

Top Decays with Flavor Changing Neutral Higgs Interactions at the LHC

Chung Kao
University of Oklahoma

†Presented at the 2012 Phenomenology Symposium at the University of Pittsburgh, Pittsburgh, Pennsylvania, May 7-9, 2012.

Top Decays with Flavor Changing Neutral Higgs Interactions at the LHC

Kao, Cheng, Hou, and Sayre, [arXiv:1112.1707](#)

- ~ A Special Two Higgs Doublet Model for the Top
- ~ Flavor Changing Neutral Higgs Boson in Top Decays
- ~ The Physics Background
- ~ Realistic Acceptance Cuts
- ~ The Discovery Potential at the LHC
- ~ Comparison of Production Rates
- ~ Conclusions

A Special Two Higgs Doublet Model for the Top Quark

Das and Kao (1996)

- ~ We propose that the top quark is the only elementary fermion getting a mass from a much larger VEV of a second Higgs doublet.
- ~ The top quark is naturally heavier than other quarks and leptons in the 3 known generations.
- ~ The ratio of the Higgs VEVs, $\tan(\beta) = |v_2|/|v_1|$, is naturally large, which enhances the Yukawa couplings of the lighter quarks and leptons with the Higgs bosons.
- ~ There are flavor changing neutral Higgs (FCNH) interactions among the up type quarks.

A Special Two Higgs Doublet Model for the Top Quark

We choose the Lagrangian density of Yukawa interactions be of the following form:

$$\begin{aligned}\mathcal{L}_Y = & - \sum_{m,n=1}^3 \bar{L}_L^m \phi_1 E_{mn} l_R^n - \sum_{m,n=1}^3 \bar{Q}_L^m \phi_1 F_{mn} d_R^n \\ & - \sum_{\alpha=1}^2 \sum_{m=1}^3 \bar{Q}_L^m \tilde{\phi}_1 G_{m\alpha} u_R^\alpha - \sum_{m=1}^3 \bar{Q}_L^m \tilde{\phi}_2 G_{m3} u_R^3 \\ & + \text{H.c.}\end{aligned}$$

Flavor Changing Neutral Higgs Interactions in Top Decays

Kao, Cheng, Hou, and Sayre, [arXiv:1112.1707](#)

- Let us consider the following Lagrangian involving flavor changing neutral Higgs interactions with top and charm quarks:

$$\mathcal{L} = -\lambda_{tc}\bar{t}cH^0 - i\lambda_{tc}\bar{t}\gamma_5cA^0 + \text{H.c.}$$

where H^0 is a scalar and A^0 is a pseudoscalar.

- This is a general feature of Model III of Yukawa Interactions in Two Higgs Doublet Models.

FCNH Yukawa Coupling

Cheng and Sher (1987)

- ~ Let us consider the FCNH coupling of $t\bar{c}H$ to be the geometric mean of the Yukawa couplings of the quarks:

$$\lambda_{tc} = \frac{\sqrt{m_t m_c}}{v}$$

- ~ In general, we will take it as a free parameter.

Top Decay Width

Hou (1991)

~ The FCNH top decay width is

$$\Gamma(t \rightarrow c\phi^0) = \frac{|\lambda_{tc}|^2}{16\pi} \times (m_t) \times [(1 \pm \rho_c)^2 - \rho_\phi^2] \\ \times \sqrt{1 - (\rho_\phi + \rho_c)^2} \sqrt{1 - (\rho_\phi - \rho_c)^2}$$

$\rho_c = m_c/m_t$, $\rho_H = M_H/m_t$, + for H^0 and - for A^0 .

~ The total width is

$$\Gamma_t = \Gamma(t \rightarrow bW) + \Gamma(t \rightarrow c\phi^0)$$

FCNH Branching Fraction

As a case study, we take the FCNH Yukawa couplings to be the geometric mean of the Yukawa couplings of the quarks with $m_t = 173.3$ GeV and $m_c = 1.4$ GeV:

$$\lambda_{tc} = \frac{\sqrt{m_t m_c}}{v} \simeq 0.063$$

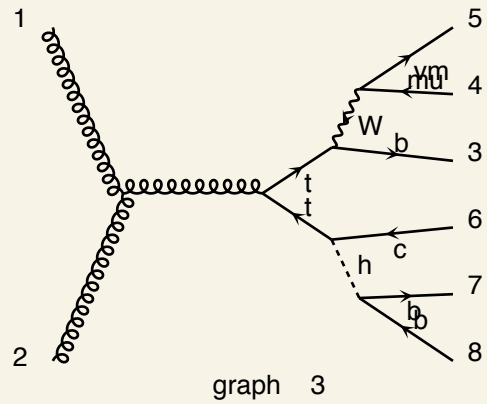
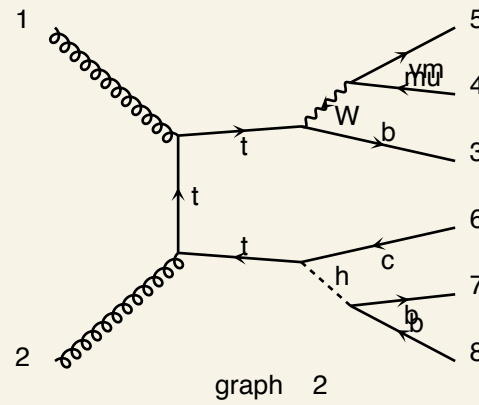
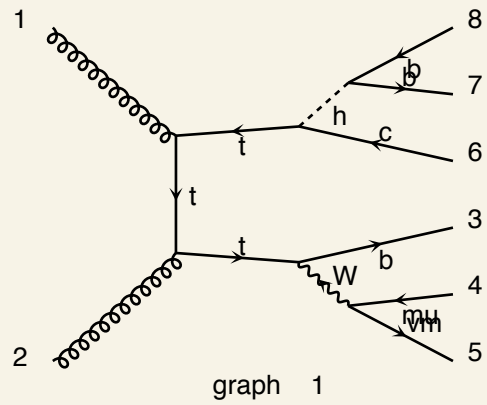
$$\mathcal{B}(t \rightarrow c\phi^0) = 2.6 \times 10^{-3} \quad \text{for } M_\phi = 120 \text{ GeV},$$

$$\mathcal{B}(t \rightarrow c\phi^0) = 6.2 \times 10^{-4} \quad \text{for } M_\phi = 150 \text{ GeV}.$$

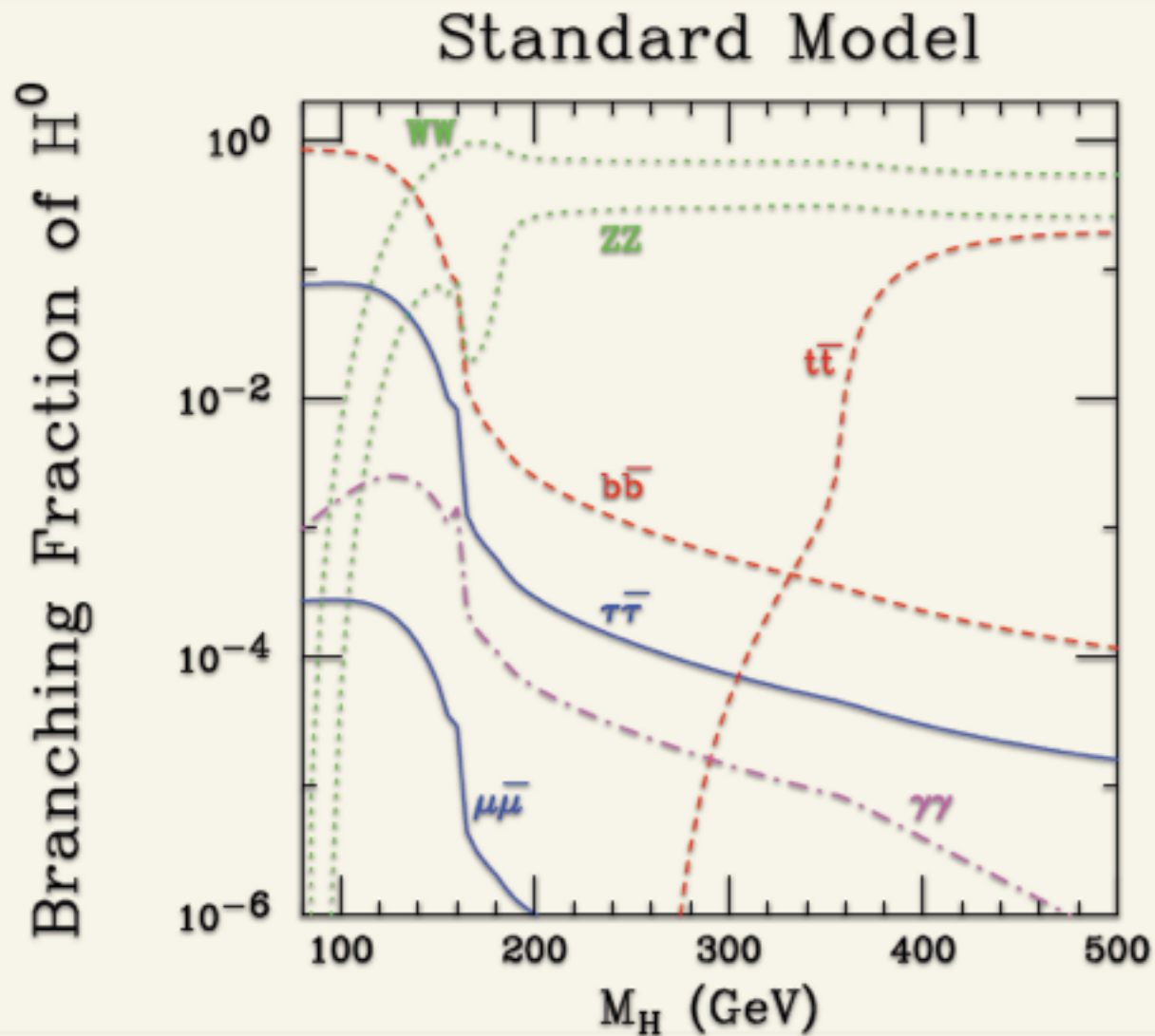
The FCNH Signal

Diagrams by MadGraph

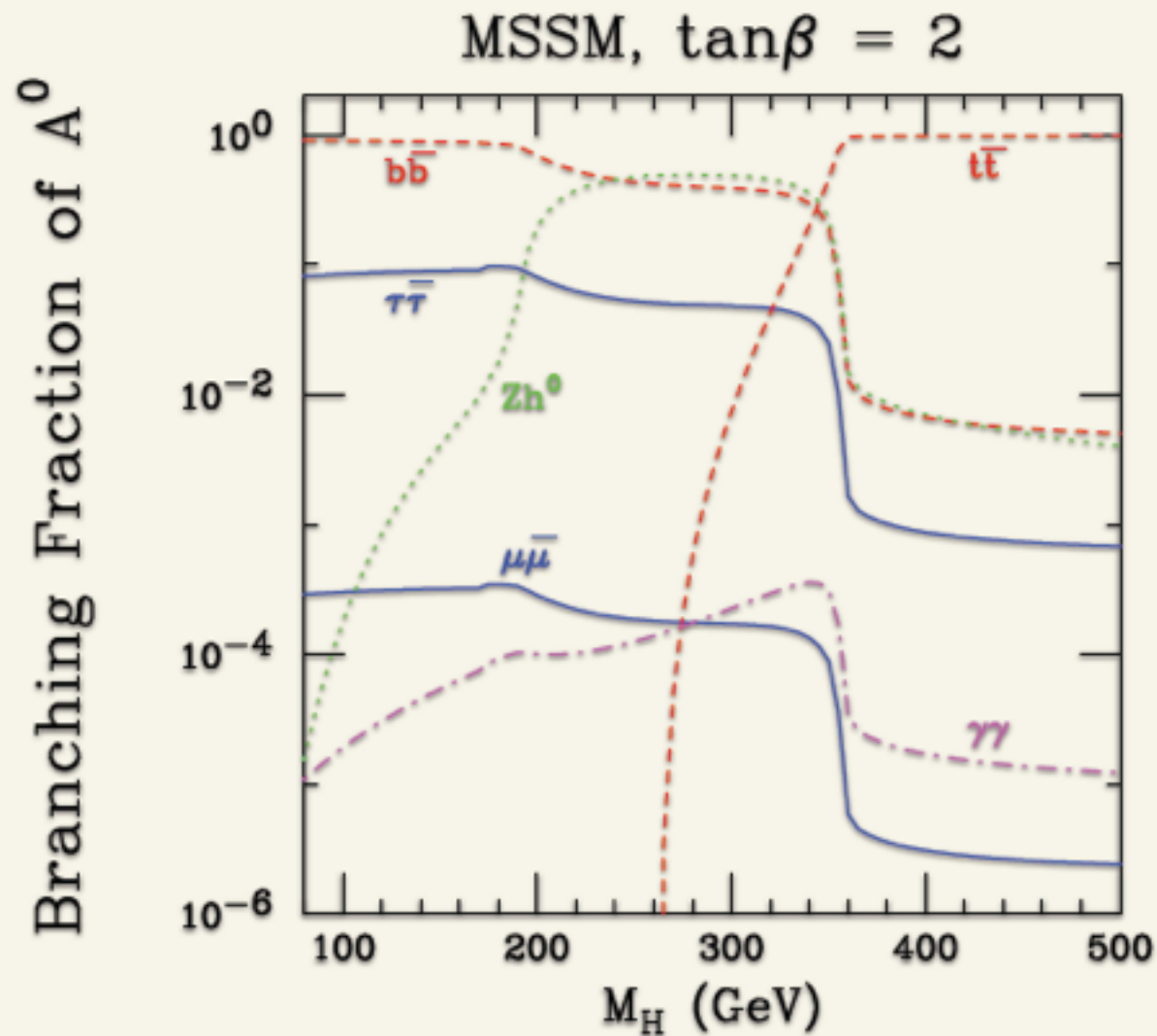
$g g \rightarrow b \mu^+ \nu_\mu c \bar{b}$



Branching Fractions of the Higgs Boson



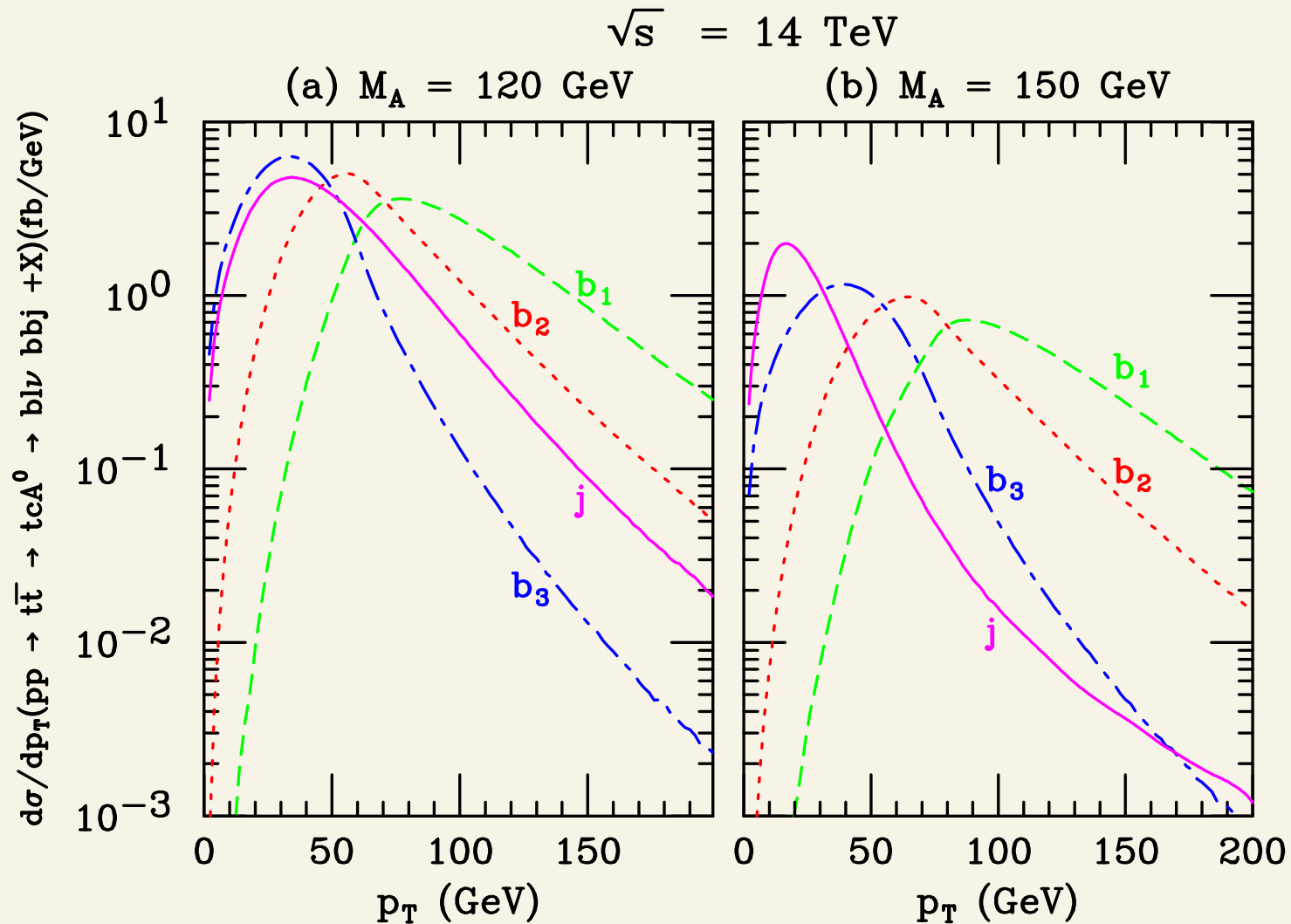
Branching Ratios of a Higgs Pseudoscalar



The FCNH Signal at the LHC

- ~ We employ the programs MadGraph and HELAS to evaluate the exact matrix element for the FCNH signal from gluon fusion and quark-antiquark annihilation in pp collisions.
Stelzer and Long (1994); Alwall et al. (2007); Murayama, Watanabe and Hagiwara (1991).
- ~ In addition, we apply narrow width approximation to check the exact results.
- ~ The cross sections are evaluated with the parton distribution functions of CTEQ6L1.

Transverse Momentum Distribution for the Higgs Signal



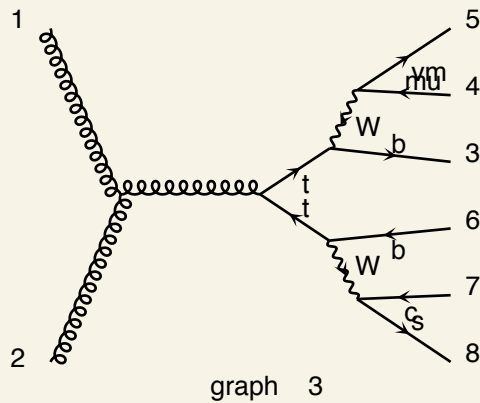
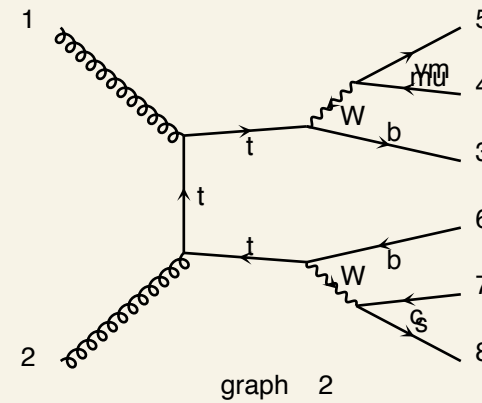
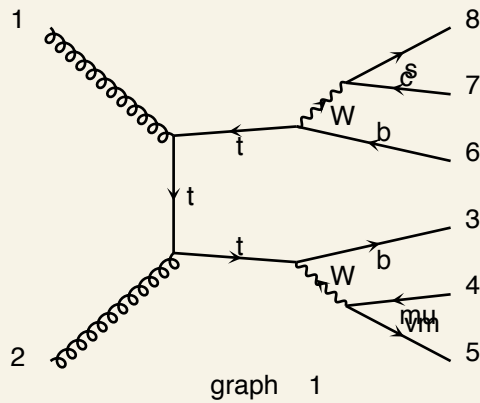
FCNH Signal Cross Section

M_A	$s(\text{bmncbb})$	$B(\text{t cH})$	$\text{Gamma}(\text{H})$	$B(\text{H bb})$
120.0	0.440E+02	0.259E-02	0.351E-02	0.728E+00
140.0	0.820E+01	0.117E-02	0.428E-02	0.677E+00
150.0	0.268E+01	0.621E-03	0.473E-02	0.649E+00

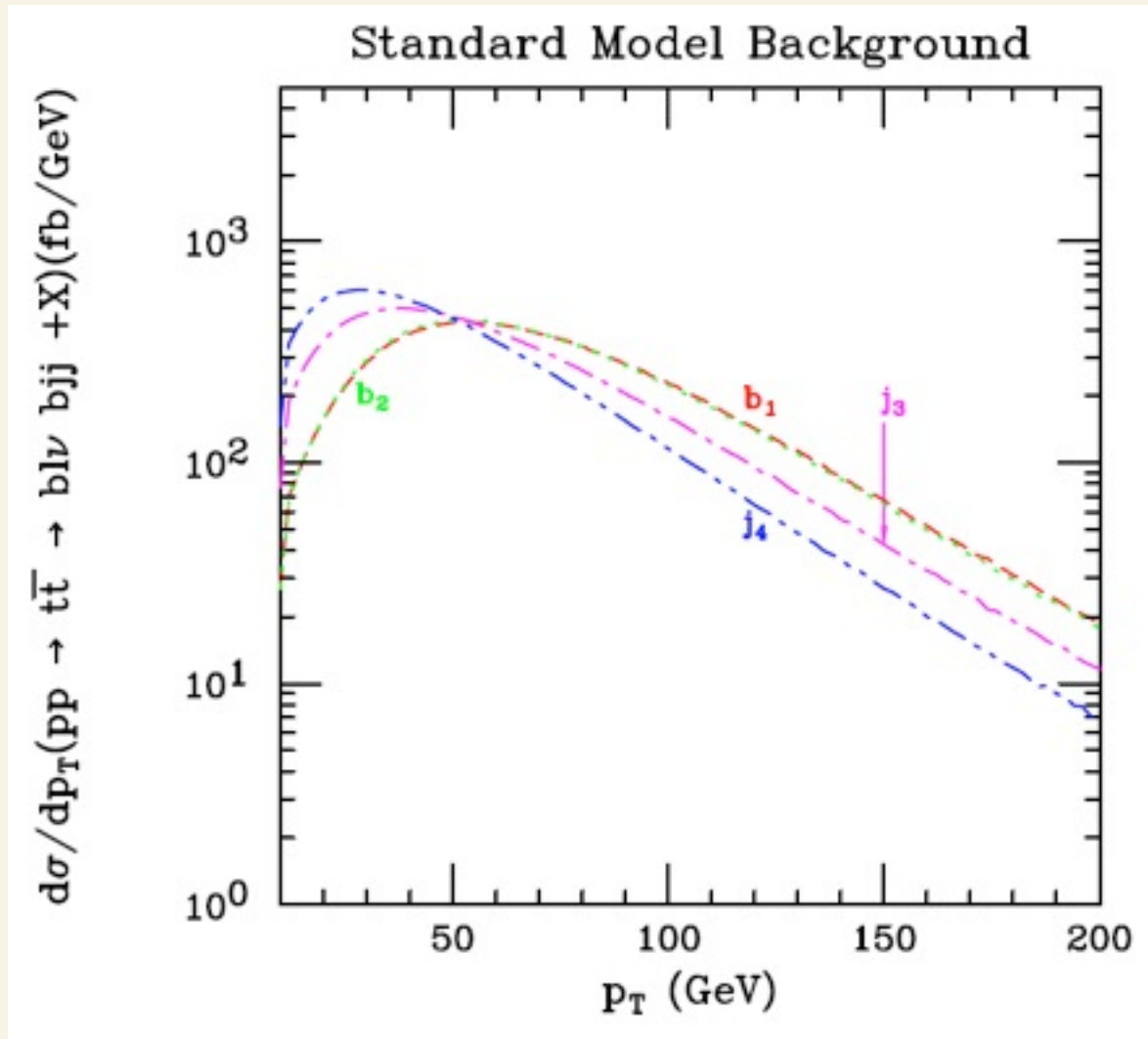
Dominant Physics Background from top quark pairs

Diagrams by MadGraph

$g g \rightarrow b \mu^+ \nu_m b \bar{c} \bar{s}$



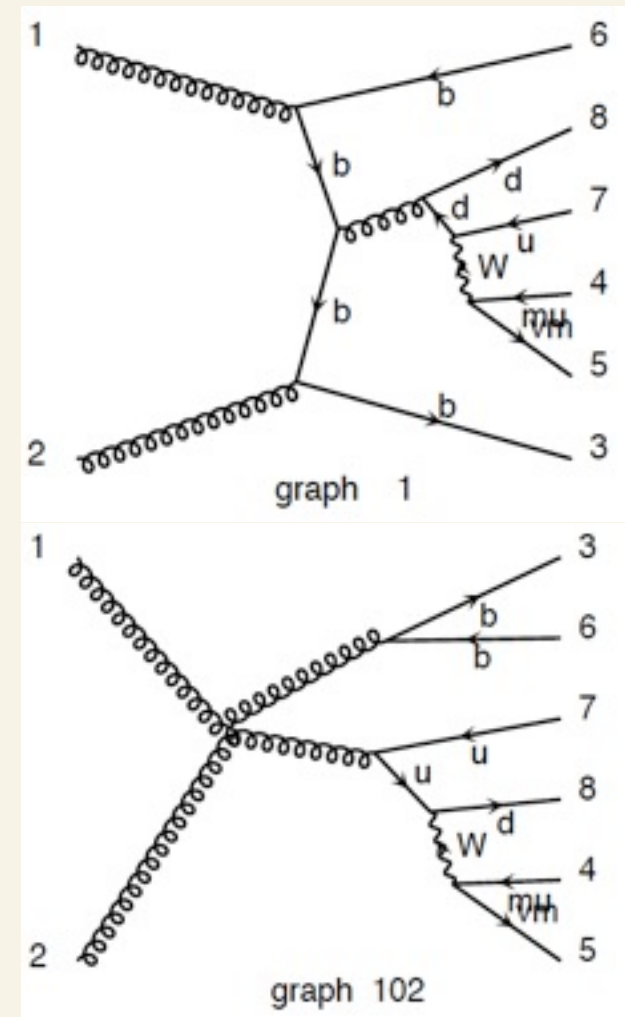
Transverse Momentum Distribution for the Physics Background



Additional Backgrounds

gg, qq to Wbbjj

- ~ We have included additional backgrounds with Wbbjj:
- ~ gg to Wbbjj
- ~ $q_v q_v$ to Wbbjj
- ~ gg to Wbbjj
- ~ $q_v q_s$ to Wbbjj

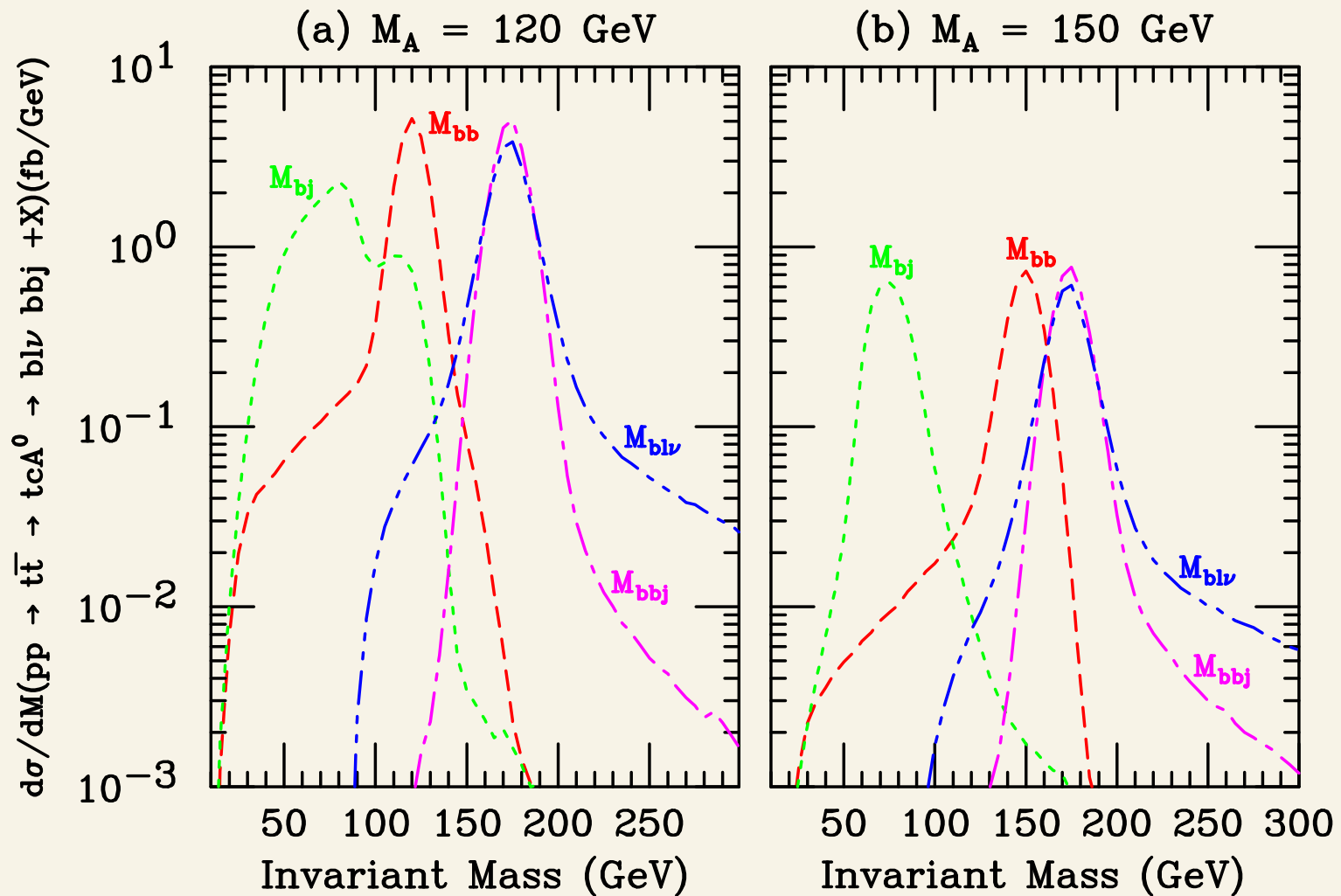


Mass Reconstruction

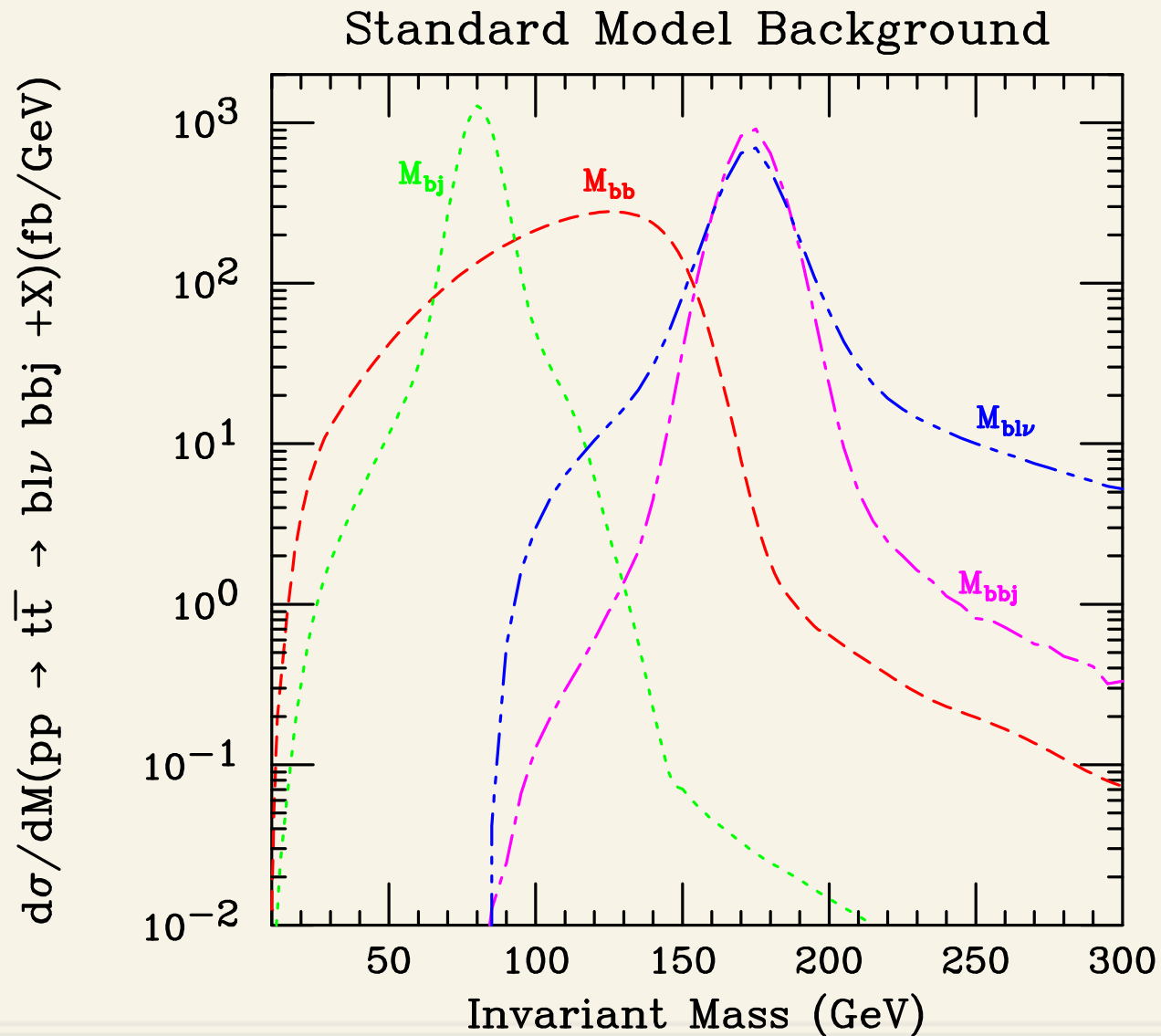
- ~ Since our FCNC signal comes from one top quark decay, we will choose the pair of b jets that minimize $|M_{bbj}-m_t|$ as b_1b_2 and label the other b jet as b_3 .
- ~ For a correctly reconstructed event, b_1 and b_2 are the products of a Higgs decay as well, such that their invariant mass has a peak near M_H .
- ~ For a background event, we identify b_2 as the member of this pair that minimizes $|M_{bj}-m_W|$.
- ~ The remaining b quark (b_3) should reproduce m_t with the charged lepton and neutrino momenta.

Invariant Mass: FCNH Signal

$$\sqrt{s} = 14 \text{ TeV}$$



Invariant Mass: Physics Background



Realistic Acceptance Cuts

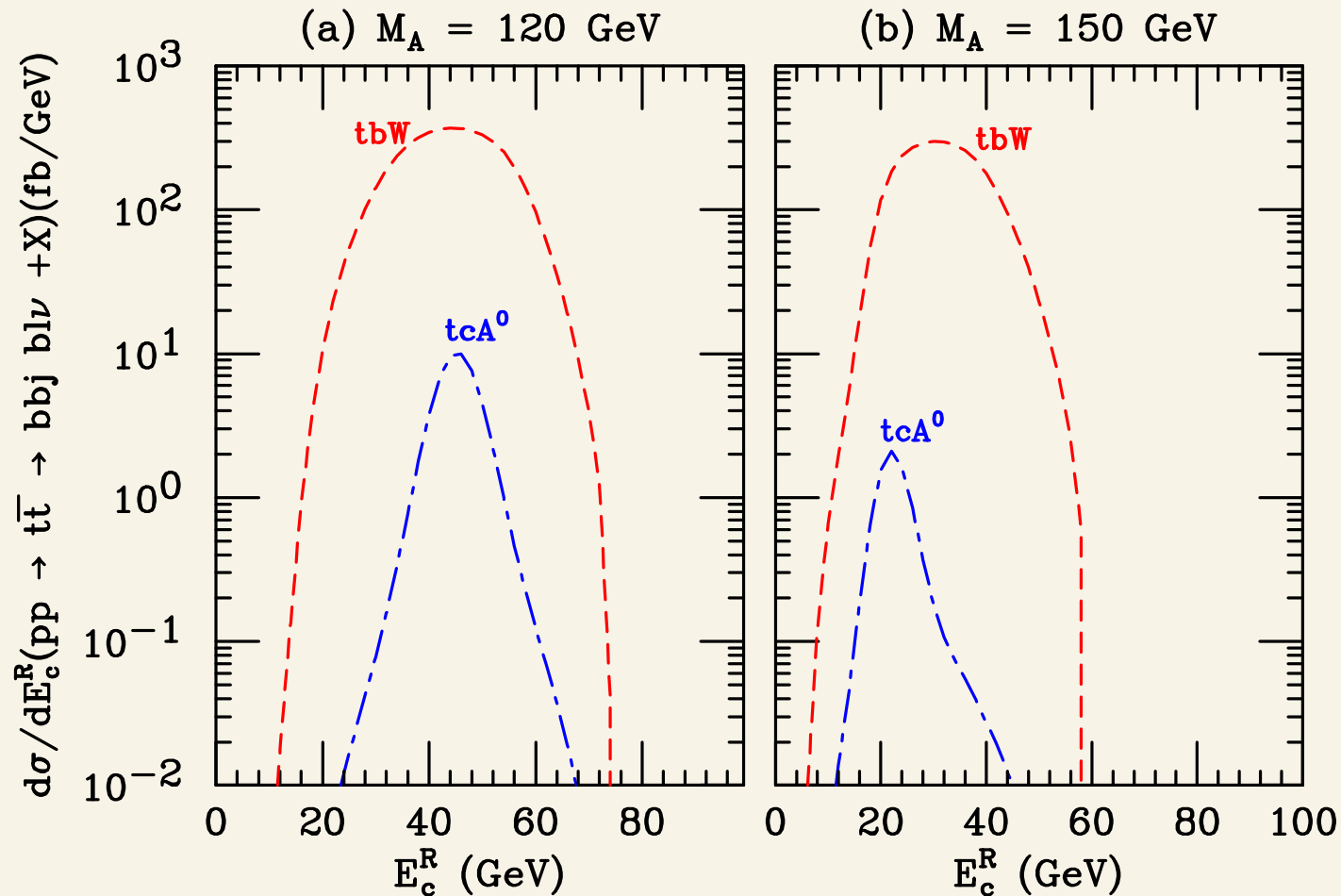
For (a) the early LHC and (b) full CM energy with low luminosity, we require that in every event there should be

- ▶ exactly 4 jets that have $p_T > 15$ GeV and $|\eta| < 2.5$, and three of them must be tagged as b -jets;
- ▶ exactly one isolated lepton with $p_T > 20$ GeV and $|\eta| < 2.5$;
- ▶ the missing transverse energy must be greater than 20 GeV;
- ▶ at least one pair of b -jets such that the invariant mass of $b_1 b_2$ should be near m_t : $|M_{b_1 b_2} - m_t| \leq 25$ GeV;
- ▶ the pair of b -jets, $b_1 b_2$, that reconstructs the hadronically decaying top should also satisfy $|M_{b_1 b_2} - M_\phi| \leq 0.15 M_\phi$;
- ▶ a third b jet such that the invariant mass of $b_3 \ell \nu$ should be near m_t : $|M_{b_3 \ell \nu} - m_t| \leq 25$ GeV;
- ▶ the reconstructed W must satisfy $|M_{\ell \nu} - m_W| \leq 15$ GeV.

Reconstructed E_{charm}

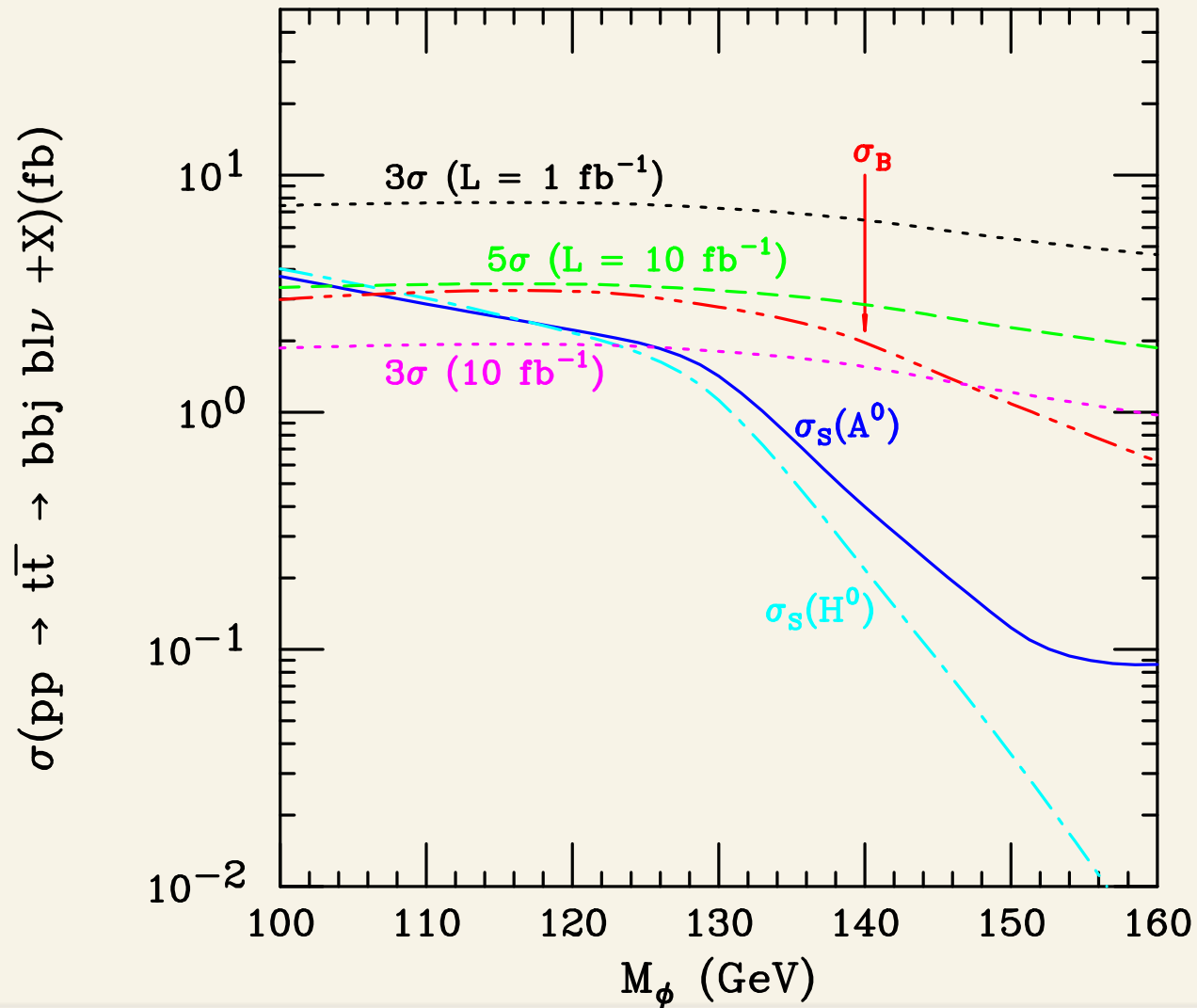
Han, Jiang, and Sher (2001)

$\sqrt{s} = 14 \text{ TeV}$



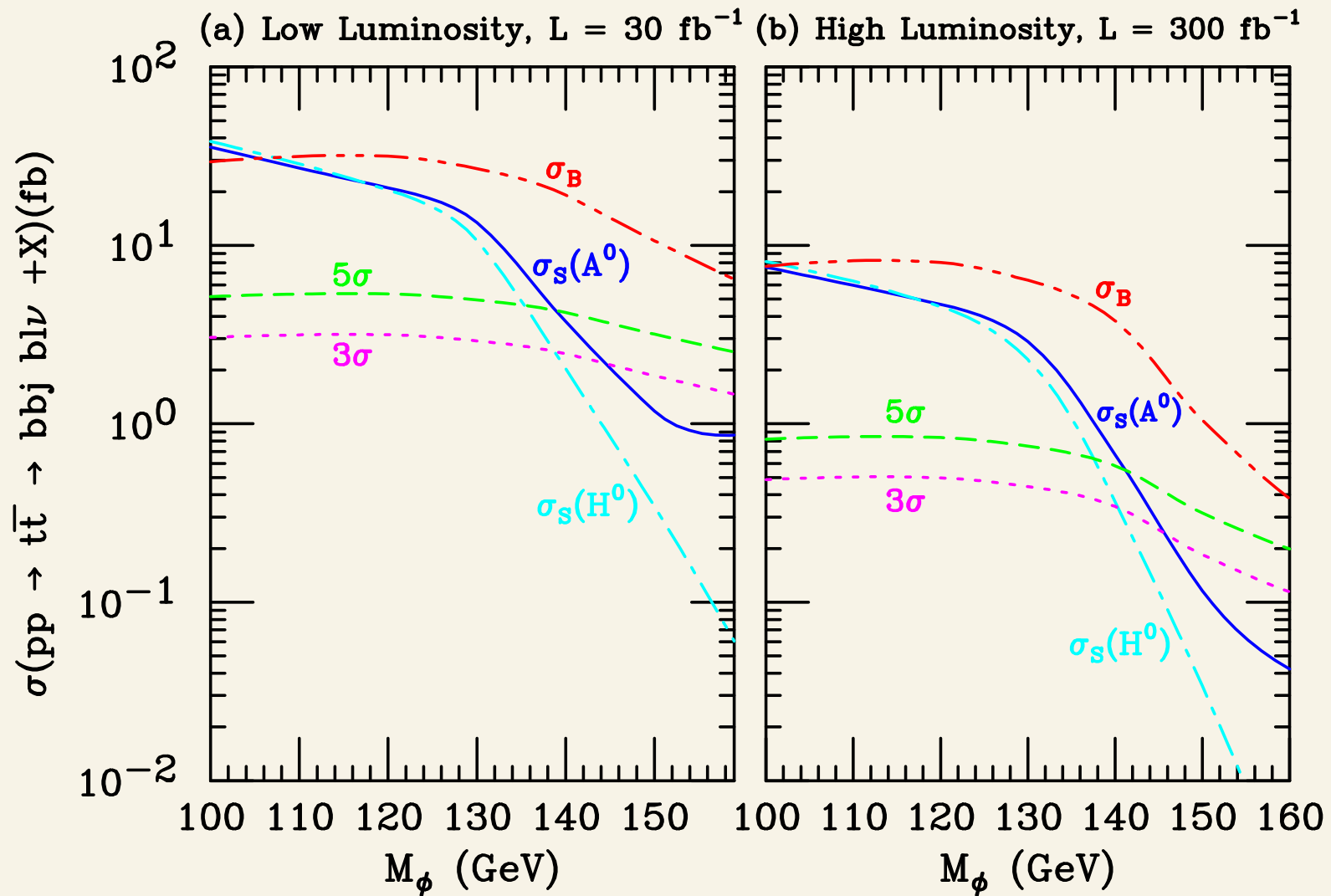
Signal versus Background

$$\sqrt{s} = 7 \text{ TeV}$$



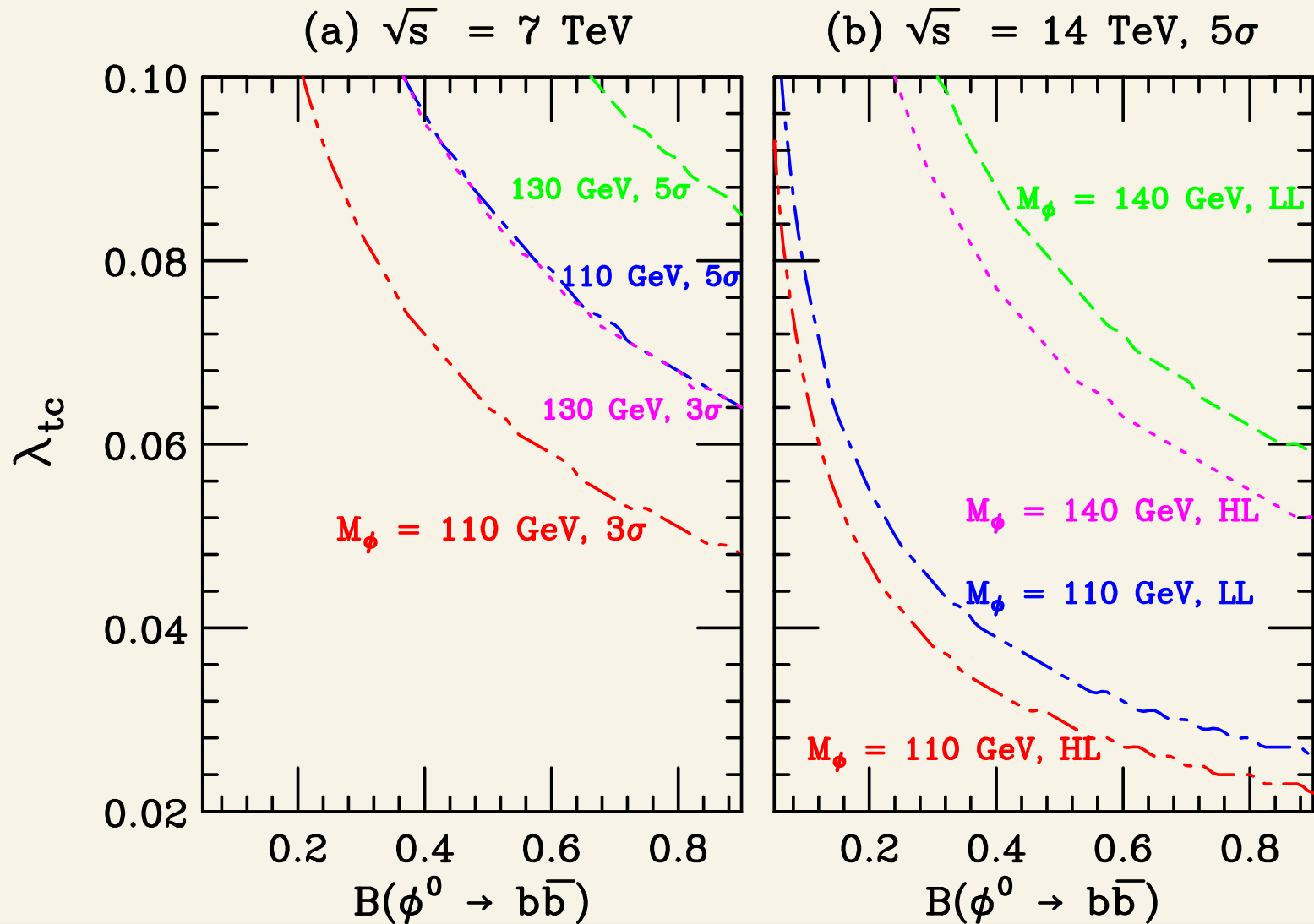
Signal versus Background

$$\sqrt{s} = 14 \text{ TeV}$$



Discovery Contours

$L = 10 \text{ fb}^{-1}$ at 7 TeV



Comparison of Production Rates with $H_T(\text{jets+leptons})$

Kao, Cheng, Hou, and Sayre (2011)

Aguilar-Saavedra and Branco (2000)

	Low Luminosity (10 fb^{-1})		High Luminosity (100 fb^{-1})	
	Before Cuts	Standard Cuts	Before Cuts	Standard Cuts
Signal	200 (267)	46.7 (98.2)	1630 (2150)	394 (797)
$t\bar{t}$	5491 (7186)	20.2 (33.2)	44540 (58230)	174 (270)
$Wbbjj$	58 (77)	0.232 (0.3)	476 (644)	2.00 (2.2)

Comparison of Production Rates with $H_T(\text{jets})$

Kao, Cheng, Hou, and Sayre (2011)
Aguilar-Saavedra and Branco (2000)

	Low Luminosity (10 fb^{-1})		High Luminosity (100 fb^{-1})	
	Before Cuts	Standard Cuts	Before Cuts	Standard Cuts
Signal	200 (267)	30.4 (98.2)	1630 (2150)	251 (797)
$t\bar{t}$	5491 (7186)	10.1 (33.2)	44540 (58230)	83.9 (270)
$Wbbjj$	58 (77)	0.085 (0.3)	476 (644)	0.680 (2.2)

Conclusions

- ~ At the LHC, it is promising to detect FCNH top decays for $\lambda_{tc} > 0.02$ and $M_H < 140$ GeV.
- ~ For $M_H > 150$ GeV, most c-jets are removed by acceptance cuts.
- ~ Higher energy and higher luminosity can improve the discovery reach significantly.
- ~ We might find out if nature chooses the same mechanisms for electroweak symmetry breaking and FCNC.