

# Multi Vector Boson Production via Gluon fusion @ LHC

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# Outline

- 1 Introduction
- 2 Details of Calculation
- 3 Numerical Checks and Results
- 4 Summary

- SM predictions at higher energies and TeV-scale physics.
- background counting for signal.
- large gluon luminosity at the LHC.

- We are interested in gluonic contribution to

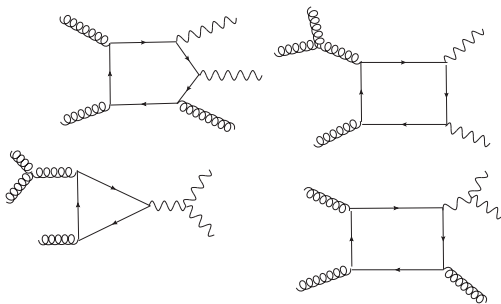
$$p + p \rightarrow V_1 + V_2 + g + X \quad (1)$$

at the LHC, where  $V_1, V_2 \in (\gamma, Z, W)$ .

- The inclusive cross-section is given by,

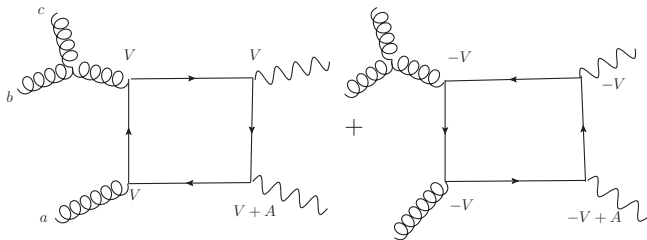
$$\sigma(p + p \rightarrow V_1 + V_2 + g + X) = \int_0^1 dx_1 dx_2 f_g(x_1) * f_g(x_2) * \hat{\sigma}(g + g \rightarrow V_1 + V_2 + g) \quad (2)$$

- At LO, these processes proceed via quark loop diagrams and can be put into following three general classes.

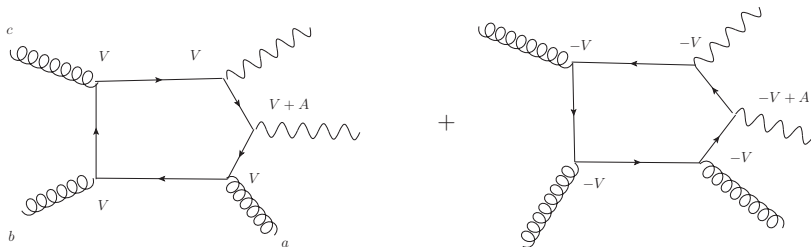


The wavy line represent any one of the  $\gamma/Z/W$ .

- structure of the  $g + g \rightarrow g + Z + \gamma$ , amplitude  
 $3 \times 6$ , Box diagrams and 24, Penta diagrams.



$$\longrightarrow c_V f^{abc} B_V$$



$$\rightarrow (c_V f^{abc} P_V + c_A d^{abc} P_A)$$

- Thus the full amplitude has following general structure.

$$M^{abc}(gg \rightarrow gZ\gamma) = i \frac{f^{abc}}{2} M_V + \frac{d^{abc}}{2} M_A \quad (3)$$

$$M_V = P_V - B_V = \frac{e^2 g_s^3}{\sin\theta_w \cos\theta_w} \left[ \left( \frac{7}{12} - \frac{11}{9} \sin^2\theta_w \right) M_V^{(0)} + \left( \frac{1}{6} - \frac{4}{9} \sin^2\theta_w \right) M_V^{(t)} \right] \quad (4)$$

$$M_A = P_A = (-) \frac{e^2 g_s^3}{\sin\theta_w \cos\theta_w} \left[ \frac{7}{12} M_A^{(0)} + \frac{1}{6} M_A^{(t)} \right] \quad (5)$$

- We will focus on vector contribution only.



- Quark loop **trace** is calculated in FORM. The amplitude at this stage involves tensor integrals, 5-tensor penta-integral being the most complicated one.

$$E_{\mu\nu\rho\sigma\alpha} = \int \frac{d^n l}{(2\pi)^n} \frac{l_\mu l_\nu l_\rho l_\sigma l_\alpha}{D_0 D_1 D_2 D_3 D_4} \quad (6)$$

- Reduction** of tensor integrals into appropriate scalars is done using methods of Oldenborgh and Vermaseren. Using Schouten Identity, we reduce penta-tensor and scalar integrals into lower rank box-tensor and scalar integrals.

$$E_0(0, 1, 2, 3, 4) = \sum_l c_l D_0^{(l)} + O(\epsilon) \quad (7)$$

(Z. Phys. C 46 (1990))

- After all above reductions, the amplitude has following general structure of any one loop amplitude in 4–dimensions.

$$M^{1-loop} = \sum_i d_i D_{0i} + \sum_i c_i C_{0i} + \sum_i b_i B_{0i} + \sum_i a_i A_{0i} + R \quad (8)$$

- Scalar Integrals with massless internal lines are evaluated analytically (up to  $D_0$  with two massive external legs) following 't Hooft and Veltman.  
 (Nucl. Phys. B 153, 365 (1979))
- Scalar integrals with massive internal lines are called from LoopTools.

- We regularize UV sing. in dimensional regularization while IR sing.(for massless quark contribution) are regularized by giving small mass to quarks in the loop.
- for massless quarks in the loop the amplitude looks,

$$M = M_{UV} \frac{1}{\epsilon_{UV}} + M_{IR^2} \log^2(m_q^2) + M_{IR} \log(m_q^2) + M_F \quad (9)$$

- For the massive quark contribution, 2<sup>nd</sup> and 3<sup>rd</sup> terms are absent.

- We have considered real polarization vectors for gauge bosons to compute the amplitude.
- Our processes being LO are expected to be UV as well as IR/ Collinear finite.

## Divergence cancellation:

- UV cancellation ( $M_{UV} = 0$ ):

For both massive and massless quark contributions, we verify that Penta and the three classes of box amplitudes are separately UV finite.

- IR cancellation ( $M_{IR^2} = M_{IR} = 0$ ):

We verify that the amplitude is finite in  $m_q \rightarrow 0$  limit. The IR/ collinear finiteness of the amplitude follows from the fact that each fermion loop diagram is collinear finite by itself.

## Gauge Invariance :

- We have checked the gauge invariance of the vector part of the amplitude with respect to the  $\gamma$ ,  $Z$  and all the three gluons. We do it numerically by replacing their polarizations with their respective 4-momenta.

- As one would expect the penta and three classes of box contributions are separately gauge invariant with respect to the  $\gamma$  and  $Z$ .
- For each gluon one of the three classes of box diagrams is gauge invariant and the cancellation takes place among penta and the other two box contributions.

## Results:

- We present the preliminary result for the five massless quark contribution to the vector part of the amplitude.
- Cuts, Scales and Parameters:

$$P_T(\gamma, Z, g) > 20 \text{ GeV} ; \quad |\eta| < 2.5$$

$$\mu_F = E_T(Z) ;$$

$$\sigma(fb) =$$

$$34.0044, \sqrt{s} = 14 \text{ TeV}$$

$$10.0519, \sqrt{s} = 7 \text{ TeV}$$



- It is important to account for all possible SM backgrounds at colliders.
- Gluon initiated SM processes are important at the LHC.
- Multi-vector bosons produced via gluon fusion can be important background to many new physics models in TeV-region.