# *tt* Forward-Backward Asymmetry from Loops of New Strongly Coupled Quarks

#### Thomas McElmurry University of Rochester

H. Davoudiasl, T.M., A. Soni, Phys. Rev. D 85, 054001

Pheno 2012 Pittsburgh, Pennsylvania 2012-05-07

◆□▶ ◆□▶ ▲□▶ ▲□▶ □ のQ@

# The top quark

Vital statistics

- Spin:  $\frac{1}{2}$
- Charge:  $+\frac{2}{3}$
- Color: 3
- Height: Zero?
- Weight: 173 GeV  $\times g/c^2$
- Age: 17 years
- Sign: Pisces
- Birthplace: Batavia, Illinois
- Current residence: Geneva, Switzerland

◆□▶ ◆□▶ ▲□▶ ▲□▶ □ のQ@

# The top quark

- The top quark is very heavy:  $m_t \approx m_W$  (W = tungsten)
- and very close to the EW scale:  $m_t \approx v/\sqrt{2}$ .
- This suggests that it may play a role in EWSB or be sensitive to other new physics.
- One observable where new physics may show up is the forward-backward asymmetry.

(ロ) (同) (三) (三) (三) (○) (○)

#### Forward-backward asymmetry

- Define  $t\bar{t}$  events as F or B based on sign( $\Delta y_{t\bar{t}}$ )
- Asymmetry:

$$A \equiv \frac{N_F - N_B}{N_F + N_B}$$

- To see this asymmetry, one needs an asymmetric initial state:
  - A = 0 at a *pp* collider due to the symmetric initial state
  - Gluon and sea quark PDFs are the same in p as in  $\overline{p}$ .
  - Nonzero A requires  $p\bar{p}$  collider and  $u\bar{u}$  or  $d\bar{d}$  initial state.
- The Tevatron is a good place to measure the asymmetry.
  - It's a pp̄ collider.
    - Protons go clockwise ( $\equiv$  Forward).
    - Antiprotons go anticlockwise ( $\equiv$  Backward).
  - $t\bar{t}$  production is dominated by the  $q\bar{q}$  channel.

# Forward-backward asymmetry

in the Standard Model

At leading order, A = 0 due to charge-conjugation symmetry



 NLO QCD permits an asymmetry due to interference between C-even and C-odd amplitudes



- Virtual corrections give a positive asymmetry
- Real radiation gives a negative asymmetry
- Net result:  $A^{t\bar{t}} \approx 0.05$ . [Kühn, Rodrigo 1998]

# **Tevatron results**

[CDF, Phys. Rev. D83, 112003 (2011); D0, arXiv:1107.4995]

- CDF measures  $A^{t\bar{t}} = 0.158 \pm 0.075 \dots$ 
  - ... with a significant dependence on invariant mass:
    - $A^{t\bar{t}}(M_{t\bar{t}} < 450 \,\text{GeV}) = -0.116 \pm 0.153$
    - $A^{t\bar{t}}(M_{t\bar{t}} > 450 \,\text{GeV}) = 0.475 \pm 0.114$
  - ... and on separation in rapidity:
    - $A_{t\bar{t}}^{t\bar{t}}(|\Delta y_{t\bar{t}}| < 1) = 0.026 \pm 0.118$
    - $A^{t\bar{t}}(|\Delta y_{t\bar{t}}| > 1) = 0.611 \pm 0.256$
- D0 measures  $A^{t\bar{t}} = 0.196 \pm 0.065 \dots$ 
  - ... without the strong enhancement at high  $M_{t\bar{t}}$  and  $|\Delta y_{t\bar{t}}|$ .

### The theorists' challenge

How to explain the observed asymmetry without changing the total cross section or violating other constraints?

- Various ideas:
  - ▶ New particle in the *s*-channel (axigluons, ...)



New particle in the *t*-channel (flavor-changing Z', new scalar, ...)



◆□▶ ◆□▶ ▲□▶ ▲□▶ ■ ののの

## The theorists' challenge

How to explain the observed asymmetry without changing the total cross section or violating other constraints?

 A new idea: A<sup>tt</sup> arises from loop diagrams involving new heavy quarks.



◆□▶ ◆□▶ ▲□▶ ▲□▶ ■ ののの

- ▶ New heavy quarks Q = t', b', with mass  $m_Q$
- New scalars  $\phi^0, \phi^{\pm}$ , with mass  $m_{\phi}$
- Flavor-changing interactions  $\phi \bar{q} Q$

$$\mathcal{L} \supset \lambda_{ut'} \phi^0 \bar{u}t' + \lambda_{ub'} \phi^+ \bar{u}b' + \lambda_{dt'} \phi^- \bar{d}t' + \lambda_{db'} \phi^0 \bar{d}b' + \lambda_{tt'} \phi^0 \bar{t}t' + \lambda_{tb'} \phi^+ \bar{t}b' + \text{H.C.}$$

▲□▶ ▲□▶ ▲ 三▶ ▲ 三▶ - 三 - のへぐ

- New heavy quarks Q = t', b', with mass  $m_Q$
- New scalars  $\phi^0, \phi^{\pm}$ , with mass  $m_{\phi}$
- Flavor-changing interactions  $\phi \bar{q} Q$

$$\mathcal{L} \supset \lambda_{\boldsymbol{u}}(\phi^{0}\bar{\boldsymbol{u}}t' + \phi^{+}\bar{\boldsymbol{u}}b') + \lambda_{\boldsymbol{d}}(\phi^{-}\bar{\boldsymbol{d}}t' + \phi^{0}\bar{\boldsymbol{d}}b') + \lambda_{\boldsymbol{t}}(\phi^{0}\bar{\boldsymbol{t}}t' + \phi^{+}\bar{\boldsymbol{t}}b') + \text{H.C.}$$

▲□▶ ▲□▶ ▲ 三▶ ▲ 三▶ - 三 - のへぐ

- New heavy quarks Q = t', b', with mass  $m_Q$
- New scalars  $\phi^0, \phi^{\pm}$ , with mass  $m_{\phi}$
- Flavor-changing interactions  $\phi \bar{q} Q$

$$\mathcal{L} \supset \lambda_{\boldsymbol{u}}(\phi^{0}\bar{\boldsymbol{u}}t' + \phi^{+}\bar{\boldsymbol{u}}b') + \lambda_{\boldsymbol{d}}(\phi^{-}\bar{\boldsymbol{d}}t' + \phi^{0}\bar{\boldsymbol{d}}b') + \lambda_{\boldsymbol{t}}(\phi^{0}\bar{\boldsymbol{t}}t' + \phi^{+}\bar{\boldsymbol{t}}b') + \text{H.C.}$$

- Some of these Yukawa couplings will need to be strong.
- Couplings to s, c, b may exist, but are constrained by flavor physics.

- New heavy quarks Q = t', b', with mass  $m_Q$
- New scalars  $\phi^0, \phi^{\pm}$ , with mass  $m_{\phi}$
- Flavor-changing interactions  $\phi \bar{q} Q$

$$\mathcal{L} \supset \lambda_{\boldsymbol{u}}(\phi^{0}\bar{\boldsymbol{u}}t' + \phi^{+}\bar{\boldsymbol{u}}b') + \lambda_{\boldsymbol{d}}(\phi^{-}\bar{\boldsymbol{d}}t' + \phi^{0}\bar{\boldsymbol{d}}b') + \lambda_{\boldsymbol{t}}(\phi^{0}\bar{\boldsymbol{t}}t' + \phi^{+}\bar{\boldsymbol{t}}b') + \text{H.C.}$$

< □ > < 同 > < 三 > < 三 > < 三 > < ○ < ○ </p>

- ▶ φ may have additional couplings to light quarks, leptons, ...
- Decay modes of  $\phi$  are unspecified.

#### New diagrams Boxes!

The new particles appear in box diagrams ...



 ... which interfere with the LO diagram and produce an asymmetry.

▲□▶ ▲□▶ ▲□▶ ▲□▶ = 三 のへで

# New diagrams

Triangles!

We also have vertex corrections ....



... analogous to SM diagrams involving the Higgs.

[Stange, Willenbrock 1993]

These diagrams do not produce an asymmetry, but do affect the total cross section.

・ コット (雪) ( 小田) ( コット 日)

# Asymmetry

Can we match the CDF measurement?

For what values of  $\lambda_u$ ,  $\lambda_d$ ,  $\lambda_t$ ,  $m_Q$ ,  $m_\phi$  do we pass these six criteria?

Agreement with the CDF inclusive asymmetry, within 2σ

< □ > < 同 > < 三 > < 三 > < 三 > < ○ < ○ </p>

- and with the high and low  $M_{t\bar{t}}$  bins, within  $2\sigma$
- and with the high and low  $|\Delta y_{t\bar{t}}|$  bins, within  $2\sigma$
- and with the SM total cross section, within 30%

#### Asymmetry

Can we match the CDF measurement?



Cross-hatched:  $\lambda_u = 1, \lambda_d = 3, \lambda_t = 5$ Shaded:  $\lambda_u = 1, \lambda_d = 2, \lambda_t = 6$ Hatched:  $\lambda_u = 0, \lambda_d = 3.5, \lambda_t = 4.5$ 

# Asymmetry

Can we match the D0 measurement?

- Generally, the limiting factor in matching the CDF measurement is the high-M<sub>tī</sub> bin.
- Therefore the D0 measurement can be accommodated more easily (i.e. with smaller couplings).

(ロ) (同) (三) (三) (三) (三) (○) (○)

#### What about constraints?

- This model predicts new heavy quarks t' and b'.
  - Experimental searches rule out e.g.  $t' \rightarrow Wb$  with  $m'_t \lesssim 550 \,\text{GeV}$ . [CMS, arXiv:1203.5410]
  - These bounds do not directly apply to our model, since we expect decay modes such as ...

• 
$$t' \to t\phi^0 \to Wbjj$$

• 
$$t' \rightarrow d\phi^+ \rightarrow 3j$$

- etc.
- We can produce same-sign top pairs:  $uu \rightarrow tt$ .
  - ► We estimate that recent constraints [CMS, JHEP 1108, 005 (2011)] do not exclude our model if  $m_Q \gtrsim 400$  GeV.
  - We can also evade the bounds by choosing couplings such that λ<sub>u</sub> is small.

(日) (日) (日) (日) (日) (日) (日)

#### What about constraints?

Our model allows exotic single-top production.



▲□▶ ▲□▶ ▲□▶ ▲□▶ □ のQで

- This can be suppressed by  $\lambda_{u} \rightarrow 0$ .
- Or, if  $\lambda_{ut'} \approx -\lambda_{ub'}$ , the two diagrams cancel.

#### Conclusions

- Strongly coupled heavy quarks appearing in loops can produce an A<sup>tt̄</sup> consistent with the CDF results.
- The D0 results can also be accommodated with smaller couplings.
- Various predictions and possibilities for the LHC:
  - Discovery of t', b'
  - Decay modes such as  $t' 
    ightarrow t \phi$
  - Single t', b' production
  - Forward-backward asymmetry in  $t'\bar{t}'$  and  $b'\bar{b}'$  production

(ロ) (同) (三) (三) (三) (三) (○) (○)