

Constraining Split Universal Extra Dimensions

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arXiv:1204.0522, with K.C. Kong and S.C. Park

MUED: only parameter R

- Features: all fields in bulk, flat, KK parity
- Constrained: $700 \text{ GeV} \lesssim 1/R \lesssim 850 \text{ GeV}$

SUED: Fermion Bulk Mass term

$$S_{\text{SUED}} \ni - \int d^4x \int_{-L}^L dy M_{\Psi}(y) \bar{\Psi} \Psi ,$$

Preserves 5D Lorentz symmetry, gauge symmetry.

$M_{\Psi}(y)$ Odd to preserve KK parity. Simplest: $M(y) \sim \theta(y)$

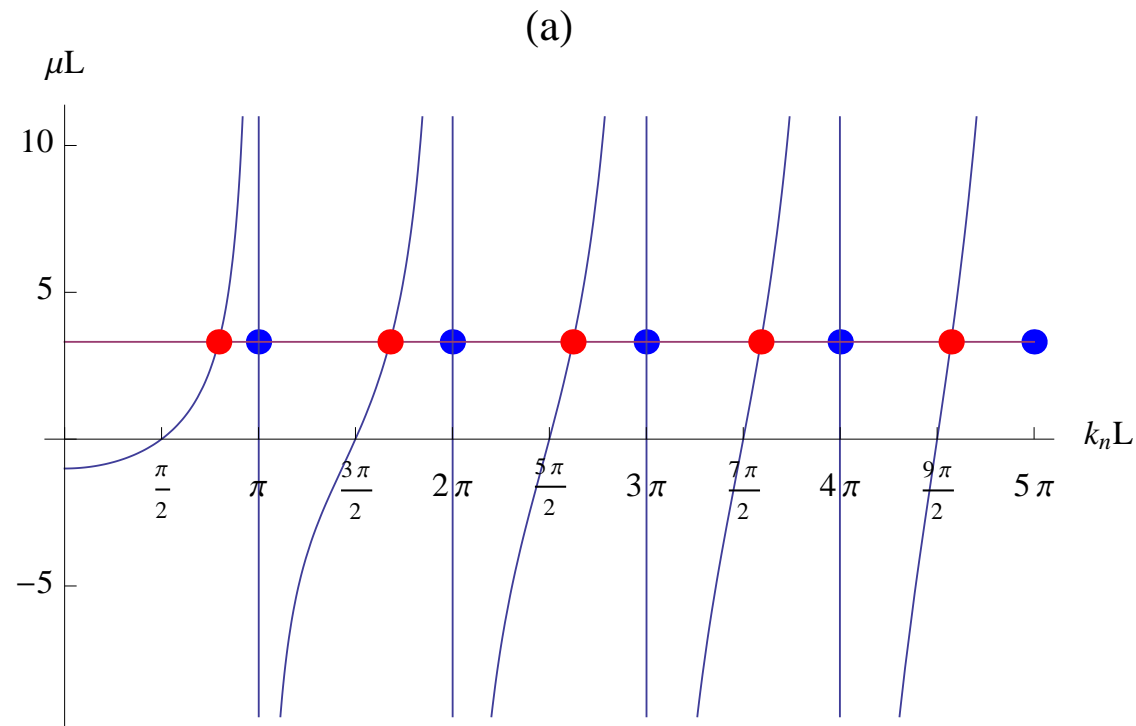
$$\text{Flavor Bound: } \rightarrow \begin{cases} -M_Q = M_U = M_D = \mu_Q \theta(y) \\ -M_L = M_E = \mu_L \theta(y) \end{cases}$$

Universal bulk mass limit: $\mu_L = \mu_Q \equiv \mu$

Split Spectrum: KK Fermion Masses

$$m_{\Psi^{(n)}}^2 = \begin{cases} \lambda_{\Psi}^2 v^2 & \text{if } n = 0 \\ \mu^2 + k_n^2 + \lambda_{\Psi}^2 v^2 & \text{if } n \geq 1 \end{cases}$$

$\mu = -k_n \cot(k_n L)$ for odd n , possible imaginary k_1



Recovering the MUED limit at $\mu = 0$

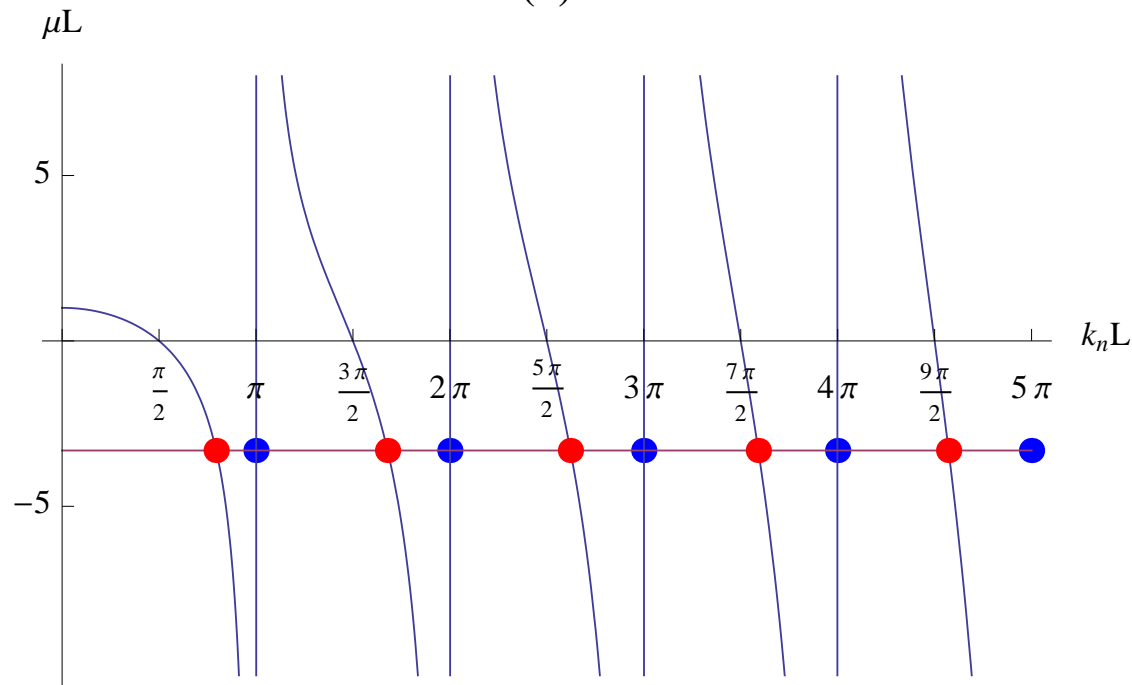
LKP = γ_1 for $\mu > 0$

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$\mu = +k_n \cot(k_n L)$ for odd n , possible imaginary k_1

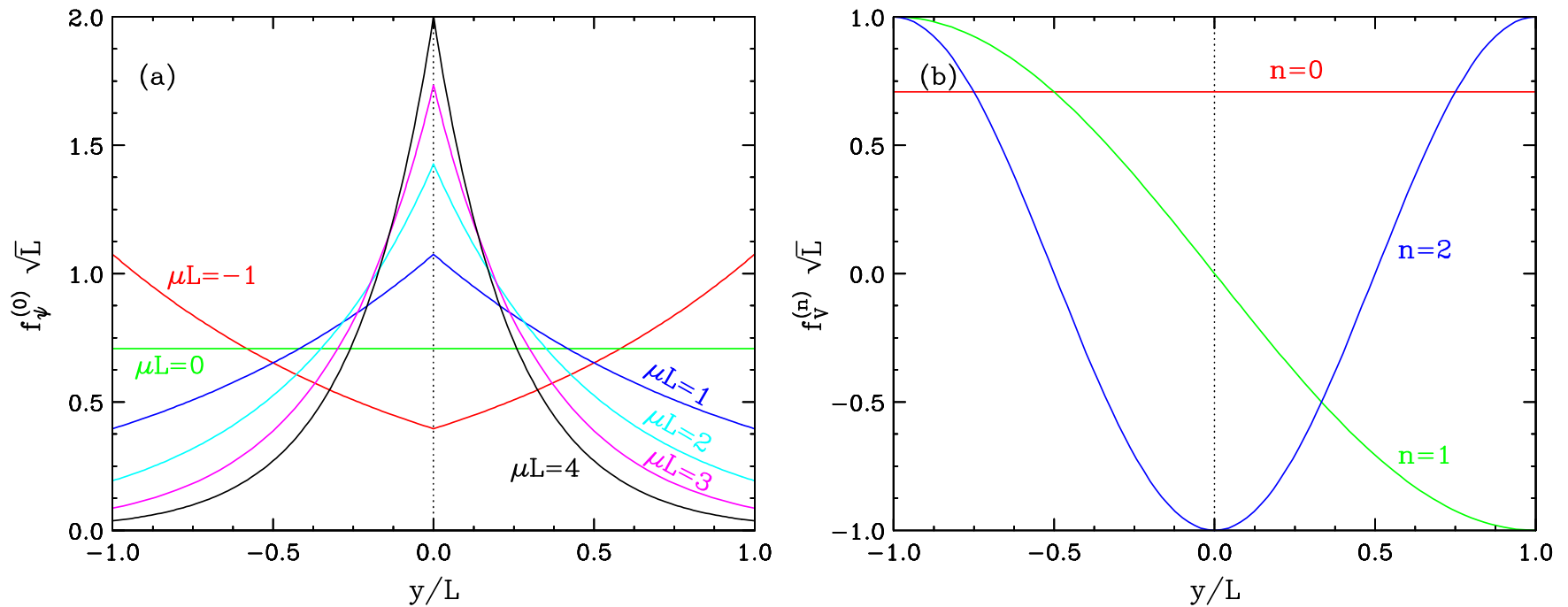
(b)



Recovering the MUED limit at $\mu = 0$

LKP = γ_1 for $\mu > 0$

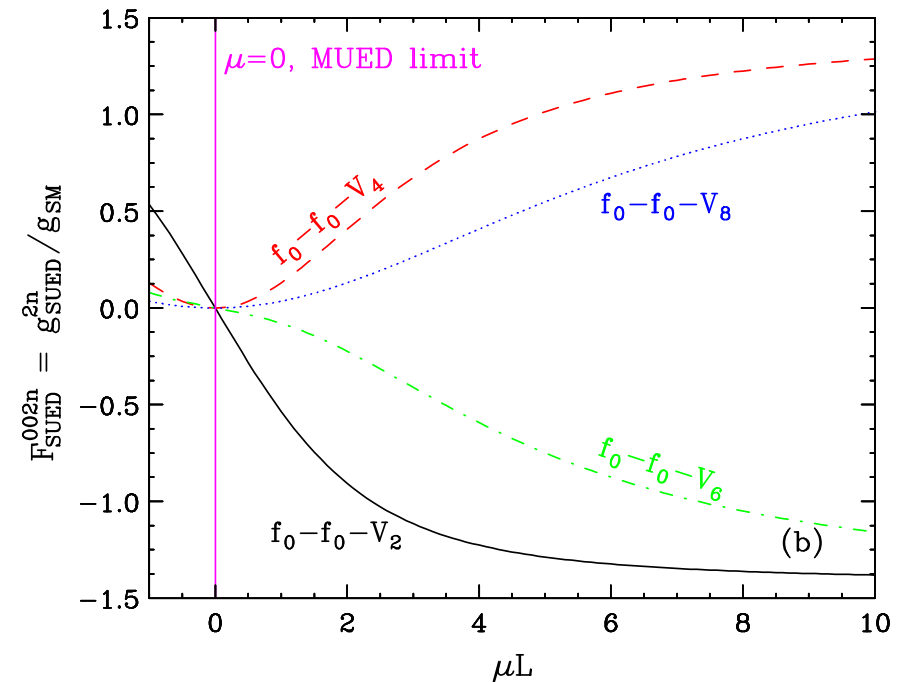
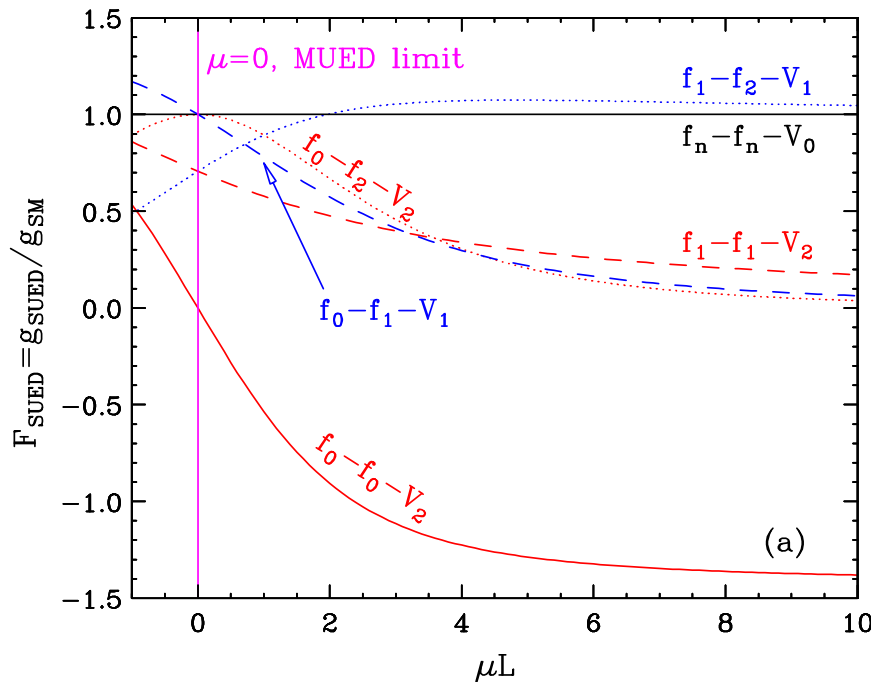
Fermion and Gauge Boson Bulk Profiles



KK couplings, from overlap integrals

$$g_{mln} = \frac{g_5}{\sqrt{2L}} \int_{-L}^L dy \psi_m(y) \psi_l^*(y) f_V^n(y) \equiv g_{\text{SM}} \mathcal{F}_{ml}^n(\mu_\Psi L)$$

$$\mathcal{F}_{00}^{2n} = \frac{x_\Psi^2 [1 - (-1)^n e^{2x_\Psi}] [1 - \coth(x_\Psi)]}{\sqrt{2(1 + \delta_{0n})} [x_\Psi^2 + n^2 \pi^2 / 4]}, \quad x_\Psi = \mu_\Psi L$$



Large μ , asymptotic behavior

Constraining the parameter space

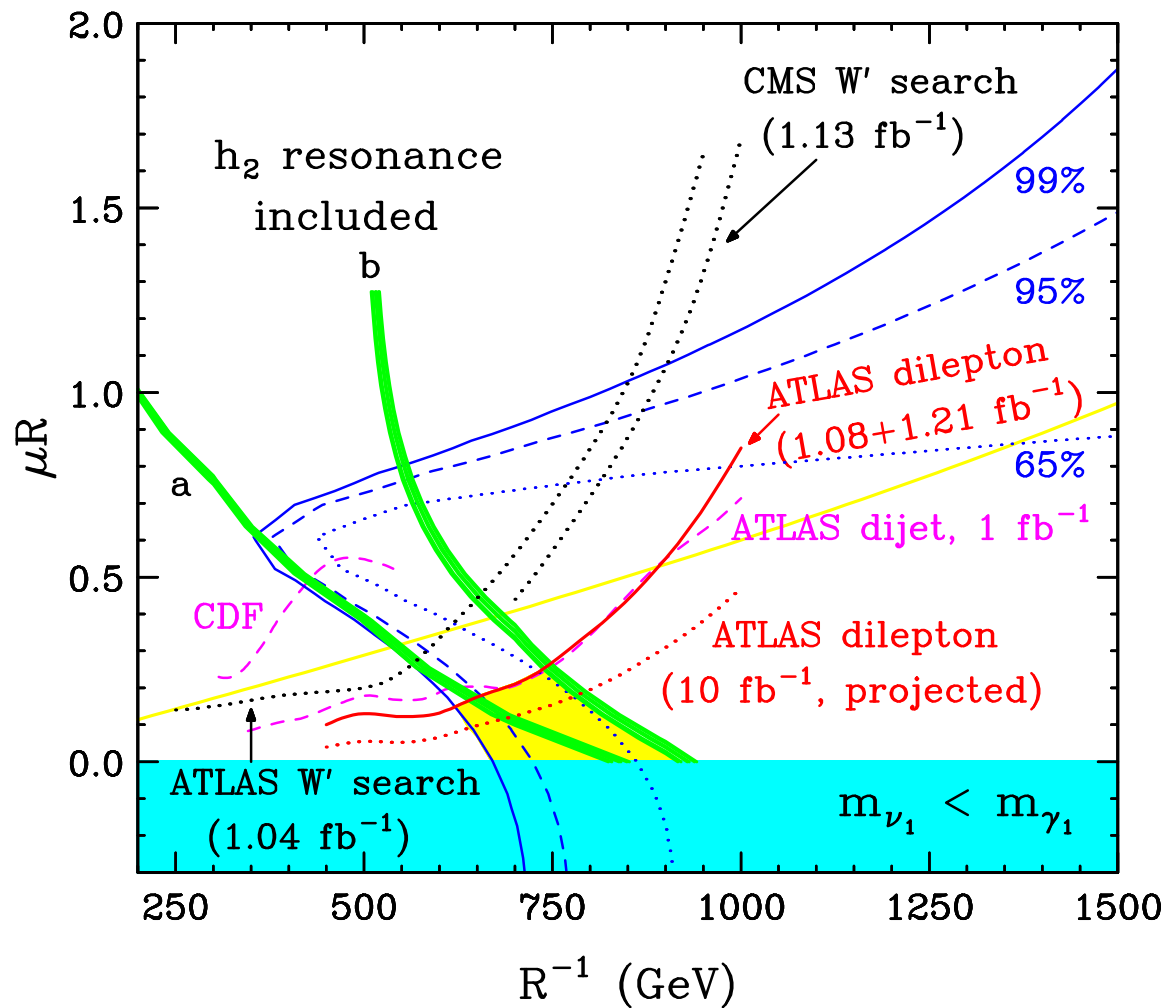
- Relic Density
- Collider
- Electroweak (Oblique, 4-Fermi, $g - 2\dots$)

Special Cases:

- Universal Bulk Mass (μ, R^{-1})
- Non-Universal Bulk Mass (μ_Q, μ_L, R^{-1})

The Universal Case: $\mu_L = \mu_Q$

Summary Plot of Constraints



Relic Density: $\Omega h^2 \propto n \cdot m_{DM} \sim 0.112$, $n \propto \frac{1}{\langle \sigma v \rangle}$

Annihilation

via KK-1 fermions (t-, u-) and KK-2 bosons (s-channel)

$\rightarrow ff, WW, hh$ final states

$m_{f1} = \sqrt{m_{f0}^2 + k_1^2 + \mu^2} > 1/R$ for $\mu > 0$, raised from MUED

Smaller cross-section, higher relic density, lower DM mass

Resonance through h_2

Effect evident when mass difference within a few percent

Detailed one-loop calculation needed

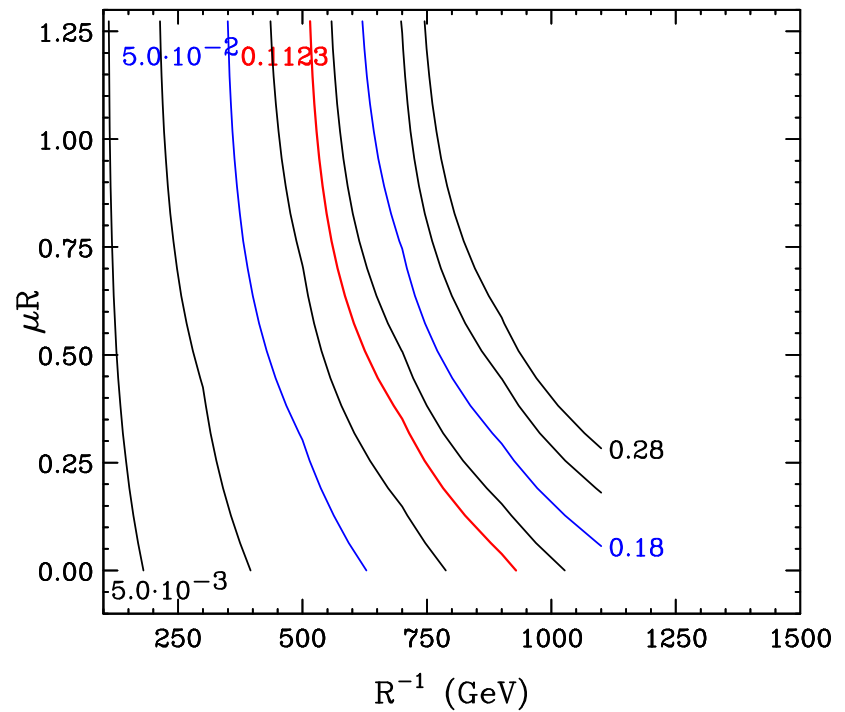
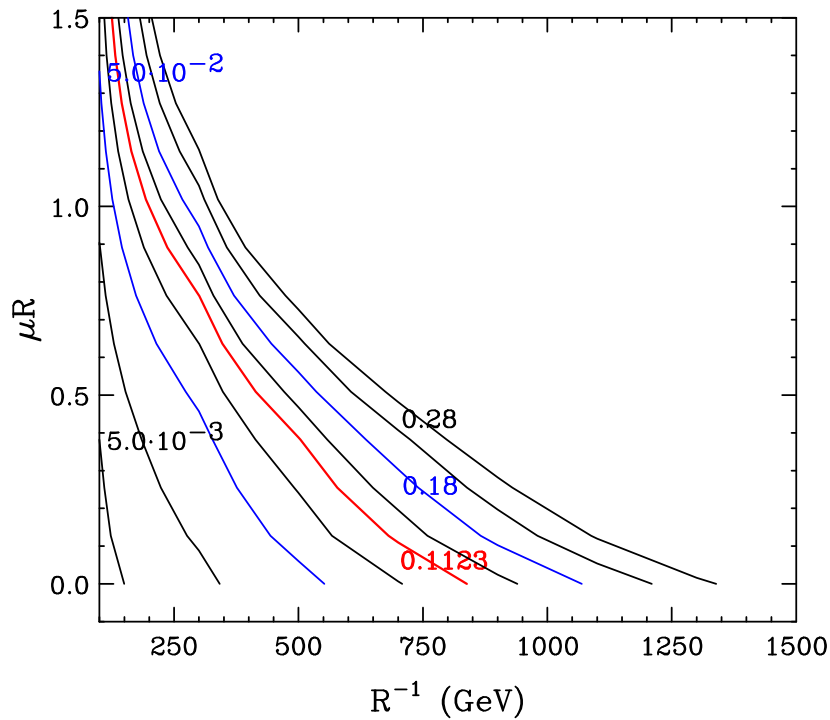
Coannihilation with KK-1 particles

Need near degeneracy \sim few %

Broken by μ for $\mu R > 0.01$ and radiative correction

Relic density contours

Without and with h_2 coannihilation



Collider Searches of KK excitations

MUED

KK-1 particles can be pair-produced

(KK-parity, \cancel{E} signature)

$1/R \gtrsim 700$ GeV from first year LHC data

KK-2 production loop suppressed.

SUED

\mathcal{F}_{00}^2 at tree-level

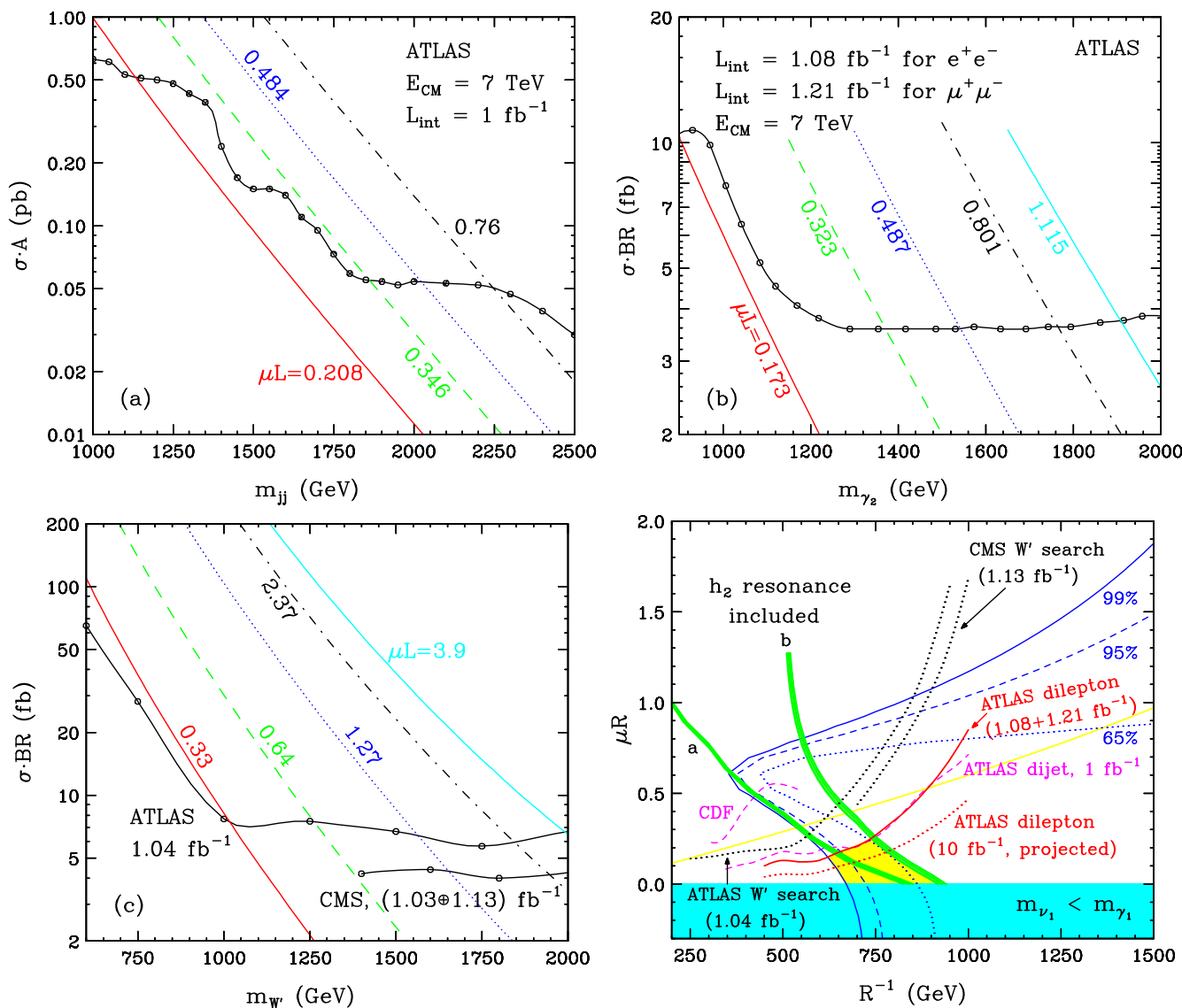
Resonance/Single production of KK-2 bosons (LHC 2011)

$jj, \ell\ell, \ell\nu$ channels

Bounds on masses of KK resonances

jj channel: $G_2 (\gamma_2, Z_2, W_2)$, $\delta m \sim 0.3m_{\gamma_2}$

ll channel: $\gamma_2 + Z_2$, $\delta m \sim 0.07m_{\gamma_2}$



S, T, U in UED

$$\begin{aligned}
S &= \frac{4s_W^2}{\alpha} \left[\frac{3g^2}{4(4\pi)^2} \left(\frac{2}{9} \sum_n \frac{m_t^2}{(n/R)^2} \right) + \frac{g^2}{4(4\pi)^2} \left(\frac{1}{6} \frac{m_h^2}{1/R} \right) \zeta(2) \right] \\
T &= \frac{1}{\alpha} \left[\frac{3g^2}{2(4\pi)^2} \frac{m_t^2}{m_W^2} \left(\frac{2}{3} \sum_n \frac{m_t^2}{(n/R)^2} \right) + \frac{g^2 s_W^2}{(4\pi)^2 c_W^2} \left(-\frac{5}{12} \frac{m_h^2}{1/R} \right) \zeta(2) \right] \\
U &= -\frac{4s_W^2}{\alpha} \left[\frac{g^2 s_W^2}{(4\pi)^2} \frac{m_W^2}{(1/R)^2} \left(\frac{1}{6} \zeta(2) - \frac{1}{15} \frac{m_h^2}{(1/R)^2} \zeta(4) \right) \right]
\end{aligned}$$

Riemann zeta functions: $\zeta(m) = \sum_n \frac{1}{n^m}$

Contributions from top, gauge and Higgs loops

 S, T, U in SUED

- Corrections from Fermion Mass
- Corrections from KK W contributions to G_F

S, T, U in SUED

- Corrections from Fermion Mass

$$\sum_n \frac{m_t^2}{(n/R)^2} \rightarrow \sum_n \frac{m_t^2}{m_t^2 + \mu^2 + (n/R)^2}$$

- Corrections from KK W contributions to G_F

$$S_{SUED} = S_{UED}$$

$$T_{SUED} = T_{UED} - \frac{1}{\alpha} \frac{\delta G_F}{G_F}$$

$$U_{SUED} = U_{UED} + \frac{4s_W^2}{\alpha} \frac{\delta G_F}{G_F}$$

$$G_F = G_F^0 + \delta G_F = \frac{g^2}{\sqrt{32}m_W^2} + \frac{1}{\sqrt{32}} \sum_n \frac{g_{002n}^2}{m_W^2 + \left(\frac{2n}{R}\right)^2}$$

S, T, U fitting contours

from Gfitter

- fitted values

$$S_{NP} = 0.04 \pm 0.10,$$

$$T_{NP} = 0.05 \pm 0.11,$$

$$U_{NP} = 0.08 \pm 0.11,$$

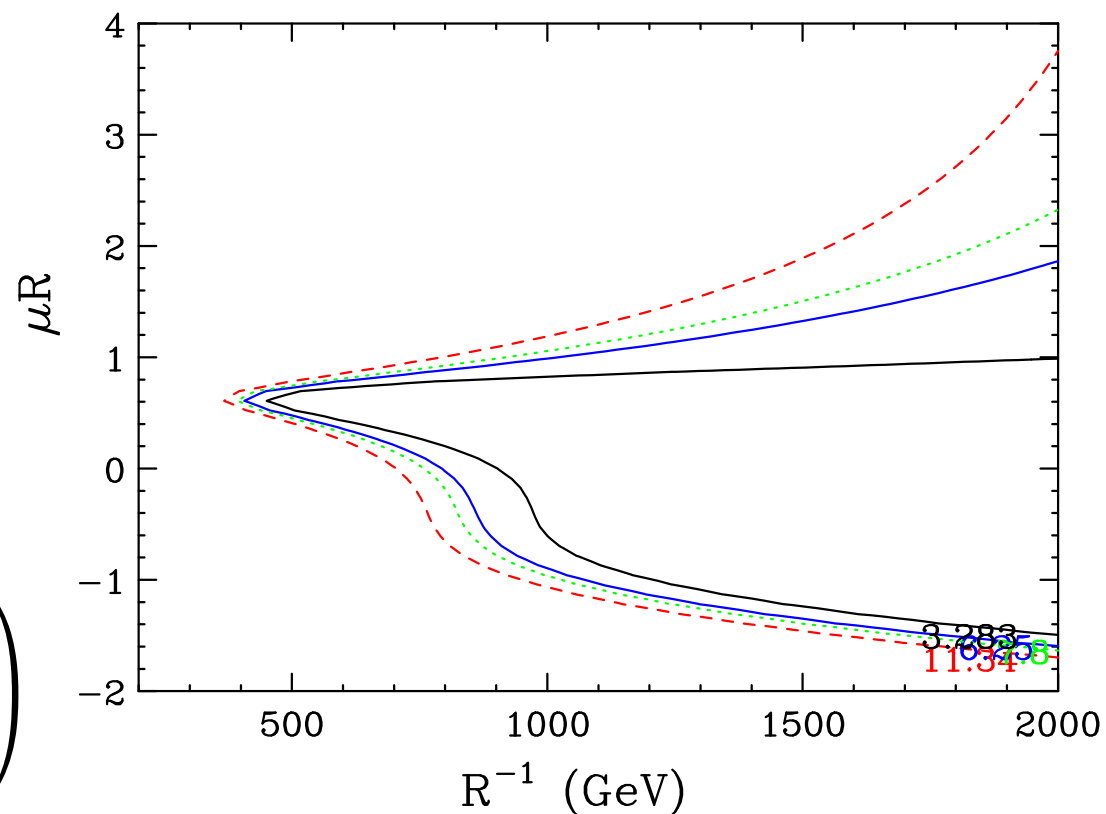
- reference point

$$m_h = 120 \text{ GeV},$$

$$m_t = 173 \text{ GeV}$$

- correlation coeffs

$$\begin{pmatrix} 1 & 0.89 & -0.45 \\ 0.89 & 1 & -0.69 \\ -0.45 & -0.69 & 1 \end{pmatrix}$$



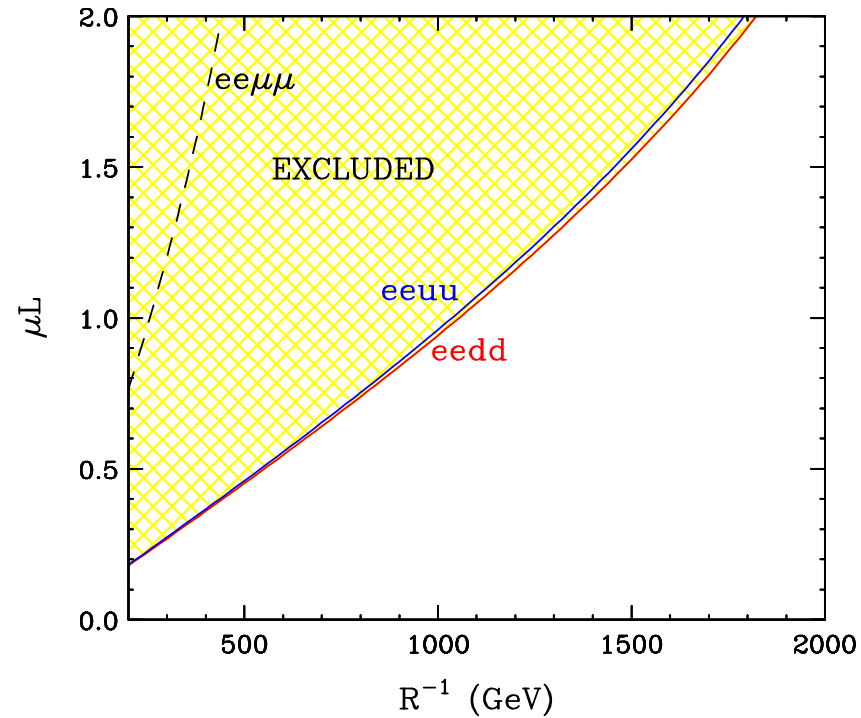
KK EW gauge boson contribute to 4-point interactions.
 PDG bounds for quark lepton compositeness

TeV	$eeee$	$e\epsilon\mu\mu$	$e\epsilon\tau\tau$	$llll$	$qqqq$	$eeuu$	$eedd$
Λ_{LL}^+	> 8.3	> 8.5	> 7.9	> 9.1	> 2.7	> 23.3	> 11.1
Λ_{LL}^-	> 10.3	> 9.5	> 7.2	> 10.3	2.4	> 12.5	> 26.4

Most relevant operators:

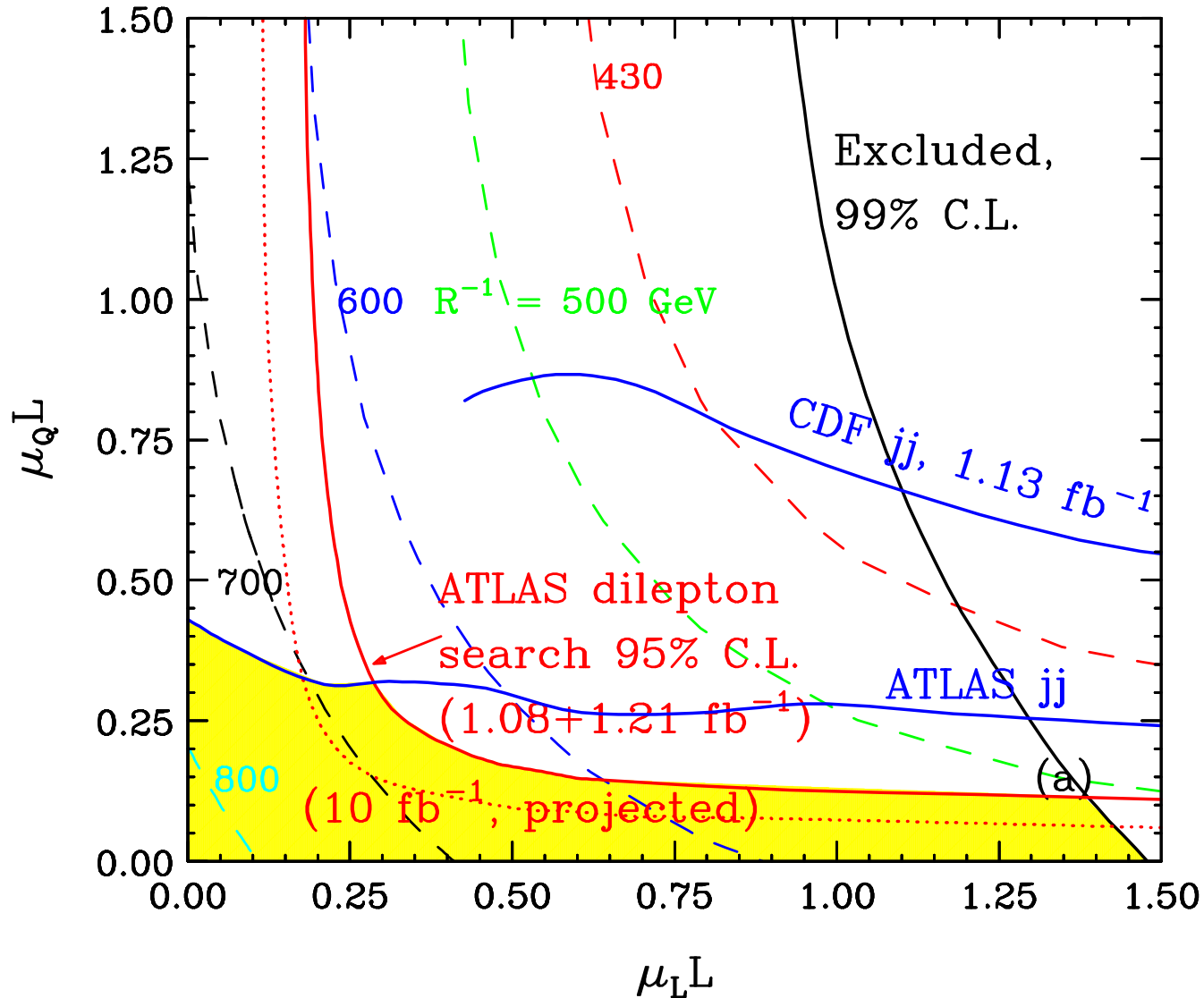
$$\mathcal{L}_{\text{eff}}^{eq} \ni \sum_{q=u,d} \sum_{\{A,B\}=\{L,R\}} \frac{4\pi}{\Lambda_{q,AB}^2} \eta_{AB}^q \bar{e}_A \gamma^\mu e_A \bar{q}_B \gamma_\mu q_B$$

$$\begin{aligned} \frac{4\pi}{\Lambda_{q,AB}^2} \eta_{AB}^q &= 4\pi N_c \sum_{n=1}^{\infty} (\mathcal{F}_{00}^{2n}(\mu R))^2 \times \left[\frac{3}{5} \frac{\alpha_1 Y_{e_A} Y_{q_B}}{Q^2 - M_{B_{2n}}^2} + \frac{\alpha_2 T_{e_A}^3 T_{q_B}^3}{Q^2 - M_{W_{2n}^3}^2} \right] \\ &\approx -\pi N_c R^2 \left(\frac{3}{5} \alpha_1 Y_{e_A} Y_{q_B} + \alpha_2 T_{e_A}^3 T_{q_B}^3 \right) \times \sum_{n=1}^{\infty} \frac{(\mathcal{F}_{00}^{2n}(\mu R))^2}{n^2} \end{aligned}$$

Four-Fermi constraints from $ee\mu\mu$, $eeuu$ and $eedd$ 

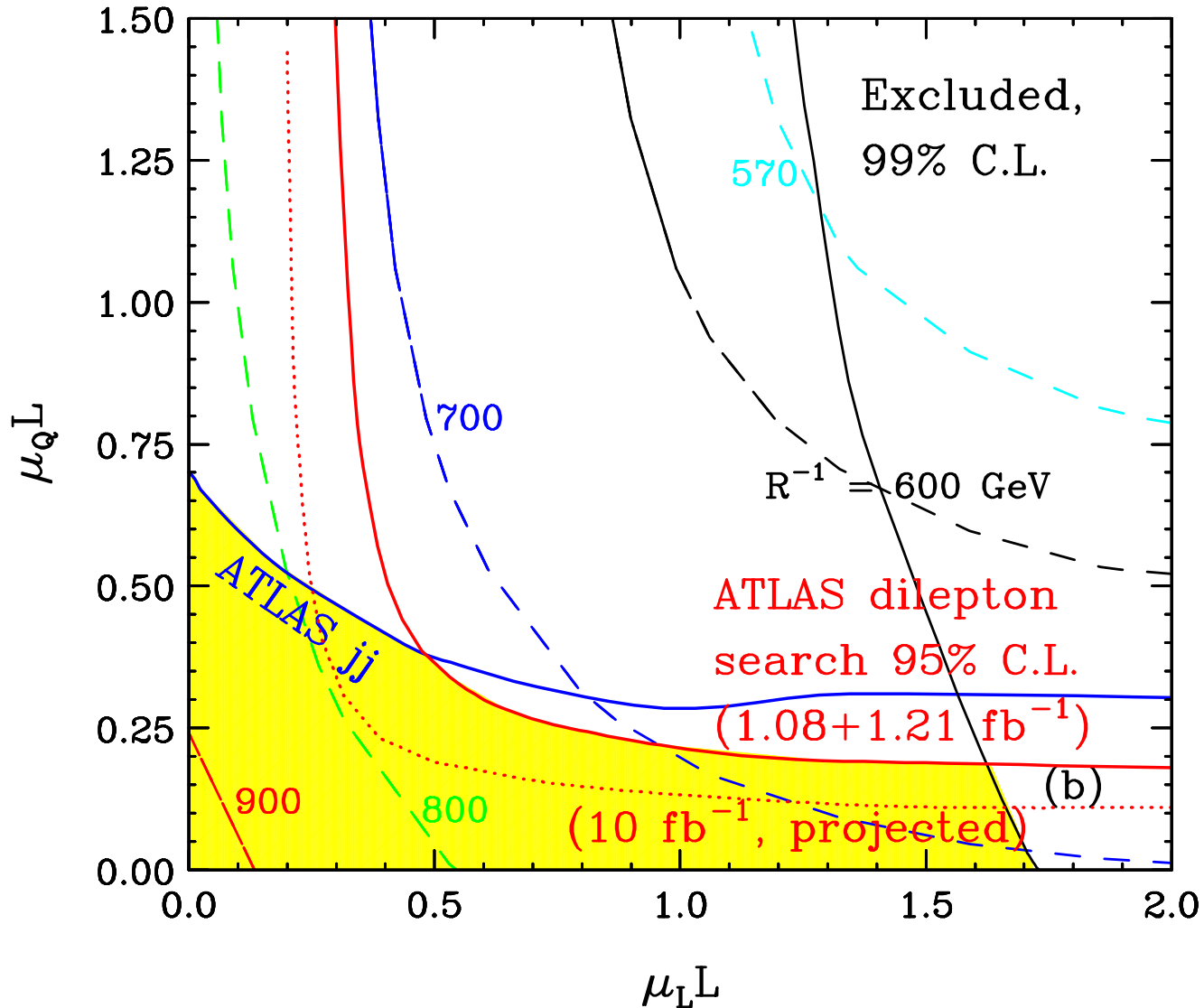
A Non-Universal Case: $\mu_L \neq \mu_Q$

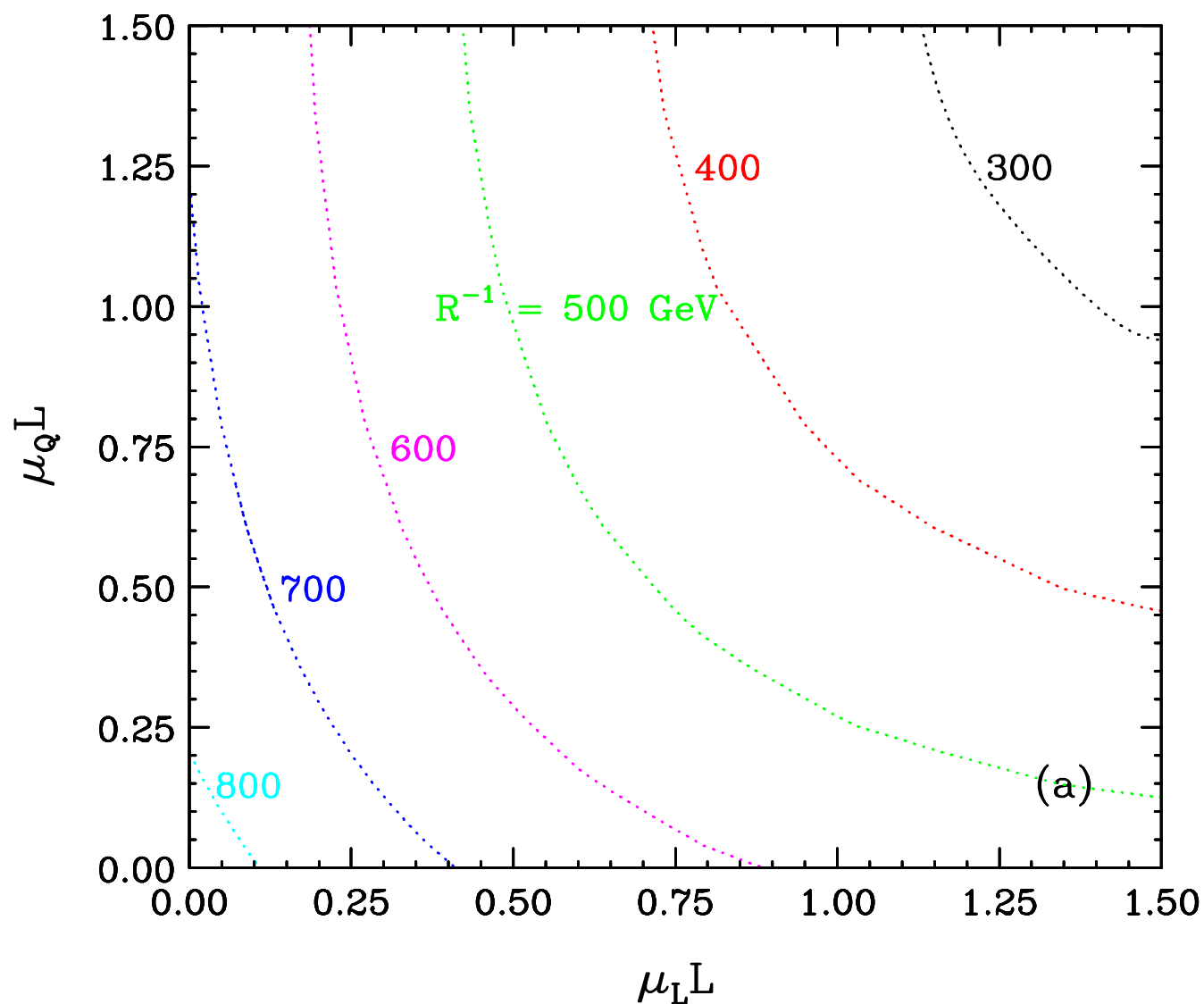
Summary of Constraints (without h_2 resonance)



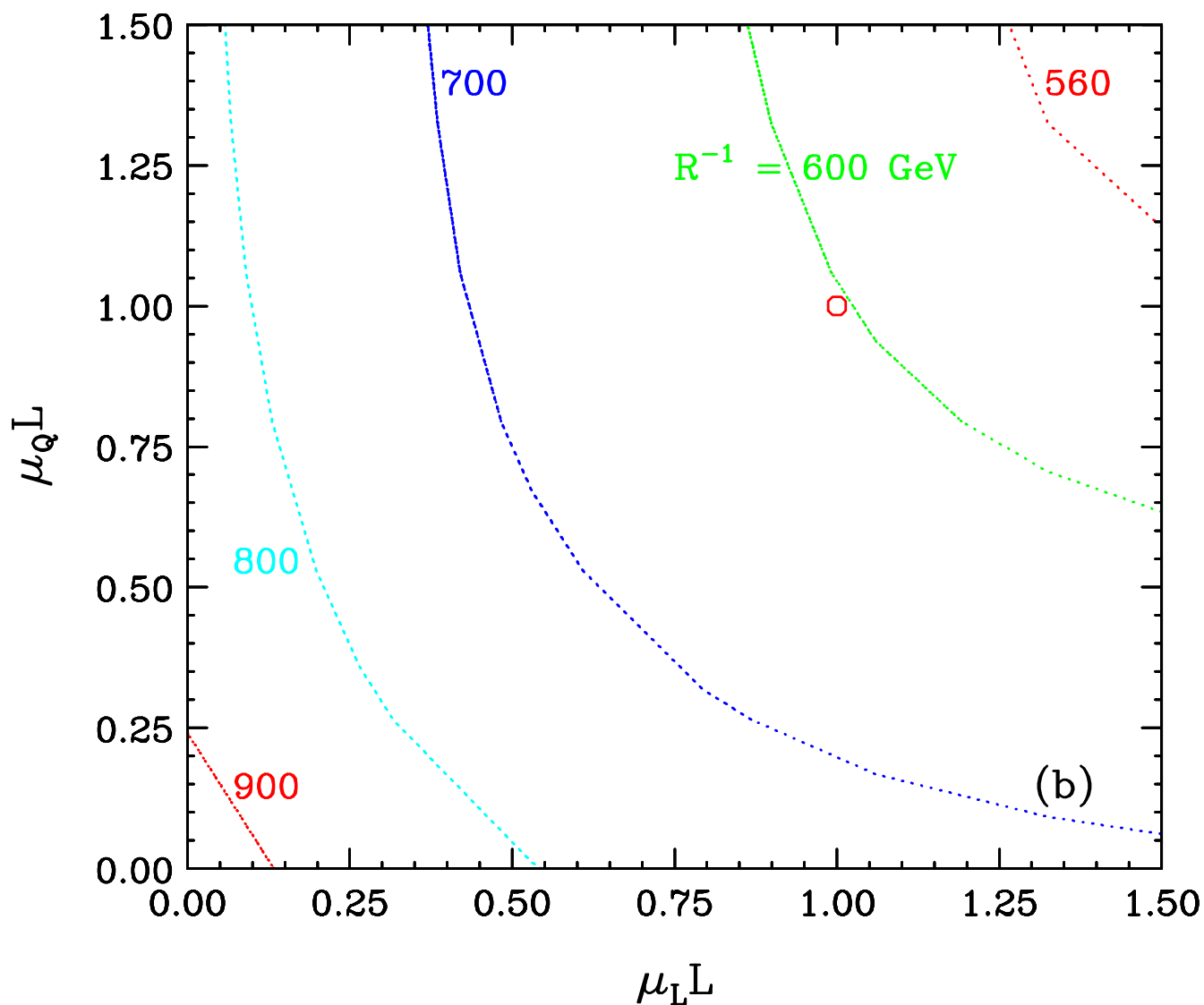
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Summary of Constraints (with h_2 resonance)



Relic abundance (without h_2 resonance)

μ_Q and μ_L factor in similarly with some difference
Contours 'roughly' symmetric about the diagonal

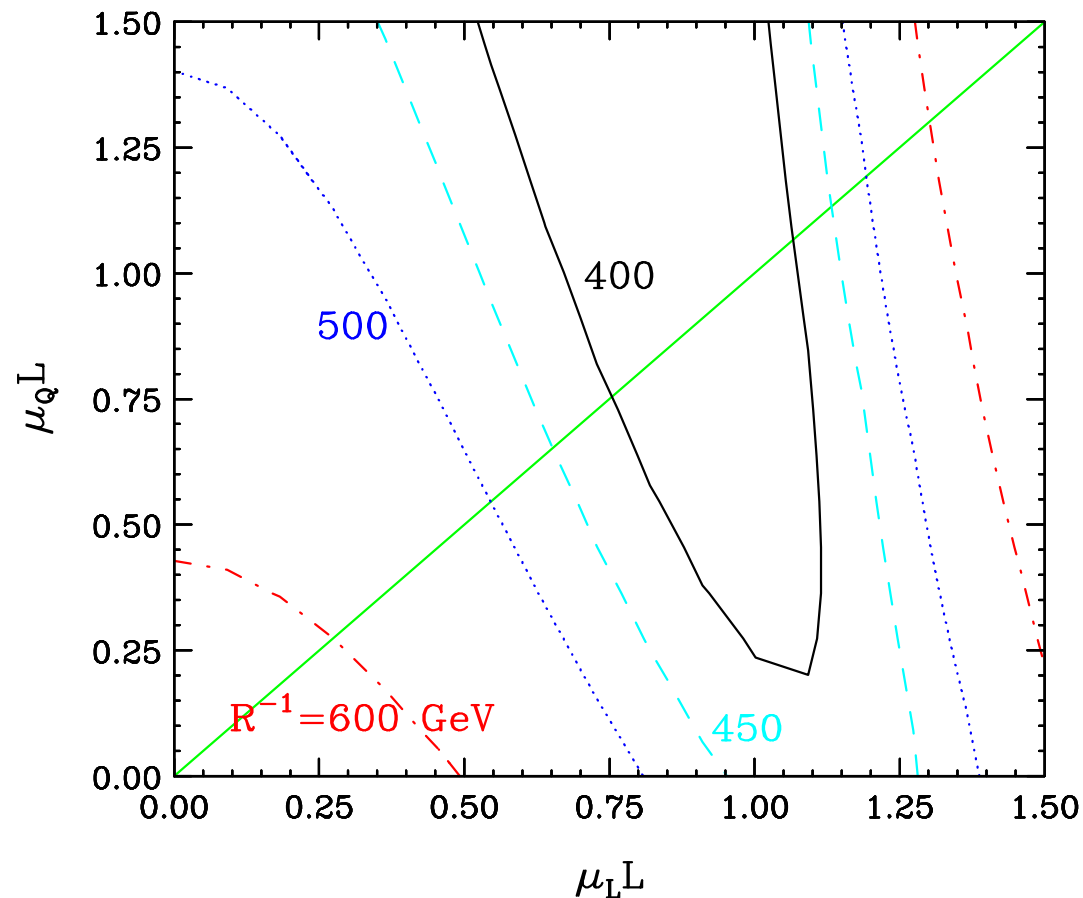
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Simple substitutions

$\mu \rightarrow \mu_Q$ in KK-top loops, $\mu \rightarrow \mu_L$ in δG_F

99% C.L. fit surfaces in (μ_Q, μ_L, R^{-1}) projected down to fit contours in (μ_Q, μ_L)

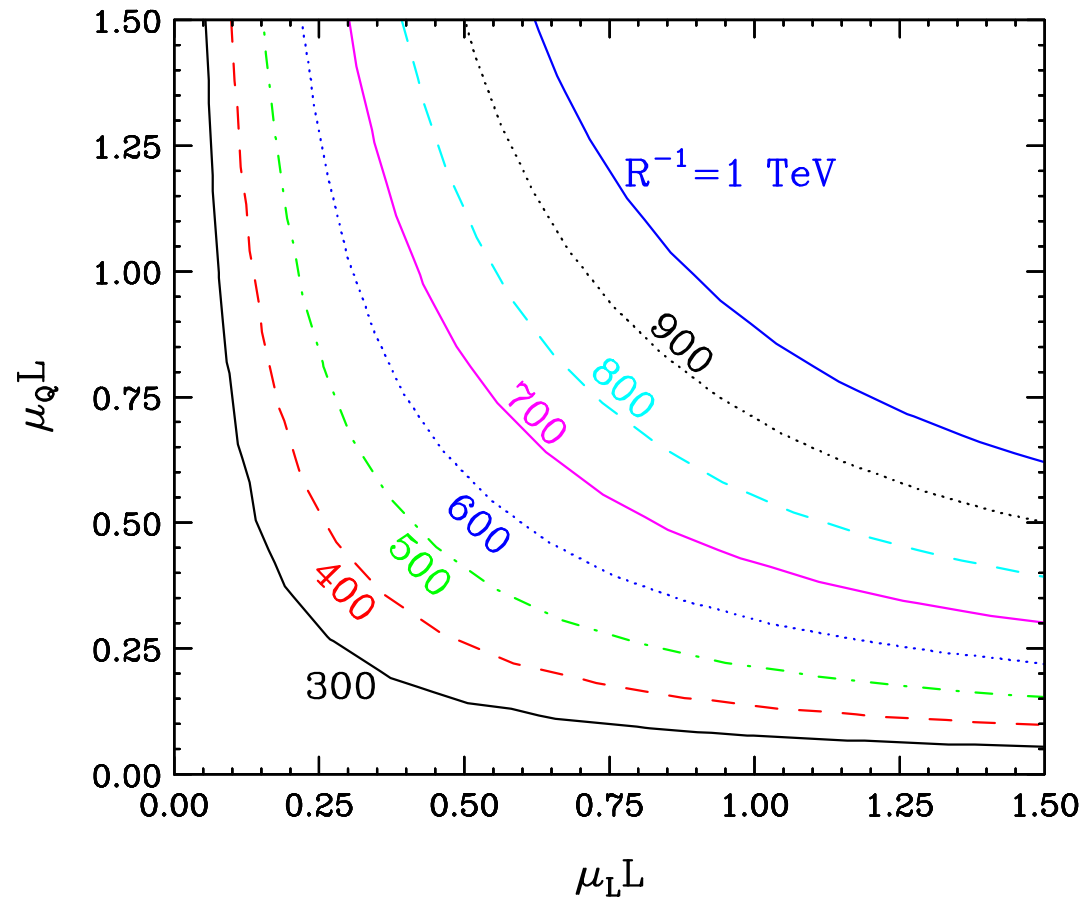


Leptons/Quarks treated differently, (S, T, U) may be insufficient

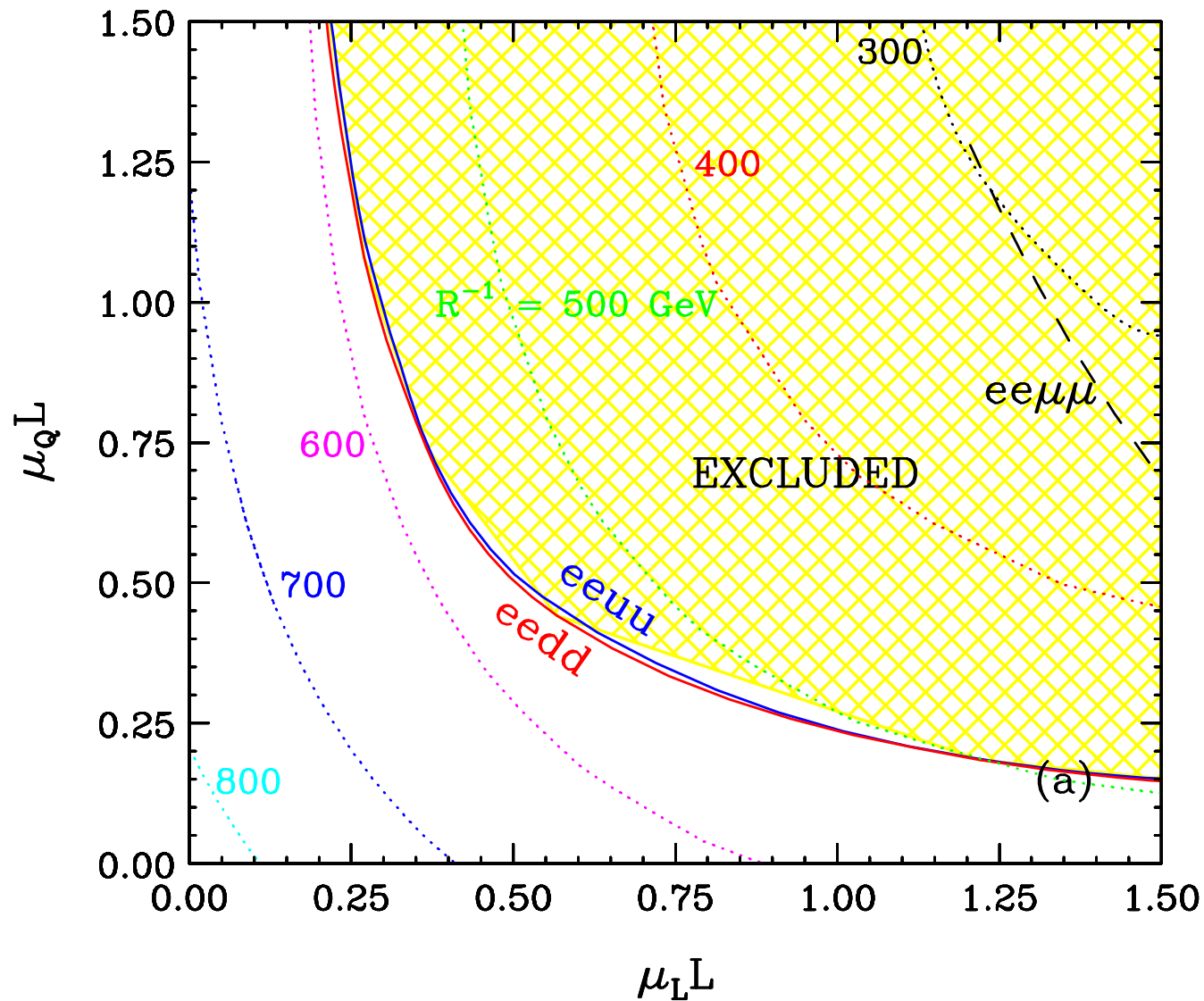
Different dependence on μ_Q, μ_L

$$(\mathcal{F}_{00}^{2n}(\mu R))^2 \rightarrow \mathcal{F}_{00}^{2n}(\mu_Q R) \mathcal{F}_{00}^{2n}(\mu_L R)$$

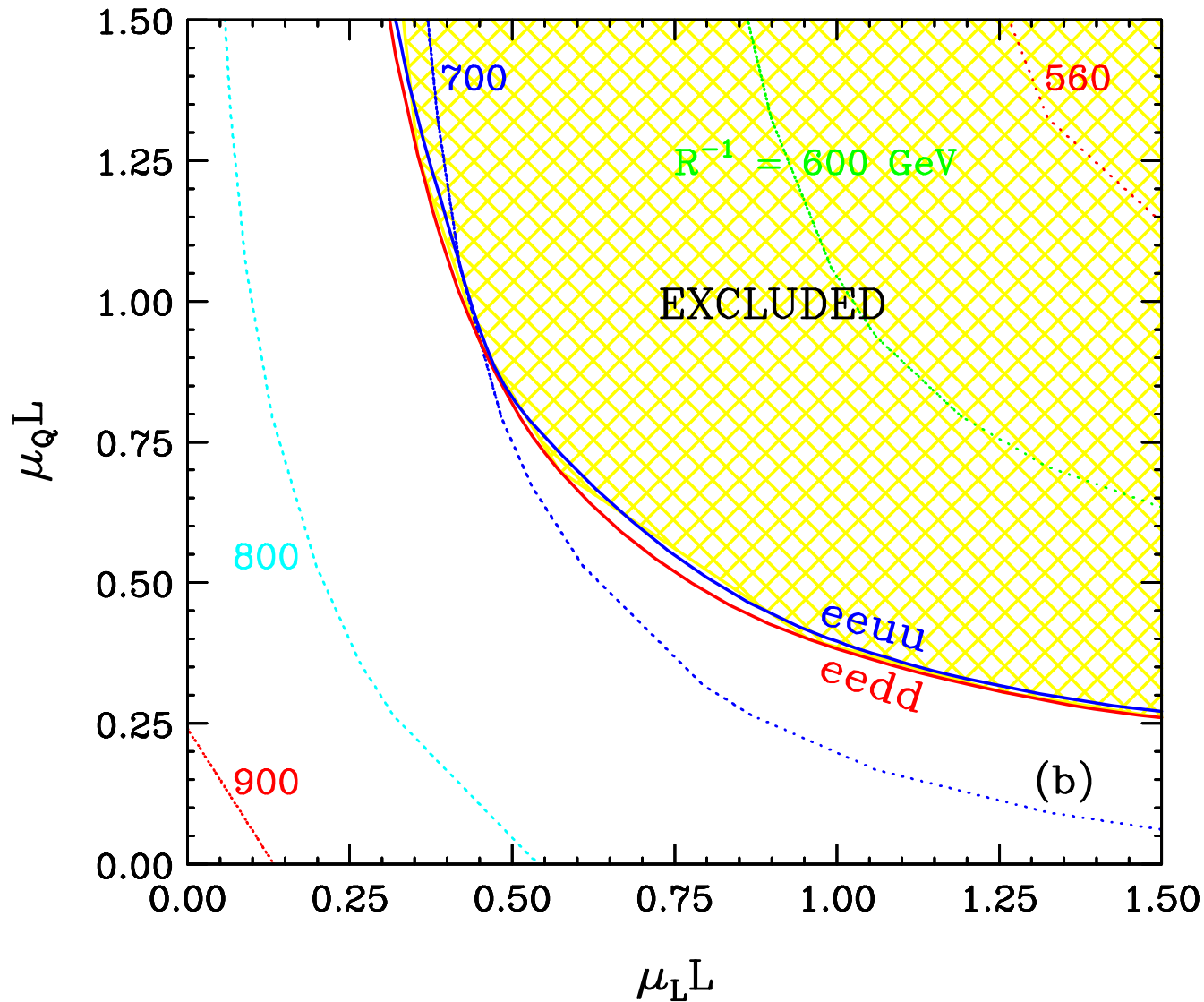
eedd constraint surface projected down



Four-Fermi combined/intersected with relic abundance (without h_2 resonance)



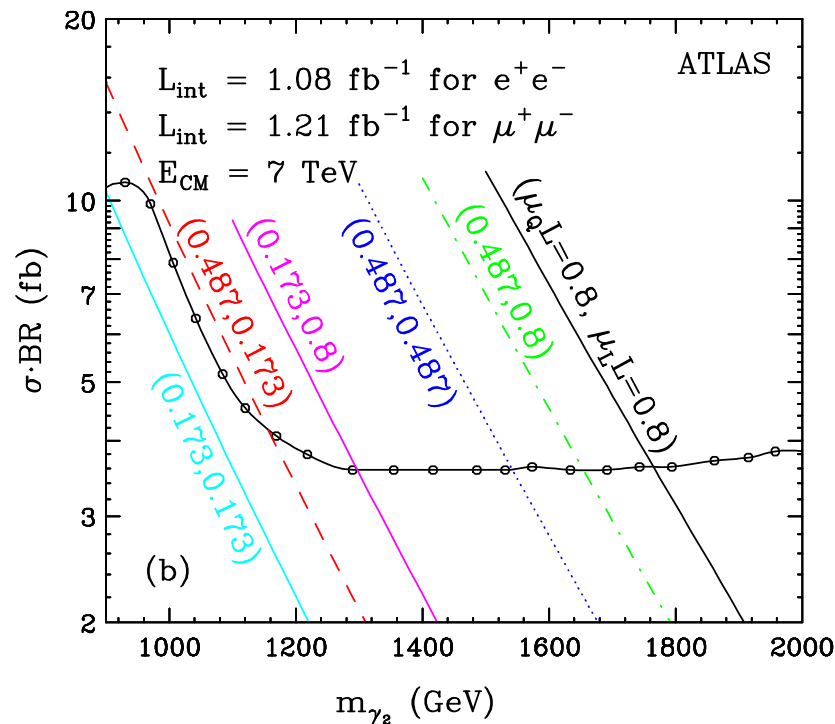
Four-Fermi combined/intersected with relic abundance (with h_2 resonance)



jj channel, only μ_Q relevant

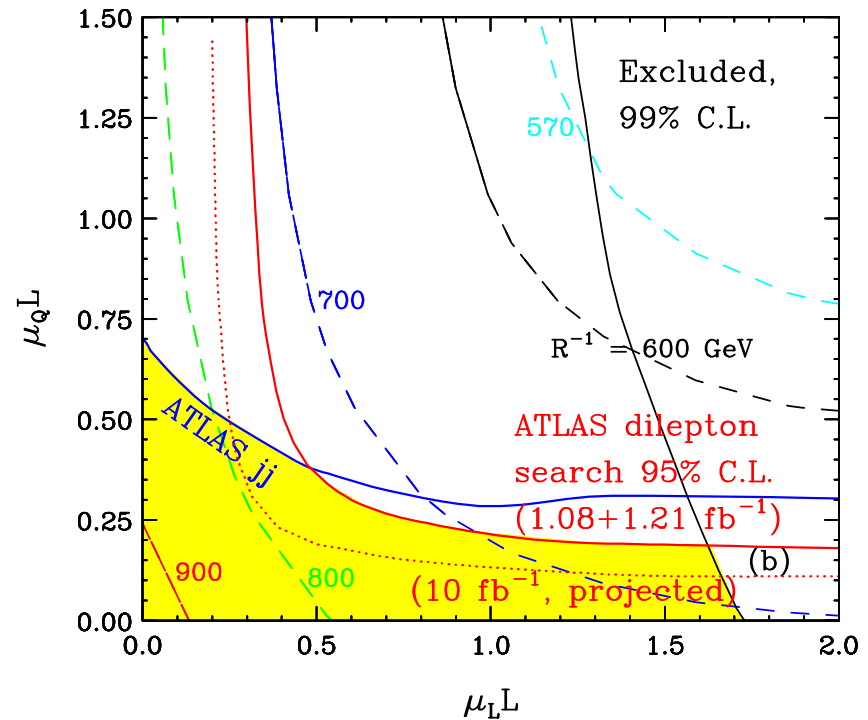
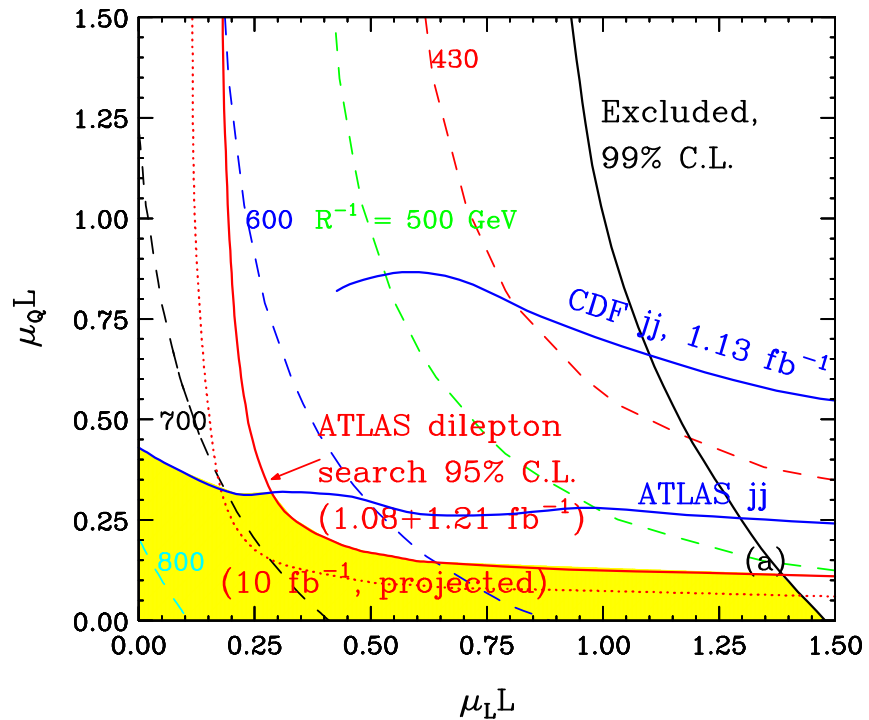
$ll, l\nu$ channels, both μ_Q and μ_L

ll channel search limit



jj channel search limit

Identical to the Universal case, with $\mu \rightarrow \mu_Q$



without h_2 resonance

$$500\text{GeV} \lesssim 1/R \lesssim 850\text{GeV}, \mu_Q L \lesssim 0.4, \mu_L L \lesssim 1.5$$

with h_2 resonance

$$650\text{GeV} \lesssim 1/R \lesssim 950\text{GeV}, \mu_Q L \lesssim 0.7, \mu_L L \lesssim 1.7$$

Lower bound on $1/R$ relaxed, μ_L less constrained than μ_Q

SUED parameter space

- Tightly Constrained

R^{-1} within a couple hundred GeV range, below TeV
prefers small μ

- Slightly Relaxed

UED \rightarrow SUED with $\mu \rightarrow$ SUED with μ_Q, μ_L

- LHC

Key role in remaining parameter space
Used $\sim 1 \text{ fb}^{-1}$ at 7 TeV, has $> 5 \text{ fb}^{-1}$
 15 fb^{-1} at 8 TeV by 2012

- SUED $S=?$