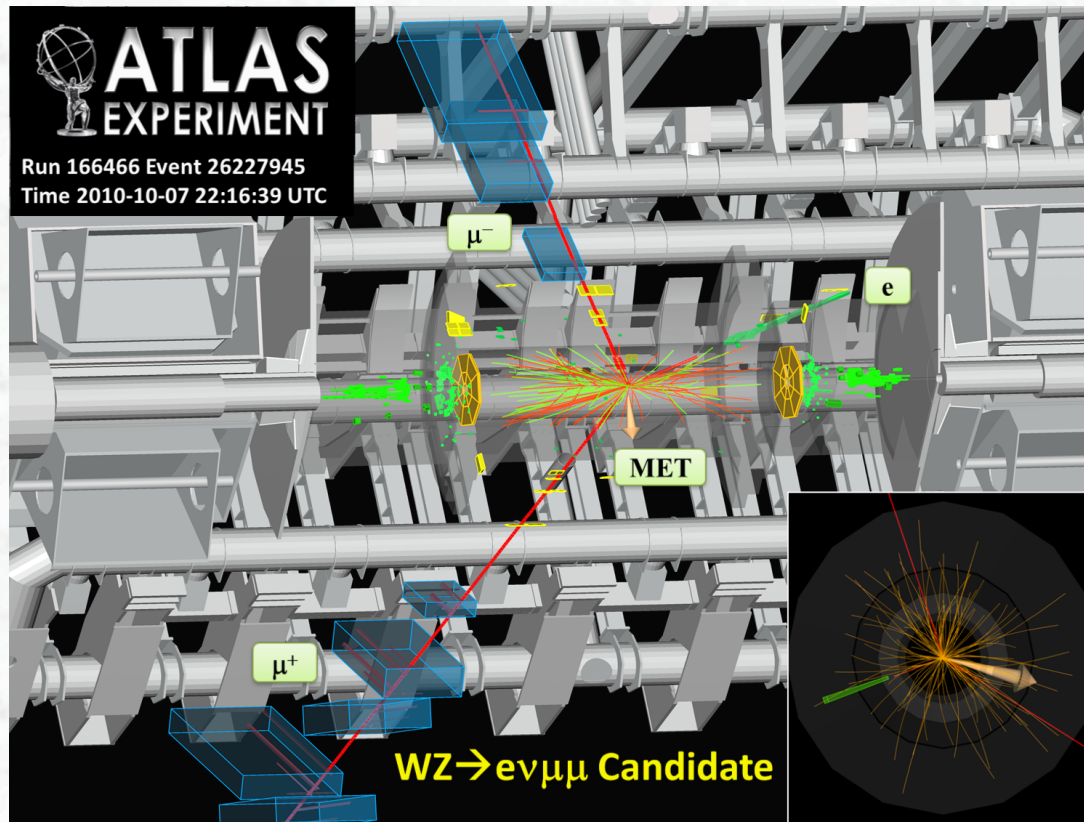


Physics at the LHC

- From the Standard Model to Searches for New Physics -



Karl Jakobs
Physikalisches Institut
Universität Freiburg

Outline of the lectures

1. Introduction
(LHC, detector performance)
2. Test of perturbative QCD
(Jet production, W/Z production, tt production)
3. Electroweak parameters
(m_W , m_t , gauge couplings, ..)
4. Summary of the search for the Higgs Boson (short → C. Mariotti)
5. Search for Physics Beyond the Standard Model
(Supersymmetry, a few other selected examples (short → M. Narain))

Disclaimer: I will try to highlight important physics measurements and results on searches for new physics. The coverage is not complete, i.e. not all results available are presented; Results from both general purpose experiments, ATLAS and CMS, plus a few from LHCb, are shown, but there might still be a bias towards the experiment I am working on. This bias is not linked to the scientific quality of the results.

The role of the LHC

1. Explore the **TeV mass scale**

- What is the origin of the electroweak symmetry breaking ?
Does the Higgs boson exist?
- Search for physics Beyond the Standard Model
(Low energy supersymmetry, other scenarios...)

Look for the “expected”, but we need to be open for surprises

→ perform as many searches (inclusive, exclusive...) for as many final states as possible

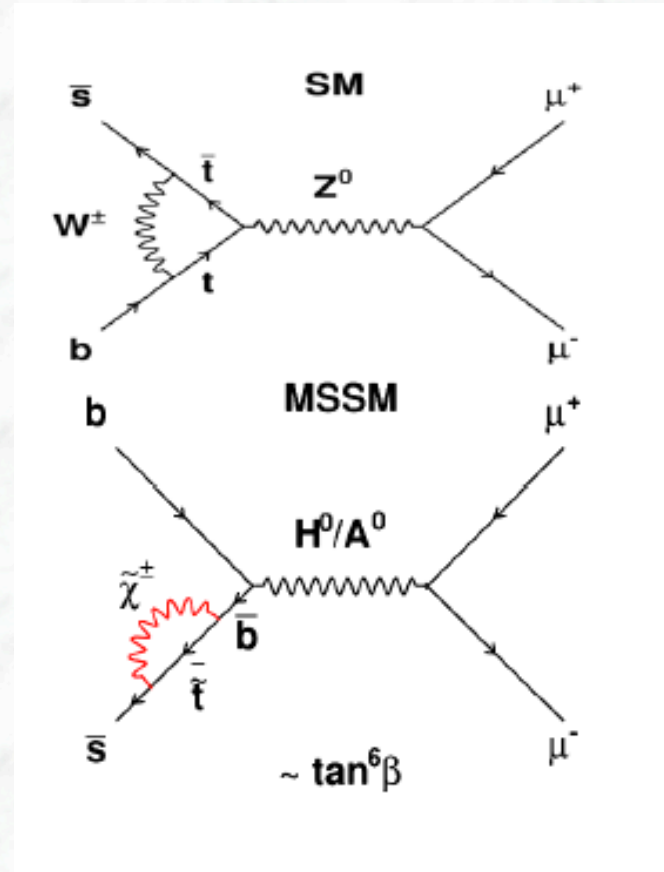
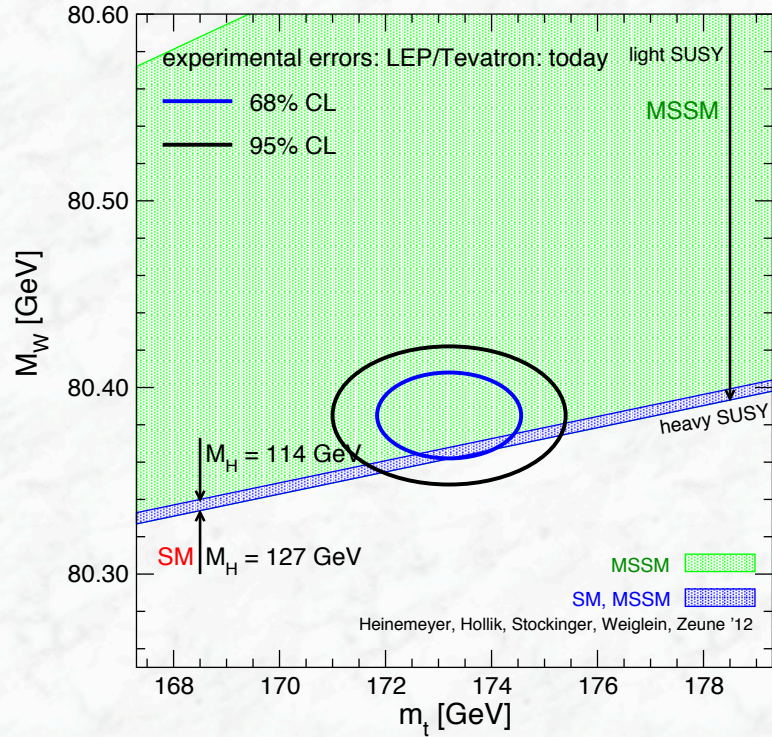
2. Precise tests of the Standard Model

- There is much sensitivity to physics beyond the Standard Model in the precision area (loop-induced effects, probe energy scales far beyond direct reach)
→ precise measurements, search for rare processes

→ Guidance to theory and Future Experiments

Two important examples:

2012



Ultimate test of the Standard Model:

Compare indirect prediction of the Higgs boson mass with direct observation

Many theoretical models
for physics Beyond the
Standard Model



An aerial photograph of a valley with a patchwork of green and brown fields. In the distance, there are blue mountains with snow-capped peaks under a clear blue sky. A red oval is drawn over the valley, with small red circles at its top and bottom points, representing the LHC tunnel. The text "The LHC" and "- a new era in particle physics-" is overlaid in yellow.

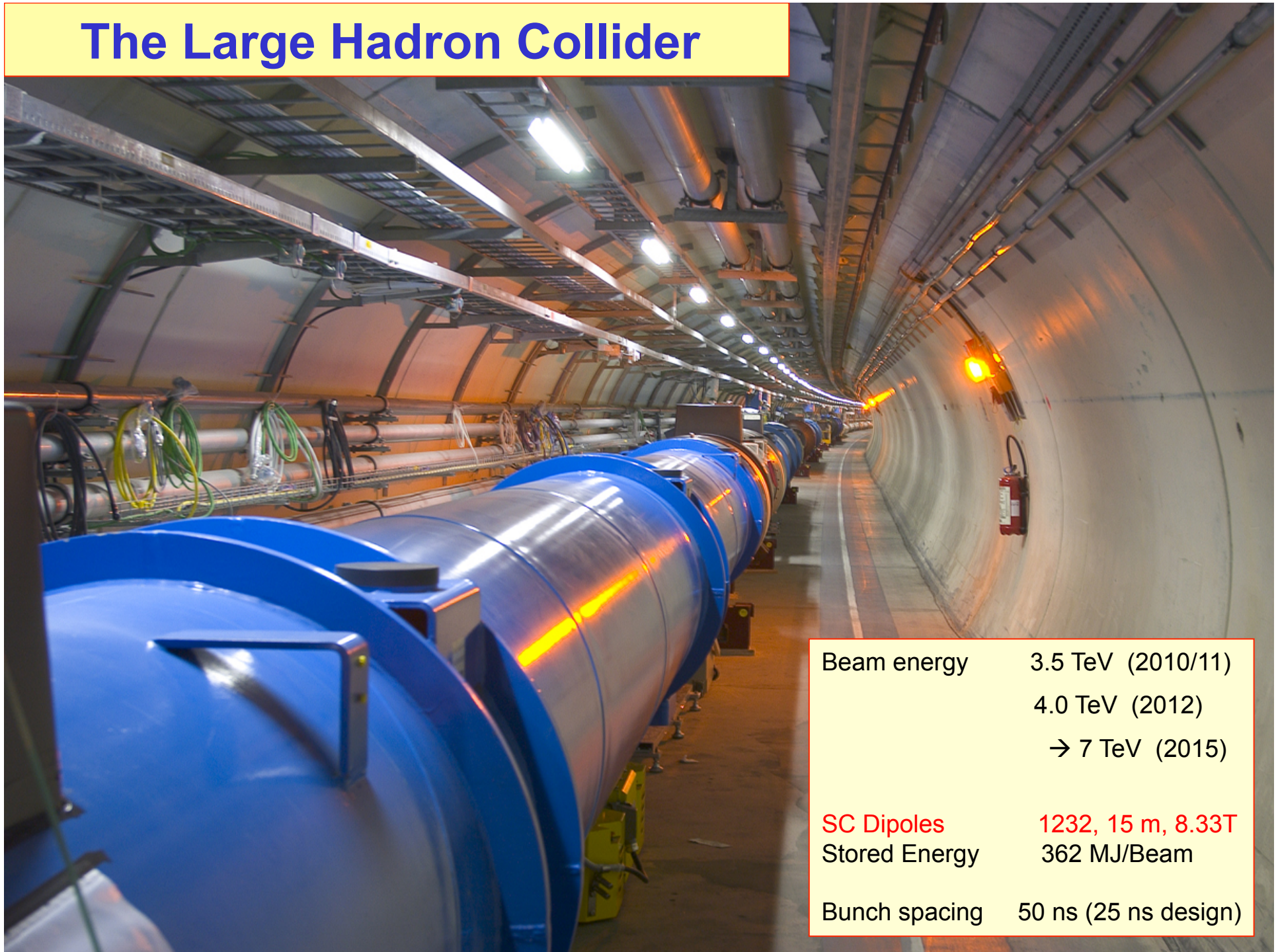
The LHC

- a new era in particle physics-

Steve Meyers at "Physics at LHC 2012":

"The first two years of LHC operation have produced sensational performance: well beyond our wildest expectations. The combination of the performance of the LHC machine, the detectors and the GRID have proven to be a terrific success story in particle physics."

The Large Hadron Collider

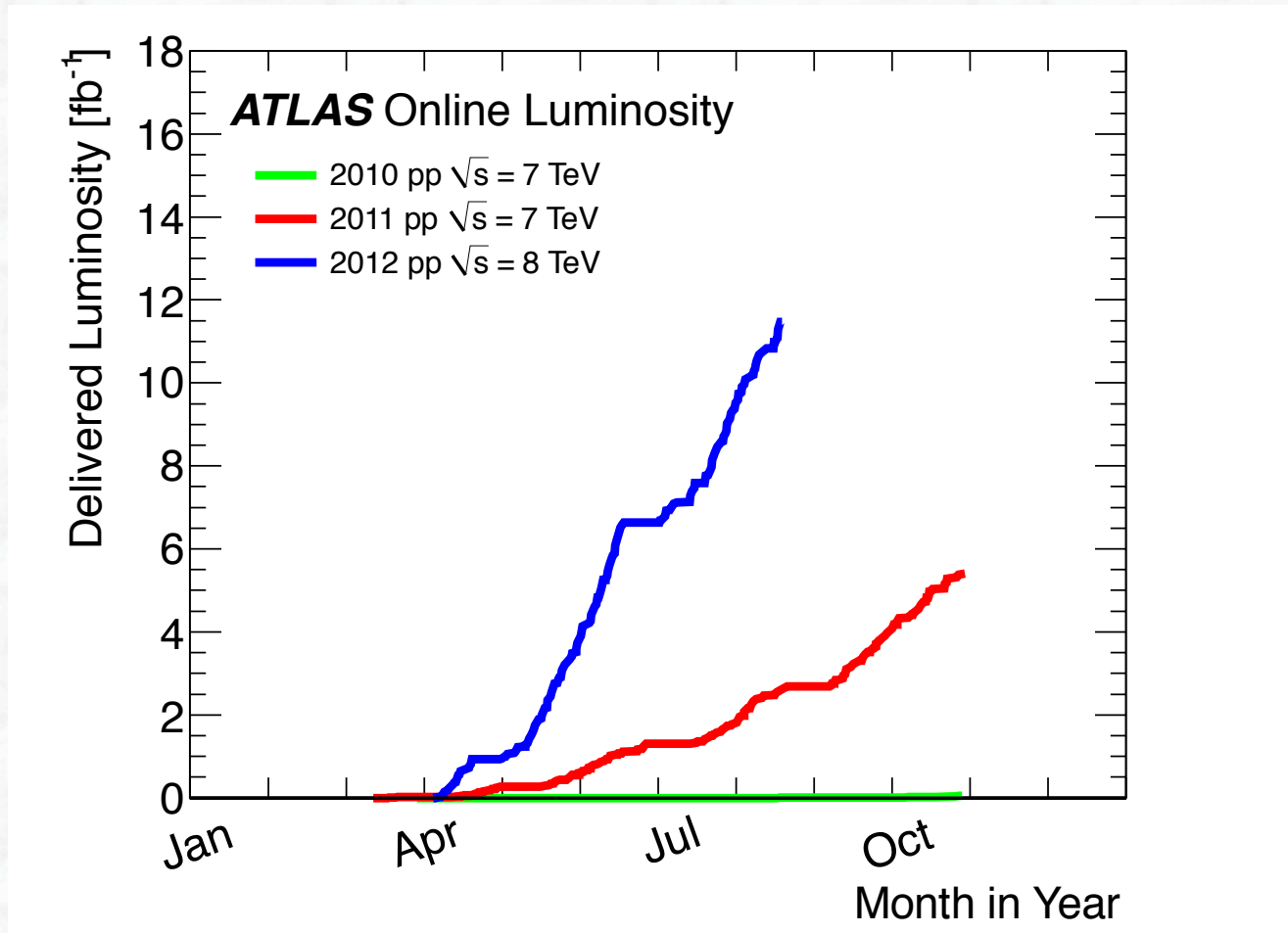


Beam energy	3.5 TeV (2010/11)
	4.0 TeV (2012)
	→ 7 TeV (2015)

SC Dipoles	1232, 15 m, 8.33T
Stored Energy	362 MJ/Beam

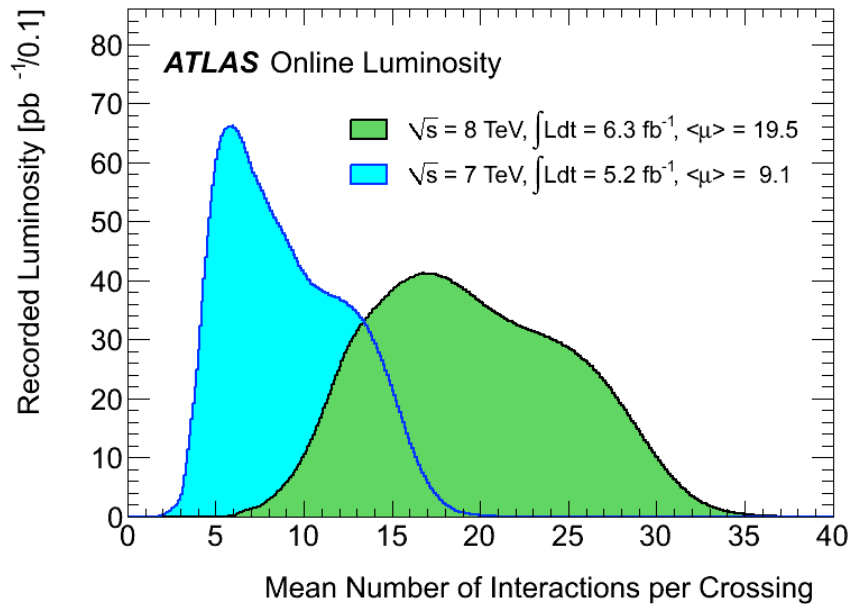
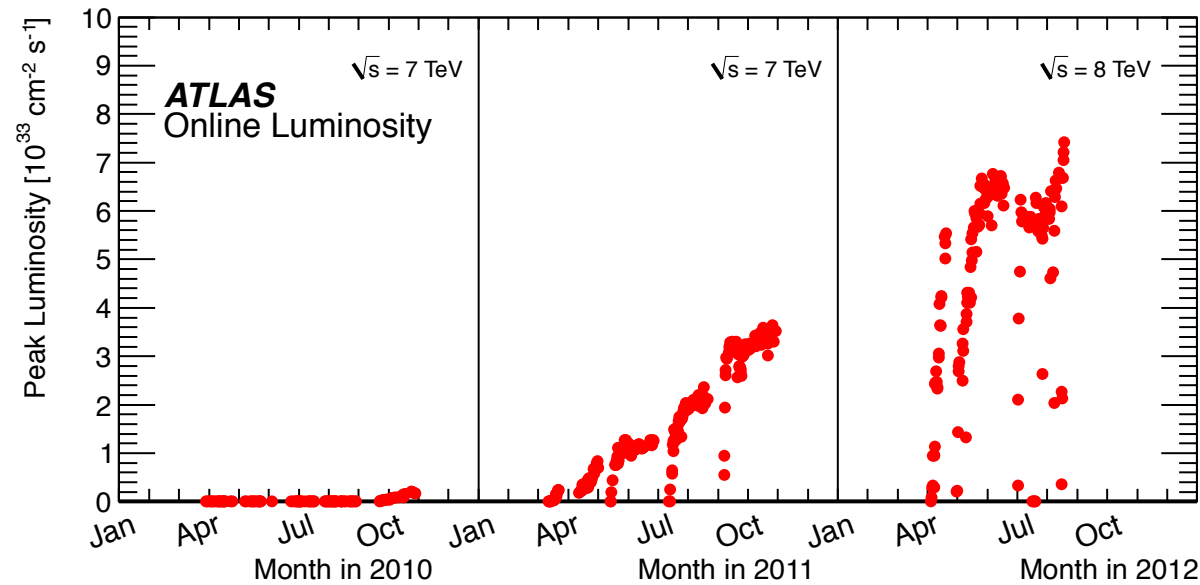
Bunch spacing	50 ns (25 ns design)
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The LHC integrated luminosity



Very rapid rise in luminosity + good machine stability
→ high integrated luminosities

The LHC instantaneous luminosity



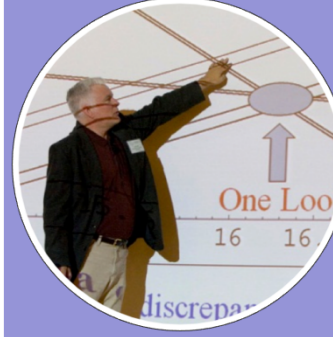
- World record on instantaneous luminosity on 22. April 2011:
 $4.67 \cdot 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$
 (Tevatron record: $4.02 \cdot 10^{32} \text{ cm}^{-2} \text{ s}^{-1}$)
- 2011: collect per day as much integrated luminosity as in 2010
- 2012: now regularly above $6 \cdot 10^{33} \text{ cm}^{-2} \text{ s}^{-1}$

Completion of an era: Tevatron



Accelerator Innovations

- First major SC synchrotron
- Industrial production of SC cable (MRI)
- Electron cooling
- New RF manipulation techniques



But Tevatron is still in the game:

- W mass
- $H \rightarrow bb$
- B physics
- ...



and developing
• GRID pioneers

Major discoveries

- Top quark
- B_s mixing
- Precision W and Top mass \rightarrow
- Higgs mass prediction
- Direct Higgs searches
- Ruled out many exotica

The next generation

- Fantastic training ground for next generation
- More than 500 Ph.D.s
- Produced critical personnel for the next steps, especially LHC



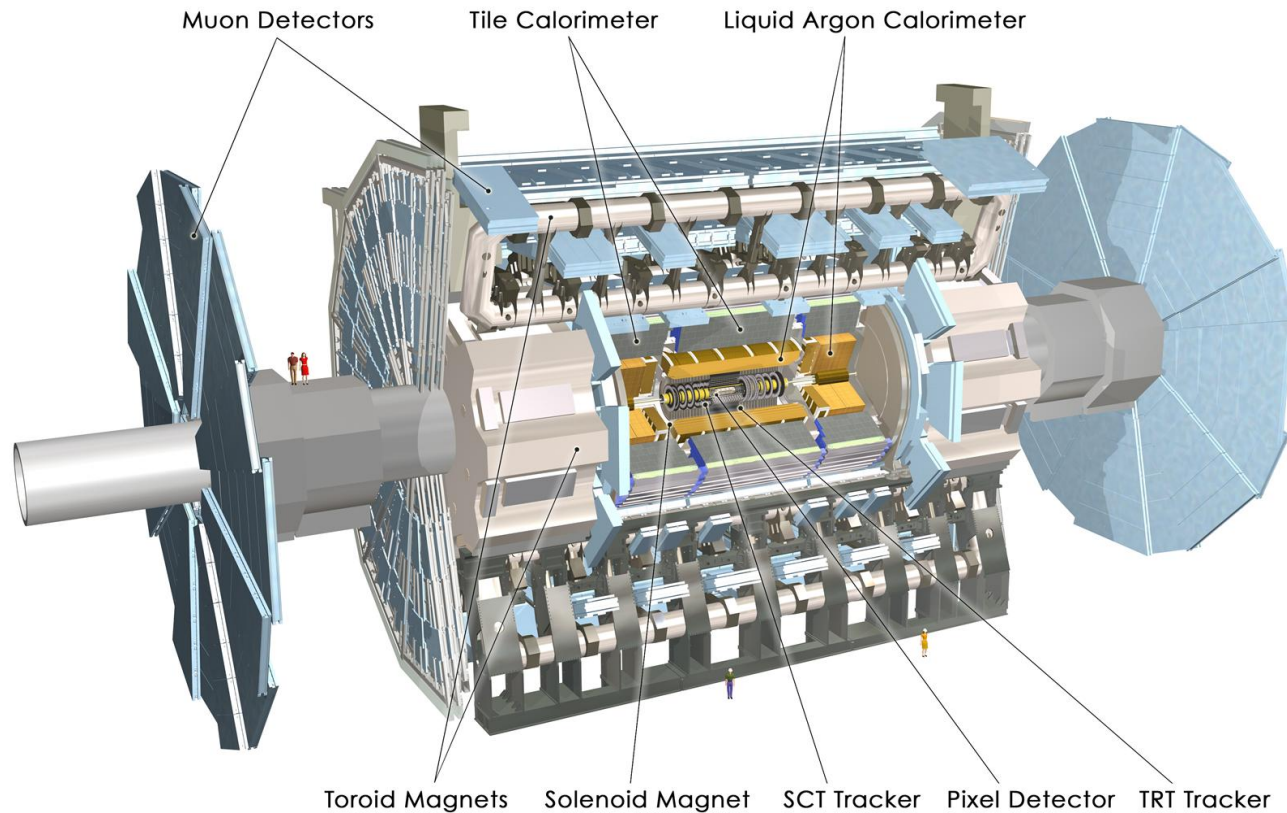
CMS

LHCb

ALICE

ATLAS

The ATLAS experiment



- Solenoidal magnetic field (2T) in the central region (momentum measurement)

High resolution silicon detectors:

- 6 Mio. channels (80 μm x 12 cm)
 - 100 Mio. channels (50 μm x 400 μm)
- space resolution: $\sim 15 \mu\text{m}$

- Energy measurement down to 1° to the beam line
- Independent muon spectrometer (supercond. toroid system)

Diameter	25 m
Barrel toroid length	26 m
End-cap end-wall chamber span	46 m
Overall weight	7000 Tons

CMS

Superconducting
Coil, 4 Tesla

CALORIMETERS

ECAL

76k scintillating
PbWO₄ crystals

HCAL

Plastic scintillator/brass
sandwich

IRON YOKE

TRACKER

Pixels
Silicon Microstrips
210 m² of silicon sensors
9.6M channels

MUON BARREL

Drift Tube
Chambers (**DT**)

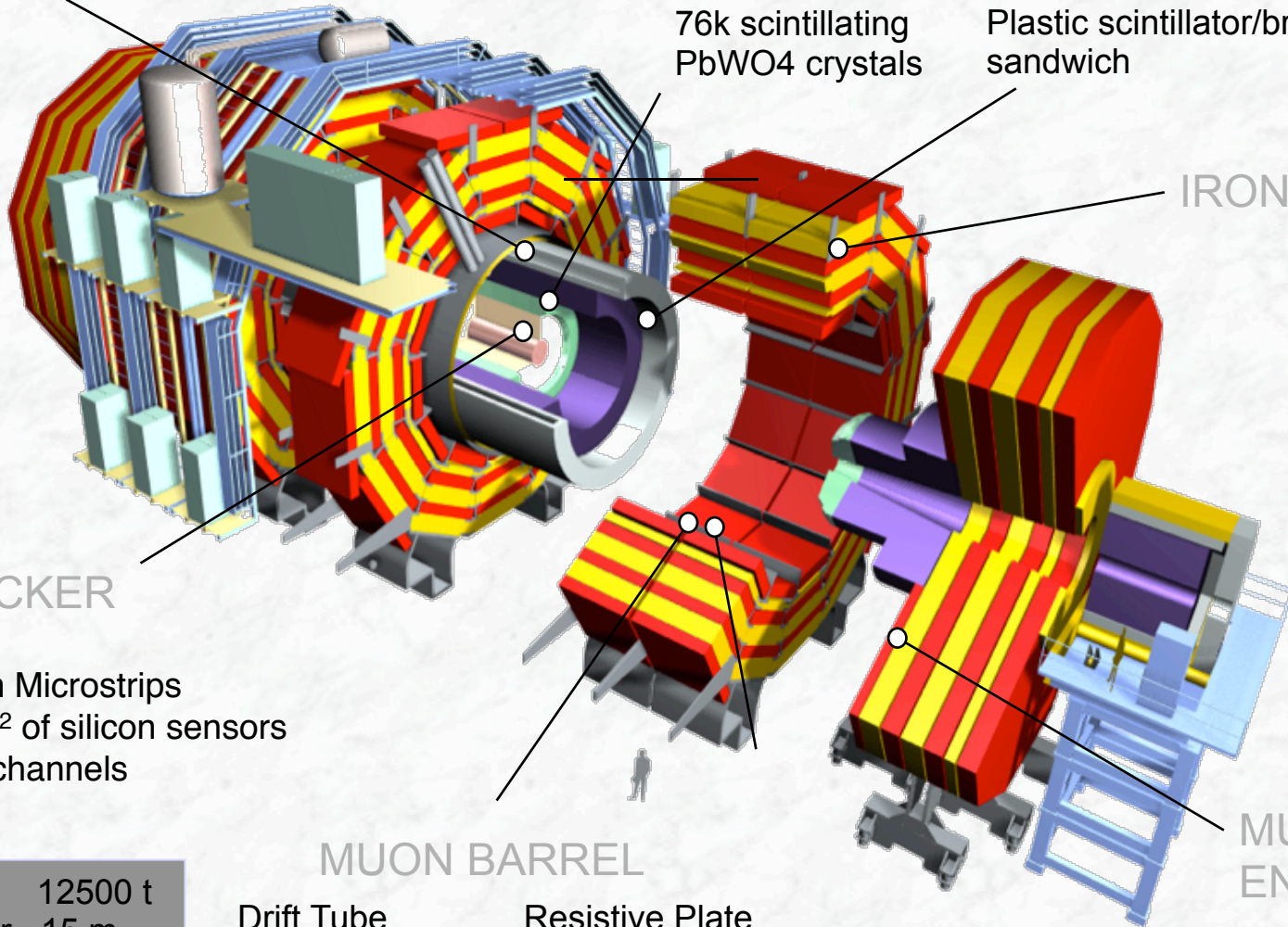
Resistive Plate
Chambers (**RPC**)

MUON
ENDCAPS

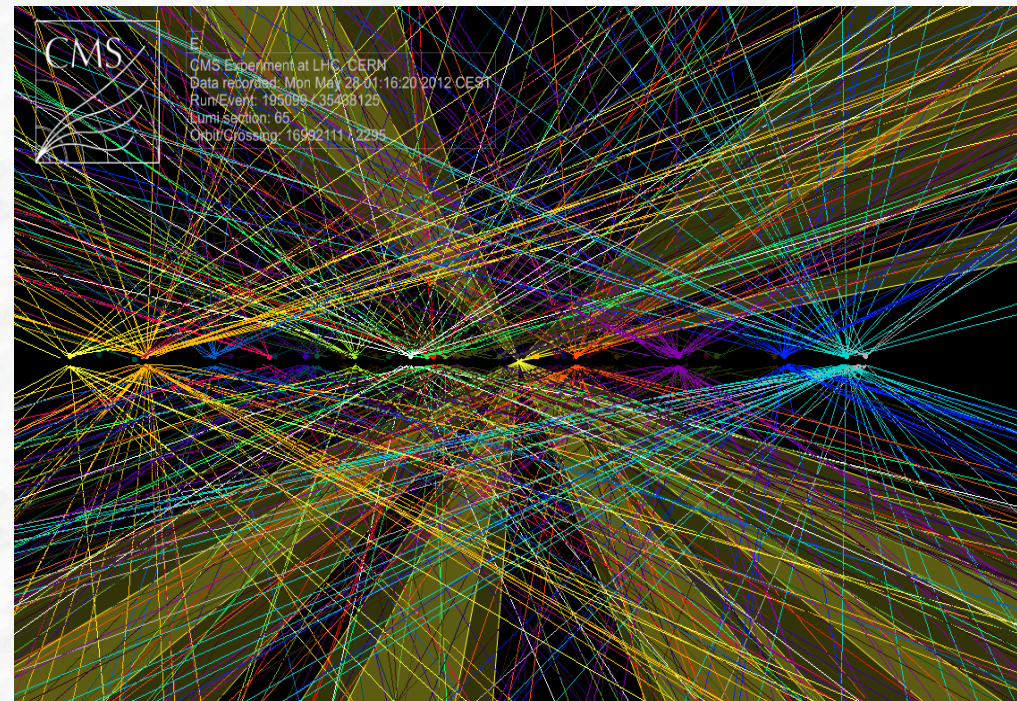
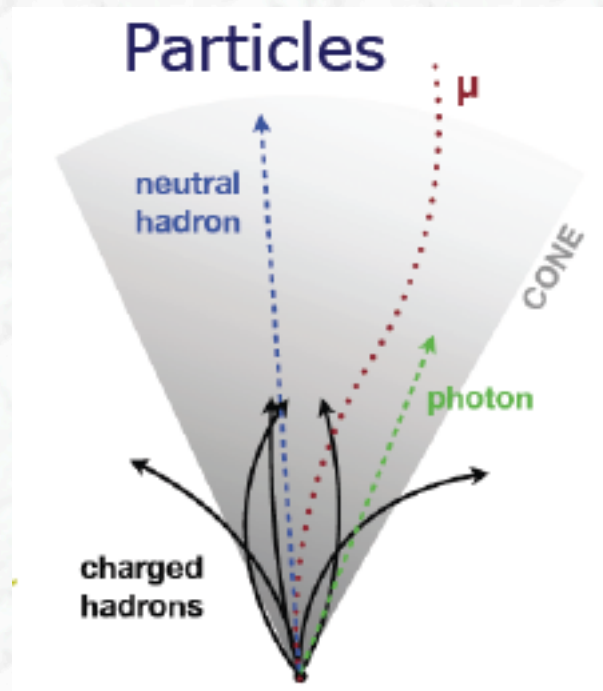
Cathode Strip Chambers (**CSC**)
Resistive Plate Chambers (**RPC**)

15

Total weight	12500 t
Overall diameter	15 m
Overall length	21.6 m



1.2 Detector Performance



Some bonus slides on

“Important kinematic variables

in pp collisions”

(i) Rapidity y

Usually the beam direction is defined as the z axis (Transverse plane: x - y plane).

The rapidity y is defined as:

$$y = \frac{1}{2} \ln \left(\frac{E + p_z}{E - p_z} \right) = \tanh^{-1} \left(\frac{p_z}{E} \right)$$

Under a **Lorentz boost** in the z -direction to a frame with velocity β

the rapidity y transforms as: $y \rightarrow y - \tanh^{-1} \beta$

Hence the shape of the rapidity distribution dN/dy is invariant, as are differences in rapidity.

(ii) Pseudorapidity η

Rapidity:
$$y = \frac{1}{2} \ln \left(\frac{E + p_z}{E - p_z} \right) = \tanh^{-1} \left(\frac{p_z}{E} \right)$$

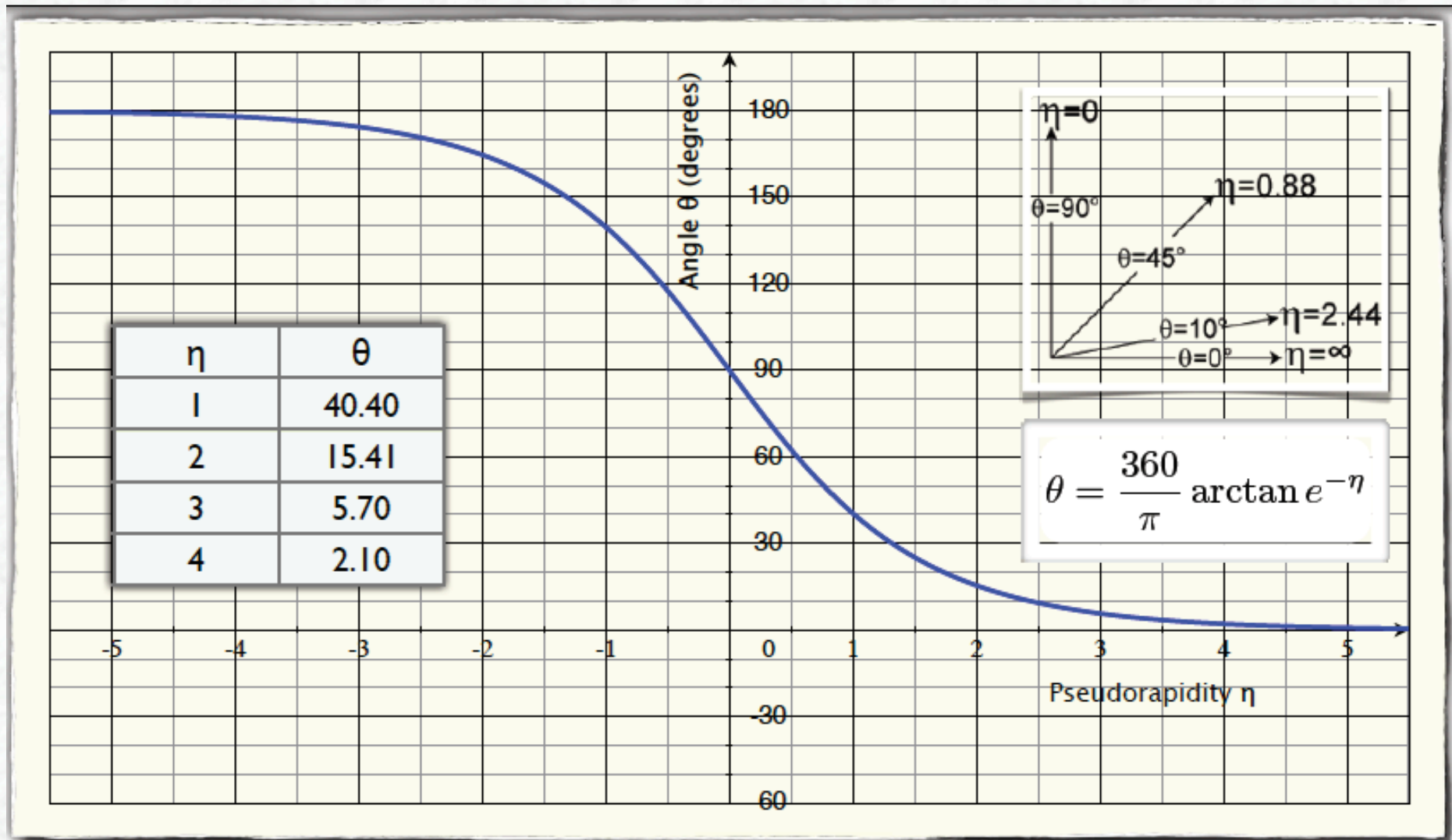
For $p \gg m$, the rapidity may be expanded to obtain

$$y = \frac{1}{2} \ln \frac{\cos^2(\theta/2) + m^2/4p^2 + \dots}{\sin^2(\theta/2) + m^2/4p^2 + \dots}$$
$$\approx -\ln \tan(\theta/2) \equiv \eta$$

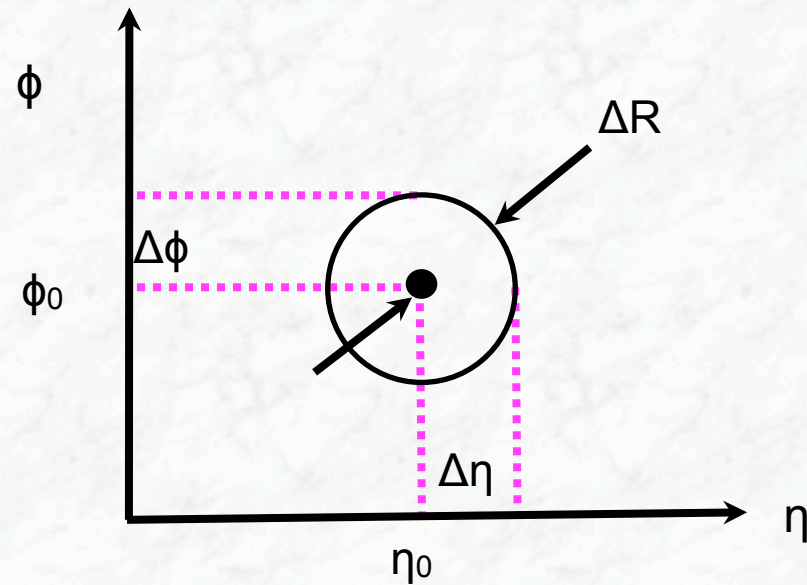
where $\cos \theta = p_z/p$.

Identities: $\sinh \eta = \cot \theta$, $\cosh \eta = 1/\sin \theta$, $\tanh \eta = \cos \theta$

Relation between pseudorapidity η and polar angle θ



(iii) Distance in $\eta - \phi$ space:



Rapidity y :

$$y = 1/2 \ln[(E + p_z)/(E - p_z)]$$

Pseudorapidity η :

$$\eta = -\ln \tan(\theta/2)$$

Distance in η - ϕ :

$$\Delta R = \sqrt{\Delta\eta^2 + \Delta\phi^2}$$

(iv) Transverse Energy

At hadron colliders, a significant and unknown proportion of the energy of the incoming hadrons in each event escapes down the beam-pipe. Consequently if invisible particles are created in the final state, their net momentum can only be constrained in the plane transverse to the beam direction. Defining the z-axis as the beam direction, this net momentum is equal to the missing transverse energy vector

missing transverse energy

$$\mathbf{E}_T^{\text{miss}} = - \sum_i \mathbf{p}_T(i)$$

where the sum runs over the transverse momenta of all visible final state particles.

(v) Transverse mass (invisible particles)

Consider a single heavy particle of mass M which decays to two particles, of which one (labelled particle 1) is invisible. The mass of the parent particle can be constrained with the quantity M_T defined by

$$M_T^2 \equiv [E_T(1) + E_T(2)]^2 - [\mathbf{p}_T(1) + \mathbf{p}_T(2)]^2$$

$$= m_1^2 + m_2^2 + 2[E_T(1)E_T(2) - \mathbf{p}_T(1) \cdot \mathbf{p}_T(2)]$$

where $\mathbf{p}_T(1) = -\mathbf{P}_T^{\text{miss}}$

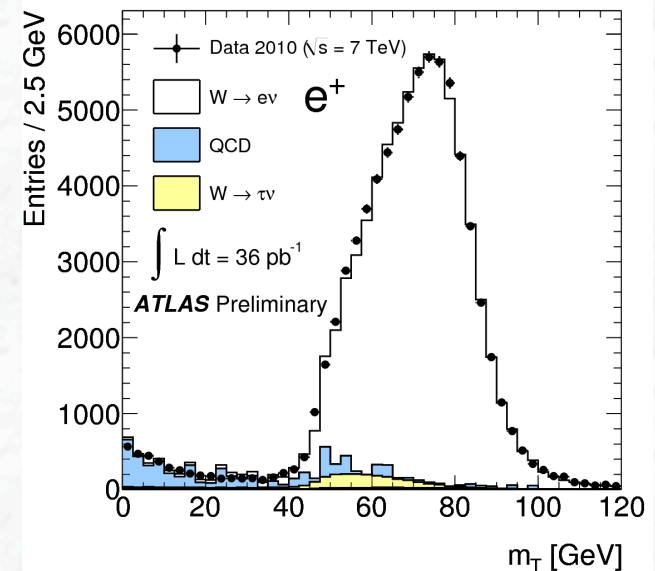
This quantity is called the **transverse mass**. Its distribution possesses an endpoint at $M_T^{\text{max}} = M$.

For $m_1 = m_2 = 0 \rightarrow$

$$M_T^2 = 2|\mathbf{p}_T(1)||\mathbf{p}_T(2)|(1 - \cos \phi_{12})$$

where ϕ_{ij} is defined as the angle between particles i and j in the transverse plane.

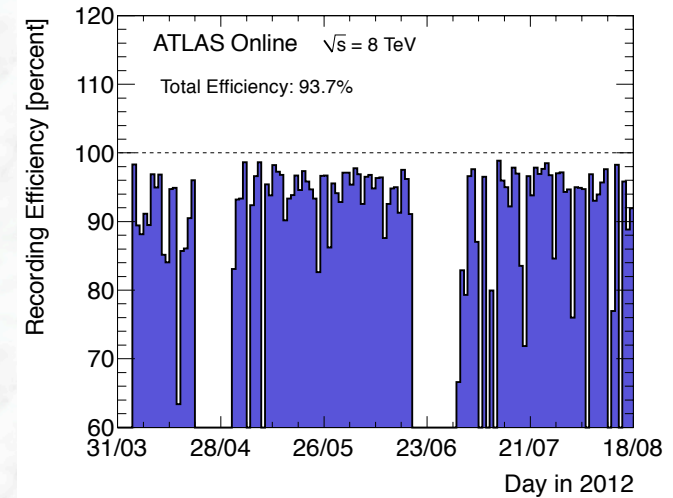
Transverse mass



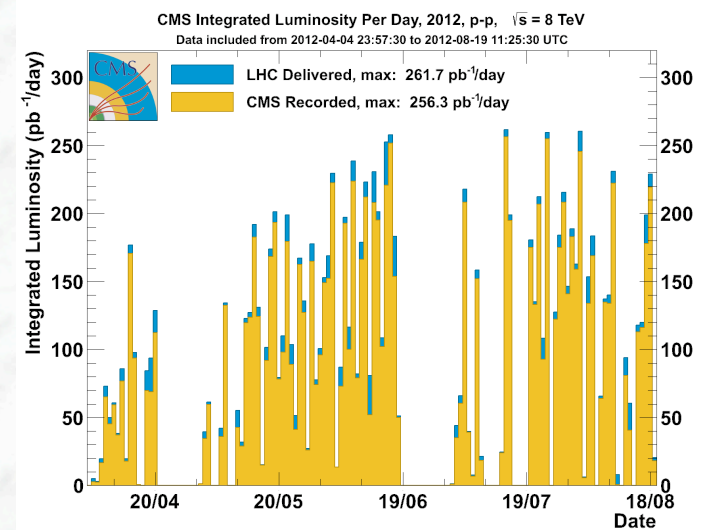
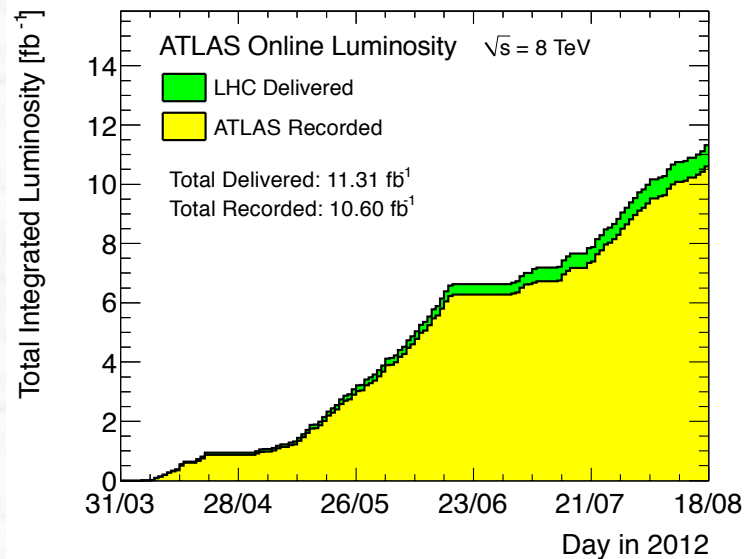
$$m_T = \sqrt{2P_T(e)E_T^{\text{miss}}(1 - \cos \Delta\phi)}$$

Detector performance is impressive:

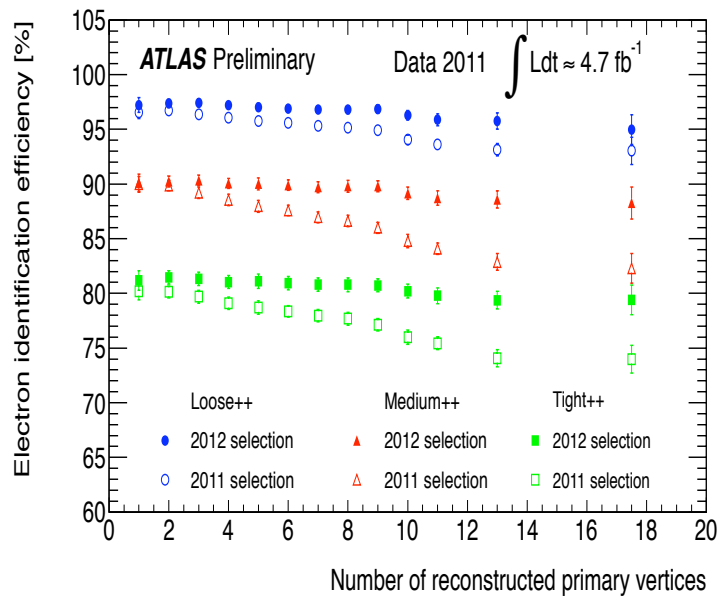
- Very high number of working channels (> 99% for many sub-systems) in all experiments;
- Data taking efficiency is high (> 94%)
- Impressive reconstruction capabilities for physics objects ($e, \gamma, \mu, \tau, \text{jets}, b\text{-tagging}, E_T^{\text{miss}}$)



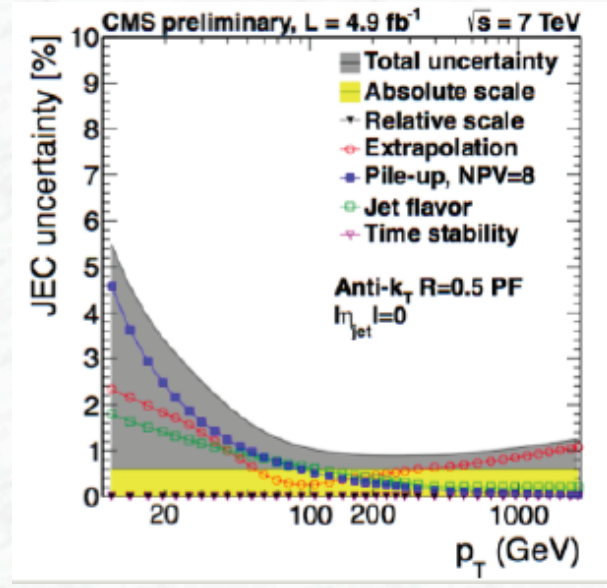
Have been optimized to cope with the ever increasing number of pile-up interactions



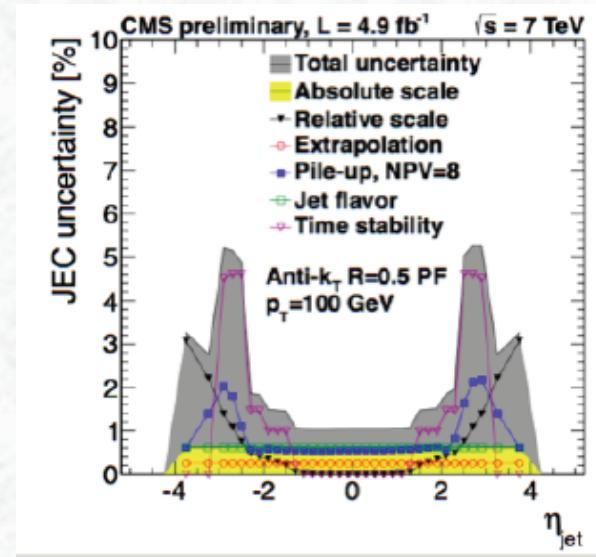
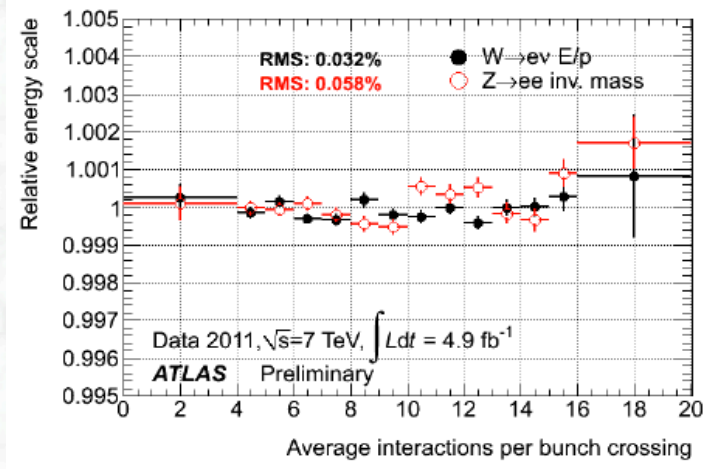
Some performance figures from 2011 data:



Electron ID efficiency in ATLAS

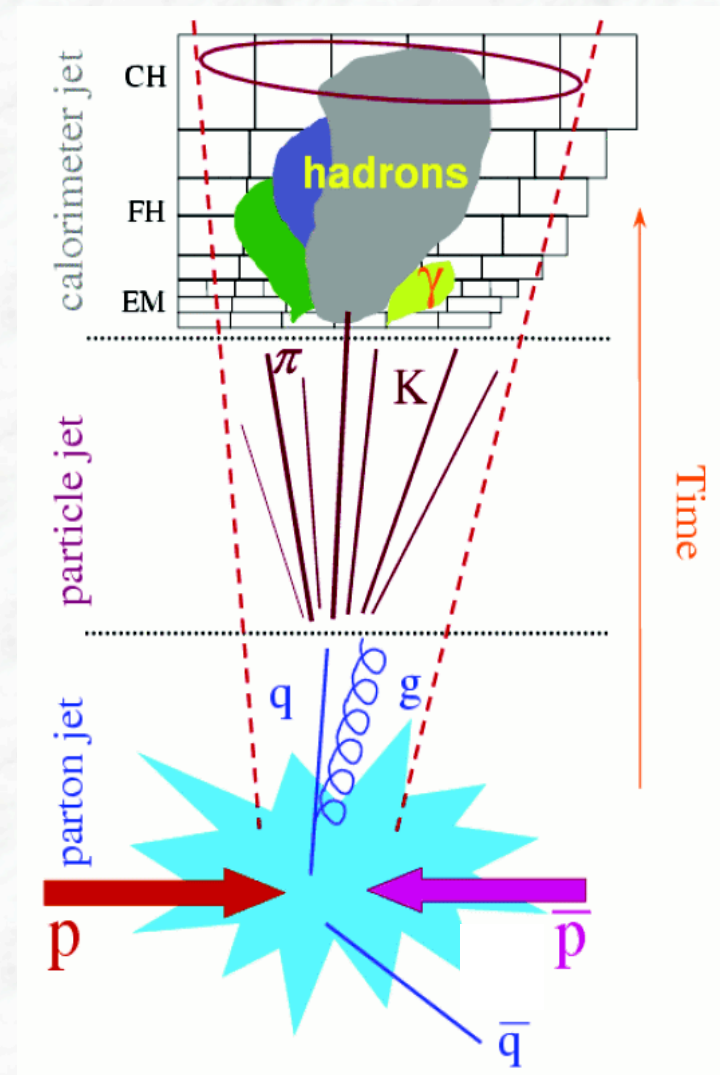


Jet energy scale, E-flow in CMS

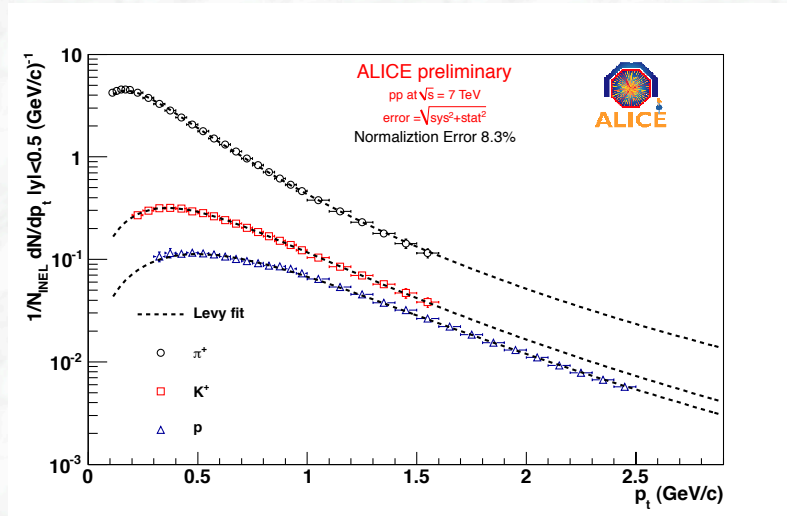
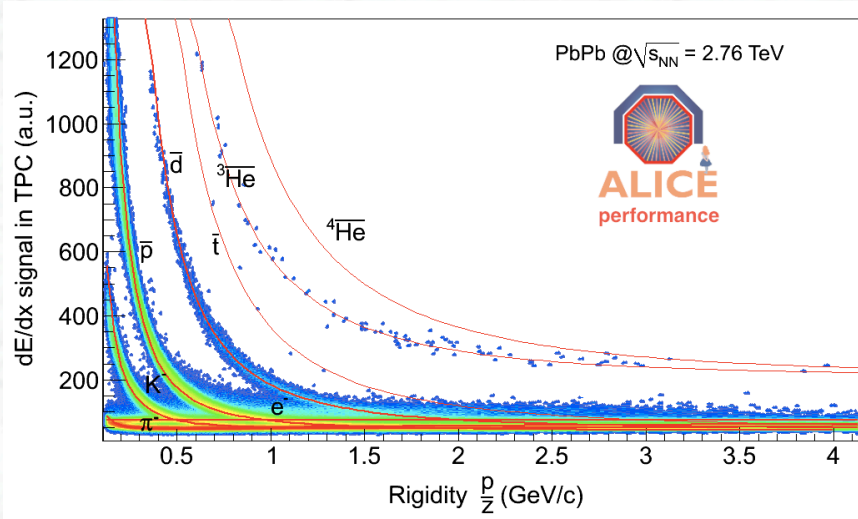


Jet reconstruction and energy measurement

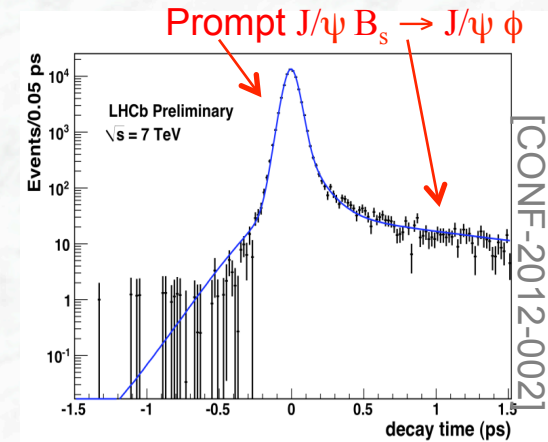
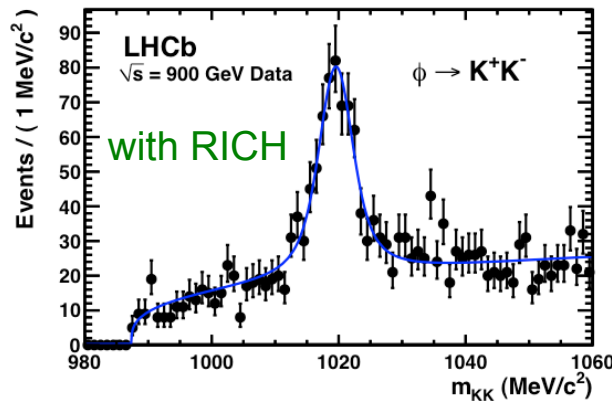
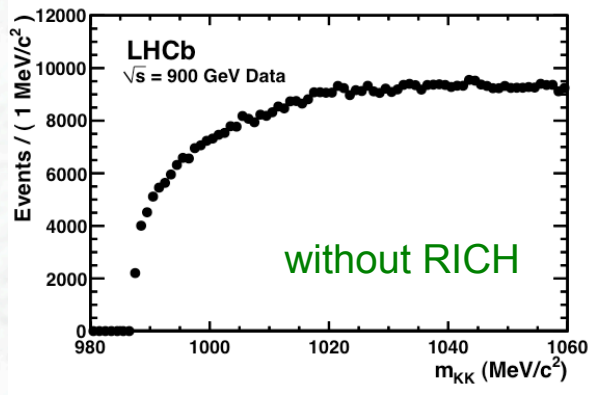
- A jet is **NOT** a well defined object
(fragmentation, gluon radiation, detector response)
- The detector response is different for particles interacting electromagnetically (e, γ) and for hadrons
→ for comparisons with theory, one needs to correct back the calorimeter energies to the „particle level“ (particle jet)
Common ground between theory and experiment
- One needs an algorithm to define a jet and to measure its energy
conflicting requirements between experiment and theory (exp. simple, e.g. cone algorithm, vs. theoretically sound (no infrared divergencies))
- Energy corrections for losses of fragmentation products outside jet definition and underlying event or pileup energy inside



Particle Identification in ALICE and LHCb:

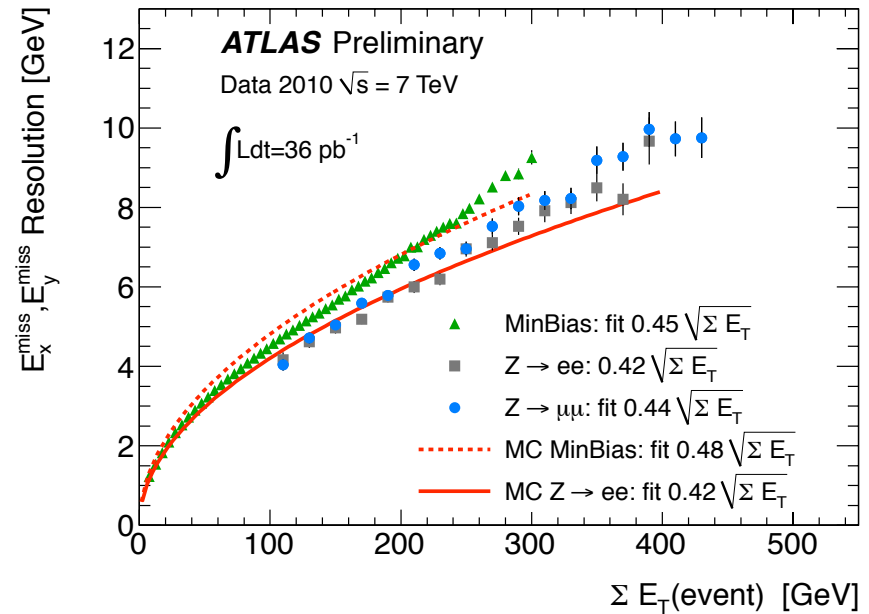
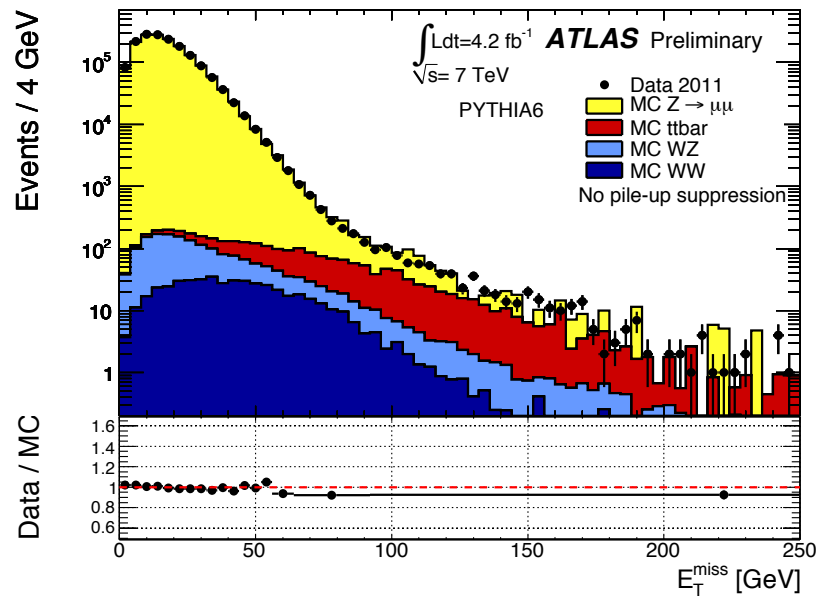


LHCb: Search for $\phi \rightarrow K^+K^-$



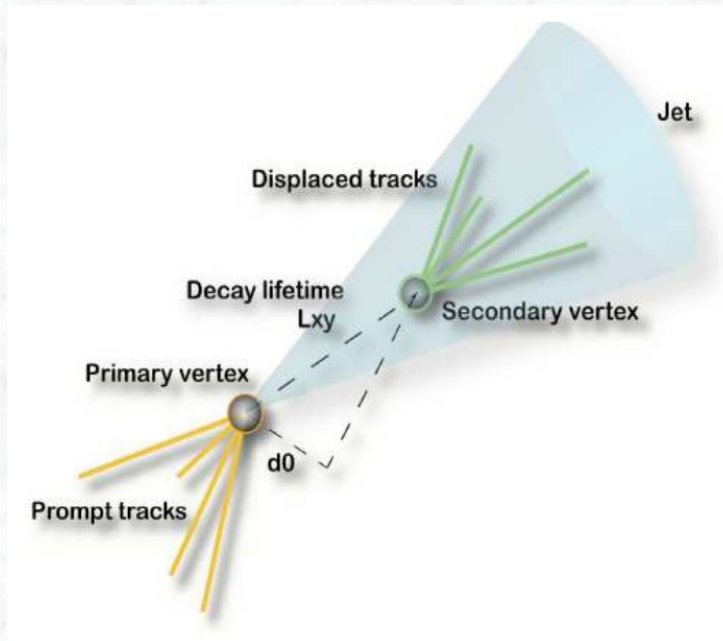
Proper time resolution: 45 fs

Measurement of the missing transverse energy E_T^{miss}



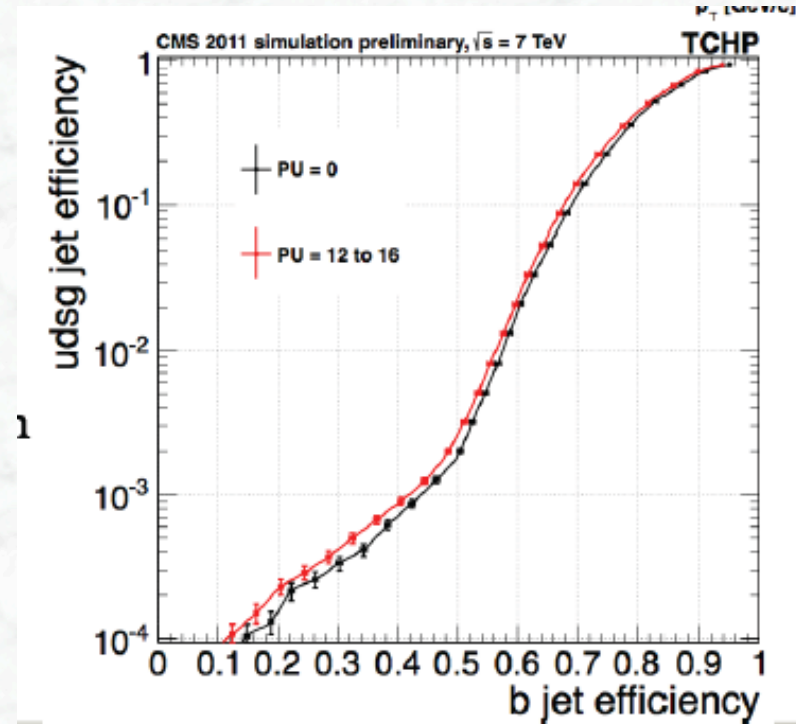
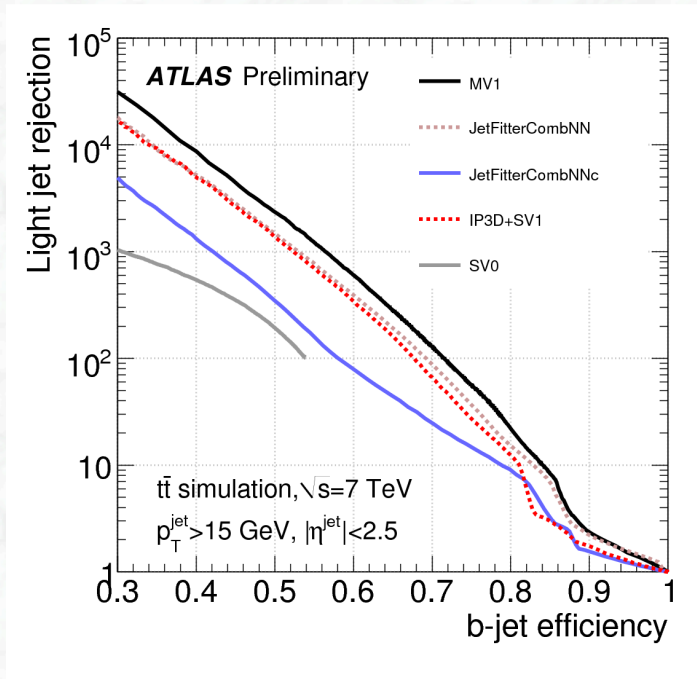
Resolution of E_x^{miss} and E_y^{miss} as a function of the total transverse energy in the event calculated by summing the p_T of muons and the total calorimeter energy. The resolution in $Z \rightarrow ee$ and $Z \rightarrow \mu\mu$ events is compared with the resolution in minimum bias for data taken at $\sqrt{s} = 7 \text{ TeV}$. The fit to the resolution in Monte Carlo minimum bias and $Z \rightarrow ee$ events are superposed.

How well can b-quarks be tagged ?

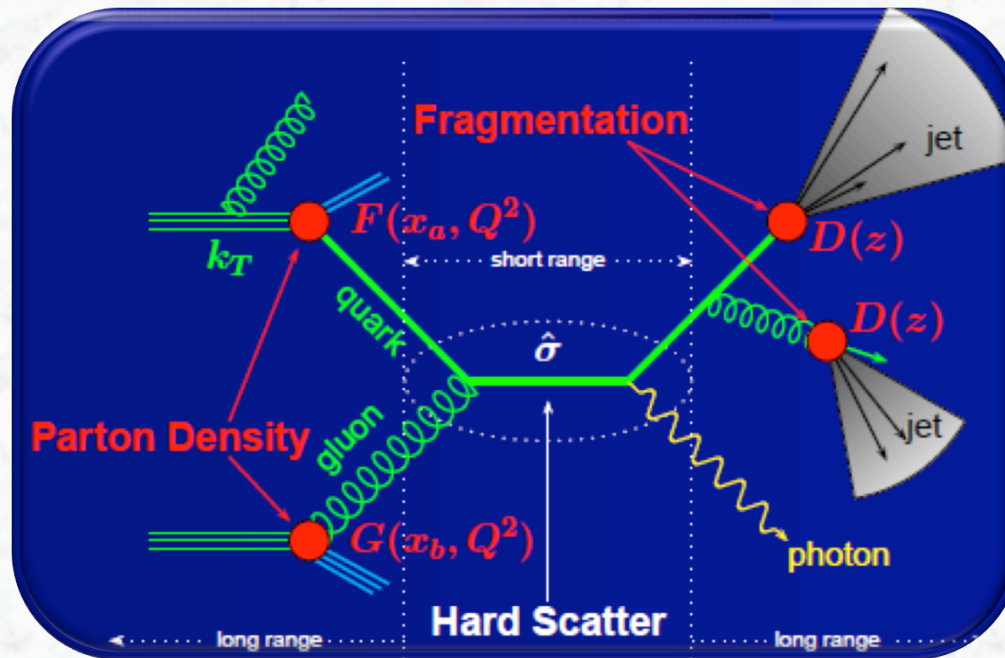


- b quarks fragment into B hadrons (mesons and baryons)
 - B mesons have a lifetime of ~ 1.5 ps
- They fly in the detector about 2-3 mm before they decay
- reconstruction of a secondary vertex possible
(requires high granularity silicon pixel and strip detectors close to the interaction point)
 - tracks from B meson decays have a large impact parameter w.r.t. the primary vertex

b-tagging performances in ATLAS and CMS: extremely important for many physics analyses (Higgs, SUSY, SM,)

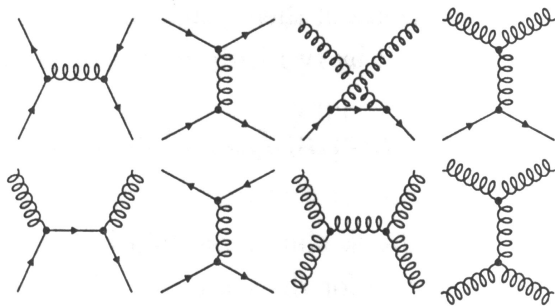


1.3 Scattering processes at a hadron collider

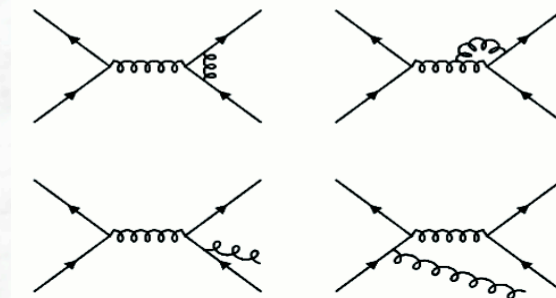


Dominant hard scattering processes: qq, qg and gg “scattering”

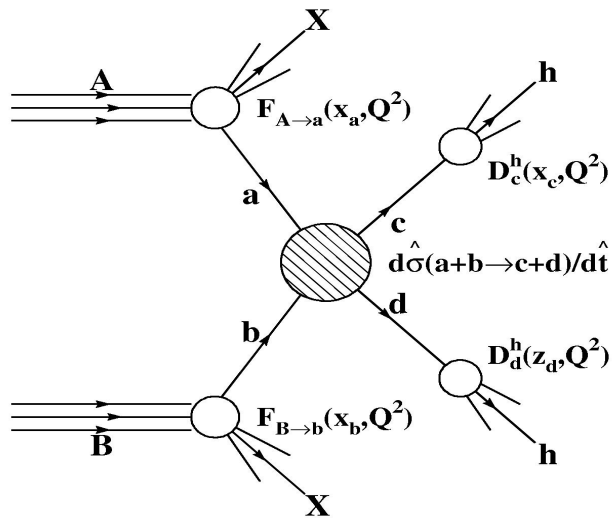
Leading order



...some NLO contributions



Calculation of cross sections



$$\sigma = \sum_{a,b} \int dx_a dx_b f_a(x_a, Q^2) f_b(x_b, Q^2) \hat{\sigma}_{ab}(x_a, x_b, \alpha_s)$$

Sum over initial partonic states a, b

$\hat{\sigma}_{ab} \equiv$ hard scattering cross section

$f_i(x, Q^2) \equiv$ parton density function

... + higher order QCD corrections (perturbation theory)

meanwhile available for many signal and background processes !

Huge theoretical effort

which for some processes turn out to be large

(e.g. Higgs production via gg fusion)

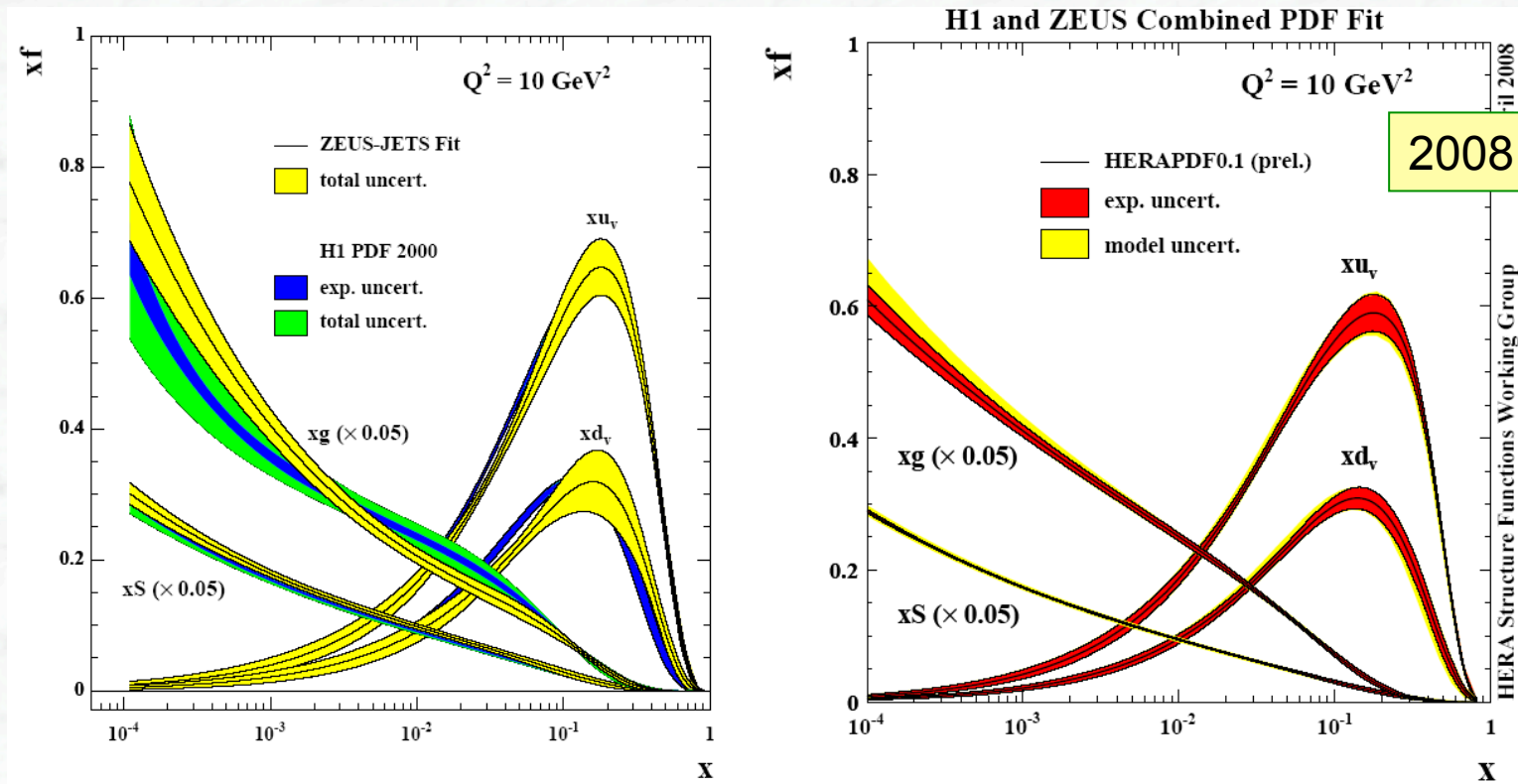
usually introduced as K-factors: $K_{[n]} = \sigma_{[n]} / \sigma_{[LO]}$

a few examples: Drell-Yan production of W/Z: $K_{NLO} \sim 1.2$

Higgs production via gg fusion: $K_{NLO} \sim 1.8$

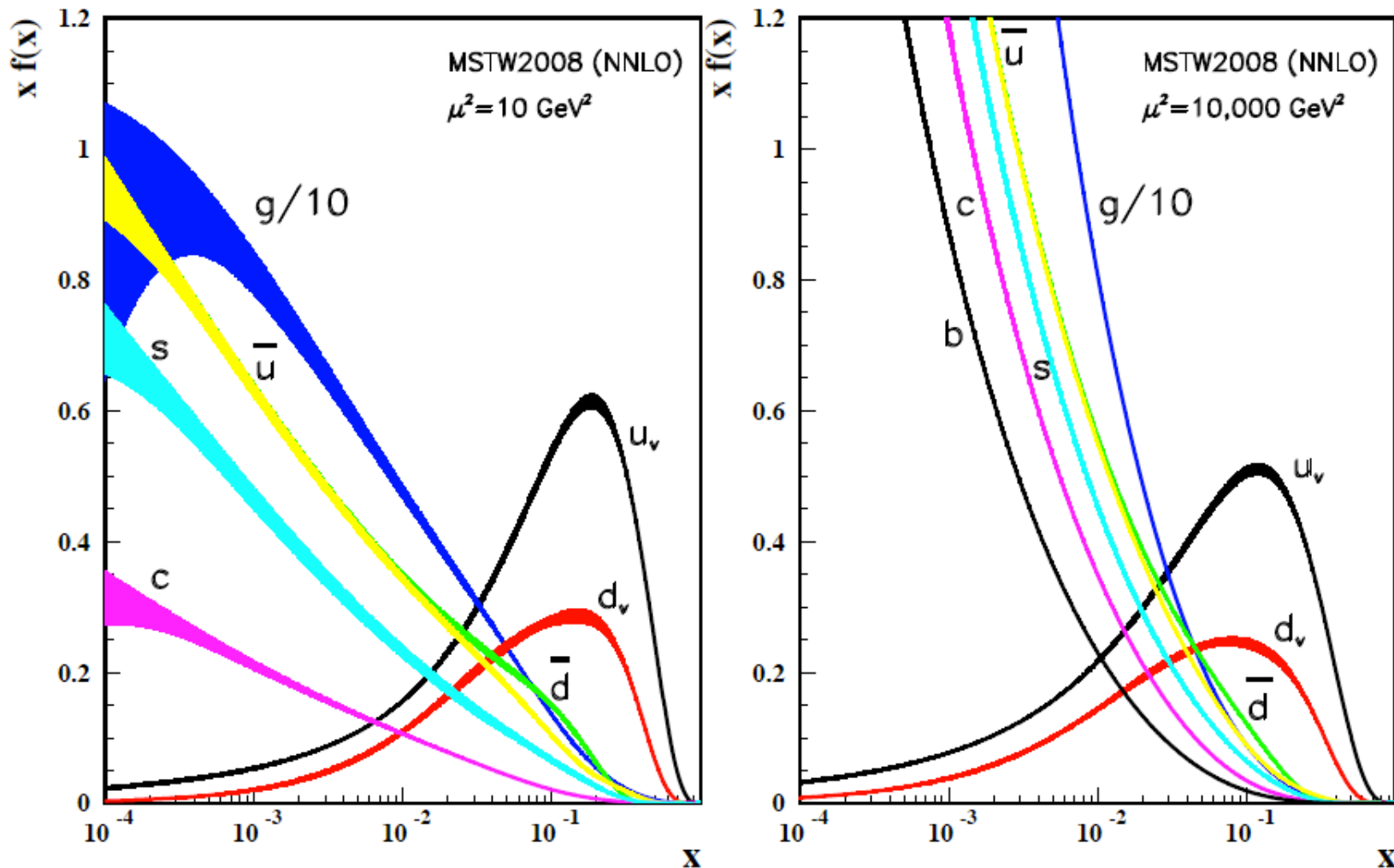
Results from HERA on the proton structure

- Large data sets and combination of the two HERA experiments (H1 and ZEUS) improve the precision on the parton distribution functions

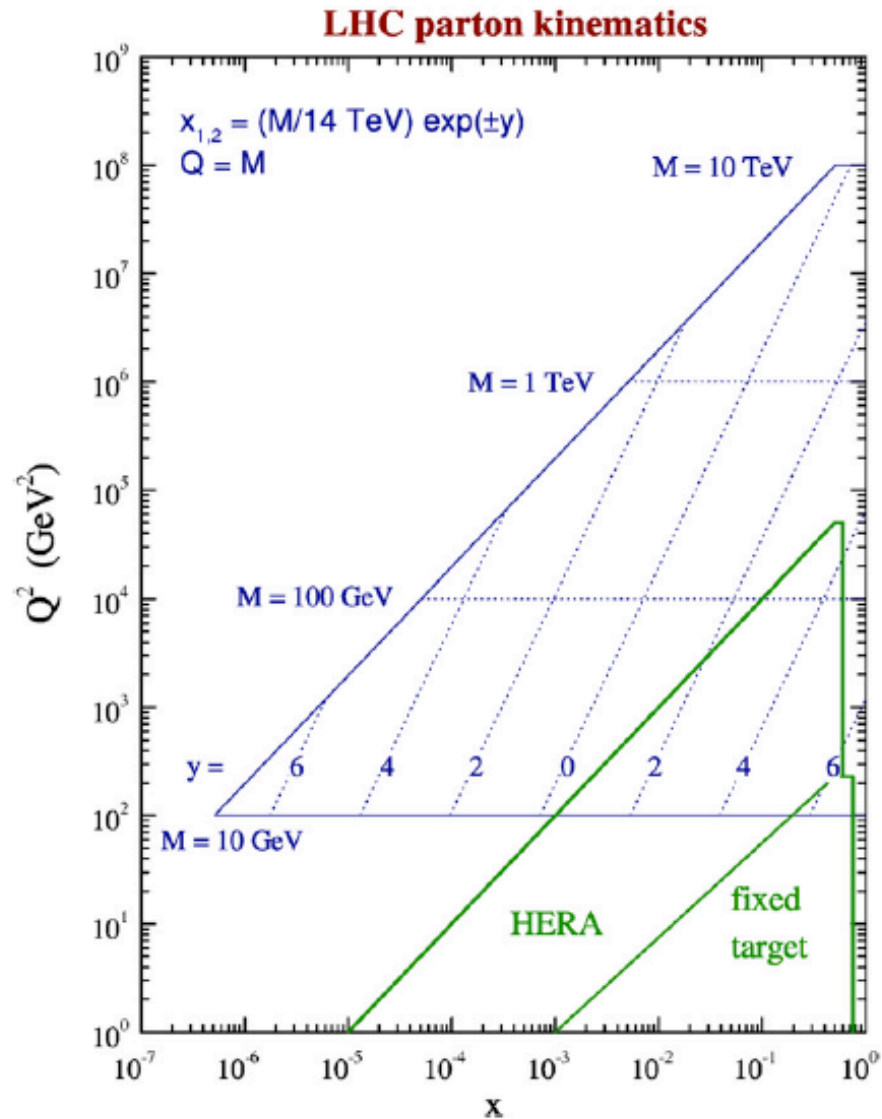


- Very important to reduce cross section uncertainties at hadron colliders; but still not good enough ($\sim 10\%$ errors for LHC cross sections)

Q^2 evolution following the DGLAP equation



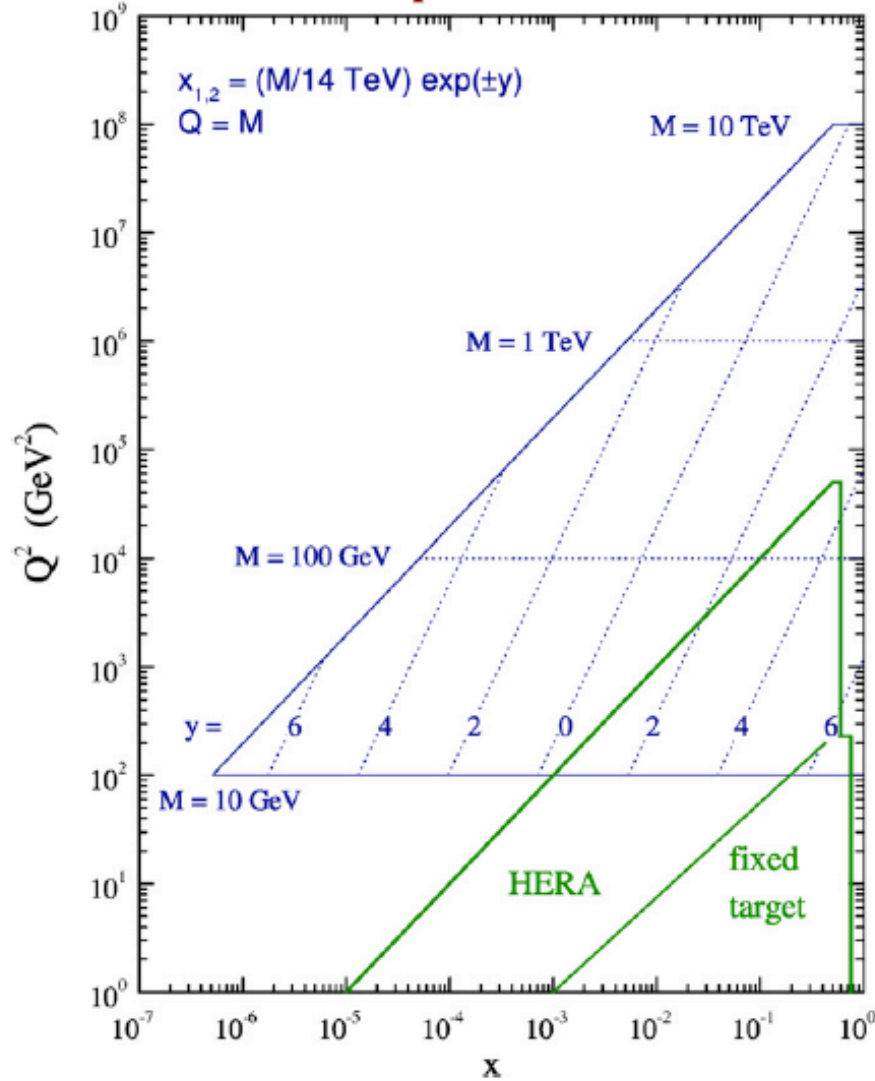
Distributions of x times the unpolarized parton distributions $f(x)$, where $f = u_v, d_v, \bar{u}, \bar{d}, s, b, g$ and their associated uncertainties using the NNLO MRST2008 parametrization at a scale $\mu^2 = 10 \text{ GeV}^2$ and $\mu^2 = 10,000 \text{ GeV}^2$.



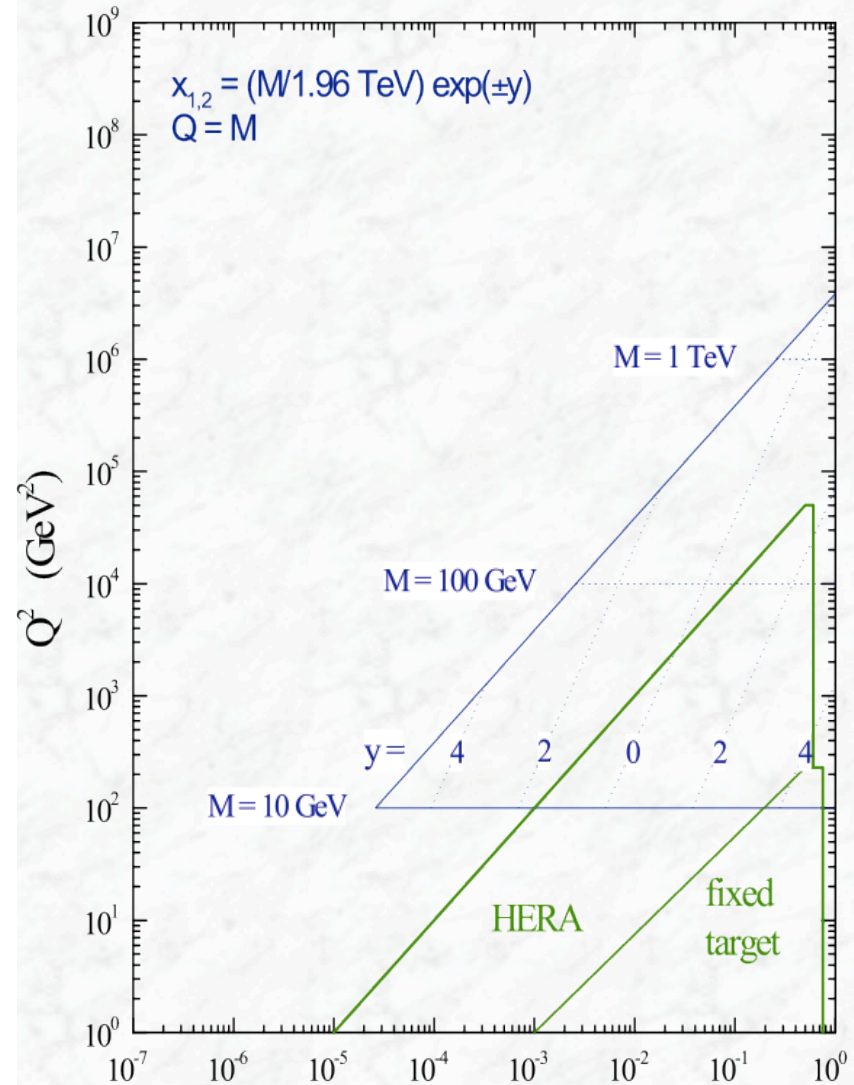
Graphical representation of the relationship between parton (x , Q^2) variables and the kinematic variables corresponding to a final state of mass M with rapidity y at the LHC with $\sqrt{s} = 14 \text{ TeV}$

Comparison between the Tevatron and the LHC (14 TeV)

LHC parton kinematics

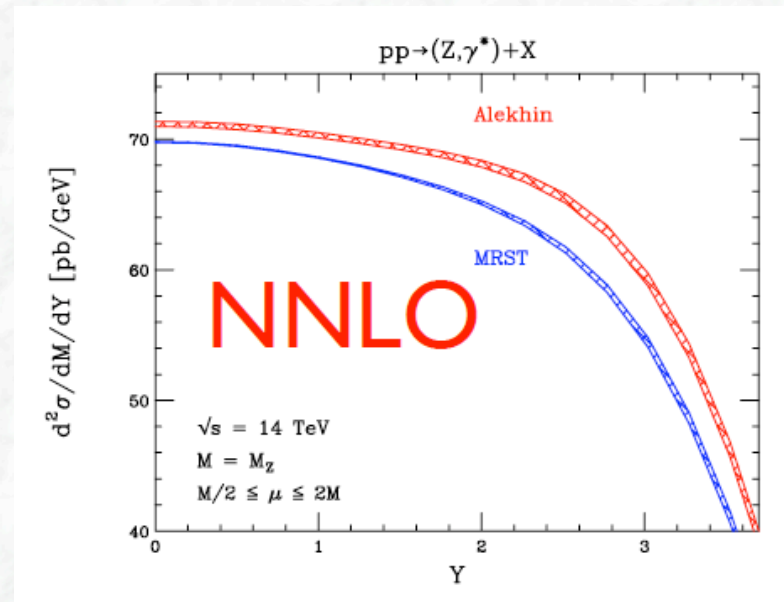
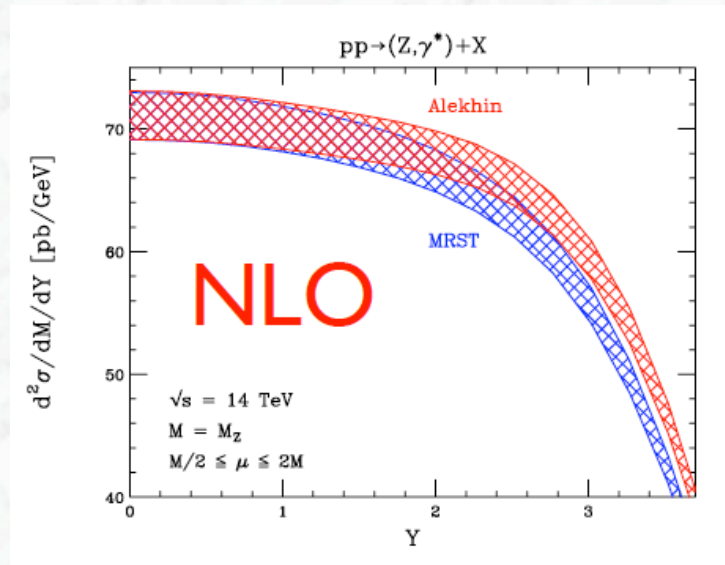
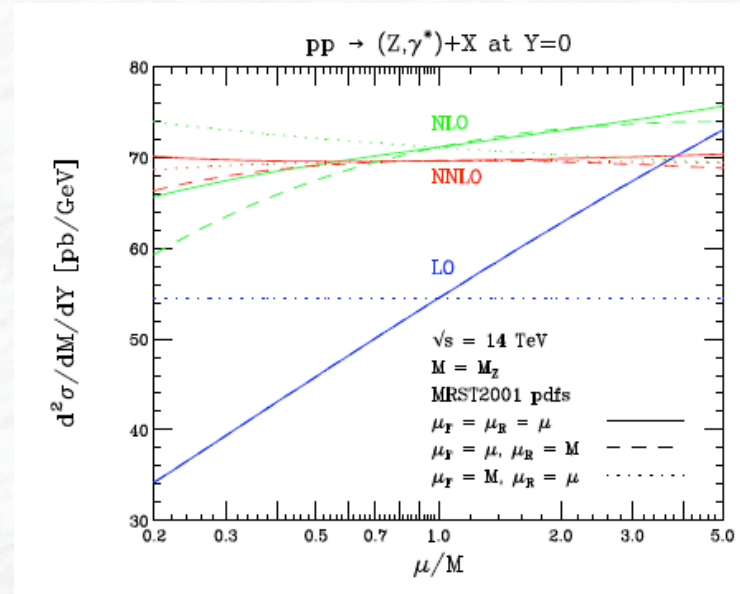
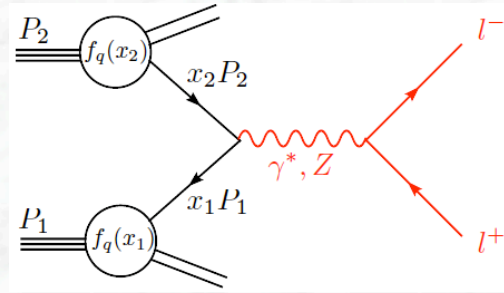


Tevatron parton kinematics



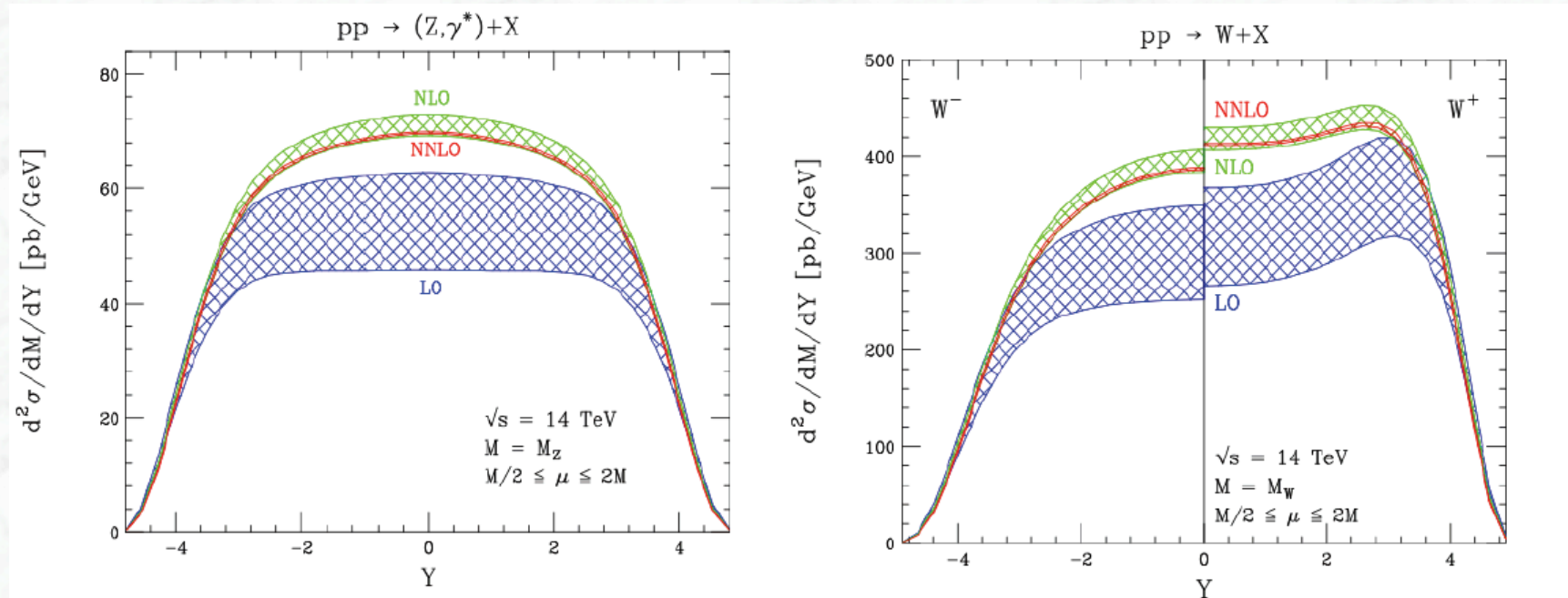
For the same masses (e.g. 100 GeV): x -values about 10 times lower at the LHC

Example: Drell-Yan production of W/Z bosons



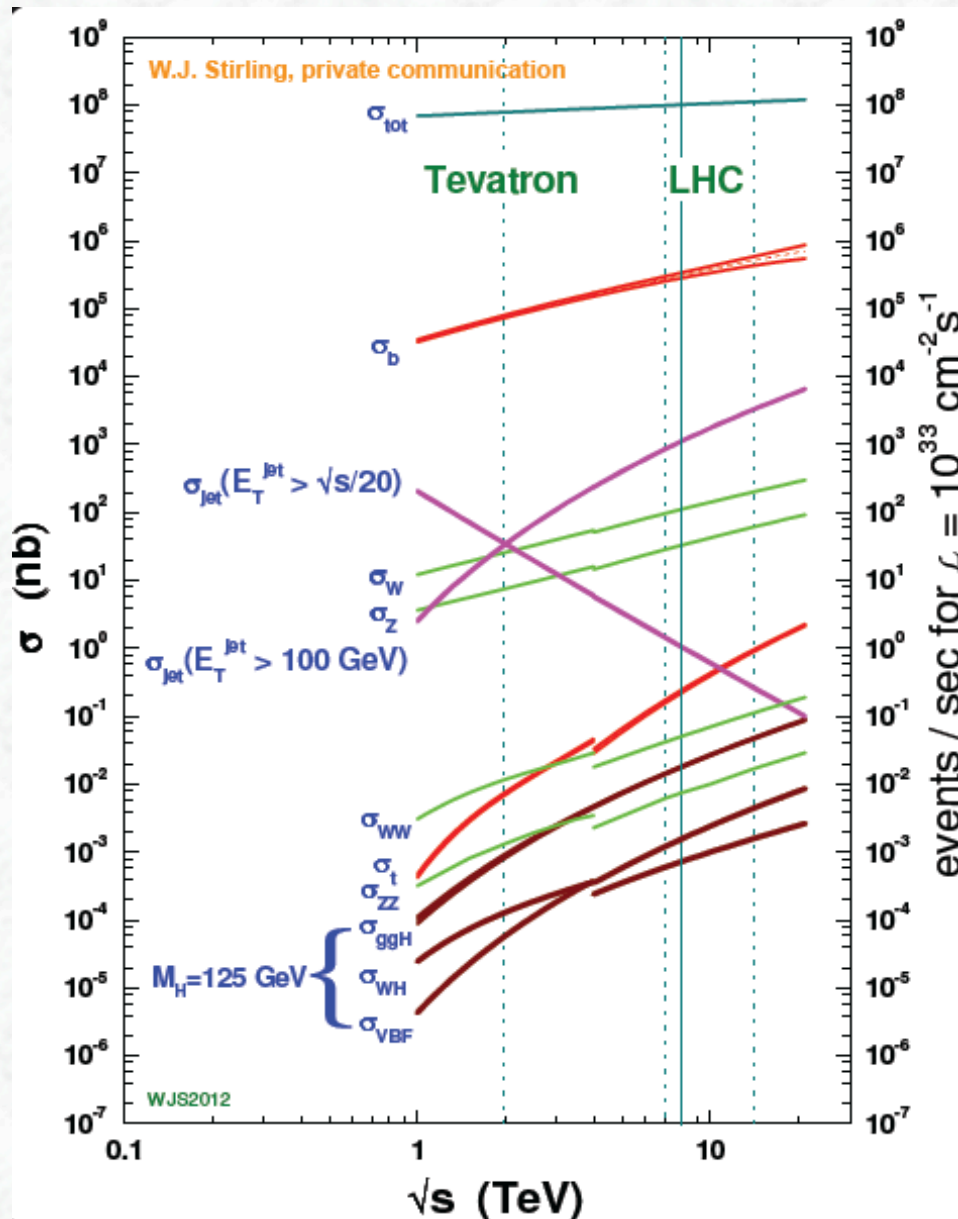
Example: Drell-Yan production of W/Z bosons (cont.)

Rapidity distributions for Z and W^\pm production at LO, NLO, and NNLO



Note: LHC data will be used in the future to further constrain the parton densities

Cross Sections and Production Rates

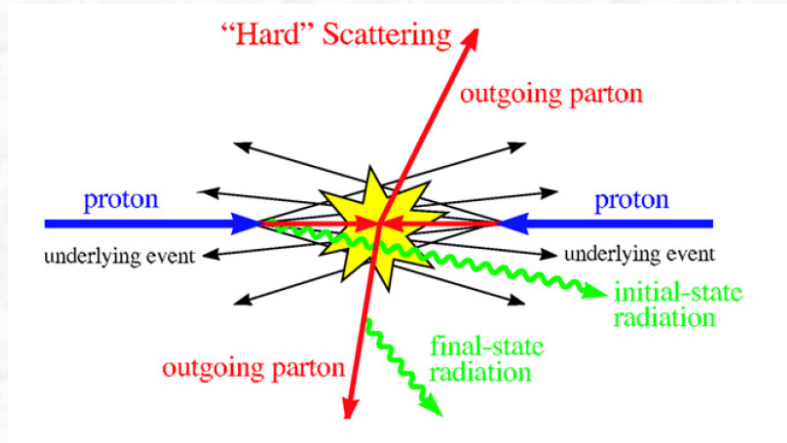


LHC is a factory for:
top-quarks, b-quarks, W, Z, ..., Higgs, ...

but other more prominent processes
dominate the production rates:

- Jet production via QCD scattering
- Soft pp collisions ($\sigma \sim 100 \text{ mb}$)

1.4 Soft proton-proton interactions



- First physics at the LHC was dominated by large cross section of inelastic hadronic interactions
- Most interactions are due to **interactions at large distance** between incoming protons → **small momentum transfer**, particles in the final state have large longitudinal, but small transverse momentum

$\langle p_T \rangle \approx 600 \text{ MeV}$ (of charged particles in the final state)



- Measurements necessary to constrain phenomenological models of soft-hadronic interactions and to predict properties at higher centre-of-mass energies (underlying event, pile-up of minimum bias events at high luminosity,)



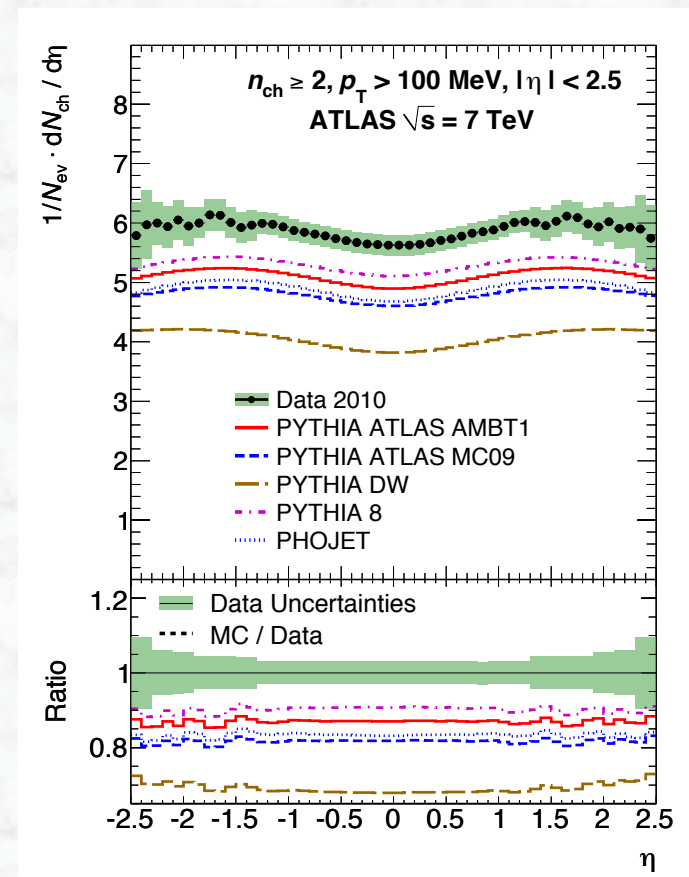
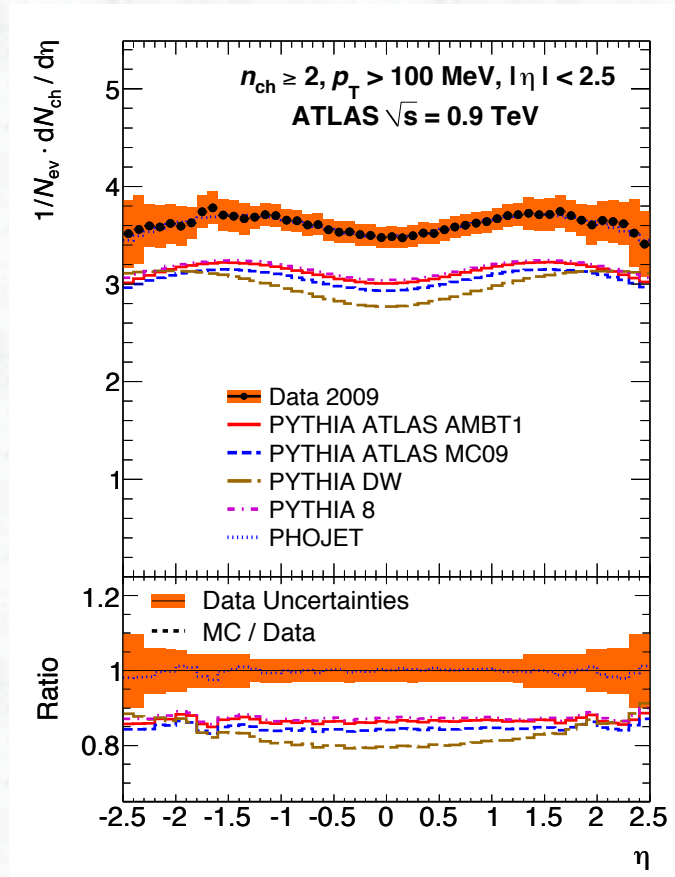
Charged particle density versus η

N_{ch} : number of primary charged particles corrected to particle level, normalized to the number of selected events N_{ev}

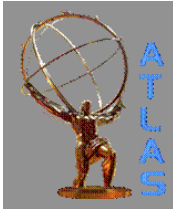
0.9 TeV

and

7 TeV data

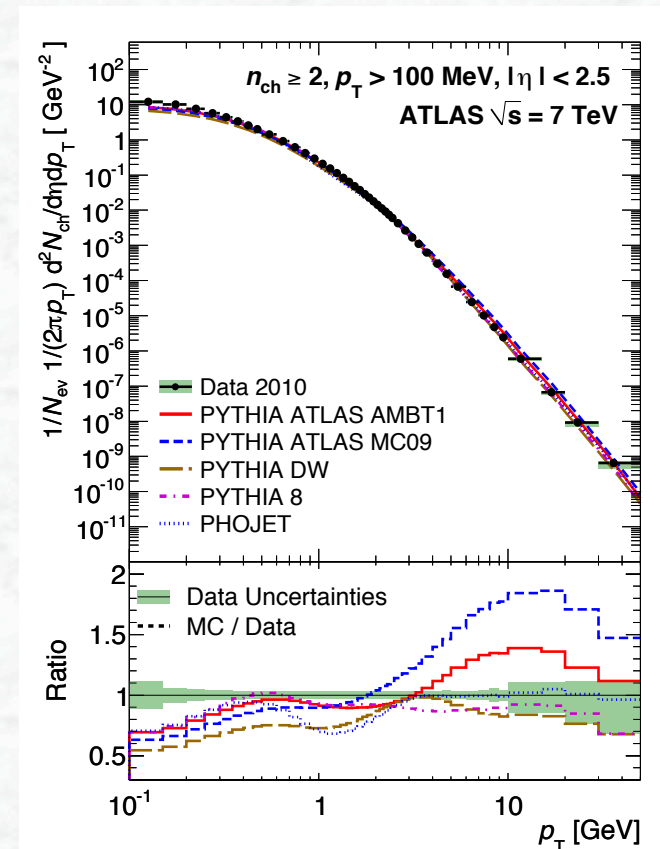
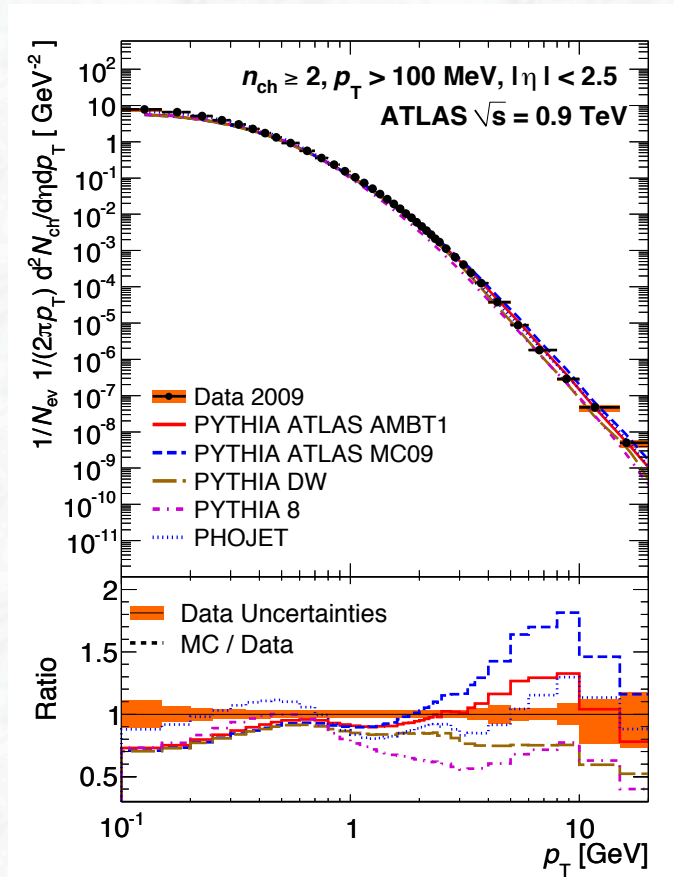


Various Monte Carlo models fail to describe the ATLAS data at both collider energies \rightarrow tuning of Monte Carlo parameters needed



Charged particle multiplicities as function of p_T

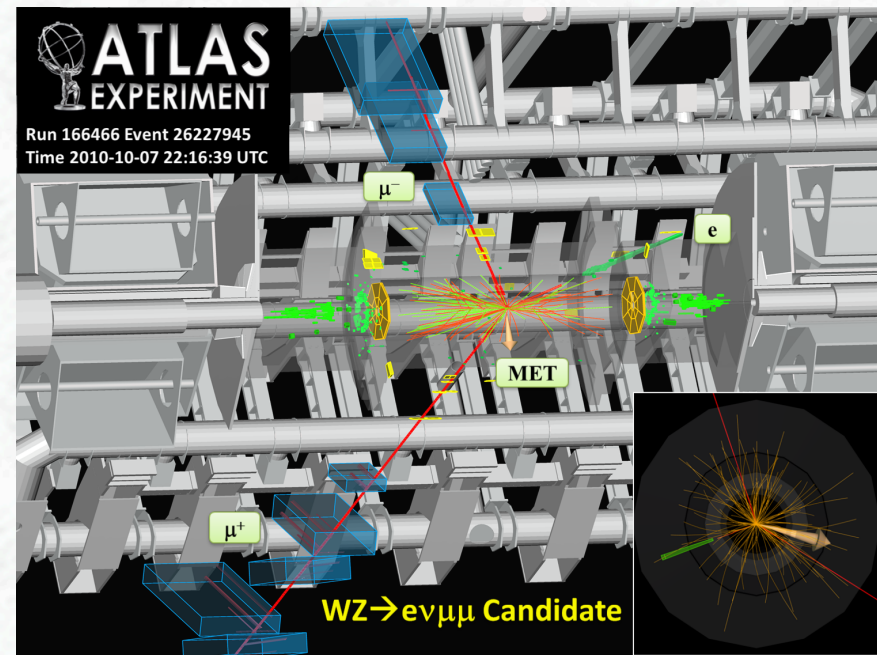
N_{ch} : number of primary charged particles corrected to particle level, normalized to the number of selected events N_{ev}



Monte Carlo models also fail to describe the p_T spectrum

Part 2: Test of perturbative QCD

- Jet production
- W/Z production
- Production of top quarks



It is important to establish the Standard Model reference processes:

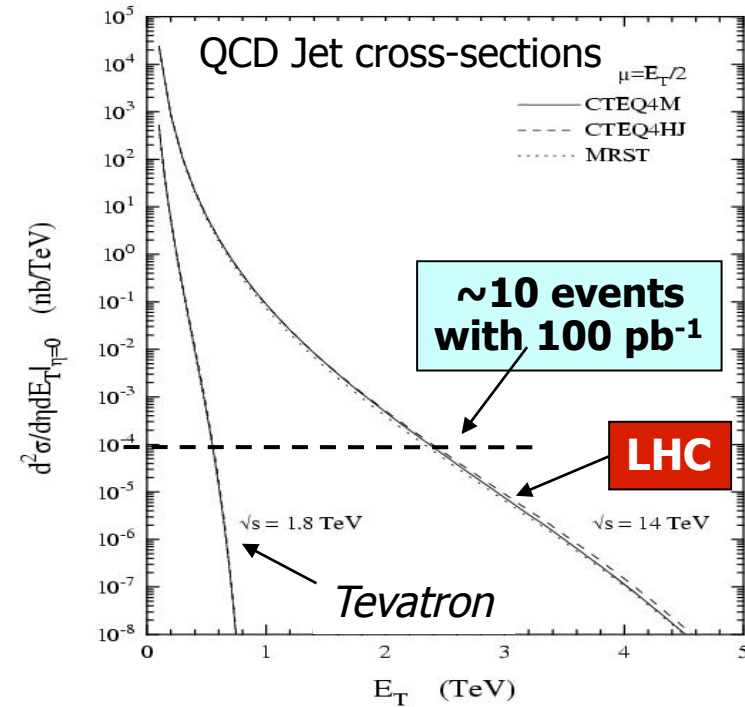
- Test of the theory itself
Deviations → evidence for Physics beyond the Standard Model
- Important to understand the detector performance
→ understand the so called “Fake” or “instrumental” background,
in particular for leptons (e,μ) and E_T^{miss}
- Standard Model processes are important background processes for many searches for Physics Beyond the Standard Model
“Physics Background”

Typical selections require: leptons, jets, E_T^{miss} ,

→ W/Z + jets and tt productions are omnipresent !

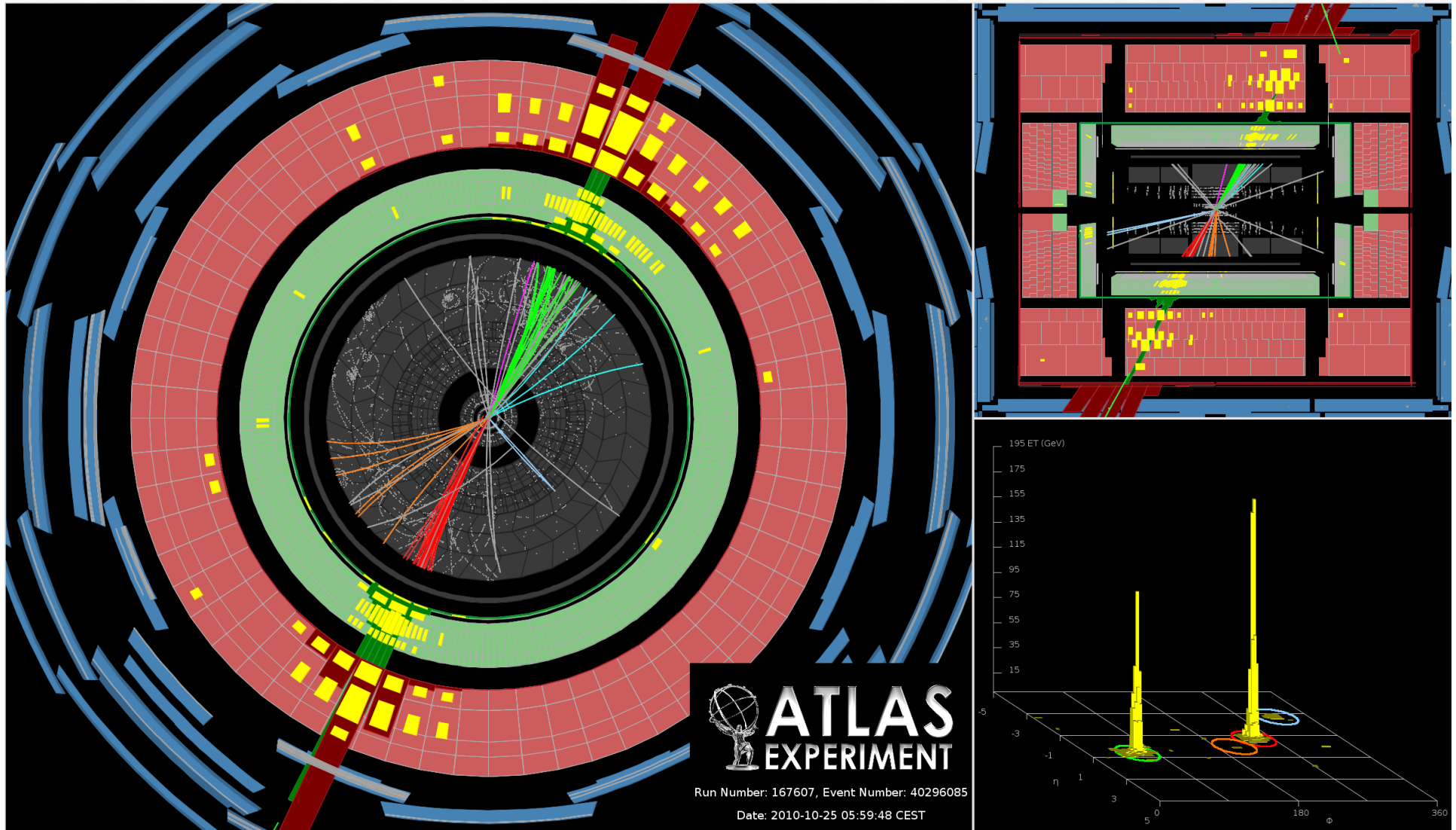
2.1 Jets from QCD production

- Rapidly probe perturbative QCD in a new energy regime (at a scale above the Tevatron, large cross sections)
- **Experimental challenge:** understanding of the detector
 - main focus on **jet energy scale**
 - resolution
- **Theory challenge:**
 - improved calculations... (renormalization and factorization scale uncertainties)
 - pdf uncertainties



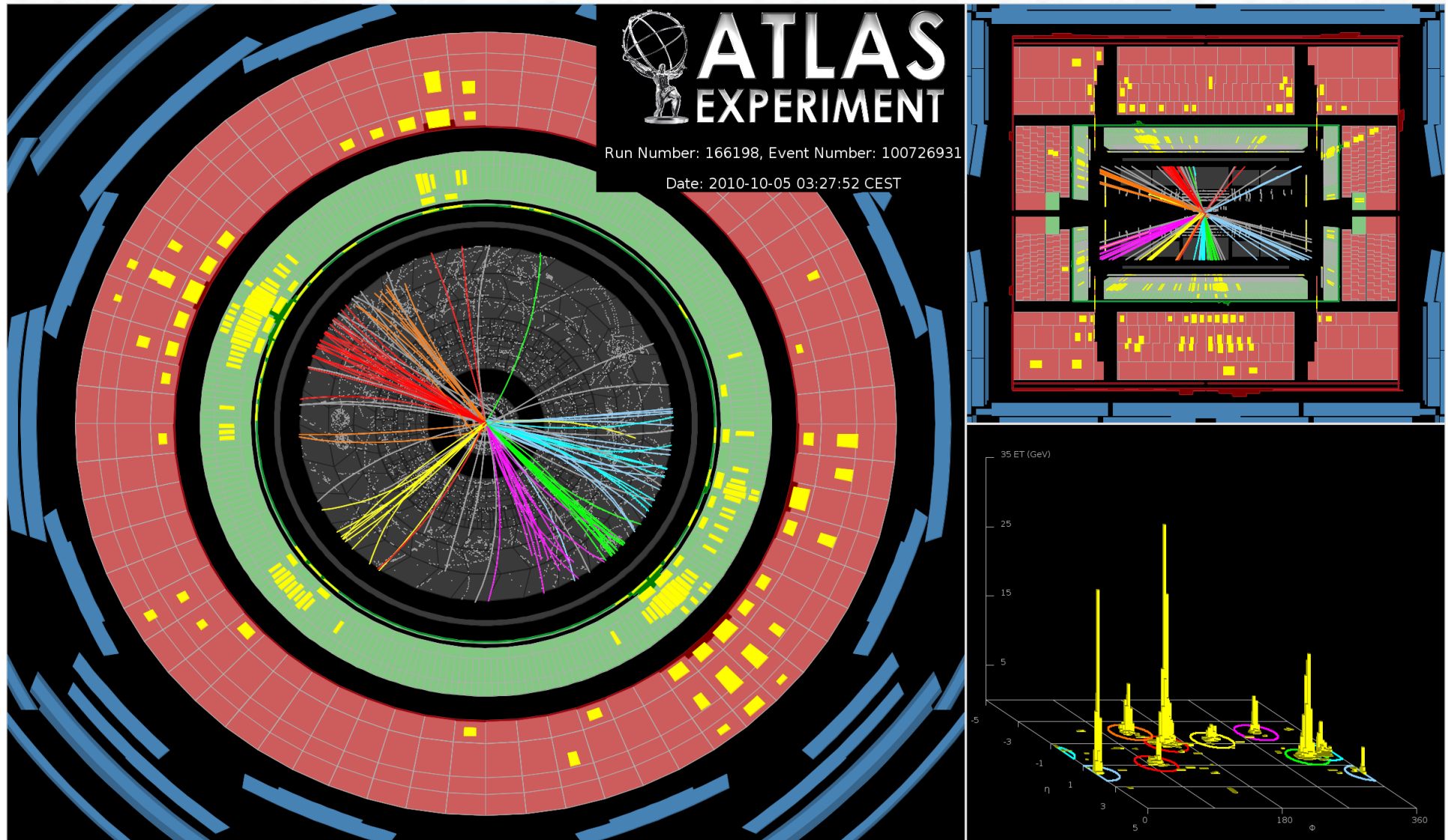
A comparison between the Tevatron and the LHC (14 TeV)

High p_T jet events at the LHC



Event display that shows the highest-mass central dijet event collected during 2010, where the two leading jets have an invariant mass of 3.1 TeV. The two leading jets have (p_T, y) of (1.3 TeV, -0.68) and (1.2 TeV, 0.64), respectively. The missing E_T in the event is 46 GeV. From [ATLAS-CONF-2011-047](#).

An event with a high jet multiplicity at the LHC

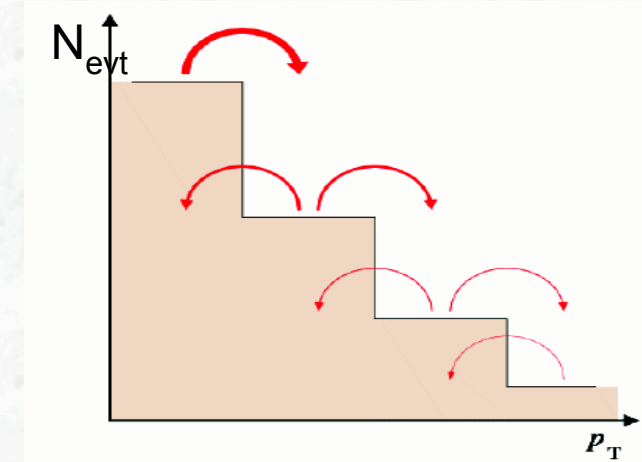
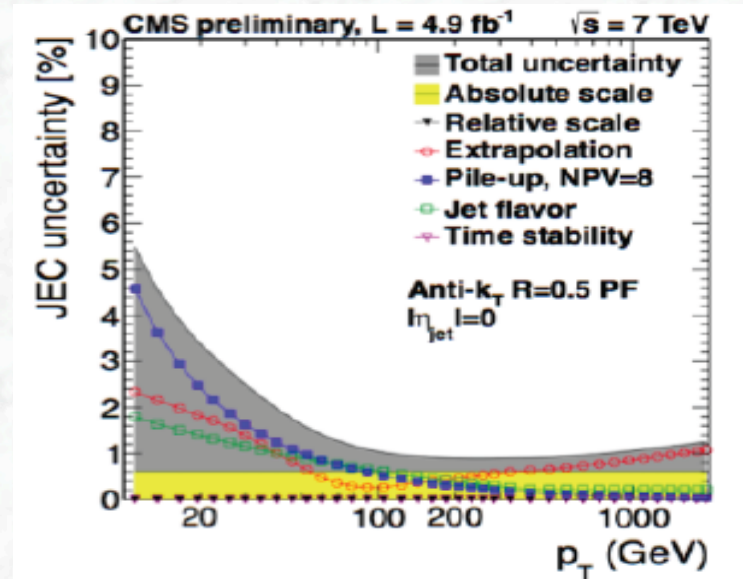
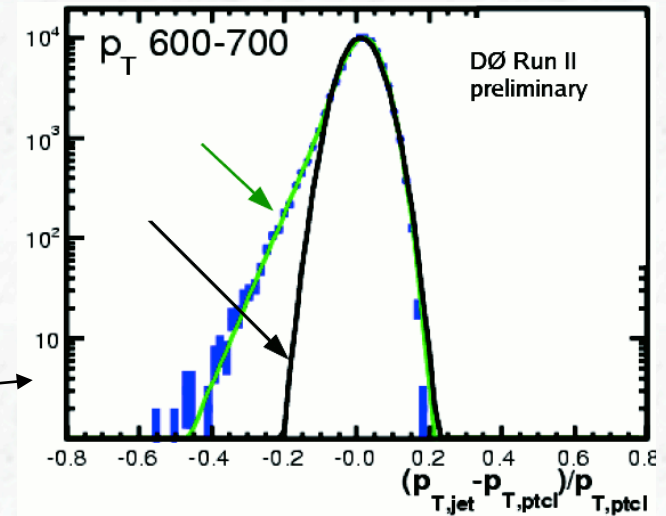


The highest jet multiplicity event collected, counting jets with p_T greater than 60 GeV: this event has eight. 1st jet (ordered by p_T): $p_T = 290$ GeV, $\eta = -0.9$, $\phi = 2.7$; 2nd jet: $p_T = 220$ GeV, $\eta = 0.3$, $\phi = -0.7$ Missing $E_T = 21$ GeV, $\phi = -1.9$, Sum $E_T = 890$ GeV.

Jet measurements

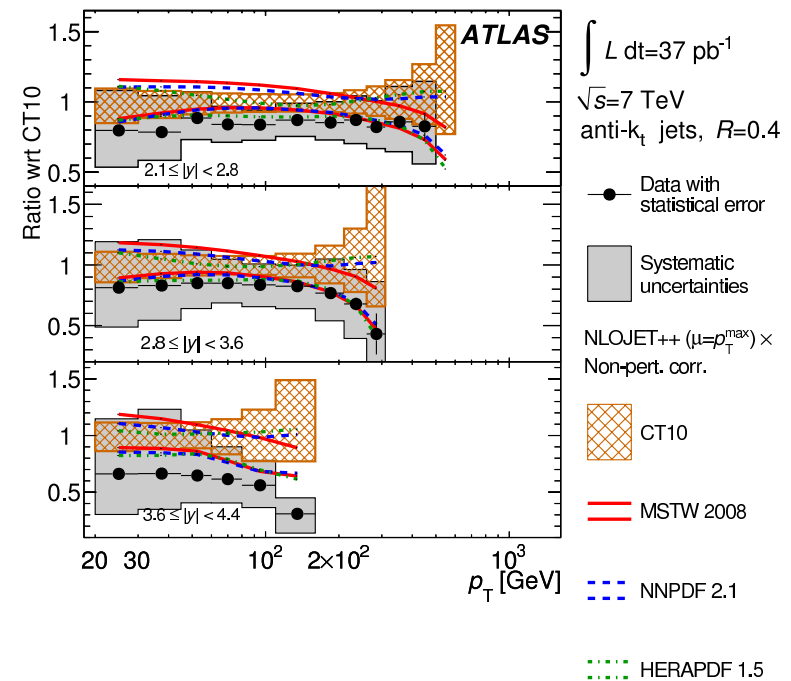
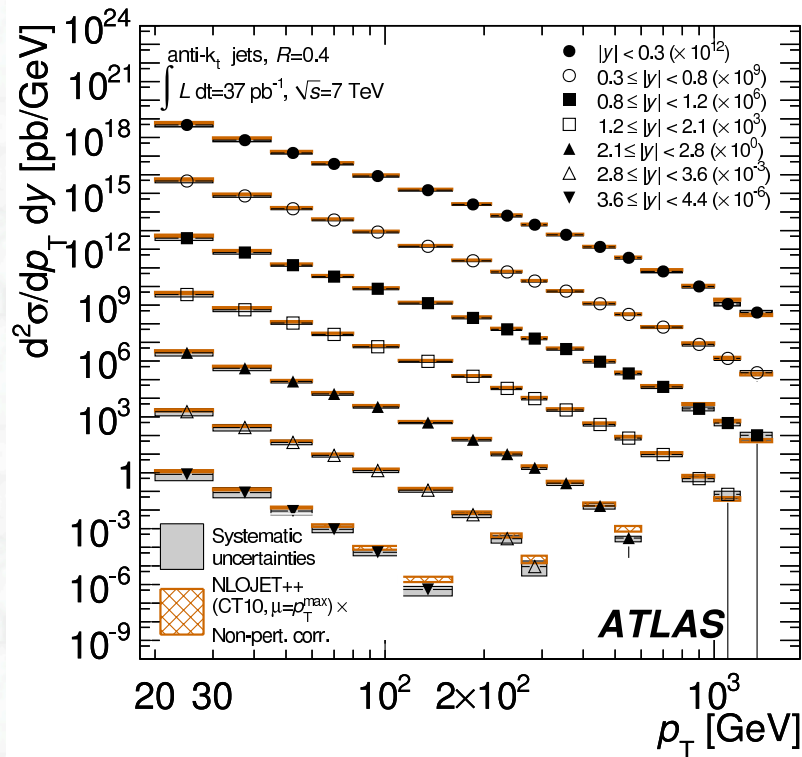
$$d^2\sigma / dp_T d\eta = N / (\epsilon \cdot L \cdot \Delta p_T \cdot \Delta\eta)$$

- In principle a simple counting experiment
- However, steeply falling p_T spectra are sensitive to jet energy scale uncertainties and resolution effects (migration between bins) → corrections (unfolding) to be applied
- Jet energy scale uncertainty:
 CMS: ~1.5 - 3% (after two years)
 (similar for ATLAS, impressive achievements)





Double differential cross sections, as function of p_T and rapidity y (full 2010 data set)

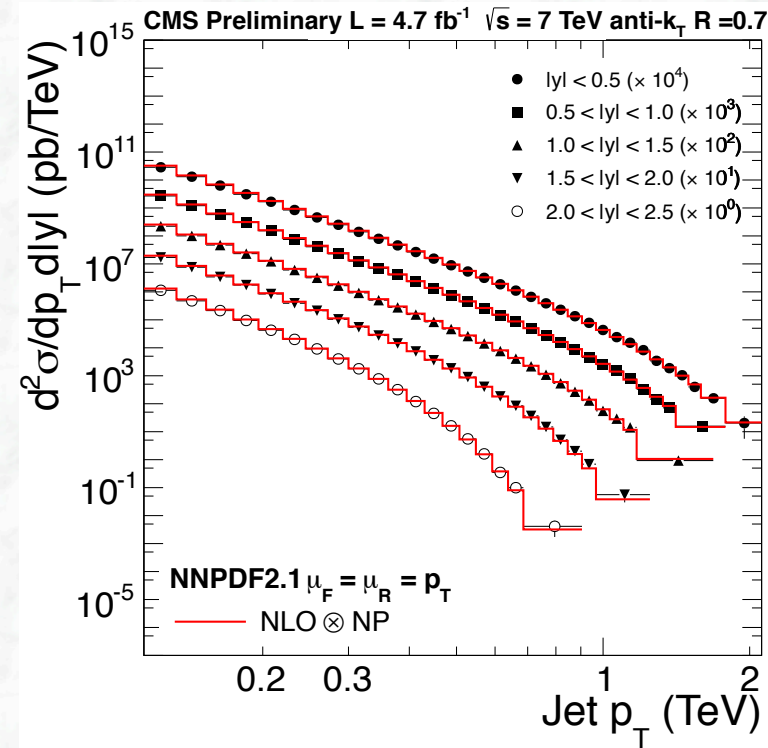
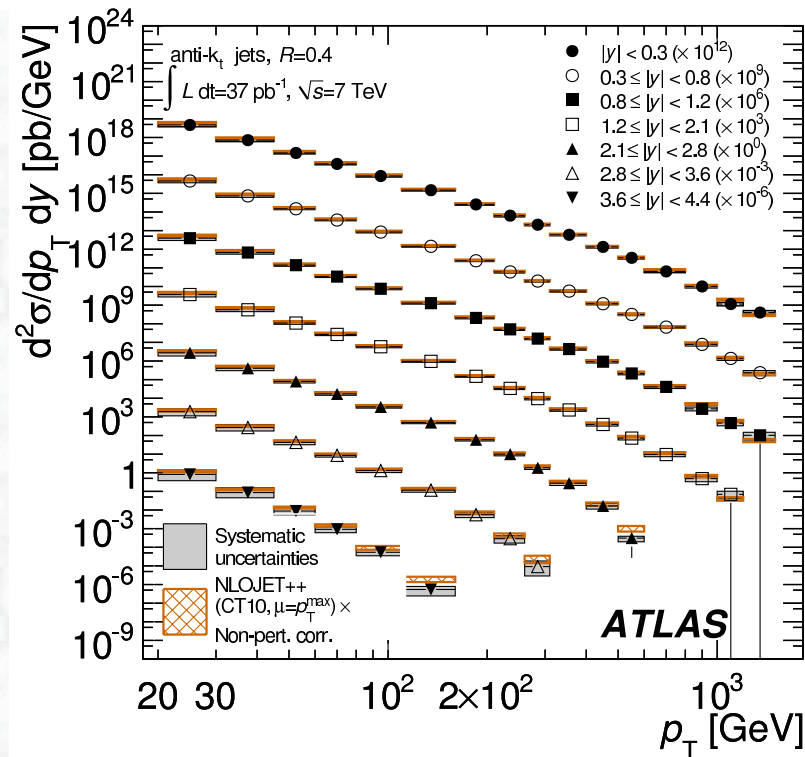


somewhat larger deviations in the forward region

- Data are well described by NLO pert. QCD calculations (NLOJet++)
- Experimental systematic uncertainty is dominated by jet energy scale uncertainty
- Theoretical uncertainties: renormalization/ factorization scale, pdfs, α_s , ..., uncertainties from non-perturbative effects



Double differential cross sections, as function of p_T and rapidity y : (full 2010 data set)

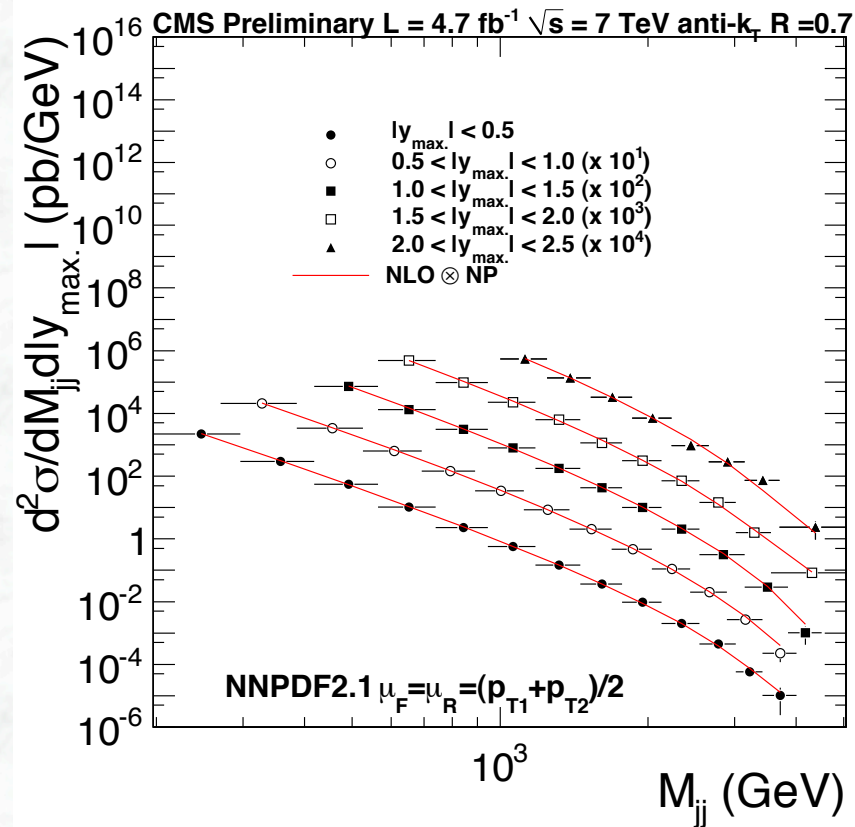
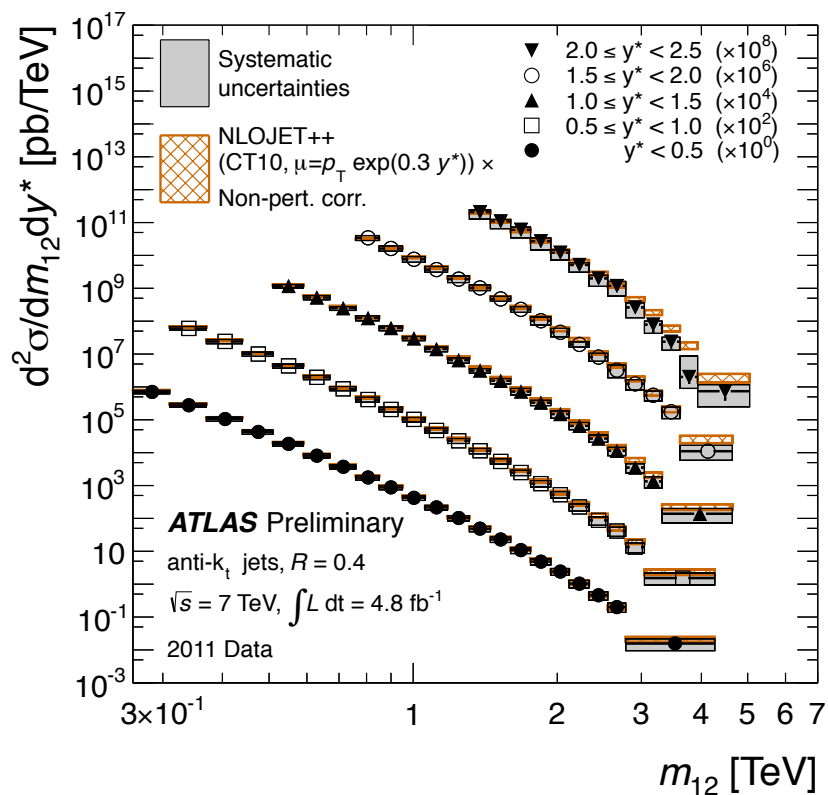
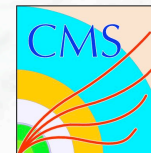


CMS: include full 2011 data set;
comparison up to 2 TeV (central rapidities)

- Data are well described by NLO pert. QCD calculations (NLOJet++)
- Experimental systematic uncertainty is dominated by jet energy scale uncertainty
- Theoretical uncertainties: renormalization/ factorization scale, pdfs, α_s , ..., uncertainties from non-perturbative effects



Invariant di-jet mass spectra

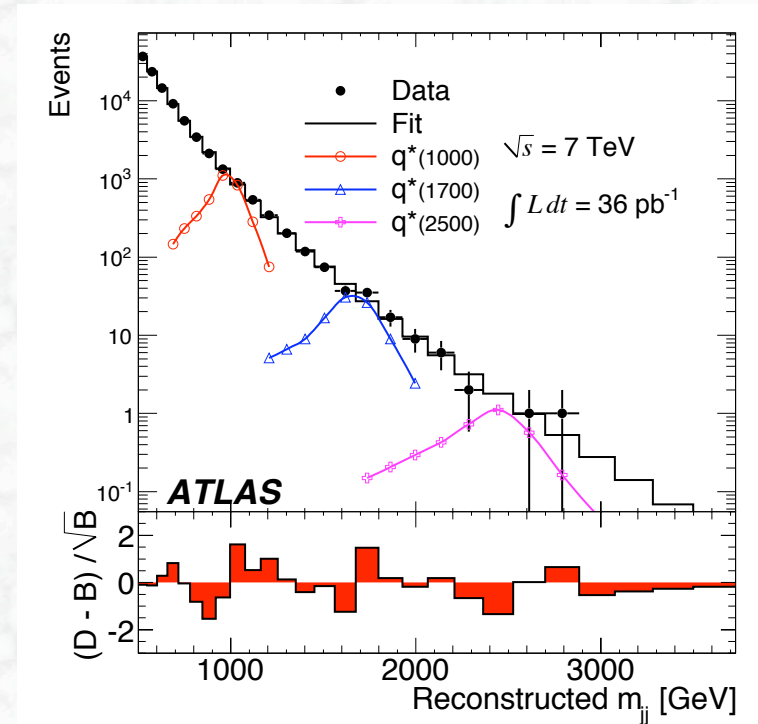


- Important for:
- Test of QCD
 - Search for new resonances decaying into two jets (\rightarrow next slide)



In addition to QCD test: Sensitivity to New Physics

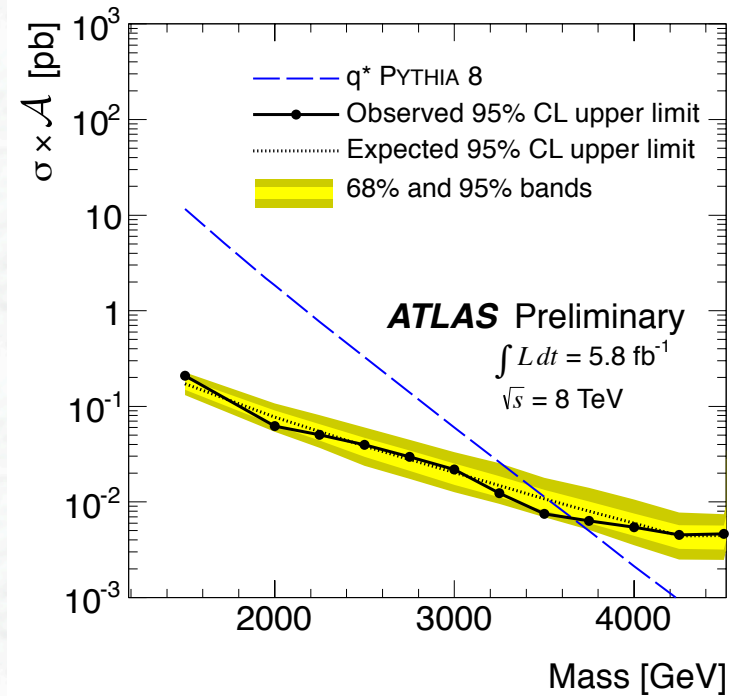
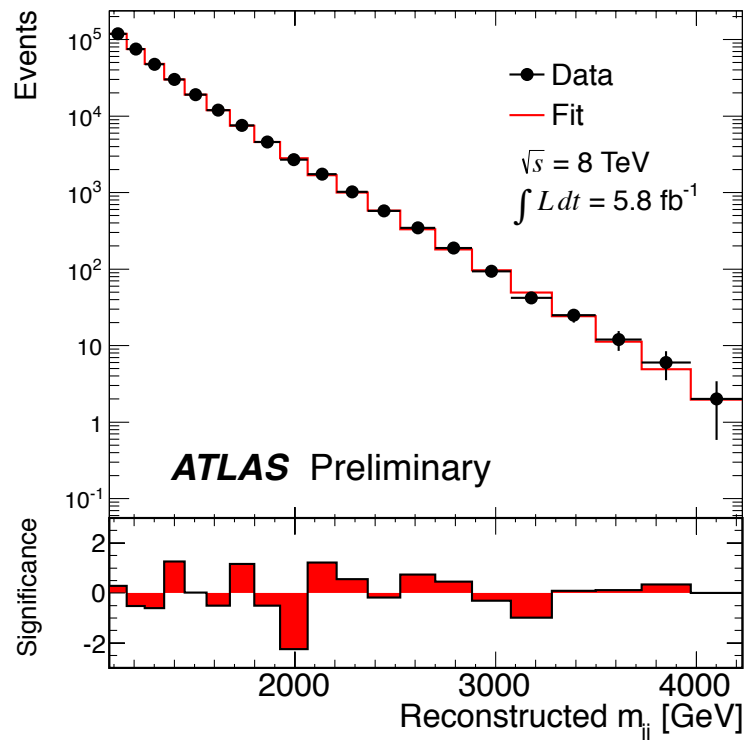
- Di-jet mass spectrum provides large sensitivity to new physics
e.g. Resonances decaying into qq , excited quarks q^* ,
- Search for resonant structures in the di-jet invariant mass spectrum



CDF (Tevatron), $L = 1.13 \text{ fb}^{-1}$:	$0.26 < m_{q^*} < 0.87 \text{ TeV}$
ATLAS (LHC), $L = 0.000315 \text{ fb}^{-1}$	exclude (95% C.L) q^* mass interval
$L = 0.036 \text{ fb}^{-1}$:	$0.30 < m_{q^*} < 1.26 \text{ TeV}$
	$0.60 < m_{q^*} < 2.64 \text{ TeV}$



- Include new data at $\sqrt{s} = 8$ TeV (2012)
- Invariant di-jet masses up to 4.1 TeV



CDF (Tevatron), $L = 1.13 \text{ fb}^{-1}$:

$0.26 < m_{q^*} < 0.87 \text{ TeV}$

ATLAS (LHC), $L = 0.000315 \text{ fb}^{-1}$

exclude (95% C.L) q^* mass interval

$0.30 < m_{q^*} < 1.26 \text{ TeV}$

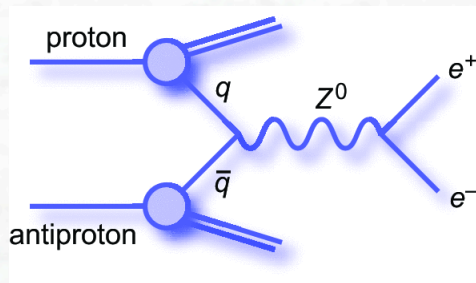
$L = 0.036 \text{ fb}^{-1}$:

$0.60 < m_{q^*} < 2.64 \text{ TeV}$

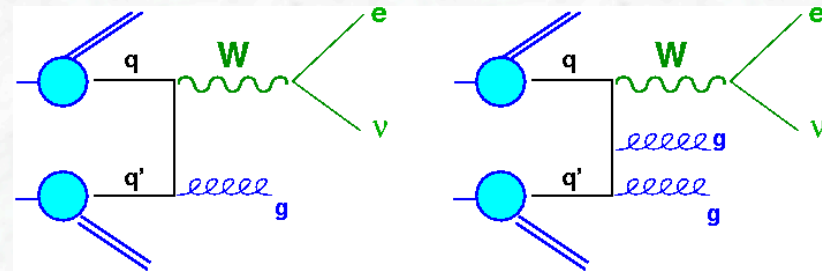
ATLAS (LHC), $L = 5.8 \text{ fb}^{-1}, 8 \text{ TeV}$:

$m_{q^*} < 3.66 \text{ TeV}$

2.2 QCD aspects in W/Z (+ jet) production



QCD at work

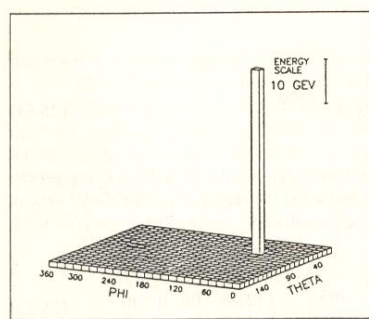


- Important test of NNLO Drell-Yan QCD prediction for the total cross section
- Test of perturbative QCD in high p_T region (jet multiplicities, p_T spectra,....)
- Tuning and „calibration“ of Monte Carlos for background predictions in searches at the LHC

How do W and Z events look like ?

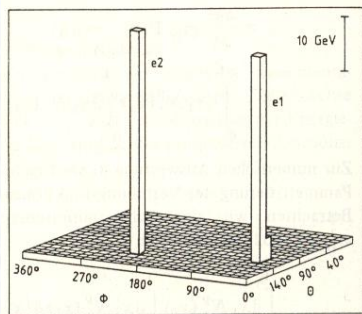
As explained, leptons, photons and missing transverse energy are key signatures at hadron colliders

→ Search for leptonic decays: $W \rightarrow \ell \nu$ (large $P_T(\ell)$, large E_T^{miss})
 $Z \rightarrow \ell \ell$

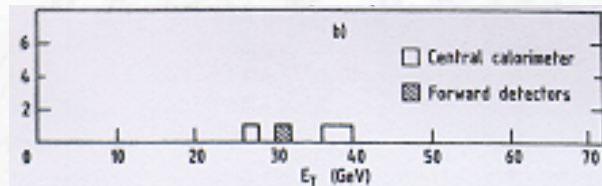


A bit of history: one of the first W events seen; UA2 experiment

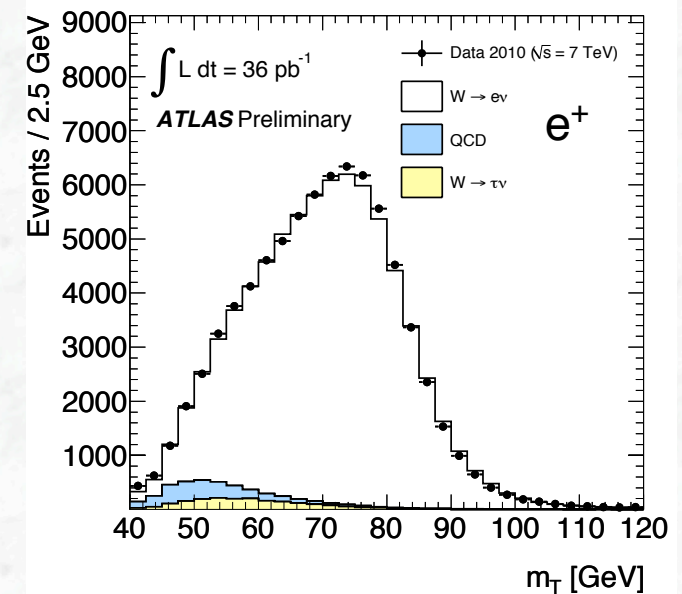
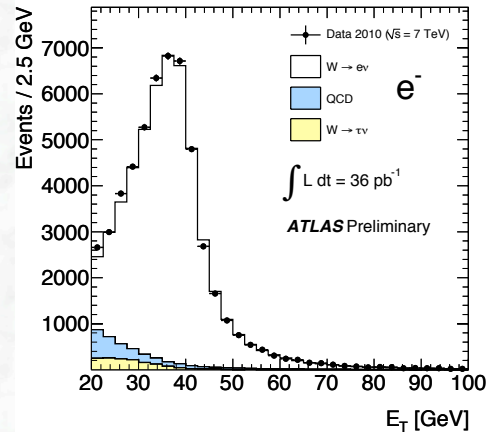
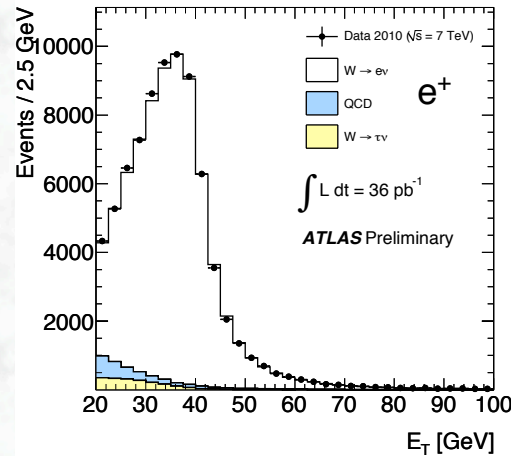
W/Z discovery by the UA1 and UA2 experiments at CERN (1983/84)



Transverse momentum of the electrons



W/Z selections in the ATLAS / CMS experiments



Electrons:

- Trigger: high p_T electron candidate in calorimeter
- Isolated el.magn. cluster in the calorimeter
- $P_T > 25 \text{ GeV}/c$
- Shower shape consistent with expectation for electrons
- Matched with tracks

Z $\rightarrow ee$

- $76 \text{ GeV}/c^2 < m_{ee} < 106 \text{ GeV}/c^2$

W $\rightarrow e\nu$

- Missing transverse momentum $> 25 \text{ GeV}/c$
- Transverse mass cut $M_T > 50 \text{ GeV}$

$$M_W^T = \sqrt{2 \cdot P_T^l \cdot P_T^\nu \cdot (1 - \cos \Delta\phi^{l,\nu})}$$

Transverse mass
(longitudinal component of the
neutrino cannot be measured)

Ingredients for cross-section measurements

$$\sigma_{W(Z)}^{\text{tot}} \cdot BR(W(Z) \rightarrow \ell\nu (\ell\ell)) = \frac{N_{W(Z)}^{\text{sig}}}{A_{W(Z)} \cdot C_{W(Z)} \cdot L_{W(Z)}}$$

- Number of W/Z signal candidates $N^{\text{sig}} = N^{\text{evt}} - N^{\text{back}}$
Estimated background (Physics background, “fake” background,...)
- $C_{W(Z)}$: reconstruction efficiencies, detector effects, ...
- $A_{W(Z)}$: acceptance (usually the final state products are measured in a so called fiducial region of the detector,
e.g. η coverage of the muon detector, p_T threshold of the reconstruction)

This last quantity can only be calculated with Monte Carlo, using theoretical inputs !!

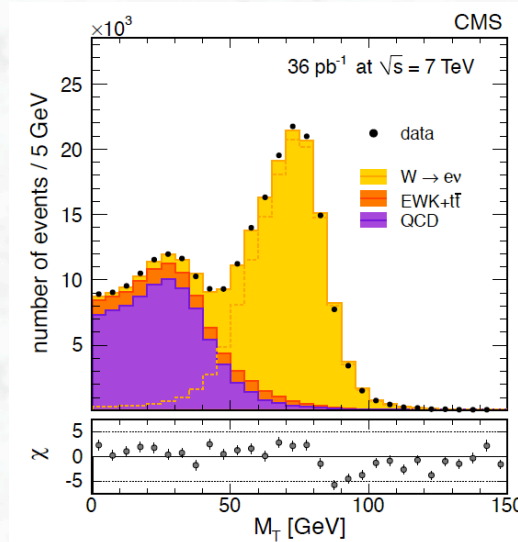
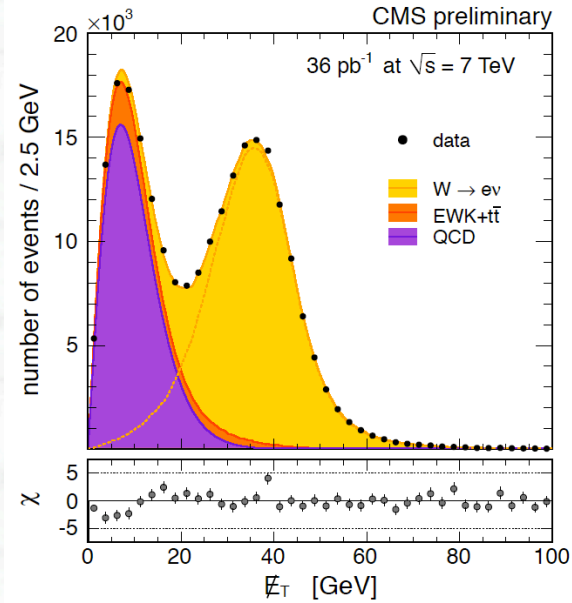
(N)NLO calculations, parton density functions,

- Cross sections for $A_{W(Z)} = 1$ are called “fiducial cross sections”
- Less affected by theoretical / pdf uncertainties...

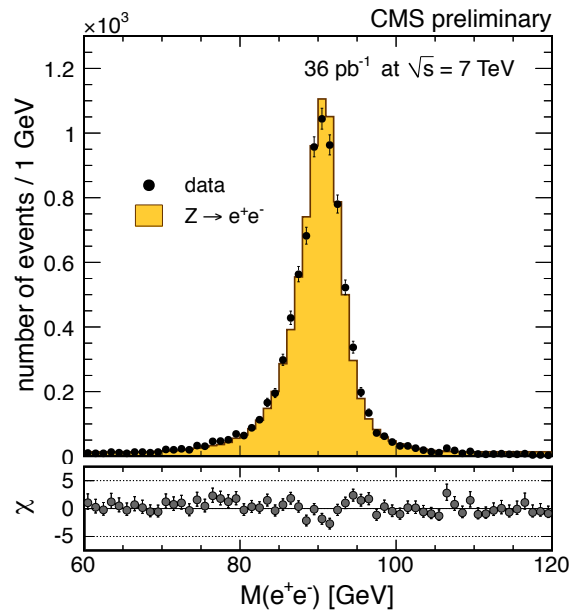
- $L_{W(Z)}$: integrated luminosity



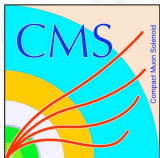
An example: CMS data from 2010: 36 pb⁻¹



Distributions of the missing transverse energy, E_T^{miss} , (left) and transverse mass m_T (right) of electron candidates for data and Monte Carlo simulation, broken down into the signal and various background components.

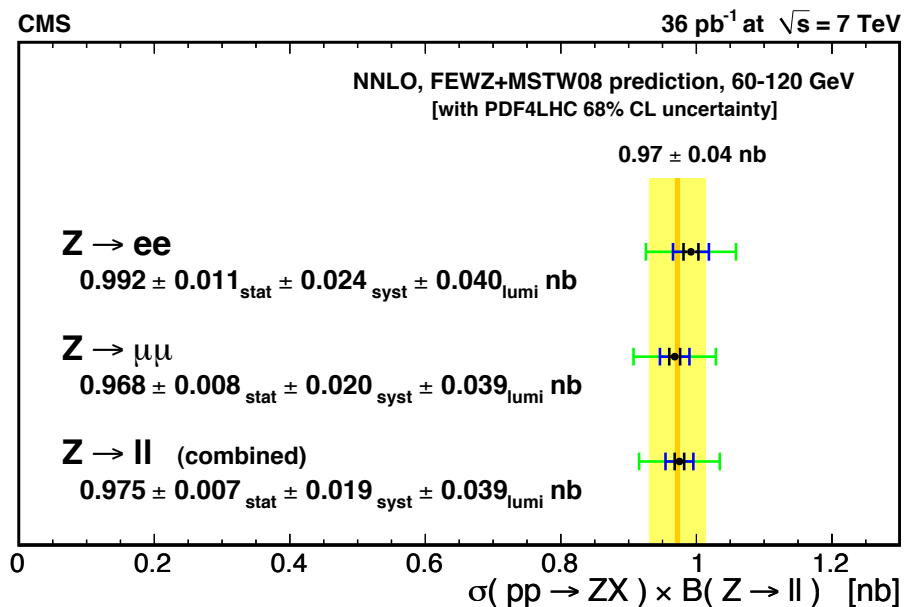
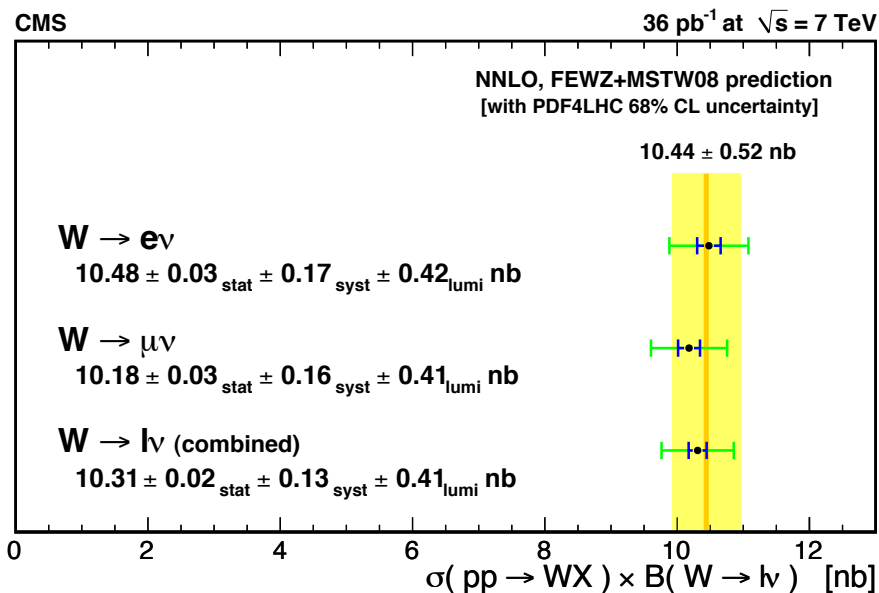


Distributions of the invariant di-electron mass, m_{ee} , for events passing the Z selection. The data are compared to Monte-Carlo simulation, the background is very small.



W and Z production cross sections at the LHC

Measured cross section values in comparison to NNLO QCD predictions:



Data are well described by NNLO QCD calculations

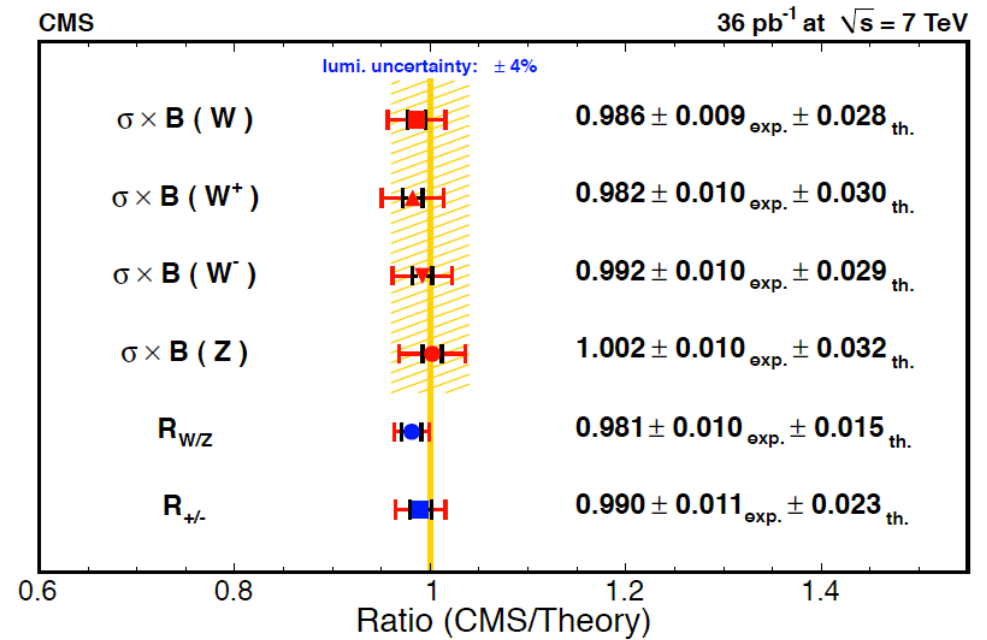
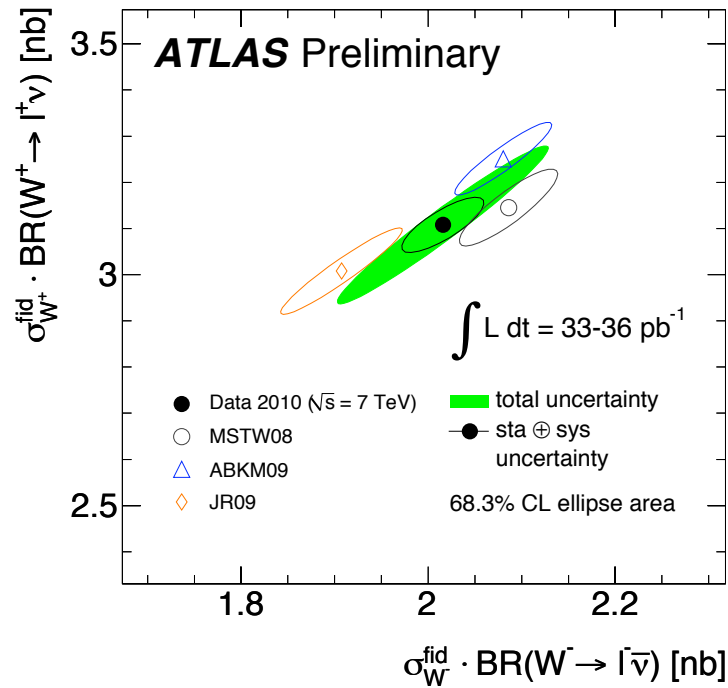
C.R.Hamberg et al, Nucl. Phys. B359 (1991) 343.

Precision is already dominated by systematic uncertainties

[The error bars represent successively the statistical, the statistical plus systematic and the total uncertainties (statistical, systematic and luminosity). All uncertainties are added in quadrature.]



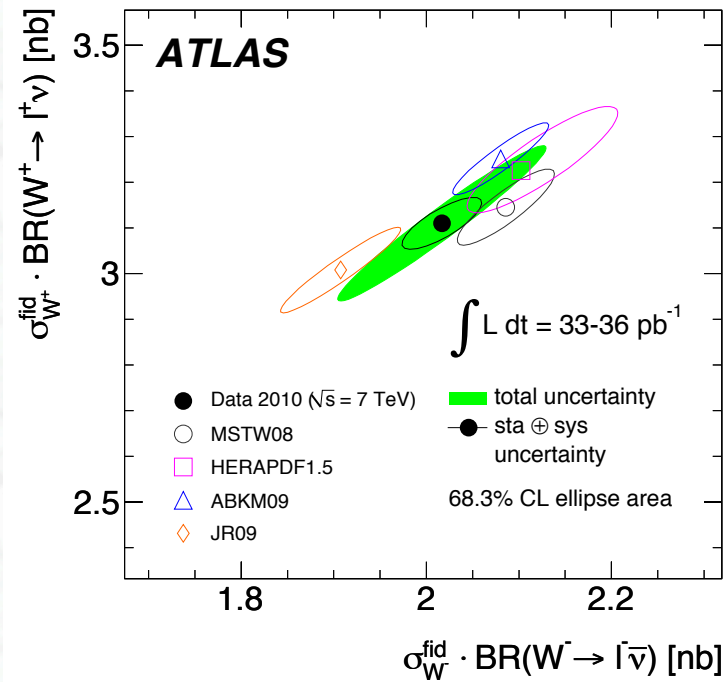
W cross sections at the LHC -charge separated, e/ μ universality



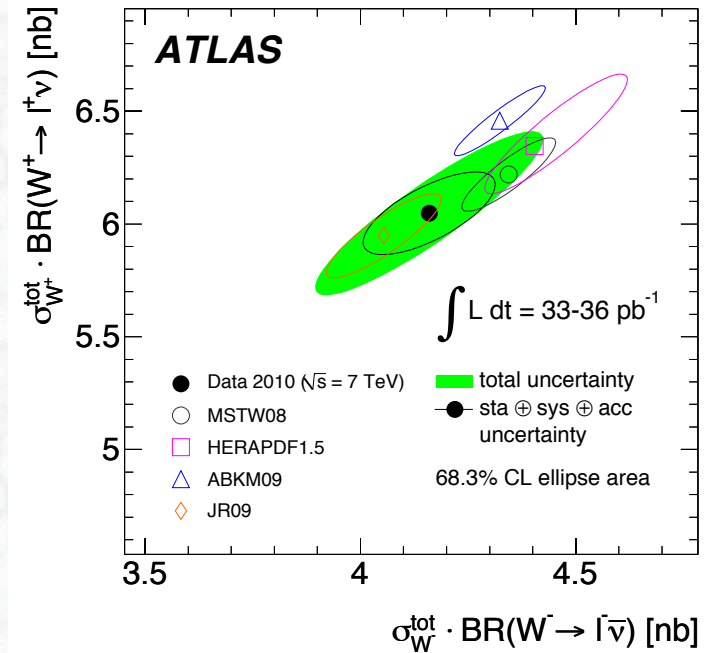
Good agreement between data and NNLO QCD predictions for all measurements

From fiducial cross sections to total cross sections

Fiducial cross sections



Total cross sections



$P_T(e) > 20 \text{ GeV}, \quad \eta < 2.5$
 $P_T(\nu) > 25 \text{ GeV}$
 $m_T(e \nu) > 40 \text{ GeV}$

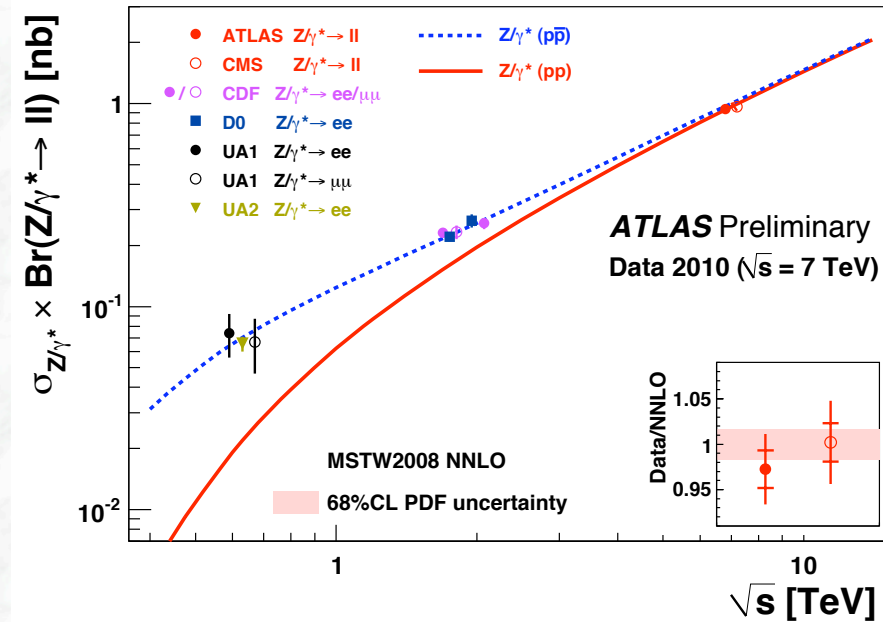
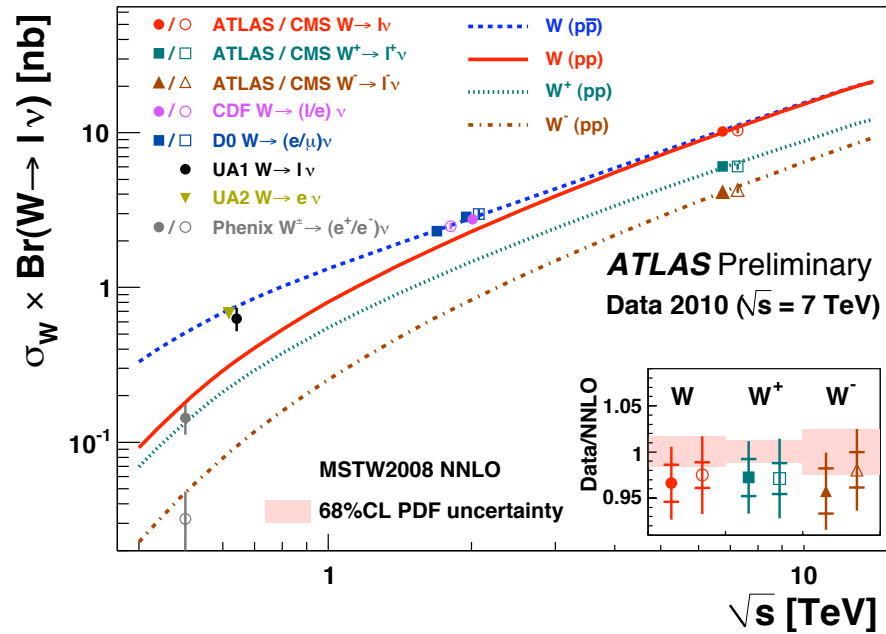
Uncertainties in W/Z cross section measurements

Electron channels (%)	W^\pm	W^+	W^-	Z
Trigger	0.4	0.4	0.4	<0.1
Reconstruction	0.8	0.8	0.8	1.6
Identification	0.9	0.8	1.1	1.8
Isolation	0.3	0.3	0.3	—
Energy scale and resolution	0.5	0.5	0.5	0.2
Defective LAr channels	0.4	0.4	0.4	0.8
Charge misidentification	<0.1	0.1	0.1	0.6
E_T^{miss}	0.8	0.7	1.0	—
Pile-up	0.3	0.3	0.3	0.3
Vertex position	0.1	0.1	0.1	0.1
QCD Background	0.4	0.4	0.4	0.7
EWK+ $t\bar{t}$ Background	0.2	0.2	0.2	<0.1
$C_{W/Z}$ Theor. uncertainty	0.6	0.6	0.6	0.3
Total Exp. uncertainty	1.8	1.8	2.0	2.7
$A_{W/Z}$ Theor. uncertainty	1.4	1.6	1.9	1.9
Total excluding Luminosity	2.3	2.4	2.8	3.3

In addition: luminosity uncertainty $\pm 3.4\%$
 (now better known in both experiments, better than $\pm 2\%$)



W and Z production cross sections at hadron colliders

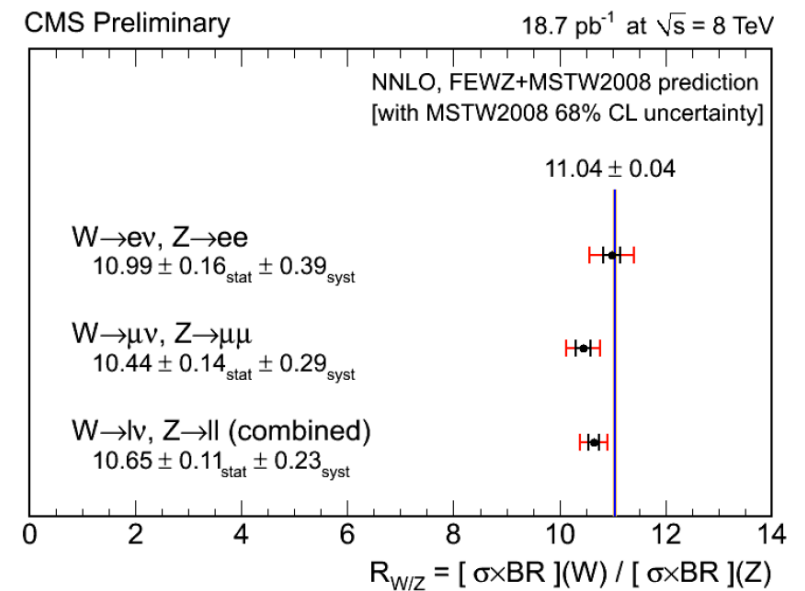
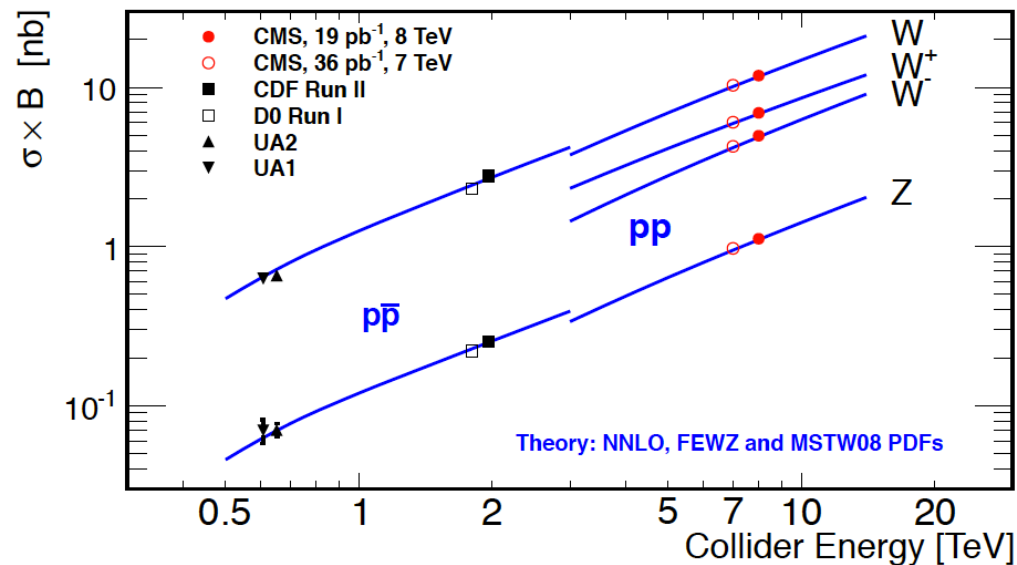


- Theoretical NNLO predictions in very good agreement with the experimental measurements (for pp, ppbar and as a function of energy)
- Good agreement as well between the ATLAS and CMS experiments



W and Z production cross sections at $\sqrt{s} = 8$ TeV

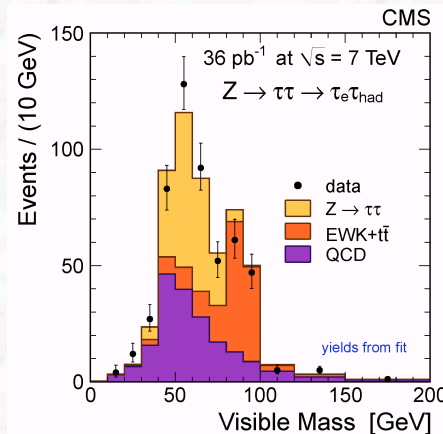
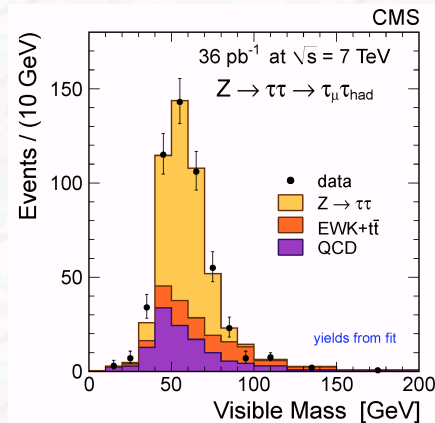
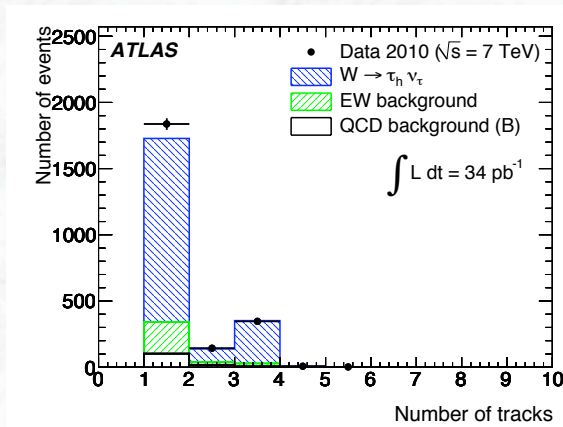
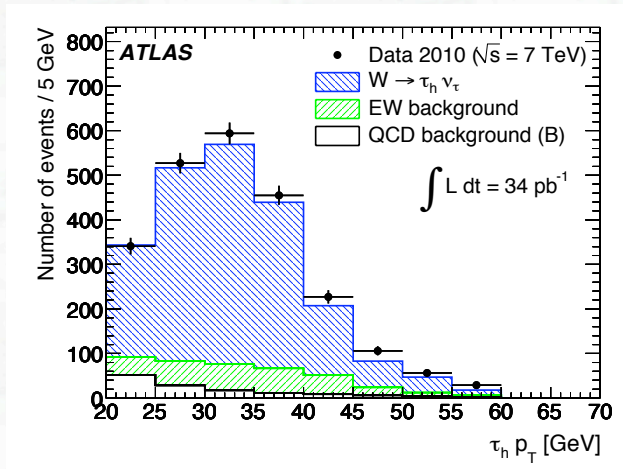
- CMS has already presented first results at 8 TeV (the first 18.7 pb^{-1})
About 75.000 $W \rightarrow e\nu$ and 4.800 $Z \rightarrow ee$ candidates



- No surprise at the new energy, theoretical predictions in good agreement with the measurements
- W/Z cross-section ratio remains a bit high, but consistent within uncertainties

First physics signals with hadronic tau final states

- Taus are more difficult to detect
- They decay with a short lifetime (0.3 ps) into 1 or 3 charged hadrons (65%) and a neutrino
- Taus have to be separated from hadronic jets



- First tau signals established in both ATLAS and CMS
- Important reference signals for searches with taus in Higgs and SUSY areas

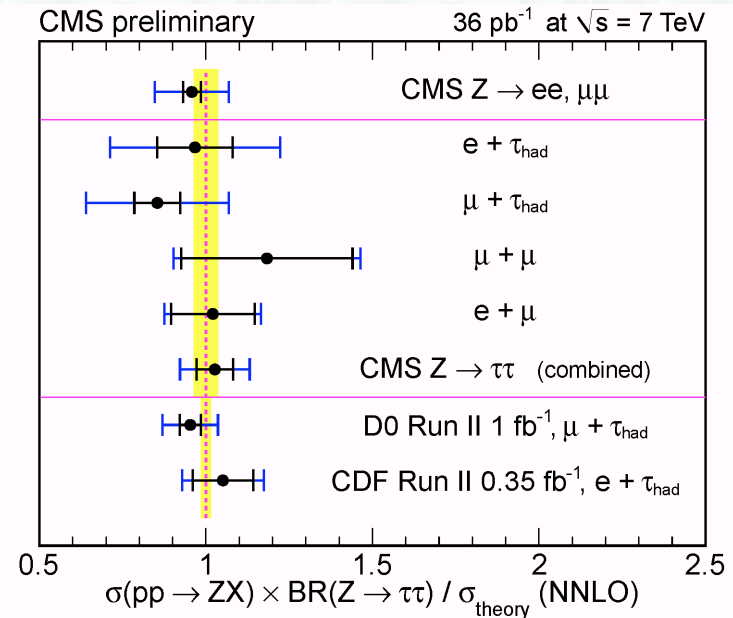
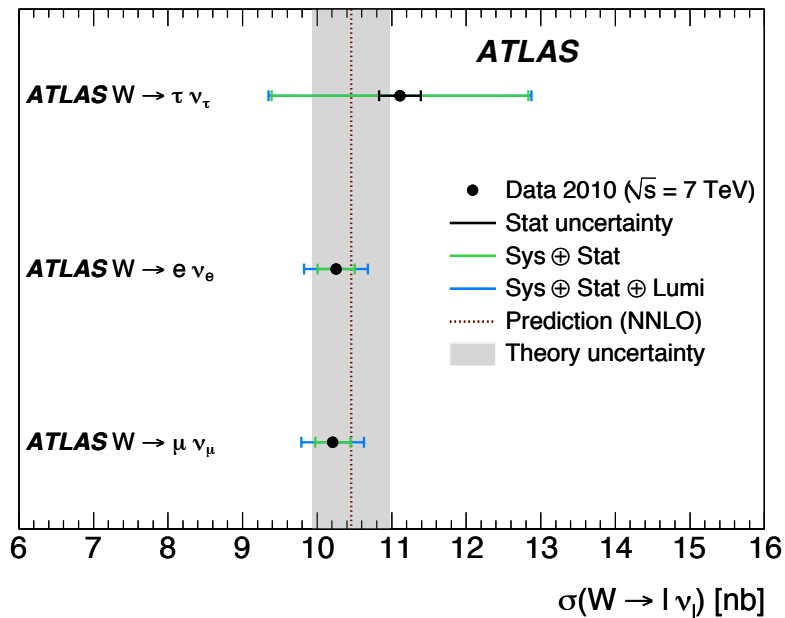


First physics signals with hadronic tau final states



$$W \rightarrow \tau \nu$$

$$Z \rightarrow \tau \tau$$



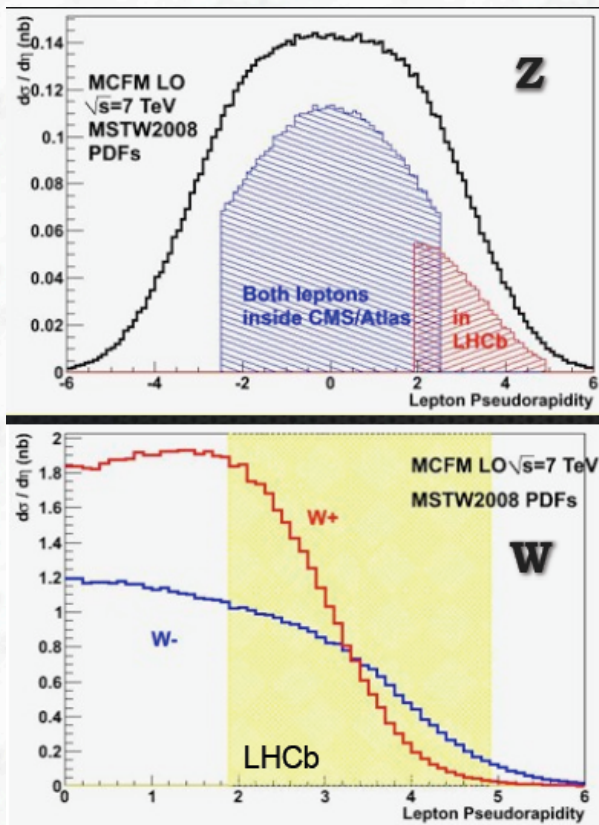
- Good agreement between the measured cross sections in the three lepton flavours
- Experimental uncertainties ($Z \rightarrow \tau\tau$) already comparable to Tevatron measurements

Can the parton distribution functions be constrained?

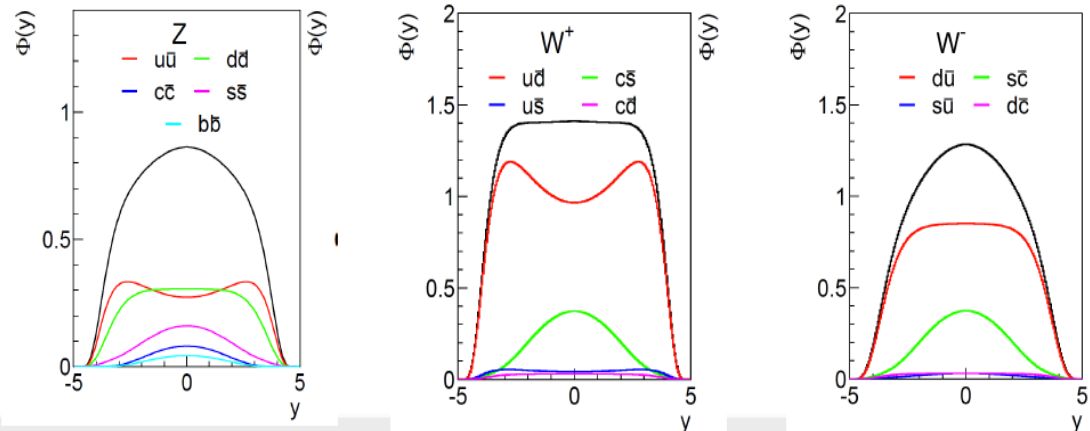
- Sensitive measurements: differential W and Z production cross sections as function of lepton or boson rapidity, charge separated for W^+ and W^-

LHCb experiment can contribute significantly in the forward region:
 η coverage from 1.9 – 4.9

- Derived quantity: charge asymmetry: $\sigma(W^+) - \sigma(W^-) / [\sigma(W^+) + \sigma(W^-)]$

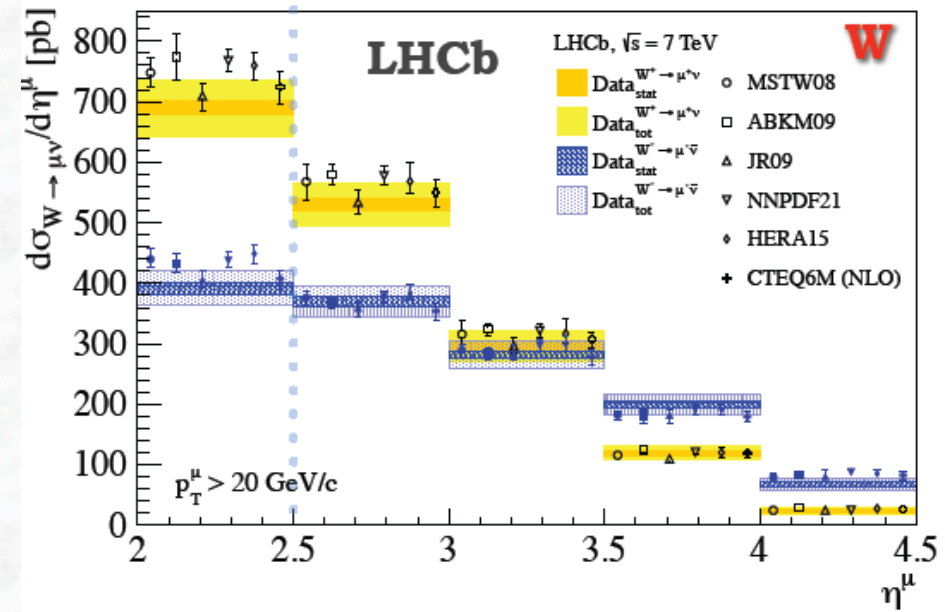
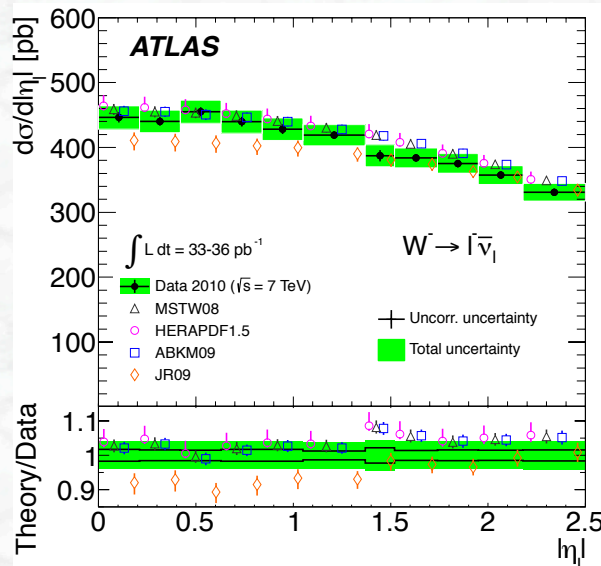
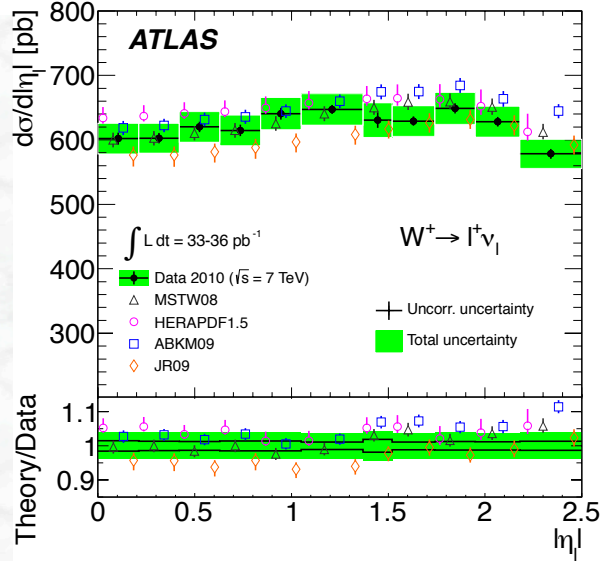


Leading order (tree level) contributions to W/Z production





Differential cross section measurements

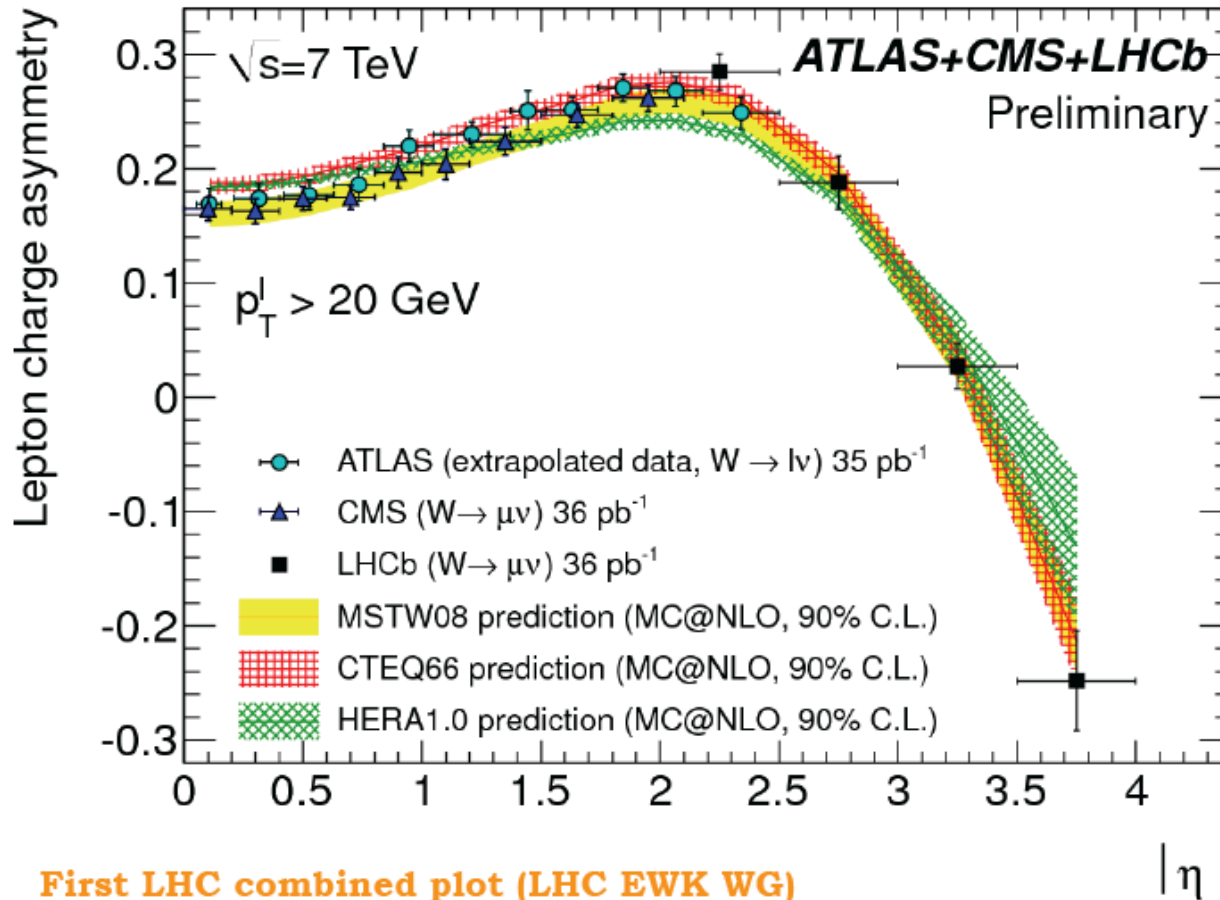


- Rough features of the measured differential cross sections are well described; (some tension at intermediate η region)
- Data start to be discriminating between pdf models;

These data will have impact on pdf uncertainties



W charge asymmetries



$$A(\eta_e) = \frac{d\sigma_{W^+}(\eta_e) - d\sigma_{W^-}(\eta_e)}{d\sigma_{W^+}(\eta_e) + d\sigma_{W^-}(\eta_e)}$$

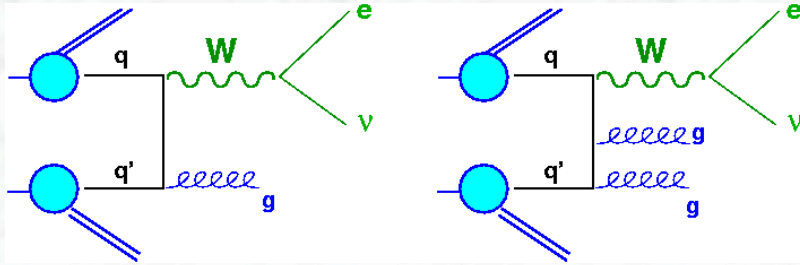
All data are unfolded,
 $P_T(l) > 20 \text{ GeV}$

First LHC combined plot (LHC EWK WG)

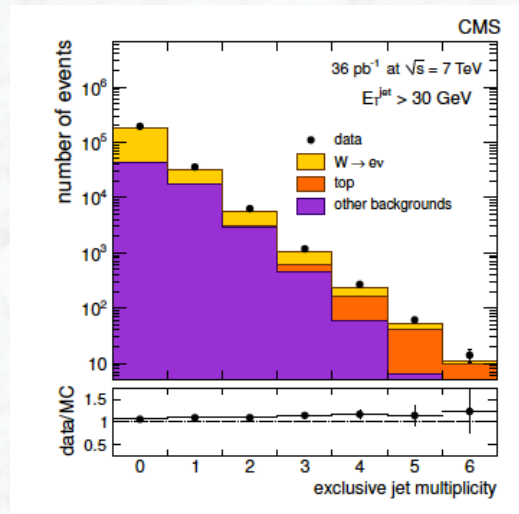


- Combination of the LHC experiments leads to large η coverage interesting constraints already today

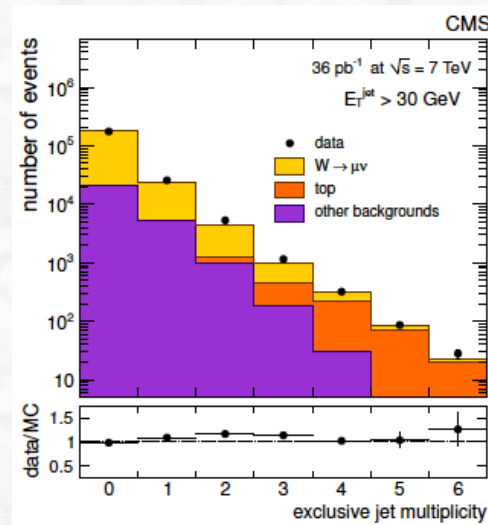
QCD Test in W/Z + jet production



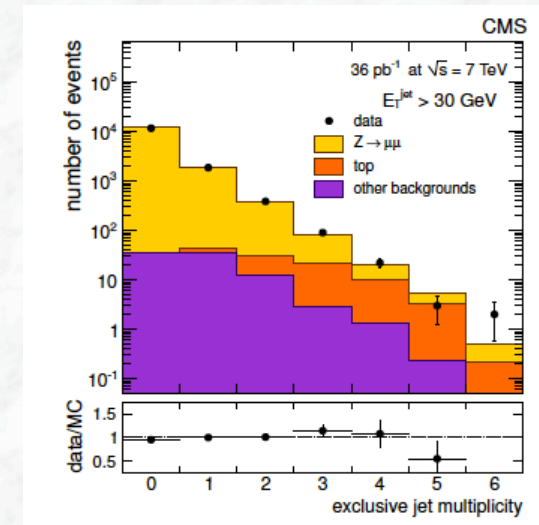
- CMS inclusive spectra of jets associated to W/Z production (36 pb^{-1});
- At detector level, compared to Monte Carlo Simulation (Madgraph + PYTHIA) (normalized to (N)NLO calculations)



$W \rightarrow e\nu$



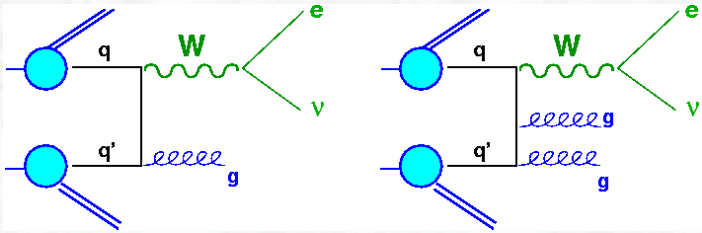
$W \rightarrow \mu\nu$



$Z \rightarrow \mu\mu$

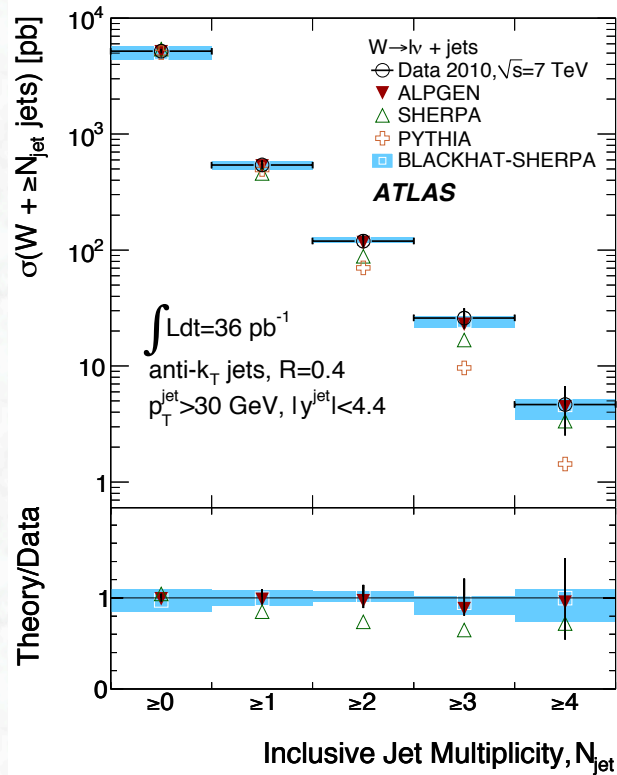
- Good agreement at that stage (jets with $p_T > 30 \text{ GeV}$),
- Top contribution clearly visible in high multiplicity bins of W + jet production

W/Z + jet cross section measurements

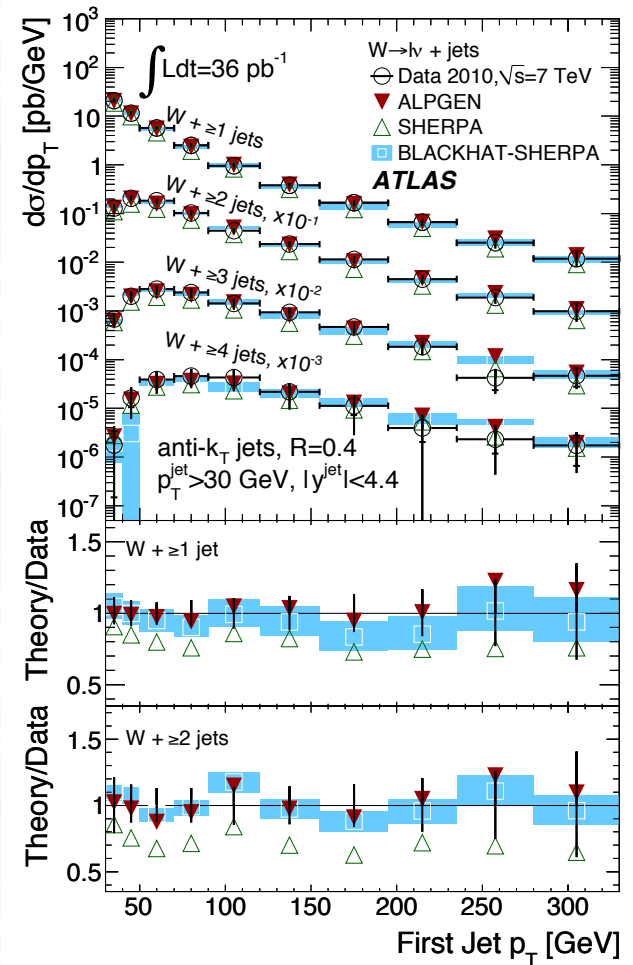


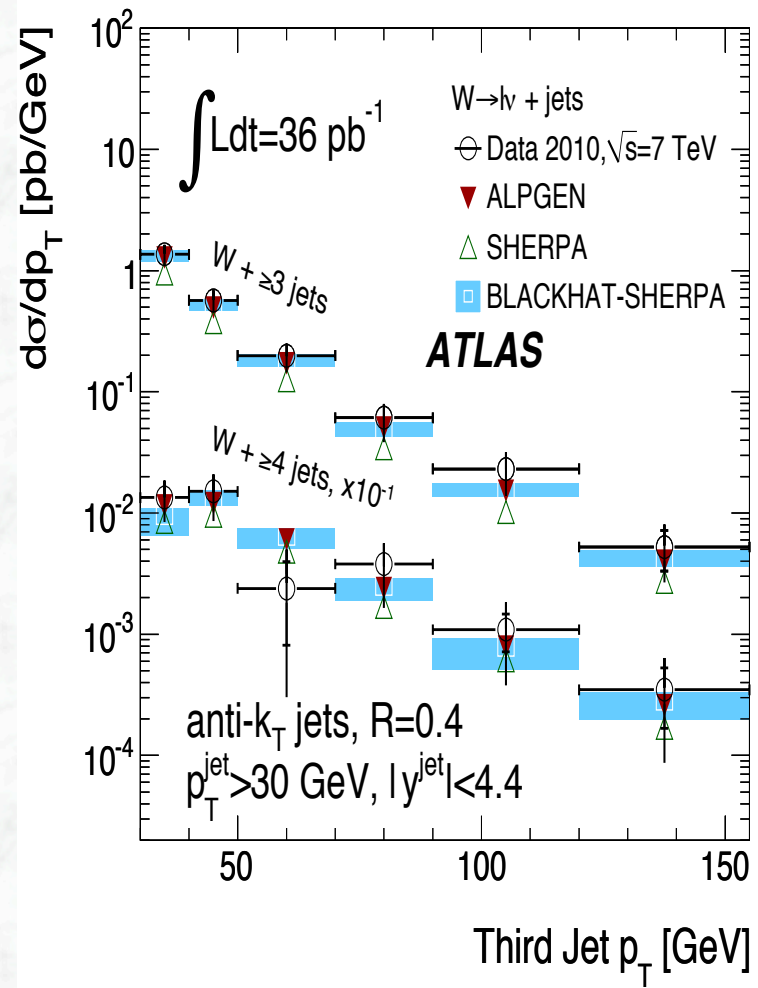
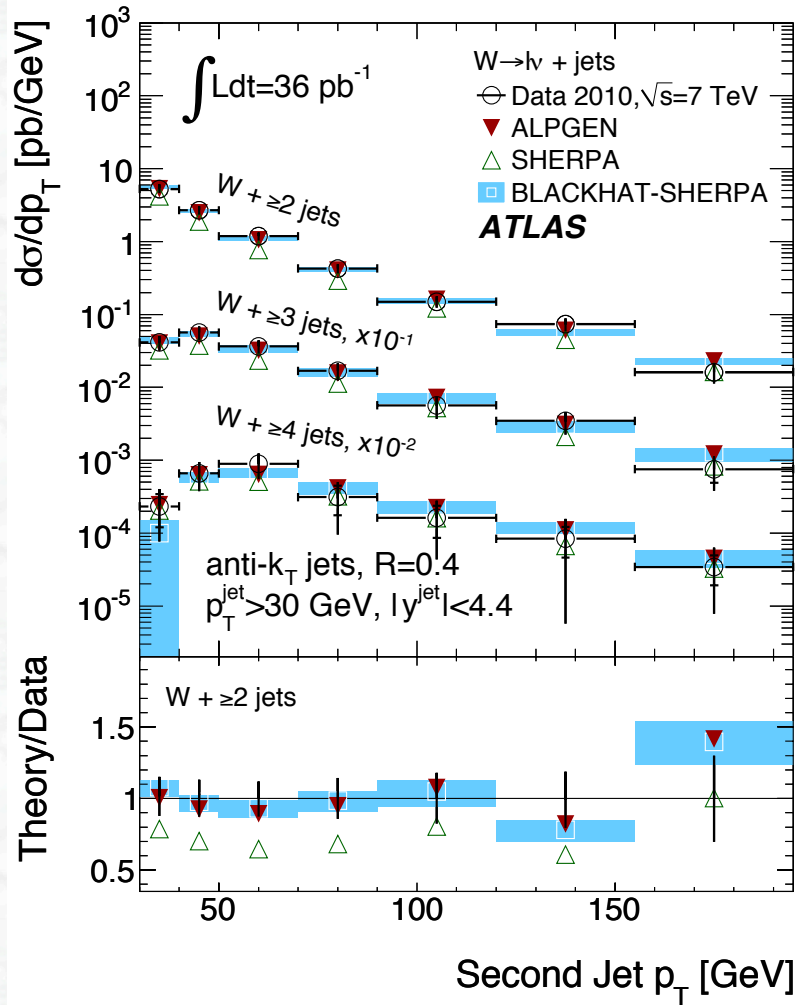
- LO predictions fail to describe the data;
- Jet multiplicities and p_T spectra in agreement with NLO predictions within errors;

Jet multiplicities in W+jet production



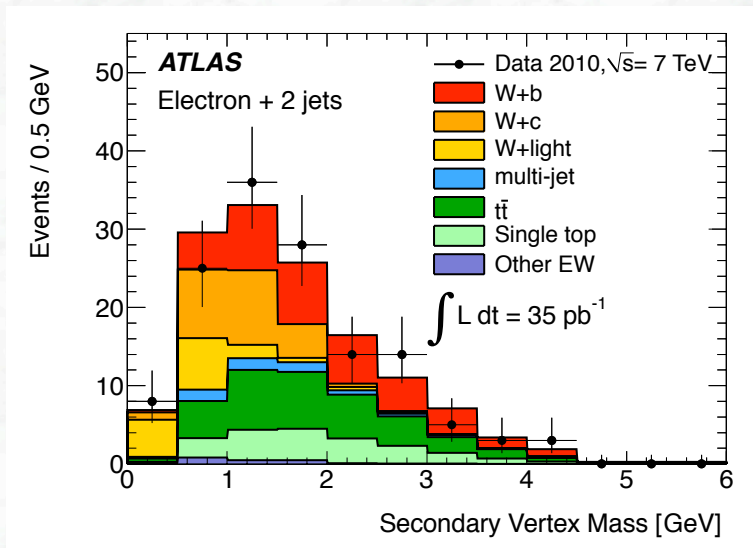
p_T spectrum of leading jet





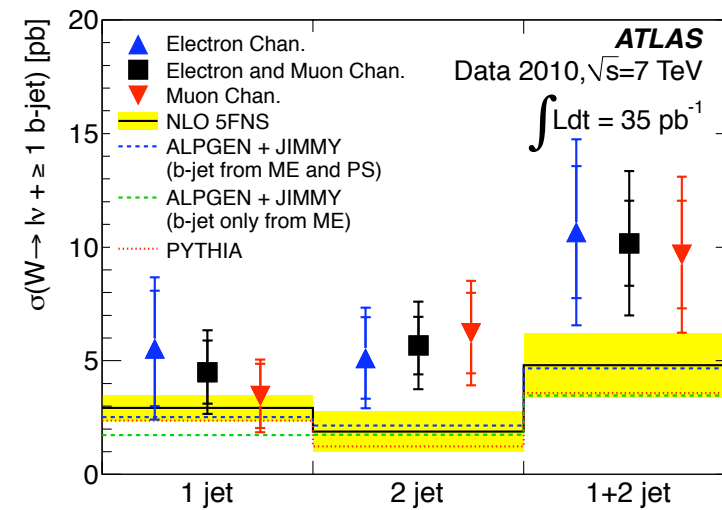
W + b jets

- Important background for many studies (Higgs, SUSY, top)
- Measurements at the Tevatron exceed NLO prediction
- Measured by ATLAS using 2010 data sample
 - studied W + 1 jet and W + 2 jets
 - require at least one b-tagged jet



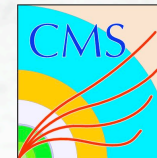
$W \rightarrow e \nu + 2 \text{ jets}$

Distribution of the mass of the particles associated to the secondary vertex for b-tagged jets

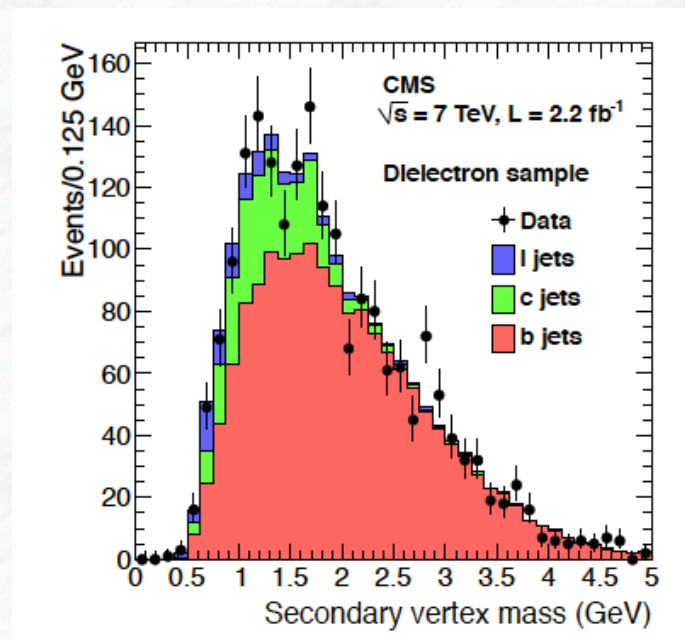
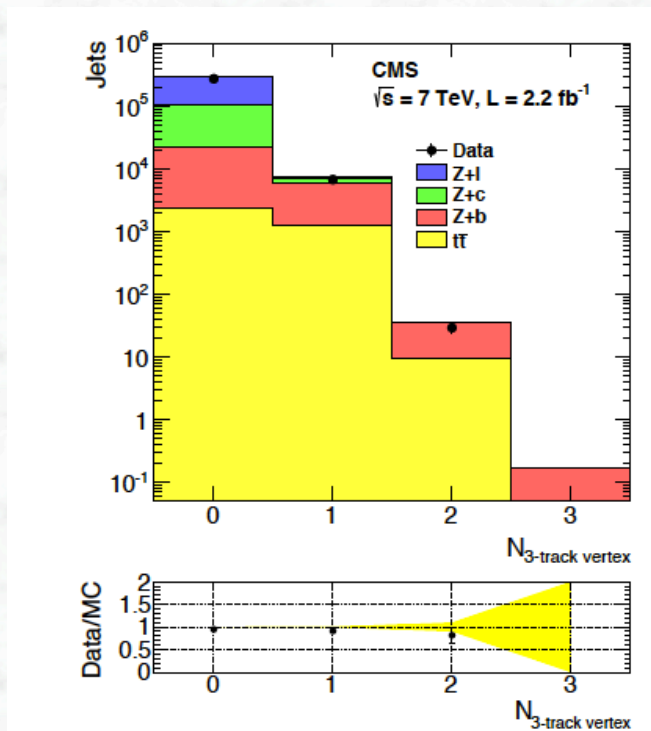
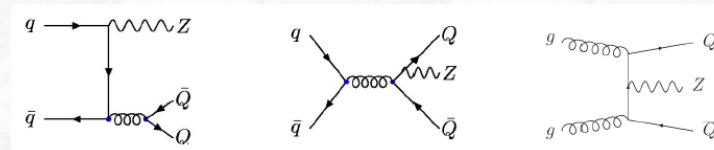


Results from e and μ combined.
Measurements $\sim 1.5\sigma$ above NLO prediction, but still consistent within uncertainties

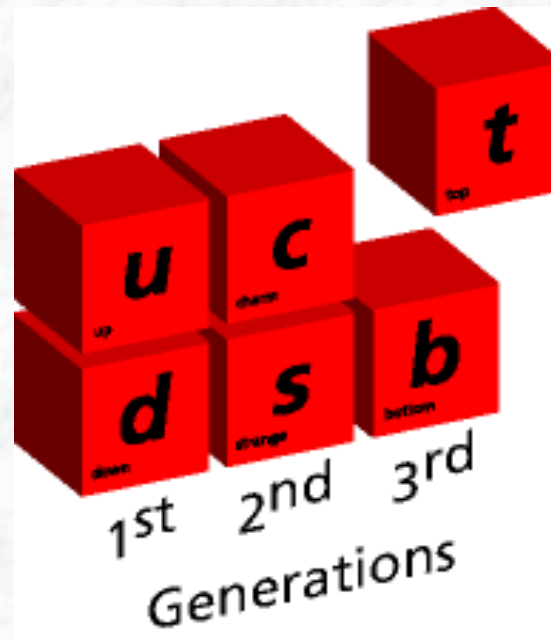
Z + b jets



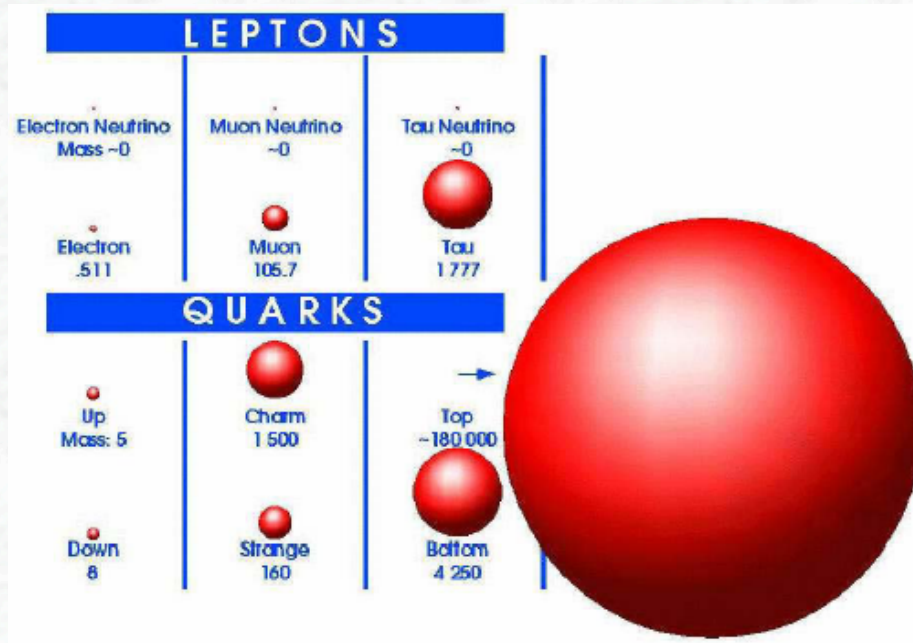
- Important background for many studies (Higgs, SUSY, top)
- Measured by CMS using 2011 data sample



Top Quark Physics



Why is Top-Quark so important ?



The top quark may serve as a window to **New Physics** related to the electroweak symmetry breaking;

Why is its Yukawa coupling ~ 1 ??

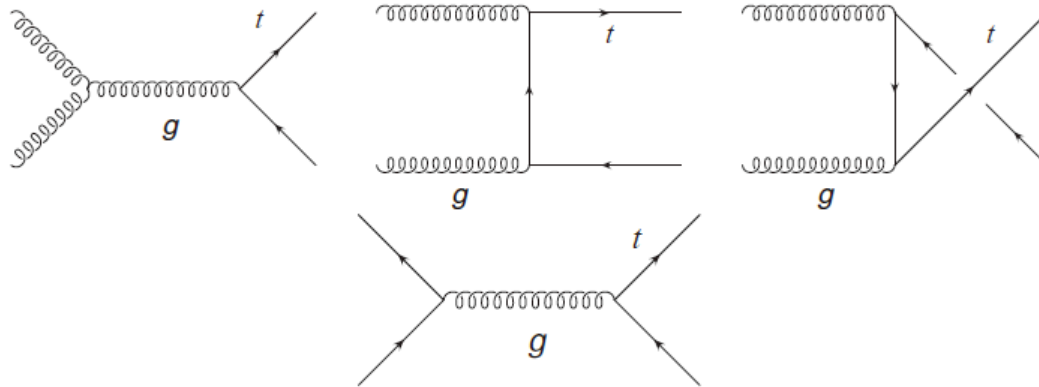
$$M_t = \frac{1}{\sqrt{2}} \lambda_t v$$

$$\Rightarrow \lambda_t = \frac{M_t}{173.9 \text{ GeV} / c^2}$$

- A unique quark: decays before it hadronizes, lifetime $\sim 10^{-25}$ s
no “toponium states”
remember: bb, bd, bs..... cc, cs..... bound states (mesons)
- **We still know little about the properties of the top quark:**
mass, spin, charge, lifetime, decay properties (rare decays), gauge couplings, Yukawa coupling,...

Top Quark Production

Pair production: qq and gg-fusion



Top-quark pair production in the Born approximation.

- NLO corrections completely known
- NNLO partly known

approximate NNLO results:

$$\sigma_{\text{LHC}} = (887_{-33}^{+9} (\text{scale})_{-15}^{+15} (\text{PDF})) \text{ pb} \quad (14 \text{ TeV}),$$

$$\sigma_{\text{Tev}} = (7.04_{-0.36}^{+0.24} (\text{scale})_{-0.14}^{+0.14} (\text{PDF})) \text{ pb} \quad (1.96 \text{ TeV}).$$

	Tevatron 1.96 TeV	LHC 14 TeV
qq	85%	5%
gg	15%	95%
σ (pb)	7.0 pb	887 pb

For LHC running at $\sqrt{s} = 7 \text{ TeV}$, the cross section is reduced by a factor of ~ 5 , but it is still a factor 25 larger than the cross section at the Tevatron

Top Quark Decays

BR ($t \rightarrow Wb$) $\sim 100\%$

Dilepton channel:

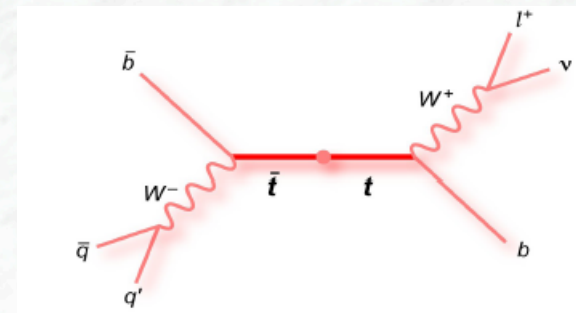
Both W 's decay via $W \rightarrow \ell\nu$ ($\ell=e$ or μ ; 4%)

Lepton + jet channel:

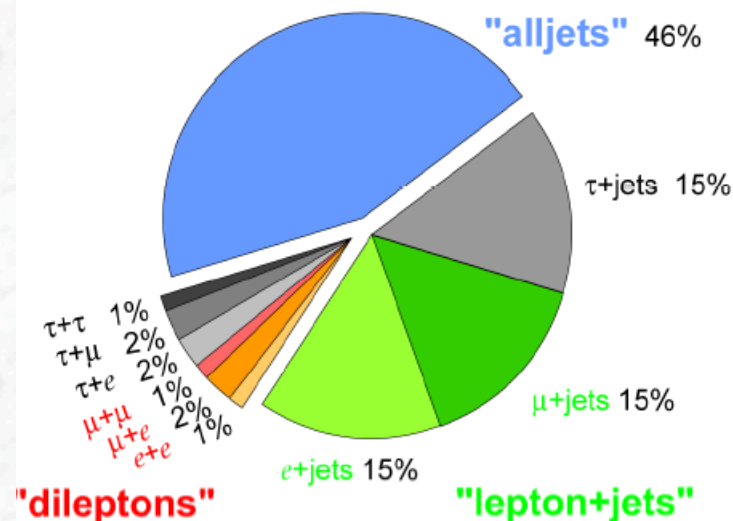
One W decays via $W \rightarrow \ell\nu$ ($\ell=e$ or μ ; 30%)

Full hadronic channel:

Both W 's decay via $W \rightarrow qq$ (46%)



Top Pair Branching Fractions



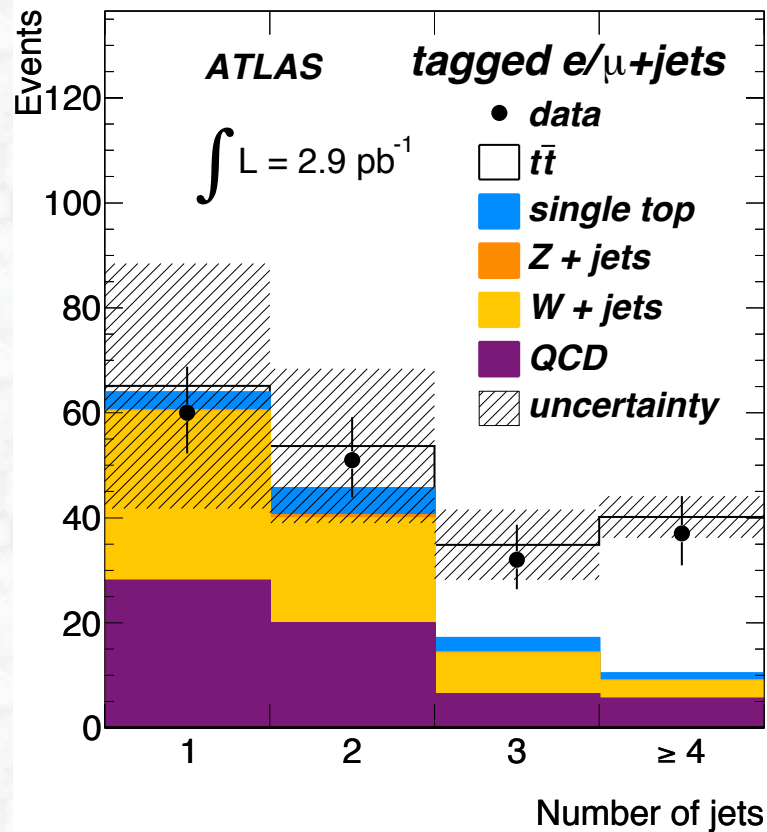
Important experimental signatures: - Lepton(s)

- Missing transverse momentum

- b-jet(s)



First results on top production from the LHC

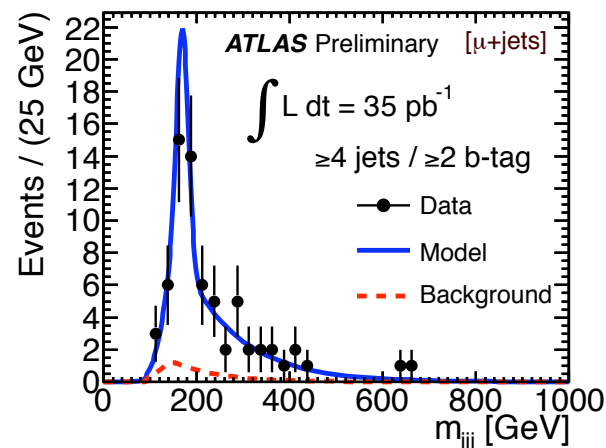
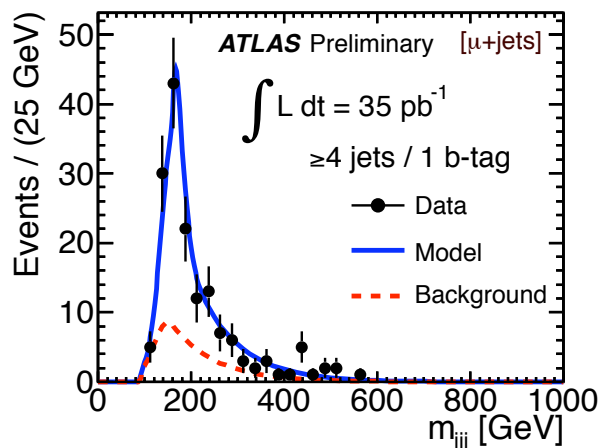
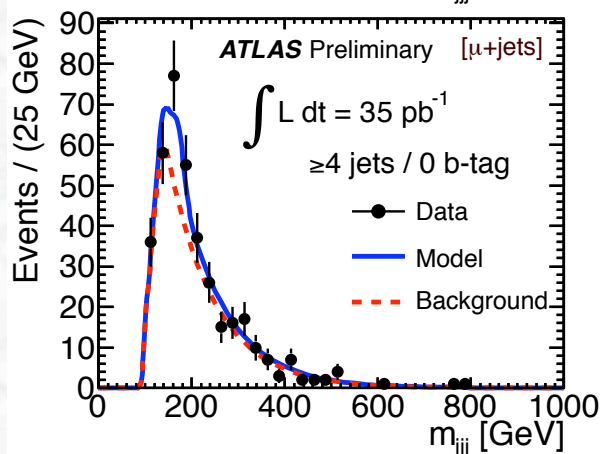
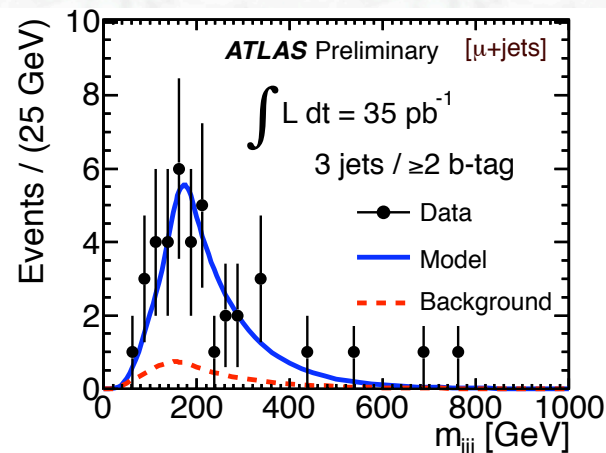
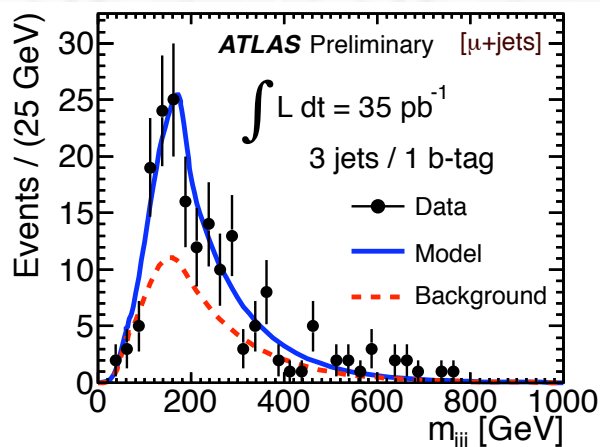
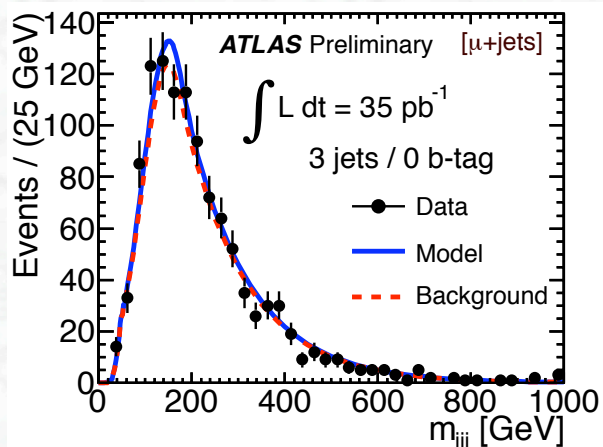
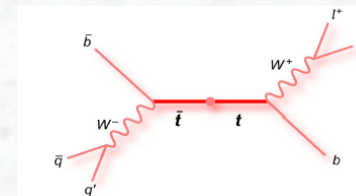


Event Selection:

- Lepton trigger
- One identified lepton (e, μ) with $p_T > 20 \text{ GeV}$
- Missing transverse energy: $E_T^{\text{miss}} > 35 \text{ GeV}$ (significant rejection against QCD events)
- Transverse mass: $M_T(l, \nu) > 25 \text{ GeV}$ (lepton from W decay in event)
- One or more jets with $p_T > 25 \text{ GeV}$ and $\eta < 2.5$



Invariant mass distributions in the l-had channel



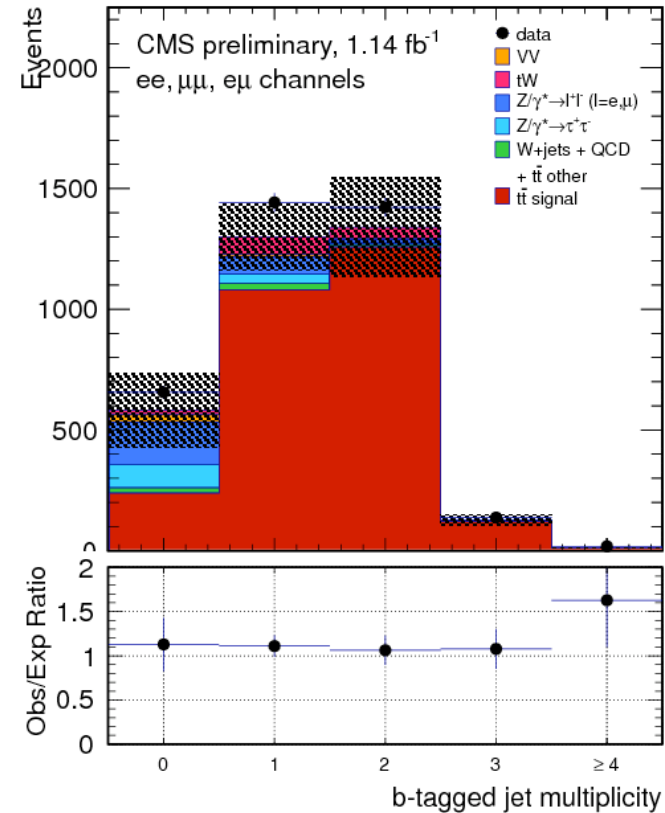
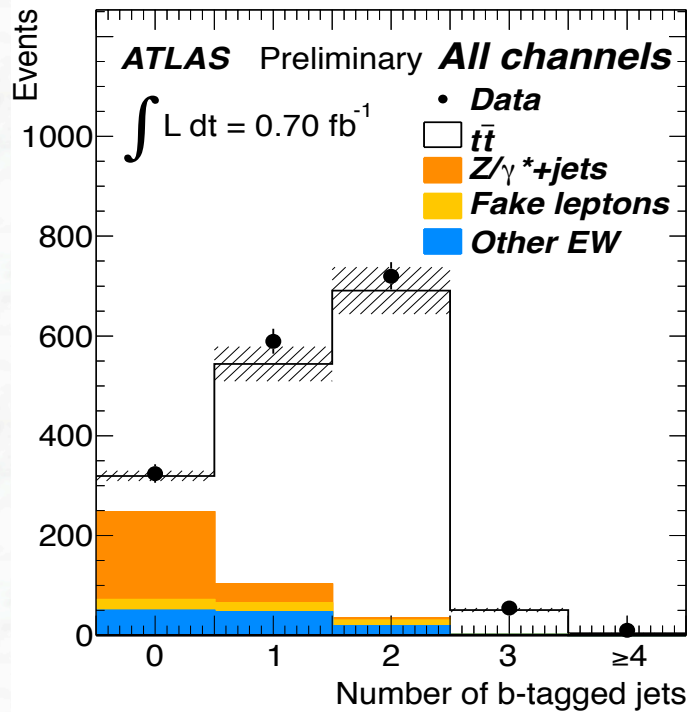
- Top fractions increase with number of b-tags
- Good description for all jet-multiplicity and b-tag combinations
- Data are consistent with top quark production with mass of 173 GeV

Top-quark production measured in many different decay modes

(i) Di-lepton selection in both ATLAS and CMS (0.7 fb⁻¹ – 1.14 fb⁻¹)



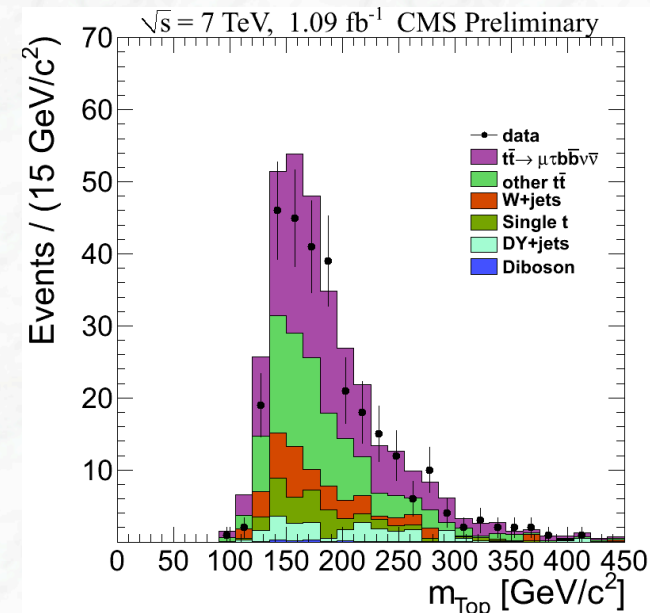
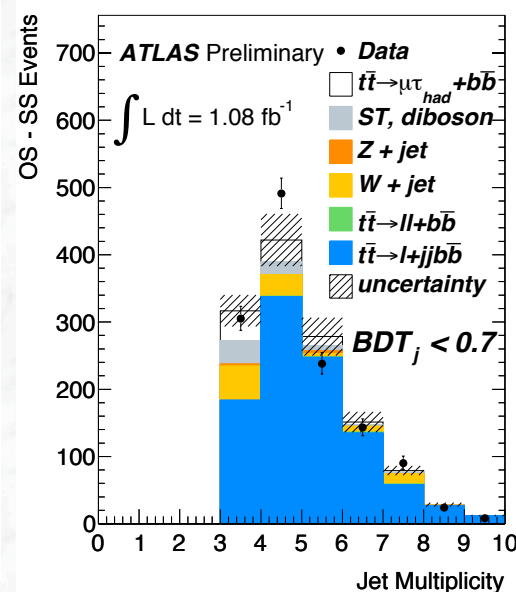
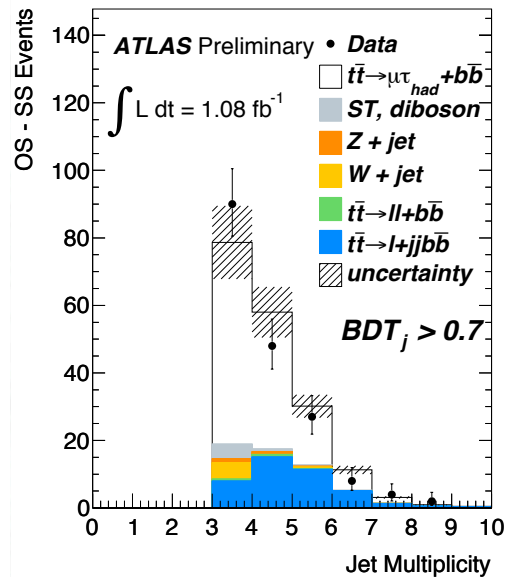
Multiplicity distributions of b-tagged jets
(small backgrounds, mainly from Z+jet production)





(ii) $\mu + \tau$ final states in both ATLAS and CMS ($0.7 \text{ fb}^{-1} - 1.14 \text{ fb}^{-1}$)

Require: $\mu +$ hadronically decaying τ , $E_T^{\text{miss}} + \text{b-jets}$
(significant backgrounds, but signal contribution needed)



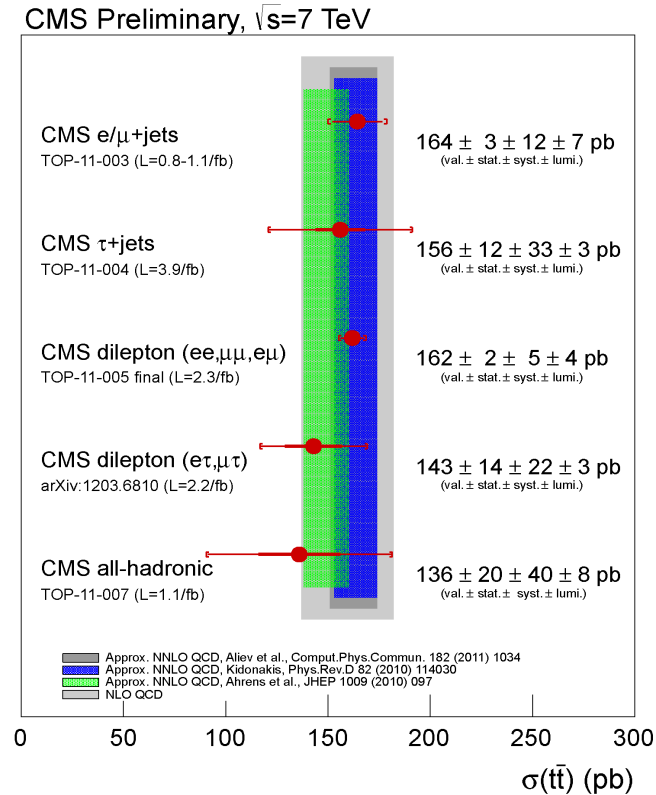
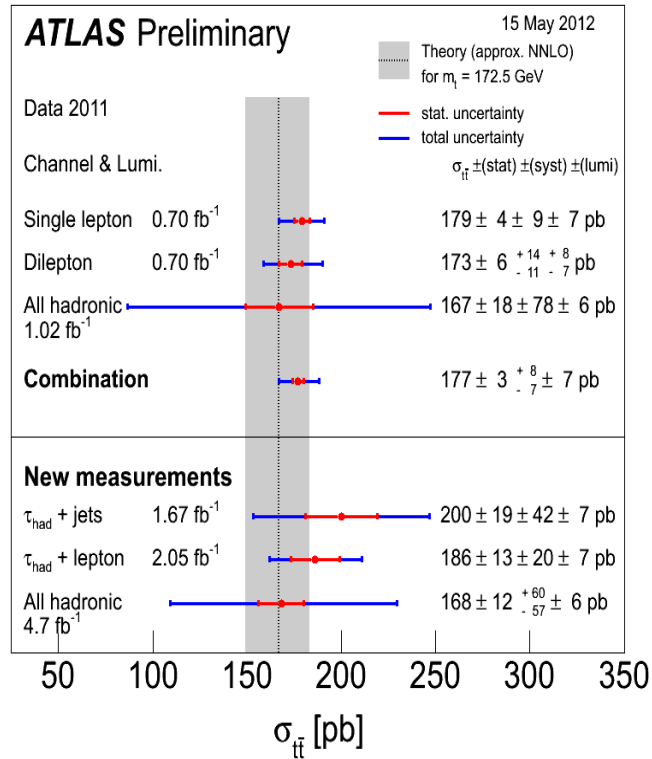
ATLAS: Multivariate analysis
Jet multiplicity distribution in signal (left)
and background (right) regions

reconstructed mass in CMS



Top pair production cross section measurements

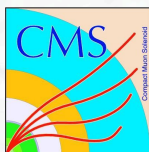
-likelihood combination of all channels-



$\sigma = 177 \pm 3$ (stat) ± 7 (syst) ± 7 (lum) pb

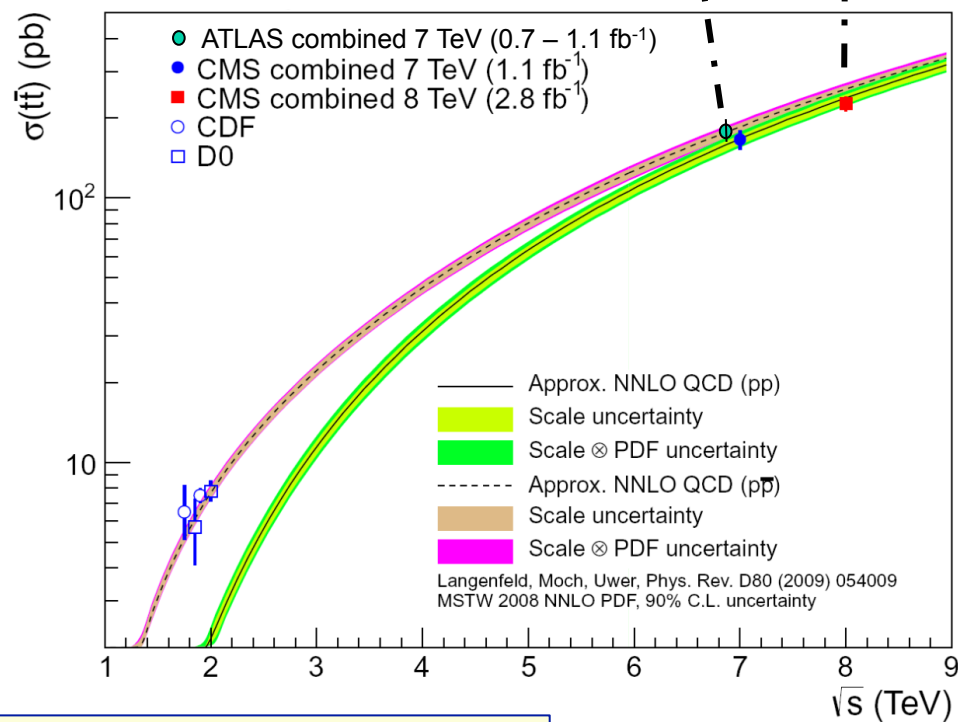
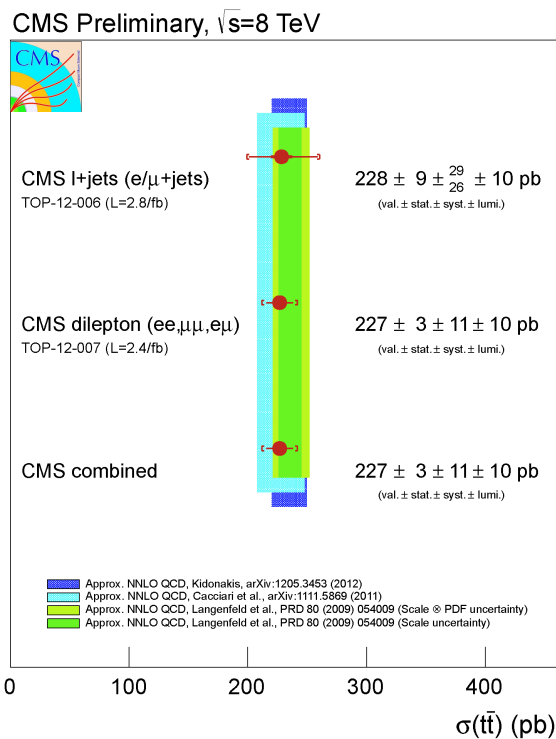
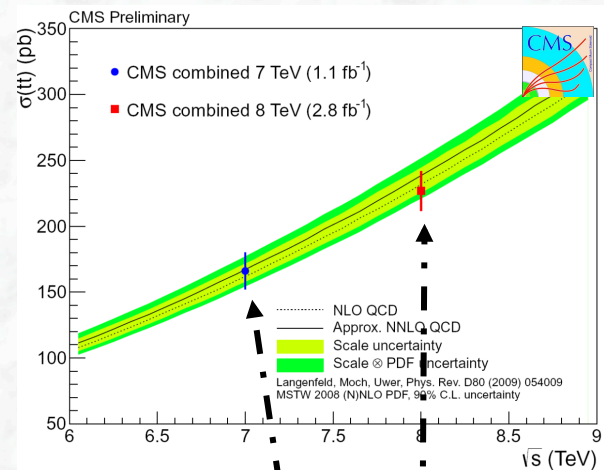
$\sigma = 165.8 \pm 2.2$ (stat) ± 10.6 (syst) ± 7.8 (lum) pb

- Perturbative QCD calculations (approx. NNLO) describe the data well;
- The two LHC experiments agree within the systematic uncertainties
- Total uncertainty already at the level of $\pm 6\%$



- **CMS: new measurement at 8 TeV !**
- Lepton + jets and di-lepton channels combined:

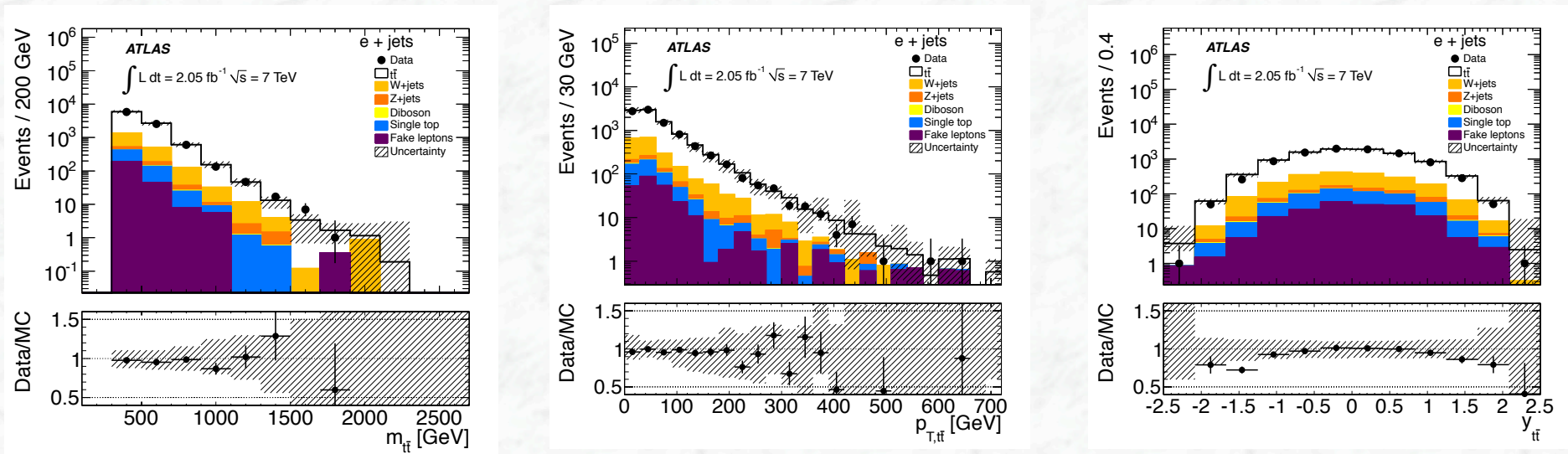
$$\sigma = 227 \pm 3 \text{ (stat)} \pm 11 \text{ (syst.)} \pm 10 \text{ (lum.) pb}$$



$$\sigma(8\text{TeV})/\sigma(7\text{TeV}) = 1.41 \pm 0.11; \text{ no correlation assumed}$$

Top-antitop differential cross sections

- Important test of the Standard Model (perturbative QCD), deviations may indicate new physics
e.g. new particles (resonances) decaying into $t\bar{t}$, or other new/unexpected effects (\rightarrow Tevatron charge asymmetry)
- Important variables studied:
 - $t\bar{t}$ mass distribution
 - Rapidity y and p_T of the $t\bar{t}$ system



ATLAS comparison on detector level shows good agreement in all variables (background partially extracted from data)



\rightarrow not much room left / no signs yet of Physics beyond the Standard Model (more in the lecture of M. Narain)



- Both collaborations have unfolded the detector effects and have extracted differential cross-section measurements (normalized to the $t\bar{t}$ cross section \rightarrow sensitivity in shapes of distributions)

