



PHENO 2012, Pittsburgh, May 7 2012



Low-energy imprints of extended gauge symmetries

Michal Malinský

AHEP group of IFIC, CSIC/University of Valencia

based on

Phys.Rev. D 84, 053012 (2011)
JHEP02 (2012) 084
Nucl.Phys. B 854 (2012) 28

in collaboration with

Renato Fonseca, Martin Hirsch, Valentina de Romeri, Werner Porod, Lazslo Reichert, Florian Staub

Extended gauge symmetries

EW scale

Extended gauge symmetries

Planck scale

EW scale

Extended gauge symmetries

Planck scale

His sunt dragones



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Extended gauge symmetries

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His sunt dragones



EW scale



Extended gauge symmetries

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Low-scale extended
gauge symmetries?

EW scale

See talks by

M.Bishara
D.Duffy
R.Ruiz

and others

Extended gauge symmetries

Planck scale



low-scale extended
gauge symmetries?

SUSY @ TeV scale

EW scale

See talks by
almost anyone in Higgs & SUSY
sessions of PHENO'12

e.g.

R.Huo
and many others

Extended gauge symmetries



Planck scale

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A. Elsayed, S. Khalil and S. Moretti
arXiv:1106.2130 [hep-ph]

M.Hirsch, MM, W.Porod, L.Reichert
JHEP02 (2012) 084

Extended gauge symmetries

Planck scale

high-scale extended
gauge symmetries?



EW scale

Extended gauge symmetries

unified dynamics ?

Planck scale

GUT scale ?



EW scale

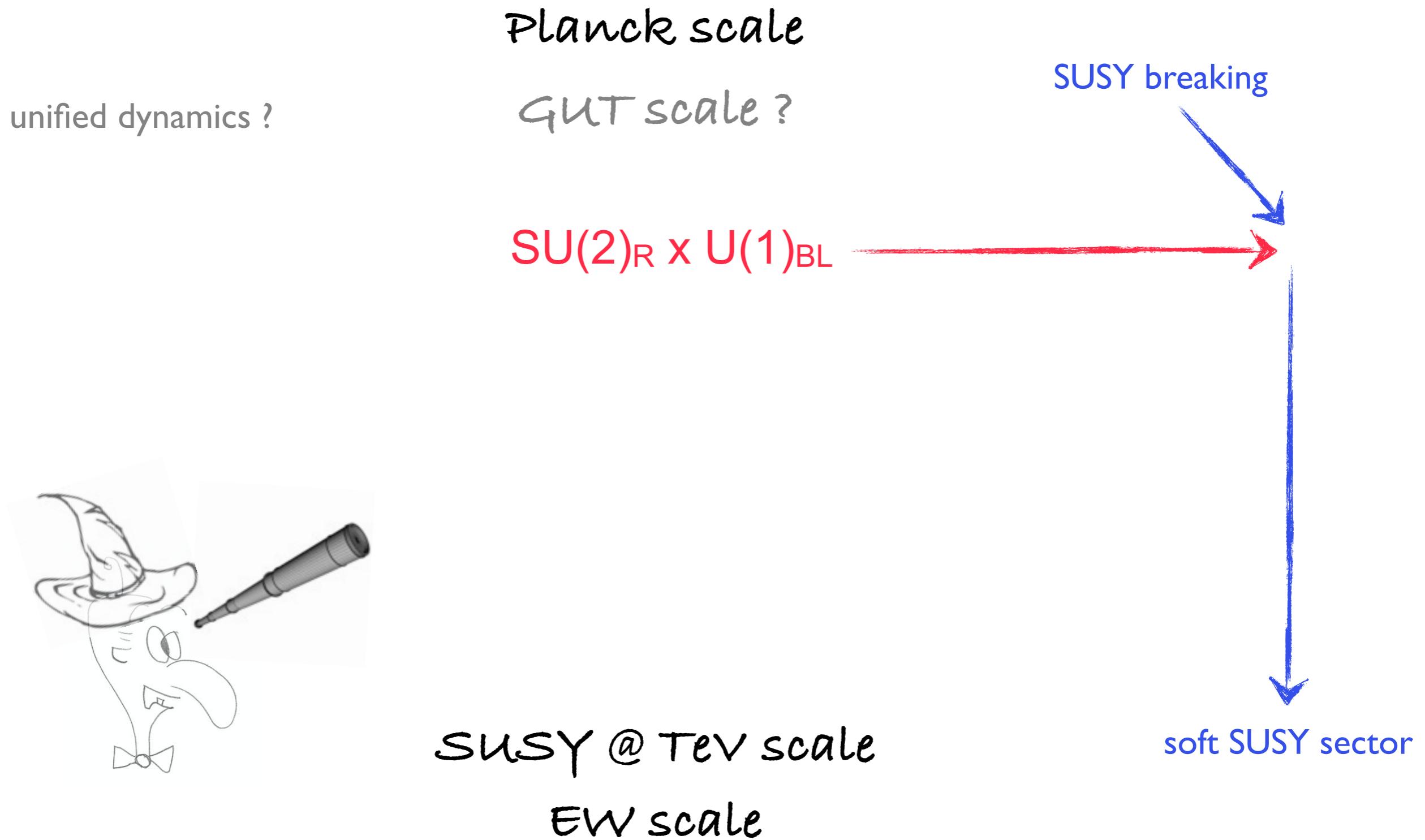
Extended gauge symmetries



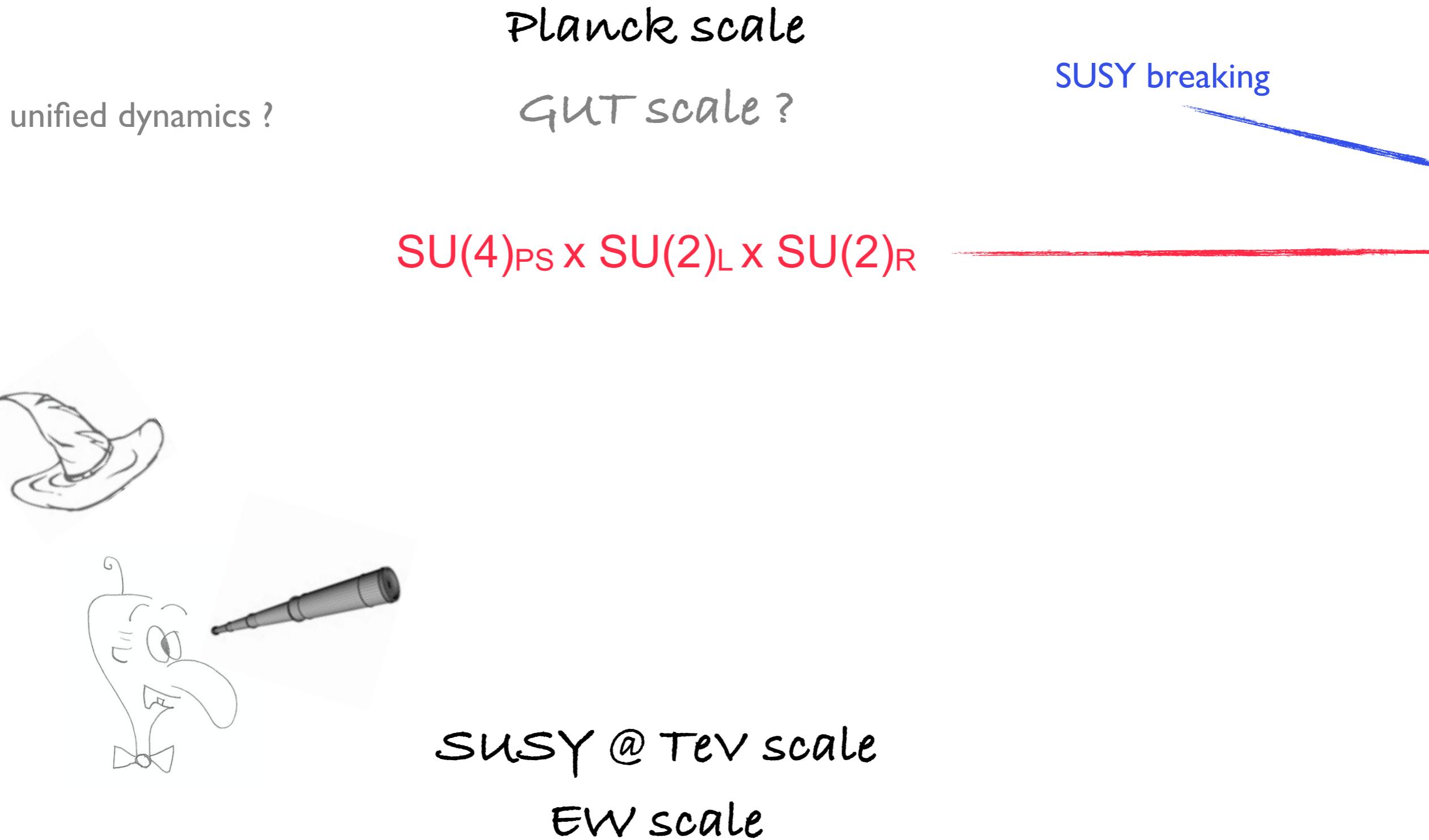
Extended gauge symmetries



Extended gauge symmetries



Extended gauge symmetries



Leading-log mSUGRA RG invariants

$$m_{\tilde{f}}^2(M_{\text{SUSY}}) = \textcolor{red}{m}_0^2 + \frac{\textcolor{red}{M}_{1/2}}{\alpha(M_G)^2} \sum_S \sum_{i=1}^{N_S} \textcolor{blue}{F}_{\tilde{f},S}^i \alpha_i (v_S^<)^2$$

$$M_i(M_{\text{SUSY}}) = \textcolor{red}{M}_{1/2} \frac{\alpha_i(M_{\text{SUSY}})}{\alpha(M_G)}$$

$$\textcolor{blue}{F}_{\tilde{f},S}^i = \frac{c_{\tilde{f},S}^i}{b_i} \left\{ 1 - \left[\alpha_i(v_S^<)/\alpha_i(v_S^>) \right]^2 \right\}$$

$$c_{\tilde{f},S}^i = 2C_{G_i^S}(R_f)$$

Leading-log mSUGRA RG invariants

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$$(m_{\tilde{A}}^2 - m_{\tilde{B}}^2)/M_i^2$$

independent on mSUGRA parameters at the leading log level

S. P. Martin and P. Ramond, Phys. Rev. D 48, 5365 (1993).

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G.A. Blair, W. Porod, and P.M. Zerwas,
 Phys. Rev. D 63, 017703 (2000)

P. Bechtle, K. Desch, W. Porod, and P. Wienemann,
 Eur. Phys. J. C 46, 533 (2006)

R. Lafaye, T. Plehn, M. Rauch, and D. Zerwas,
 Eur. Phys. J. C 54, 617 (2008)

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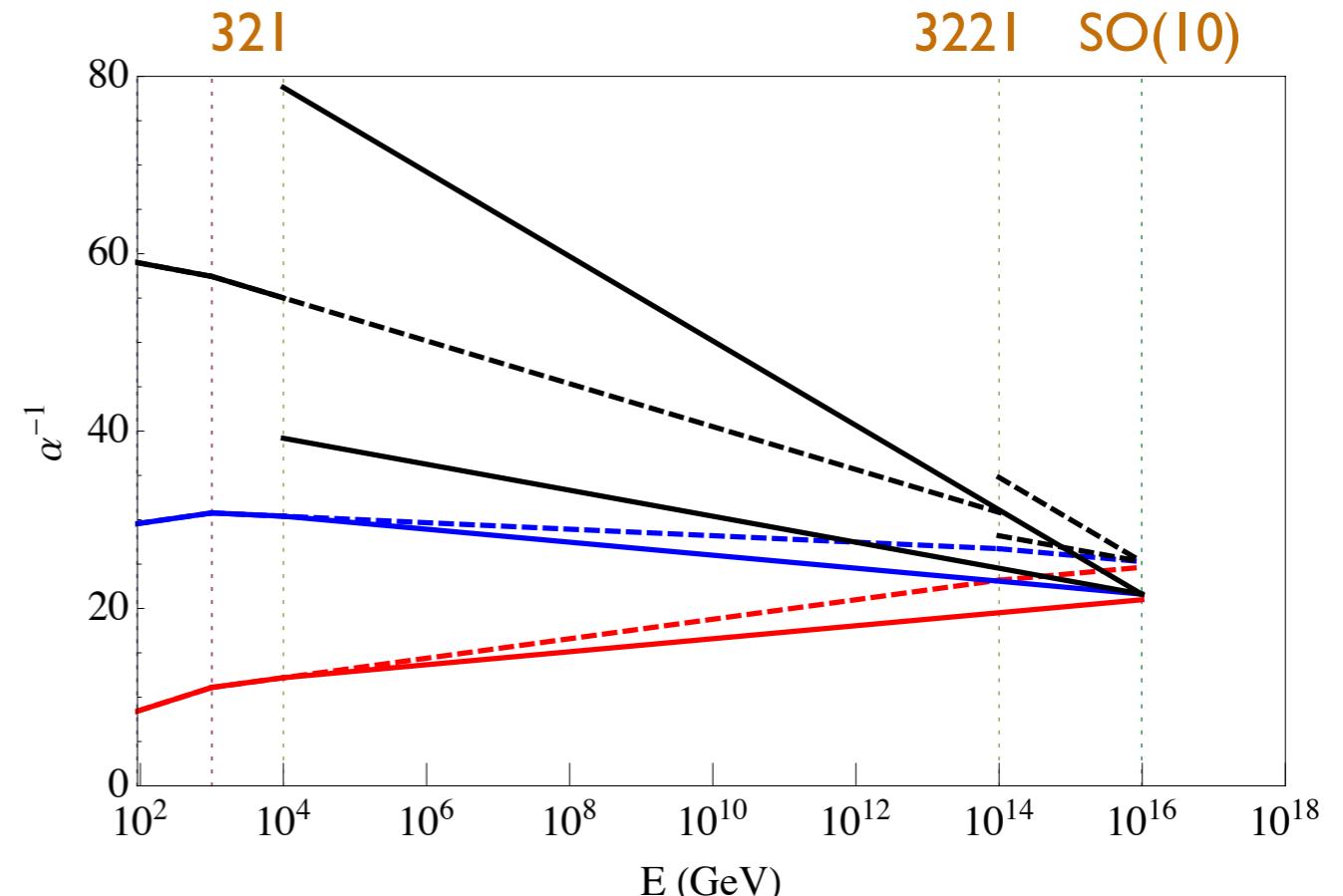
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Example: sliding left-right scale

extra vector-like down-type quarks

Matter Higgs

#	$3_c 2_L 2_R 1_{B-L}$	$SO(10)$
3	$(3, 2, 1, +\frac{1}{3})$	16
3	$(\bar{3}, 1, 2, -\frac{1}{3})$	16
3	$(1, 2, 1, -1)$	16
3	$(1, 1, 2, +1)$	16
3	$(1, 1, 1, 0)$	1
1	$(3, 1, 1, -\frac{2}{3}), (\bar{3}, 1, 1, +\frac{2}{3})$	10
1	$(1, 2, 2, 0)$	10, 120
1	$(1, 2, 1, \pm 1)$	$\overline{16}, 16$
3	$(1, 1, 2, \mp 1)$	$\overline{16}, 16$



V. DeRomeri, M. Hirsch, MM, L. Reichert
 Phys. Rev. D 84, 053012 (2011)

a variant of

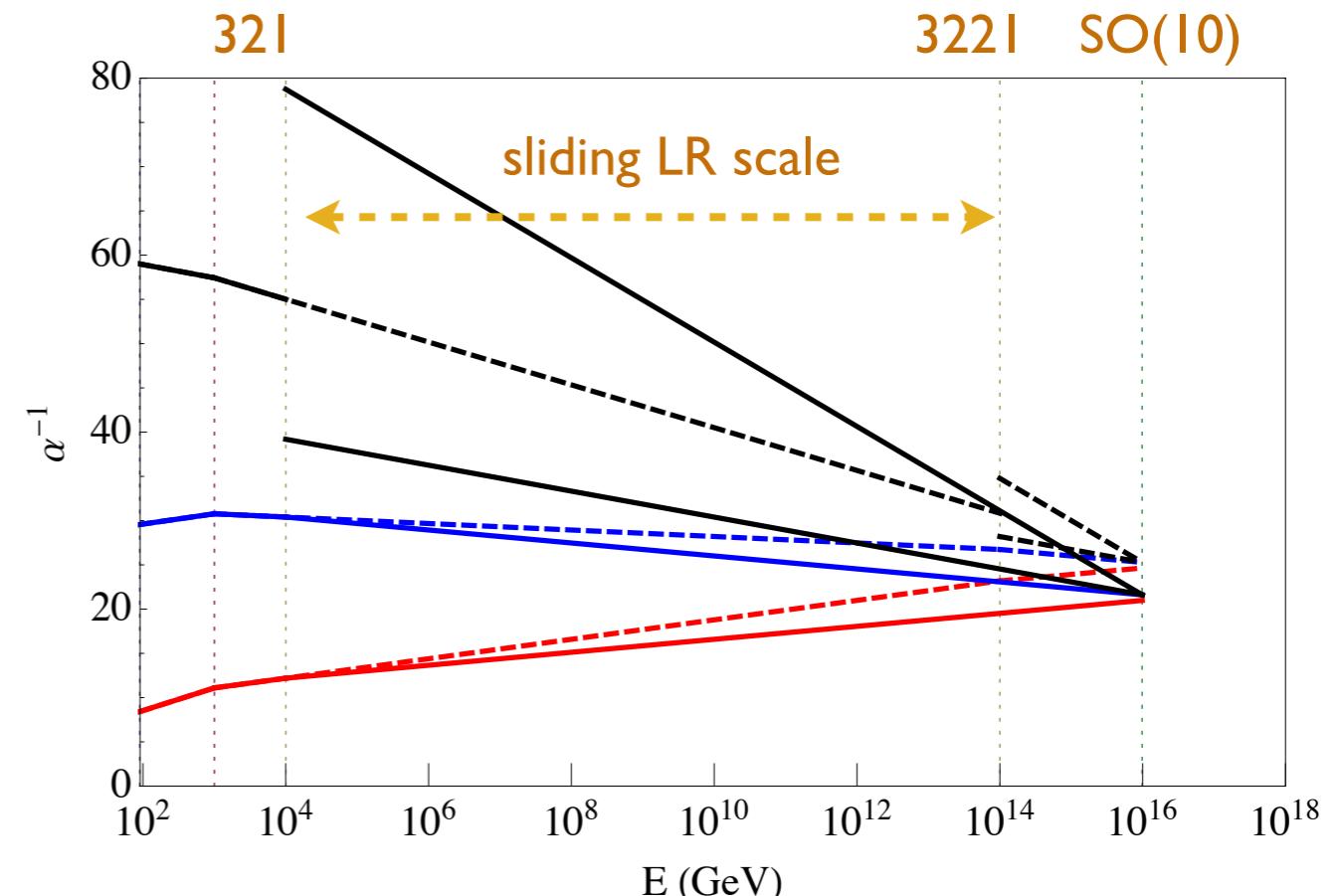
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Matter

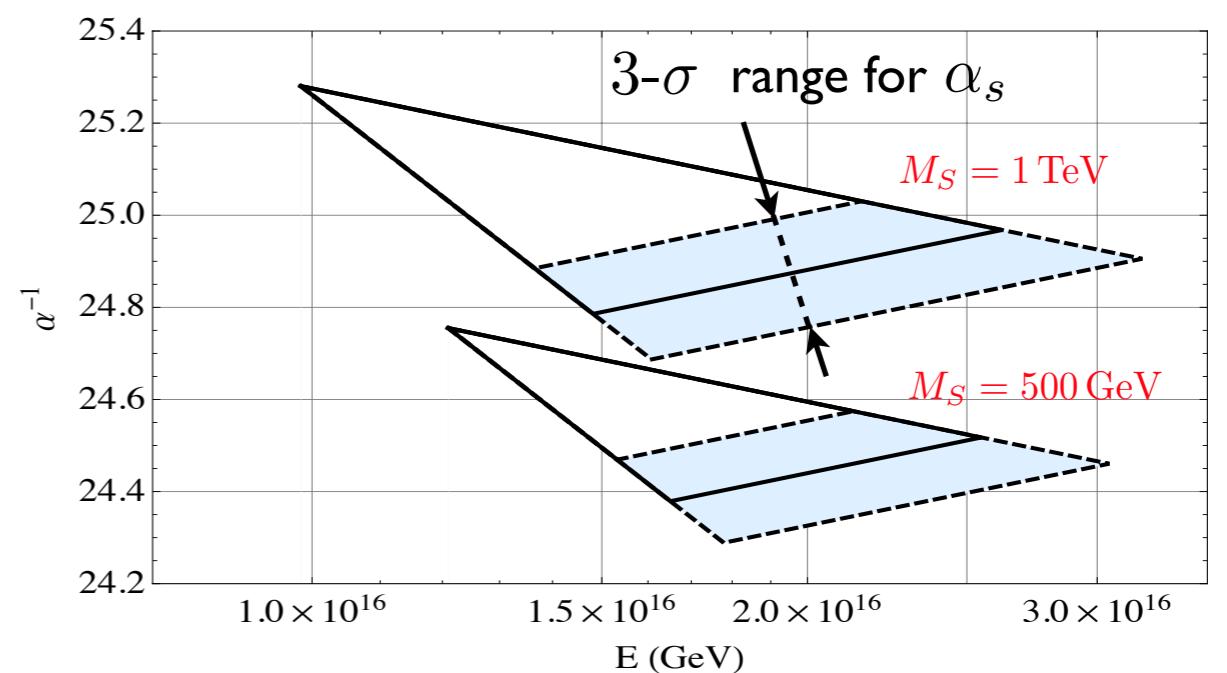
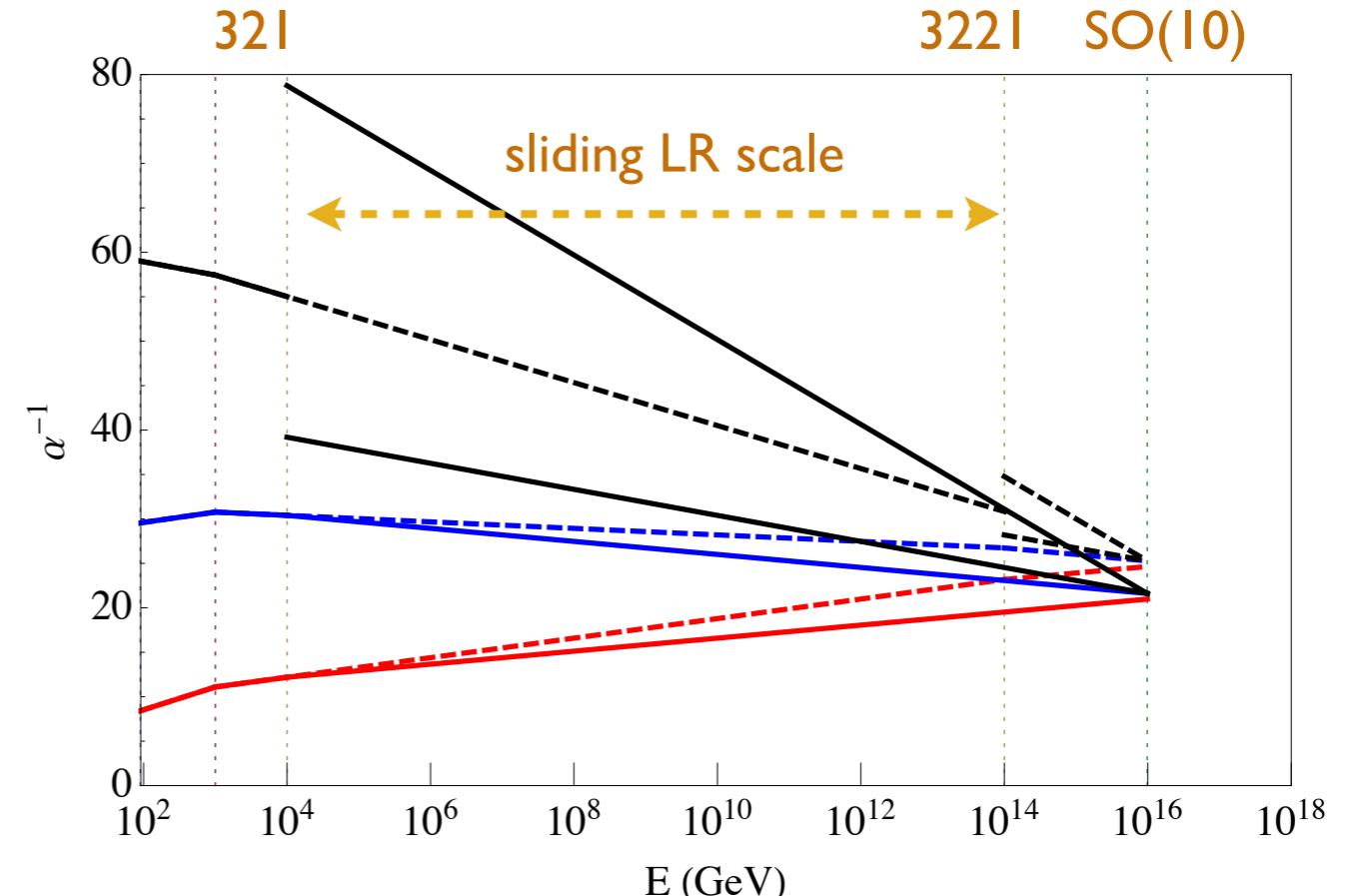
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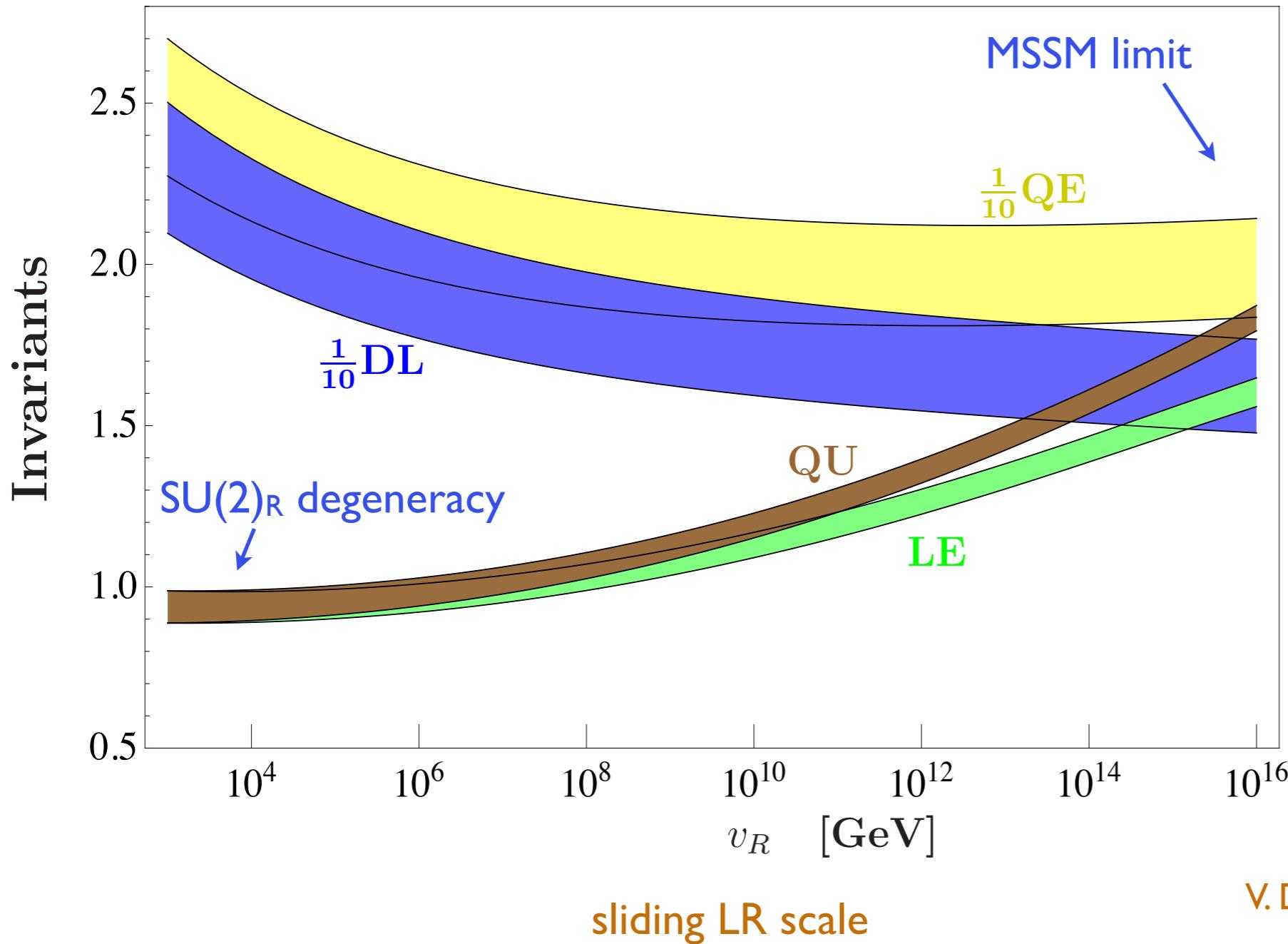
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Example: sliding left-right scale



$$LE \equiv (m_{\tilde{L}}^2 - m_{\tilde{E}}^2)/M_1^2$$

$$QE \equiv (m_{\tilde{Q}}^2 - m_{\tilde{E}}^2)/M_1^2$$

$$DL \equiv (m_{\tilde{D}}^2 - m_{\tilde{L}}^2)/M_1^2$$

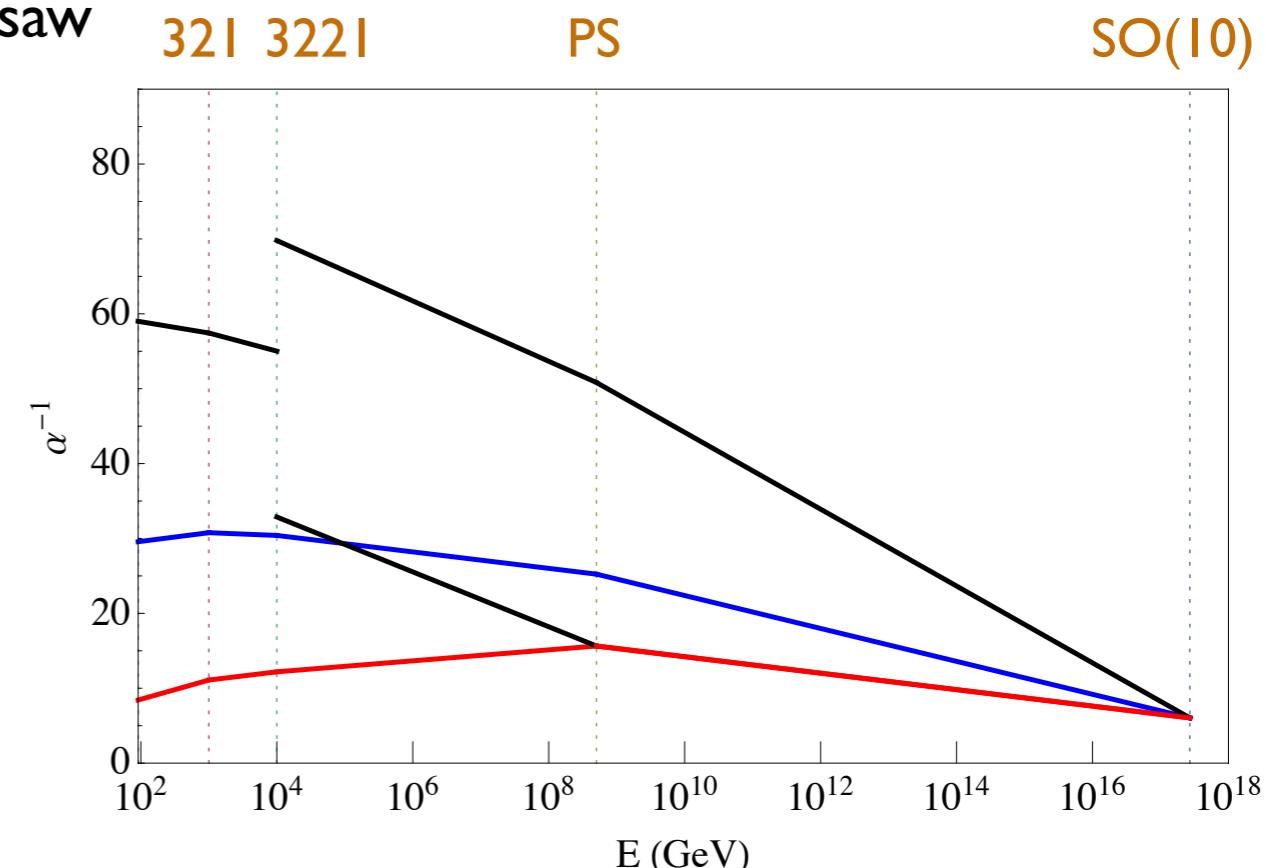
$$QU \equiv (m_{\tilde{Q}}^2 - m_{\tilde{U}}^2)/M_1^2$$

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Example: sliding Pati-Salam & LR scales

extra vector-like down-type quarks, “exotic” seesaw

#	$3_c 2_L 2_R 1_{B-L}$	$4_C 2_L 2_R$	$SO(10)$
Matter	$(3, 2, 1, +\frac{1}{3})$	$(4, 2, 1)$	16
	$(\bar{3}, 1, 2, -\frac{1}{3})$	$(\bar{4}, 1, 2)$	16
	$(1, 2, 1, -1)$	$(4, 2, 1)$	16
	$(1, 1, 2, +1)$	$(\bar{4}, 1, 2)$	16
	$(1, 1, 3, 0)$	$(1, 1, 3)$	45
	$(3, 1, 1, \mp\frac{2}{3})$	$(6, 1, 1)$	10
Higgs	$(1, 2, 2, 0)$	$(1, 2, 2)$	10
	$(1, 1, 3, 0)$	$(1, 1, 3)$	45
	$(1, 2, 1, \pm 1)$	$(\bar{4}, 2, 1), (4, 2, 1)$	$\overline{16}, 16$
	$(1, 1, 2, \mp 1)$	$(4, 1, 2), (\bar{4}, 1, 2)$	$\overline{16}, 16$
	absent	$(15, 1, 1)$	45



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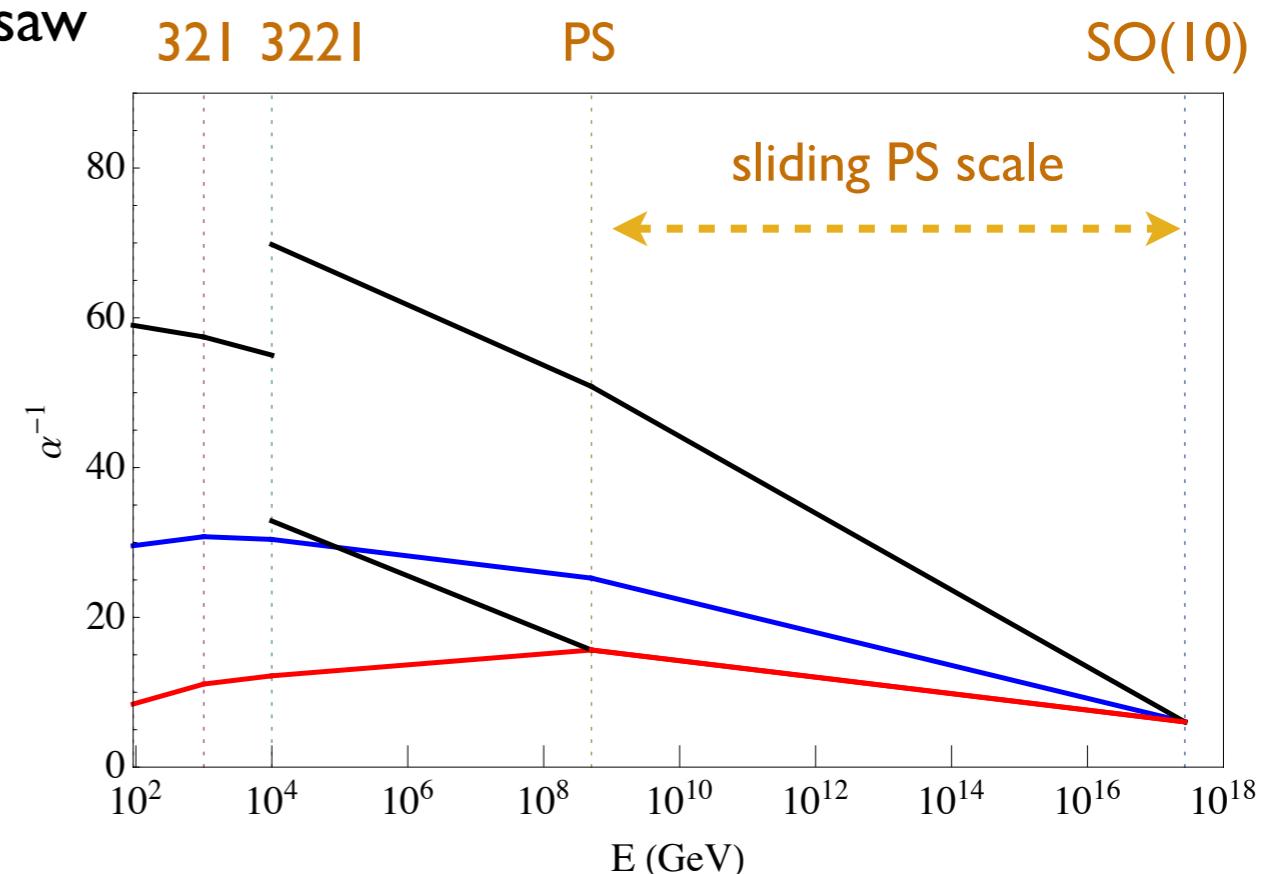
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Matter

Higgs

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3	$(3, 2, 1, +\frac{1}{3})$	$(4, 2, 1)$	16
3	$(\bar{3}, 1, 2, -\frac{1}{3})$	$(\bar{4}, 1, 2)$	16
3	$(1, 2, 1, -1)$	$(4, 2, 1)$	16
3	$(1, 1, 2, +1)$	$(\bar{4}, 1, 2)$	16
3	$(1, 1, 3, 0)$	$(1, 1, 3)$	45
1	$(3, 1, 1, \mp\frac{2}{3})$	$(6, 1, 1)$	10
2	$(1, 2, 2, 0)$	$(1, 2, 2)$	10
1	$(1, 1, 3, 0)$	$(1, 1, 3)$	45
1	$(1, 2, 1, \pm 1)$	$(\bar{4}, 2, 1), (4, 2, 1)$	$\overline{16}, 16$
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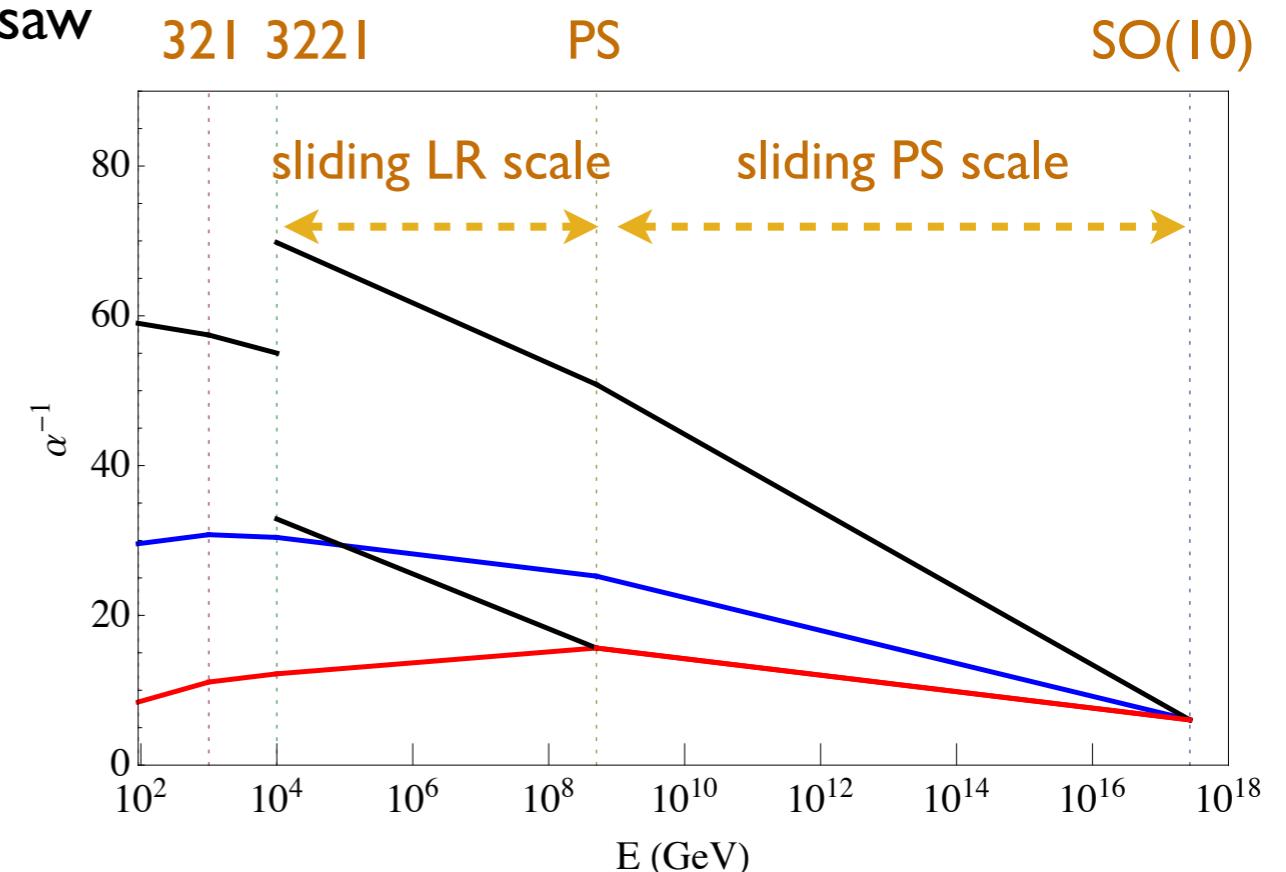


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3	(3, 2, 1, $+\frac{1}{3}$)	(4, 2, 1)	16
	($\bar{3}$, 1, 2, $-\frac{1}{3}$)	($\bar{4}$, 1, 2)	16
	(1, 2, 1, -1)	(4, 2, 1)	16
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Higgs	(1, 2, 2, 0)	(1, 2, 2)	10
	(1, 1, 3, 0)	(1, 1, 3)	45
	(1, 2, 1, ± 1)	($\bar{4}$, 2, 1), (4, 2, 1)	$\overline{16}$, 16
	(1, 1, 2, ∓ 1)	(4, 1, 2), ($\bar{4}$, 1, 2)	$\overline{16}$, 16
	absent	(15, 1, 1)	45



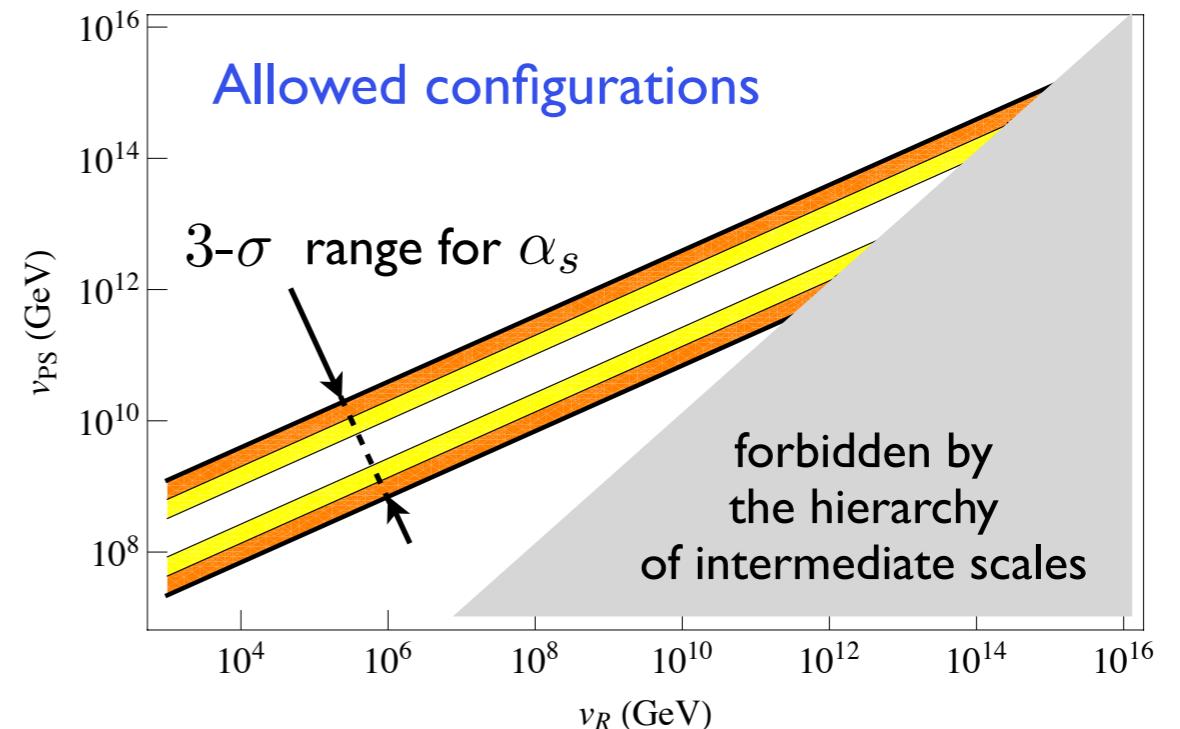
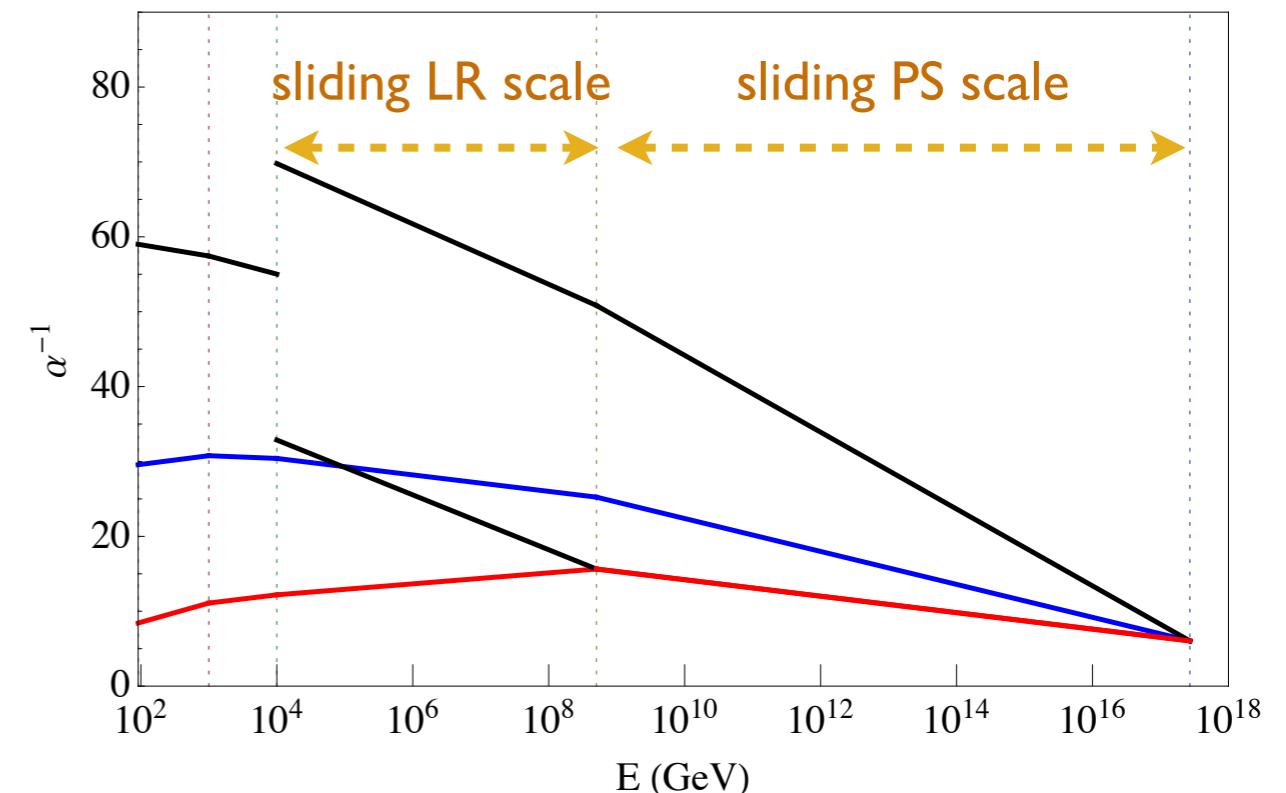
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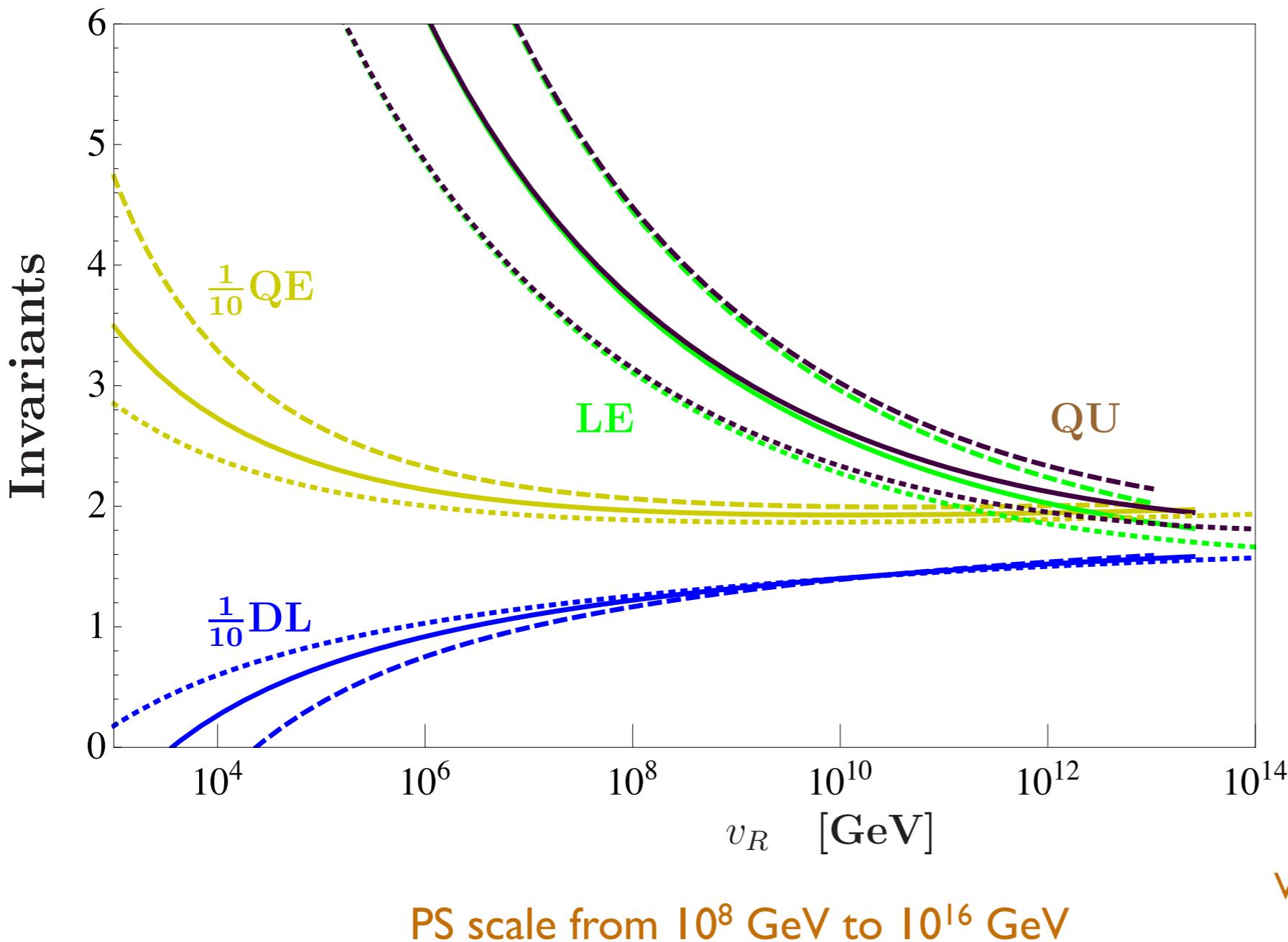
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1	(1, 1, 3, 0)	(1, 1, 3)	45
1	(1, 2, 1, ± 1)	($\bar{4}$, 2, 1), (4, 2, 1)	$\overline{16}$, 16
1	(1, 1, 2, ∓ 1)	(4, 1, 2), ($\bar{4}$, 1, 2)	$\overline{16}$, 16
1	absent	(15, 1, 1)	45

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V. DeRomeri, M.Hirsch, MM, L.Reichert
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Example: sliding Pati-Salam & LR scales



$$LE \equiv (m_{\tilde{L}}^2 - m_{\tilde{E}}^2)/M_1^2$$

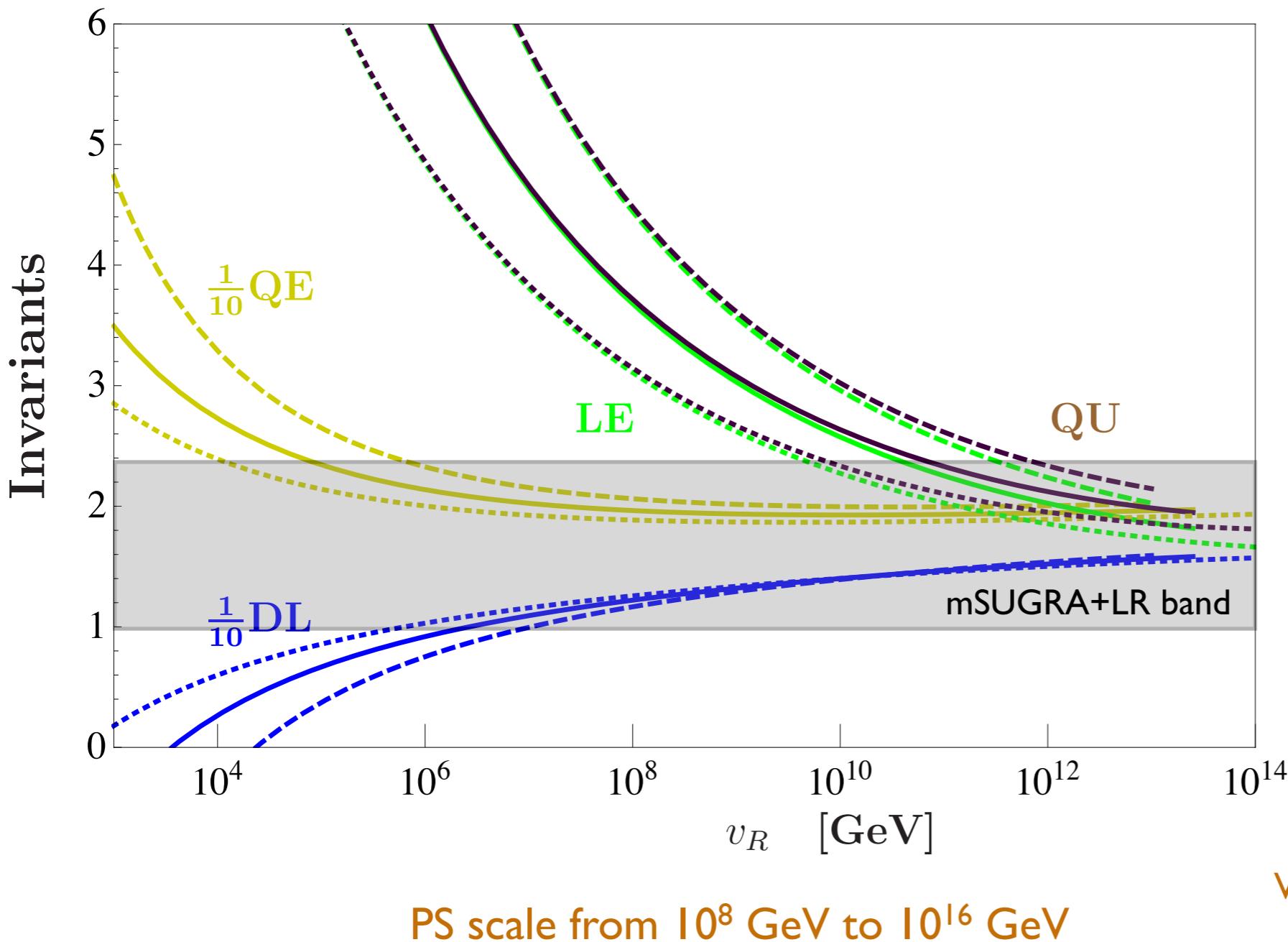
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Anti-example: sliding $U(1)_R \times U(1)_{BL}$

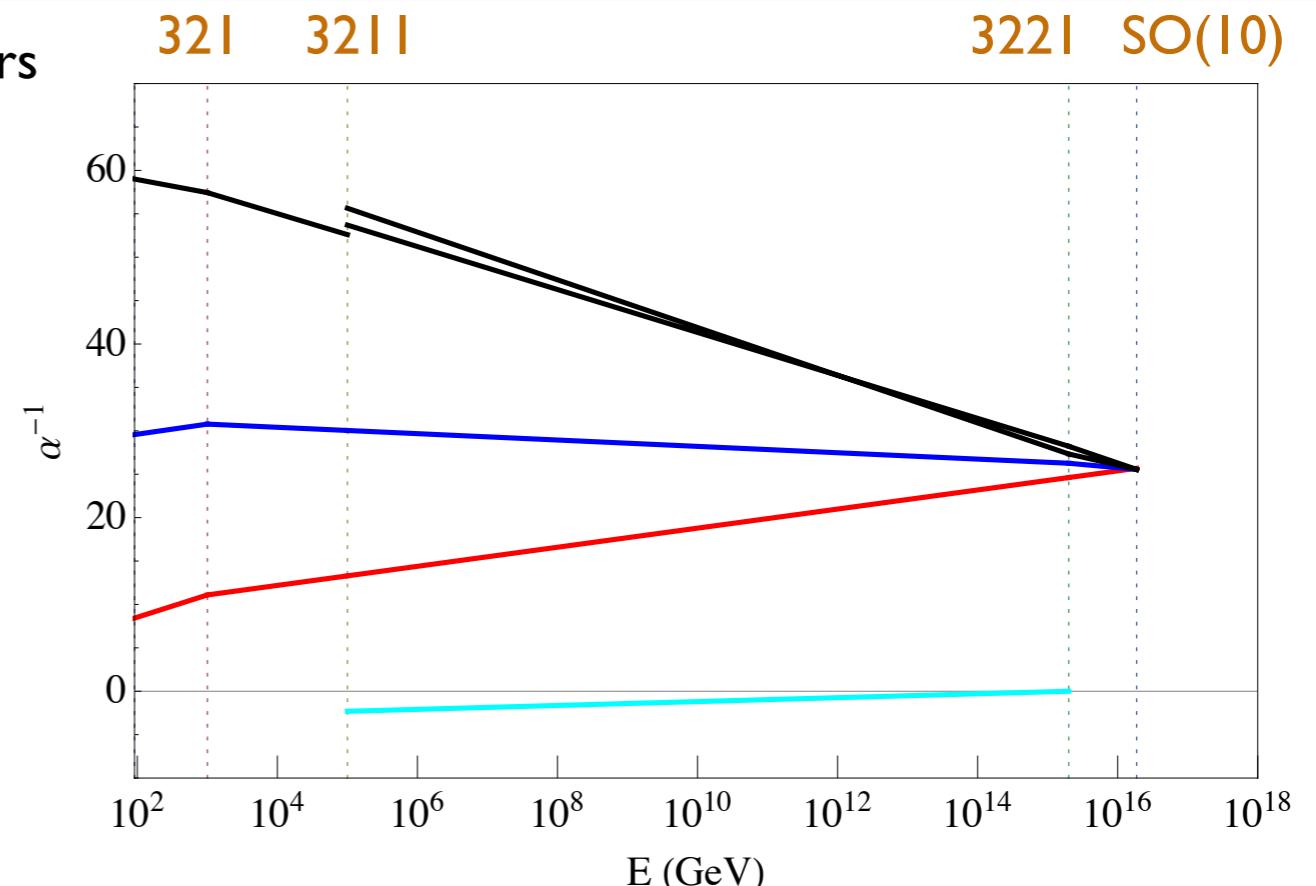
minimally finetuned setting with a pair of Abelian factors

#	$3_c 2_L 1_R 1_{B-L}$	$3_c 2_L 2_R 1_{B-L}$	$SO(10)$
3	$(3, 2, 0, +\frac{1}{3})$	$(3, 2, 1, +\frac{1}{3})$	16
3	$(\bar{3}, 1, \pm\frac{1}{2}, -\frac{1}{3})$	$(\bar{3}, 1, 2, -\frac{1}{3})$	16
3	$(1, 2, 0, -1)$	$(1, 2, 1, -1)$	16
3	$(1, 1, \pm\frac{1}{2}, +1)$	$(1, 1, 2, +1)$	16
3	$(1, 1, 0, 0)$	$(1, 1, 1, 0)$	1
2	$(1, 2, \pm\frac{1}{2}, 0)$	$(1, 2, 2, 0)$	10
1	absent	$(1, 1, 3, 0)$	45
1	absent	$(1, 2, 1, \pm 1)$	$\overline{16}, 16$
1	$(1, 1, \pm\frac{1}{2}, \mp 1)$	$(1, 1, 2, \mp 1)$	$\overline{16}, 16$

Matter

Higgs

MM, J. C. Romao, and J.W. F. Valle
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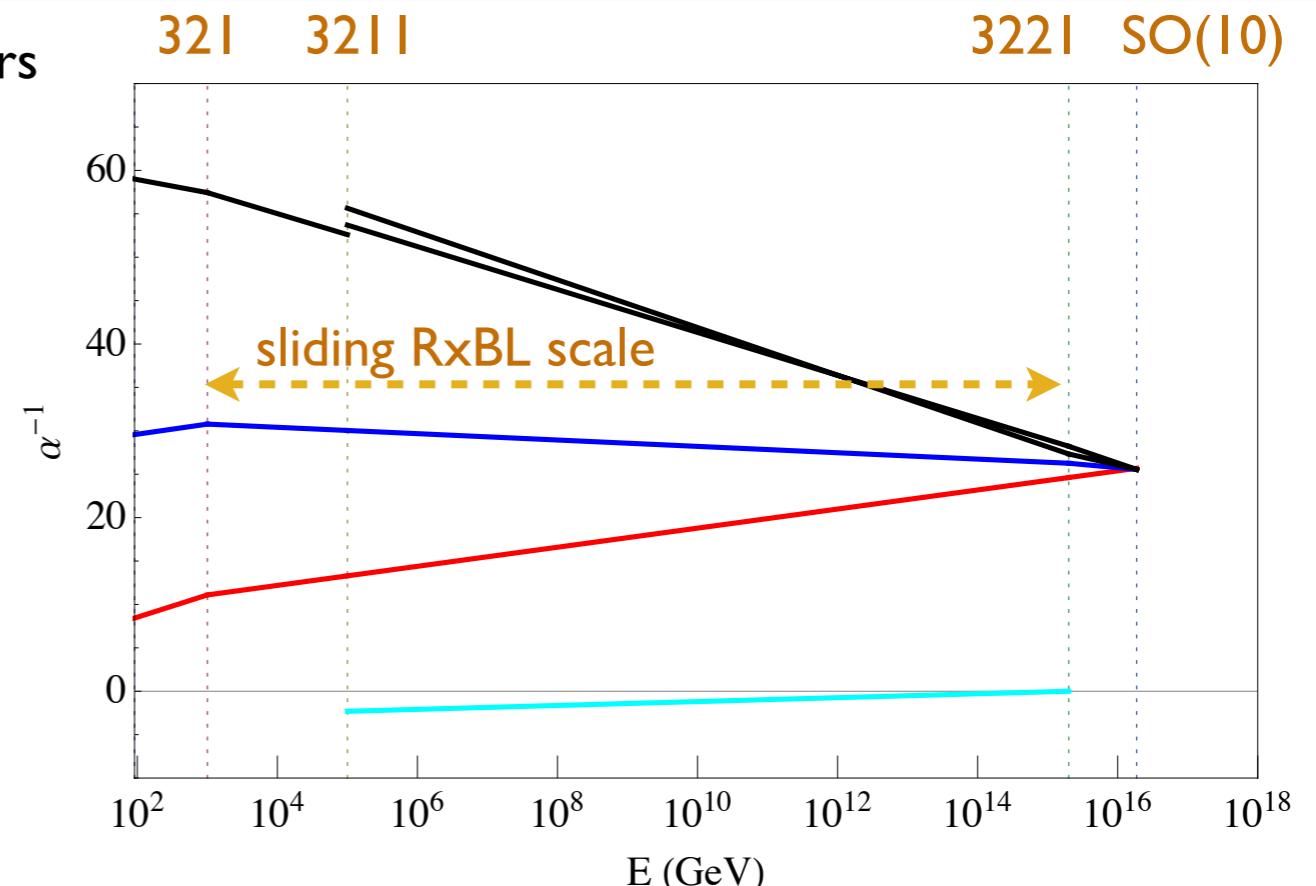
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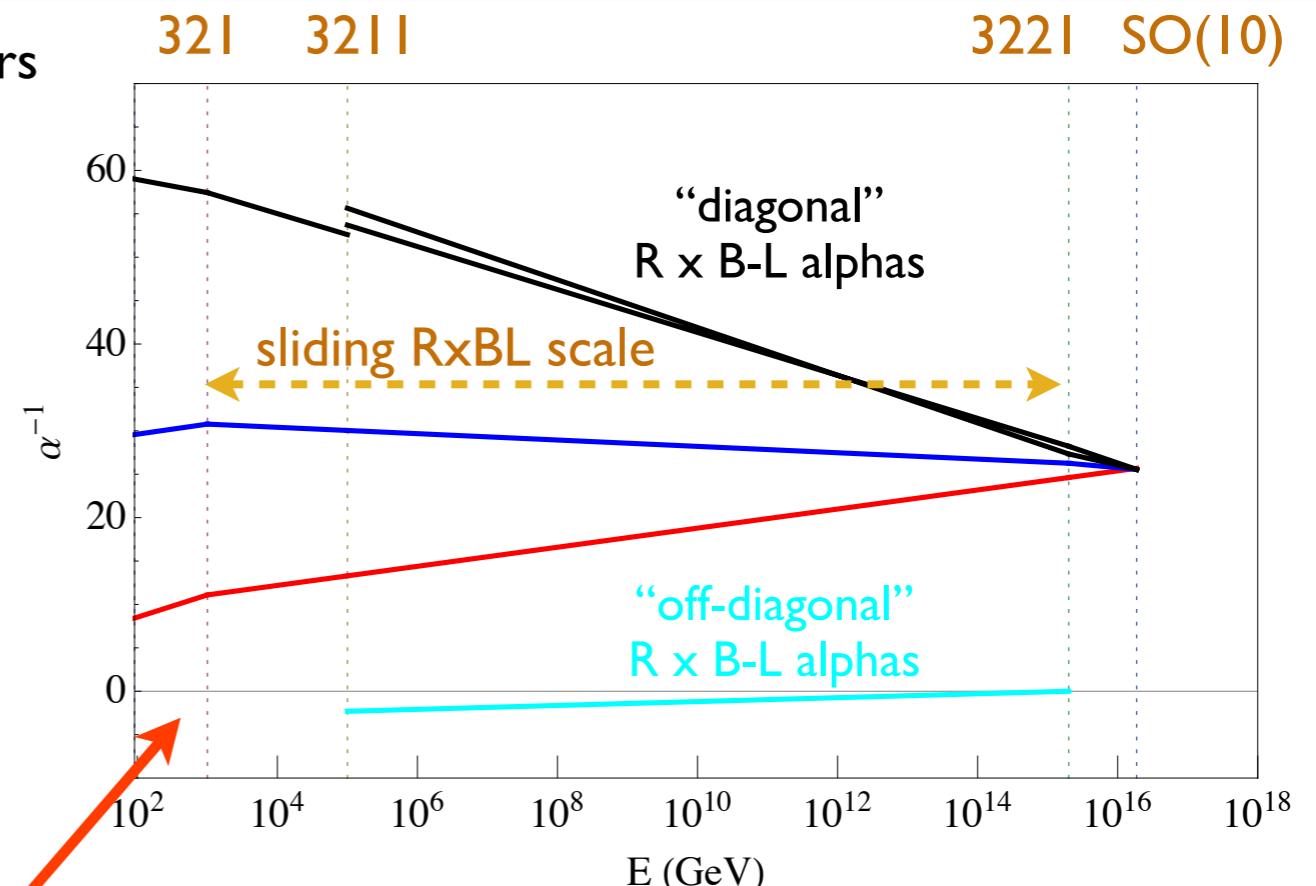
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B. Holdom, F. del Aguila, G.D. Coughlan, M. Quiros
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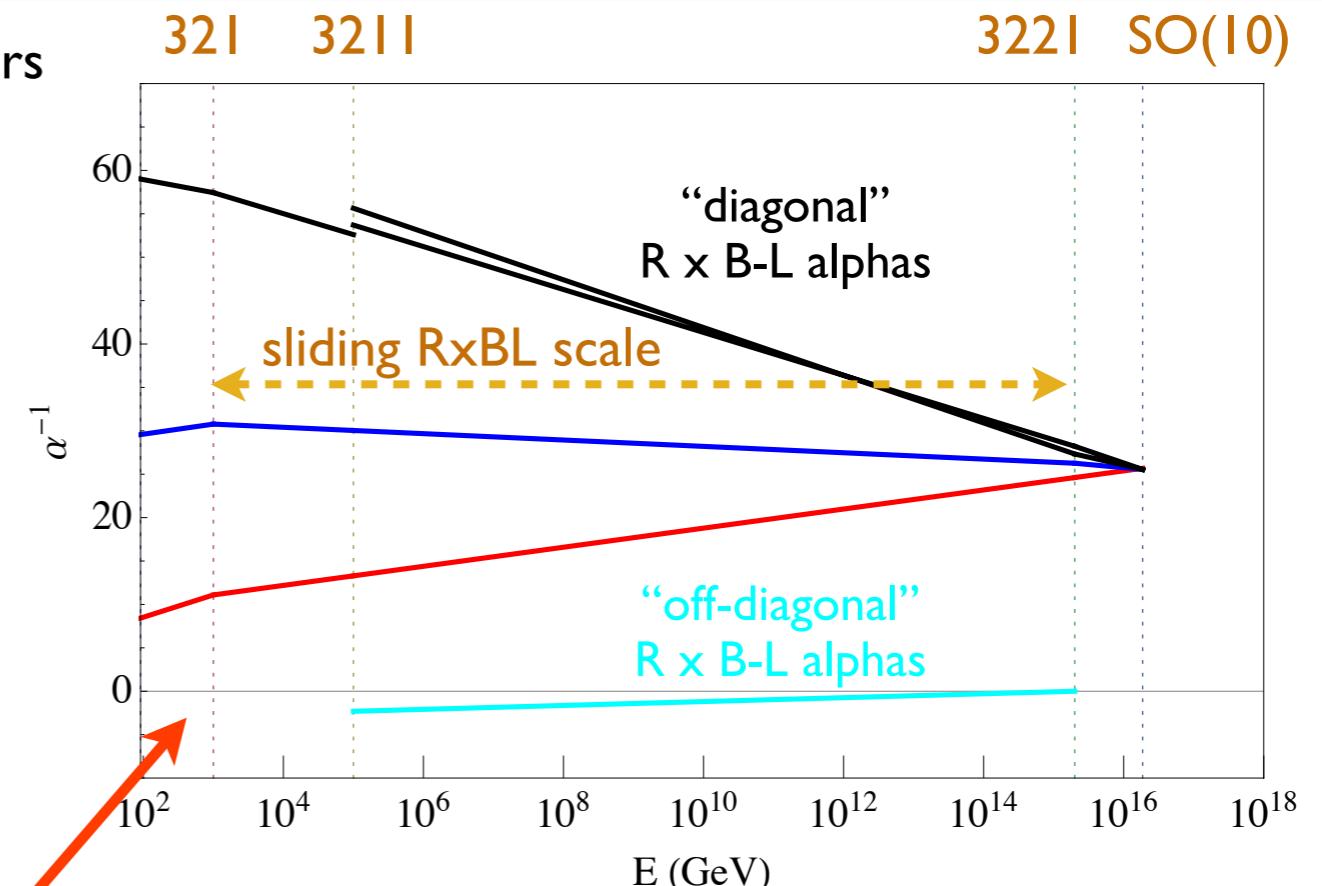
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#	$3_c 2_L 1_R 1_{B-L}$	$3_c 2_L 2_R 1_{B-L}$	$SO(10)$
3	$(3, 2, 0, +\frac{1}{3})$	$(3, 2, 1, +\frac{1}{3})$	16
3	$(\bar{3}, 1, \pm\frac{1}{2}, -\frac{1}{3})$	$(\bar{3}, 1, 2, -\frac{1}{3})$	16
3	$(1, 2, 0, -1)$	$(1, 2, 1, -1)$	16
3	$(1, 1, \pm\frac{1}{2}, +1)$	$(1, 1, 2, +1)$	16
3	$(1, 1, 0, 0)$	$(1, 1, 1, 0)$	1
2	$(1, 2, \pm\frac{1}{2}, 0)$	$(1, 2, 2, 0)$	10
1	absent	$(1, 1, 3, 0)$	45
1	absent	$(1, 2, 1, \pm 1)$	$\overline{16}, 16$
1	$(1, 1, \pm\frac{1}{2}, \mp 1)$	$(1, 1, 2, \mp 1)$	$\overline{16}, 16$

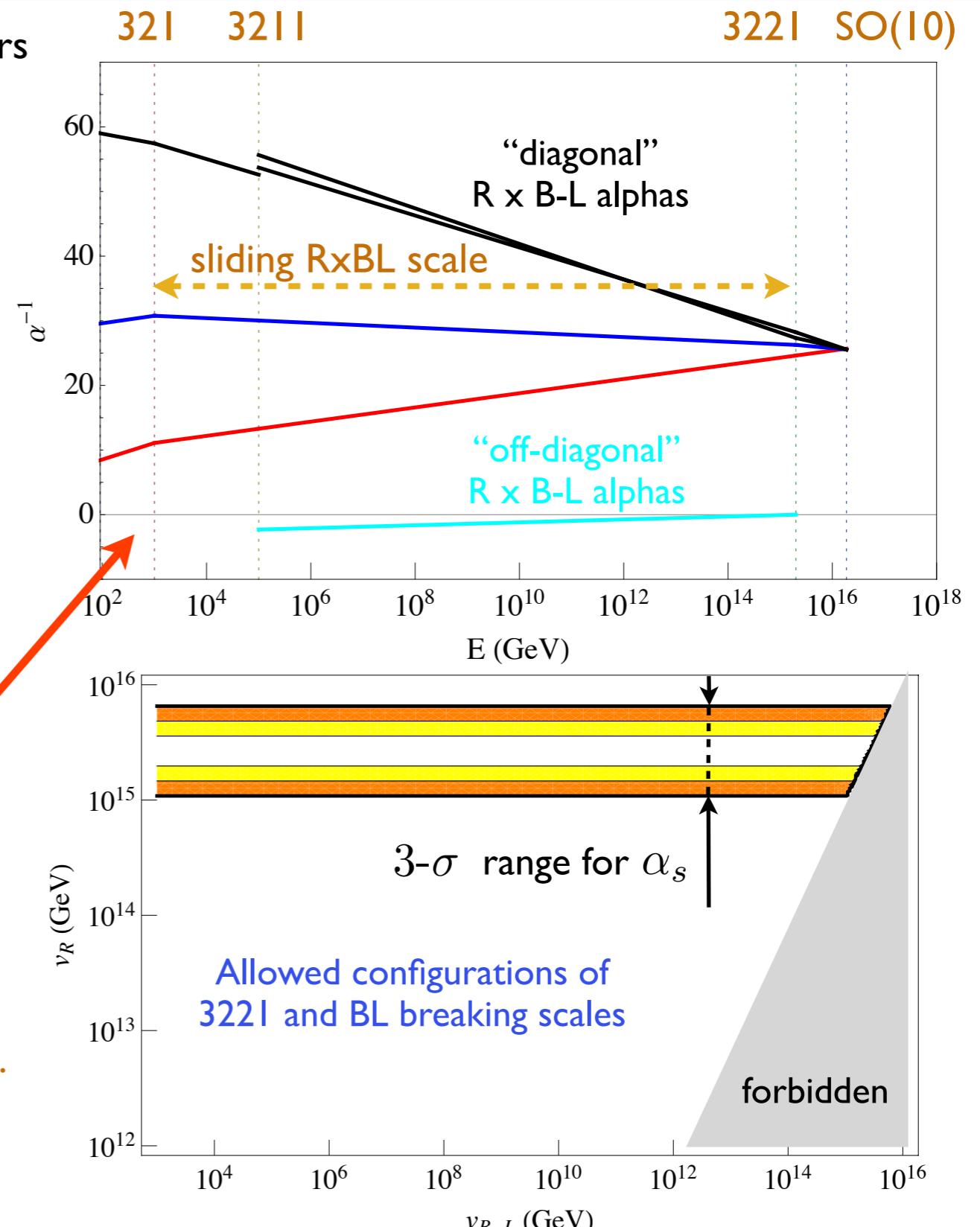
MM, J. C. Romao, and J.W. F. Valle
 Phys. Rev. Lett. 95, 161801 (2005)

The $U(1)$ mixing is important here!

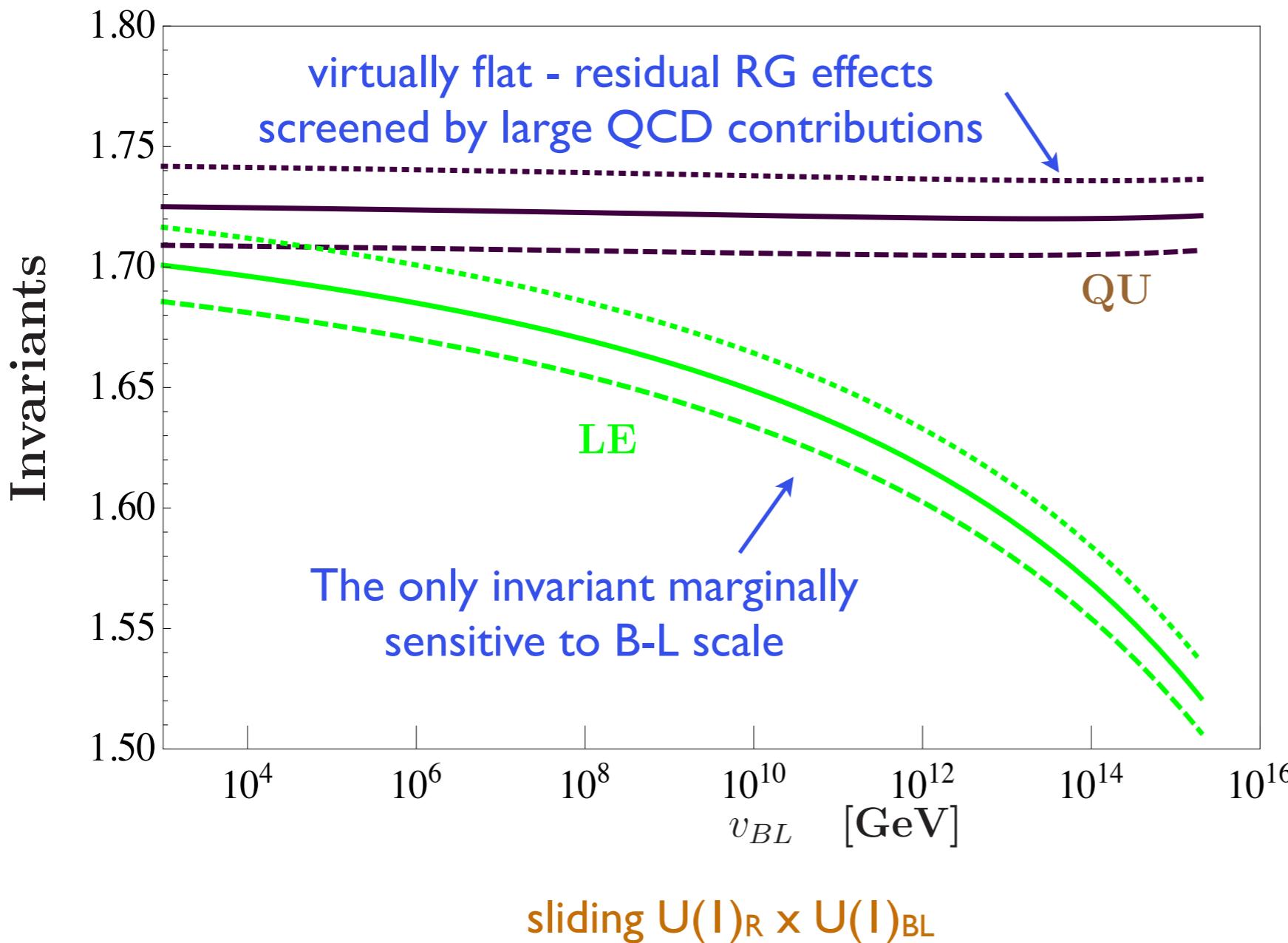
B. Holdom, F. del Aguila, G.D. Coughlan, M. Quiros
 et al. in 1980's

M.E. Machacek, M.T. Vaughn, Nucl. Phys. B 222 (1983) etc.

R. Fonseca, MM, W. Porod, and F. Staub
 Nucl.Phys. B 854 (2012) 28

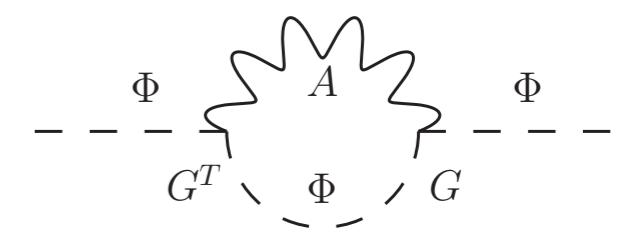
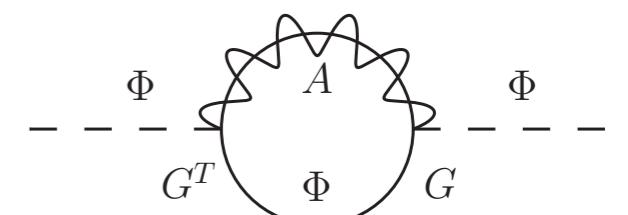


Anti-example: sliding $U(1)_R \times U(1)_{BL}$

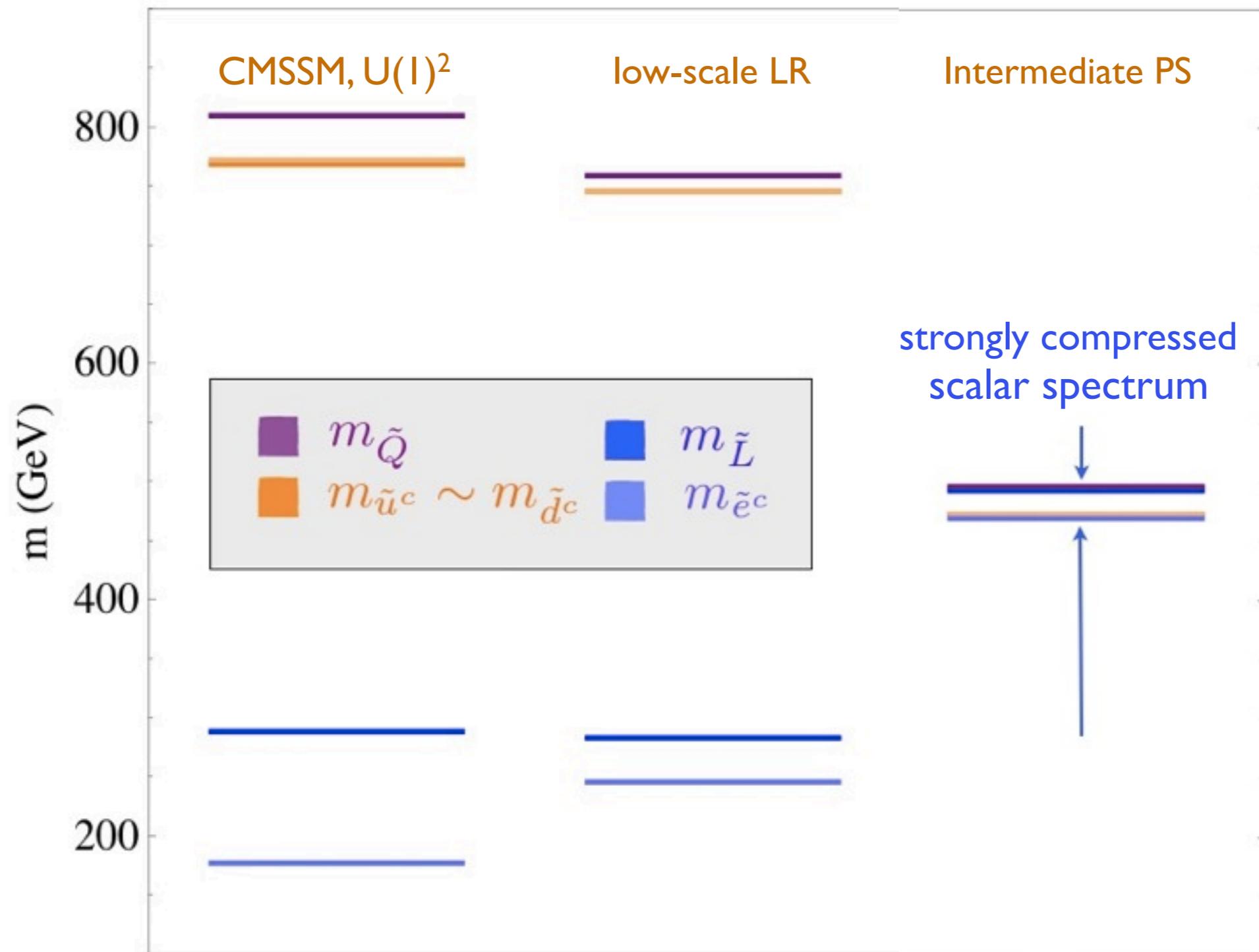


$$QU \equiv (m_{\tilde{Q}}^2 - m_{\tilde{U}}^2)/M_1^2$$

$$LE \equiv (m_{\tilde{L}}^2 - m_{\tilde{E}}^2)/M_1^2$$



The spectra



The MSSM soft squark and slepton spectra calculated for the SPS3 point ($m_0 = 90\text{GeV}$, $M_{1/2} = 400\text{GeV}$) and $v_R = 10^3\text{GeV}$.

Few remarks

- * Softs reconstruction is not straightforward - further inputs needed
- * Sleptons and gauginos (if not too heavy) @ ILC - fantastic precision
- * Squarks may be too heavy for ILC, be happy with LHC (**if seen**)
 - long decay chains can be very helpful (few percent precision)

$$\tilde{q} \rightarrow \chi_2^0 q, \quad \chi_2^0 \rightarrow \tilde{l}^\pm l^\mp \rightarrow l^\pm l^\mp \chi_1^0$$

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* **Disclaimer: it is extremely difficult to do all this in full glory!!!**