



MAKE NEW STANDARDS.

東海国立大学機構



NAGOYA
UNIVERSITY

Top-quark mass measurement results and prospects

26/12/25 (Fri.)

Haruka Asada

From Nagoya University

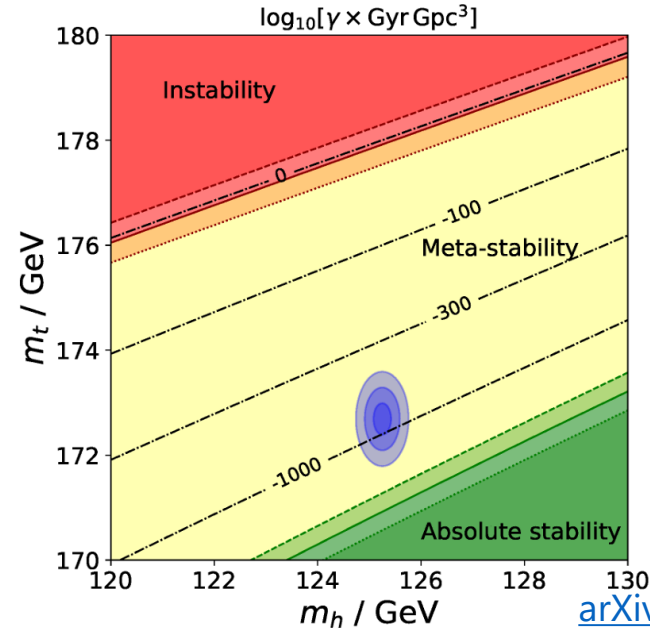
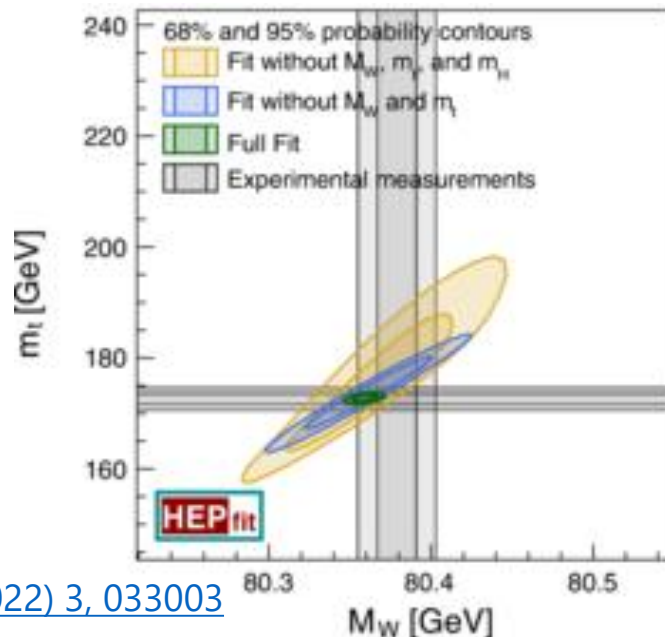
Tera-Scale workshop 2025 @University of Osaka

Motivation of top-quark mass measurement

Top quark is the heaviest elementary particle among the observed ones.

- Top-quark mass (m_{top}) is a fundamental parameter of the SM.
 - Relationship among the Higgs-boson mass, W -boson mass, and m_{top} helps to test the SM by comparing to the EW fit predictions.
 - The Higgs quartic coupling λ at high scales depends on the m_{top} and determine the EW vacuum stability.

→ Precise m_{top} measurements are crucial for the validation of the SM.



[PRD 106 \(2022\) 3, 033003](#)

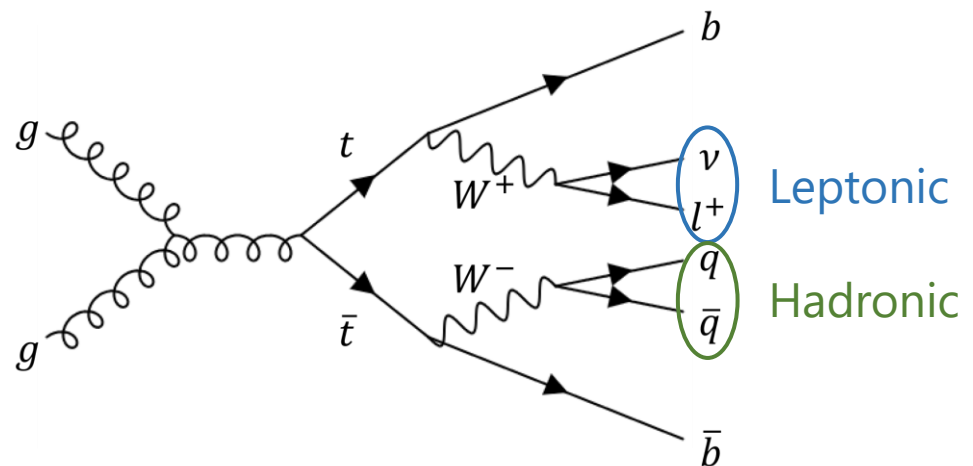
[arXiv:1707.09301](#)

Top-quark mass measurements

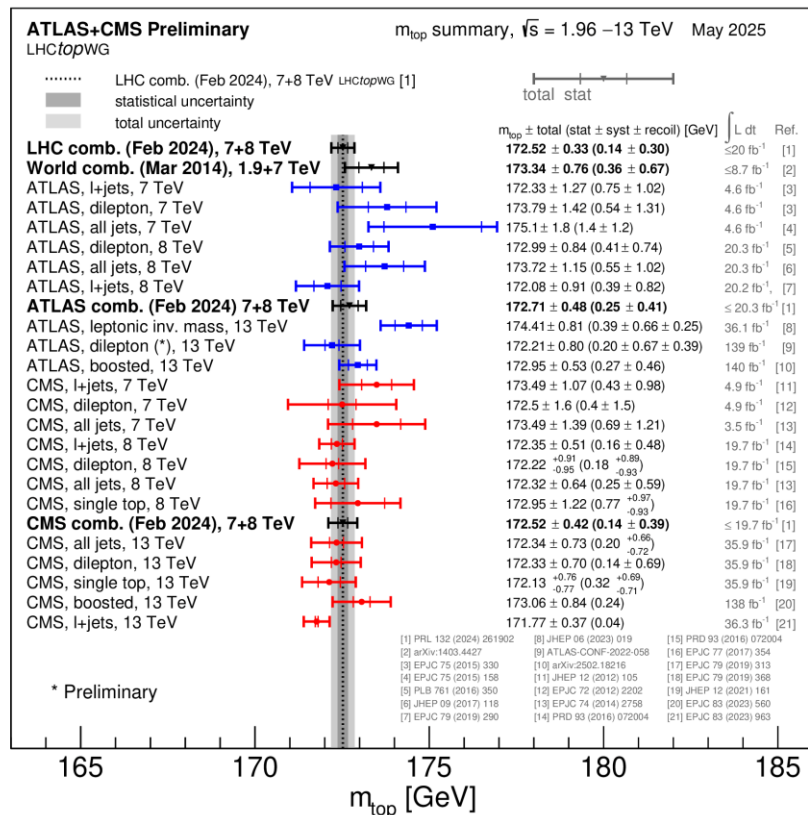
- Direct measurements
 - Reconstruct the top quark from final-state kinematics and extract the mass by comparing distributions with MC predictions.
 - Measure the MC mass: input to MC simulations, including non-perturbative QCD effects modeled by the MC framework.
- Indirect measurements
 - Derived from measurements of the $t\bar{t}$ production cross section.
 - Measure the pole mass (or similar): parameter in quantum field theory and more straightforward to interpret theoretically.

$t\bar{t}$ events are classified into three decay channels based on the W -boson decay modes:

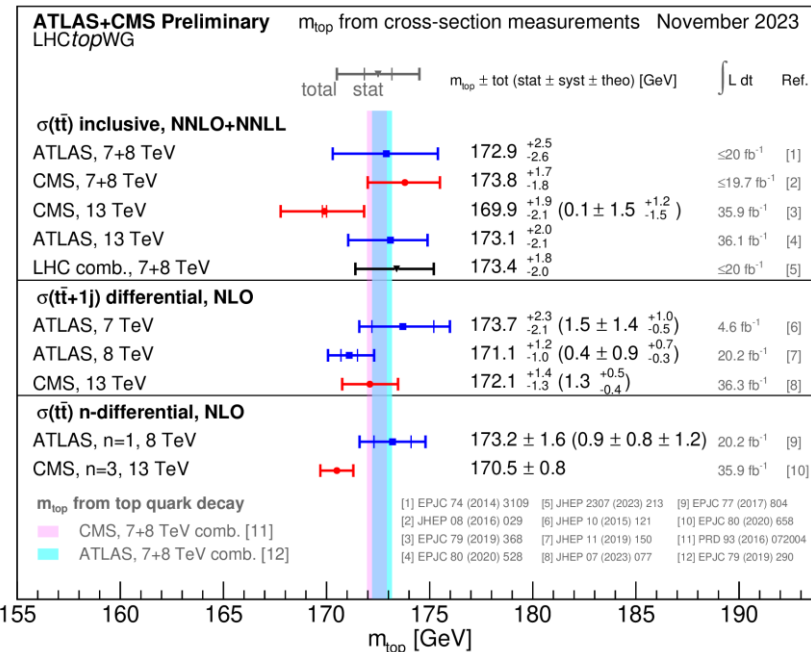
- Dilepton ($\sim 7\%$)
- All hadronic ($\sim 55\%$)
- ℓ +jets ($\sim 38\%$)



Direct measurements



Indirect measurements



LHC Run 1 combined results:

$172.52 \pm 0.33 \text{ GeV}$

$173.4_{-2.0}^{+1.8} \text{ GeV}$ ($\sigma(t\bar{t})$ inclusive)

The relation between MC mass and the pole mass remains an open question.

- The conversion is estimated to have an ambiguity (0.5—1 GeV).

[Ann. Rev. Nucl. Part. Sci. 70, 225 \(2020\)](#)

Top mass measurement in 2025

Using Run 1 dataset, m_{top} has been already measured with high precision.

→ Further improvements require reducing systematic uncertainties, e.g. by exploiting rare phase space regions.

This year, LHC reported three m_{top} measurements.

- Full Run 2 dataset ($\sqrt{s}=13$ TeV, 140 fb^{-1}) is used.

- **Measurement of the top quark mass with the ATLAS detector using $t\bar{t}$ events with a high transverse momentum top quark**
- [Physics Letters B 867 \(2025\) 139608](#)
- **Measurement of the top-quark mass using decays with a J/ψ meson at $\sqrt{s}=13$ TeV with the ATLAS detector**
- [arXiv:2511.23091](#) (Submitted to JHEP)
- **Measurement of the top-quark pole mass in dileptonic $t\bar{t} + 1\text{-jet}$ events at $\sqrt{s}=13$ TeV with the ATLAS experiment**
- [JHEP 12 \(2025\) 023](#)

Direct measurements

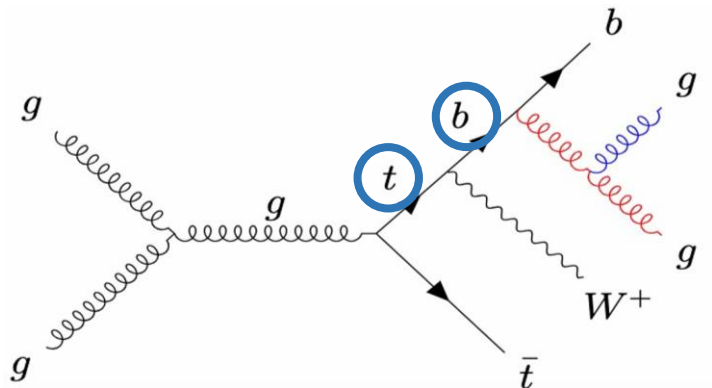
Indirect measurements

Significant uncertainties in MC mass

Uncertainties in the Jet Energy Scale (JES)

- JES calibration restores the jet energy to that of jets reconstructed at the particle level.
- One of the dominant uncertainties in Run 1 measurements.

Uncertainty from the choice of secondary gluon radiation models.



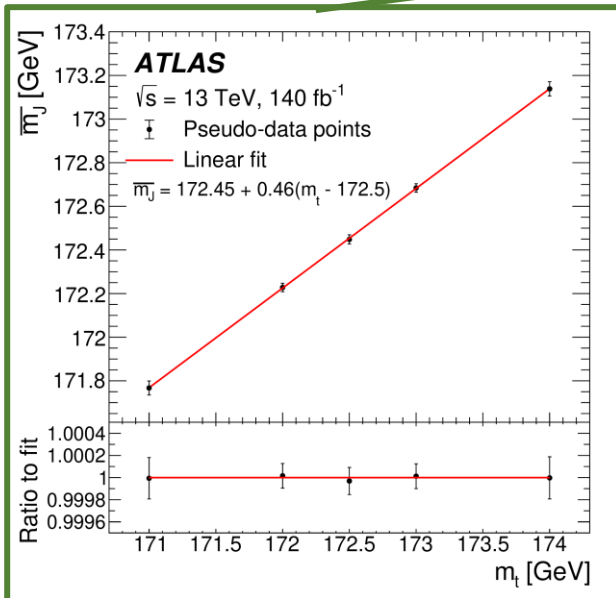
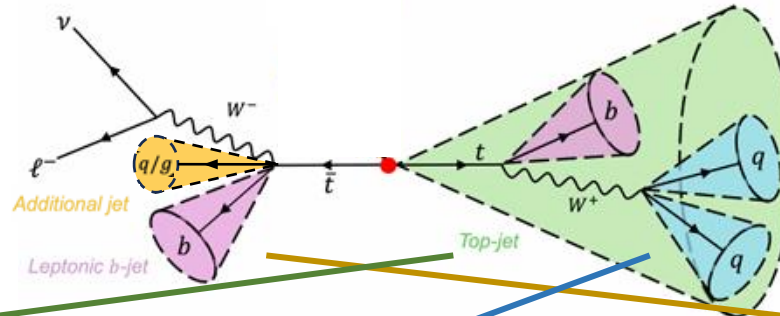
Model in Pythia8	Recoiler
Recoil-To-Colour=ON	b
Recoil-To-Top (new)	t

- Difference between the two models taken as uncertainty.
 - A new model was proposed by the Pythia8 authors a few years ago.
- Affects $p_T(b \text{ jet})$, $p_T(\ell)$, $\Delta R(\ell, \text{jet})$, and b -quark fragmentation modelling.
- Currently, only ATLAS accounts for this uncertainty.

Observables

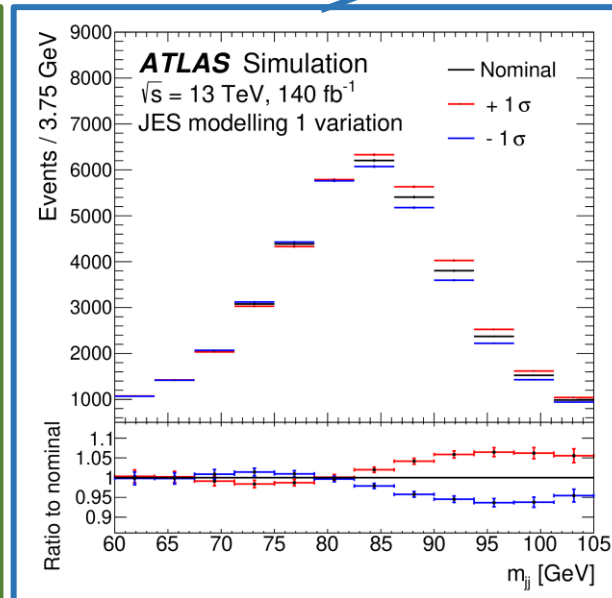
PLB 867 (2025) 139608

Tuned to be mutually uncorrelated.



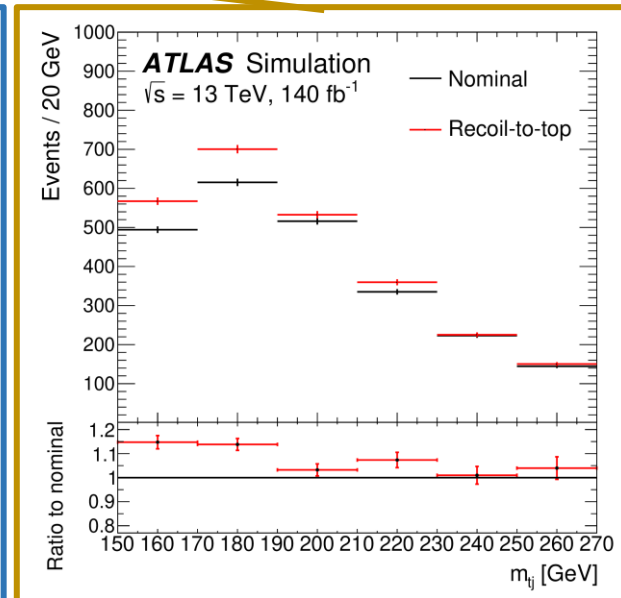
Mean value of m_j (\bar{m}_j)

- Calculated with $145 < m_j < 205 \text{ GeV}$ events.
- Sensitive to m_{top} .



Inv. mass of the two non b -tagged constituent jets (m_{jj})

- Sensitive to JES.



Inv. mass of the semi-leptonically decaying top quark and the closest additional jet (m_{tj})

- Sensitive to QCD radiation.

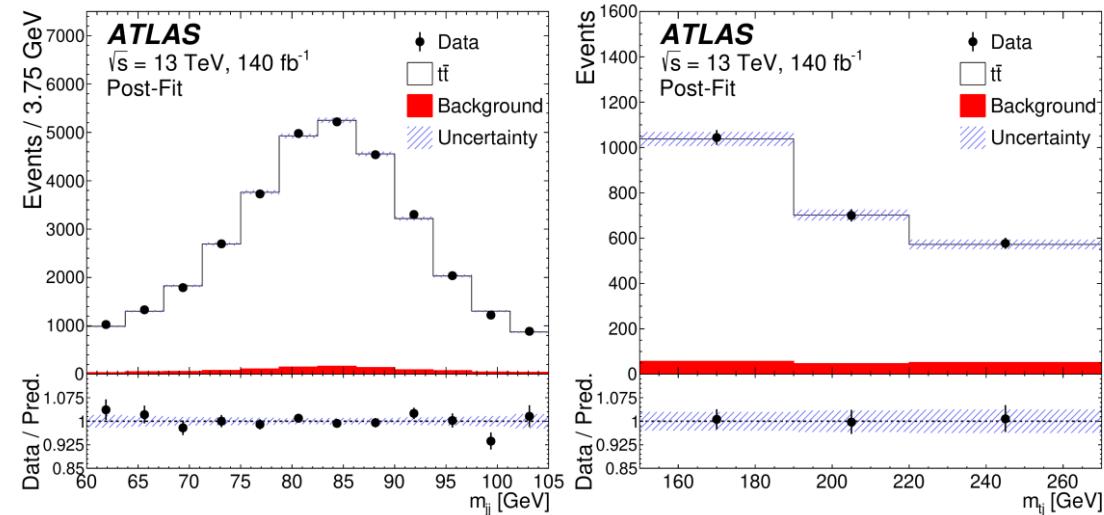
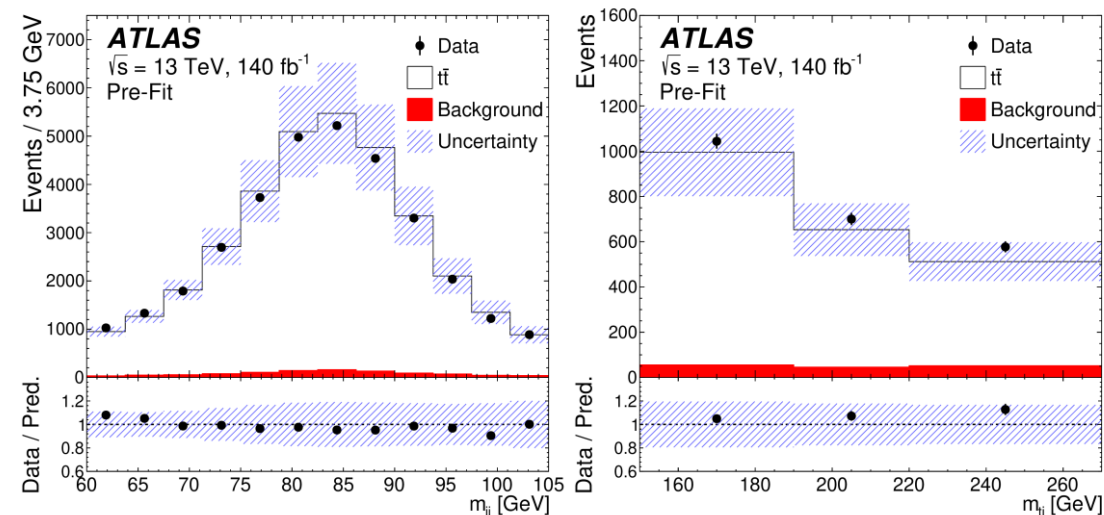
Impact on the uncertainty estimation

[PLB 867 \(2025\) 139608](#)

$m_{\text{top}} = 172.5 \text{ GeV}$:

Source	Uncertainty [GeV]	
	Pre fit	Post fit
JES	1.4	0.26
ISR & FSR	0.82	0.14
Recoil	0.36	0.08
Total	1.7	0.51

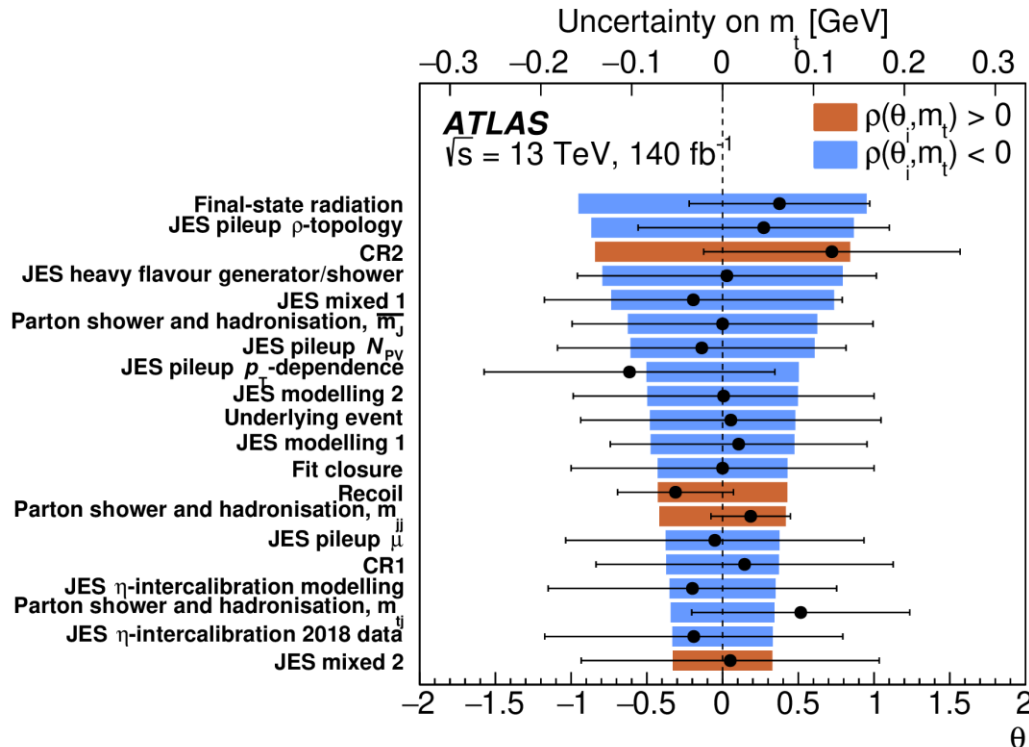
Significant improvement on the precision by adding m_{jj} and m_{tj} to the fit.



$m_{\text{top}} = 172.95 \pm 0.53 \text{ GeV}$

The most precise ATLAS m_{top} measurement in a single channel.

The profile likelihood fit with three observables successfully reduced the uncertainties due to JES and recoil effect by $\sim 80\%$.



Source	Uncertainty [GeV]
JES	± 0.29
Radiation (ISR and FSR)	± 0.17
Colour reconnection (CR1 and CR2)	± 0.15
JES heavy flavour	± 0.14
Parton shower and hadronisation model	± 0.14
JER	± 0.10
MC statistics	± 0.08
Underlying event	± 0.08
Recoil	± 0.07
Fit closure	± 0.07
Background modelling	± 0.05
Matrix element matching ($p_T^{\text{hard}} = 1$)	± 0.04
b -tagging	± 0.04
Higher-order corrections	± 0.02
E_T^{miss}	± 0.02
Pileup	± 0.01
JVT	± 0.01
PDF	± 0.01
Leptons	± 0.01
Luminosity	< 0.01
Total statistical	± 0.27
Total systematic	± 0.46
Total	± 0.53

MC mass using J/ψ meson

[arXiv:2511.23091](https://arxiv.org/abs/2511.23091)

Motivations:

- Purely leptonic observables.
→ Less sensitive to JES than the ones from jet reconstruction.
- Partial reconstruction of decay.
→ Less sensitive to top p_T modelling than pure prompt lepton kinematics.

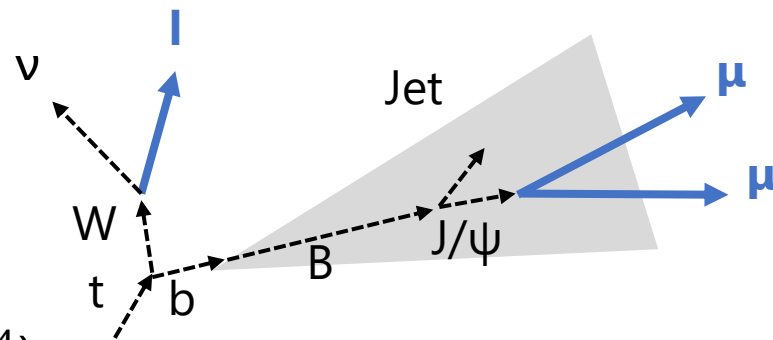
Complementary measurement with different syst. uncertainty sources.
→ Help to reduce the uncertainties in combination.

$t \rightarrow (W \rightarrow l\nu)(b \rightarrow J/\psi + X \rightarrow \mu^+\mu^- + X)$ process ($l = e, \mu$) is used.

- The first measurement using this phase space in ATLAS.

Challenges:

- Small BR of the signal process ($\sim 2.7 \times 10^{-4}$)
→ Expected large stat. uncertainty.
- Still sensitive to parton shower, hadronization, and b -fragmentation effects.
→ Must be controlled well.

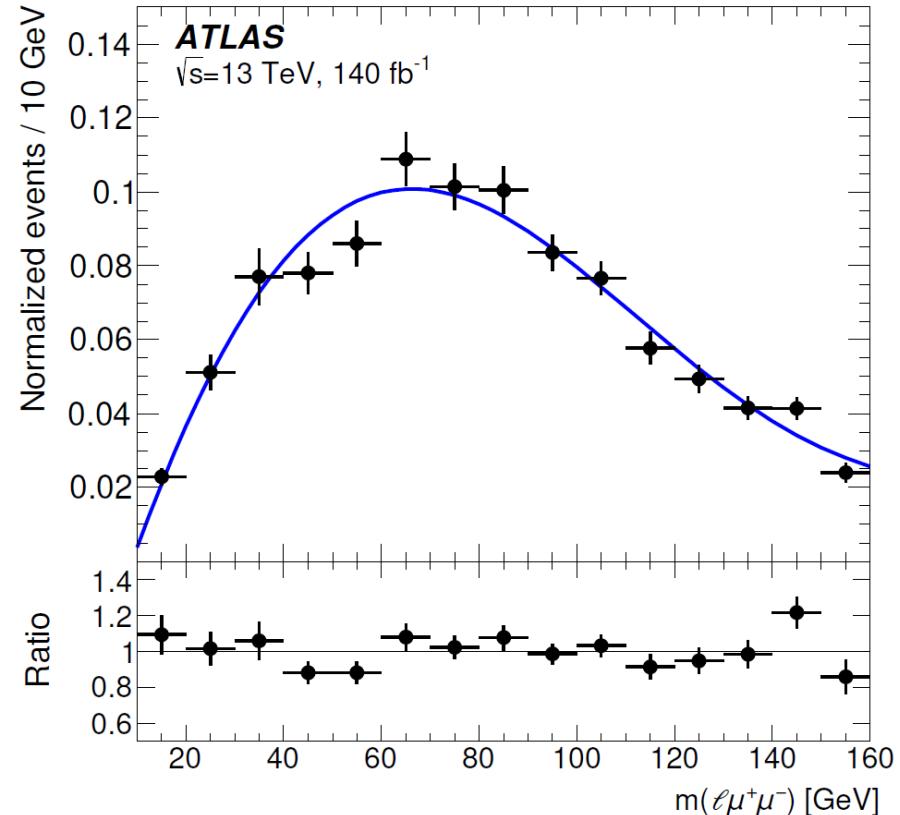
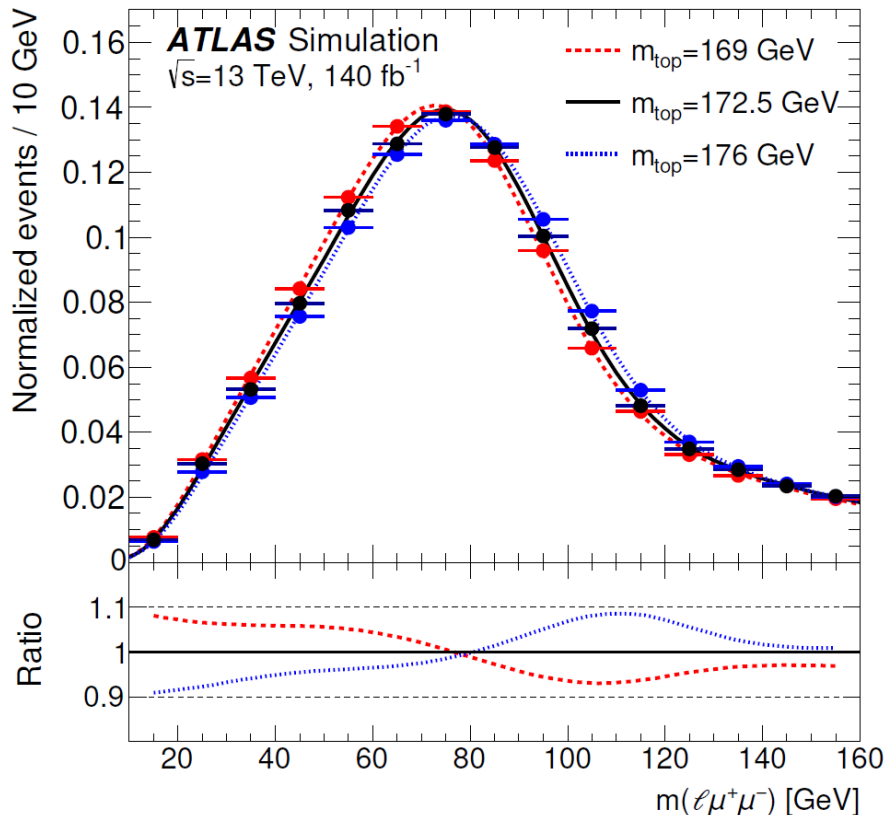


PDF for likelihood fit

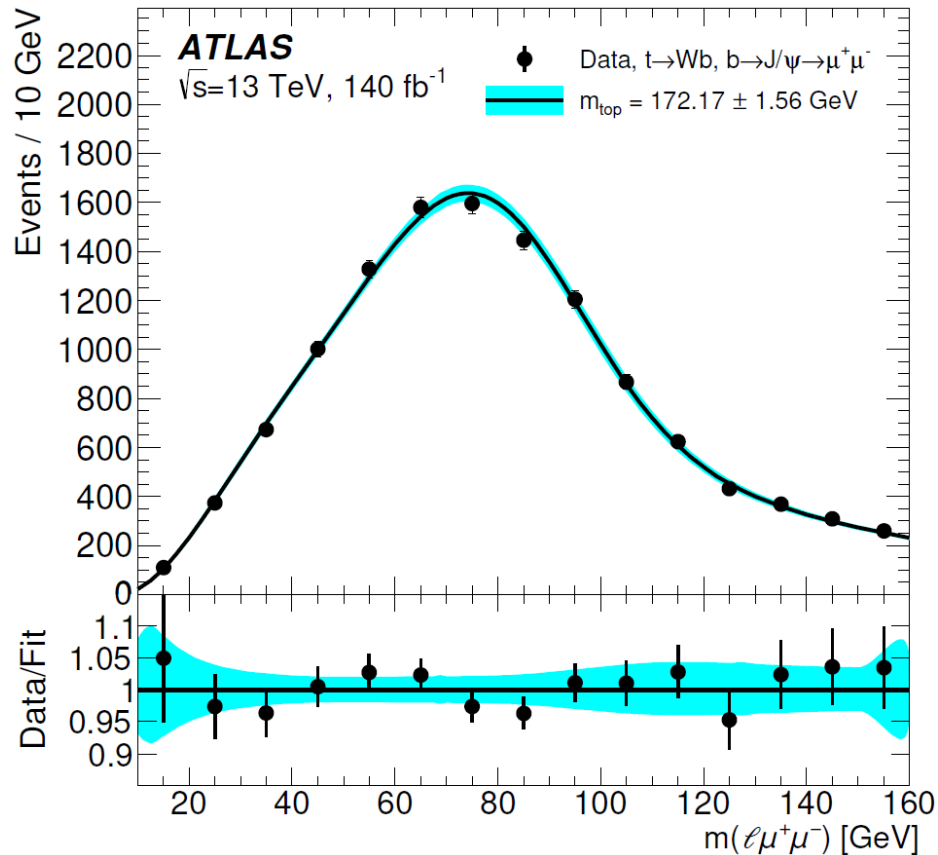
[arXiv:2511.23091](https://arxiv.org/abs/2511.23091)

m_{top} is obtained by the unbinned likelihood fit to **the $m(\ell\mu^+\mu^-)$ distribution**.
The distribution shape is parametrized by the following functions:

- $t\bar{t}$ and single-top events:
 m_{top} -dependent component.
 - PDF=Gaussian + Lognormal
- V +jets, fake lepton events, etc:
 m_{top} -independent component.
 - PDF=Chebyshev polynomial



$$m_{\text{top}} = 172.17 \pm 0.80 \text{ (stat)} \pm 0.81 \text{ (syst)} \pm 1.07 \text{ (recoil) GeV}$$
$$= 172.17 \pm 1.56 \text{ GeV}$$



Stat. uncertainty is of the same order as in other measurements.

Jet-related uncertainties are smaller than in any measurements using other channels (<100 MeV).

The recoil modeling uncertainty is dominant: 1.07 ± 0.22 GeV.

Impact on b -fragmentation modelling is significant.

- b -frag. model in new recoil scheme has not been tuned yet.

→ The uncertainty is conservatively included, separately from other systematic modelling uncertainties.

Other syst. uncertainties are well controlled and reduced to a level comparable to the stat uncertainty.

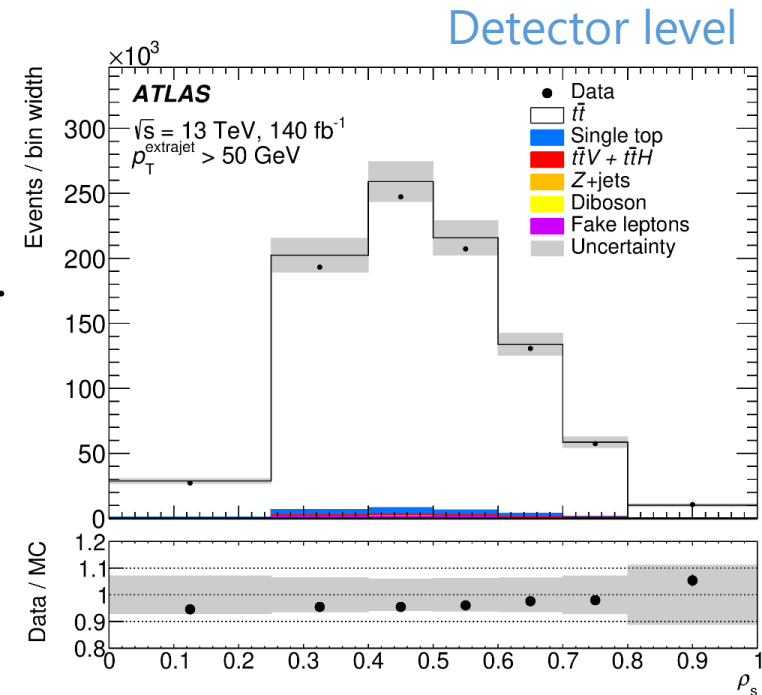
Pole mass using dileptonic $t\bar{t}+1$ jet

[JHEP 12 \(2025\) 023](#)

The normalized differential cross-section defined as

$$\mathcal{R}(\rho_s; m_t^{pole}) = \frac{1}{\sigma_{t\bar{t}+1\text{-jet}}} \cdot \frac{d\sigma_{t\bar{t}+1\text{-jet}}}{d\rho_s}, \text{ with } \rho_s = \frac{2m_0}{\sqrt{s_{t\bar{t}+1\text{-jets}}}} \quad (m_0=170 \text{ GeV})$$

- $\sqrt{s_{t\bar{t}+1\text{-jets}}}$ is the invariant mass of the reconstructed $t\bar{t}+1$ -jet system.
- Existence of an extra 1 jet enhances sensitivity of cross-section to the pole mass.
- Dilepton channel, especially required $e^\pm\mu^\mp$.
- 98% efficiency in $t\bar{t}$ reconstruction by combined two dedicated methods to estimate the kinematics of neutrinos.
- 95% purity of final selection with extra jet of $p_T > 50 \text{ GeV}$, $|\eta| < 2.5$.



Unfolding and χ^2 fit

The ρ_s distribution is unfolded to the parton level by using Iterative Bayesian Unfolding.

Measurement is performed for the two theoretical predictions separately:

- **$pp \rightarrow t\bar{t}j$ (2 \rightarrow 3) process**

- [JHEP 05 \(2022\) 146](#)

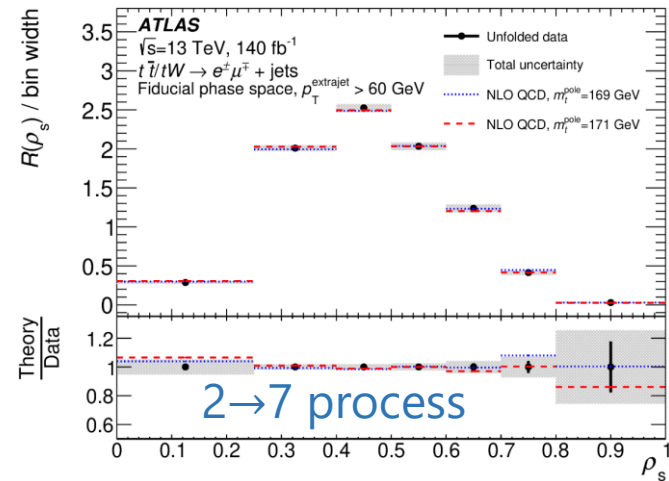
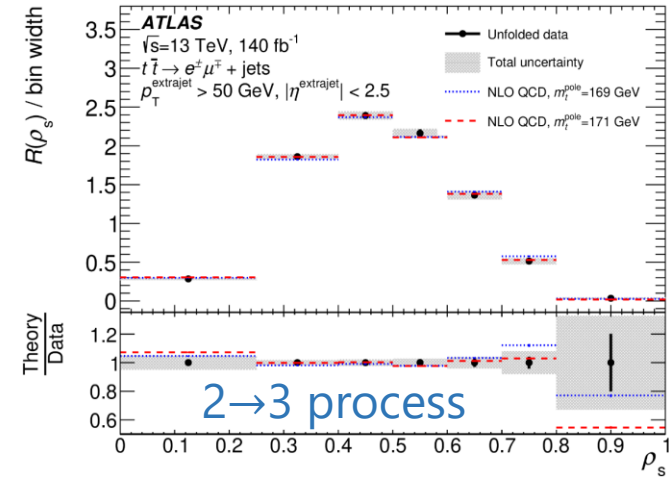
- Top quarks are taken as stable particles.

- **$pp \rightarrow b\bar{b}\ell^+\nu\ell^-\bar{\nu}j$ (2 \rightarrow 7) process**

- [JHEP 11 \(2016\) 098](#)

- Includes top-quark decays to the dilepton final state and off-shell effects.

χ^2 fit of the unfolded normalized differential cross-section distribution.



Measurement using 2→3 process shows the best precision:

$$m_t^{\text{pole}} = 170.73 \pm 0.33 \text{ (stat.)} \pm 1.36 \text{ (syst.)} \begin{matrix} +0.34 \\ -0.28 \end{matrix} \text{ (scale)} \pm 0.24 \text{ (PDF} \oplus \alpha_s \text{) GeV}$$

(PDF set : PDF4LHC21)

→ Main result of this analysis.

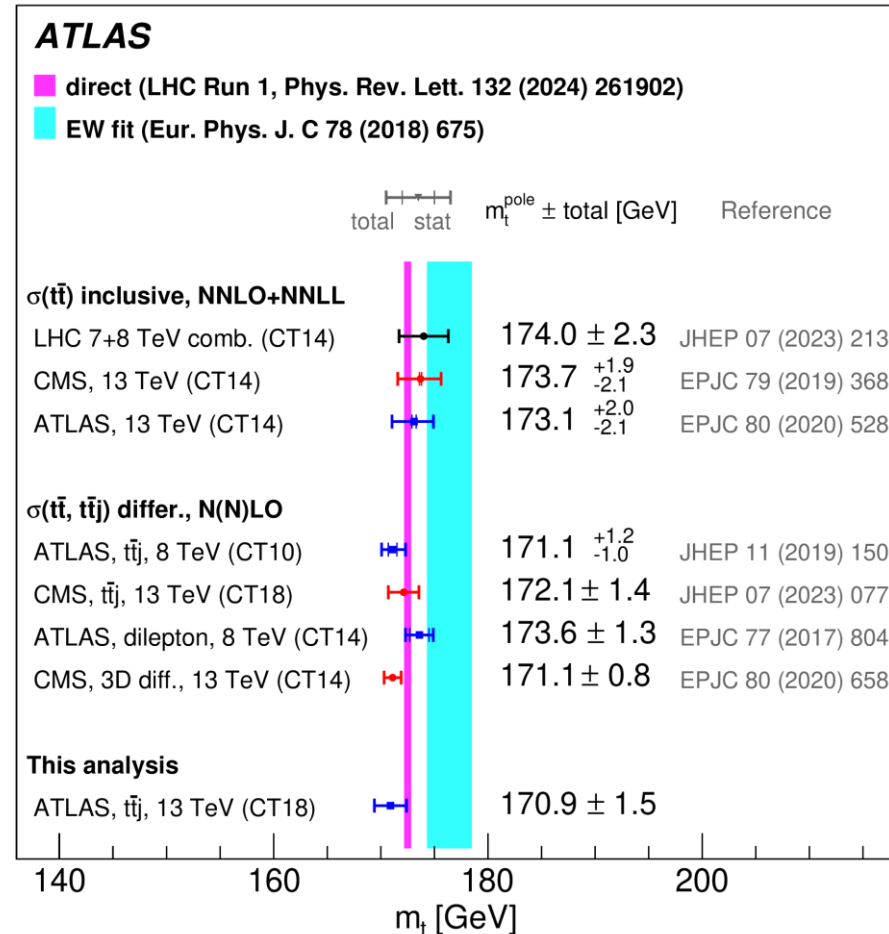
- $t\bar{t}$ modelling, JES and b -tagging are dominant source of uncertainties.

Compatible with 2 → 7 one.

- 2 → 7 has reduced sensitivity to m_t^{pole} and suffers from larger theoretical scale uncertainties (expected).

Overall good compatibility with other determinations.

- A difference of ~5 GeV from EW precision observables.



Summary

The LHC has pursued top-quark mass measurements intensively.

- Run 1: combined results already achieved high precision.
- **Run 2: further reduction of systematic uncertainties using advanced analysis**, e.g. simultaneous fits, rare phase space.
 - Central values are broadly consistent across analyses.

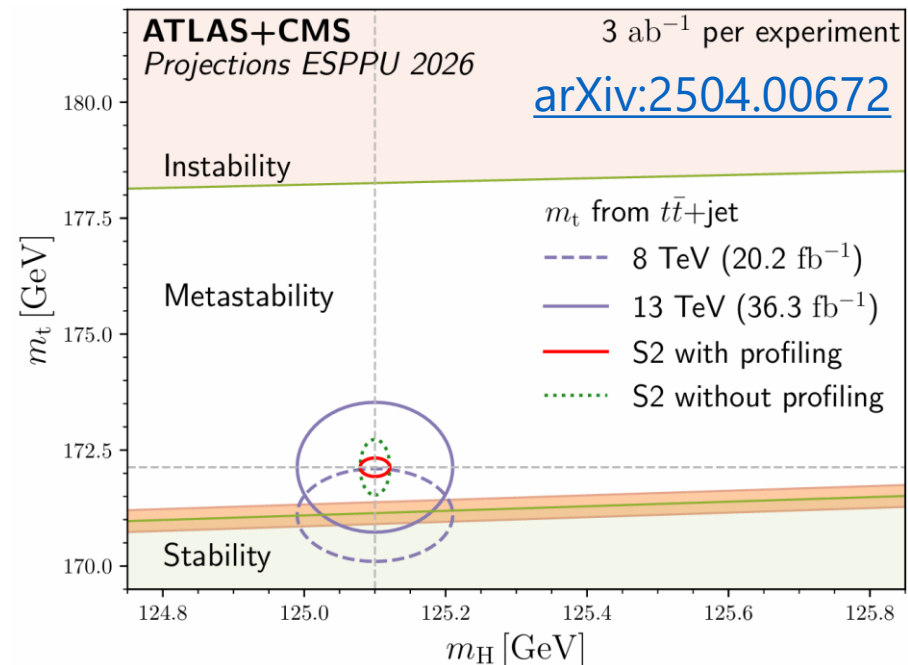
Projection of the m_{top} measurements for **HL-LHC** has been performed.

- **Profile likelihood fits** with large datasets are expected to significantly reduce measurement uncertainties

m_{top} measurements using LHC and HL-LHC dataset will show further improvements and help advance validation of the SM.

- E.g. EW fit, vacuum stability.

Stay tuned for future results!



With profiling, precision will be ~ 200 MeV.

Backup