

Flavor and new Physics

Andreas Weiler
(DESY)

Cargese 2012: Across the TeV frontier with the LHC

Lecture I

blackboard

Determining the CKM matrix

$$V_{\text{CKM}} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$$

Parameterisations

V_{CKM} 3x3 unitary matrix ($V^\dagger V = I$) with 1 phase.

PDG parameterisation

$$\begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta} & c_{23}c_{13} \end{pmatrix}$$

$$(c_{12} = \cos\theta_{12}, s_{12} = \sin\theta_{12}, \dots)$$

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$$(c_{12} = \cos\theta_{12}, s_{12} = \sin\theta_{12}, \dots)$$

From experiment we find a strong hierarchy:

$$s_{12} = 0.22 \gg s_{23} = \mathcal{O}(10^{-2}) \gg s_{13} = \mathcal{O}(10^{-3})$$

Reveals an almost diagonal ($\lambda = 0.22$) structure:

$$|V_{\text{CKM}}| = \begin{pmatrix} |V_{ud}| & |V_{us}| & |V_{ub}| \\ |V_{cd}| & |V_{cs}| & |V_{cb}| \\ |V_{td}| & |V_{ts}| & |V_{tb}| \end{pmatrix} \sim \begin{pmatrix} 1 & \lambda & \lambda^3 \\ \lambda & 1 & \lambda^2 \\ \lambda^3 & \lambda^2 & 1 \end{pmatrix}$$

Wolfenstein Parameterisation

Wolfenstein '83

$$\begin{pmatrix} 1 - \frac{1}{2}\lambda^2 & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda & 1 - \frac{1}{2}\lambda^2 & A\lambda^2 \\ A\lambda^3(1 - \rho - i\eta) & -A\lambda^2 & 1 \end{pmatrix} + \mathcal{O}(\lambda^4)$$

$$A, \rho, \eta \sim 0.8, 0.2, 0.3$$

η measures CPV

Wolfenstein Parameterisation

Wolfenstein '83

Choose parameters that make hierarchy obvious

$$s_{12} \equiv \lambda \quad s_{23} \equiv A\lambda^2 \quad s_{13} e^{-i\delta} \equiv A\lambda^3(\rho - i\eta)$$

$$\begin{pmatrix} 1 - \frac{1}{2}\lambda^2 & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda & 1 - \frac{1}{2}\lambda^2 & A\lambda^2 \\ A\lambda^3(1 - \rho - i\eta) & -A\lambda^2 & 1 \end{pmatrix} + \mathcal{O}(\lambda^4)$$

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Unitarity triangles

CKM Matrix is unitary $\mathbf{V}^\dagger \mathbf{V} = \mathbf{V} \mathbf{V}^\dagger = \mathbf{I}$

$$\mathbf{V}_{i1} (\mathbf{V}_{j1})^* + \mathbf{V}_{i2} (\mathbf{V}_{j2})^* + \mathbf{V}_{i3} (\mathbf{V}_{j3})^* = 0 \quad \text{for } i \neq j$$

\vec{a} \vec{b} \vec{c}

vectors in complex plane: $\vec{a} + \vec{b} + \vec{c} = 0$

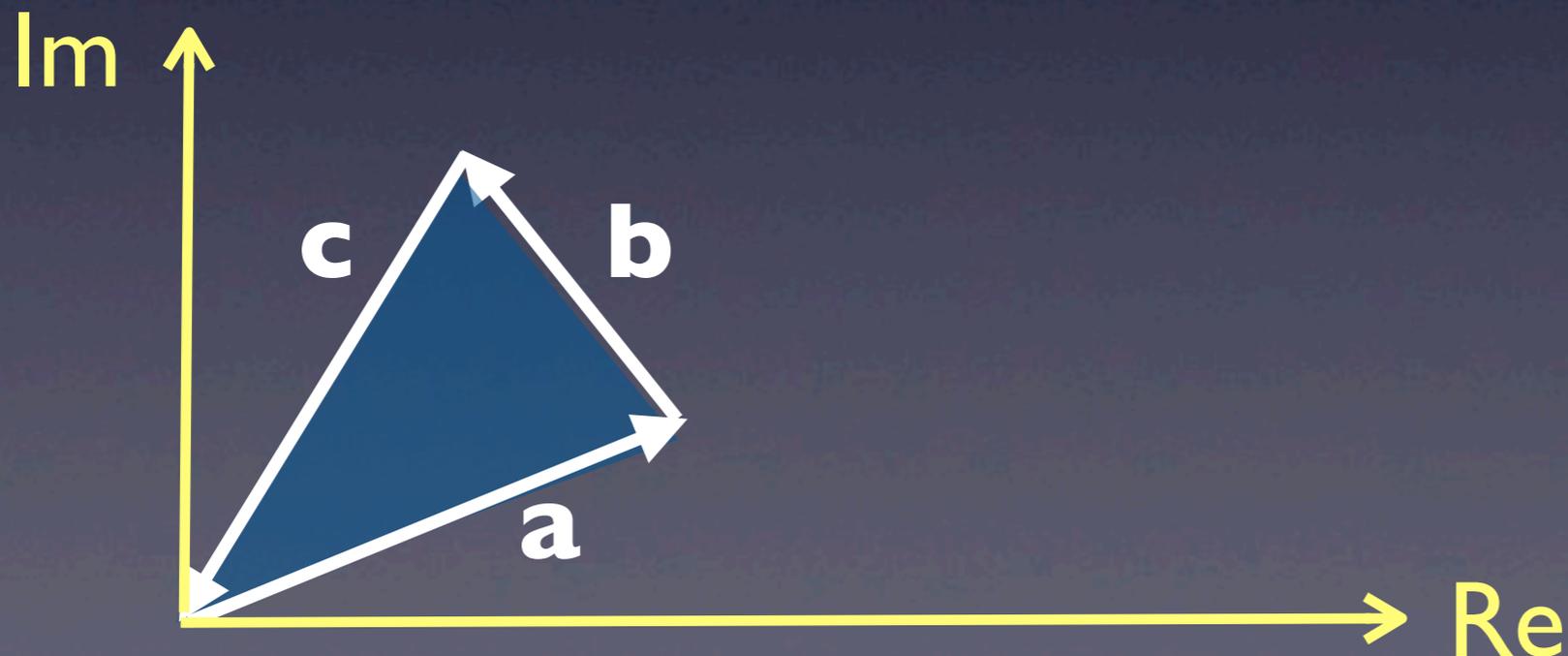
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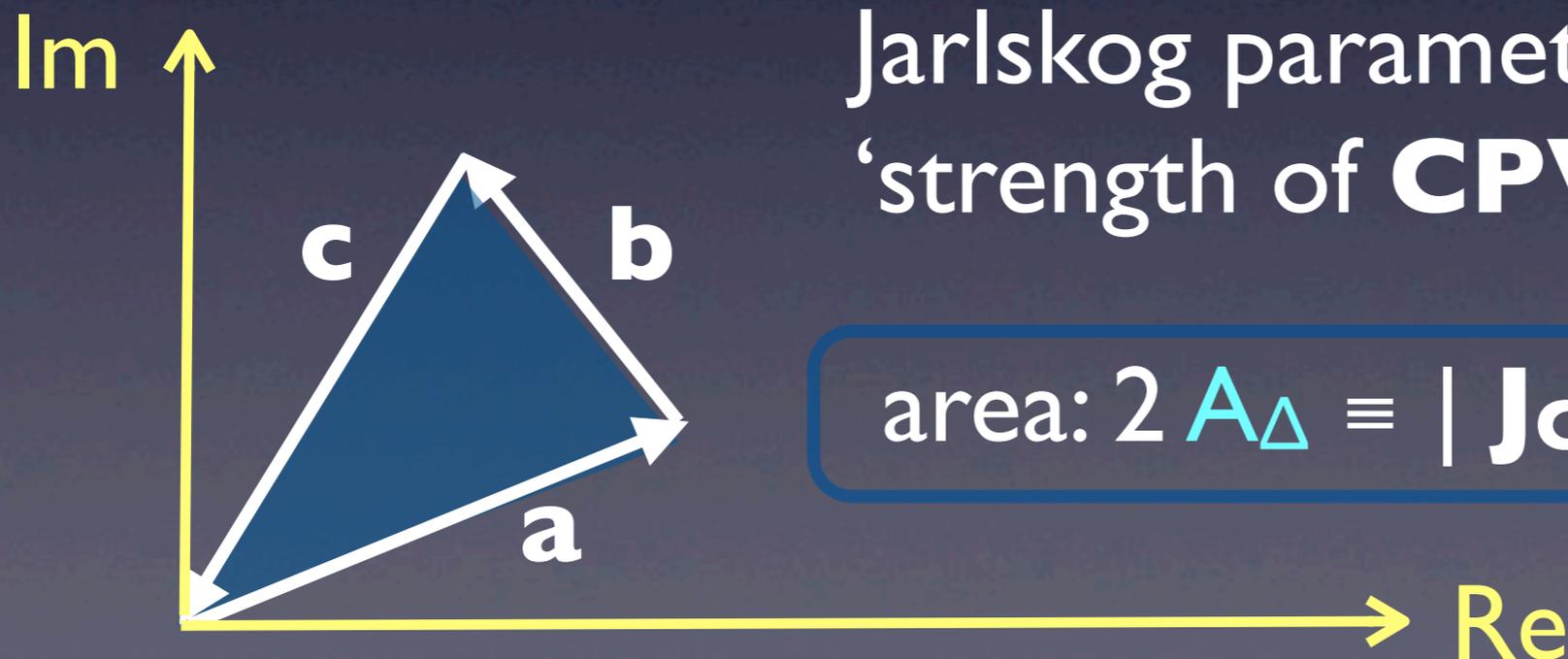
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vectors in complex plane: $\vec{a} + \vec{b} + \vec{c} = 0$



Jarlskog parameter \mathbf{J}_{CP} :
'strength of **CPV**'

area: $2 A_{\Delta} \equiv |\mathbf{J}_{\text{CP}}|$

The Unitarity triangle

$$\mathbf{V}_{i1} (\mathbf{V}_{j1})^* + \mathbf{V}_{i2} (\mathbf{V}_{j2})^* + \mathbf{V}_{i3} (\mathbf{V}_{j3})^* = \mathbf{0} \quad \text{for } i \neq j$$

of the **3x2 = 6** possible triangles pick one which has sides a,b,c of similar length (try!).

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The Unitarity triangle

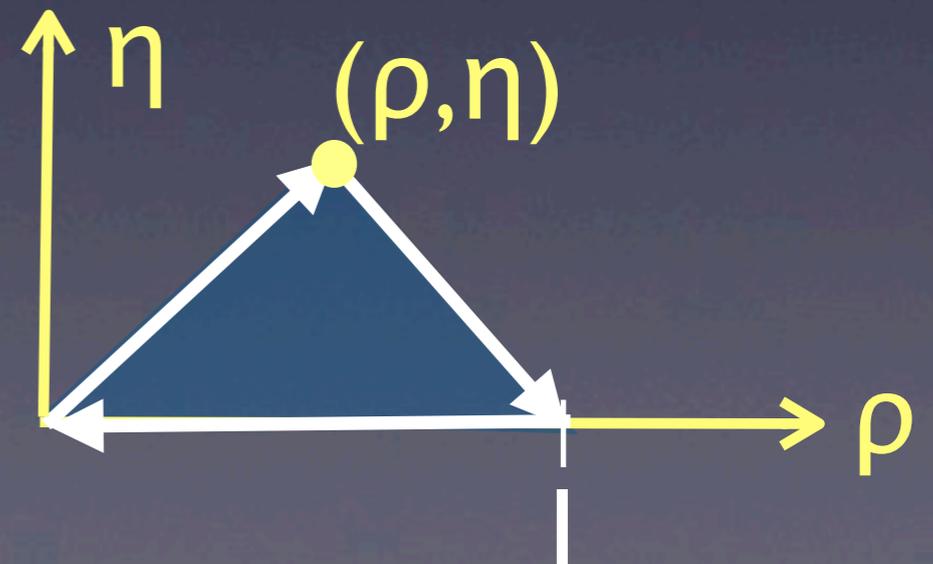
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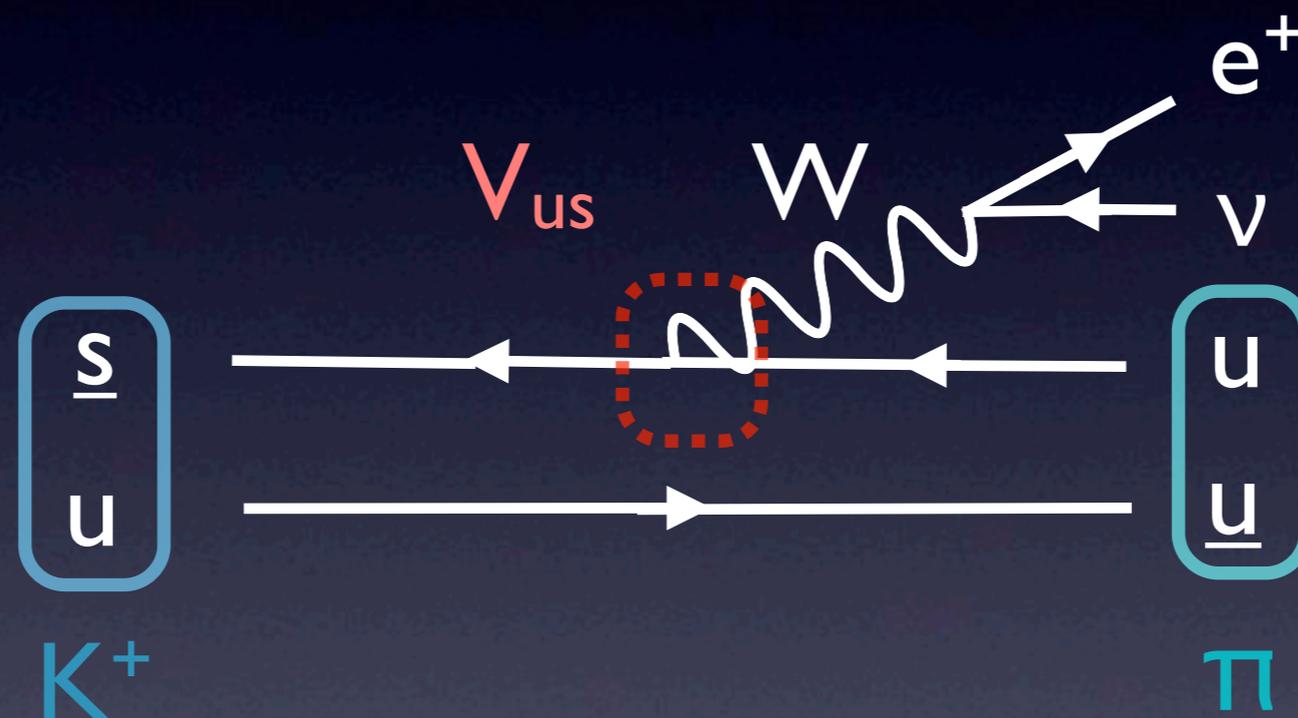
Divide by $\lambda^3 A = |\mathbf{V}_{cd} (\mathbf{V}_{cb})^*|$:

$$\rho + i\eta + (1 - \rho - i\eta) - 1 = 0$$



Direct Determination of the CKM matrix

I) $|V_{us}| = \lambda$ from $K^+ \rightarrow \pi^0 e^+ \nu_e$

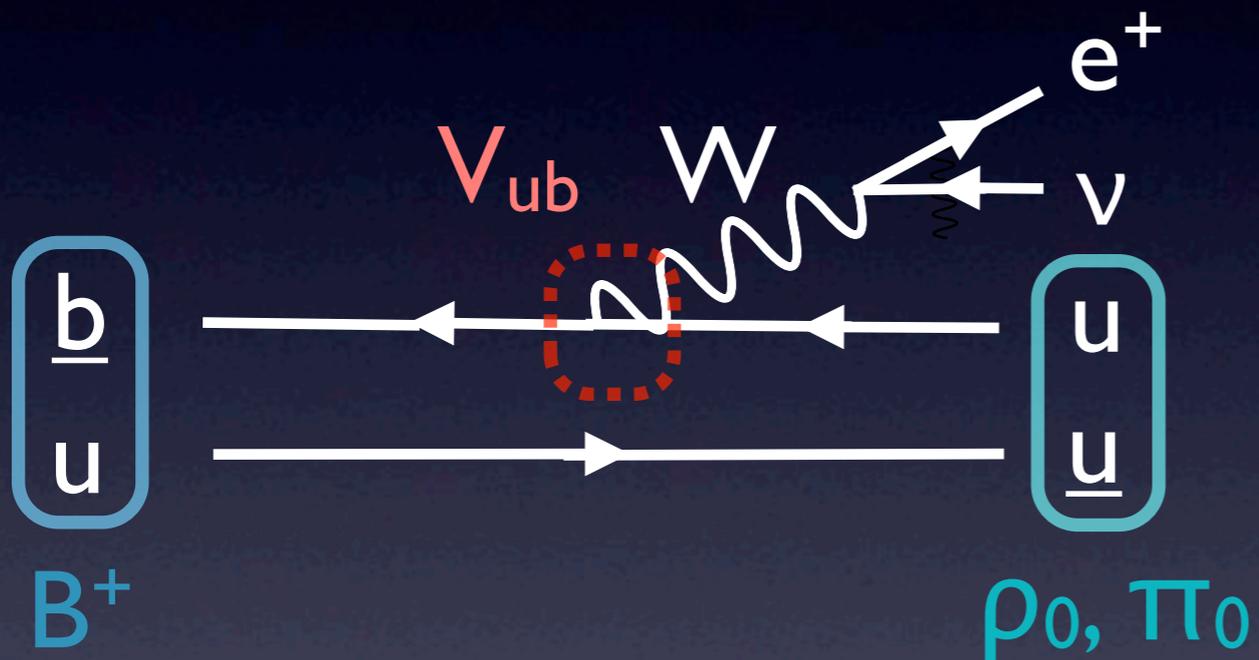


Result: $|V_{us}| = 0.2245 \pm 0.0016$

from: Moulson '07, KLOE, ISTRA+,
NA48, KTEV

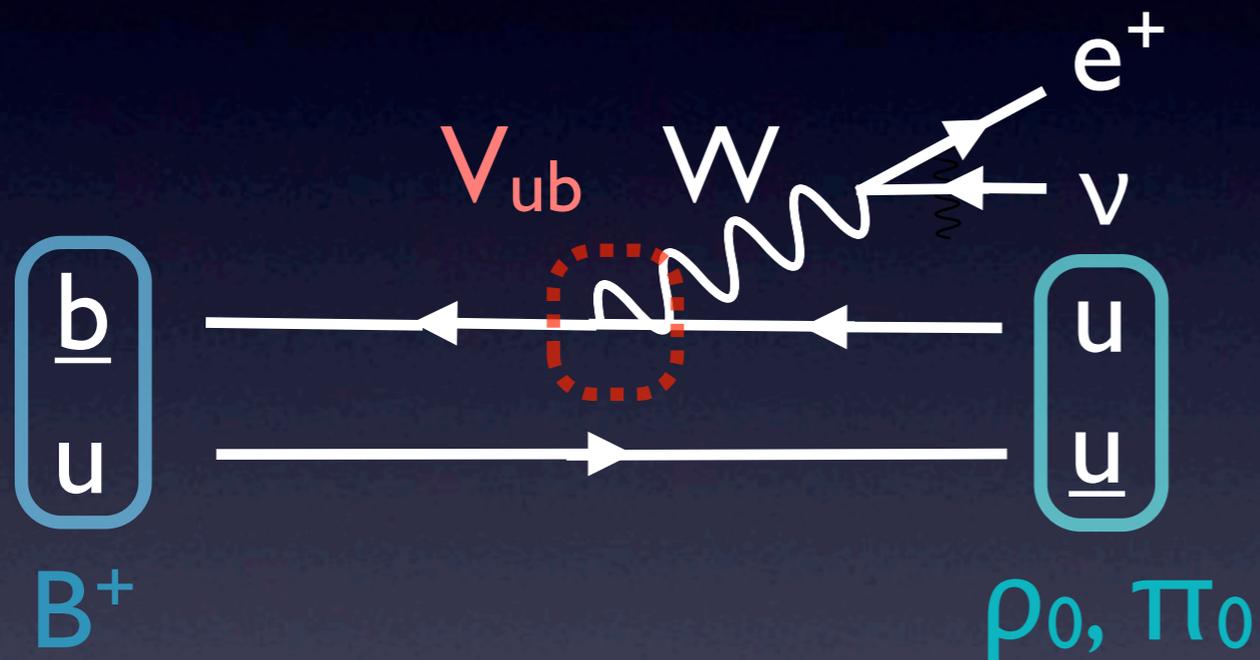
Side remark: Partonic vs Hadronic

Free quarks are easy...



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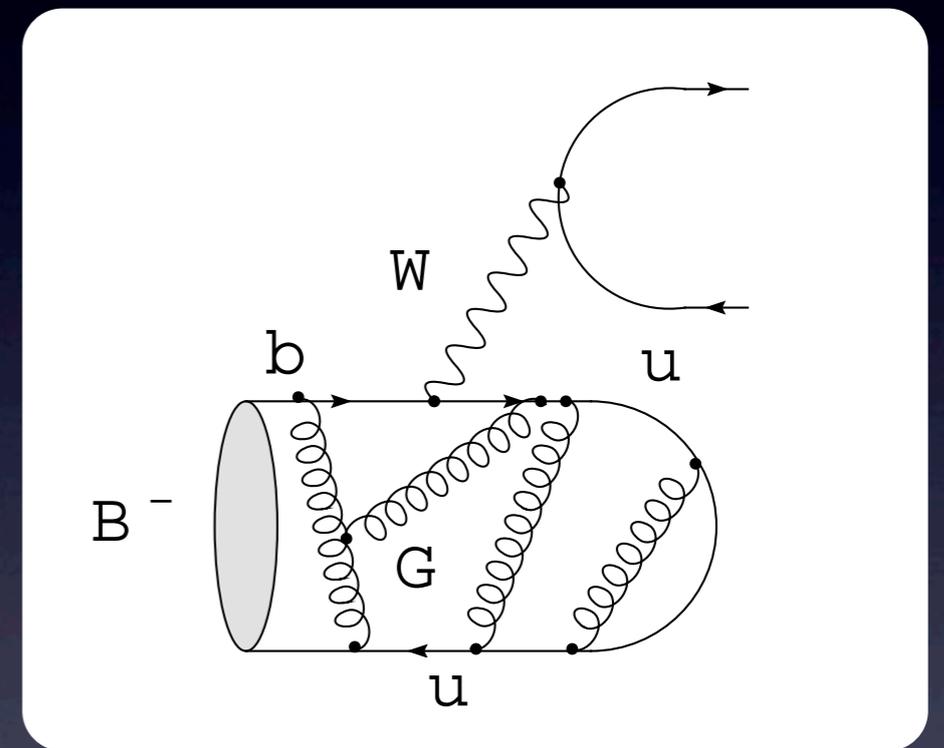
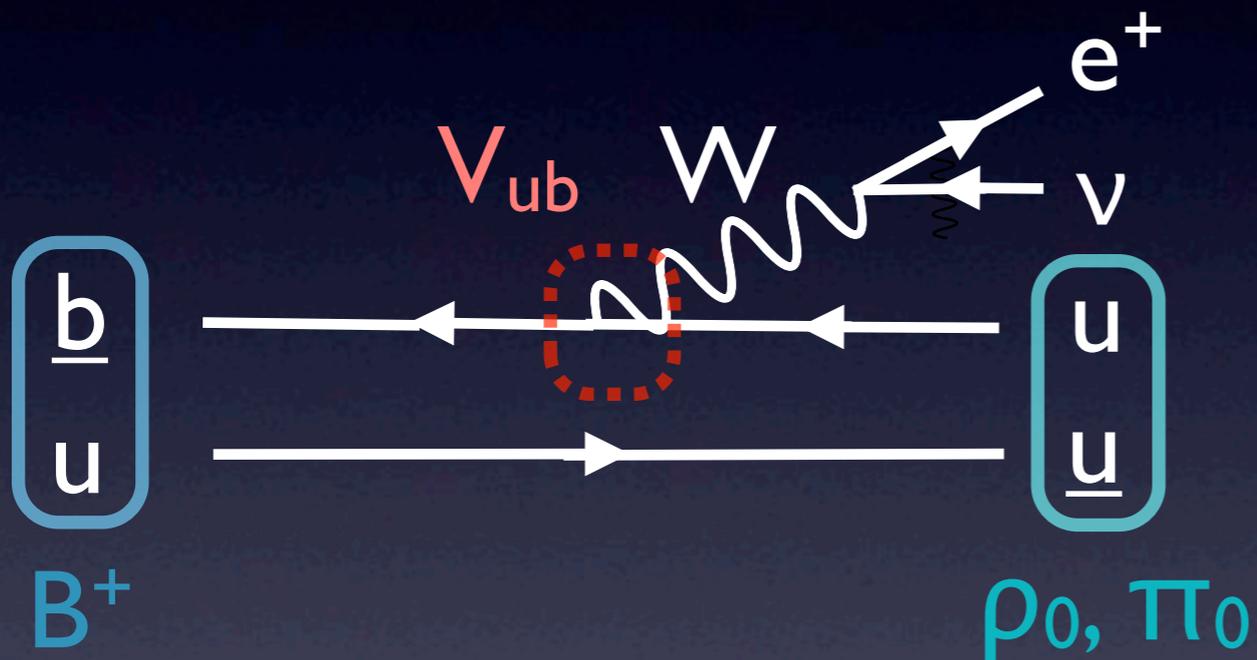


$$\Gamma(b \rightarrow ul\bar{\nu}) = \frac{G_F^2}{192\pi^2} |V_{ub}|^2 m_b^5$$

Side remark: Partonic vs Hadronic

Free quarks are easy...

Hadrons not so much

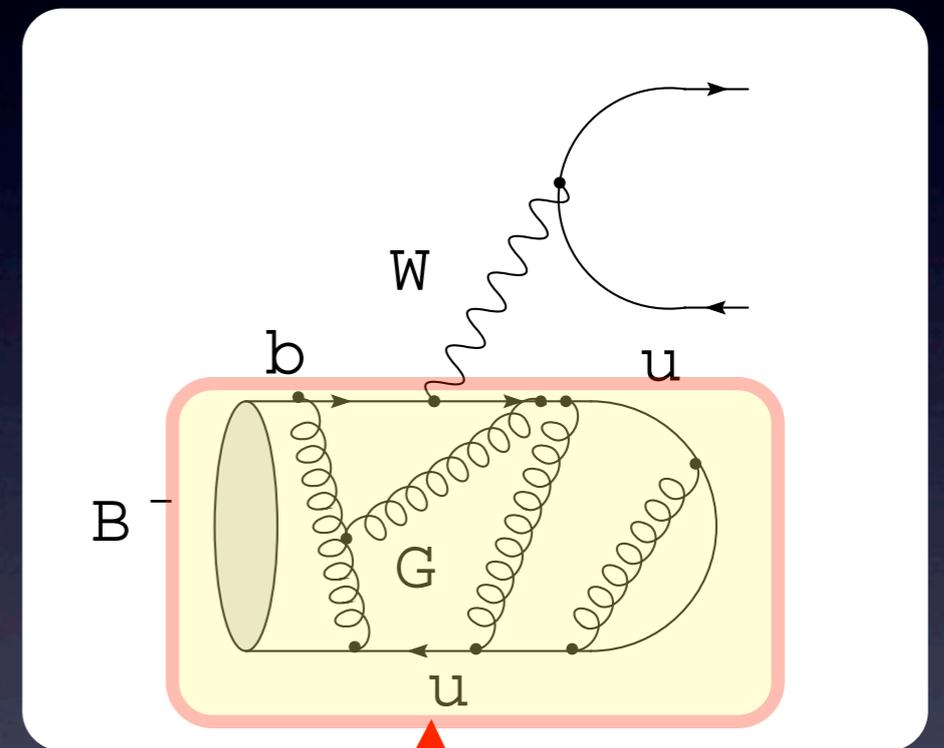
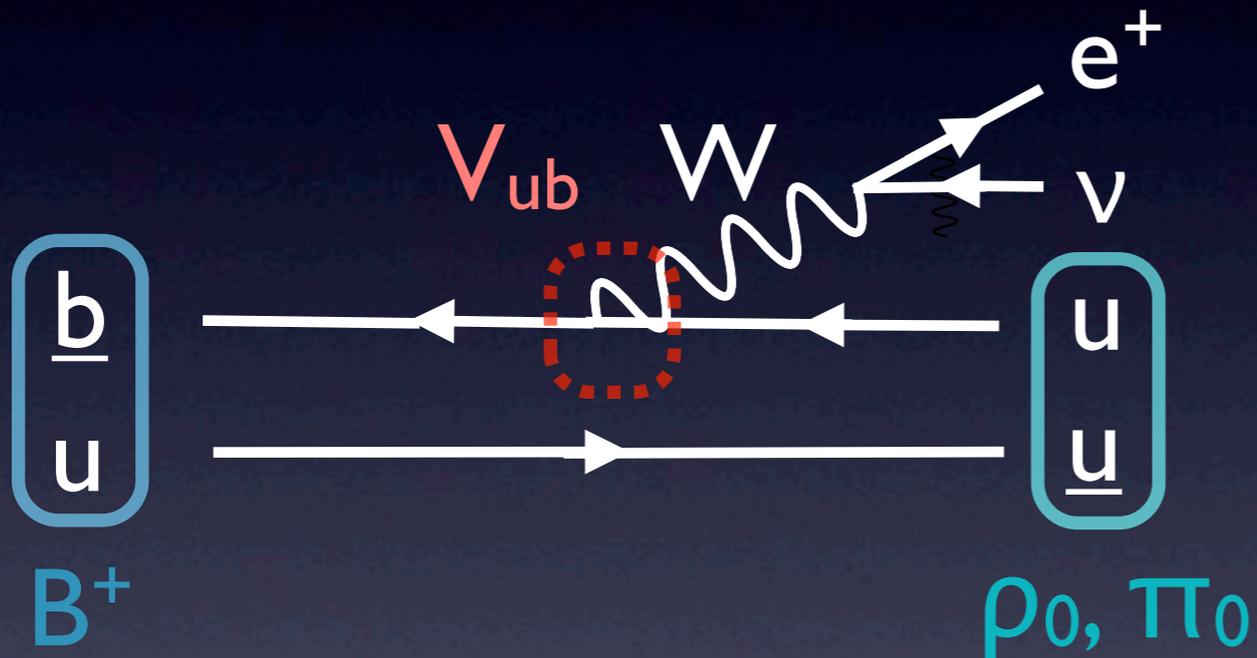


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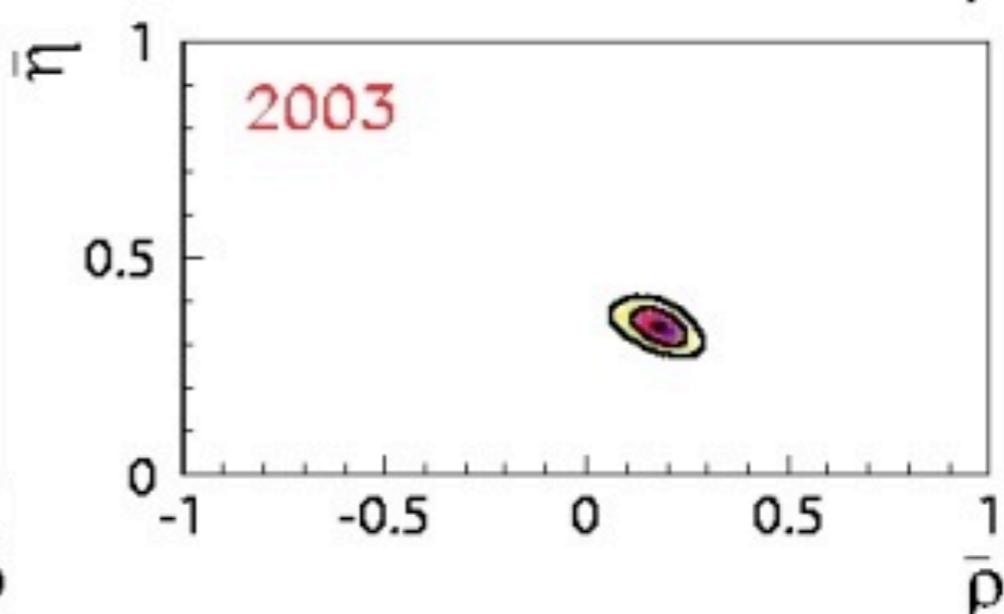
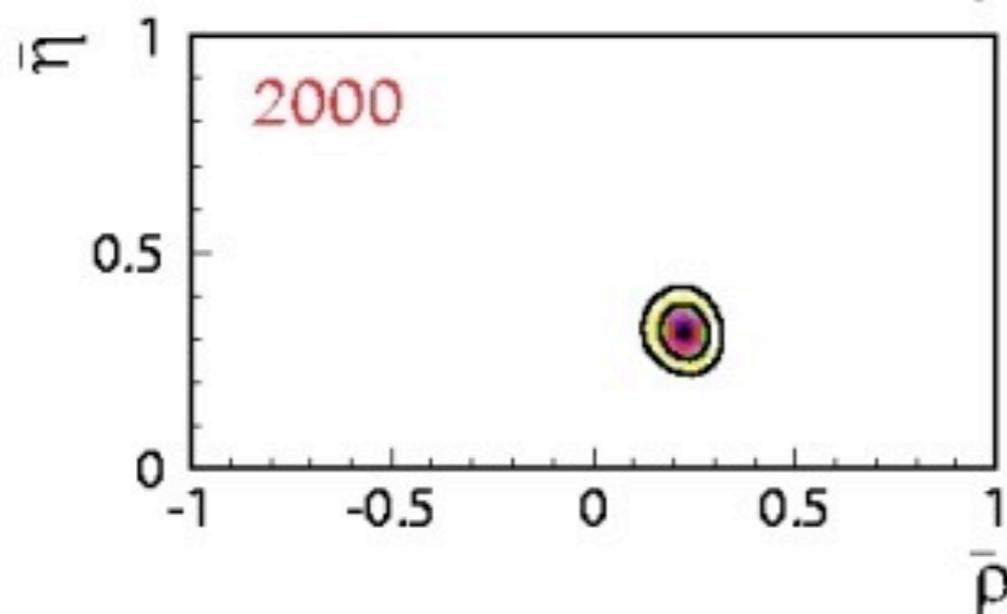
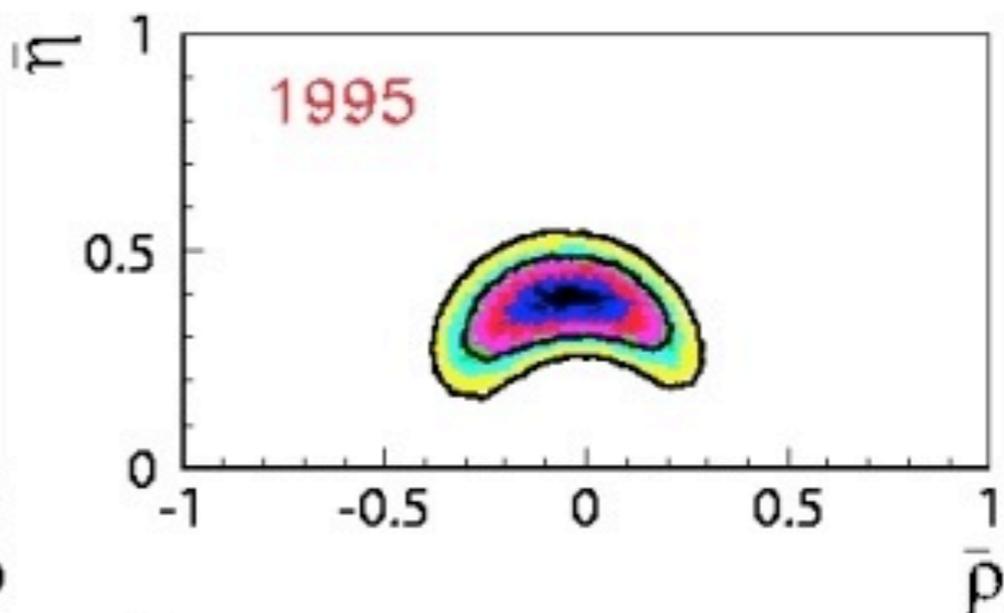
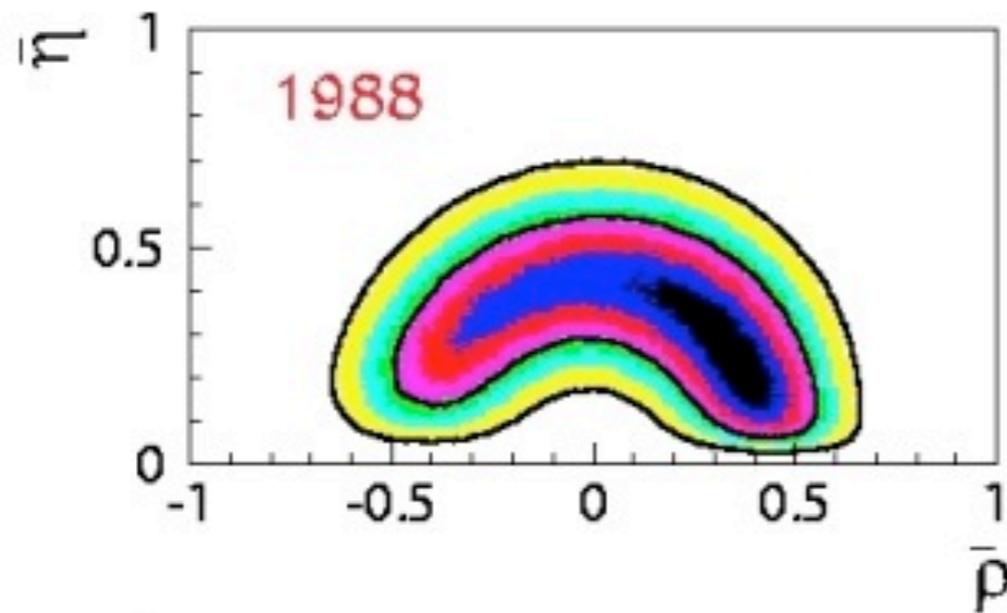
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hadronic matrix element

$$\langle 0 | \bar{u} \gamma^\mu (1 - \gamma_5) b | B^- \rangle$$

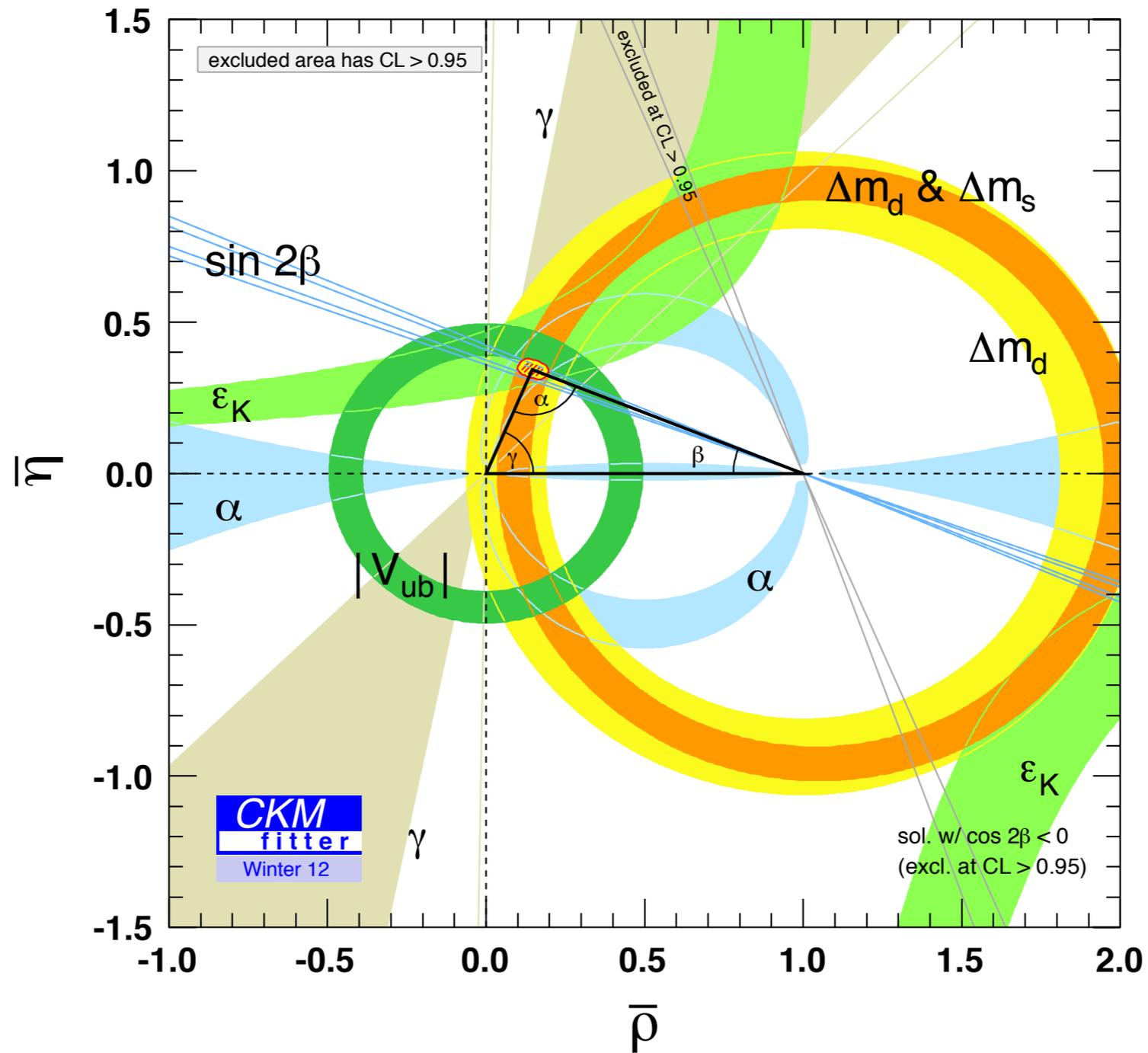
lattice QCD or inclusive decays



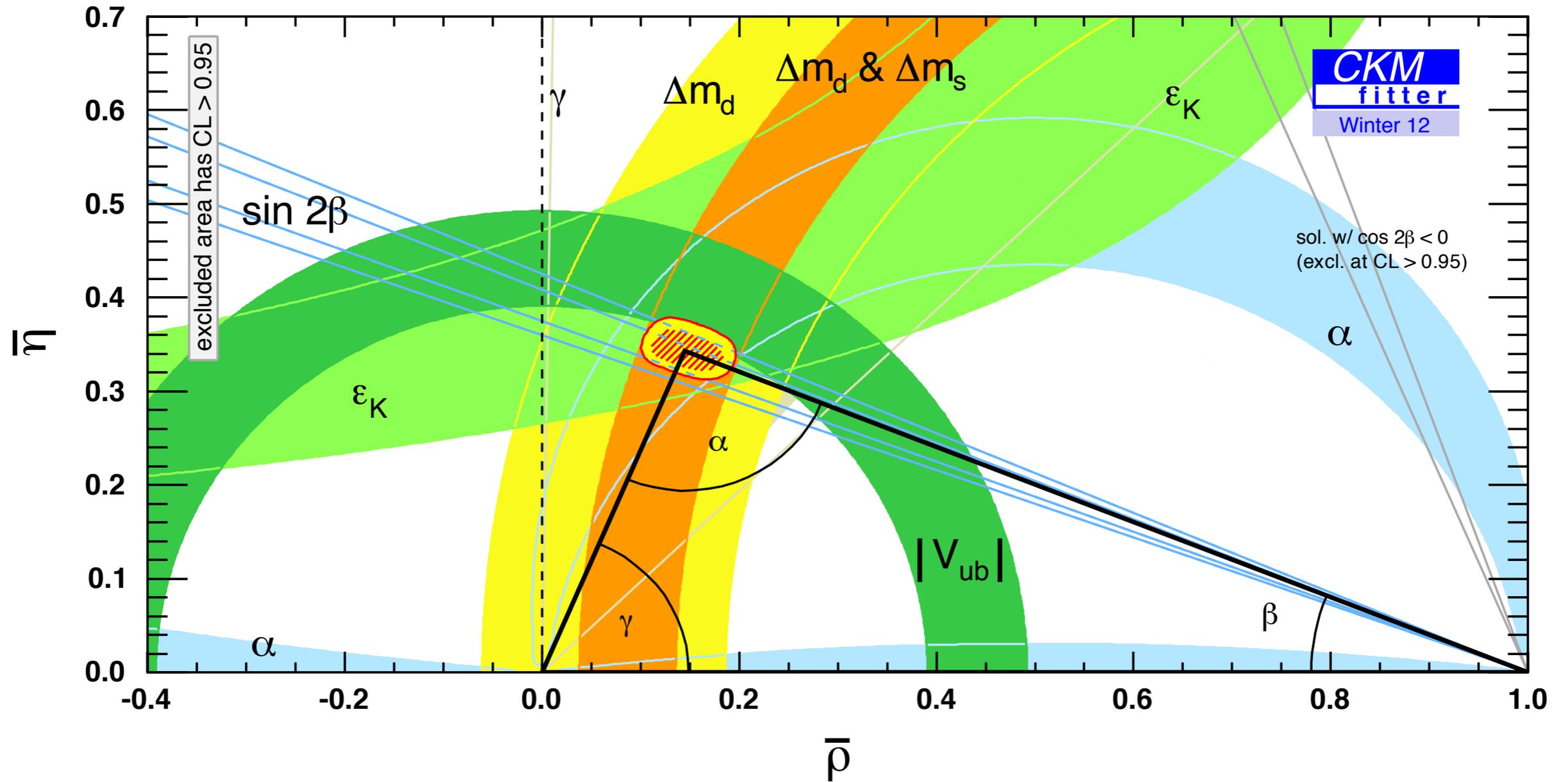


source: utfit.org

unitarity triangle '12



Zoom



- The SM is the dominant source for the observed flavor and CP violation.



Flavor changing neutral currents

blackboard

FCNC limits

UTfit 08, Isidori, Perez, Nir '10

Operator	Bounds on Λ in TeV ($c_{ij} = 1$)		Observables
	Re	Im	
$(\bar{s}_L \gamma^\mu d_L)^2$	9.8×10^2	1.6×10^4	$\Delta m_K; \epsilon_K$
$(\bar{s}_R d_L)(\bar{s}_L d_R)$	1.8×10^4	3.2×10^5	$\Delta m_K; \epsilon_K$
$(\bar{c}_L \gamma^\mu u_L)^2$	1.2×10^3	2.9×10^3	$\Delta m_D; q/p , \phi_D$
$(\bar{c}_R u_L)(\bar{c}_L u_R)$	6.2×10^3	1.5×10^4	$\Delta m_D; q/p , \phi_D$
$(\bar{b}_L \gamma^\mu d_L)^2$	5.1×10^2	9.3×10^2	$\Delta m_{B_d}; S_{\psi K_S}$
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Very strong suppression! New flavor violation must either approximately (exactly?) follow SM structure...

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Very strong suppression! New flavor violation must either **approximately (exactly?) follow SM structure...**

... or exist only at **very high scales ($10^2 - 10^5$ TeV)**

Lecture 2

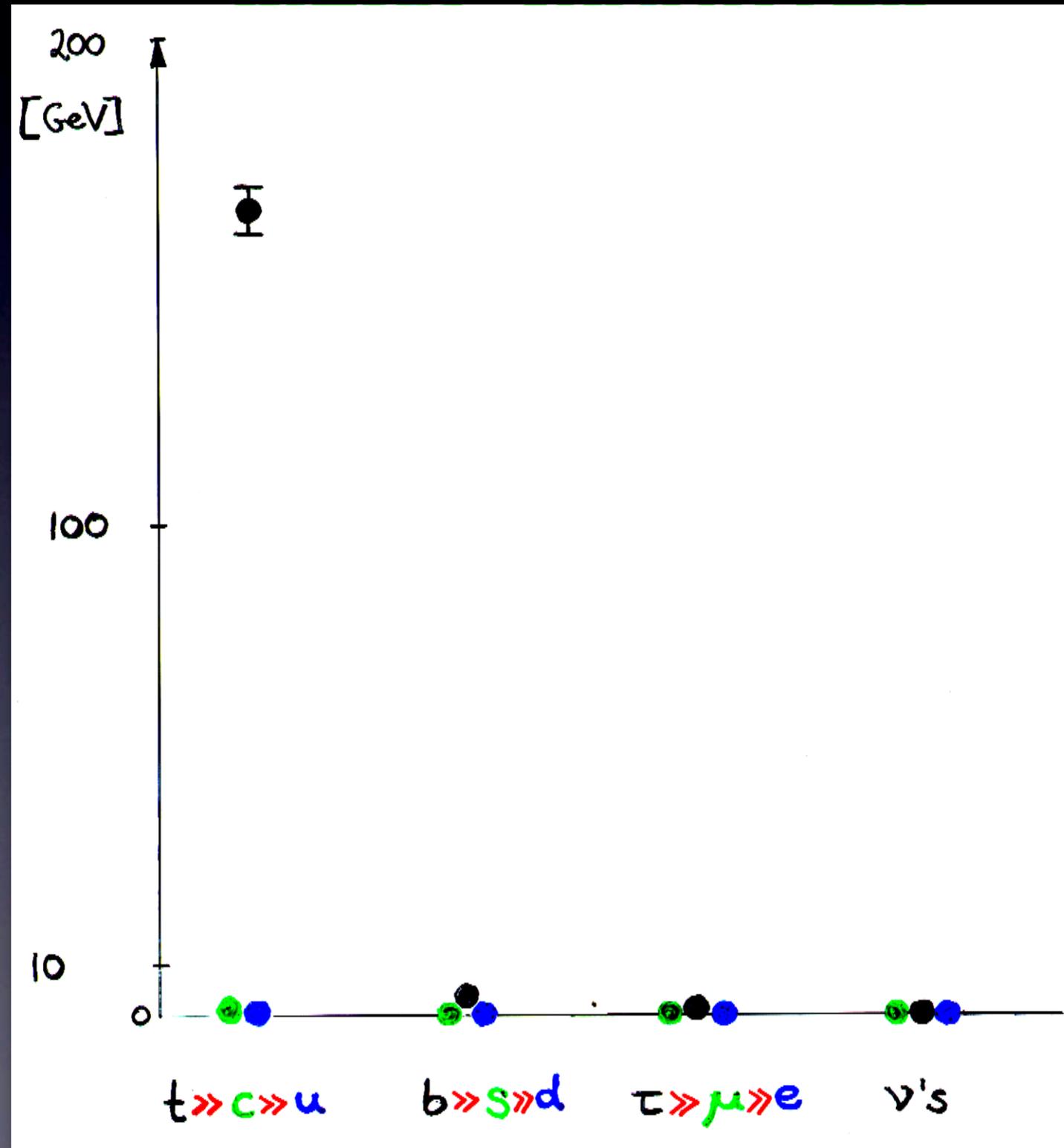
The flavor puzzle

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UTfit 08, Isidori, Perez, Nir '10

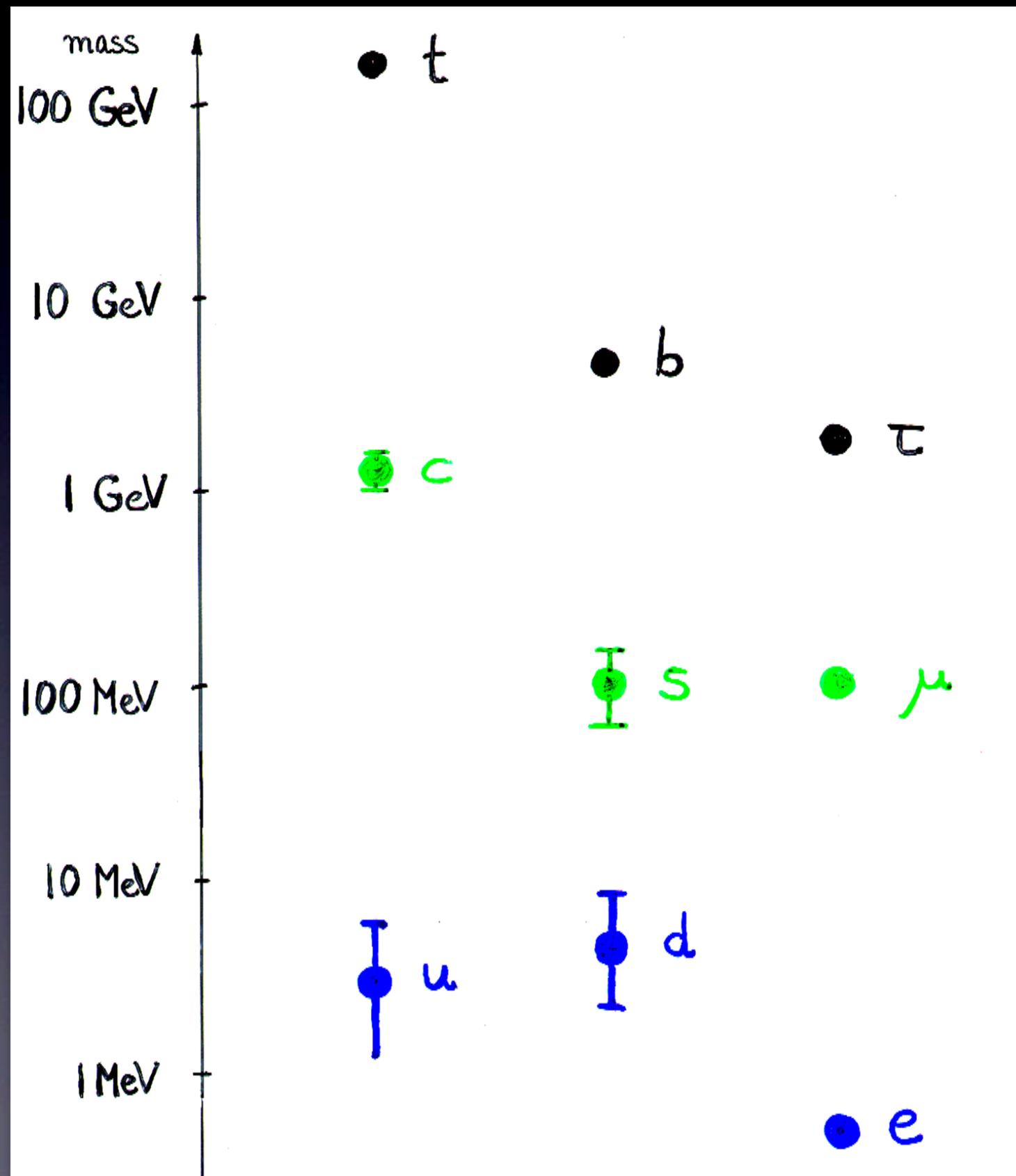
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Quark and Lepton mass hierarchy



LHC view

Masses on a Log-scale



The SM flavor puzzle

$$Y_D \approx \text{diag} (2 \cdot 10^{-5} \quad 0.0005 \quad 0.02)$$

$$Y_U \approx \begin{pmatrix} 6 \cdot 10^{-6} & -0.001 & 0.008 + 0.004i \\ 1 \cdot 10^{-6} & 0.004 & -0.04 + 0.001i \\ 8 \cdot 10^{-9} + 2 \cdot 10^{-8}i & 0.0002 & 0.98 \end{pmatrix}$$

Why this structure?

Other dimensionless parameters of the SM:

$$g_s \sim 1, \quad g \sim 0.6, \quad g' \sim 0.3, \quad \lambda_{\text{Higgs}} \sim 1,$$

The SM flavor puzzle

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Why this structure?

Other dimensionless parameters of the SM:

$$g_s \sim 1, \quad g \sim 0.6, \quad g' \sim 0.3, \quad \lambda_{\text{Higgs}} \sim 1, \quad |\theta| < 10^{-9}$$

Log(SM flavor puzzle)

$$-\log |Y_D| \approx \text{diag} (11 \quad 8 \quad 4)$$

$$-\log |Y_U| \approx \begin{pmatrix} 12 & 7 & 5 \\ 14 & 6 & 3 \\ 18 & 9 & 0 \end{pmatrix}$$

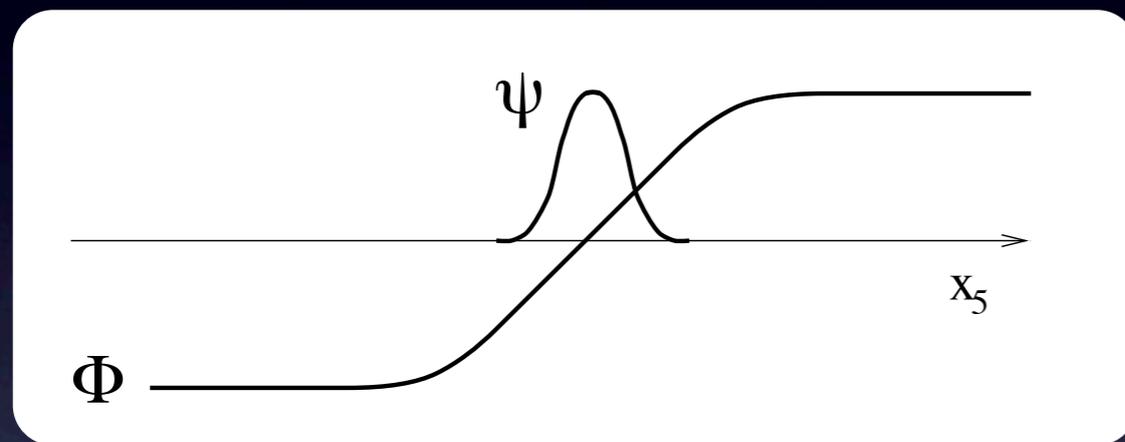
If $Y = e^{-\Delta}$, then the Δ don't look crazy.

Flavor from Geometry

Hierarchies w/o Symmetries

Arkani-Hamed, Schmaltz

SM on thick brane & domain wall \Rightarrow chiral localization



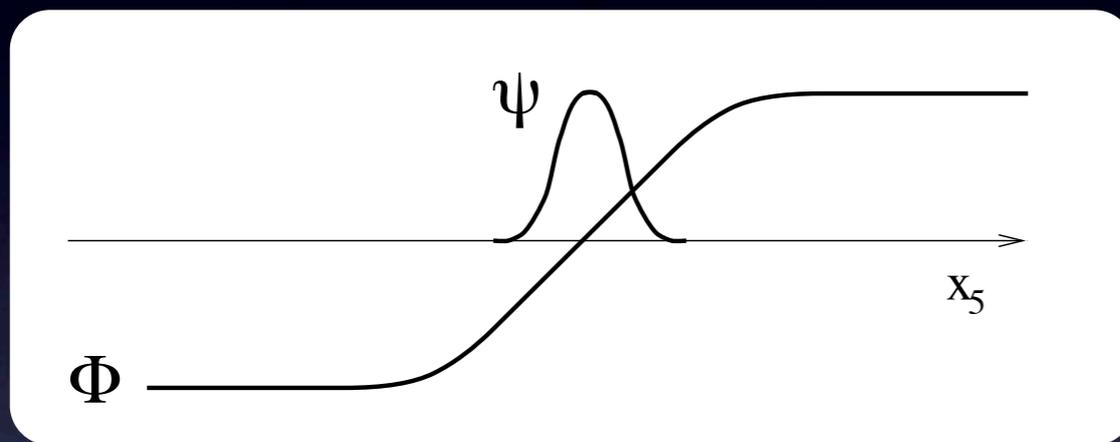
$$\mathcal{S} = \int d^5x \sum_{i,j} \bar{\Psi}_i [i \not{\partial}_5 + \lambda \Phi(x_5) - m]_{ij} \Psi_j$$

$$\Psi = \begin{pmatrix} \Psi_L \\ \Psi_R \end{pmatrix} = \begin{pmatrix} \psi_L^0 \\ 0 \end{pmatrix} + \text{KK modes}$$

Hierarchies w/o Symmetries

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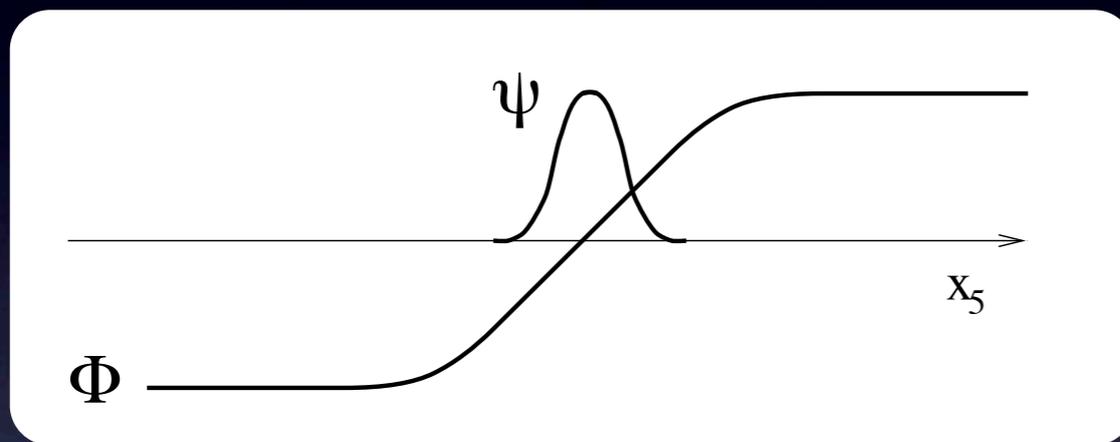
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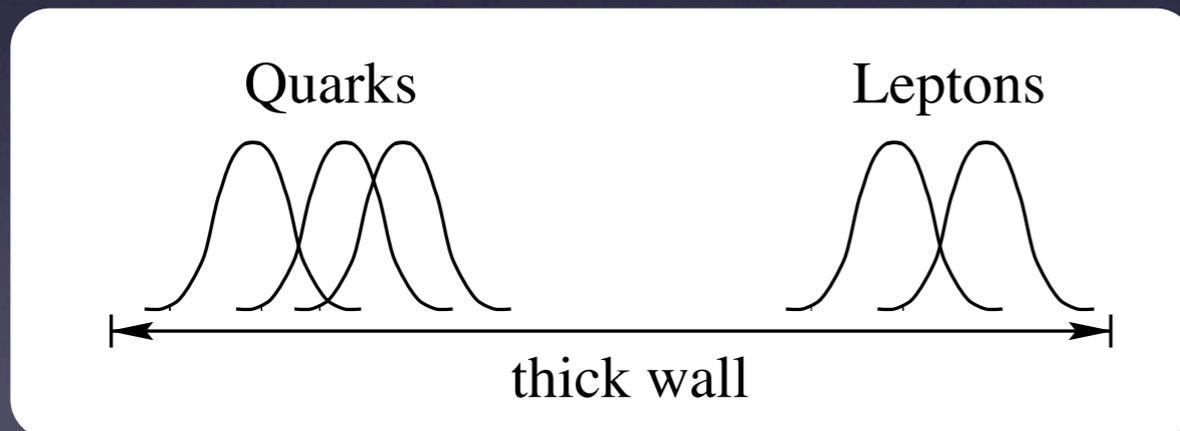
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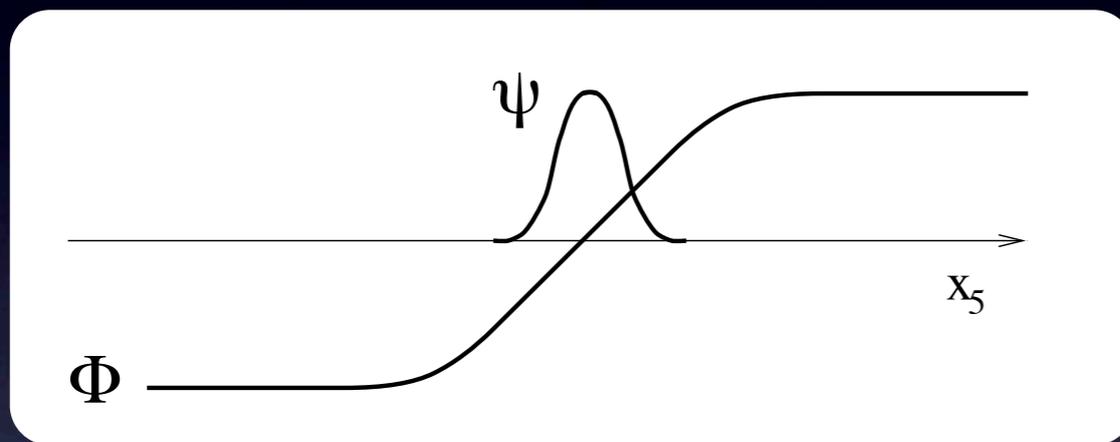
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Hierarchies w/o Symmetries

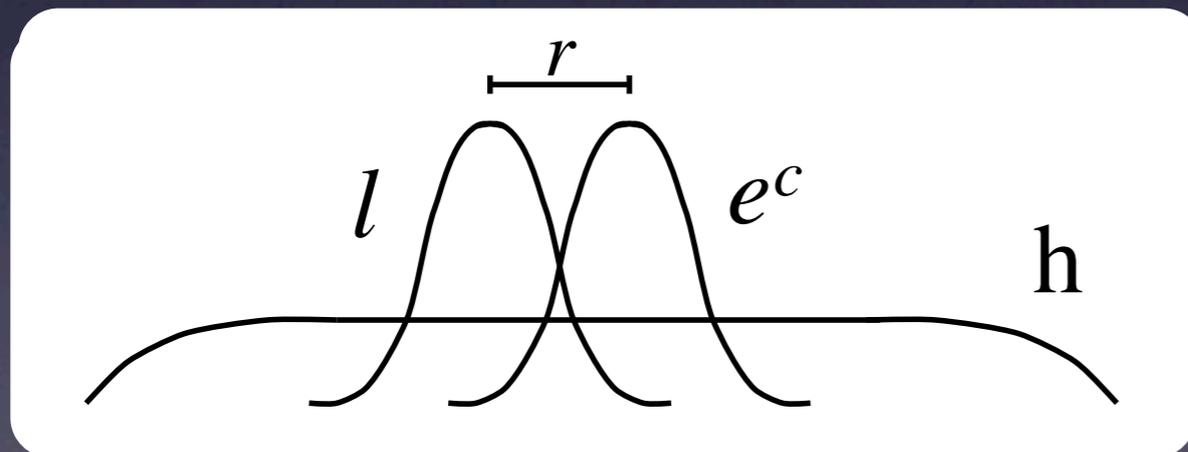
Arkani-Hamed, Schmaltz

SM on thick brane & domain wall \Rightarrow chiral localization



$$\mathcal{S} = \int d^5x \sum_{i,j} \bar{\Psi}_i [i \not{\partial}_5 + \lambda \Phi(x_5) - m]_{ij} \Psi_j$$

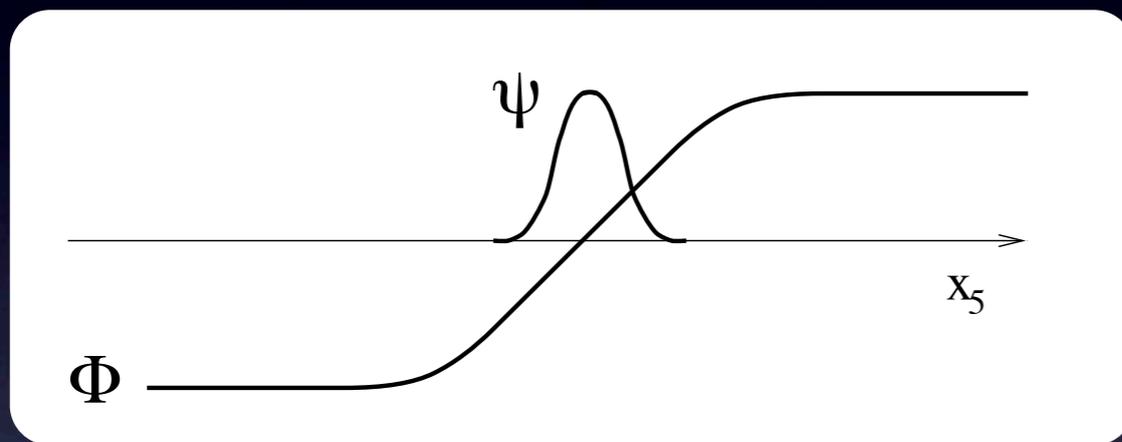
$$\Psi = \begin{pmatrix} \Psi_L \\ \Psi_R \end{pmatrix} = \begin{pmatrix} \psi_L^0 \\ 0 \end{pmatrix} + \text{KK modes}$$



Hierarchies w/o Symmetries

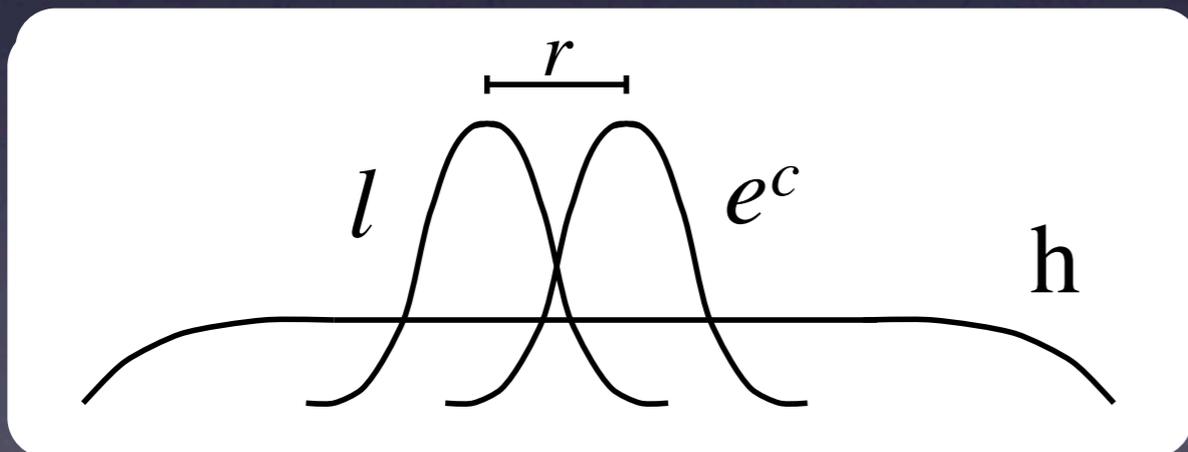
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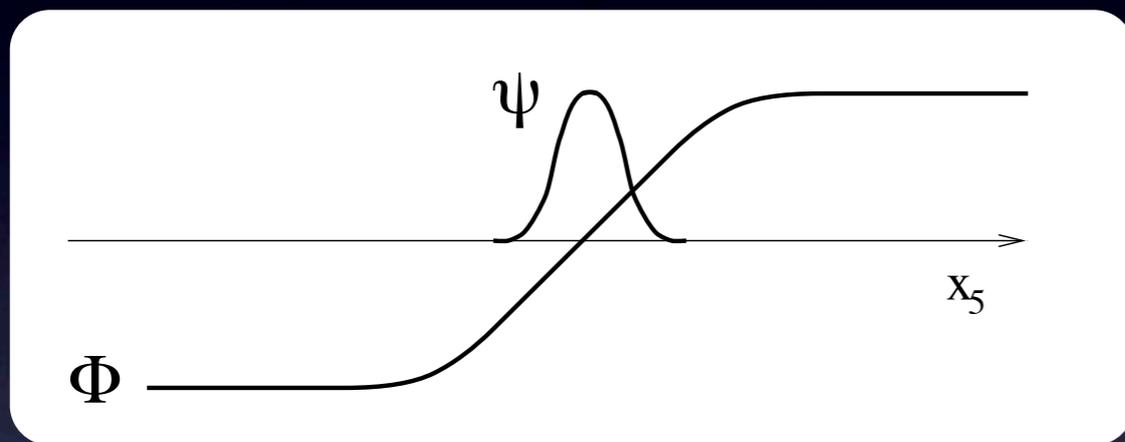


$$\int dx_5 \phi_l(x_5) \phi_{ec}(x_5) = \frac{\sqrt{2}\mu}{\sqrt{\pi}} \int dx_5 e^{-\mu^2 x_5^2} e^{-\mu^2 (x_5 - r)^2} = e^{-\mu^2 r^2 / 2}$$

Hierarchies w/o Symmetries

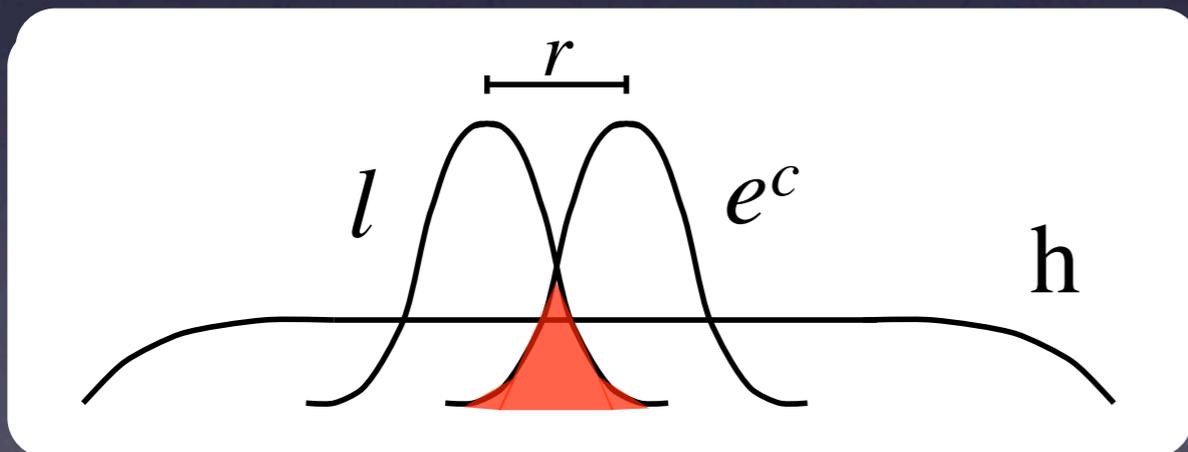
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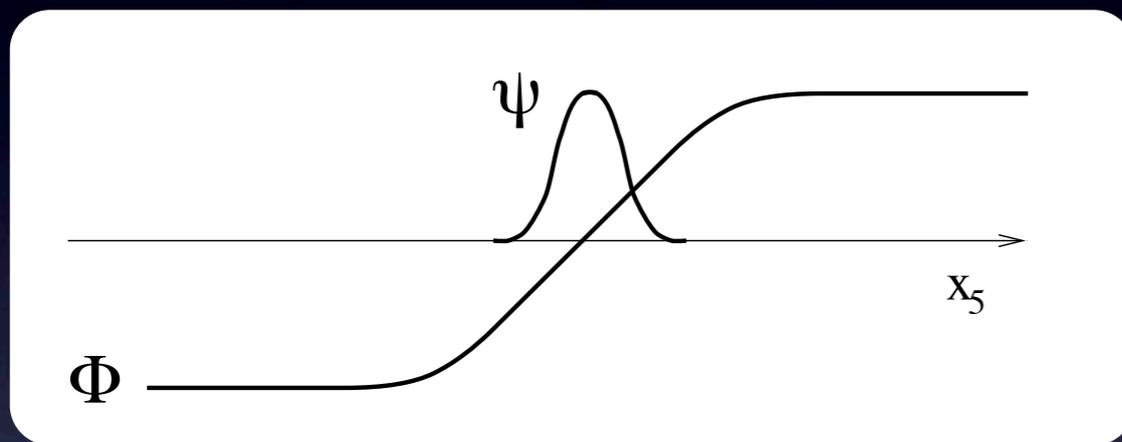


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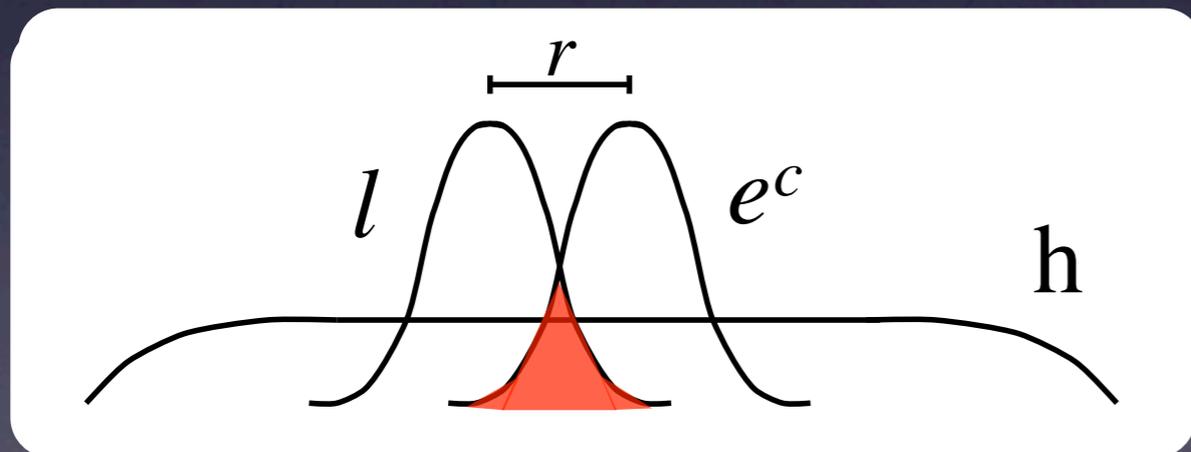
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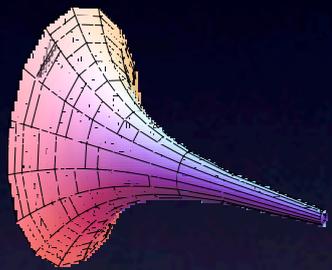
Log(flavor hierarchy)!

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$$ds^2 = dx_\mu dx_\nu - dy^2$$



Randall, Sundrum

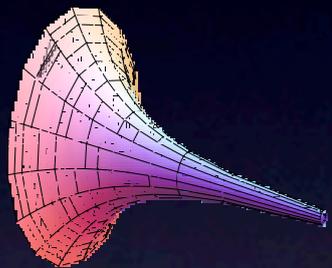


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Randall, Sundrum



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✓ AdS/CFT intuition: reappraisal of strong EW symmetry breaking (composite Higgs, technicolor,...)

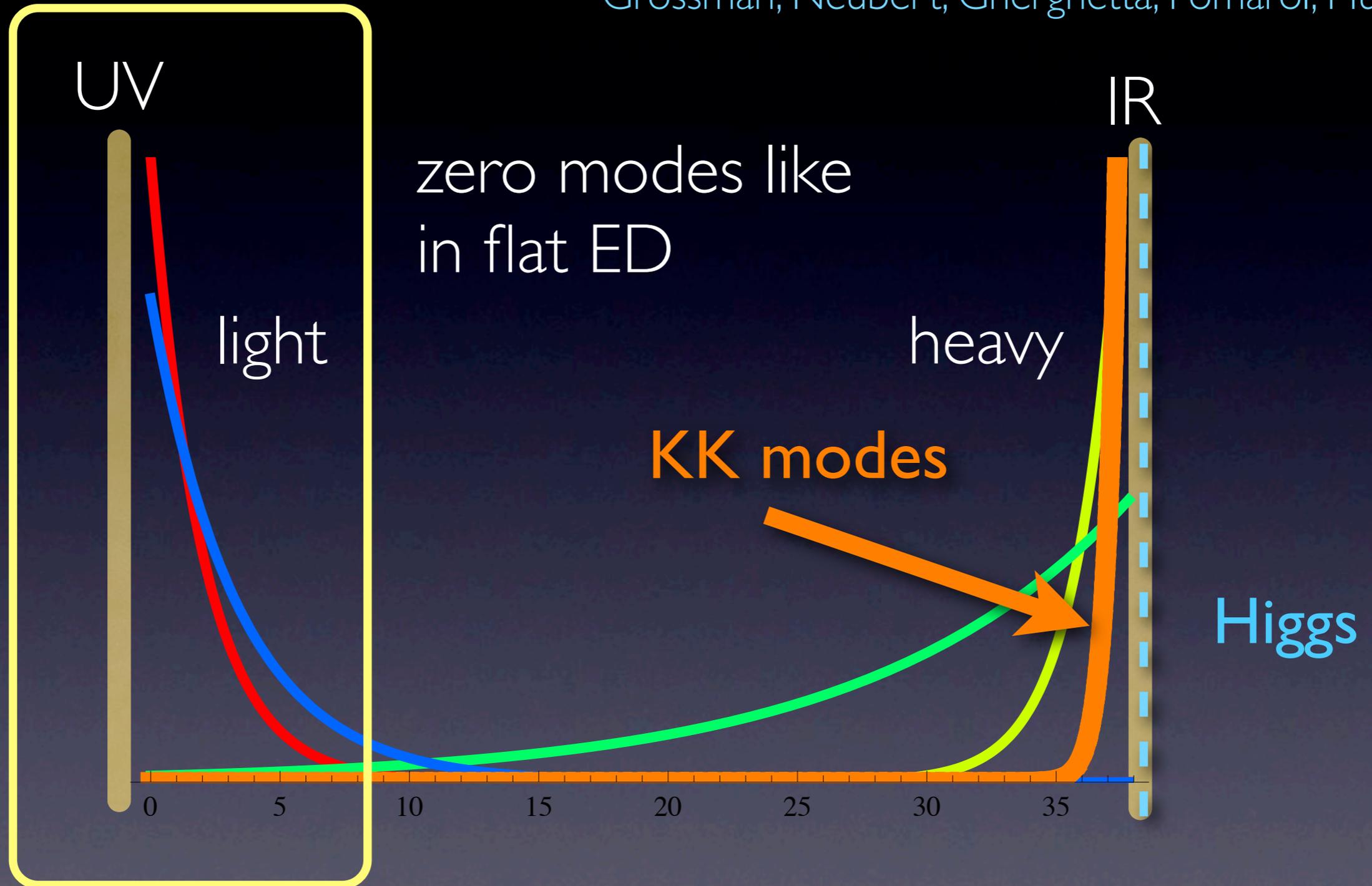
Flavor in RS

Grossman, Neubert; Gherghetta, Pomarol; Huber;



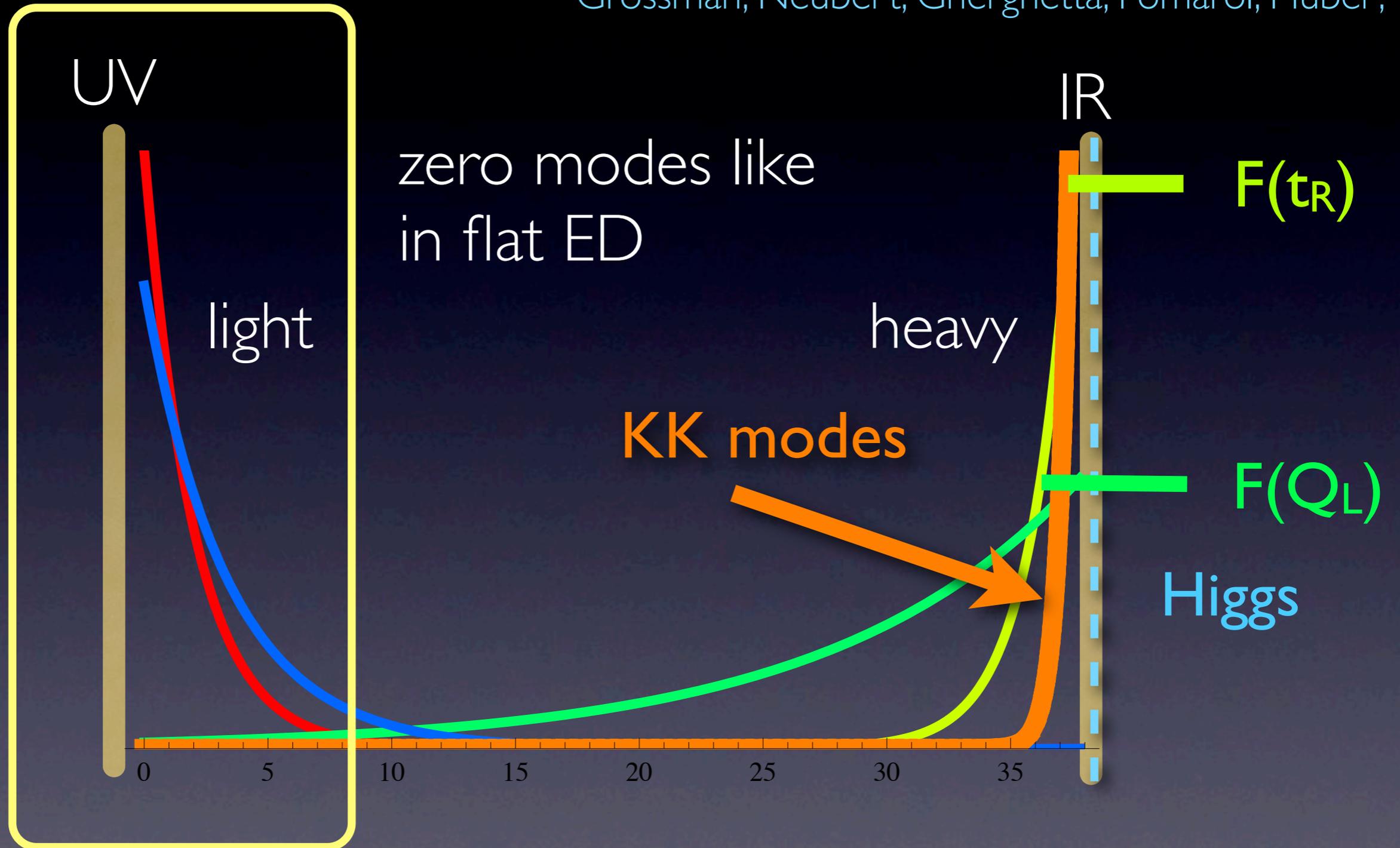
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Two ways of giving mass to fermions...

Bi-linear:

$$\mathcal{L} = y f_L \mathcal{O}_H f_R, \quad \mathcal{O}_H \sim (1, 2)_{\frac{1}{2}}$$

Linear:

D.B. Kaplan '91

$$\mathcal{L} = y f_L \mathcal{O}_R + y_R f_R \mathcal{O}_L + m \mathcal{O}_L \mathcal{O}_R, \quad \mathcal{O}_R \sim (3, 2)_{\frac{1}{6}}$$



Partial compositeness

$$\mathcal{L} = \mathcal{L}_{elem}(g_{elem}) + \mathcal{L}_{comp}(g_*) + \mathcal{L}_{mix}$$

$$1 \lesssim g_* \lesssim 4\pi$$

$$|SM\rangle = \cos \phi |elem.\rangle + \sin \phi |comp.\rangle$$

$$|heavy\rangle = -\sin \phi |elem.\rangle + \cos \phi |comp.\rangle$$

1) Linear coupling of SM fields to composites

$$\mathcal{L}_{UV} \supset \lambda \bar{\mathcal{O}}_R \psi_L$$

Contino, Pomarol

2) Strong sector conformal over some energy range

$$\mu \frac{d\lambda}{d\mu} = \gamma \lambda \quad \gamma = \dim[\mathcal{O}_R] + 3/2 - 4$$


$$\lambda \sim \left(\frac{\text{TeV}}{M_{Pl}} \right)^\gamma$$

AdS/CFT translation:

$$\gamma = c - \frac{1}{2}$$

Partial compositeness

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Degree of compositeness:

$$\sin \phi = F(c) \sim \left(\frac{\text{TeV}}{M_{pl}} \right)^{c - \frac{1}{2}}$$

Flavor from symmetries

Hierarchies from symmetries

Froggatt, Nielsen '79

Add horizontal $U(1)_F$, flavon Φ_F ($m_\Phi \sim \Lambda$, $q_F = -1$)

$$Y_d^{ij} \left(\frac{\Phi_F}{\Lambda} \right)^{-q_i + h + d_j} \bar{Q}_L^i H D_R^j$$

Hierarchies from symmetries

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$U(1)_F$ broken by $F = \langle \Phi_F \rangle$, $F < \Lambda$

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$U(1)_F$ broken by $F = \langle \Phi_F \rangle$, $F < \Lambda$

$$Y_{eff,d}^{ij} = Y_d^{ij} \left(\frac{F}{\Lambda} \right)^{-q_i + h + d_j} \Rightarrow \text{hierarchies}$$

Similar results...

hierarchical
Yukawa



Mass_{ij}

anarchic (“structure-less”)



$$\propto Y_{ij} e^{-MR(c_i + c_j)}$$

split fermions/RS

$$\propto Y_{ij} \left(\frac{\mu_{\text{low}}}{\mu_{\text{high}}} \right)^{\gamma^i + \gamma^j}$$

strong dynamics

$$\propto Y_{ij} \left(\frac{\langle \Phi \rangle}{M_{\text{mess}}} \right)^{Q^i - Q^j}$$

Froggatt-Nielsen

‘Tutte le strade portano a Roma’

