

# QCD resummation of jet mass distributions



Simone Marzani  
IPPP / Durham University



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*A. Banfi, M. Dasgupta, K. Khelifa-Kerfa and M. Spannowsky*

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# Jet sub-structure

- ◆ Understanding the substructure of jets is crucial for LHC phenomenology
- ◆ It is important for new physics searches
  - ◆ we need tools to distinguish jets coming from decays of boosted resonances from QCD jets
- ◆ It has also interest on its own as a test of pQCD
- ◆ Jet shapes enable us to look at energy distributions inside a jet
- ◆ Grooming techniques enable us to improve our ability of distinguishing signal from background by getting rid of soft radiation

◆ e.g. Butterworth et al (2008), Ellis et al. (2009), Krohn et al. (2010) and several others

# The Jet Mass

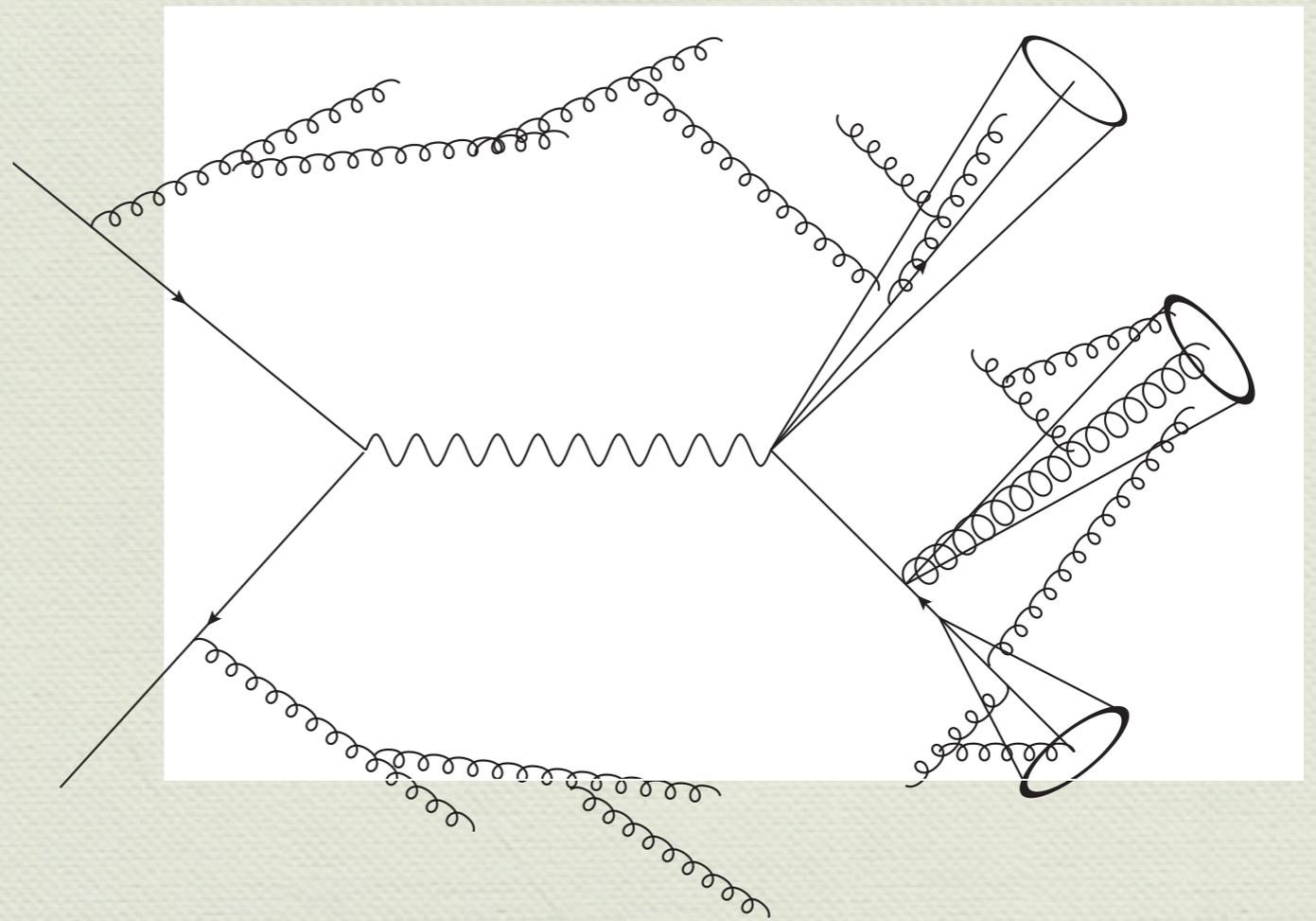
- ◆ In our studies we concentrate on the jet-mass as an example of a jet shape
- ◆ Jet mass distributions are affected by large logarithms

$$\frac{1}{\sigma} \frac{d\sigma}{dm_J^2} = \frac{1}{m_J^2} \left[ \alpha_s A_1 \ln \frac{m_J^2}{p_T^2} + \alpha_s^2 A_2 \ln^3 \frac{m_J^2}{p_T^2} + \dots \right]$$

- ◆ Reliable calculations in pQCD require resummation
- ◆ In  $e^+e^-$  resummation exists CTTW (1993), Burby and Glover (2001)
- ◆ Can we extend this to hadron colliders ?

# The picture at the LHC

- ◆ Can we deal with a more complicated environment ?
- ◆ Complex colour structure and geometry ?
- ◆ Does the choice of the jet algorithm play a role ?
- ◆ How do we treat ISR ?
- ◆ Can we reduce the underlying event ?



# MC vs Analytical Approach

- ◆ MC simulations using parton showers
  - ◆ they are powerful general purpose tools
  - ◆ they are often interfaced with hadronisation and UE models to give a realistic description
  - ◆ they are formally LL (although they contain many sub-leading terms)
- ◆ Analytical resummations
  - ◆ well defined and improvable accuracy (LL, NLL, etc)
  - ◆ the dependence on various parameters is manifest
- ◆ The two approaches are complementary

# One gluon calculation

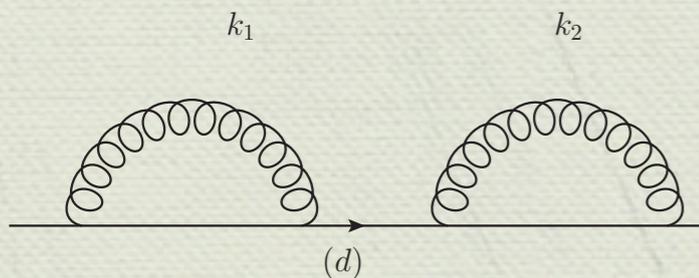
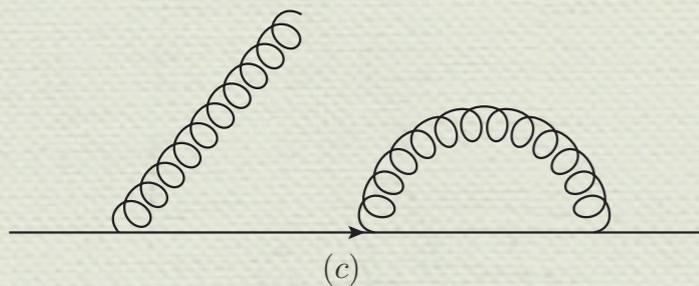
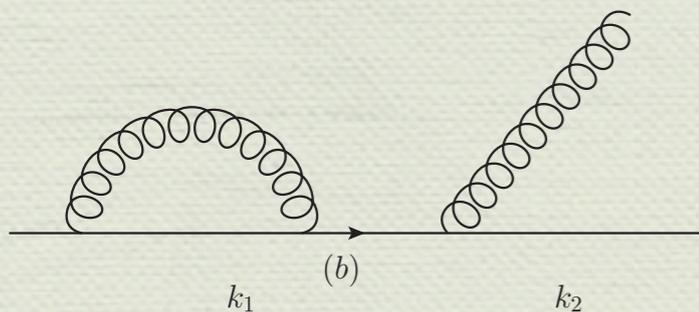
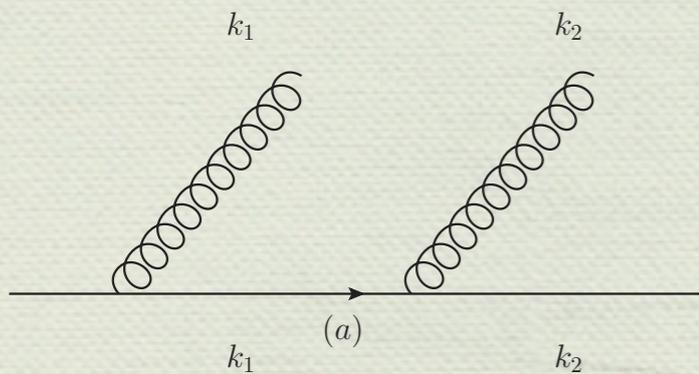
- ◆ We wish to study the mass of one or more high- $p_T$  jets
- ◆ We follow a traditional approach to QCD resummation
- ◆ There are many results also from SCET (e.g. Horning, Lee, Stewart, Walsh and Zuberi; Schwartz et al.)
- ◆ We consider the emission of one soft gluon off a quark-antiquark dipole (as in  $e^+e^-$ )
- ◆ At this level all IRC jet algorithm will give the same answer

$$\frac{1}{\sigma} \frac{d\sigma^{(1)}}{d\rho} = \frac{d}{d\rho} \left[ \frac{\alpha_s}{2\pi} C_F \ln^2 \frac{R^2}{\rho} \right], \quad \rho = \frac{m_J^2}{p_T^2}$$

- ◆ The naïve expectation is that the resummed result is obtained by exponentiation of the 1-loop contribution

# Two gluons ... $n$ gluons (?)

$E_1 \gg E_2$



- ◆ Let's look at the two-gluon calculation
- ◆ Double virtual (d) never contributes
- ◆  $k_1$  in,  $k_2$  out: (a) and (c) cancel; (b) is zero
- ◆  $k_1$  out,  $k_2$  in: (a) and (b) cancel; (c) is zero
- ◆ Non-vanishing contributions only when both gluons are inside the jet:

$$\frac{1}{\sigma} \frac{d\sigma^{(2)}}{d\rho} = \frac{d}{d\rho} \left[ \frac{1}{2} \left( \frac{\alpha_s}{2\pi} \right)^2 C_F^2 \ln^4 \frac{R^2}{\rho} \right]$$

- ◆ Consistent with exponentiation
- ◆ Is this result the same with any jet algorithm ?

# The role of the jet algorithm

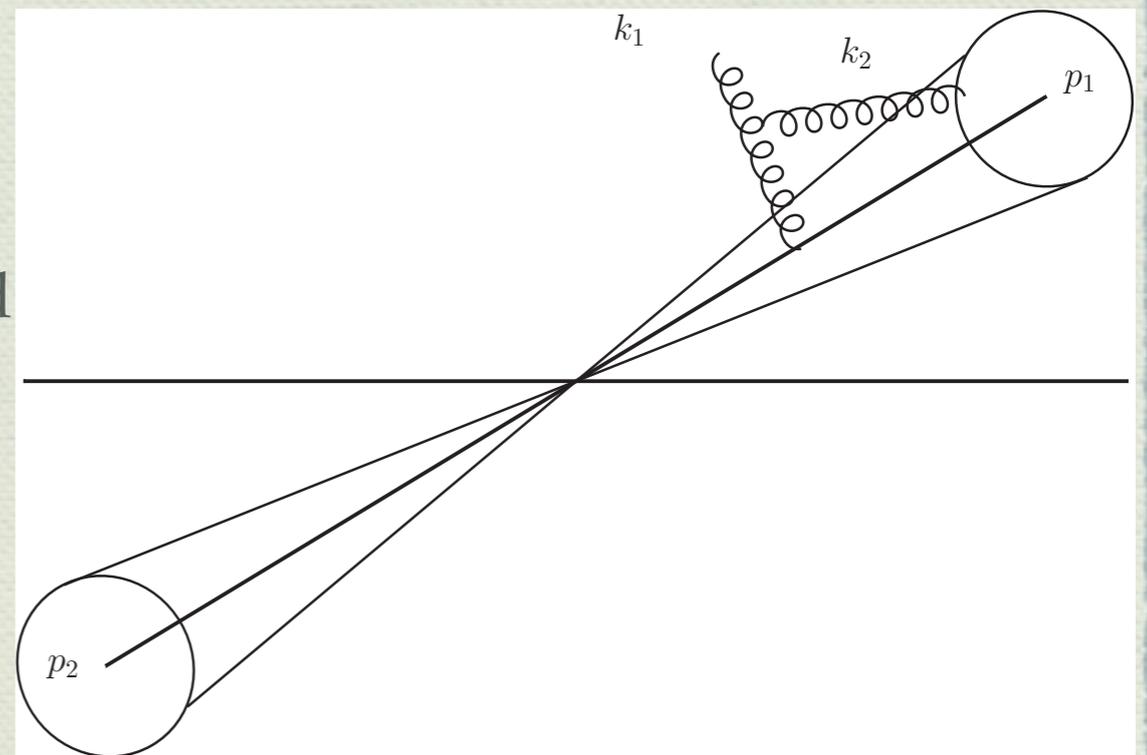
- ◆ **NO!** Because of soft gluon self-clustering
- ◆ Two gluons are recombined with each other if their distance is smaller than  $R$
- ◆ The recombined momentum essentially lies along the harder one
- ◆ As a result a hard gluon can pull a softer one out of the jet
- ◆ This does not happen if we use anti- $k_t$  algorithm: two soft gluons are always far apart with this measure
- ◆ The anti- $k_t$  algorithm in the soft limit works as a perfect cone
- ◆ Exponentiation of the independent emission is OK in this case, but what happens with different algorithms ?

# C/A and $k_t$ algorithm

- ◆ With C/A or  $k_t$  algorithm two soft gluons can be recombined together
- ◆ If  $k_1$  is outside the jet, it can pull  $k_2$  outside (if they are close enough)
- ◆ This spoils the cancellation between diagrams (a) and (b) (virtual gluons are unaffected by clustering)
- ◆ This results into a new single-log contribution wrt to the anti- $k_t$  case
- ◆ Resummation possible for gaps between jets Delenda, Appleby, Banfi, Dasgupta (2006)
- ◆ These effects received recently a lot of attention Kelley, Walsh and Zuberi (2012)
- ◆ From now on we work with the anti- $k_t$  algorithm

# Non-global logarithms

- ◆ BUT, even if we use anti- $k_t$ , exponentiation of the independent emission is not the whole story
- ◆ The jet-mass is a non-global observable: it receives single log corrections from correlated emission
- ◆ This is a  $C_F C_A$  term and it's missed by single gluon exponentiation
- ◆ In principle we need to consider any number of gluons outside the jet
- ◆ Colour structure becomes intractable, so the resummation is performed in the large  $N_c$  limit



Dasgupta and Salam (2001)

# From $e^+e^-$ to $pp$ collisions

- ◆ So far we have discussed the resummation for a quark-antiquark dipole (appropriate for  $e^+e^-$ )
- ◆ In the soft limit, the emission probability for a gluon off an ensemble of hard partons can be written as

$$|\mathcal{M}|^2 = \mathcal{M}_0^\dagger \exp \left[ \frac{\alpha_s}{2\pi} \left( -2 \sum_{i < j} W_{ij} t_i \cdot t_j \right) \right] \mathcal{M}_0$$

- ◆ The sum runs over all the possible dipoles in the hard scattering
- ◆ Dipoles involving the measured jet(s) will produce double logs (soft and collinear gluons), while the other ones only single logs (soft gluons at large angle)

# The NLL resummed result

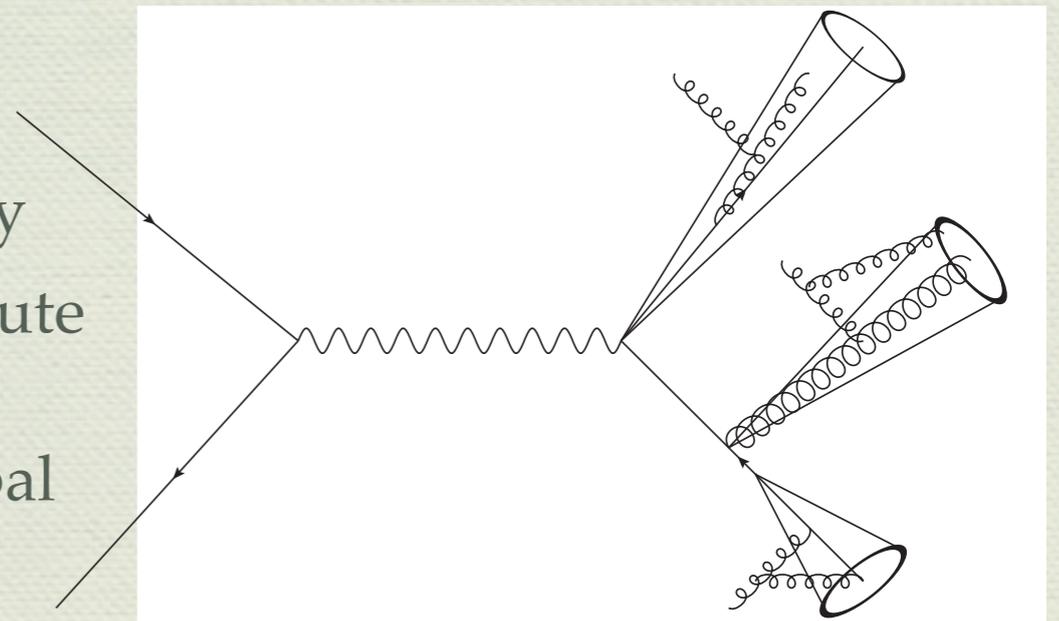
- ◆ For each partonic sub-process we can write a NLL resummed expression
- ◆ Once a basis is specified the colour factors are represented by matrices

$$\frac{1}{\sigma} \frac{d\sigma}{d\rho} = \frac{d}{d\rho} \text{tr} \left[ H \left( 1 + \frac{\alpha_s}{2\pi} C_1 \right) e^{-\mathcal{R}^\dagger} S e^{-\mathcal{R}} \right]$$

- ◆ H contains the Born contributions and  $S = 1$  for our choice of bases
- ◆ The radiator is also a matrix in colour space. To NLL it contains
  - ◆ soft and / or collinear radiation
  - ◆ two-loop running coupling
  - ◆ non-global logs in the large  $N_c$  limit

# Small R and beyond

- ◆ In the small-R limit a simpler picture emerges
- ◆ Large-angle emissions are R-suppressed and only the dipoles involving the measured jet(s) contribute
- ◆ Each jet evolves independently and the non-global part is the same as in the hemisphere case
- ◆ Finite R-corrections are at the percent level if  $R < 0.4$ , but they reach  $O(40\%)$  for  $R=1.0$
- ◆ It turns out that they are dominated by  $O(R^2)$  terms, while  $O(R^4)$  are below 1% (even at  $R=1.0$ )



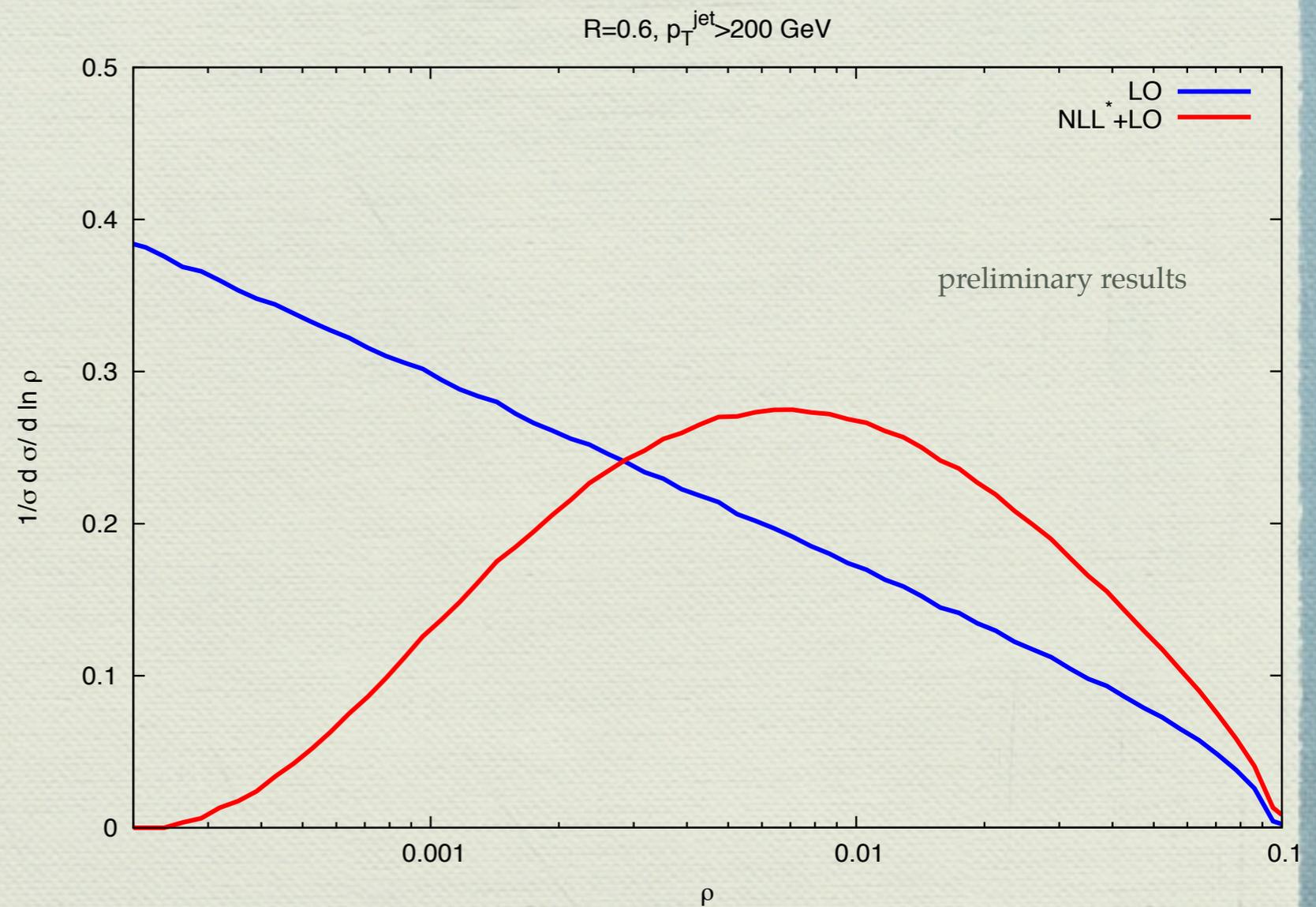
- ◆ For phenomenological studies we then keep  $O(R^2)$

# An example: Z + jet

- ◆ We study in detail the mass distribution of a QCD jet produced in association with a Z boson
- ◆ Background to associated Higgs production with a Z
- ◆ If the Higgs is boosted its decay products are likely to end up in a fat jet  
[Butterworth et al. \(2008\)](#)
- ◆ Interesting correlations between grooming techniques found in this channel  
[Soper and Spannowsky \(2010\)](#)
- ◆ QCD resummation
  - ◆ colour structure is relative simple because only 3 partons in the Born
  - ◆ we are differential in the jet (or Z boson) kinematics

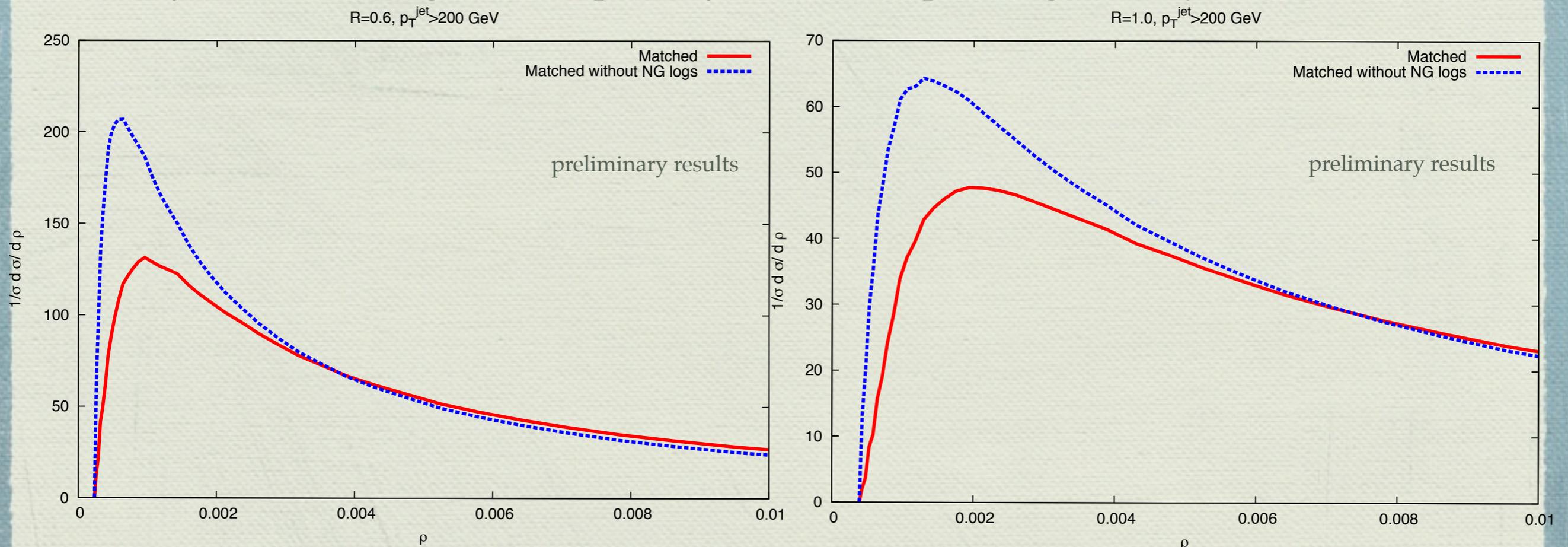
# Matching: $NLL^* + LO$

- ◆ We perform an (almost)  $NLL$  resummation (without the coefficient  $C_1$ )
- ◆ We match it a fixed-order computation of the mass distribution (at LO)
- ◆ A complete pheno study will require  $C_1$  and NLO matching together with a faithful estimate of the perturbative uncertainties



# Effect of NG logs

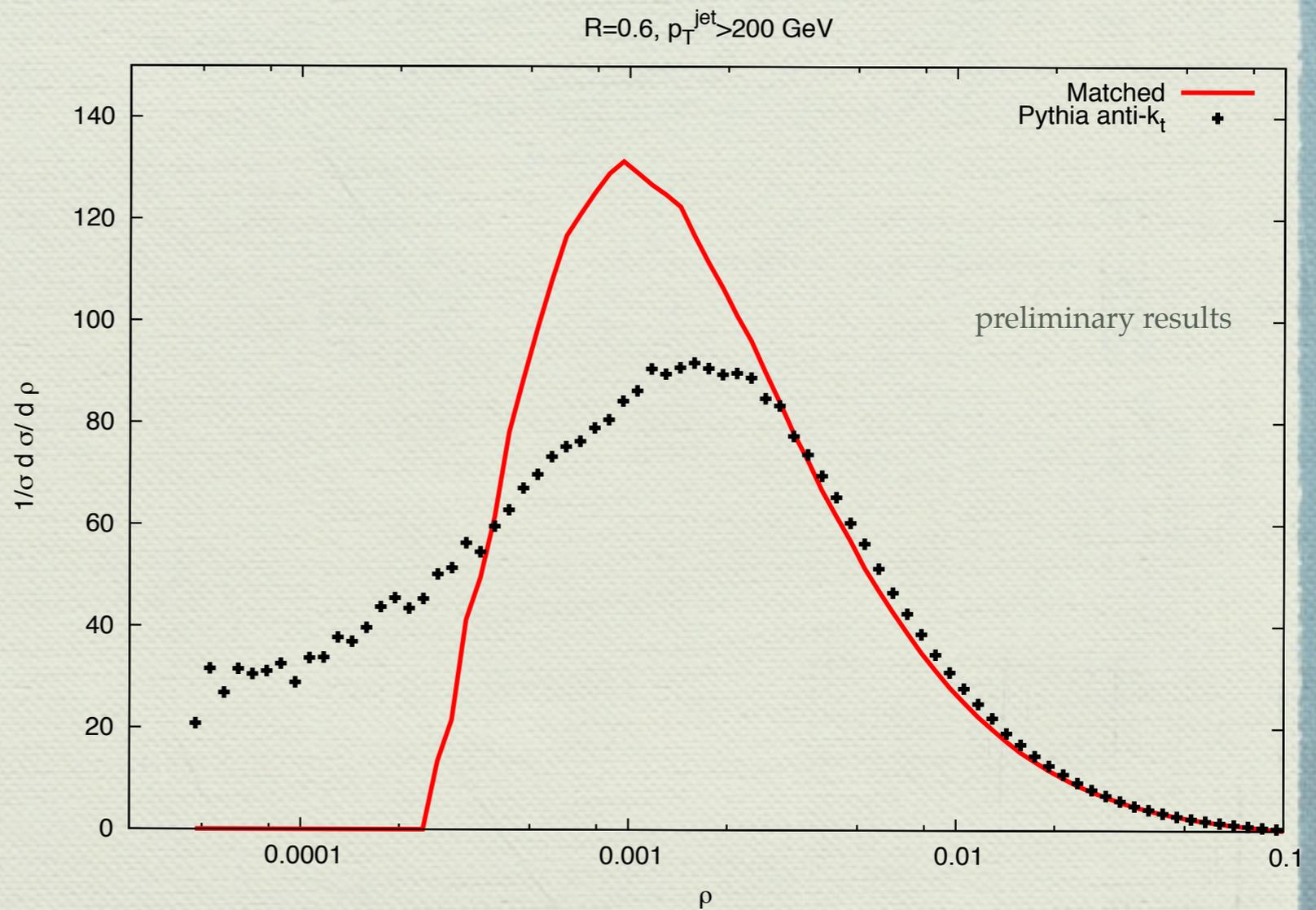
- ◆ Non-global logs have a significant impact on the peak of the distribution
- ◆ They lower the height of the peak by  $O(35\%)$ , depending on the jet radius



- ◆ These effects were not considered in previous studies and by fairly large NP effects were introduced

# Comparison to MC

- ◆ We compare our  $NLL^* + LO$  result to Pythia8 (shower only)
- ◆ Still a difference in the peak height
- ◆ Actually the resummation and the shower populate the low mass region quite differently (peak  $\sim 6$  GeV for  $p_T=200$  GeV)
- ◆ It would be interesting to go to higher  $p_T$  and move away from the NP region



sharp cut-off because of Landau pole

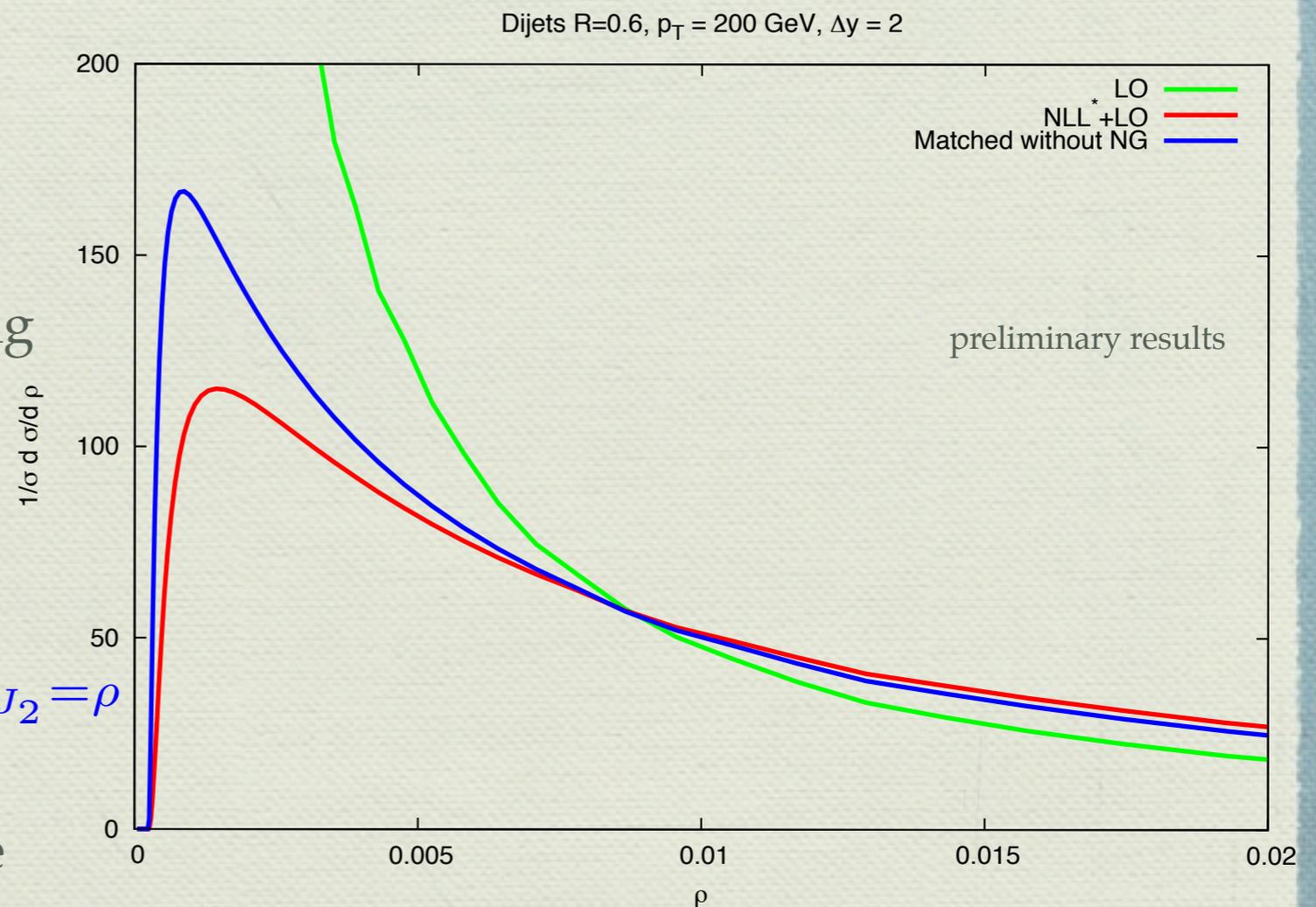
# Another example: dijets

◆ We also study the jet-mass distribution in dijets

◆ We measure the mass of the leading jets

$$\frac{d\sigma}{d\rho} = \left[ \frac{d\sigma}{d\rho_{J_1}} + \frac{d\sigma}{d\rho_{J_2}} \right]_{\rho_{J_1} = \rho_{J_2} = \rho}$$

◆ More complicated colour structure



# Analytic studies of grooming

- ◆ We are trying to better understand the discrepancy with MC
- ◆ Precise phenomenology will require full NLL+NLO to reduce perturbative uncertainties
- ◆ We are using analytic methods to study and compare different grooming techniques (filtering, pruning and trimming)
- ◆ This study will shed light on relations and differences between these procedures

# Conclusions

- ◆ I have presented analytical predictions for jet mass distributions in hadron collisions
- ◆ I have discussed on the different effects one should consider to perform a NLL resummation, with particular emphasis on NG logs
- ◆ I have applied this formalism to two processes, Z+jet and dijets
- ◆ We are comparing our results to standard MC in different kinematical regions