

# NNPDF determination of polarized PDFs at NLO

XX International Workshop on Deep-Inelastic Scattering  
and Related Subjects

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DI MILANO



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  - Issues in Standard PDF determination
- 2 NNPDF fitting approach
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- 3 Towards NNPDFpol1.0
  - Experimental dataset and PDF parametrization
  - Results: polarized PDFs and the spin content of the proton
- 4 Conclusions
  - Summary and outlook

# 1. Introduction

# Why do we need polarized PDFs?

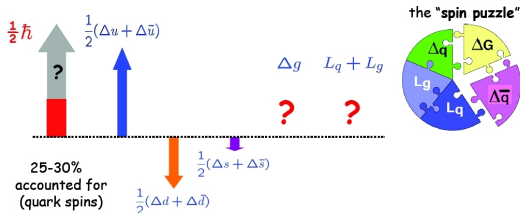
- 1 How do quarks and gluons carry proton's spin?

$$a_0 = \Delta\Sigma \equiv \int_0^1 \Delta\Sigma(x) dx = 2\langle S_z^{\text{quarks}} \rangle \sim 1 \quad \Delta\Sigma(x) = \sum_{i=u,d,s} (\Delta q_i + \Delta \bar{q}_i)$$

EMC experiment (1988):  $a_0 = 0.114 \pm 0.012 \pm 0.026$  "SPIN CRISIS"

$$\langle S_z \rangle = \frac{1}{2} \langle \Delta\Sigma \rangle \longrightarrow \langle S_z \rangle = \frac{1}{2} \langle \Delta\Sigma \rangle + \langle \Delta g \rangle + L_q + L_g$$

- 2 Focus on quark and gluon pieces of the "spin puzzle"



- 3 Rich phenomenology, explore QCD beyond helicity-averaged case

# Issues in standard PDF determination

- Extraction of a set of functions with error bands from a set of data points.
- We need an error band, i.e. a **probability density**  $\mathcal{P}[\Delta q(x)]$  in the space of PDFs:

$$\langle \mathcal{O} \rangle = \int \mathcal{D}\Delta q \mathcal{P}[\Delta q] \mathcal{O}[\Delta q]$$

$$\sigma_{\mathcal{O}}^2 = \int \mathcal{D}\Delta q \mathcal{P}[\Delta q] (\mathcal{O}[\Delta q] - \langle \mathcal{O} \rangle)^2$$

## Standard approach

- 1 Choose a fixed functional form like
$$\Delta q_i(x, Q_0^2) = A_i x^{b_i} (1-x)^{c_i} (1+\dots)$$
- 2 Determine best-fit parameters
- 3 Errors determined via Gaussian linear error propagation

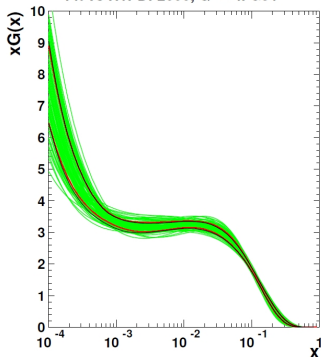
## But...

- 1 Is the parametrization flexible enough?
- 2 What is the error associated to any particular choice?
- 3 Need to rely on linear error propagation

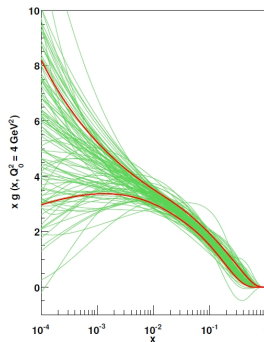
# Simple functional forms vs Neural Networks

## HERA-LHC 2009 PDF benchmarks

Fit vs H1PDF2000,  $Q^2 = 4. \text{ GeV}^2$



simple functional forms



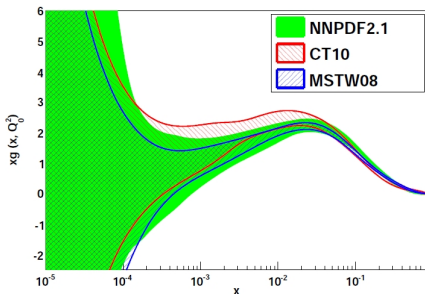
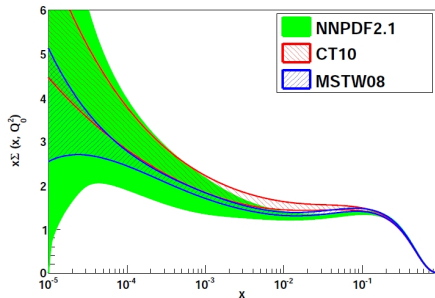
Neural Networks

- Simple functional forms  $\Delta q(x) = Ax^b(1-x)^c P(x)$   
→ systematic underestimation of uncertainties  $\Rightarrow$  tolerance
- Artificial Neural Networks as universal interpolants  
→ reduce theoretical bias from choice of PDF functional form

# PDF fitting: a new approach

NNPDF: a new approach to PDF fitting based on  
**Monte Carlo** sampling and **Neural Networks**

The NNPDF Collaboration, Nucl.Phys. B849 (2011) 296, 1101.1300

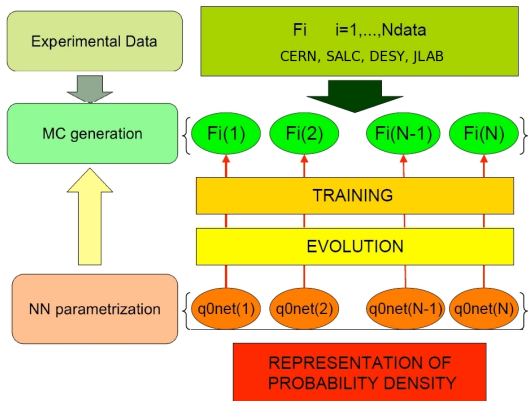


Successfully applied in the unpolarized case: most recent global fit NNPDF2.1  
Routinely used in LHC data analysis and theory prediction

## 2. NNPDF fitting approach



# A general overview on the recipe



Ingredients:

Monte Carlo sampling and Neural Networks

# Ingredient 1: Monte Carlo sampling of experimental data

## MONTE CARLO SAMPLING

- Sample the probability density  $\mathcal{P}[\Delta q]$  in the space of functions assuming **multi-Gaussian** data probability distribution

$$g_{1,p}^{(\text{art}),k}(x, Q^2) = (1 + r_{k,N}\sigma_N) \left[ g_{1,p}^{(\text{exp})}(x, Q^2) + r_{k,t}\sigma_t(x, Q^2) \right]$$

$r_k$ : Gaussian random numbers       $\sigma_N$  quadratic sum of normalization errors  
 $\sigma_t$ : total error (summing in quadrature statistical and systematic errors)

- Generate MC ensemble of  $N_{\text{rep}}$  replicas with the data probability distribution

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## MAIN FEATURES

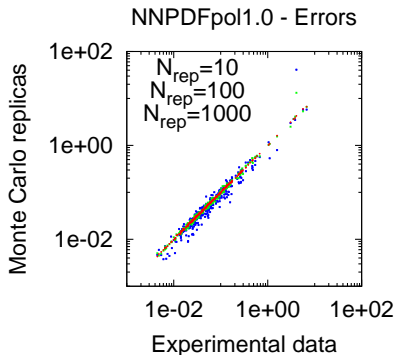
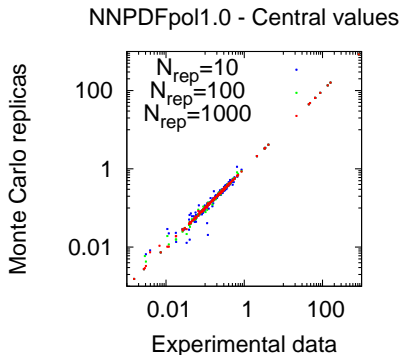
- **Expectation values** for observables are **Monte Carlo integrals**

$$\langle \mathcal{O}[\Delta q] \rangle = \frac{1}{N_{\text{rep}}} \sum_{k=1}^{N_{\text{rep}}} \mathcal{O}[\Delta q_k]$$

... and the same is true for errors, correlations etc.

- No need to rely on **linear propagation** of errors
- Possibility to test for **non-Gaussian** behaviour in fitted PDFs

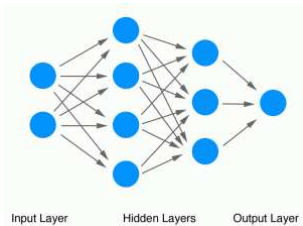
# Ingredient 1: Monte Carlo sampling of experimental data



- Require the average over the replicas reproduces central values and errors of the original experimental data to desired accuracy → determine size of the sample
- Accuracy of few % requires  $\sim 1000$  replicas

## Ingredient 2: Neural Networks

A convenient **functional form**  
providing **redundant** and **flexible** parametrization  
used as a generator of random functions in the PDF space



$$\xi_i^{(l)} = g \left( \sum_j^{n_{l-1}} \omega_{ij}^{(l-1)} \xi_j^{(l-1)} - \theta_i^{(l)} \right)$$

$$g(x) = \frac{1}{1 + e^{-x}}$$

- made of neurons grouped into layers (define the architecture)
- each neuron receives input from neurons in preceding layer (feed-forward NN)
- activation determined by parameters (**weights** and **thresholds**)
- activation determined according to a **non-linear function**

# Ingredient 2: Neural Networks

## NEURAL NETWORKS

- **Parametrize** each polarized PDF replica with flexible Neural Network

DSSV, AAC, LSS, BB

$\mathcal{O}(10 - 20)$  parameters

NNPDFpol

$\mathcal{O}(200)$  parameters

- **Train** NN to determine the best fit for each replica
- Compute an ensemble of observables and compare to experimental data

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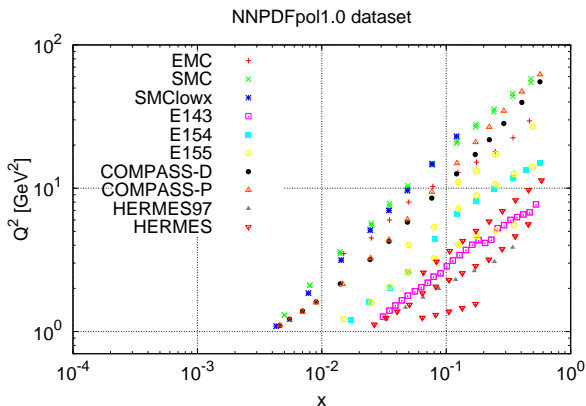
## MAIN FEATURES

- Only require **smoothness** of the fitted function
- Do not require any other prejudice on *a priori* functional form
- **Reduce** the **bias** associated to the choice of some functional form
- Given smoothness, the algorithm provided by NN is efficient and can be easily implemented with other algorithms (e.g. genetic algorithms)

### 3. Towards NNPDFpol1.0



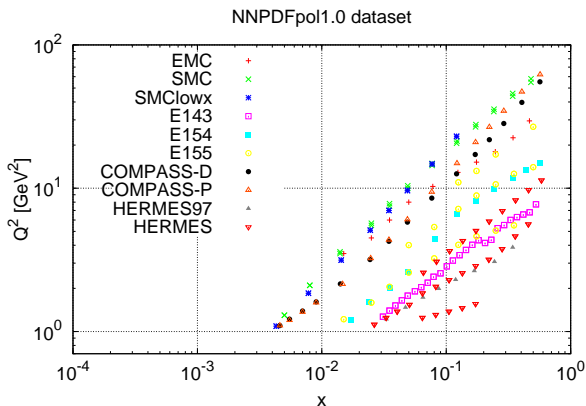
# Experimental dataset



$$g_1(x, Q^2) = A_1(x, Q^2) \frac{F_2(x, Q^2)}{2x(1 + R(x, Q^2))} (1 + \gamma^2)$$

$$\gamma^2 = \frac{4M_N^2 x^2}{Q^2}$$

# Experimental dataset



- 1 All relevant polarized DIS data on proton, neutron and deuteron targets
- 2 Kinematical cuts to remove the sensitivity to dynamical higher-twist
  - $Q^2 > 1 \text{ GeV}^2$
  - $W^2 = Q^2(1-x)/x \geq 6.25 \text{ GeV}^2$  (C. Simolo, Ph.D. Thesis. [arXiv:0807.1501](https://arxiv.org/abs/0807.1501))

- 1 **Four polarized PDFs** (gluon + linear combinations of light quarks)
  - singlet  $\Delta\Sigma(x) \equiv \sum_{i=1}^{n_f} (\Delta q_i(x) + \Delta \bar{q}_i(x))$
  - gluon  $\Delta g(x)$
  - triplet  $\Delta T_3(x) \equiv (\Delta u(x) + \Delta \bar{u}(x)) - (\Delta d(x) + \Delta \bar{d}(x))$
  - octet  $\Delta T_8(x) \equiv (\Delta u(x) + \Delta \bar{u}(x)) + (\Delta d(x) + \Delta \bar{d}(x)) - 2(\Delta s(x) + \Delta \bar{s}(x))$
- 2 At **initial scale**  $Q_0^2 = 1\text{GeV}^2$  and using  $\alpha_s(M_Z^2) = 0.119$
- 3 Assume all heavy quarks are generated radiatively
- 4 Must satisfy **theoretical constraints**:
  - Sum rules

$$[\Delta T_3(Q_0^2)] \equiv \int_0^1 dx \Delta T_3(x, Q_0^2) = a_3 \quad [\Delta T_8(Q_0^2)] \equiv \int_0^1 dx \Delta T_8(x, Q_0^2) = a_8$$

- Positivity bound for all proton, neutron and deuteron targets

$$|g_1(x, Q^2)| \leq F_1(x, Q^2)$$

# Preprocessing: basic idea

- 1 Each polarized PDF parametrized with a multi-layer feed-forward NN. All NN have the same **architecture (2-5-3-1)**.
- 2 Parametrization supplemented with a **preprocessing polynomial**: exponents **m** and **n randomly chosen** in fixed intervals; the NN only fits the deviation from this function.

$$\Delta\Sigma(x, Q_0^2) = (1-x)^{m_{\Delta\Sigma}} x^{-n_{\Delta\Sigma}} NN_{\Delta\Sigma}(x)$$

$$\Delta g(x, Q_0^2) = (1-x)^{m_{\Delta g}} x^{-n_{\Delta g}} NN_{\Delta g}(x)$$

$$\Delta T_3(x, Q_0^2) = A_{\Delta T_3} (1-x)^{m_{\Delta T_3}} x^{-n_{\Delta T_3}} NN_{\Delta T_3}(x)$$

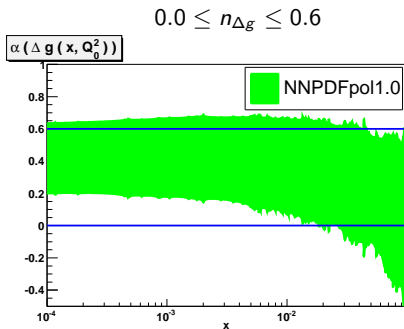
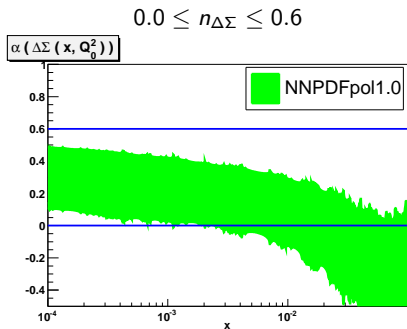
$$\Delta T_8(x, Q_0^2) = A_{\Delta T_8} (1-x)^{m_{\Delta T_8}} x^{-n_{\Delta T_8}} NN_{\Delta T_8}(x)$$

- 3 **Overall normalization constant** factored out for **triplet** and **octet**. Determined by imposing the sum rules.

$$A_{\Delta T_3} = \frac{a_3}{\int_0^1 dx [(1-x)^{m_{\Delta T_3}} x^{-n_{\Delta T_3}} NN_{\Delta T_3}(x)]}$$

$$A_{\Delta T_8} = \frac{a_8}{\int_0^1 dx [(1-x)^{m_{\Delta T_8}} x^{-n_{\Delta T_8}} NN_{\Delta T_8}(x)]}$$

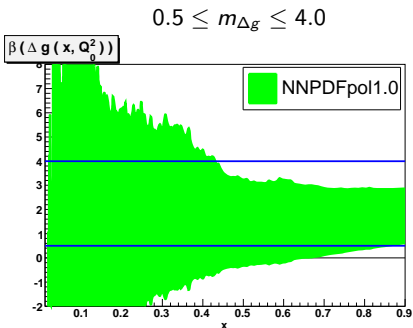
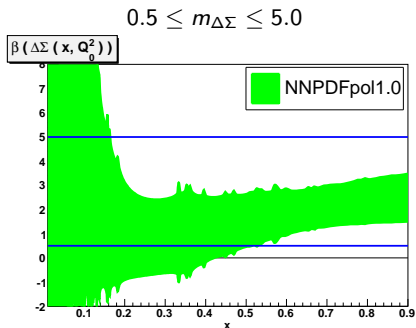
# Preprocessing: effective asymptotic exponents



$$\alpha_{\text{eff}}(x, Q^2) \equiv \ln(\Delta q(x, Q^2)) / \ln(1/x) \quad \text{at } Q^2 = Q_0^2 = 1\text{GeV}^2$$

Effective exponents always contained in the preprocessing exponents range  
The polarized PDF is driven only by experimental data

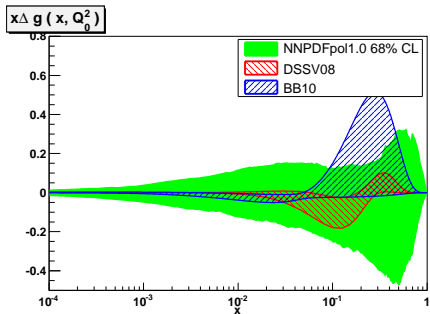
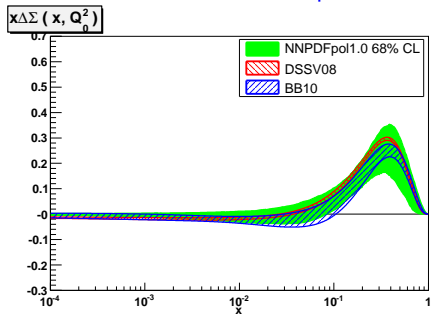
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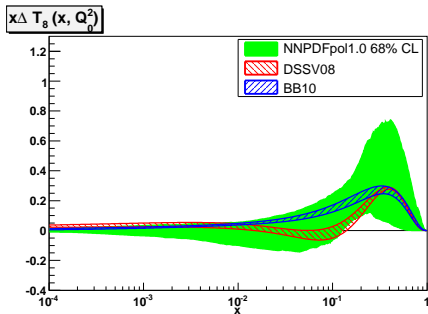
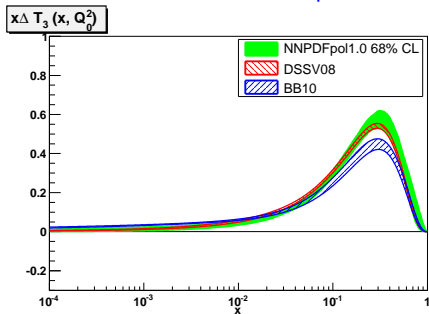
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## Comparison with DSSV08 and BB10



Much larger error bands for Singlet and Gluon

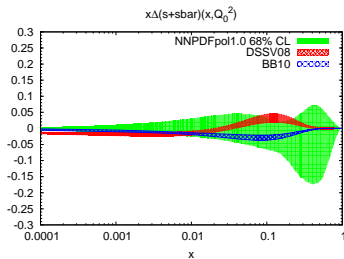
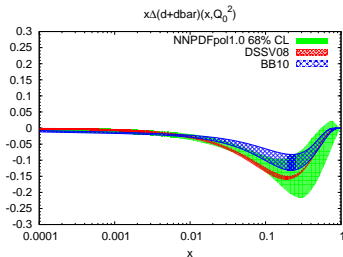
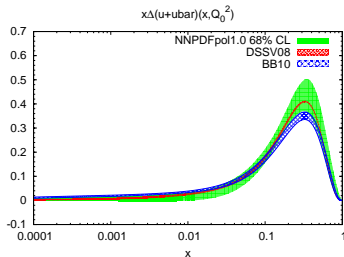
## Comparison with DSSV08 and BB10



Triplet agrees with DSSV08 fit, but does not with BB10 fit  
Much larger error band for Octet



## Comparison with DSSV08 and BB10



Much larger error bands for all light quark combinations

NNPDFpol1.0 does not agree with BB10 for the  $x\Delta(u + \bar{u})$  combination

# The spin content of the proton

Singlet and Gluon first moments in  $\overline{\text{MS}}$  scheme at  $Q_0^2 = 1 \text{ GeV}^2$

	NNPDFpol1.0	DSSV08	AAC08
$[\Delta\Sigma]$	$0.32 \pm 0.11$	$0.26 \pm 0.03$	$0.26 \pm 0.06$
$[\Delta g]$	$-0.2 \pm 1.1$	$-0.12 \pm 0.12$	$0.40 \pm 0.28$

Notice the large uncertainty on the first moments:

**Singlet** between two and four times

**Gluon** almost one order of magnitude

$$\langle S_z \rangle = \frac{1}{2} \langle \Delta\Sigma \rangle + \langle \Delta g \rangle + L_q + L_g$$

$$\frac{1}{2} = (-0.1 \pm 1.1) + L_q + L_g$$

Gluon uncertainty dominates the contribution to proton's spin

# 4. Conclusions

## Summary

- 1 The NNPDF technology provides a statistically sound procedure for PDF fitting
- 2 NNPDFpol1.0 is the first polarized parton determination using NNPDF approach
- 3 The analysis from inclusive DIS data leads to
  - able to discriminate Triplet (agreement with DSSV08, not with BB10)
  - large uncertainties on Singlet and Octet and very large on Gluon
  - uncertainty on Singlet first moment between two and four times bigger
  - uncertainty on Gluon first moment almost one order of magnitude bigger

## Outlook

- 1 Include data sets from other processes (open charm and jet production with fixed target, inclusive jet production, W boson production at RHIC, ...)
- 2 Determine the strong-coupling constant from polarized DIS data
- 3 Investigate the sensitivity of polarized data to  $a_3$  and  $a_8$  axial constants

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**Thank you for your attention!**

# 5. Backup

# PDF fitting: state of the art

- 1 First stage: first **moments** of polarized PDFs and polarized **sum rules** (last 25 years)  
→ “historical” experimental collaborations (at CERN, SLAC, DESY, JLAB):  
EMC, SMC, E142, E143, E154, E155, COMPASS, HERMES, CLAS, ...
- 2 Second stage: polarized **PDF fits** from **global NLO QCD analysis** (last ~15 years)  
→ different choice of datasets, parton parametrization, treatment of higher twists, ...  
ABFR ([arXiv:hep-ph/9803237](https://arxiv.org/abs/hep-ph/9803237), 1998), BB ([arXiv:1005.3113](https://arxiv.org/abs/1005.3113), 2010) (DIS only);  
AAC ([arXiv:0808.0413](https://arxiv.org/abs/0808.0413), 2008), LSS ([arXiv:1010.0574](https://arxiv.org/abs/1010.0574), 2010) (DIS+SIDIS);  
DSSV ([arXiv:0904.3821](https://arxiv.org/abs/0904.3821), 2009) (DIS+SIDIS+pp)
- 3 Third stage: provide **uncertainties** on polarized PDFs (last ~10 years)  
→ Gaussian error propagation, Lagrange multiplier + Hessian method; fit with orthogonal polynomials ([arXiv:1011.4873](https://arxiv.org/abs/1011.4873), 2010)

# Monte Carlo sampling: more detail

- The  $k$ -th MC replica is generated assuming a multi-Gaussian distribution

$$g_{1,i}^{(\text{art}),k}(x, Q^2) = (1 + r_{k,N}\sigma_N) \left[ g_{1,p}^{(\text{exp})}(x, Q^2) + r_{k,t}\sigma_t(x, Q^2) \right]$$

$r_k$ : Gaussian random numbers       $\sigma_N$  quadratic sum of normalization errors  
 $\sigma_t$ : total error (summing in quadrature statistical and systematic errors)

- Require the average over the replicas reproduces central values and errors of the original experimental data to desired accuracy  $\rightarrow$  determine size of the sample

$r [g_1]$	1.00
$\langle \sigma^{(\text{exp})} \rangle_{\text{dat}} (\%)$	0.13E+03
$\langle \sigma^{(\text{gen})} \rangle_{\text{dat}}$	0.11E+03
$r [\sigma^{(\text{gen})}]_{\text{dat}} (\%)$	0.10E+01
$\langle \rho^{(\text{exp})} \rangle_{\text{dat}}$	0.64E-01
$\langle \rho^{(\text{gen})} \rangle_{\text{dat}}$	0.64E-01
$r [\rho^{(\text{gen})}]_{\text{dat}}$	0.98E+00
$\langle \text{cov}^{(\text{exp})} \rangle_{\text{dat}}$	0.25E-01
$\langle \text{cov}^{(\text{gen})} \rangle_{\text{dat}}$	0.25E-01
$r [\text{cov}^{(\text{gen})}]_{\text{dat}}$	0.10E+01

- Accuracy of few % requires  $\sim 1000$  replicas



# Monte Carlo sampling: more detail

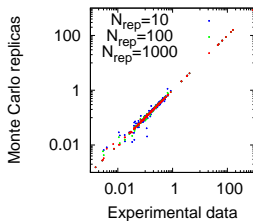
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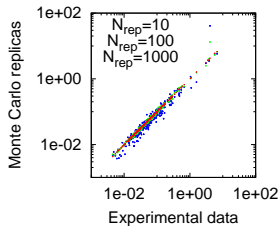
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NNPDFpol1.0 - Central values

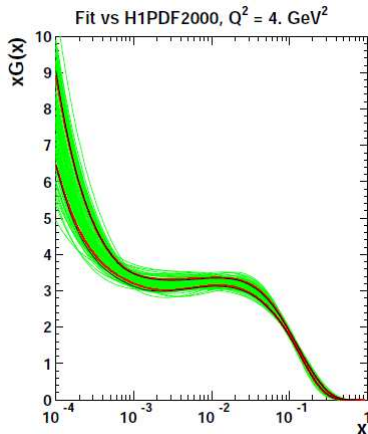


NNPDFpol1.0 - Errors



- Accuracy of few % requires  $\sim 1000$  replicas

# Monte Carlo vs Hessian PDF uncertainties



## HERA-LHC 2009 PDF benchmarks

- H1PDF2000 fit done with Hessian method and with Monte Carlo method
- The standard deviation of the 100 PDF replicas (MC method) is in perfect agreement with Hessian errors with  $\Delta\chi^2 = 1$
- The MC method to estimate PDF uncertainties reproduces Hessian result when global  $\chi^2$  is quadratic

In Mellin space the DGLAP equations

$$\begin{aligned} \mu^2 \frac{\partial}{\partial \mu^2} \Delta q_{NS}^{\pm, \nu}(N, \mu^2) &= \Delta \gamma_{NS}^{\pm, \nu} q_{NS}^{\pm, \nu}(N, \mu^2) \\ \mu^2 \frac{\partial}{\partial \mu^2} \begin{pmatrix} \Delta \Sigma \\ \Delta g \end{pmatrix}(N, \mu^2) &= \begin{pmatrix} \Delta \gamma_{qq}(N, \alpha_s(Q^2)) & \Delta \gamma_{qg}(N, \alpha_s(Q^2)) \\ \Delta \gamma_{gq}(N, \alpha_s(Q^2)) & \Delta \gamma_{gg}(N, \alpha_s(Q^2)) \end{pmatrix} \begin{pmatrix} \Delta \Sigma \\ \Delta g \end{pmatrix} \end{aligned}$$

can be solved analytically

$$\Delta q_{NS}^{\pm, \nu}(N, Q^2) = \Gamma_{NS}^{\pm, \nu}(N, a_s, a_0) \Delta q_{NS}^{\pm, \nu}(N, Q_0^2), \quad a_s \equiv \alpha_s/2\pi$$

where, at NLO,

$$\Gamma_{NS, NLO}^{\pm, \nu}(N, a_s, a_0) = \exp \left\{ \frac{U_1^{\pm, \nu}}{b_1} \ln \left( \frac{1 + b_1 a_s}{1 + b_1 a_0} \right) \right\} \left( \frac{a_s}{a_0} \right)^{-R_0^{NS}}$$

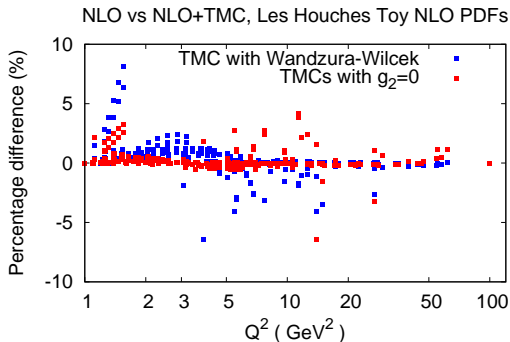
# Polarized PDF evolution

NNPDF NLO polarized PDF evolution (**Fast Kernel method**) benchmarked with the Les Houches PDF benchmarks ([G. Salam and a. Vogt, hep-ph/0511119](#))

$x$	$\epsilon_{\text{rel}}(\Delta u_V)$	$\epsilon_{\text{rel}}(\Delta d_V)$	$\epsilon_{\text{rel}}(\Delta \Sigma)$	$\epsilon_{\text{rel}}(\Delta g)$
$10^{-3}$	$1.1 \cdot 10^{-4}$	$9.2 \cdot 10^{-5}$	$9.9 \cdot 10^{-5}$	$1.1 \cdot 10^{-4}$
$10^{-2}$	$1.4 \cdot 10^{-4}$	$1.9 \cdot 10^{-4}$	$3.5 \cdot 10^{-4}$	$9.3 \cdot 10^{-5}$
0.1	$1.2 \cdot 10^{-4}$	$1.6 \cdot 10^{-4}$	$5.4 \cdot 10^{-6}$	$1.7 \cdot 10^{-4}$
0.3	$2.3 \cdot 10^{-6}$	$1.1 \cdot 10^{-5}$	$7.5 \cdot 10^{-6}$	$1.7 \cdot 10^{-5}$
0.5	$5.6 \cdot 10^{-6}$	$9.6 \cdot 10^{-6}$	$1.6 \cdot 10^{-5}$	$2.5 \cdot 10^{-5}$
0.7	$1.2 \cdot 10^{-4}$	$9.2 \cdot 10^{-7}$	$1.6 \cdot 10^{-4}$	$7.8 \cdot 10^{-5}$
0.9	$3.5 \cdot 10^{-3}$	$1.1 \cdot 10^{-2}$	$4.1 \cdot 10^{-3}$	$7.8 \cdot 10^{-3}$

**Very accurate evolution!**

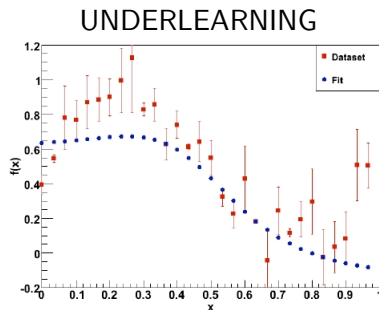
# Target mass corrections



Kinematical cuts exclude the largest- $x$  and smallest- $Q^2$  data region,  
where the TMC effects are most important  
Moderate impact of TMC corrections (small percent at small  $Q^2$ )

# One more ingredient: Stopping

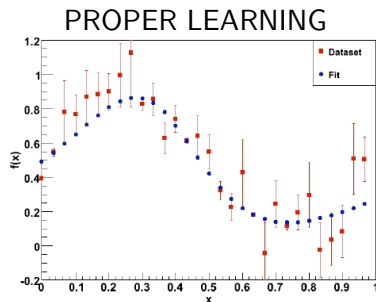
- 1 NN are flexible tools
  - can learn fluctuations
- 2 Cross-validation method
  - divide data into two subsets (training & validation)
  - train the NN on training subset
  - compute  $\chi^2$  for each subset
  - stop when  $\chi^2$  of validation subset no longer decreases (NN are learning fluctuations!)



The best fit does not coincide with the  $\chi^2$  absolute minimum

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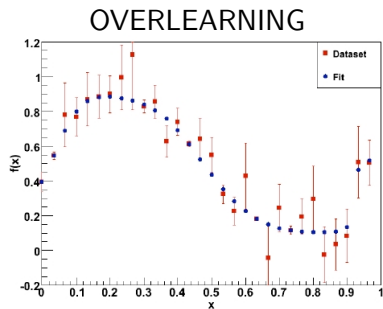
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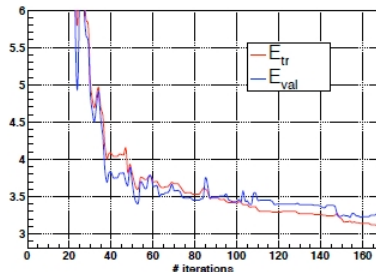


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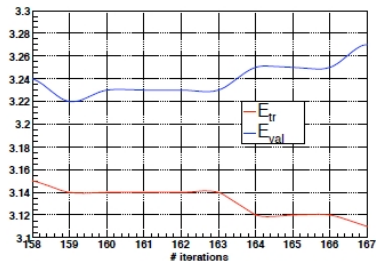
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The best fit does not coincide with the  $\chi^2$  absolute minimum

- Set Neural Network parameters randomly
- Make clones of the parameter vector and mutate them
- Define a **figure of merit** or error function for the  $k$ -th replica

$$E^{(k)} = \frac{1}{N_{\text{rep}}} \sum_{i,j=1}^{N_{\text{rep}}} \left( F_i^{(\text{art})(k)} - F_i^{(\text{net})(k)} \right) \left( (\text{cov}_{t_0})^{-1} \right)_{ij} \left( F_j^{(\text{art})(k)} - F_j^{(\text{net})(k)} \right)$$

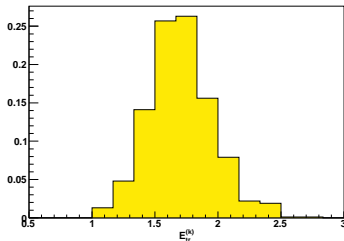
$$(\text{cov}_{t_0})_{ij} = \sigma_{N,i} \sigma_{N,j} \mathbf{g}_{1i}^{(0)} \mathbf{g}_{1j}^{(0)} + \delta_{ij} \sigma_{i,tot}^2$$

is the covariance matrix including normalization errors using the  $t_0$  method  
([arXiv:0912.2276](https://arxiv.org/abs/0912.2276))

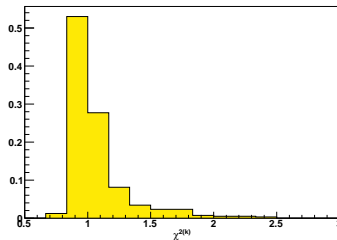
- Select the best ones and perform other manipulations (crossing, mutating, ...) until stability is reached

# NNPDFpol1.0: global $\chi^2$

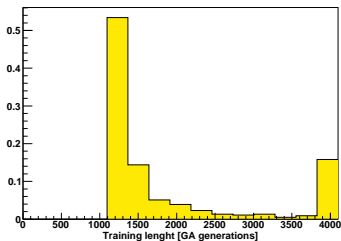
$E_{tr}$  distribution for MC replicas



$\chi^2(k)$  distribution for MC replicas



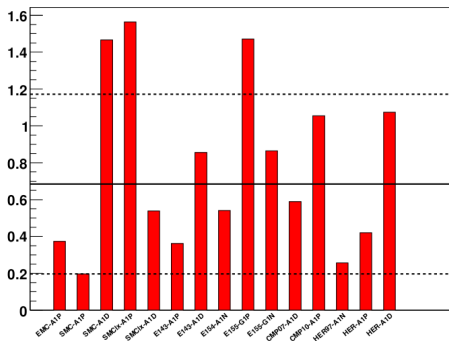
Distribution of training lengths



$\chi_{tot}^2$	0.75
$\langle E \rangle \pm \sigma_E$	$1.89 \pm 0.19$
$\langle E_{tr} \rangle \pm \sigma_{E_{tr}}$	$1.64 \pm 0.22$
$\langle E_{val} \rangle \pm \sigma_{E_{val}}$	$2.14 \pm 0.33$
$\langle \chi^{2(k)} \rangle \pm \sigma_{\chi^2}$	$0.96 \pm 0.12$
$\langle \sigma^{(exp)} \rangle_{dat} (\%)$	130
$\langle \sigma^{(net)} \rangle_{dat} (\%)$	48
$\langle TL \rangle$	1825

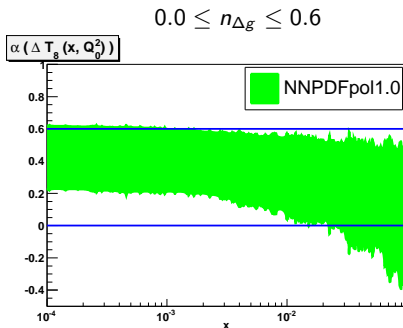
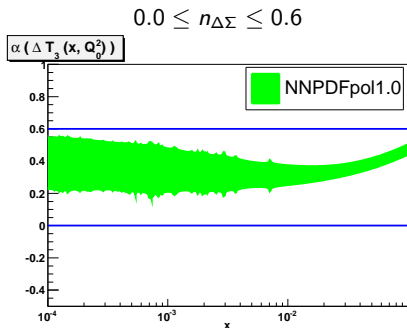
# NNPDFpol1.0: individual experiments $\chi^2$

Distribution of  $\chi^2$  for sets



- No evidence of any specific dataset being inconsistent with each other
- Distribution of individual  $\chi^2$  values broadly consistent with statistical expectations

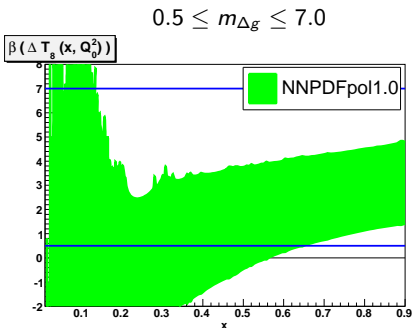
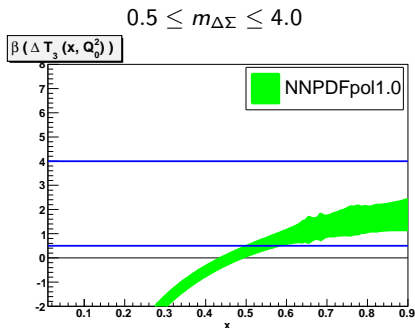
# Preprocessing: effective asymptotic exponents



$$\alpha_{\text{eff}}(x, Q^2) \equiv \ln(\Delta q(x, Q^2)) / \ln(1/x) \quad \text{at } Q^2 = Q_0^2 = 1\text{GeV}^2$$

Effective exponents always contained in in the preprocessing exponents range  
The polarized PDF is driven only by experimental data

# Preprocessing: effective asymptotic exponents

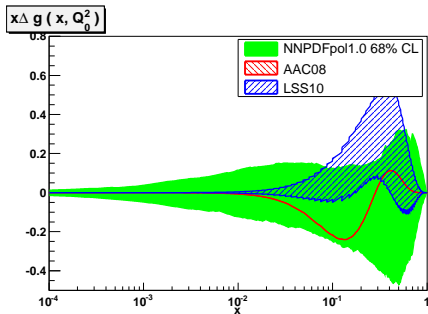
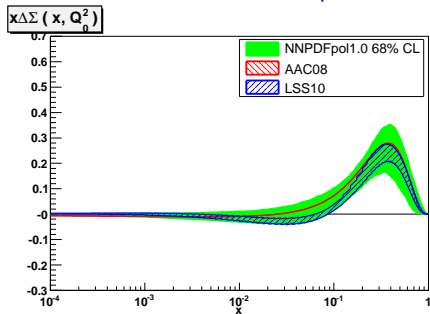


$$\beta_{\text{eff}}(x, Q^2) \equiv \ln(\Delta q(x, Q^2)) / \ln(1-x) \quad \text{at } Q^2 = Q_0^2 = 1\text{GeV}^2$$

Effective exponents always contained in in the preprocessing exponents range  
The polarized PDF is driven only by experimental data

# NNPDFpol1.0: more results

## Comparison with AAC08 and LSS10

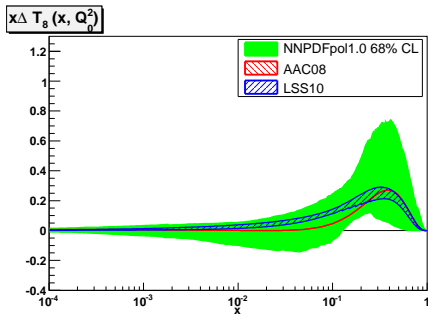
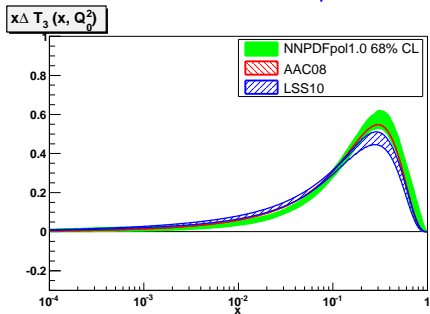


Substantial agreement for Singlet  
Gluon essentially unconstrained



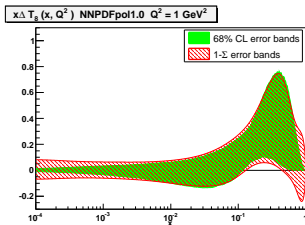
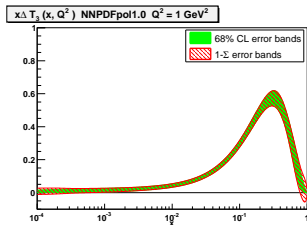
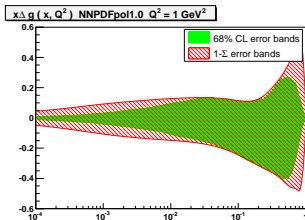
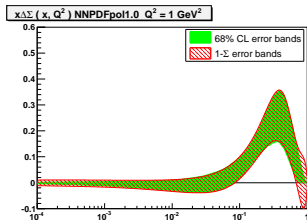
# NNPDFpol1.0: more results

## Comparison with AAC08 and LSS10



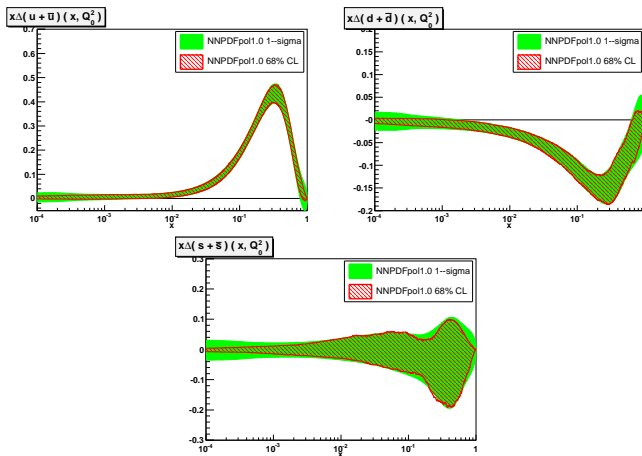
Substantial agreement for Octet, but with larger error band  
Triplet slightly bigger

# NNPDFpol1.0: 68% confidence levels



comparison between  $1\sigma$  error bands and 68% confidence level  
test for non-Gaussian behaviour  
sizeable deviations from Gaussian behaviour in the extrapolation region

# NNPDFpol1.0: 68% confidence levels



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