

On the NNLO QCD corrections to DIS heavy-quark production

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Plan

- Talk based on results on ...
 - ... improved heavy-quark coefficient functions at NNLO
H. Kawamura , S.M., A. Io Presti and A. Vogt to appear
 - ... a determination of the running charm-quark mass
S. Alekhin , K. Daum , K. Lipka and S.M. to appear
 - ... a spin-off on top-quark hadro-production
S.M., P. Uwer and A. Vogt to appear

Heavy-quark production in DIS

QCD factorization

- Structure functions for DIS heavy-quark production
- Perturbative expansion at NLO, e.g., neutral current

$$F_2^c(x, Q^2, m^2) = \frac{\alpha_s e_q^2}{\pi^2} \sum_{i=q, \bar{q}, g} f_i(z, \mu^2) \otimes \left\{ c_{i,k}^{(0)}(\eta, \xi) + (4\pi\alpha_s) c_{i,k}^{(1)}(\eta, \xi, \mu^2/m^2) \right\}$$

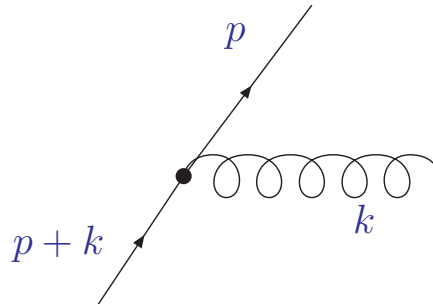
- Kinematic variables
 - distance from threshold $\eta = s/(4m^2) - 1$
 - virtuality $\xi = Q^2/m^2$

NLO

- NLO for coefficient functions $c_{i,k}$
 - neutral current Laenen, Riemersma, Smith, van Neerven '93
 - charged current Gottschalk '81; Glück, Kretzer, Reya '96

Strategy beyond NLO

- Use universal features of soft/collinear regions of phase space
 - double logarithms from singular regions in Feynman diagrams
 - propagator vanishes for: $E_g = 0$, soft $\theta_{qg} = 0$ collinear

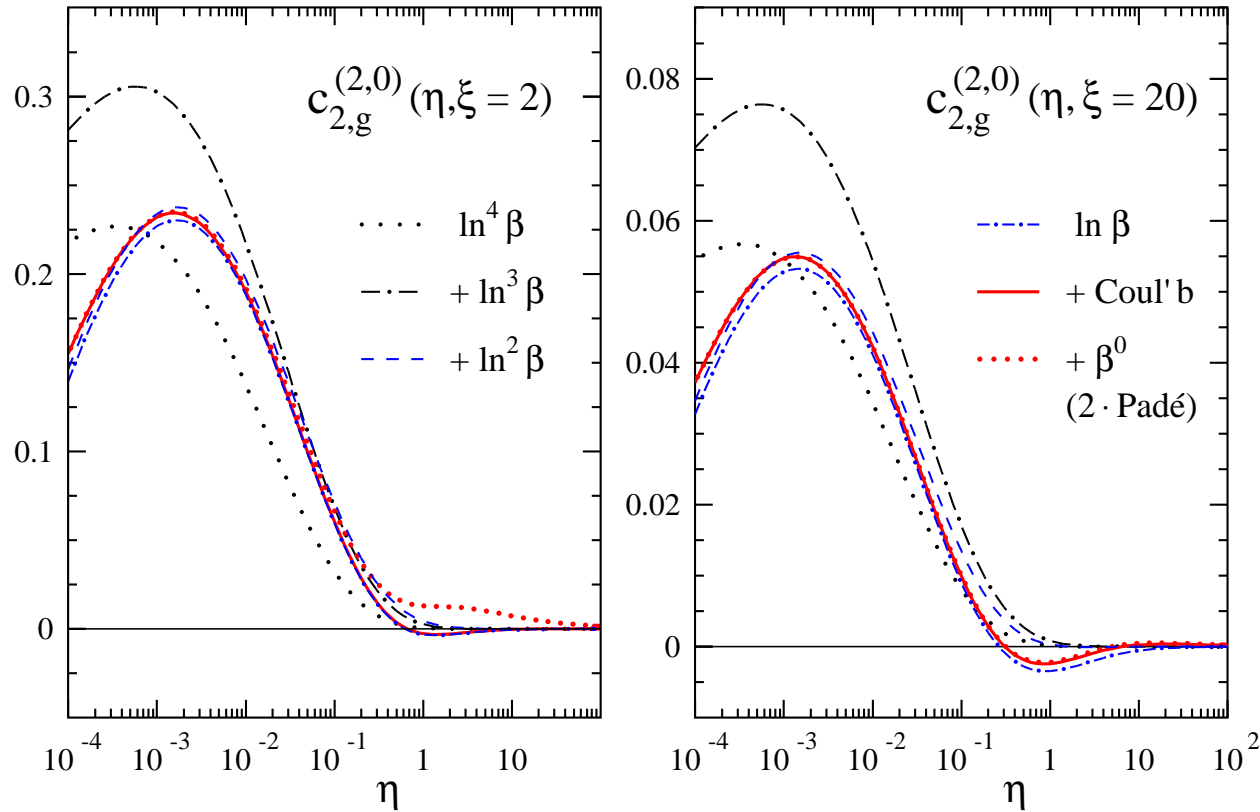


$$\alpha_s \int d^4k \frac{1}{(p+k)^2} \quad \begin{aligned} &= \frac{1}{2p \cdot k} = \frac{1}{2E_q E_g (1 - \cos \theta_{qg})} \\ &\longrightarrow \alpha_s \int dE_g d\sin \theta_{qg} \frac{1}{2E_q E_g (1 - \cos \theta_{qg})} \\ &\longrightarrow \alpha_s \ln^2(\dots) \end{aligned}$$

Threshold improvements beyond NLO

- Parton cross section close to threshold $s \simeq 4m^2$
 - Sudakov logarithms $\ln(\beta)$ with velocity of heavy quark
 $\beta = \sqrt{1 - 4m^2/s}$ at n^{th} -order: $\alpha_s^n \ln^{2n}(\beta) \longleftrightarrow \alpha_s^n \ln^{2n}(N)$
- Resummation in Mellin space (renormalization group equation) predicts fixed orders in perturbation theory
 - approximate expressions to NNLO
 Laenen, S.M. '98; Alekhin, S.M. '08; Lo Presti, Kawamura, S.M., Vogt '10

Approximate coefficient functions at NNLO



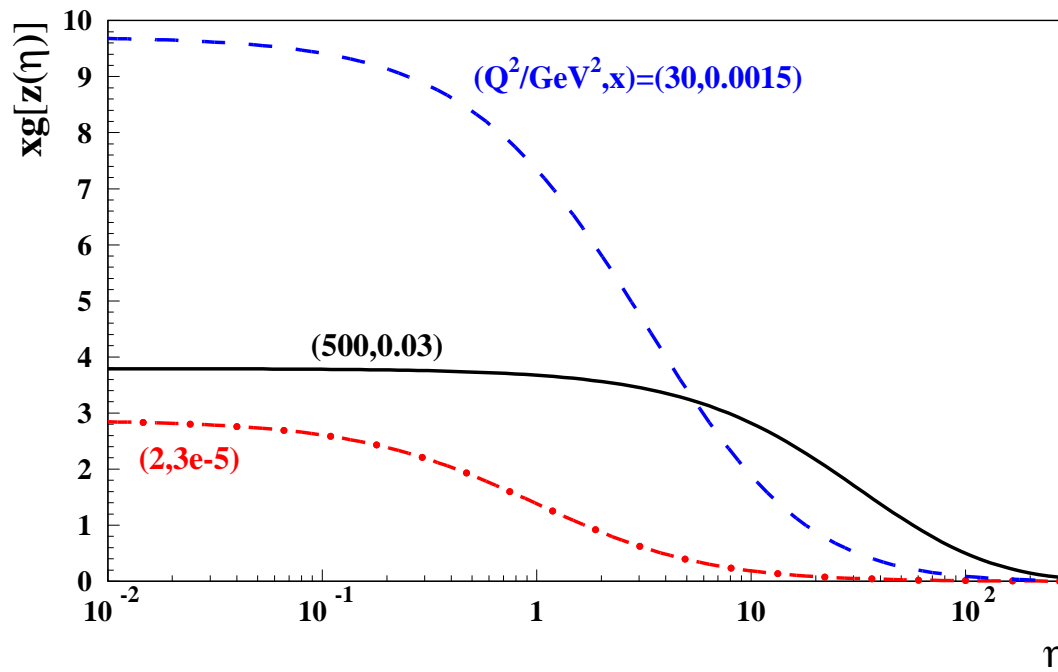
- Distance from threshold $\eta = s/(4m^2) - 1$
 - Sudakov logarithms dominant for $\eta \ll 1$ now known to NNLL
Lo Presti, Kawamura, S.M., Vogt '10
- Small impact of subleading terms $\mathcal{O}(\beta^0)$ (Padé estimate)

Convolution with gluon PDF

- Recall QCD factorization

$$x^{-1} F_2^c(x, Q^2, m^2) = \frac{\alpha_s e_q^2}{\pi^2} \sum_{i=q, \bar{q}, g} \int_0^{\eta_{max}} d\eta f_i(z(\eta), \mu^2) c_{i,k}(\eta, \xi, \mu^2)$$

- gluon PDF gives large weight to parton dynamics near threshold for $Q^2 \lesssim 10 - 30 \text{ GeV}^2$



Asymptotics beyond NLO

- Exact factorization at large $Q^2 \gg m^2$ ($\xi \rightarrow \infty$)

Buza, Matiounine, Smith, Migneron, van Neerven '95

- decompose heavy-quark $c_{k,i}$ into light-quark $C_{k,j}^{\text{light}}$ and operator matrix elements (OMEs) A_{ji}

$$c_{k,i}(\eta, \xi, \mu^2) \rightarrow \left[A_{ji} \left(\frac{m^2}{\mu^2} \right) \otimes C_{k,j}^{\text{light}} \left(\frac{Q^2}{\mu^2} \right) \right] (x)$$

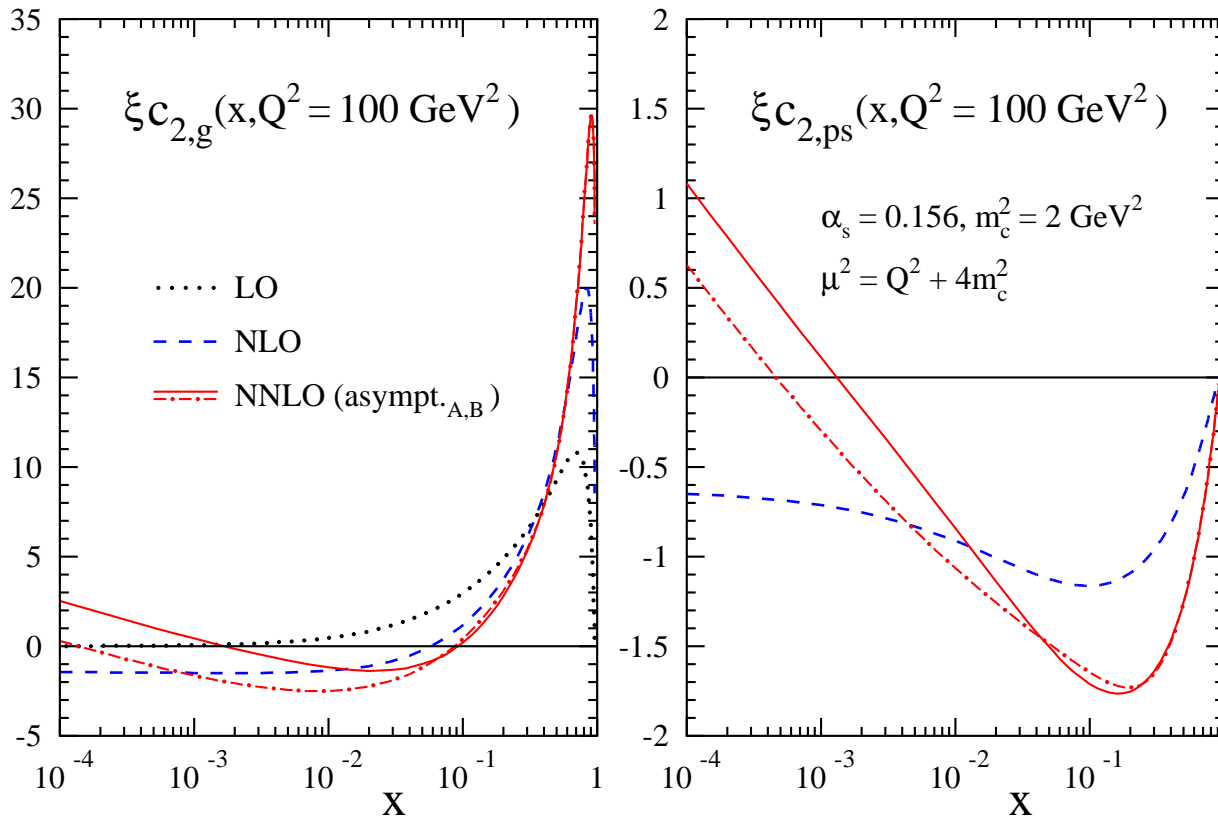
- NNLO corrections require three-loop results

- massless coefficient functions Vermaseren, S.M., Vogt '05

- heavy-quark OMEs; currently only fixed Mellin moments available

Bierenbaum, Blümlein, Klein '09

Approximate coefficient functions at NNLO



- Asymptotically large $\xi = Q^2/m^2$
 - complete functional dependence on x now known to NNLO
Lo Presti, Kawamura, S.M., Vogt to appear
- Sizable impact of parametrization uncertainties in OMEs at small x
 - significant improvements with more Mellin moments

High-energy limit beyond NLO

- Well-known k_t -factorization at large η (small x)

Catani, Ciafaloni, Hautmann '91

- decompose heavy-quark structure function $F_2^c(x, Q^2, m^2)$ into gluon PDF un-integrated in k_t and heavy-quark impact factor

- Schematic expansion up to NNLO

$$c_{k,i}^{(1)}(\eta, \xi, \mu^2) \approx c_{LL_x}^{(1)} \text{const}_\eta$$

$$c_{k,i}^{(2)}(\eta, \xi, \mu^2) \approx c_{LL_x}^{(2)} \ln \eta + c_{NLL_x}^{(2)} \text{const}_\eta$$

- Small- x dependence governed by BFKL anomalous dimension

- prediction of $c_{LL_x}^{(l)}$

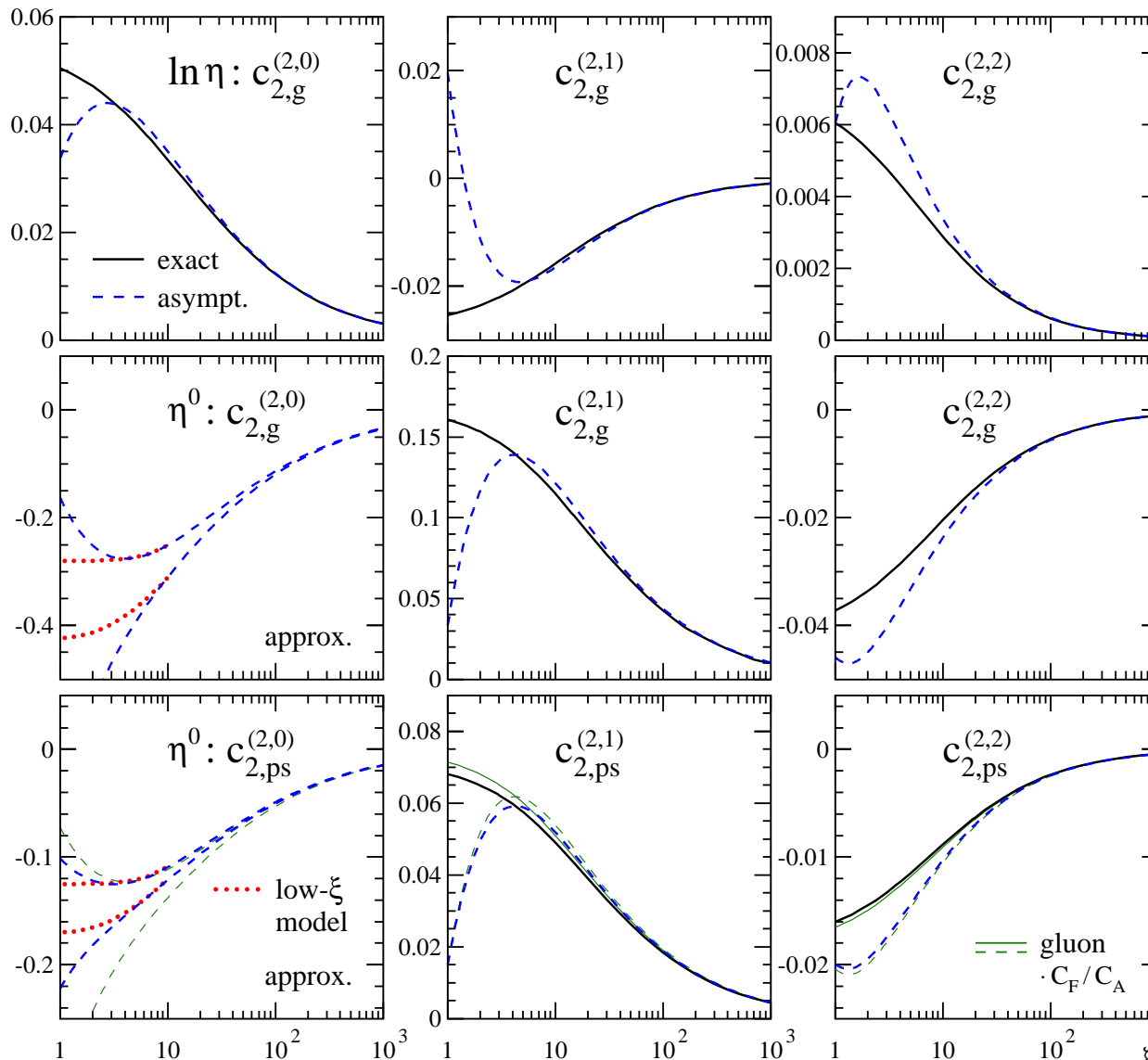
- Result valid for all ξ ;

- for $c_{LL_x}^{(2)}$ known exactly Catani, Ciafaloni, Hautmann '91

- for $c_{NLL_x}^{(2)}$ at large ξ from asymptotic factorization

High-energy limit

- Coefficients of large η (small x) expansion at NNLO



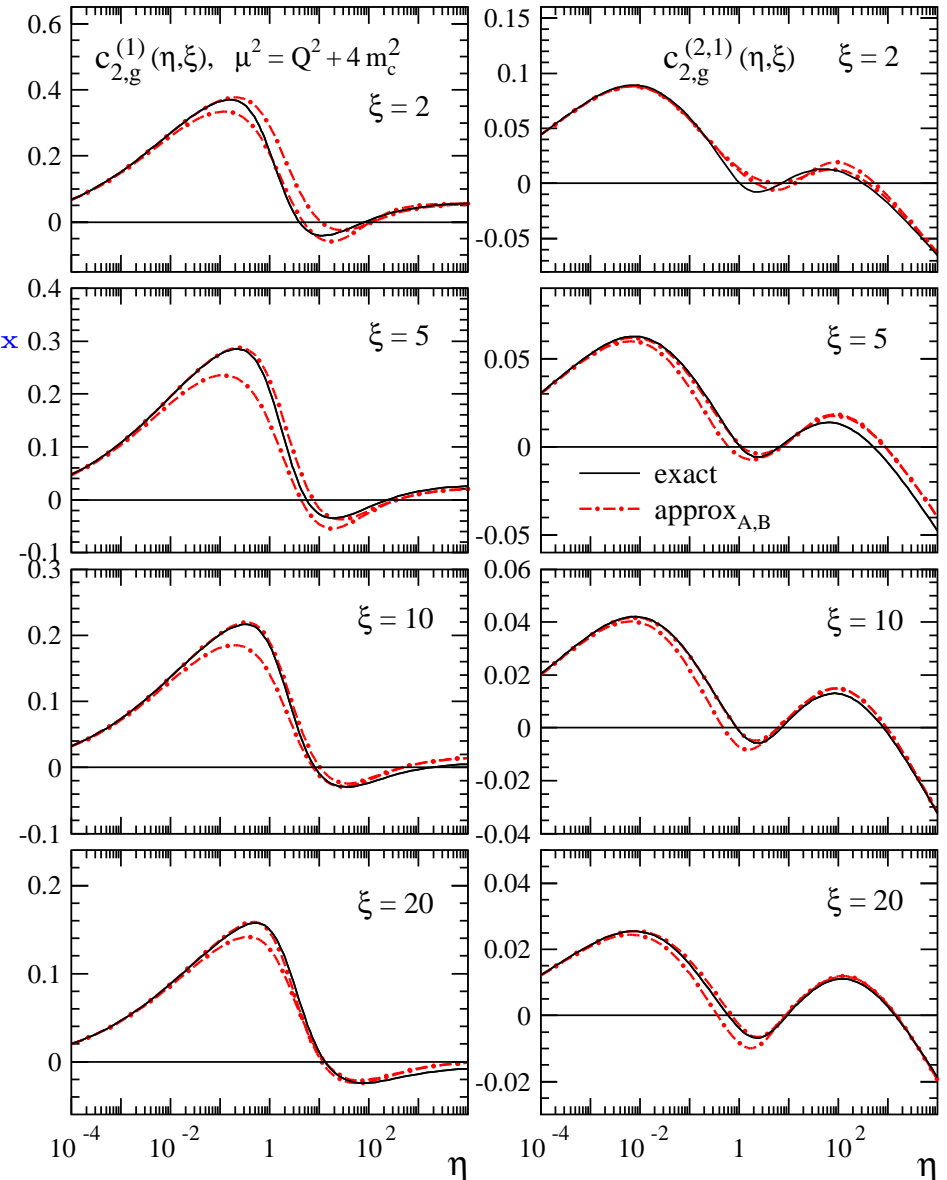
Quality check

- Combine all limits
- Construct approximate result $c_{2,k}^{\text{approx}}$ for full kinematic range

- merge $c_{2,k}^{\text{approx}} = c_{2,k}^{\text{thresh}} + c_{2,k}^{\text{LLx}} + \eta^\gamma / (C + \eta^\gamma) c_{2,k}^{\text{NLLx}}$

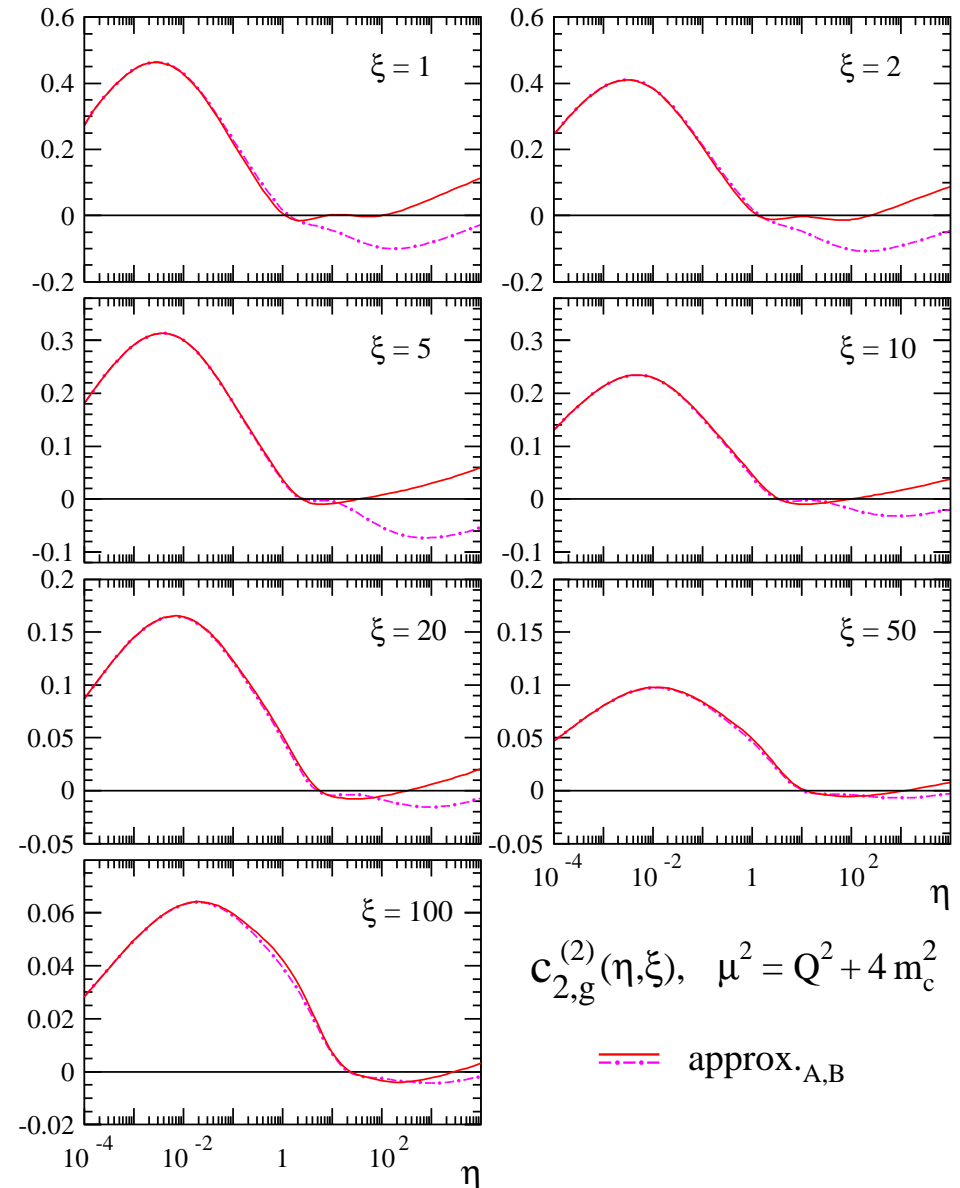
Tests

- at NLO for $c_{2,g}^{(1)}$
- at NNLO for $c_{2,g}^{(2,1)}$
(scale dependent part)



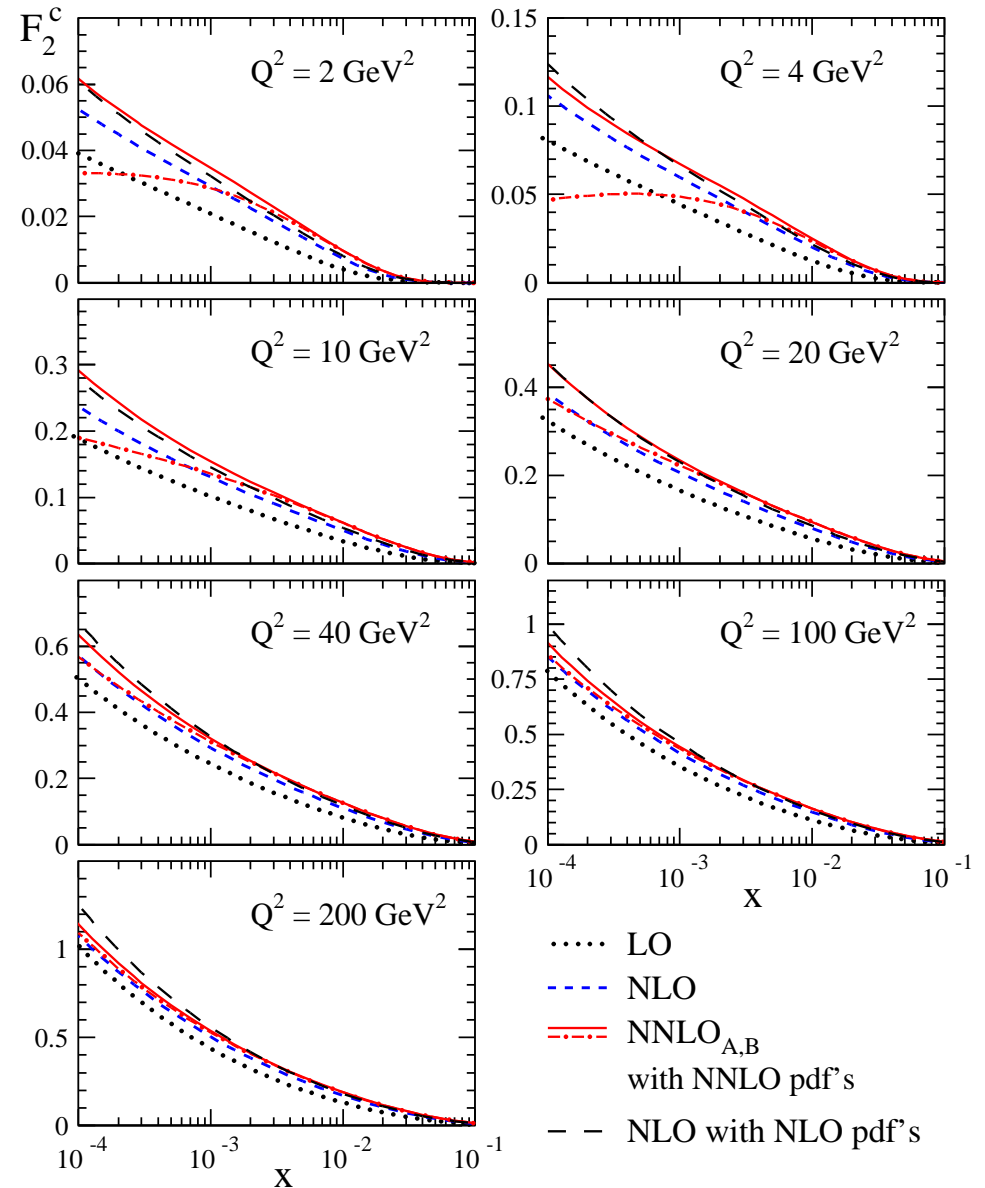
Approximate coefficient functions at NNLO

- New approximate NNLO result
 - very well constrained at large ξ
 - sizable uncertainties for $\xi \lesssim 10$
- Realistic estimate of uncertainty through combination of all limits
- Progress would require
 - more Mellin moments for A_{ji}
 - computation of $c_{2,k}^{\text{NLL}_x}$ (NLO impact factor for k_t -factorization)

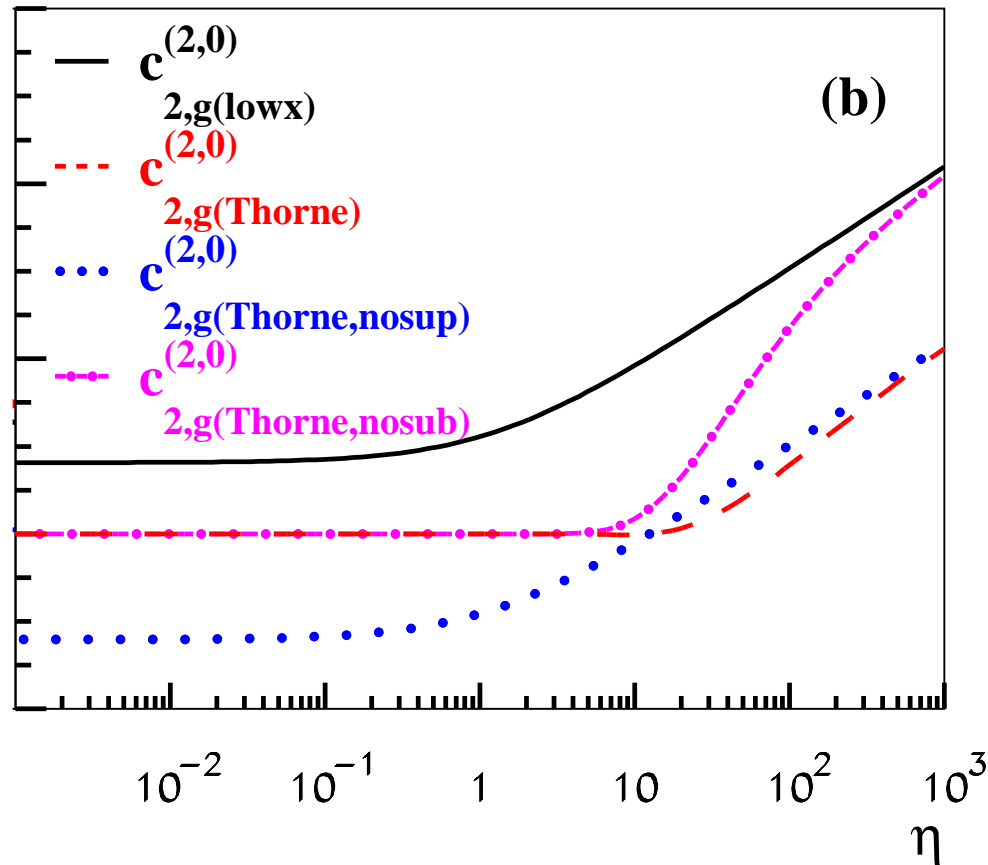


Approximate charm structure functions at NNLO

- New approximate NNLO result
 - convolution with ABM11 PDF set
- Uncertainty on NNLO dominates kinematic region in Q^2 relevant for HERA
 - very well constrained at large Q^2
 - sizable uncertainties for $Q^2 \lesssim 20 \text{ GeV}^2$ (even at not so small x)
 - gluon PDF does not fall fast enough as η grows larger



Charm coefficient functions at NNLO in TR' scheme



- Coefficient function $c_{2,g}^{(2)}$ at NNLO only modelled at high- x scale

Thorne '06

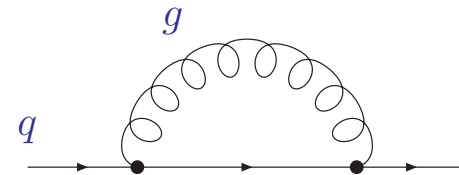
- no valid description in threshold region
- no uncertainty at high energies attached

Mass measurement

Pole mass

- Based on (unphysical) concept of heavy-quark being a free parton

$$\not{p} - m_q - \Sigma(p, m_q) \Big|_{p^2 = m_q^2}$$



- heavy-quark self-energy $\Sigma(p, m_q)$ receives contributions from regions of all loop momenta – also from momenta of $\mathcal{O}(\Lambda_{QCD})$
- Definition of pole mass ambiguous up to corrections $\mathcal{O}(\Lambda_{QCD})$

Running quark masses

- \overline{MS} mass definition $m(\mu_R)$ realizes running mass (scale dependence)
 - short distance mass probes at scale of hard scattering
 $m_{\text{pole}} = m_{\text{short distance}} + \delta m$
 - conversion between pole mass and \overline{MS} mass definition in perturbation theory: $m = m(\mu_R) \left(1 + a_s(\mu_R) d^{(1)} + a_s(\mu_R)^2 d^{(2)} \right)$

Quark masses in the PDF fits

		m_c/m_b (GeV)	
MSTW	DIS+DY+jets	1.40 / 4.75 (GMVFN)	EPJC 63, 189 (2009)
CTEQ	DIS+DY+jets	1.3 / 4.5 (GMVFN)	PRD 82, 074024 (2010)
NNPDF	DIS+DY+jets	$\sqrt{2}$ / 4.75 (GMVFN)	NPB 849, 296 (2011)
JR	DIS+DY	1.3 / 4.2 (FFN)	PRD 70, 074023 (2009)
HERAPDF	DIS	1.4 / 4.75 (GMVFN)	JHEP 1001, 109 (2010)
PDG:		1.66 / 4.79	

PDF fits assume pole mass scheme

The values of masses are systematically lower than those from PDG

In the GMVFN scheme the mass is used as a matching parameter as well

Additional uncertainty in the c.s. predictions

Alekhin

NNLO value of m_c from DIS (theory)

Alekhin

The NNLO* theoretical accuracy:

- the NNLO PDF evolution
- approximate NNLO Wilson coefficients with a linear interpolation between two margins A and B and interpolation coefficient d_N is a parameter of fit

$$A(1-d_N) + d_N B$$

The PDFs are fixed at the ABM11 NNLO central values with m_c -dependence included

- variation of PDFs within the ABM11 uncertainties makes little effect on the value of m_c

The experimental data are confronted to:

- the structure function $F_2^{\text{cc}}(m_c)$ (charm production c.s. $\sigma_{\text{full}}^{\text{cc}}$)
- the charm-production c.s with account of the visible phase space cuts $\sigma_{\text{full}}^{\text{cc}}(m_c)\varepsilon(m_c)$. The correction factor

$$\varepsilon(m_c) = \sigma_{\text{vis}}^{\text{cc}}(m_c) / \sigma_{\text{full}}^{\text{cc}}(m_c)$$

is calculated with the fully exclusive code HVQDIS; the NNLO contribution partially cancel in the ratio. The correction is sensitive to $m_c \rightarrow$ recalculated iteratively in the fit

Harris, Simth NPB 452, 109 (1995)

Fitting m_c (data)

The data on D^* c.s.

H1 Collaboration EPJC 71, 1769 (2011)

D^* cross section transformed into charm cross section using fragmentation function from [H1 Collaboration EPJC59:589(2009)], with code *hvqdispawutil* [K. Daum]

correction for the unmeasured region is necessary

consistent treatment of the F_L contribution

The data on structure function obtained with the shifted vertex

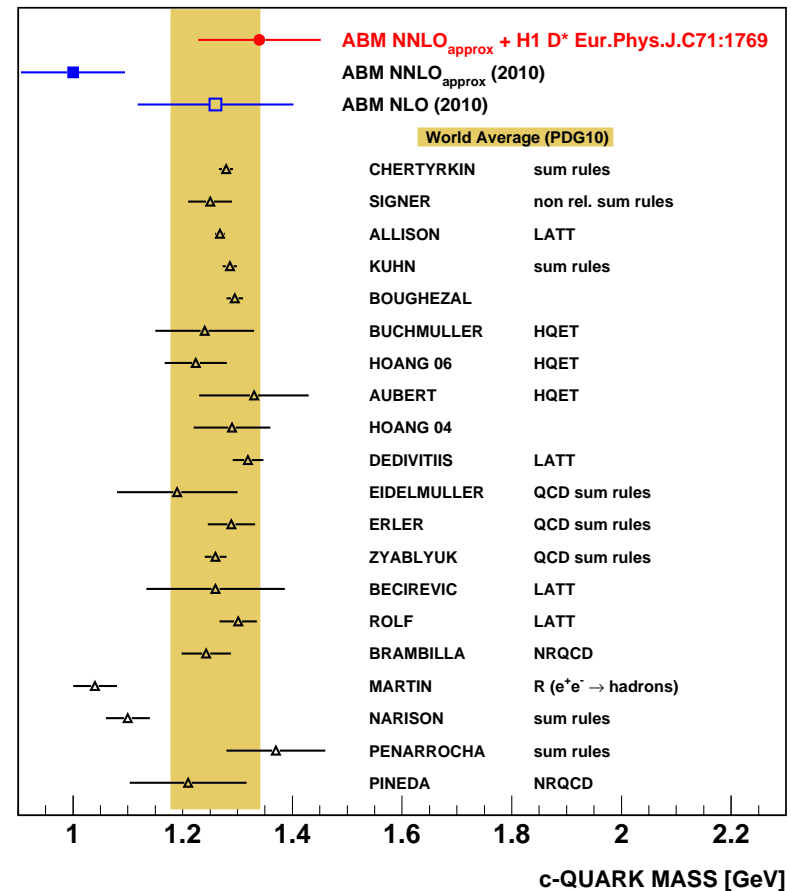
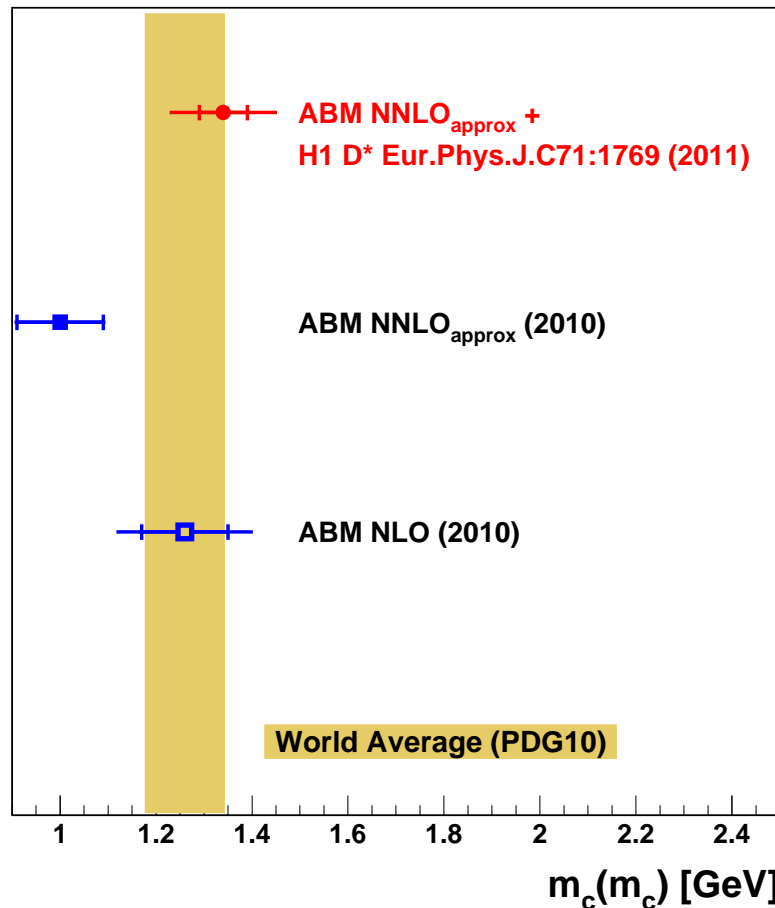
H1 Collaboration EPJC 65, 89 (2010)

inessential cut on $p_T \rightarrow$ the correction for the unmeasured region is small

Alekhin

Running charm mass from HERA

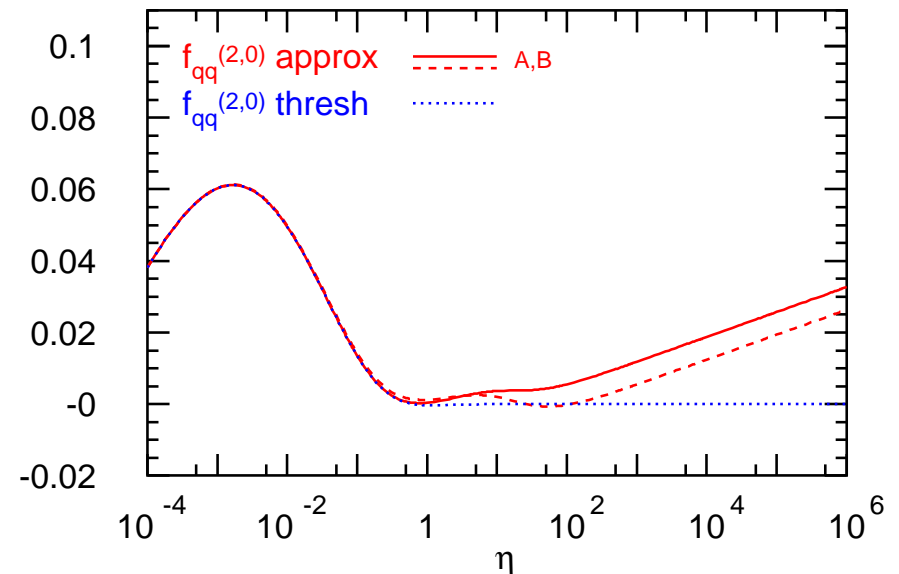
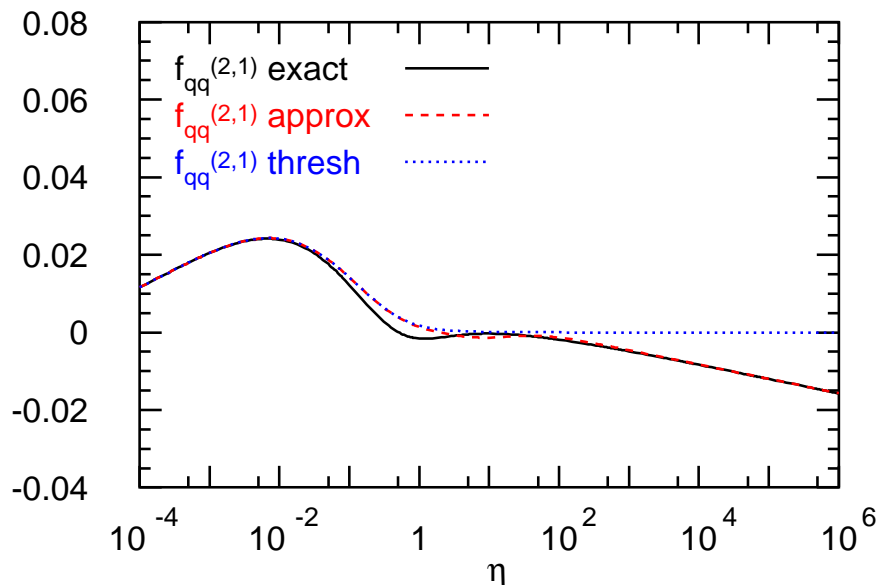
- Determination of $\overline{\text{MS}}$ -mass $m_c(\bar{m}_c)$ in DIS
 - significant impact of data on F_2^c from HERA II
 - ongoing work (Alekhin, Daum, Lipka, S.M.)



Top-pair hadro-production

- Spin-Off from work on heavy-quark DIS [S.M, Uwer, Vogt to appear](#)
- Threshold at $s \simeq 4m_t^2$ with logarithms $\ln(\beta)$ in velocity of heavy quark

$$\beta = \sqrt{1 - 4m_t^2/s}$$
 at n^{th} -order
[S.M, Uwer '08; Beneke, Czakon, Falgari, Mitov, Schwinn '09](#)
- High-energy limit for $\rho = 4m_t^2/s \rightarrow 1$
[Catani, Ciafaloni, Hautmann '91; Ball, Ellis '01](#)



Summary

Heavy-quark DIS

- Charm structure function F_2^c
 - progress on approximation at NNLO QCD
 - large uncertainty due to poorly constrained NLL_x term at NNLO

Heavy quark masses

- Determination of running charm-quark mass from DIS data
- Outlook:
 - tuning of ABM11 PDFs with updated Wilson coefficients
 - comparison to $m_c(m_c)$ at NLO scale variation uncertainty
 - impact of the combined HERA charm data

Top-quark hadro-production

- Improved approximate NNLO cross sections with small uncertainty