

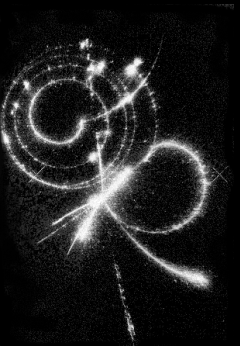
$H_2 \rightarrow H_1 H_1$ decays in the cMSSM

Hunting for the Higgs in a Complex World



Karina Williams, IPPP, Durham

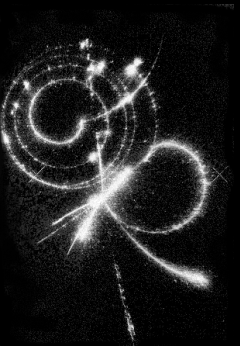
In collaboration with Georg Weiglein plus others



Minimal Supersymmetric Standard Model



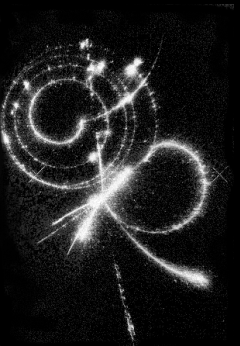
- **Minimum** number of superpartners
- **2 Higgs Doublets**
- No assumption about SUSY breaking mechanism - All terms that **break SUSY softly** are added to Lagrangian
- Has **more than 100 free parameters** (in addition to those in the SM)
- Predicts that the **lightest Higgs Mass < 140 GeV** – within reach of the LHC and a LC



complex MSSM (cMSSM)



- Not to be confused with the **CMSSM** – the **Constrained** MSSM
- The extra complex phases give rise to **CP violation** in the Higgs Sector
- A theory with CP violation is needed to describe the **matter-antimatter asymmetry** in the universe



Stops in the cMSSM



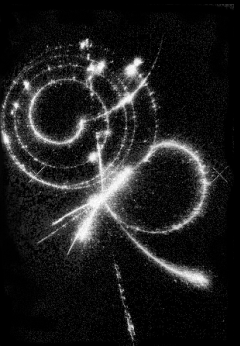
Stop mass mixing matrix can have 2 complex parameters:

$$M_{\tilde{t}} = \begin{pmatrix} M_L^2 + m_t^2 + DT_1 & m_t X_t^* \\ m_t X_t & M_R^2 + m_t^2 + DT_2 \end{pmatrix}$$

$$X_t = A_t - \mu^* \cot \beta$$

Rotate to get physical states:

$$\begin{pmatrix} \tilde{t}_L \\ \tilde{t}_R \end{pmatrix} \rightarrow \begin{pmatrix} \tilde{t}_1 \\ \tilde{t}_2 \end{pmatrix}$$



The Higgs Potential in the cMSSM



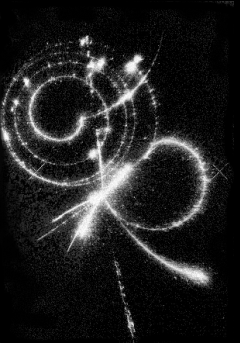
The Higgs potential also has **complex** parameters:

$$V_H = m_1^2 H_{1i}^* H_{1i} + m_2^2 H_{2i}^* H_{2i} + \epsilon_{ij} (m_{12}^2 H_{1i} H_{2j} + m_{12}^{2*} H_{1i}^* H_{2j}^*) \\ + \frac{1}{8} (g_1^2 + g_2^2) (H_{1i}^* H_{1i} - H_{2i}^* H_{2i})^2 + \frac{1}{2} g_2^2 |H_{1i}^* H_{2i}|^2.$$

One of the Higgs doublets has a **complex** phase:

$$\mathcal{H}_1 = \begin{pmatrix} H_{11} \\ H_{12} \end{pmatrix} = \begin{pmatrix} v_1 + \frac{1}{\sqrt{2}} (\phi_1 - i\chi_1) \\ -\phi_1^- \end{pmatrix},$$

$$\mathcal{H}_2 = \begin{pmatrix} H_{21} \\ H_{22} \end{pmatrix} = e^{i\xi} \begin{pmatrix} \phi_2^+ \\ v_2 + \frac{1}{\sqrt{2}} (\phi_2 + i\chi_2) \end{pmatrix}$$

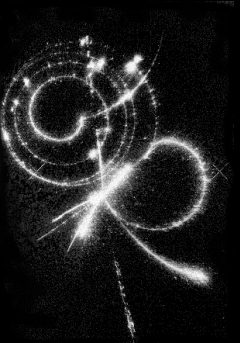


Lowest Order



Rotate to physical states h, H, A, G
Lowest order has no CP violation

$$\begin{pmatrix} h \\ H \\ A \\ G \end{pmatrix} = \begin{pmatrix} -\sin \alpha & \cos \alpha & 0 & 0 \\ \cos \alpha & \sin \alpha & 0 & 0 \\ 0 & 0 & -\sin \beta_n & \cos \beta_n \\ 0 & 0 & \cos \beta_n & \sin \beta_n \end{pmatrix} \cdot \begin{pmatrix} \phi_1 \\ \phi_2 \\ \chi_1 \\ \chi_2 \end{pmatrix}$$



Lowest Order

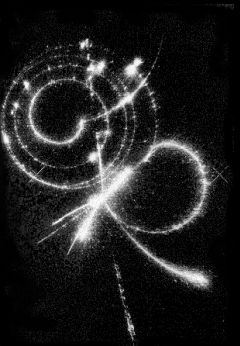


Rotate to physical states h, H, A, G
Lowest order has no CP violation

CP even

$$\begin{pmatrix} h \\ H \\ A \\ G \end{pmatrix} = \begin{pmatrix} -\sin \alpha & \cos \alpha & 0 & 0 \\ \cos \alpha & \sin \alpha & 0 & 0 \\ 0 & 0 & -\sin \beta_n & \cos \beta_n \\ 0 & 0 & \cos \beta_n & \sin \beta_n \end{pmatrix} \cdot \begin{pmatrix} \phi_1 \\ \phi_2 \\ \chi_1 \\ \chi_2 \end{pmatrix}$$

CP odd



Including loop corrections



Rotate to physical states H_1, H_2, H_3

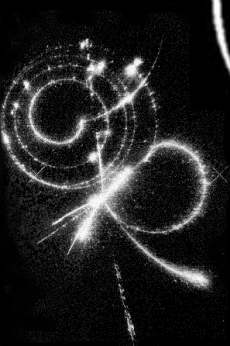
Mixed CP

$$\begin{pmatrix} H_1 \\ H_2 \\ H_3 \end{pmatrix} = \mathbf{U}_n \cdot \begin{pmatrix} h \\ H \\ A \end{pmatrix}$$

where \mathbf{U}_n diagonalises the mass matrix:

$$\begin{pmatrix} m_{h(0)}^2 - \hat{\Sigma}_{hh}(p^2) & -\hat{\Sigma}_{hH}(p^2) & -\hat{\Sigma}_{hA}(p^2) \\ -\hat{\Sigma}_{hH}(p^2) & m_{H(0)}^2 - \hat{\Sigma}_{HH}(p^2) & -\hat{\Sigma}_{HA}(p^2) \\ -\hat{\Sigma}_{hA}(p^2) & -\hat{\Sigma}_{HA}(p^2) & m_{A(0)}^2 - \hat{\Sigma}_{AA}(p^2) \end{pmatrix}$$

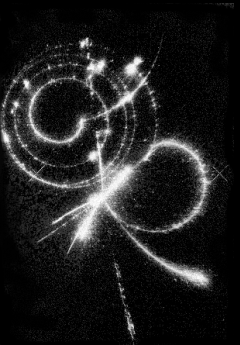
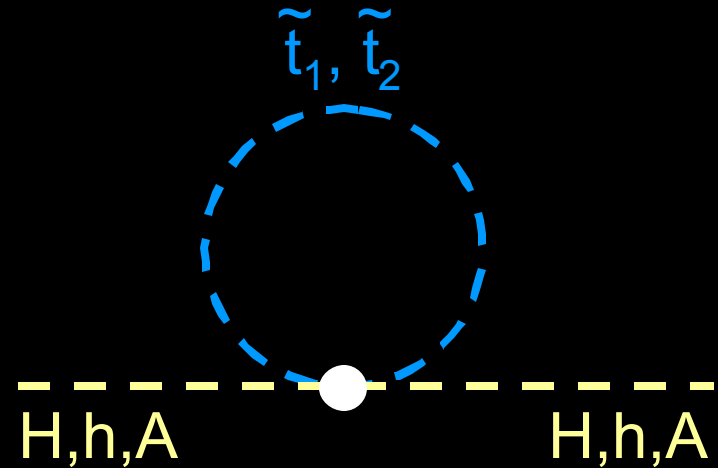
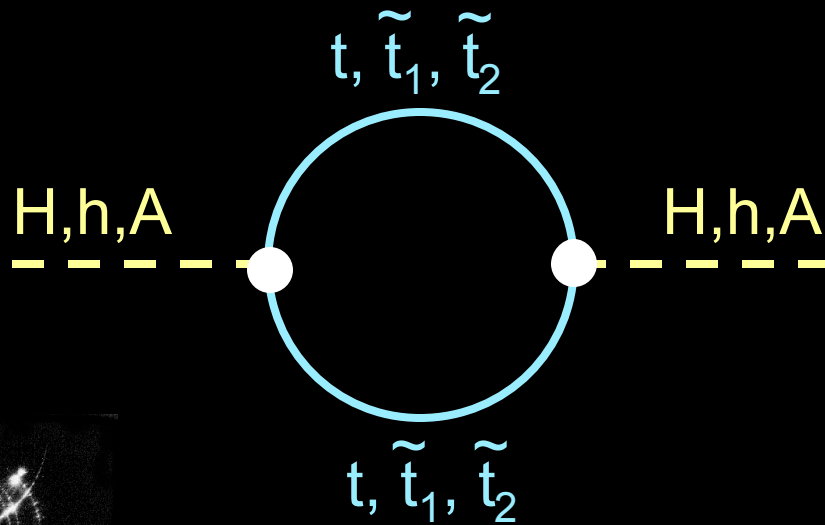
... so need expressions for the renormalised self energies $\hat{\mathbf{S}}$



Leading terms for renormalised self-energies:



- top, stop
- one loop
- terms proportional to m_t^4
- $p^2 = 0$



Renormalisation



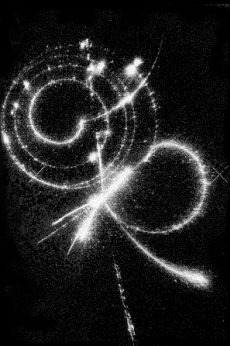
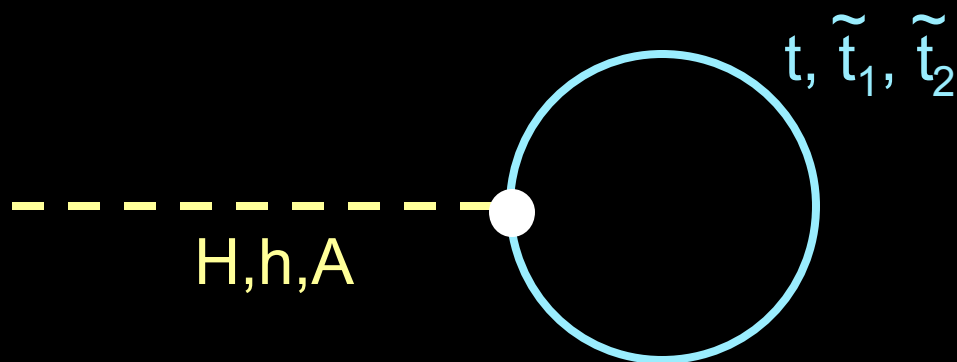
$$\delta T_h = -T_{h(1)}$$

$$\delta T_H = -T_{H(1)}$$

$$\delta T_A = -T_{A(1)}$$

$$\delta m_{H^\pm}^2 = \text{Re} \Sigma_{H^+H^-}(0)$$

No contribution from
 δM_Z^2 or $\delta \tan \beta$



2 computer programs...



FeynHiggs

by S. Heinemeyer, T. Hahn, W. Hollik and G. Weiglein

Feynman diagrammatic approach,
uses on-shell renormalization scheme

CPsuperH

by Jae Sik Lee, Apostolos Pilaftsis, M. Carena,
S. Y. Choi, M. Drees, J. Ellis and C. E. M. Wagner

Renormalisation group improved
effective potential calculation,
uses \overline{MS} renormalization scheme

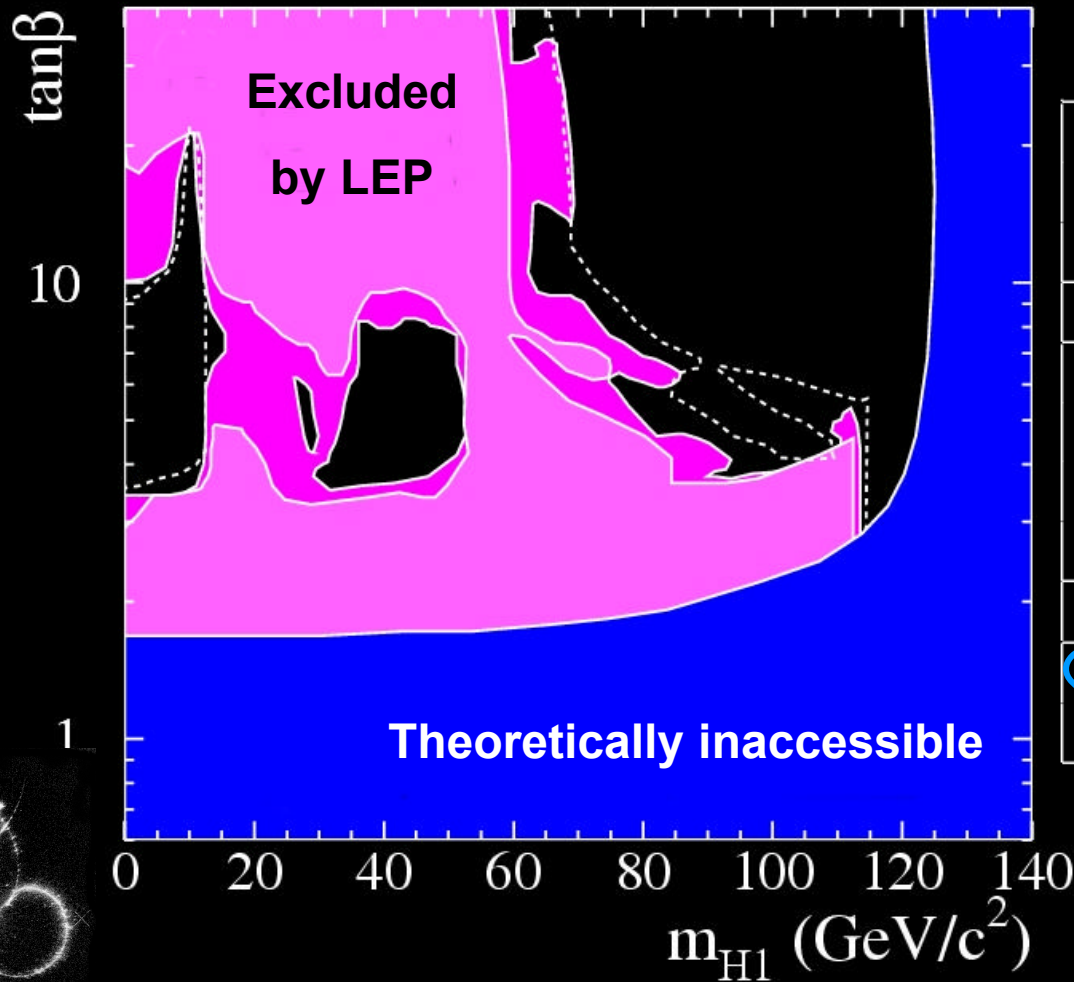
...both calculate higgs masses, couplings,
branching ratios e.t.c. for cMSSM



LEP results for CPX scenario



$$m_t = 174.3 \text{ GeV}/c^2$$

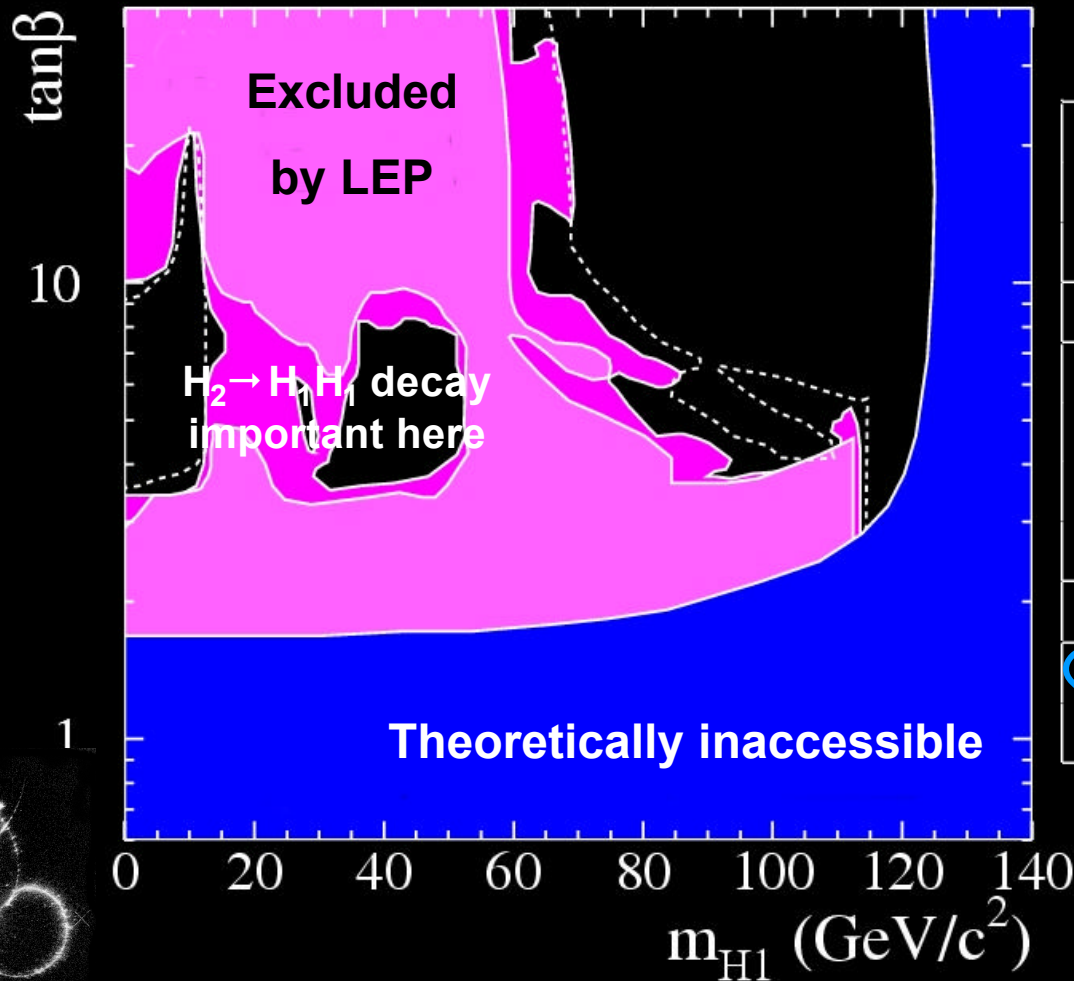


$\tan \beta$	0.6–40
m_A (GeV/ c^2)	–
m_{H^\pm} (GeV/ c^2)	4–1000
<hr/>	
M_{SUSY} (GeV)	500
M_2 (GeV)	200
μ (GeV)	2000
$m_{\tilde{g}}$ (GeV/ c^2)	1000
X_t (GeV)	$A - \mu \cot \beta$
A (GeV) \overline{MS}	1000
$\arg A, \arg m_{\tilde{g}}$	90°

LEP results for CPX scenario



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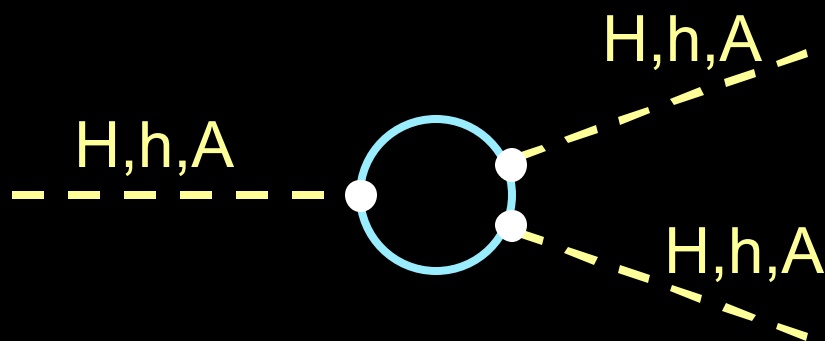


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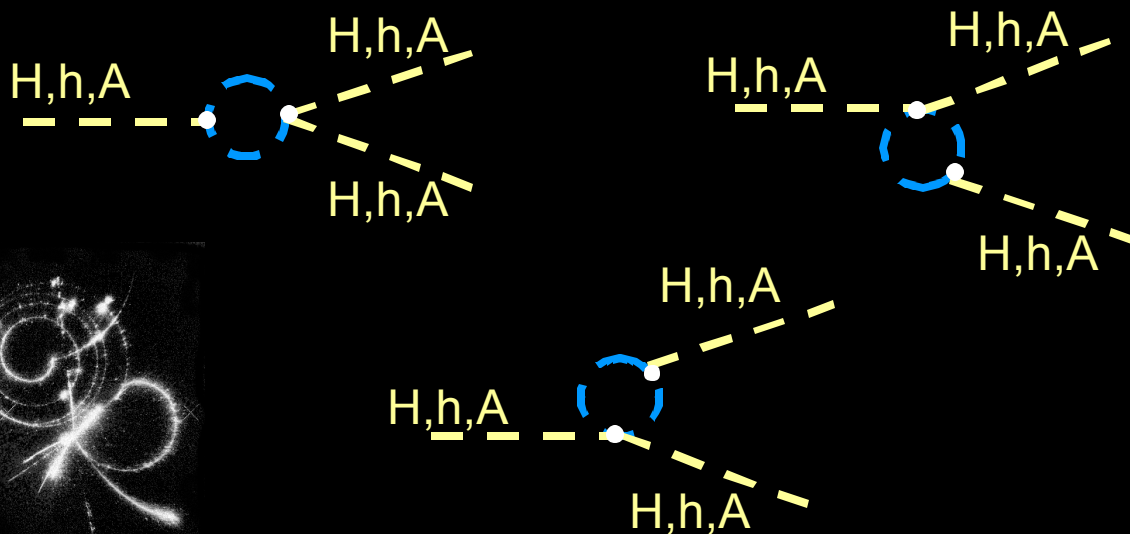


— $t, \tilde{t}_1, \tilde{t}_2$
- - - \tilde{t}_1, \tilde{t}_2

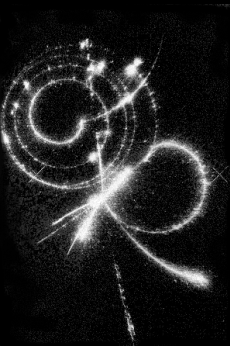
$H_2 \rightarrow H_1 H_1$ vertex corrections



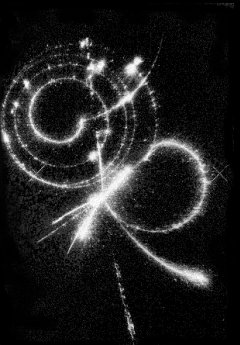
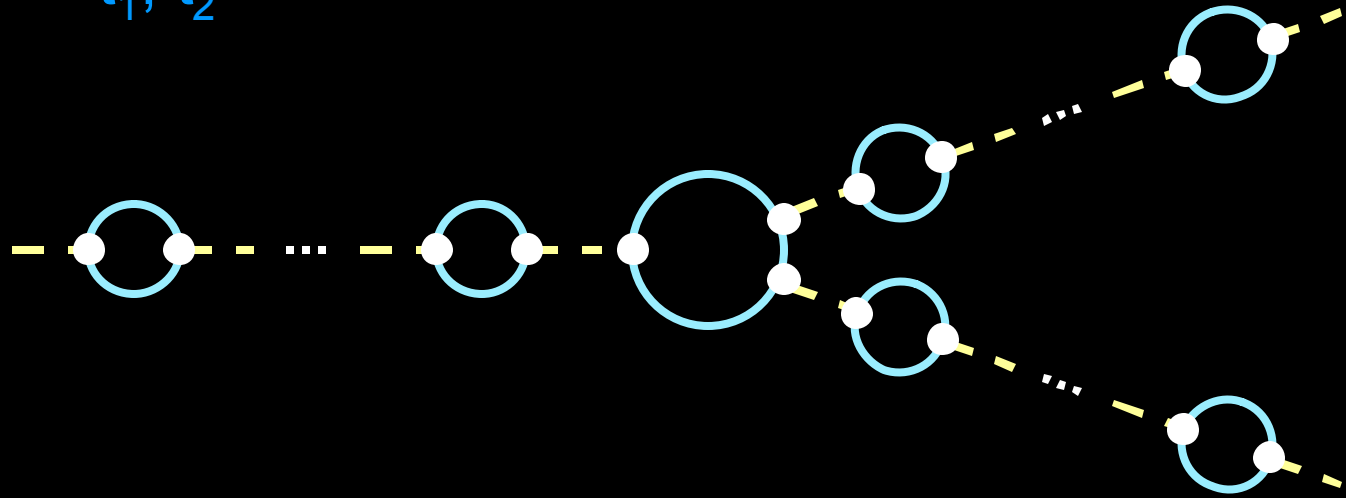
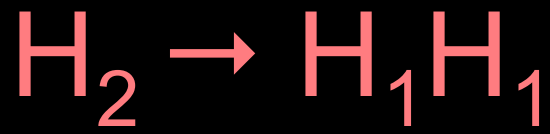
- top, stop
- one loop
- terms proportional to m_t^4
- $p^2 = 0$



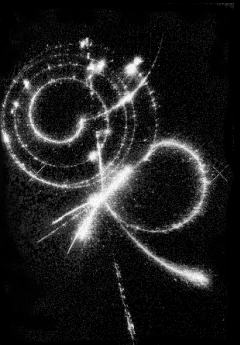
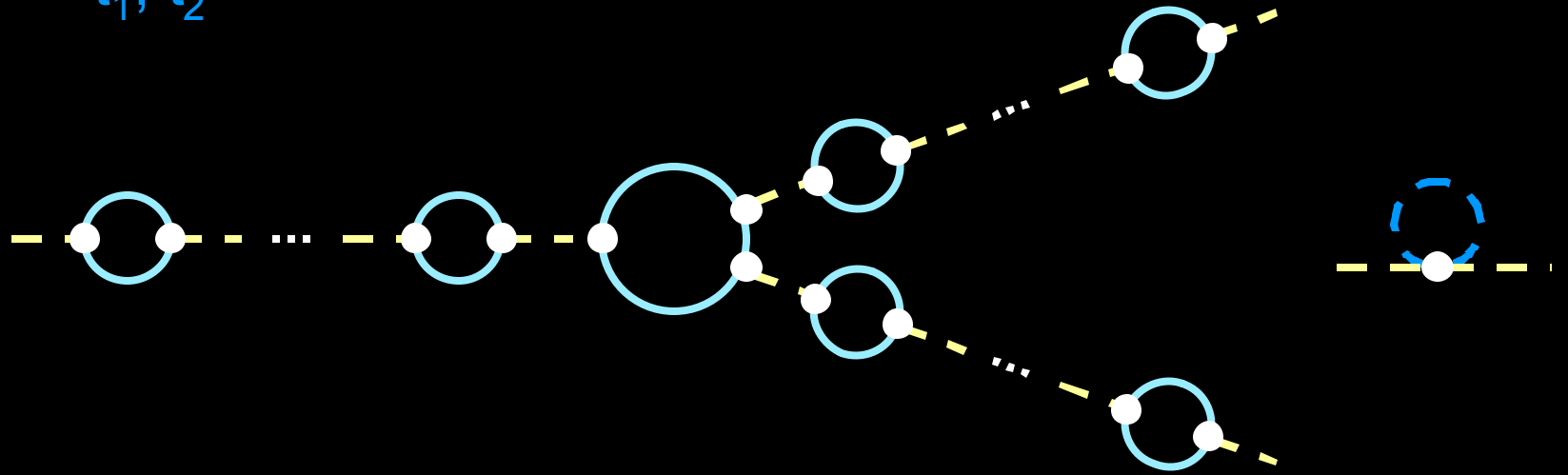
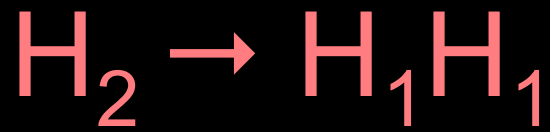
These vertex corrections are **finite** under these approximations – no counter-terms needed



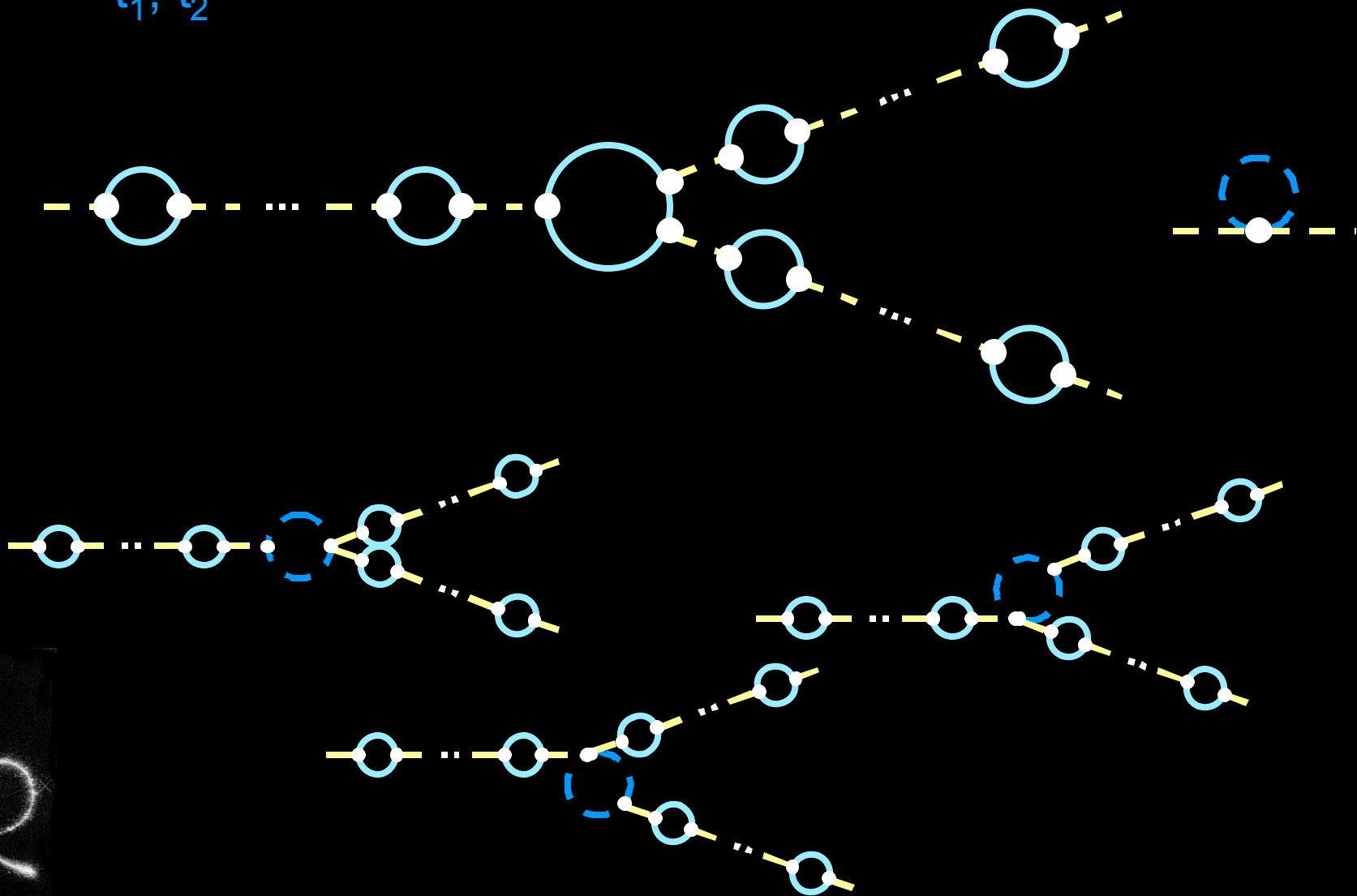
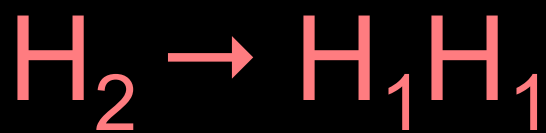
— $t, \tilde{t}_1, \tilde{t}_2$
- - - H, h, A
- - - \tilde{t}_1, \tilde{t}_2



— $t, \tilde{t}_1, \tilde{t}_2$
- - - H, h, A
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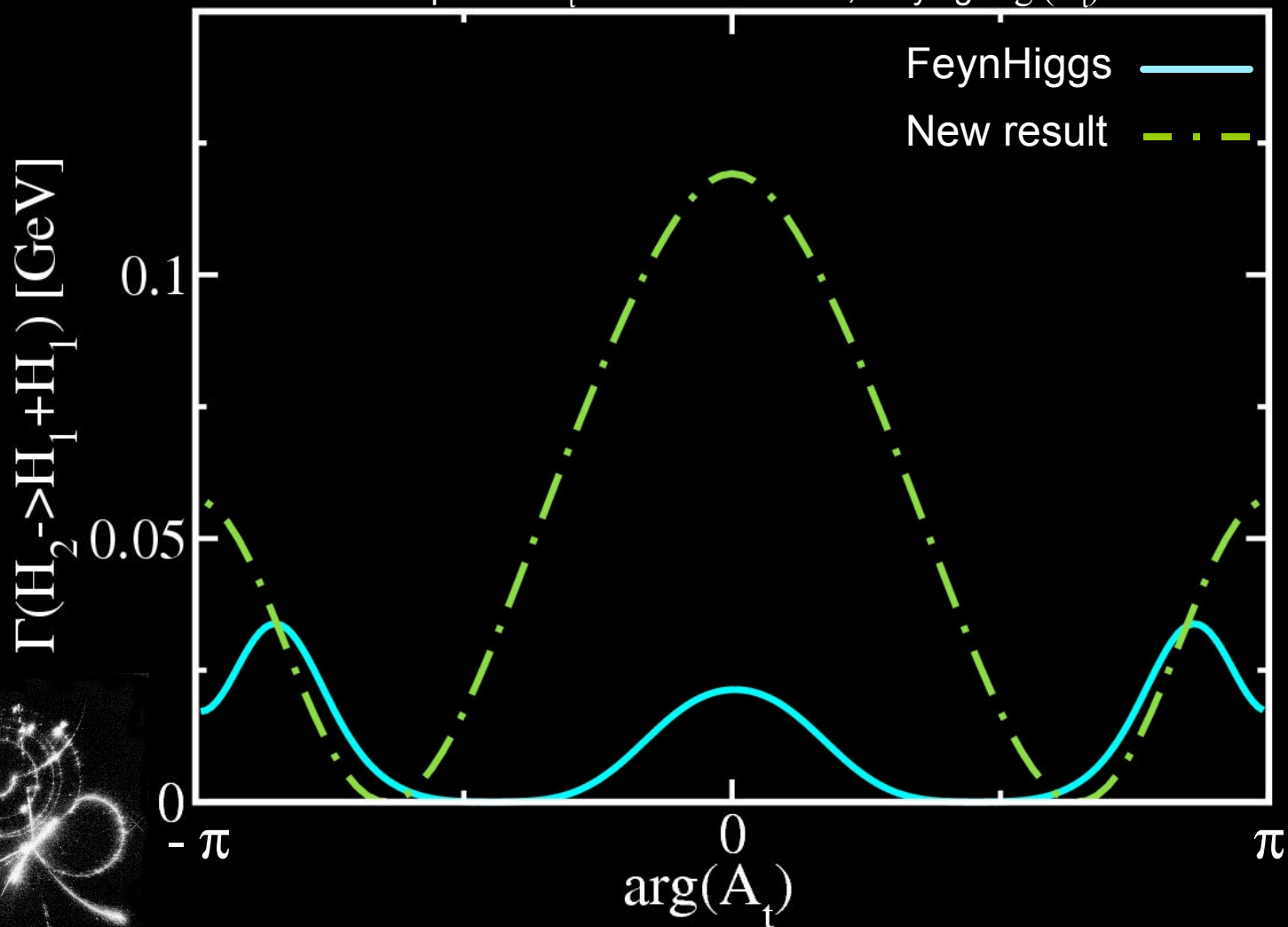
— $t, \tilde{t}_1, \tilde{t}_2$
- - - H, h, A
- - - \tilde{t}_1, \tilde{t}_2



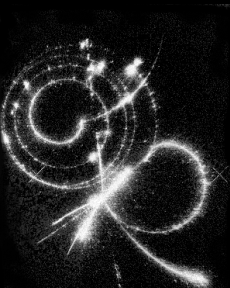
$$H_2 \rightarrow H_1 H_1$$



CPX with $\tan\beta=10$, $M_{H^\pm}=500$ GeV, $m_t=172.7$ GeV
except mod $A_t^{\text{on-shell}}=1000$ GeV, varying $\arg(A_t)$



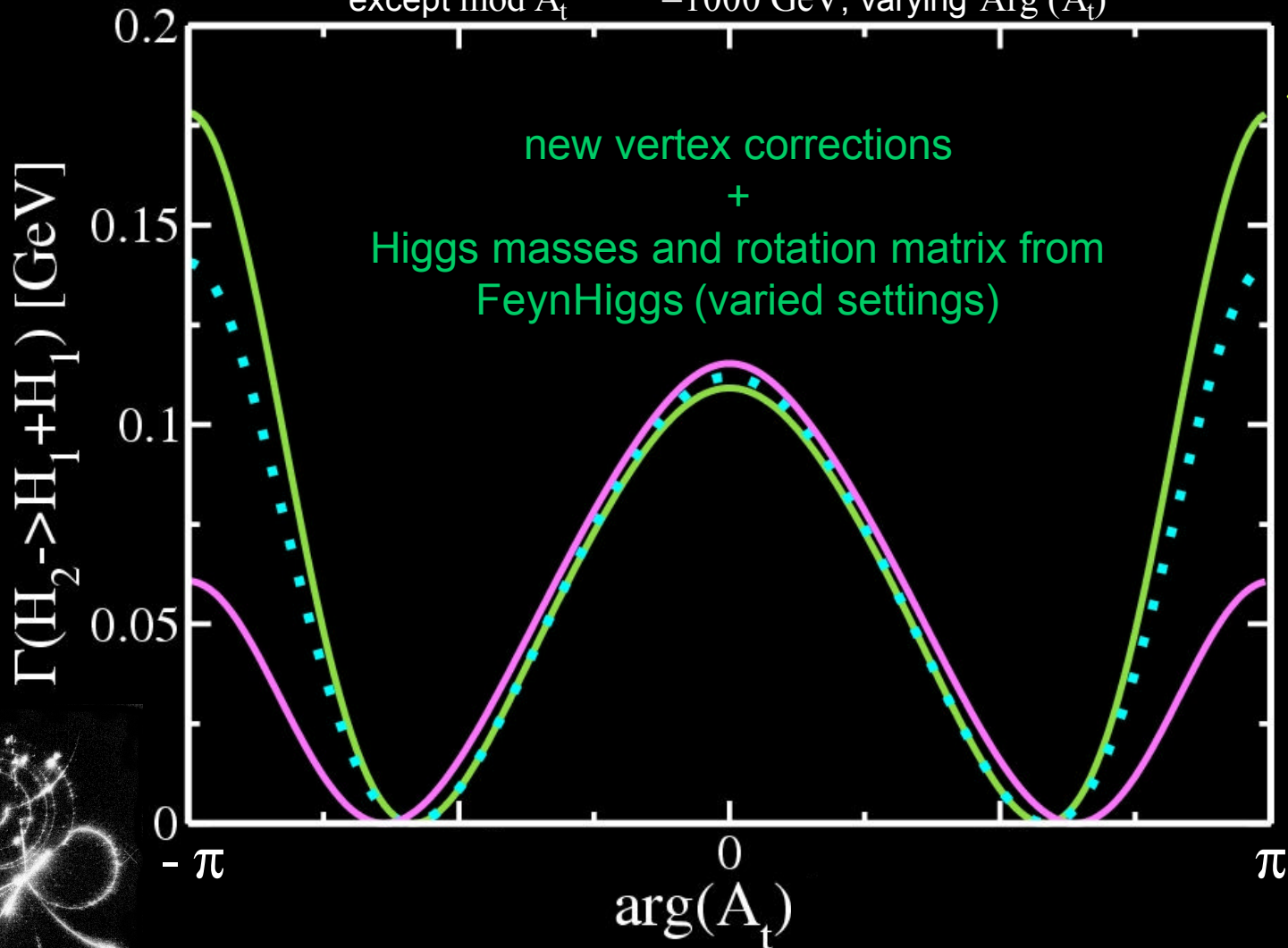
FeynHiggs settings:
top, stop,
bottom, sbottom,
 $p^2=0$,
1-loop



Preview of results



CPX with $\tan\beta = 10$, $M_{H^\pm} = 500$ GeV, $m_t = 172.7$ GeV
except mod $A_t^{\text{on-shell}} = 1000$ GeV, varying $\text{Arg}(A_t)$



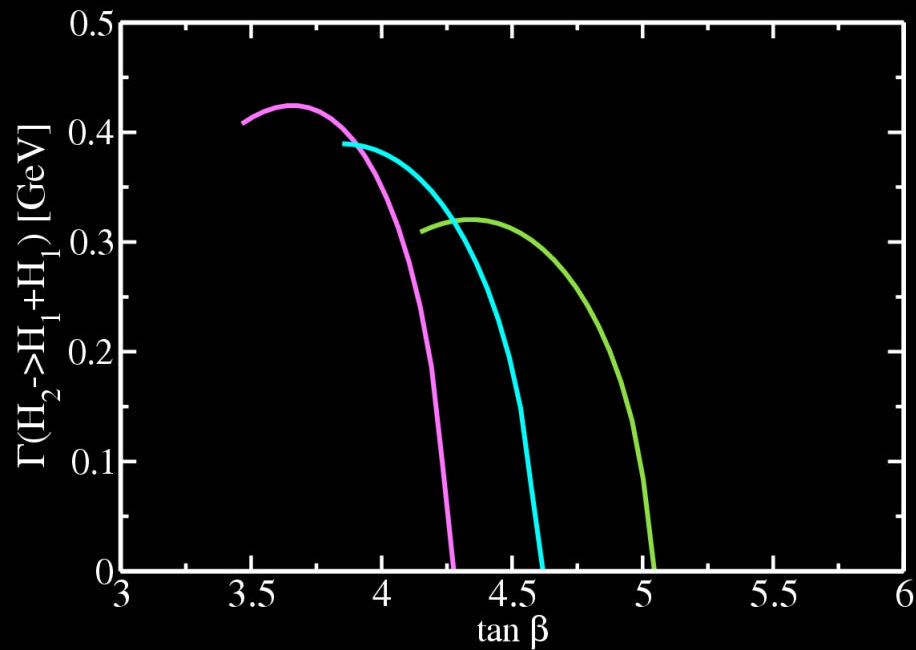
Full 1-loop,
some 2-loop

Full 1-loop

Top stop, bottom,
sbottom, $p^2=0$
(1-loop)



Near the hole in the LEP results...

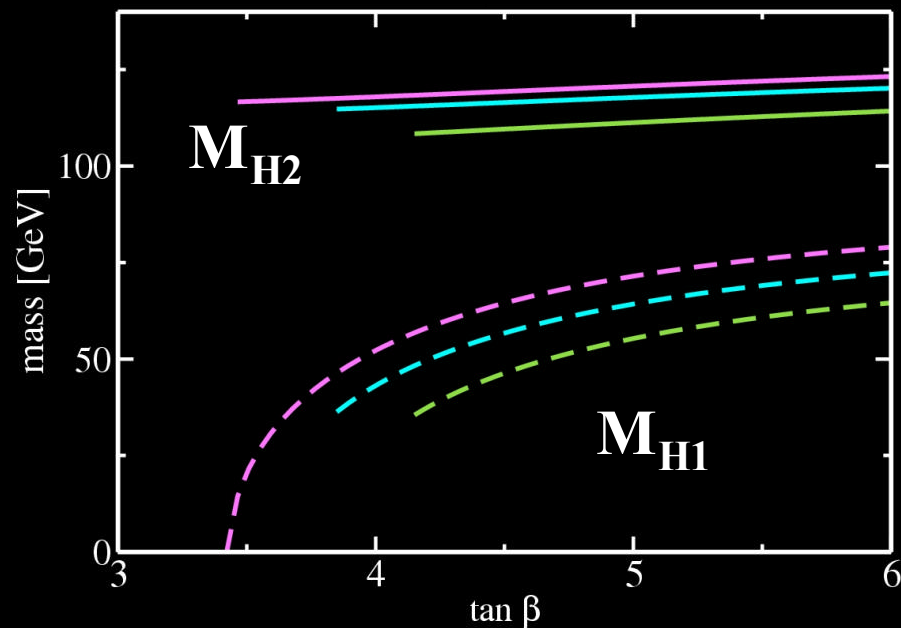


Settings for Higgs masses and rotation matrix from FeynHiggs:

Full 1-loop, some 2-loop

Full 1-loop

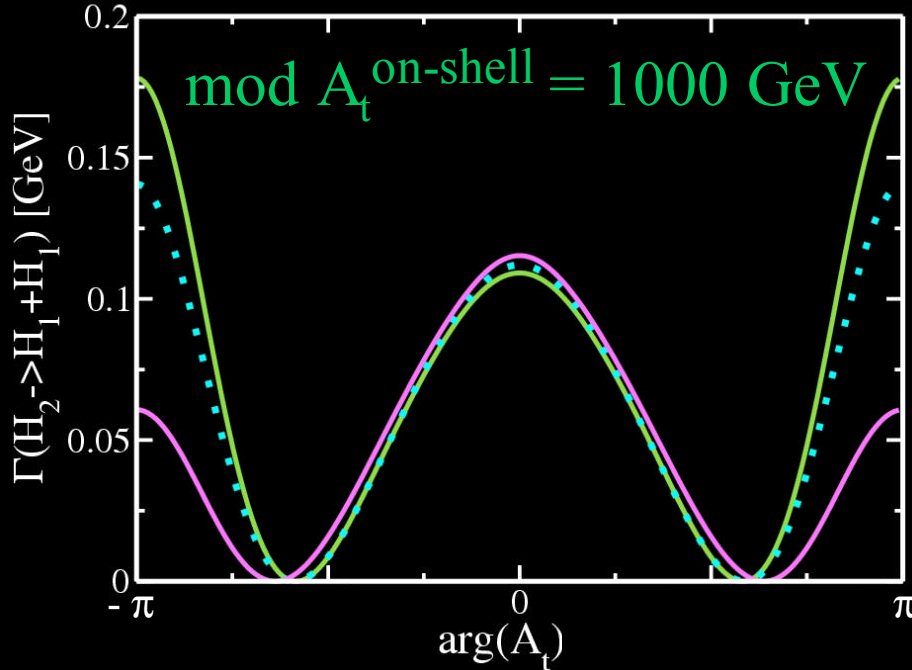
Top, stop, bottom, sbottom, $p^2=0$ (1-loop)



CPX with $M_{H^\pm} = 150$ GeV, $m_t = 172.7$ GeV
except mod $A_t^{\text{on-shell}} = 1000$ GeV



Changing A_t



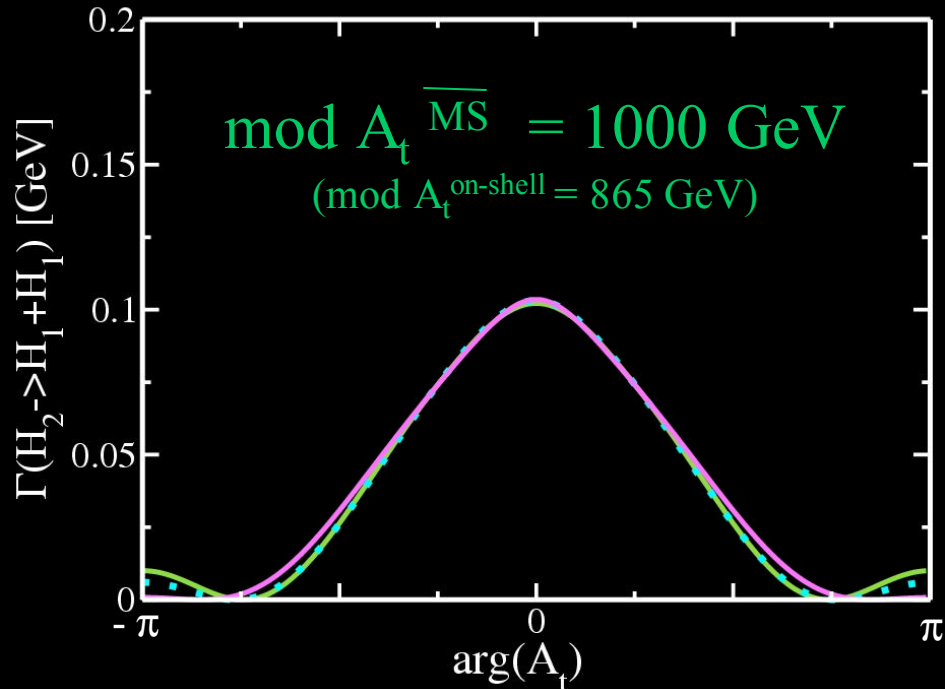
Settings for Higgs masses and rotation matrix from FeynHiggs:

Full 1-loop,
some 2-loop

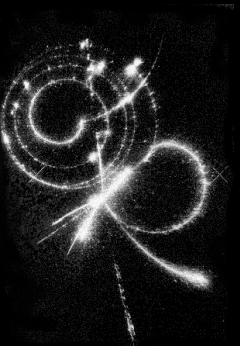
Full 1-loop

Top, stop, bottom,
sbottom, $p^2=0$ (1-loop)

using hep-ph / 0001002
(M. Carena, H.E. Haber, S. Heinemeyer,
W. Hollik, C.E.M. Wagner, G. Weiglein)



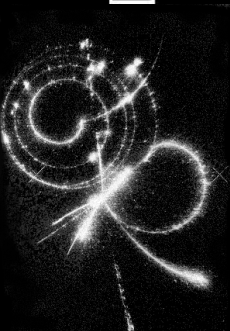
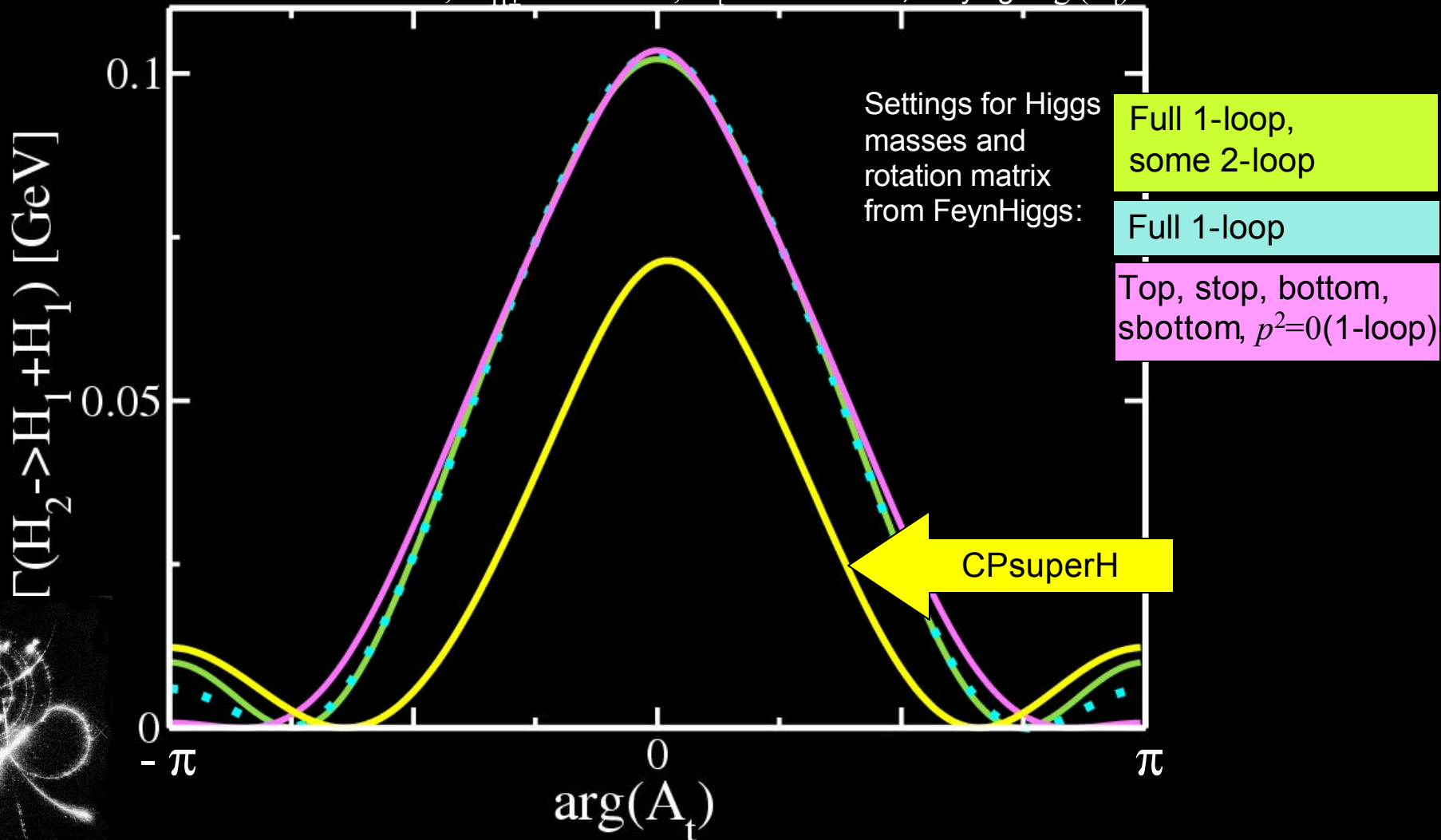
CPX with $\tan\beta = 10$, $M_{H_{\pm}} = 500 \text{ GeV}$, $m_t = 172.7 \text{ GeV}$, varying $\arg(A_t)$



Preview of results and comparison with CPsuperH



CPX with $\tan\beta = 10$, $M_{H^\pm} = 500$ GeV, $m_t = 172.7$ GeV, varying $\arg(A_t)$



Summary



- Improvement to the $H_2 \rightarrow H_1 H_1$ decay width in FeynHiggs

Next...

- full 1-loop vertex corrections
- refine conversion between different definitions of A_t
- comparison with CPsuperH

