Hawking emission from quantum black holes

Elizabeth Winstanley

Consortium for Fundamental Physics Astro-Particle Theory and Cosmology Group School of Mathematics and Statistics University of Sheffield United Kingdom

Work done in collaboration with: Marc Casals, Sam Dolan, Gavin Duffy, Panagiota Kanti and Piero Nicolini

Thanks to STFC and EU-COST network MP0905 for financial support

Outline



- 2 Semi-classical evolution
- 3 Beyond the semi-classical approximation
- 4 Conclusions

Brane worlds

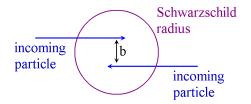
Our universe is a brane living in a higher-dimensional bulk

 ADD model - flat compactified extra dimensions Standard model physics restricted to the brane • Only gravity propagates in the bulk SM fields $(M_{\mathbf{p}})$ Fundamental higher-dimensional scale of quantum gravity, M_{*} , may be as low as the energy of Gravitons and Scalars the LHC: (M.) $M_P^2 \sim M_{\bullet}^{n+2} R^n$ v

[Figure taken from Kanti, arXiv:0802.2218 [hep-th]]

Formation of mini black holes

If $M_* \sim$ few TeV, particle collisions at LHC may produce heavy, quantum gravitational objects



Will a semi-classical black hole be formed?

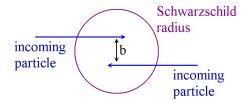
Compton wavelength of colliding particle of energy E/2 must lie within the Schwarzschild radius:

$$4\pi/E < r_h(E)$$

Therefore $E/M_* \gtrsim 10$ in order for black holes to form [Meade and Randall, arXiv:0708.3017 [hep-ph]]

Formation of mini black holes

If $M_* \sim$ few TeV, particle collisions at LHC may produce heavy, quantum gravitational objects



Will a semi-classical black hole be formed?

Compton wavelength of colliding particle of energy E/2 must lie within the Schwarzschild radius:

$$4\pi/E < r_h(E)$$

Therefore $E/M_* \gtrsim 10$ in order for black holes to form [Meade and Randall, arXiv:0708.3017 [hep-ph]]

Stages in the evolution of small black holes

Black holes formed will be rapidly rotating, highly asymmetric, and have gauge field hair

Four stages of subsequent evolution:

"Balding" stage	"Spin-down" stage	"Schwarzschild" stage	"Quantum gravity" stage	

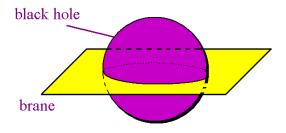
[Giddings and Thomas, hep-ph/0106219]

Elizabeth Winstanley (Sheffield)

Modelling small black holes at the end of the balding stage

Small black holes in ADD

- Metric of higher-dimensional black holes in general relativity is known [Myers and Perry, *Annals Phys.* **172**, 304 (1986)]
- Take a 'slice' through a higher-dimensional black hole to give a brane black hole



Modelling small black holes in ADD

Myers-Perry higher-dimensional black hole

$$ds^{2} = \left(1 - \frac{\mu}{\Sigma r^{n-1}}\right) dt^{2} + \frac{2a\mu\sin^{2}\theta}{\Sigma r^{n-1}} dt \, d\varphi - \frac{\Sigma}{\Delta_{n}} dr^{2} - \Sigma \, d\theta^{2}$$
$$- \left(r^{2} + a^{2} + \frac{a^{2}\mu\sin^{2}\theta}{\Sigma r^{n-1}}\right) \sin^{2}\theta \, d\varphi^{2} - r^{2}\cos^{2}\theta \, d\Omega_{n}^{2}$$

where

$$\Delta_n = r^2 + a^2 - \frac{\mu}{r^{n-1}}, \qquad \Sigma = r^2 + a^2 \cos^2 \theta$$

Black hole mass M and angular momentum J:

$$M = \frac{(n+2)A_{n+2}\mu}{16\pi G_{4+n}}, \qquad J = \frac{2aM}{n+2}$$

Modelling small black holes in ADD

Slice of Myers-Perry black hole

$$ds^{2} = \left(1 - \frac{\mu}{\Sigma r^{n-1}}\right) dt^{2} + \frac{2a\mu\sin^{2}\theta}{\Sigma r^{n-1}} dt \, d\varphi - \frac{\Sigma}{\Delta_{n}} dr^{2} - \Sigma \, d\theta^{2}$$
$$- \left(r^{2} + a^{2} + \frac{a^{2}\mu\sin^{2}\theta}{\Sigma r^{n-1}}\right) \sin^{2}\theta \, d\varphi^{2}$$

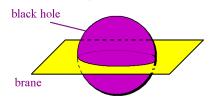
where

$$\Delta_n = r^2 + a^2 - \frac{\mu}{r^{n-1}}, \qquad \Sigma = r^2 + a^2 \cos^2 \theta$$

and n is the number of extra dimensions.

Usual Kerr black hole Set n = 0 in the above metric Elizabeth Winstanley (Sheffield) Modelling the evolution of small black holes Bonn, June 2010 8/21

Hawking radiation on the brane and in the bulk



Particles on the brane

- Standard model particles: fermions, gauge bosons, Higgs
- Also gravitons and scalars
- Live on the brane "slice" of the black hole geometry

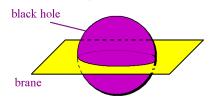
Hawking temperature

$$T_{H} = \frac{(n+1)r_{h}^{2} + (n-1)a^{2}}{4\pi(r_{h}^{2} + a^{2})r_{h}}$$

Particles in the bulk

- Gravitons and scalars
- Will be invisible
- Live on the higher-dimensional black hole geometry

Hawking radiation on the brane and in the bulk



Particles on the brane

- Standard model particles: fermions, gauge bosons, Higgs
- Also gravitons and scalars
- Live on the brane "slice" of the black hole geometry

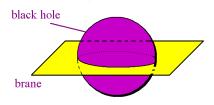
Hawking temperature

$$T_{H} = \frac{(n+1)r_{h}^{2} + (n-1)a^{2}}{4\pi(r_{h}^{2} + a^{2})r_{h}}$$

Particles in the bulk

- Gravitons and scalars
- Will be invisible
- Live on the higher-dimensional black hole geometry

Hawking radiation on the brane and in the bulk



Particles on the brane

- Standard model particles: fermions, gauge bosons, Higgs
- Also gravitons and scalars
- Live on the brane "slice" of the black hole geometry

Hawking temperature

$$T_{H} = \frac{(n+1)r_{h}^{2} + (n-1)a^{2}}{4\pi(r_{h}^{2} + a^{2})r_{h}}$$

Particles in the bulk

- Gravitons and scalars
- Will be invisible
- Live on the higher-dimensional black hole geometry

Quantum fields on black hole space-times

Quantum field theory in curved space-time

- Black hole geometry is fixed and classical
- Quantum fields (scalars, fermions, gauge bosons, gravitons) propagate on this background

Quantum field modes

- "Master" equation for fields of spin 0, ¹/₂, 1 and 2 on Kerr
 [Teukolsky, *Phys. Rev. Lett.* 29 1114 (1972); *Astrophys. J.* 185 635 (1973)]
- Expand field Ψ in terms of modes of frequency ω :

$$\Psi = \sum_{\omega \ell m} R_{s \omega \ell m}(r) S_{s \omega \ell m}(\theta) e^{-i \omega t} e^{i m \varphi}$$

- 4 同 6 4 日 6 4 日 6

Quantum fields on black hole space-times

Quantum field theory in curved space-time

- Black hole geometry is fixed and classical
- Quantum fields (scalars, fermions, gauge bosons, gravitons) propagate on this background

Quantum field modes

- "Master" equation for fields of spin 0, ¹/₂, 1 and 2 on Kerr
 [Teukolsky, *Phys. Rev. Lett.* 29 1114 (1972); *Astrophys. J.* 185 635 (1973)]
- Expand field Ψ in terms of modes of frequency ω :

$$\Psi = \sum_{\omega \ell m} R_{s \omega \ell m}(r) S_{s \omega \ell m}(\theta) e^{-i \omega t} e^{i m \varphi}$$

Computing Hawking radiation

Differential emission rates, integrated over all angles:

$$\frac{d^2}{dt \, d\omega} \begin{pmatrix} N \\ E \\ J \end{pmatrix} = \frac{1}{4\pi} \sum_{\text{modes}} \frac{\left|\mathcal{A}_{s\omega\ell m}\right|^2}{e^{\tilde{\omega}/T_H} \mp 1} \begin{pmatrix} 1 \\ \omega \\ m \end{pmatrix}$$

where $\tilde{\omega} = \omega - m\Omega_H$

Elizabeth Wi

Grey-body factor $\left|\mathcal{A}_{s\omega\ell m}\right|^2$

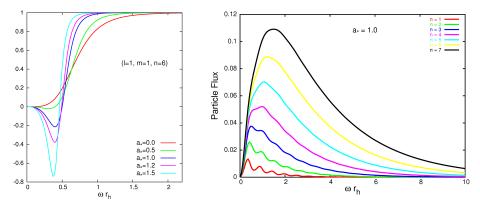
- Emitted radiation is not precisely thermal
- Interaction of emitted quanta with gravitational potential around the black hole
- For an outgoing wave from the event horizon of the black hole:

$$|\mathcal{A}_{s\omega\ell m}|^2 = 1 - |\mathcal{R}_{s\omega\ell m}|^2 = \frac{\mathcal{F}_{\text{infinity}}}{\mathcal{F}_{\text{horizon}}}$$

Grey-body factors and emission spectra

Grey-body factors for gauge boson emission and n = 6

Fermion emission spectra for a rotating black hole, integrated over all angles

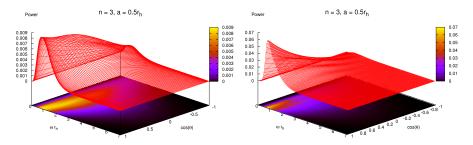


[Casals et al, hep-th/0511163] [Casals et al, hep-th/0608193]

Angular distribution of energy flux Differential energy emission rate:

$$\frac{d^{3}E}{dt \, d\omega \, d(\cos \theta)} = \frac{1}{4\pi} \sum_{\text{modes}} \frac{\omega \left| \mathcal{A}_{s\omega\ell m} \right|^{2}}{e^{\tilde{\omega}/T_{H}} \mp 1} \left[S_{|s|\omega\ell m}(\theta)^{2} + S_{-|s|\omega\ell m}(\theta)^{2} \right]$$

Energy emission for positive helicity fermions and gauge bosons for n = 3and $a_* = 0.5$



[Casals et al arXiv:0907.1511 [hep-th]

Elizabeth Winstanley (Sheffield)

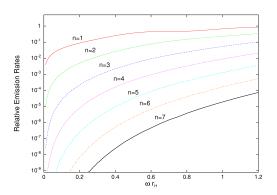
Modelling the evolution of small black hole

"Black holes radiate mainly on the brane"

[Emparan, Horowitz and Myers, hep-th/0003118]

Bulk/brane energy emission ratios for scalar fields from a non-rotating black hole

- As a function of frequency
- Total bulk energy emission/total brane energy emission



						<i>n</i> = 7
0.40	0.24	0.22	0.24	0.33	0.52	0.93

[Harris and Kanti, hep-ph/0309054]

Beyond the semi-classical approximation

The semi-classical approximation

- Treats the black hole geometry as fixed and classical
- Is tractable and yields results which can be used in simulations
- But, by definition, quantum gravity effects are important!

Non-commutative-geometry-inspired black holes - NCBHs

- Use smeared mass density inspired by non-commutative geometry as source for classical metric
- Metric:

$$ds^{2} = f(r) dt^{2} - f(r)^{-1} dr^{2} - r^{2} d\Omega_{n+2}^{2}$$
$$f(r) = 1 - \left(\frac{r_{h}}{r}\right)^{n+1} \rightarrow 1 - \frac{\mu}{r^{n+1}} \gamma \left(\frac{n+3}{2}, \frac{r^{2}}{4\vartheta}\right)$$

Nicolini et al, gr-qc/0510112]

< 177 ▶

Beyond the semi-classical approximation

The semi-classical approximation

- Treats the black hole geometry as fixed and classical
- Is tractable and yields results which can be used in simulations
- But, by definition, quantum gravity effects are important!

Non-commutative-geometry-inspired black holes - NCBHs

- Use smeared mass density inspired by non-commutative geometry as source for classical metric
- Metric:

$$ds^{2} = f(r) dt^{2} - f(r)^{-1} dr^{2} - r^{2} d\Omega_{n+2}^{2}$$
$$f(r) = 1 - \left(\frac{r_{h}}{r}\right)^{n+1} \rightarrow 1 - \frac{\mu}{r^{n+1}} \gamma \left(\frac{n+3}{2}, \frac{r^{2}}{4\vartheta}\right)$$

[Nicolini et al, gr-qc/0510112]

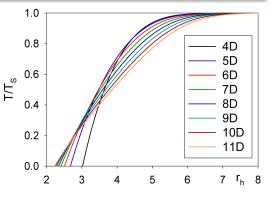
A 🖓

Hawking radiation from non-rotating, neutral NCBHs

Fluxes of scalar particles and energy

$$\frac{d^2}{dt \, d\omega} \left(\begin{array}{c} N\\ E \end{array}\right) = \frac{1}{4\pi} \sum_{\text{modes}} \frac{\left|\mathcal{A}_{\omega\ell m}\right|^2 e^{-\omega^2/2}}{e^{\omega/T_H} - 1} \left(\begin{array}{c} 1\\ \omega \end{array}\right)$$

- Temperature of NCBHs much lower than standard BHs
- Quantum back-reaction negligible
- Decay time $\sim 10^{-16}s$ [Casadio and Nicolini, arXiv:0809.2471]

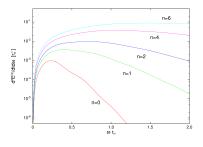


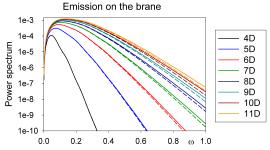
16 / 21

Brane scalar field energy emission









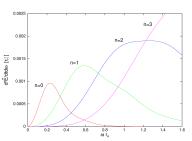
Brane emission from a Schwarzschild-Tangherlini BH, for n = 0...6

[Harris and Kanti, hep-ph/0309054] Brane emission from a higher-dimensional NCBH, for n = 0...7

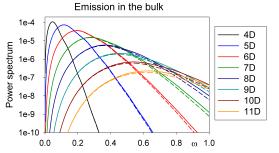
[Nicolini and EW, arXiv:1108.4419 [hep-th]]

THE 1 1

Bulk scalar field energy emission



NCBH



Bulk emission from a Schwarzschild-Tangherlini BH, for n = 0...3

[Harris and Kanti, hep-ph/0309054]

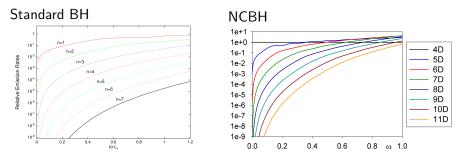
Standard BH

Bulk emission from a higher-dimensional NCBH, for n = 0...7

[Nicolini and EW, arXiv:1108.4419 [hep-th]]

Bulk/brane emission

Ratios of bulk/brane energy emission for scalar fields



	<i>n</i> = 1	<i>n</i> = 2	<i>n</i> = 3	<i>n</i> = 4	<i>n</i> = 5	<i>n</i> = 6	<i>n</i> = 7
Standard BH	0.40	0.24	0.22	0.24	0.33	0.52	0.93
NCBH	0.27	0.08	0.03	0.009	0.003	0.001	0.0003

Harris and Kanti, hep-ph/0309054] Nicolini and EW, arXiv:1108.4419 [hep-th]]

Elizabeth Winstanley (Sheffield)

Conclusions

Standard Hawking radiation scenario

- Treats black hole geometry as purely classical, with quantum radiation on it
- Two phases of evolution:
 - Spin-down phase
 - Schwarzschild phase
- Both studied in great detail and well understood
- Notable exception is graviton emission in the spin-down phase
- Used in simulations of semi-classical BHs at the LHC

Conclusions

Beyond the semi-classical approximation?

- Use an effective metric incorporating some features of quantum BHs
- Many different approaches to this, including NCBHs
- Much lower temperatures than standard BHs
- Distinctive signatures in emission spectra
- NCBHs radiate almost entirely on the brane!