Parity violation constraints on top A_{FB}

Sean Tulin (U Michigan)

With Moira Gresham, Ian-Woo Kim, Kathryn Zurek arXiv:1203.1320 [hep-ph]

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Overview

- . Top forward-backward asymmetry
 - Measured by CDF and D0 at Tevatron
 - 2–3σ disagreement with prediction from Standard Model
- 2. A hint for new physics in top sector?
- 3. Low-energy precision tests of parity violation
 - Strong constraints on many new physics models
 - Especially "collider-safe" models



A hint for new physics?

New physics models:

s-channel models

Sehgal, Wanninger (1988), Bagger, Schmidt, King (1988), Ferrario, Rodrigo (2009), Frampton, Shu, Wang (2009), Chivukula, Simmons, Yuan (2010), ...

• t-channel models

Jung, Murayama, Pierce, Wells (2009), Cheung, Keung, Yuan (2009), Shu, Tait, Wang (2010), ...

• Other scenarios



New top-like partners (stops) Isidori & Kemanik (2011)

Modification to top interactions from new physics at one-loop

Davoudiasl, McElmurry, Soni (2011), Gabrielli & Raidal (2011)

t-channel models: brief summary

New mediators with u-t or d-t flavor-violating coupling

Color singlet, weak doublet scalar

Shu et al (2009), Blum et al (2011), Nelson et al (2011), ...

- Color triplet scalar diquark Dorsner et al (2009), Shu et al (2009), Ligeti et al (2011),...
- Color sextet scalar diquark Shu et al (2009), Stone et al (2011),...
- Color singlet vector Z' Jung et al (2009),...
- Color singlet vector W'

Cheung et al (2009), Barger et al (2010), Shelton & Zurek (2011), Duffty et al (2012)

Massive spin-2 flavor-changing "graviton" Grinstein et al (2012)



Constraints on top A_{FB}

- Direct constraints: ttbar cross section, invariant mass distribution, top spin polarization, charge asymmetry
 - Light mediators (m_M < m_t) are more "collider-safe" (inv. mass distr.)
- Indirect constraints: precision tests of parity violation
 - Anomalous coupling of Z boson to first generation quarks at I -loop



New physics correction ~ $(\lambda^2/(4\pi)^2) (m_t^2/m_M^2)$

• Light mediators give $O(10^{-2})$ correction (complementary to colliders)

Weak nuclear charge

$$\mathscr{L}_{eq}^{PV} = \frac{G_F}{\sqrt{2}} \sum_{q=u,d} \left(C_{1q} \bar{e} \gamma^{\mu} \gamma_5 e \, \bar{q} \gamma_{\mu} q + C_{2q} \bar{e} \gamma^{\mu} e \, \bar{q} \gamma_{\mu} \gamma_5 q \right)$$

 $Q_W(Z,N) \equiv -2[(2Z+N)C_{1u} + (2N+Z)C_{1d}]$

 $\approx -N + (1 - \sin^2 \theta_W) Z$ (+ SM radiative corrections)



Marciano, Sirlin (1983), Erler, Ramsey-Musolf (2005)

Measurements of Q_W sensitive to:

- Running of $\sin^2 \theta_W$
- New physics for fixed $\sin^2 \theta_W$ •

extracted at Z-pole

Lepton, bottom, charm couplings to Z Inclusive hadronic couplings diluted by $1/N_F$ Weak nuclear charge

New physics contribution:

$$\begin{aligned} \mathcal{L}_{\text{eff}} &= -\frac{g_2}{c_W} Z^{\mu} \left(a_R^{\text{NP}}(q) \, \bar{q}_R \gamma_{\mu} q_R + a_L^{\text{NP}}(q) \, \bar{q}_L \gamma_{\mu} q_L \right) \\ & \swarrow & \downarrow \\ t & \downarrow \\ t & \downarrow \\ u, d & \downarrow \\$$

Parity violation constraints on top $A_{\rm FB}$

 Atomic parity violation measurements in cesium provide strongest constraints on Q_W



Parity-violating atomic transitions due to inference between Z (due to Q_W) and γ (due to external E field) Bouchiat and Bouchiat (1974)

Current bounds:

 $Q_W(\text{Cs}) = -73.20 \pm 0.35$ Wood et al. (1997), Guena et al. (2005) $Q_W^{\text{SM}}(\text{Cs}) = -73.15 \pm 0.02$ PDG

Sensitive to C_{Ig} at 0.3% level!

Parity violation constraints on top $A_{\rm FB}$

• Proton weak charge $Q_W(p) = -4 C_{Iu} - 2 C_{Id}$ $Q_W^{SM}(p) = 0.0713(8)$ SM value

 $Q_W(p) = 0.054(17)$ Current exp value Young et al (2007)

Polarized e-p elastic scattering data

Q_W(p) measurement at 4% level by Q-Weak experiment at JLab (in progress)
 Sensitive to C_{1a} at ~0.1% level!

Parity violation constraints on top $A_{\rm FB}$

Neutrino deep inelastic scattering

$$\mathscr{L}_{\nu q}^{PV} = -\frac{G_F}{\sqrt{2}} \sum_{q=u,d} \bar{\nu} \gamma^{\mu} (1-\gamma_5) \nu$$
$$\times \left(\epsilon_L(q) \, \bar{q} \gamma_{\mu} (1-\gamma_5) q + \epsilon_R(q) \, \bar{q} \gamma_{\mu} (1+\gamma_5) q \right)$$

New physics contribution: $\epsilon_{L,R}^{\rm NP}(q) = -a_{L,R}^{\rm NP}(q)$

 $\begin{array}{ll} \mbox{Measured values (PDG)} & \mbox{SM prediction} \\ g_L^2 \ \equiv \sum_q \epsilon_L^2(q) = 0.3025(14) & (g_L^2)_{\rm SM} = 0.30499(17) \\ g_R^2 \ \equiv \sum_q \epsilon_R^2(q) = 0.0309(10) & (g_R^2)_{\rm SM} = 0.03001(2) \end{array}$

Weak scalar doublet model

Interaction: $\lambda \left(\bar{u}_R V_{ib} u_L^i \phi^0 - \bar{u}_R b_L \phi^+ \right) + \text{h.c.}$ Provides good fit for A_{FB} for $m_{\phi 0} < 130 \text{ GeV}$

Blum, Hochberg, Nir (2011), Hochberg, Nir (2011)



Collider constraints: A_{FB} (parton level, NP only): $A_{FB}^{\text{high}} > 20\%, \quad A_{FB}^{\text{low}} < 20\%$ Total cross section: $\sigma(t\bar{t})_{\ell j}/\sigma(t\bar{t})_{\ell j}^{\rm SM}\big|_{\rm LO} = 1\pm 30\%$ $\sigma(t\bar{t})_{\ell\ell}/\sigma(t\bar{t})_{\ell\ell}^{\rm SM}\big|_{\rm LO} = 1\pm 30\%$ Parity violation at one-loop: $a_R^{\rm NP}(u) = \frac{\lambda^2 |V_{tb}|^2 m_t^2}{32\pi^2} \frac{|m_t^2|^2}{m_t^2} F(m_t^2/m_M^2)$

Ruled out by atomic parity violation at >4 σ

Horizontal SU(2)_X symmetry acting on $(u,t)_R$. Jung, Pierce, Wells (2011).

$$\mathscr{L} = \frac{g_X}{\sqrt{2}} V'_{\mu} \left[\bar{u}_R \gamma^{\mu} t_R + \varepsilon (\bar{u}_R \gamma^{\mu} u_R - t_R \gamma^{\mu} t_R) \right] + \text{h.c.} \qquad \text{Flavor-changing V' (like W boson)}$$

 $+ \frac{g_X}{2} Z'_{\mu} \left[\bar{t}_R \gamma^{\mu} t_R - \bar{u}_R \gamma^{\mu} u_R + 2\varepsilon (\bar{u}_R \gamma^{\mu} t_R + \bar{t}_R \gamma^{\mu} u_R) \right]$ Flavor-diagonal Z' (like Z boson)



To avoid large anomalous contribution to ttbar invariant mass distribution:

- Light $m_{V'} (m_{V'} < m_t)$
- Flavor-diagonal V'coupling ε
- $V' \rightarrow jet jet$, not $V' \rightarrow t + jet$

What about Z'?Intermediate mass Z' (130 GeV < $m_{Z'}$ < TeV) ruled out by dijet bounds</th>Heavy mass Z' is not pretty: requires O(100)-dimensional SU(2)_X-
breaking Higgs rep to have TeV-scale splitting with V'Low mass Z'? Hides under large QCD background

Parity violation at one-loop:



Massive vector theories are nonrenormalizable unless we embed them within a consistent gauge theory

Naïve estimate: regulate divergences with fixed cut-off Λ $a_R^{\rm NP}(u) = -\frac{\lambda^2}{16\pi^2} \frac{m_t^2}{m_{V'}^2} \left(F\left(\frac{m_t^2}{m_{V'}^2}\right) + \frac{1}{4} \log\left(\frac{\Lambda^2}{m_t^2}\right) \right)$ $+ \frac{N_C \lambda^2}{32\pi^2} \frac{m_t^2}{m_{Z'}^2} \log\left(\frac{\Lambda^2}{m_t^2}\right), \quad g_X \equiv \sqrt{2}\lambda$

Specify UV completion for $SU(2)_X$ model.

- 1) Introduce vector quark $t' \sim (3, 1, 2/3)$
- 2) Yukawa interactions $\mathscr{L} = y_1(\bar{u}_R, \bar{t}_R)t'_L S y_2 \bar{t}'_R(t_L, b_L)\epsilon H + h.c.$

to generate top mass $m_t = y_1 y_2 v_S v / m_{t'}$

- 3) Introduce SU(2)_X-breaking via doublet S and triplet Σ $\langle S \rangle = (0, v_S) \qquad \langle \Sigma \rangle = v_{\Sigma}(-2\varepsilon, 0, 1)/\sqrt{2}.$
- 4) Gauge boson masses

$$m_{V'}^2=g_X^2(v_S^2+v_\Sigma^2)/2$$
 and $m_{Z'}^2=g_X^2v_S^2/2$

Loop calculation:

$$a_{R}^{NP}(u) = -\frac{\lambda^{2}}{16\pi^{2}} \frac{m_{t}^{2}}{m_{V'}^{2}} F_{1}\left(\frac{m_{t}^{2}}{m_{V'}^{2}}, \frac{m_{t'}^{2}}{m_{V'}^{2}}\right) \qquad F_{1}(x,y) \equiv -\frac{1}{4}\left(2 + \frac{2 - x - y}{(1 - x)(1 - y)}\right) \qquad (14a)$$

$$+ \frac{(x^{2} - 2x + 4)\log x}{(1 - x)^{2}} + \frac{(2x^{2} - 8x)\log x}{(1 - x)(x - y)} + \frac{(y^{2} - 2y + 4)\log y}{(1 - y)^{2}} + \frac{(2y^{2} - 8y)\log y}{(1 - y)(y - x)}\right)$$

$$+ \frac{M_{C}\lambda^{2}}{32\pi^{2}} \frac{m_{t}^{2}}{m_{Z'}^{2}} F_{2}\left(\frac{m_{t}^{2}}{m_{t'}^{2}}\right), \qquad F_{2}(x) \equiv \frac{2(x - 1) - (1 + x)\log x}{1 - x}. \qquad (14b)$$



 $\Lambda = m_{t'} = 600 \text{ GeV}$ (Constraints stronger for higher mass)

Conclusions

- Possible new physics in top A_{FB} must satisfy constraints from colliders and low-energy precision tests
- t-channel models contribute to nuclear weak charge
 - Atomic parity violation (cesium)
 - Neutrino DIS
 - Proton weak charge
 - Qweak experiment: conclude data taking May 2012
- Simplified models:
 - Light color-singlet, weak doublet scalar disfavored
 - Light Z' [horizontal SU(2)_X] disfavored
 - Color triplet scalar diquark disfavored
 - Strong constaints on W' and color sextet scalar also

Backup slides



number of forward tops – backward tops $A_{FB} =$ number of forward tops + backward tops •0000

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CDF Run II Preliminary L = 8.7 fb⁻¹

Recent CDF result

Lepton $+ \geq 4$ jets (semileptonic)

Parton-level A_{FB} ("unfolded") – accounts for detector resolution & acceptances

t-tbar rest frame (top momenta reconstructed from 4 leading jets)



Parton Level	NLO (QCD+EW) $t\bar{t}$	$5.3 { m fb^{-1}}$	$8.7 { m ~fb^{-1}}$
$M_{t\bar{t}}$	$A_{\rm FB}$	$A_{\rm FB}$ (±[stat.+syst.])	$A_{\rm FB} \ (\pm [{\rm stat.+syst.}])$
$< 450 \mathrm{GeV/c^2}$	0.047	-0.116 ± 0.153	0.078 ± 0.054
$\geq 450 \mathrm{GeV/c^2}$	0.100	0.475 ± 0.112	0.296 ± 0.067

Consistent with	Subsample Data $A_{\rm FB}$ (%) MC@		MC@NLO
D0 5.4 fb ⁻ '	$m_{t\bar{t}} < 450 \text{GeV} m_{t\bar{t}} > 450 \text{GeV}$	7.8 ± 4.8 11.5 ± 6.0	1.3 ± 0.6 4.3 ± 1.3
		Data-level (not unfolded)	QCD
Confirmed in lepto	n charge asymmetr	CDF Run II	Preliminary L = 8.7 f
More forward-going e^+, μ^+ and backward-going e^-, μ^-		$A_{FB} = 0.066 \pm 0.025$ $A_{FB} = 0.016$	

	Data	NLO (QCD+EW) $t\bar{t}$
$M_{t\bar{t}}$	$A_{\rm FB}$ (± [stat.+syst.])	$A_{\rm FB}$
Inclusive	0.066 ± 0.025	0.016
$< 450 \text{GeV/c}^2$	0.037 ± 0.031	0.007
$\geq 450 \text{GeV}/\text{c}^2$	0.116 ± 0.042	0.032



Confirmed in dilepton channels



CDF 5.1 fb-1 (Conf. Note 10436)

Rapidity difference between top and antitop

Axigluon models: brief summary

Axigluon mass: 50 GeV – few TeV

Frampton et al (2009), Chivukula et al (2010), Bauer et al (2010), Cao et al (2010), Bai et al (2011), Wang et al (2011), Zerwakh et al (2011), Barcelo et al (2011), Alvarez et al (2011), Tavarez et al (2011), Gresham et al (2011), Aguilar-Saavedra et al (2011), Krnjaic (2011), ...

Phenomenological issues:

Dijet bounds, ttbar invariant mass distribution, flavor constraints, searches for additional light colored states, EW precision tests,

Still OK but not pretty



t-channel models

Phenomenological issues:

- Direct constraints: ttbar cross section, invariant mass distribution, top spin polarization, charge asymmetry
- Indirect constraints: precision tests of parity violation

Gresham, Kim, Zurek, ST (2012)

Two general classes:

- Light mediators (m_M < m_t)
 - Color-singlet, weak doublet scalar (additional Higgs)
 - Flavor-changing vector (Z',W')
- Heavy mediators (m_t < m_M)
 - Flavor-changing vector (Z',W')
 - Scalar diquarks

t-channel models: brief summary

What does the LHC have to say?

• Heavy meditors ($m_M > m_t$) contribute to $t\bar{t}$ production



On-shell mediator decays to t+jet

Affects total cross section, invariant mass distribution, number of associated jets

LHC results at ~ I fb⁻¹ see no hint for new physics in top

t-channel models for top $A_{\rm FB}$

• $t\bar{t}$ production cross section at LHC



ATLAS CONF 2011-108



t-channel models for top $A_{\rm FB}$

• Example: new vector V' with interaction $\lambda \bar{t}_R \gamma^{\mu} u_R V'_{\mu}$ + h.c.



t-channel models for top $A_{\rm FB}$

Light meditors ($m_M < m_t$) can evade LHC constraints! Assume M has (small) flavor-diagonal coupling to light quarks

 $M \rightarrow jet + jet$ dominates over $M \rightarrow t^* + jet$

No contribution to $t\bar{t}$ + jet production

