

WIMP less Dark Matter

Vikram Rantala

(w J. Feng and Z. Surujon)

arXiv: hep-ph [1108.4689](#), [1111.4479](#)

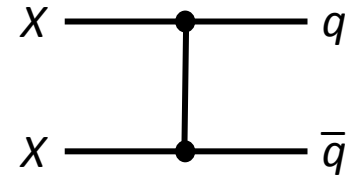
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Outline

- The WIMPless “Miracle” and AMSB
- A simple Abelian model: SQED
- Cosmological and astrophysical implications and observables

The WIMP miracle

$$\Omega_X \propto \frac{1}{\langle \sigma_{\text{an}} v \rangle} \sim \frac{m_X^2}{g_X^4}$$



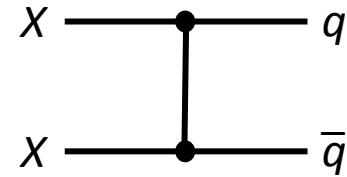
$$g_X \sim g_{\text{weak}} \sim 0.6$$

$$m_X \sim m_{\text{weak}} \sim 100 \text{ GeV} - 1 \text{ TeV}$$

$$\Omega_X \sim 0.1$$

The **WIMPLESS** miracle

$$\Omega_X \propto \frac{1}{\langle \sigma_{\text{ann}} v \rangle} \sim \frac{m_X^2}{g_X^4}$$



$$g_X \sim g_{\text{weak}} \sim 0.6$$

$$m_X \sim m_{\text{weak}} \sim 100 \text{ GeV} - 1 \text{ TeV}$$

$$\Omega_X \sim 0.1$$

WIMPLess Dark Matter with AMSB

- In AMSB, SUSY breaking gaugino masses scale like

$$m_{\chi} \sim \frac{bg^2}{16\pi^2} M_{3/2} \quad \beta(g) = \frac{b}{16\pi^2} g^3$$

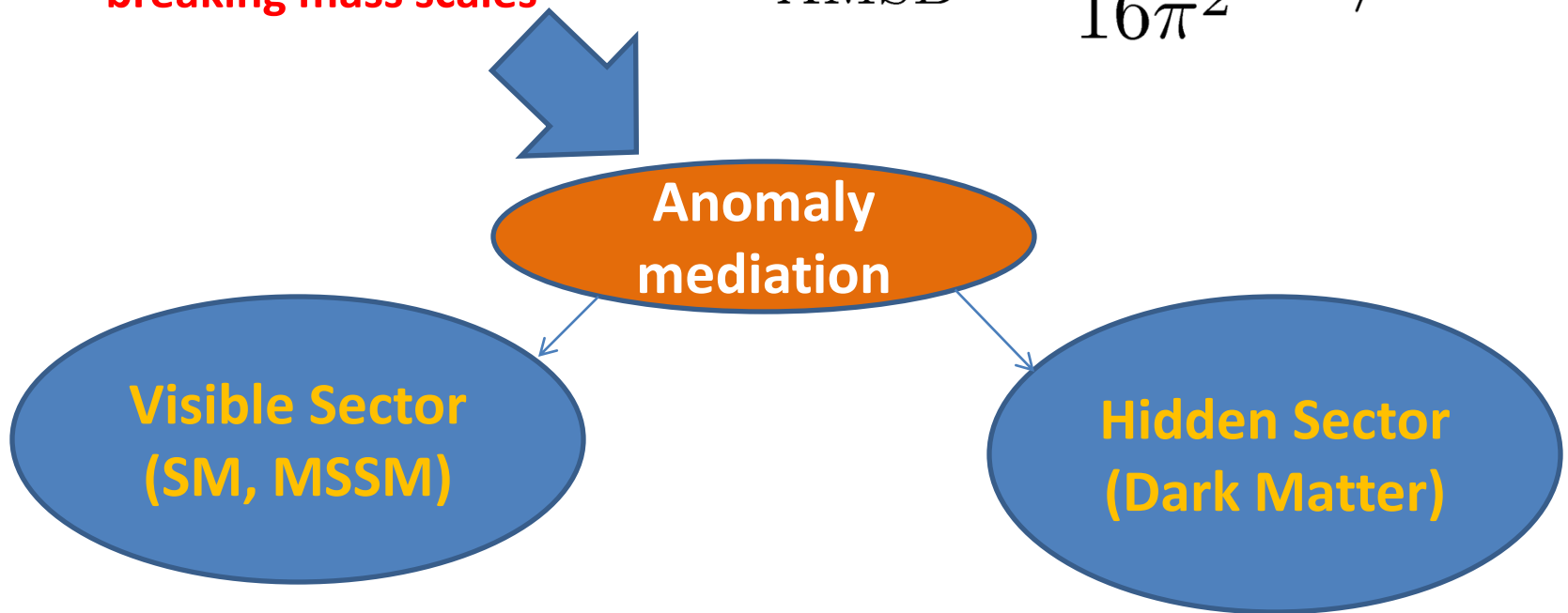
- Scalar masses scale like

$$m_{\phi}^2 = \frac{bg^4}{16\pi^2} M_{3/2}^2$$

Hidden Sectors

Dictates size of SUSY
breaking mass scales

$$m_{\text{AMSB}}^{\text{soft}} \sim \frac{g^2}{16\pi^2} M_{3/2}$$



$$M_{3/2} \sim 100 \text{ TeV}$$

Hidden Sector and Visible Sector AMSB masses

$$m_{\text{AMSB}}^{\text{soft}} \sim \frac{g^2}{16\pi^2} M_{3/2} \quad M_{3/2} \sim 100 \text{ TeV}$$

$$\Omega_X \propto \frac{1}{\langle \sigma_{\text{an}} v \rangle} \sim \frac{m_X^2}{g_X^4}$$

Hidden Sector

Visible Sector

$$\frac{m_X}{g_X^2} \sim \frac{1}{16\pi^2} M_{3/2} \sim \frac{m_{\text{weak}}}{g_{\text{weak}}^2}$$

What if you have a μ -term?

$$\tilde{m}_{\text{phys}} \sim \frac{g^2}{16\pi^2} M_{3/2} + \mu$$

We hope that... $\mu \sim \frac{g^2}{16\pi^2} M_{3/2}$

$$\tilde{m}_{\text{phys}} \sim \frac{g^2}{16\pi^2} f(R) M_{3/2} \quad R \equiv \frac{\mu}{m_{\text{AMSB}}}$$

$$R \sim \mathcal{O}(1)$$

$$\Omega_X \propto \frac{1}{\langle \sigma_{\text{an}} v \rangle} \sim \frac{m_X^2}{g_X^4} \sim f(R) M_{3/2}$$

- In addition to the right ratio of mass to couplings, the Hidden Sector needs a thermal bath to get freeze-out of the thermal relic
- This thermal bath contributes to the light degrees of freedom g^* which is constrained at CMB and BBN
- Usually parameterized as effective number of neutrinos $\Delta N_{\text{eff}} \equiv N_{\text{eff}} - N_{\text{eff}}^{\text{SM}}$

$$\Delta N_{\text{eff}} = 0.19 \pm 1.2 \text{ (95\% CL) BBN}$$

$$\Delta N_{\text{eff}} = 1.51 \pm 0.75 \text{ (68\% CL) CMB (ACT)}$$

$$\Delta N_{\text{eff}} = 0.81 \pm 0.42 \text{ (68\% CL) CMB (SPT)}$$

Model: SQED with N_f flavors

$$N_f, \mu/m_{\tilde{\gamma}}$$

$$b = 2N_f$$

$$M_{3/2}$$

Visible sector

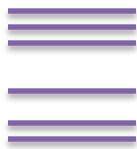
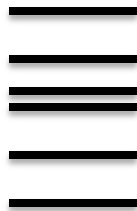
Hidden Sector

MSSM Superpartners



$$m_{\tilde{\gamma}} = b \frac{g^2}{16\pi^2} M_{3/2}$$

SM particles



$$m_{e_i} = |\mu|$$

$$m_{\tilde{e}_{i,1,2}} = \left[|\mu|^2 - 2b \left(\frac{g^2}{16\pi^2} M_{3/2} \right)^2 \right]^{1/2}$$

SM Photon



$$m_{\gamma} = 0$$

Dark Photon

Spectrum

$$N_f, \mu/m_{\tilde{\gamma}}$$

$$b = 2N_f$$

$M_{3/2}$

Visible sector

Hidden Sector

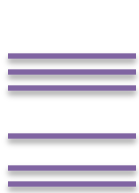
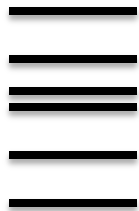
MSSM Superpartners



— $m_{\tilde{\gamma}} = b \frac{g^2}{16\pi^2} M_{3/2}$

Charged!

SM particles



$m_{e_i} = |\mu|$


$m_{\tilde{e}_{i,1,2}} = \left[|\mu|^2 - 2b \left(\frac{g^2}{16\pi^2} M_{3/2} \right)^2 \right]^{1/2}$

SM Photon



$m_{\gamma} = 0$

Dark Photon

- Dark matter is made up of both electrons and selectrons protected by charge, flavor symmetry
 - Halos are like a plasma, so the long range force is actually screened (Debye screening length)
 - Dark force! Constraints from self interacting dark matter
 - Dark photon contributes to the light degrees of freedom (g^* or ΔN_{eff}) at CMB and BBN
 - Relic Density fixed by $M_{3/2}$ at each point in the parameter space
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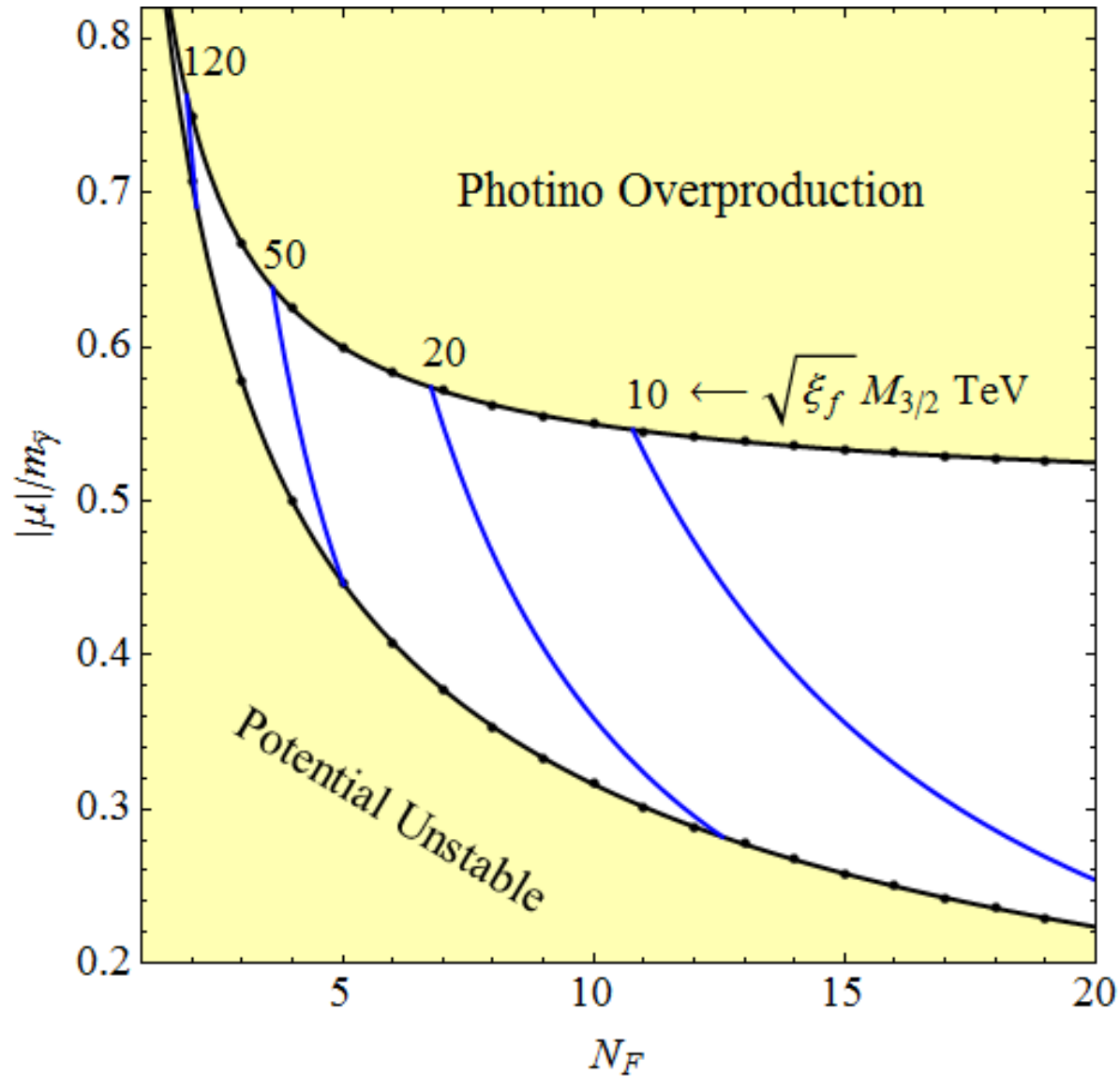
Relic Density and temperature

$$\xi_f = \frac{T_f^h}{T_f^{\text{vis}}}$$

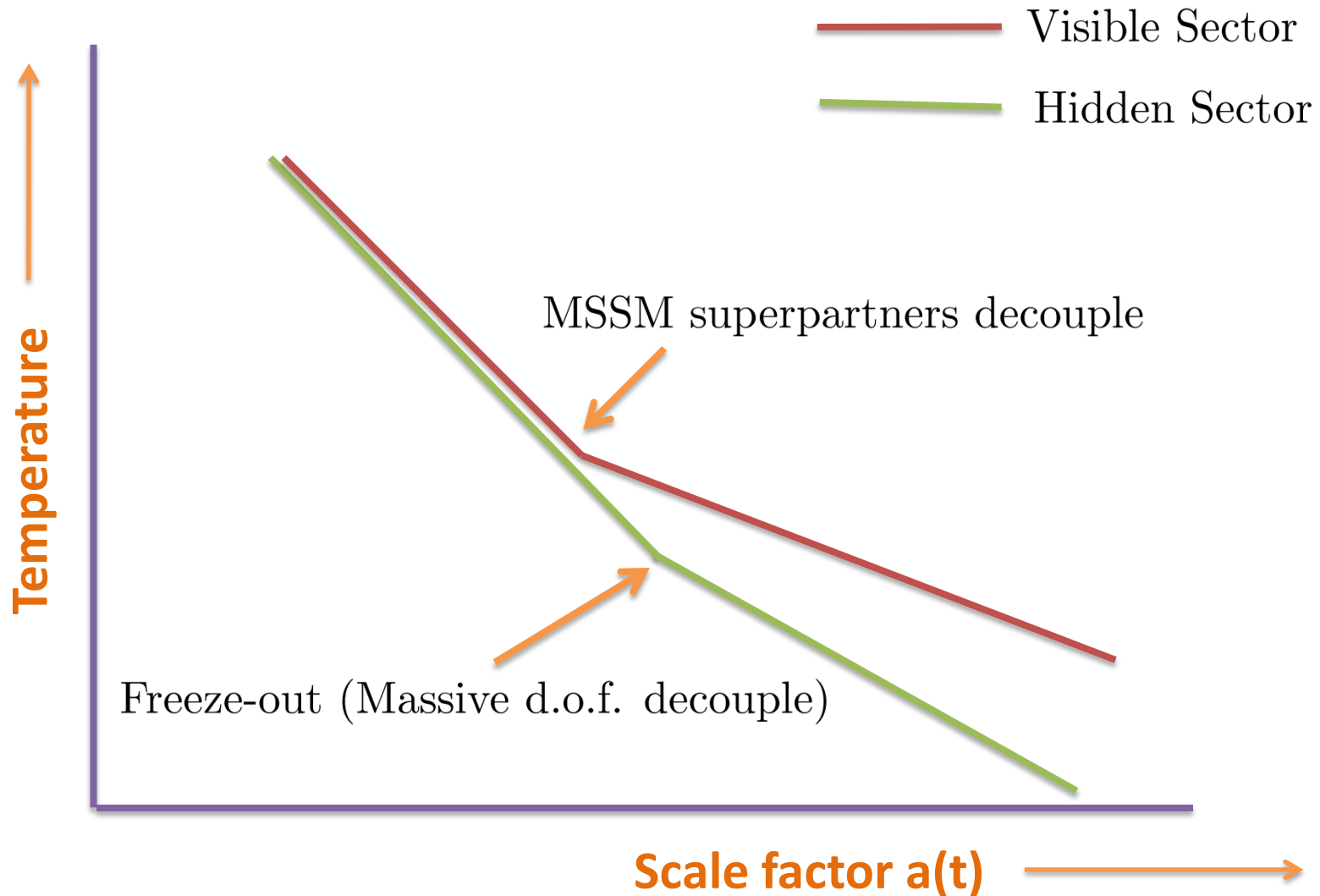
$$\Omega_X = 0.23 \mathcal{F} \left(N_F, \frac{|\mu|^2}{m_{\tilde{\gamma}}^2} \right) \left(\frac{\sqrt{\xi_f} M_{3/2}}{100 \text{ TeV}} \right)^2$$

(Assuming s-wave annihilation)

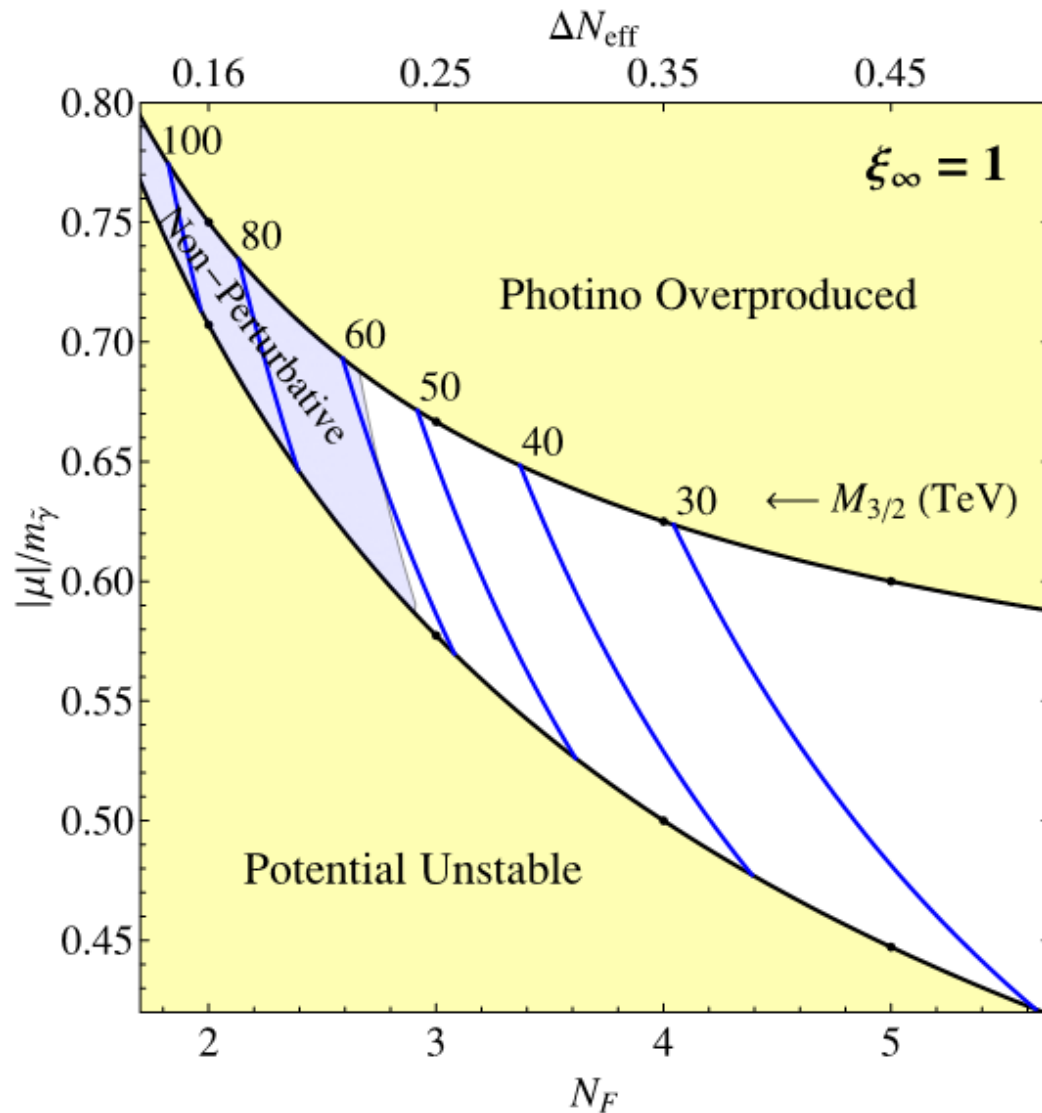
Gravitino Mass... ?



What determines $\xi_f = \frac{T_f^h}{T_f^{\text{vis}}}$?



Gravitino mass*

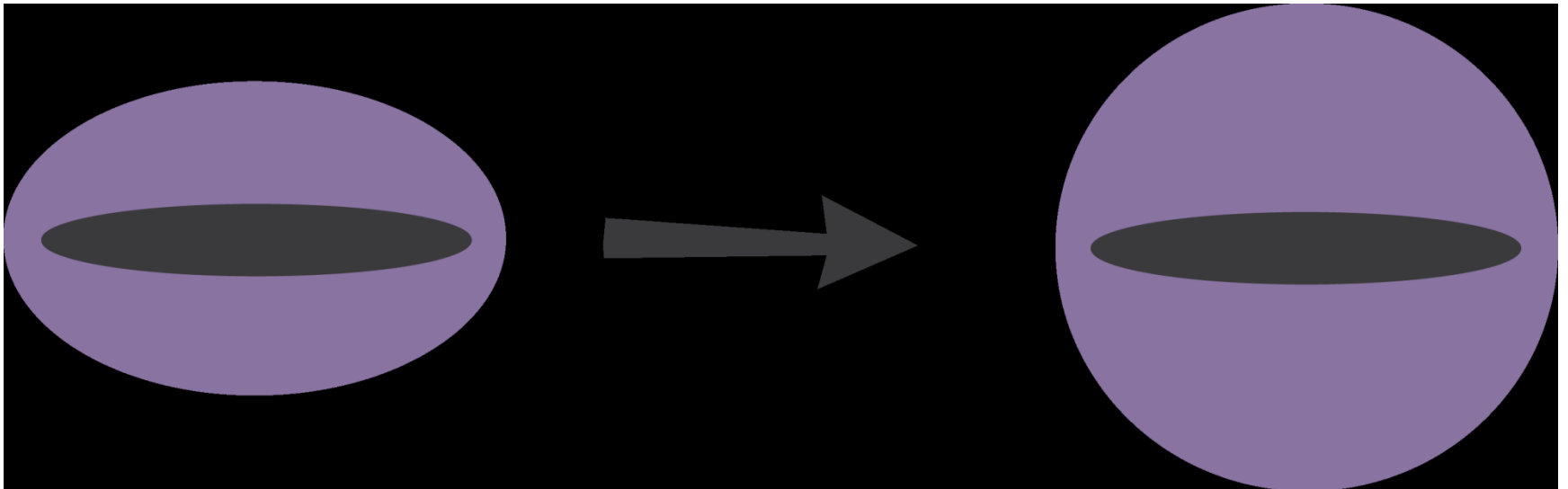


* With assumptions of a desert

Self interaction constraints

- Dark matter in halos is an ionic plasma of positively and negatively charged electrons and selectrons
- Strongest self interaction constraint is from halo shapes of clusters
- Elliptical halos would become spherical if the dark matter particles could exchange an $O(1)$ fraction of their energy

Cluster-Size Elliptical Halo Shapes



$$\tau_r \simeq \frac{1}{n_X \sigma v} = \frac{m_X^3 v_0^3}{4\sqrt{\pi} \alpha_X^2 \rho_X C}$$

(Miralda-Escude 2002, Feng+Kaplinghat+Tu+Yu 2009)

Self interaction constraint

$$\tau_r \simeq \frac{m_X^3 v_0^3}{4\sqrt{\pi}\alpha_X^2 \rho_X C}$$

Relaxation time for halos

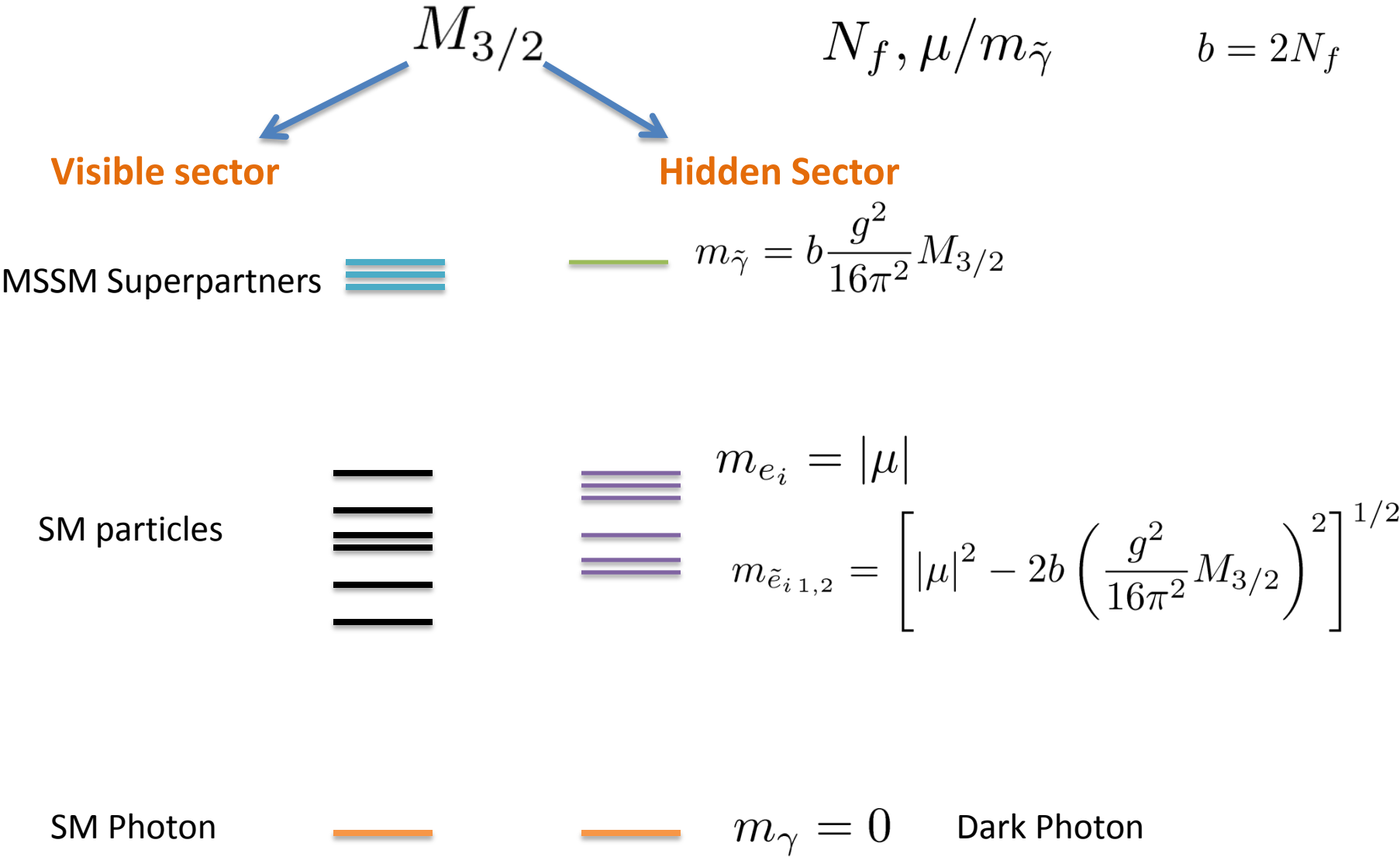
Requiring this to be greater than the age of the cluster:

$$\tau_r > 10^{10} \text{ yr}$$

$$m_{e_i} > m_{\text{DM}}^{\text{min}} \equiv \left[\frac{6.6}{N_F (\mu/m_{\tilde{\gamma}})} \right]^2 \left(\frac{100 \text{ TeV}}{M_{3/2}} \right)^2 \text{ TeV}$$

$$\text{or } \alpha > \alpha_{\text{min}}$$

Mass scale of dark matter from astrophysics!



Summary and conclusions

- WIMPless dark matter suggests a **larger range of dark matter masses and couplings** than previously assumed
- Simple model shows that AMSB WIMPless models are predictive and have a rich phenomenology
- Requirement of freeze out in the Hidden sector requires a long range force
- Leads to a **prediction for the effective number of neutrinos**
- Constrained/could be **further tested by signals of dark matter self interactions**
- **R-parity violation** allowed in the visible sector: LHC signatures
- Open to a plethora of model building possibilities

QUESTIONS, COMMENTS, SUGGESTIONS?



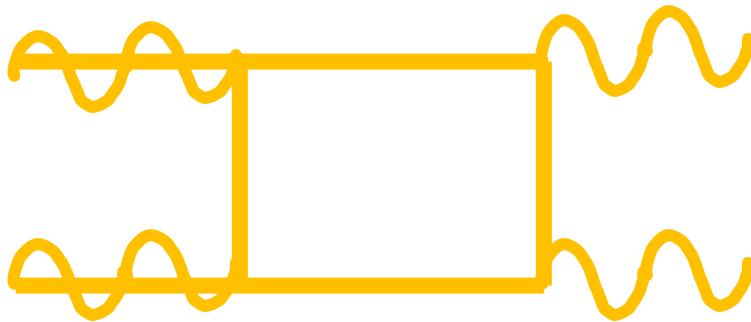
Backup Slides

Annihilation Processes

$$e_i^+ e_i^- \rightarrow \gamma\gamma$$

$$\tilde{e}_i^+ \tilde{e}_i^- \rightarrow \gamma\gamma$$

$$\tilde{\gamma}\tilde{\gamma} \rightarrow \gamma\gamma$$



Photino not a good WIMPless relic!

$$\sigma v \propto \left(\frac{g^2}{16\pi^2} \right)^2 \frac{g^4}{m^2}$$

Effective number of neutrinos

$$\Omega_{\text{DR}} \propto g_*^{\text{h}}(T_{\text{CMB}}^{\text{h}}) T_{\text{CMB}}^{\text{h}^4}$$

$$\begin{aligned} \Omega_{\nu} &\propto g_*^{\nu}(T_{\text{CMB}}^{\nu}) T_{\text{CMB}}^{\nu^4} \\ &\propto \frac{7}{8} 2N_{\nu} \left[\frac{4}{11} \right]^{\frac{4}{3}} T_{\text{CMB}}^{\text{vis}^4} \end{aligned}$$

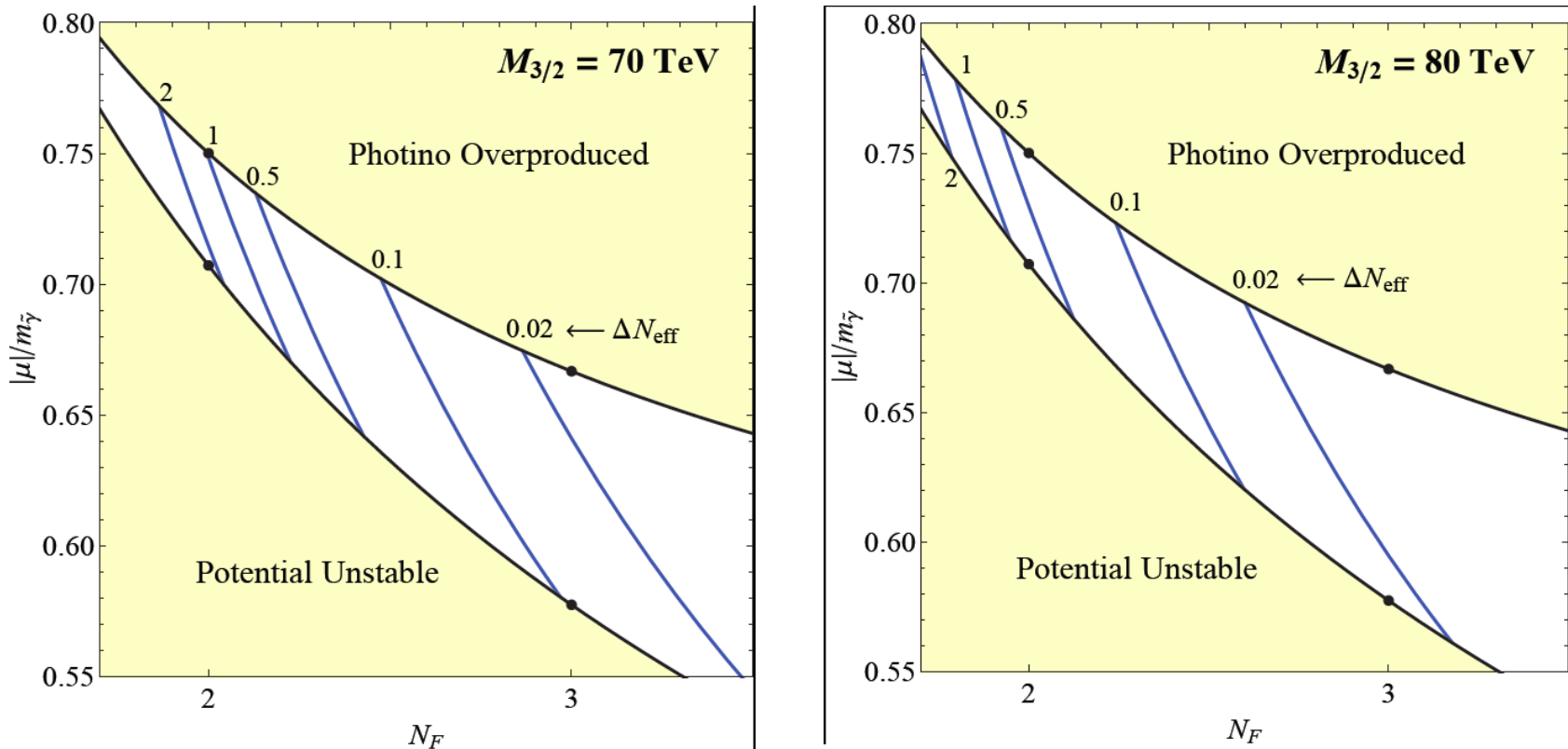
$$\Delta N_{\text{eff}} = \left(\frac{\xi_f}{2.60} \right)^4 \left[\frac{106.75}{g_*^{\nu}(T_f^{\nu})} \right]^{\frac{4}{3}}$$

(in general)

$$\Delta N_{\text{eff}} = 0.20 \left(\frac{N_F + \frac{8}{15}}{3} \right)^{\frac{4}{3}} \xi_{\infty}^4$$

(With the assumption of a desert)

Effective number of neutrinos



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