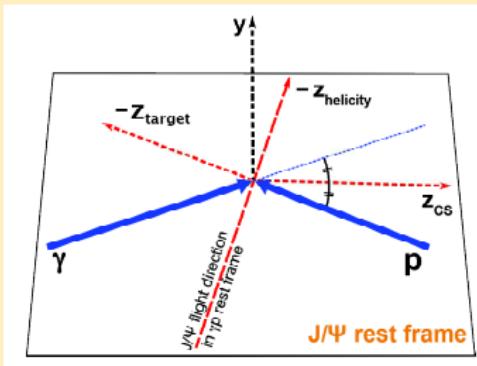
- Agreement with NRQCD at NLO worse than 2002 with LO.
- Just 16 DELPHI events with $p_T > 1 \text{ GeV}$.
- No results from ALEPH, L3, OPAL.
- Data exhausted by single-resolved contribution.

Comparison with BELLE at KEKB



- At NLO, both CSM and NRQCD agree with data.

Polarized J/ψ photoproduction



Decay angular distribution:

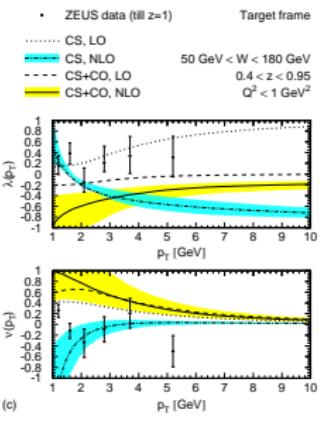
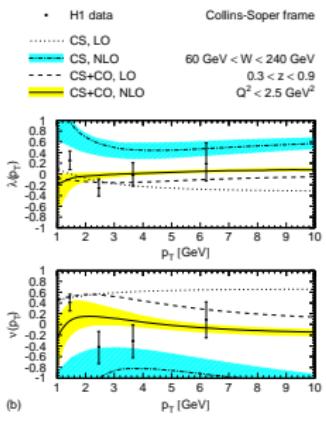
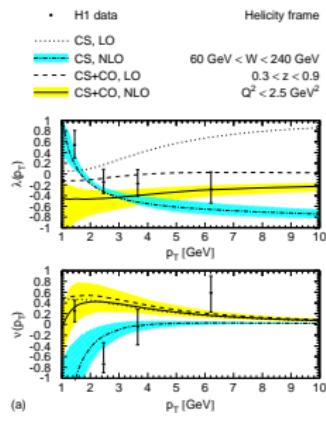
$$\frac{d\Gamma(J/\psi \rightarrow l^+ l^-)}{d\cos\theta d\phi} \propto 1 + \lambda_\theta \cos^2\theta + \lambda_\phi \sin^2\theta \cos(2\phi) + \lambda_{\theta\phi} \sin(2\theta) \cos\phi$$

Polarization observables in spin density matrix formalism:

$$\lambda_\theta = \frac{d\sigma_{11} - d\sigma_{00}}{d\sigma_{11} + d\sigma_{00}}, \quad \lambda_\phi = \frac{d\sigma_{1,-1}}{d\sigma_{11} + d\sigma_{00}}, \quad \lambda_{\theta\phi} = \frac{\sqrt{2}\operatorname{Re} d\sigma_{10}}{d\sigma_{11} + d\sigma_{00}}$$

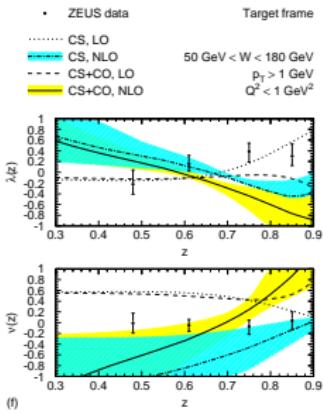
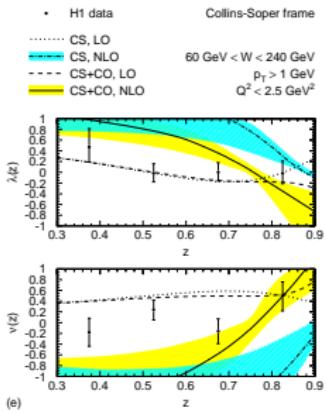
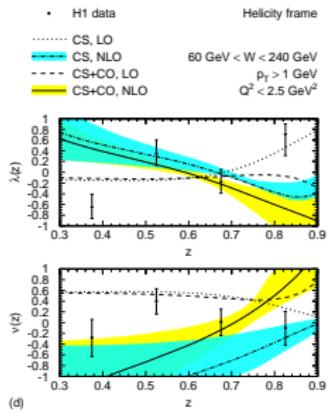
$\lambda = 0, +1, -1$: unpolarized, transversely and longitudinally polarized.

Comparison with H1 and ZEUS



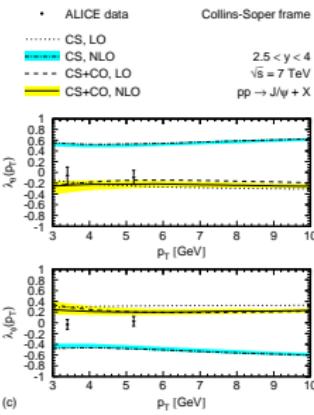
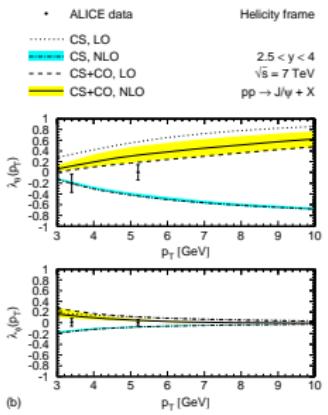
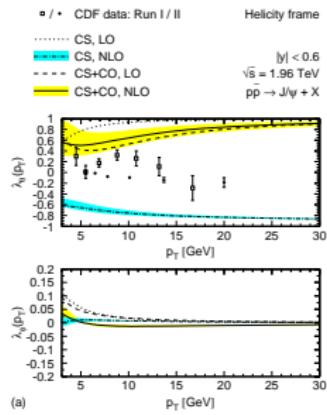
- No z cut on ZEUS data \rightsquigarrow diffractive production included.
 - Perturbative stability in NRQCD higher than in CSM.
 - J/ψ preferably unpolarized at large p_T .

Comparison with H1 and ZEUS (cont.)



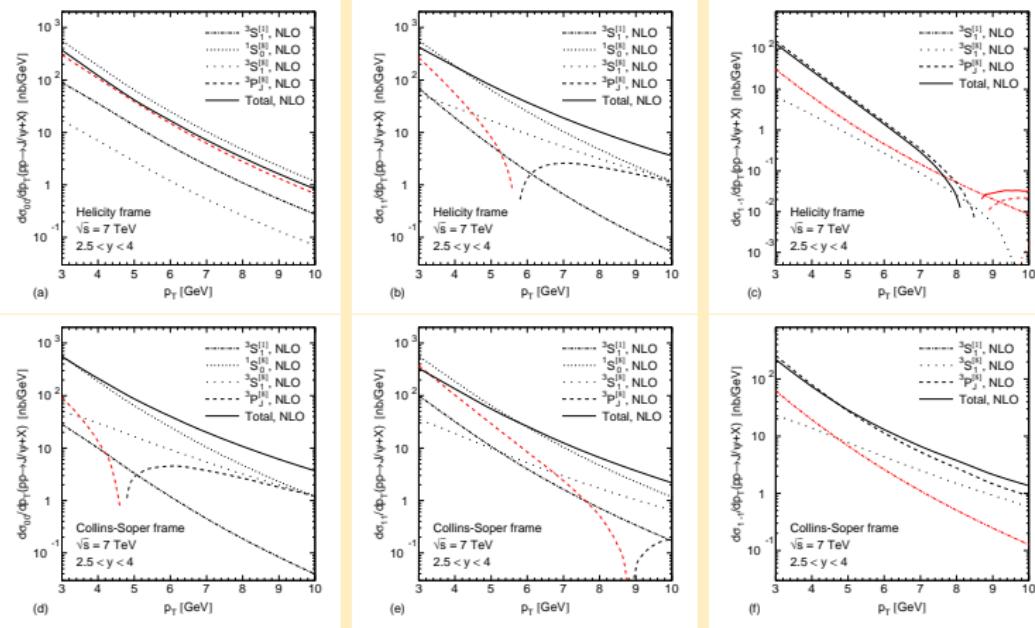
- Large scale uncertainties due to low cut $p_T > 1$.
 - Overall χ^2 w.r.t. default prediction more than halved by going from CSM to NRQCD.

Comparison with CDF and ALICE



- CDF I and II data mutually inconsistent for $p_T < 12$ GeV.
 - CDF J/ψ polarization anomaly persists at NLO ($10\text{--}20\sigma$).
 - 4/8 ALICE points agree w/ NLO NRQCD within errors, others $< 2\sigma$ away.

Decomposition for ALICE



- $d\sigma_{\text{unpol}} = d\sigma_{00} + 2d\sigma_{11}$; $d\sigma_{1,-1}$ auxiliary.
 - Previously unknown ${}^3P_J^{[8]}$ NLO correction significant.

Summary

- NRQCD provides rigorous **factorization theorem** for production and decay of heavy quarkonia; predicts:
 - existence of CO states;
 - universality of LDMEs.
 - Previous LO tests not conclusive.
 - Here: first global analysis of unpolarized- J/ψ world data at NLO.
 - Hadro- and photoproduction: striking evidence of NRQCD vs. CSM.
 - $\gamma\gamma$ scattering, e^+e^- annihilation: not conclusive yet.
 - Contributions from feed-down and B decays throughout small against theoretical uncertainties.
 - Hadroproduction data alone cannot reliably fix all 3 CO LDMEs and give misleading results for their linear combinations; cf.
Ma et al., PRL106(2011)042002; PRD84(2011)114001;
MB & BK, AIPConfProc1343(2011)409.

Summary (cont.)

- Case for NRQCD less strong in polarized J/ψ photoproduction at HERA.
 - Polarized J/ψ hadroproduction at Tevatron in severe conflict with NLO NRQCD, while first LHC data nicely agree.
 - NRQCD factorization in jeopardy! ↵ Hot topic, especially in lack of BSM signals!

Backup Slides

Dependence on low- p_T cut: Global fit

Vary low- p_T cut on pp and $p\bar{p}$ data:

Data left	$p_T > 1$ GeV 148 points	$p_T > 2$ GeV 134 points	$p_T > 3$ GeV 119 points	$p_T > 5$ GeV 86 points	$p_T > 7$ GeV 60 points
$\langle \mathcal{O}^{J/\psi}(^1S_0^{[8]}) \rangle$	5.68 ± 0.37	4.25 ± 0.43	4.97 ± 0.44	4.92 ± 0.49	3.91 ± 0.51
$\langle \mathcal{O}^{J/\psi}(^3S_1^{[8]}) \rangle$	0.90 ± 0.50	2.94 ± 0.58	2.24 ± 0.59	2.23 ± 0.62	2.96 ± 0.64
$\langle \mathcal{O}^{J/\psi}(^3P_0^{[8]}) \rangle$	-2.23 ± 0.17	-1.38 ± 0.20	-1.61 ± 0.20	-1.59 ± 0.22	-1.16 ± 0.23

↝ Global fit insensitive to low- p_T cut on pp and $p\bar{p}$ data as long as γp , $\gamma\gamma$ (74 points with $p_T > 1$ GeV), and e^+e^- data (1 point) are retained.

Vary low- p_T cut on γp and $\gamma\gamma$ data:

Data left	$p_T > 1$ GeV 74 points	$p_T > 2$ GeV 30 points	$p_T > 3$ GeV 15 points	$p_T > 5$ GeV 5 points	$p_T > 7$ GeV 1 points
$\langle \mathcal{O}^{J/\psi}(^1S_0^{[8]}) \rangle$	4.97 ± 0.44	5.10 ± 0.92	4.05 ± 1.17	5.44 ± 1.27	9.56 ± 1.59
$\langle \mathcal{O}^{J/\psi}(^3S_1^{[8]}) \rangle$	2.24 ± 0.59	2.11 ± 1.22	3.52 ± 1.56	1.73 ± 1.68	-3.66 ± 2.09
$\langle \mathcal{O}^{J/\psi}(^3P_0^{[8]}) \rangle$	-1.61 ± 0.20	-1.58 ± 0.48	-0.97 ± 0.63	-1.63 ± 0.68	-3.73 ± 0.83

↝ Global fit insensitive to **moderate** low- p_T cut on γp and $\gamma\gamma$ data as long as pp and $p\bar{p}$ data (119 points with $p_T > 3$ GeV), and e^+e^- data (1 point) are retained.

Dependence on low- p_T cut: Fit to pp and $p\bar{p}$ data only

Vary low- p_T cut:

Data left	$p_T > 1$ GeV 148 points	$p_T > 2$ GeV 134 points	$p_T > 3$ GeV 119 points	$p_T > 5$ GeV 86 points	$p_T > 7$ GeV 60 points
$\langle \mathcal{O}^{J/\psi}(^1S_0^{[8]}) \rangle$	8.54 ± 0.52	16.85 ± 1.23	11.02 ± 1.67	1.68 ± 2.20	2.18 ± 2.56
$\langle \mathcal{O}^{J/\psi}(^3S_1^{[8]}) \rangle$	-2.66 ± 0.69	-13.36 ± 1.60	-5.56 ± 2.19	8.75 ± 2.98	10.34 ± 3.55
$\langle \mathcal{O}^{J/\psi}(^3P_0^{[8]}) \rangle$	-3.63 ± 0.23	-7.70 ± 0.61	-4.46 ± 0.87	2.20 ± 1.23	3.50 ± 1.50
M_0	2.25 ± 0.12	3.51 ± 0.19	3.29 ± 0.20	5.50 ± 0.29	8.24 ± 0.58
M_1	6.37 ± 0.19	5.80 ± 0.19	5.54 ± 0.20	3.27 ± 0.29	1.63 ± 0.43

↷ Fit highly sensitive to low- p_T cut.

Comparison with fit to unpolarized, direct CDF II data with $p_T > 7$ GeV

Y.-Q. Ma, K. Wang, and K.-T. Chao, Phys. Rev. D 84, 114001 (2011):

$$M_0 = (8.54 \pm 1.02) \times 10^{-2} \text{ GeV}^3$$

$$M_1 = (1.67 \pm 1.05) \times 10^{-3} \text{ GeV}^3$$