

XX International Workshop on Deep-Inelastic Scattering and Related Subjects



26-30 March 2012, University of Bonn







University & INFN of Bologna

on Behalf of the CMS Collaboration

Motivations

- Heavy quarkonia are an excellent laboratory for understanding QCD
 - non-relativistic due to their high mass
 - nonperturbative effects can be simplified and constrained
- In the last decade, significant progress for production mechanisms
 - new experimental results
 - improved theoretical descriptions
- Definitive understanding still a challenge, several models competing for confirmation
- Renewed interest in quarkonium spectroscopy since the discovery of the XYZ exotic states:
 - search of new possible states
 - new measurements needed to understand their true nature

CMS Quarkonium Studies

Υ production cross section L_{int} = 3 pb⁻¹ *Phys. Rev. D* 83, 112004 (2011)

Prompt and non-prompt J/ ψ production L_{int} = 314 nb⁻¹ Eur.Phys. J C71, 1575 (2011)

 J/ψ and ψ (2S) production L_{int} = 37 pb⁻¹ JHEP 02, 011 (2012)

Observation of the \chi_c states L_{int}= 1.1 fb⁻¹ CERN-CMS-DP-2011-011

Measurement of the production cross section ratio of X(3872) and ψ (2S) L_{int} =40 pb⁻¹ CMS-PAS-BPH-10-018



Muons in CMS

- Quarkonium states identified in final states with di-muons
- Muon system information matched to an inner-tracker track for improved momentum resolution
- Inner Tracker:
 - Silicon pixel and strip layers
 - High p_T resolution ~1%
 - Excellent vertex reconstruction and impact parameter resolution
- Muon System
 - 3 types of gaseous detectors
 - Phase space coverage up to |η| = 2.4
 - Highly efficient muon trigger and identification
 - Resolution η dependent.



Muon Triggers



S-wave quarkonium states

- Unbinned Maximum Likelihood fit to μ⁺μ⁻ invariant mass distributions
- Signals modeled with Crystal Ball functions
- Mass differences are fixed from PDG, common resolution value (scaled by mass)
- Yields then corrected for Acceptance (from MC) and Efficiency (from data-driven methods)

 muon acceptance is strongly dependent on production polarization





Inclusive J/\psi Production $\frac{d^{2}\sigma}{dp_{T}dy}(J/\psi) \cdot \mathcal{B}(J/\psi \rightarrow \mu^{+}\mu^{-}) = \frac{N_{J/\psi}^{corr}(p_{T}, |y|)}{\int L dt \cdot \Delta p_{T} \cdot \Delta y}$ The cross section comprises 3 inclusive J/ $\psi \rightarrow \mu^{+}\mu^{-}$, corrected for acceptance



Non prompt Fraction

SV

Fraction coming from b-decay extracted with a 2-Dimensional fit of invariant mass and "pseudo-proper" decay length

$$L_{J/\psi} = rac{L_{xy} \cdot m_{J/\psi}}{p_T}$$

- I_{J/ψ} distribution components: prompt → Resolution function

 - non prompt \rightarrow Resolution function convoluted with exponential
 - background

 Pre-fitted in mass sidebands
- Decay length resolution described by "per-event uncertainty" on $I_{J/\psi}$
- Results in agreement with CDF and ATLAS







- Excellent agreement with NRQCD predictions
 - For prompt J/ ψ , feed down effect included in theory
 - $\hfill remarkable for \psi$ (2S) in absence of feed-down
- Typical uncertainties ~5 [20]% on J/ ψ [ψ (2S)] cross-sections
- Polarization uncertainty studied in 4 "extreme" scenarios, effects up to 20[30]% for J/ ψ [ψ (2S)]

 $B \rightarrow \psi(nS)X$ cross sections

Comparison with FONLL predictions:



- For J/ψ:
 - agreement below 30 GeV
 - above 30 GeV FONLL overestimate data



- For ψ (2S):
 - Shape agreement in the measurement range
 - Uniform scale discrepancy found
 - improved determination of BR

ψ (2S) to J/ ψ cross-sections ratios

Cross Section Ratio calculation:

- Systematic uncertainties largely cancel (Luminosity, Single Muon Efficiencies...)
- Direct production with same polarization
 - Residual polarization effect from J/ $\psi\,$ coming from feed-down
- No |y| dependence seen, results as function of p_T

• B $\rightarrow \psi$ (2S) X Branching Fraction

 measured fitting the non-prompt cross-section ratio with FONLL or EvtGen curves

 $\mathcal{B}(\mathrm{B}
ightarrow \psi(\mathrm{2S})X) = (3.08 \pm 0.12(\mathrm{stat.+syst.}) \pm 0.13(\mathrm{theor.}) \pm 0.42(\mathcal{B}_{\mathrm{PDG}})) \cdot 10^{-3}$

- In agreement with world average (4.8±2.4)·10^{−3}
 - Improving relative uncertainty by factor 3
 - main uncertainties from other PDG BRs



Y(nS) Cross Sections

$$\begin{split} &\sigma(\mathrm{pp} \to \mathrm{Y}(1\mathrm{S})X) \cdot \mathcal{B}(\mathrm{Y}(1\mathrm{S}) \to \mu^+\mu^-) = 7.37 \pm 0.13(\mathrm{stat.})^{+0.61}_{-0.42}(\mathrm{syst.}) \pm 0.81(\mathrm{lumi.}) \, \mathrm{nb} \\ &\sigma(\mathrm{pp} \to \mathrm{Y}(2\mathrm{S})X) \cdot \mathcal{B}(\mathrm{Y}(2\mathrm{S}) \to \mu^+\mu^-) = 1.90 \pm 0.09(\mathrm{stat.})^{+0.20}_{-0.14}(\mathrm{syst.}) \pm 0.24(\mathrm{lumi.}) \, \mathrm{nb} \\ &\sigma(\mathrm{pp} \to \mathrm{Y}(3\mathrm{S})X) \cdot \mathcal{B}(\mathrm{Y}(3\mathrm{S}) \to \mu^+\mu^-) = 1.02 \pm 0.07(\mathrm{stat.})^{+0.11}_{-0.08}(\mathrm{syst.}) \pm 0.11(\mathrm{lumi.}) \, \mathrm{nb} \end{split}$$

- Υ(1S) and Υ(2S) include feed-down from higher-mass states
- Unpolarized Y(nS) assumption
 - Extreme polarization change cross sections by about 20%



Y(1S) Cross Section



Consistent shape to PYTHIA

 PYTHIA overestimates the integrated cross section by a factor 2 Results compared to D0 and CDF measurements

 Assuming cross section uniform in rapidity an increase by a factor 3 is observed at √s = 7 TeV Differential cross-section vs rapidity:

 slight decrease towards |y|=2 consistent with PYTHIA

$\chi_{cJ} \rightarrow J/\psi \gamma$ mass distribution



The X(3872) state

- Discovered in 2003 by Belle \rightarrow its nature still unclear
- A clear signal is established in 2010 in the J/ ψ π^+ π^- decay channel

• Starting from reconstructed J/ ψ

- Searched pair of compatible good quality opposite-charged tracks in ΔR(π,J/ψ)<0.7
- Performed 4-track vertex fit with J/ ψ mass fixed to the PDG value
- Kept good quality candidates in the kinematic region

 $p_T(X) > 8 \text{ GeV and } |y(X)| < 2.2$

Unbinned maximum likelihood fit

- $m_{\Psi(2S)} = 3685.9 \pm 0.1$ (stat. only) MeV
- m_{X(3872)} = 3870.2 ± 1.9 (stat. only) MeV
- $\sigma \hat{1}_{\Psi(2S)} = 8.1 \pm 0.6 \text{ MeV}$
- $\sigma 2_{\Psi(2S)} = 3.3 \pm 0.3 \text{ MeV}$
- $\sigma_{X(3872)} = 6.3 \pm 1.3 \text{ MeV}$

PDG values

 $\begin{array}{l} m_{\Psi(2S)} = 3686.09 \pm 0.04 \mbox{ MeV} \\ m_{X(3872)} = 3871.57 \pm 0.25 \mbox{ MeV} \end{array}$

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X(3872) to ψ (2S) inclusive cross section ratio



- Acceptance and efficiency correction from simulation are applied on the yields extracted from the mass spectrum
 - Pythia 6 with mass of χ_{c1} (J^{PC}=1⁺⁺) set to 3.872 GeV
 - Null polarization assumed
 - 30% non-prompt fraction assumed
- **Ratio results**

$R = 0.087 \pm 0.017$ (stat.) ± 0.009 (syst.)



Conclusions

CMS has issued several studies on heavy quarkonia with the first LHC data:

- Measurement of J/ψ cross section from 0 to 70 GeV/c with large rapidity coverage (|y|<2.4)
- Differential cross-sections in p_T and |y| of J/ ψ and ψ (2S) mesons
 - prompt and non-prompt contributions separated
 - compatible results to NRQCD prediction up to 30 GeV/c for prompt production
 - uniform scale discrepancy found and explained for non-prompt $\psi(\text{2S})$ production w.r.t. FONLL
 - consistent results with other LHC experiments
 - improved relative uncertainty for BR(B->ψ(2S) X) of a factor 3
- Differential cross-sections in p_T for Y(nS) states
 - shape compatible to PYTHIA and results at Tevatron
- First measurement for the X(3872) to ψ (2S) cross section ratio
- χ_{cJ} peaks resolved in their radiative decay to J/ ψ , using converted photons



The CMS detector



Acceptances





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ψ (nS) Cross-section overview

$$\frac{\mathrm{d}^2\sigma}{\mathrm{d}p_{\mathrm{T}}\mathrm{d}y}(J/\psi)\cdot\mathcal{B}(J/\psi\to\mu^+\mu^-) = \frac{N_{J/\psi}^{\mathrm{corr}}(p_{\mathrm{T}},|y|)}{\int L\,\mathrm{d}t\cdot\Delta p_{\mathrm{T}}\cdot\Delta y}$$

 N_{fit} = signal yield from fit to $\mu\mu$ invariant mass $\int Ldt = integrated luminosity (4% uncertainty)$

- - Strongly dependent on production polarization
- $\boldsymbol{\varepsilon}$ = dimuon efficiency = $\varepsilon(\mu^+) \varepsilon(\mu^-) \rho \varepsilon_{\text{vertex}}$
- single muon trigger and reconstruction $\epsilon(\mu)$, data-driven via Tag & Probe
- vertexing of opposite sign dimuons (Prob>1%)

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- selection based on high quality tracks associated to muon segments: cuts on n_{hits} , χ^2 , $|d_{xy}|$, $|d_z|$
- ρ express the correlation between the two μ efficiencies

o trigger settings remove too close μ (to reduce single μ faking double μ), inducing sizable correlations \rightarrow Offline rejection of forward muons bending towards each other

A = geometrical and kinematical acceptance $|\eta^{\mu}| < 1.2 \rightarrow p_T > 4 \text{ GeV/c}$ • Strongly dependent on production polarization $1.2 < |\eta^{\mu}| < 2.4 \rightarrow p_T > 3.3 \text{ GeV/c}$

ψ (nS) Cross Sections



ψ (nS) Integrated Cross Sections

J/ψ

 $8.0 < p_T < 70.0 \text{ GeV/}c \text{ for } |y| < 0.9$ $8.0 < p_T < 45.0 \text{ GeV/}c \text{ for } 0.9 < |y| < 1.2$ $6.5 < p_T < 45.0 \text{ GeV/}c \text{ for } 1.2 < |y| < 1.6$ $6.5 < p_T < 30.0 \text{ GeV/}c \text{ for } 1.6 < |y| < 2.1$ $5.5 < p_T < 30.0 \text{ GeV/}c \text{ for } 2.1 < |y| < 2.4$ Corrected for acceptance:

ψ**(2S)**

 $6.5 < p_{\rm T} < 30.0 {
m ~GeV}/c {
m ~for} |y| < 1.2$ $5.5 < p_{\rm T} < 30.0 {
m ~GeV}/c {
m ~for} 1.2 < |y| < 2.4$

 $\begin{aligned} \mathcal{B}(J/\psi \to \mu^+ \mu^-) \cdot \sigma(\text{pp} \to \text{prompt } J/\psi) &= 54.5 \pm 0.3 \pm 2.3 \pm 2.2 \text{ nb} \\ \mathcal{B}(J/\psi \to \mu^+ \mu^-) \cdot \sigma(\text{pp} \to \text{b}X \to J/\psi X) &= 20.2 \pm 0.2 \pm 0.8 \pm 0.8 \text{ nb} \\ \mathcal{B}(\psi(2\text{S}) \to \mu^+ \mu^-) \cdot \sigma(\text{pp} \to \text{prompt } \psi(2\text{S})) &= 3.53 \pm 0.26 \pm 0.32 \pm 0.14 \text{ nb} \\ \mathcal{B}(\psi(2\text{S}) \to \mu^+ \mu^-) \cdot \sigma(\text{pp} \to \text{b}X \to \psi(2\text{S})X) &= 1.47 \pm 0.12 \pm 0.13 \pm 0.06 \text{ nb} \\ \text{Uncorrected for acceptance:} \end{aligned}$

$$\begin{split} \mathcal{B}(J/\psi \to \mu^+ \mu^-) \cdot \sigma(\text{pp} \to \text{prompt } J/\psi) &= 9.83 \pm 0.03 \pm 0.38 \pm 0.39 \text{ nb} \\ \mathcal{B}(J/\psi \to \mu^+ \mu^-) \cdot \sigma(\text{pp} \to \text{b}X \to J/\psi X) &= 4.67 \pm 0.02 \pm 0.17 \pm 0.19 \text{ nb} \\ \mathcal{B}(\psi(2\text{S}) \to \mu^+ \mu^-) \cdot \sigma(\text{pp} \to \text{prompt } \psi(2\text{S})) &= 0.410 \pm 0.009 \pm 0.023 \pm 0.016 \text{ nb} \\ \mathcal{B}(\psi(2\text{S}) \to \mu^+ \mu^-) \cdot \sigma(\text{pp} \to \text{b}X \to \psi(2\text{S})X) &= 0.235 \pm 0.006 \pm 0.013 \pm 0.009 \text{ nb} \\ \end{split}$$

B fraction results

• Above $p_T \approx 20$ GeV, more than 50% of the J/ ψ and $\psi(2S)$ mesons result from B decays

