

On the Universality of CP Violation in $\Delta F=1$ Processes

Oram Gedalia
Weizmann Institute



OG, Jernej F. Kamenik, Zoltan Ligeti and Gilad Perez, arXiv:1202.5038

Introduction

- Measurements of flavor changing neutral currents in the quark sector put strong constraints on new physics (NP).
- Models can avoid bounds by aligning the flavor structure with the standard model (SM).
- Yet flavor-breaking sources involving only quark doublets cannot be simultaneously aligned with the up and down sectors.
- Robust constraints come from the *weakest* of the bounds.

Introduction

- Example – combining constraints from K and D mixing and CP violation (CPV)*.
- We argue that for $\Delta F=1$ processes involving quark doublets CPV is universal.
- Thus, the *strongest* constraint among the up and down sectors applies.

* K. Blum, Y. Grossman, Y. Nir and G. Perez, PRL **102**, 211802 (2009) [arXiv:0903.2118]

Outline

- Two generations
- Examples
 - $\epsilon'/\epsilon \Rightarrow \Delta a_{CP}$
 - $K_L \rightarrow \pi^0 e^+ e^- \Rightarrow a_{CP}(D^+ \rightarrow \pi^+ e^+ e^-)$
- Three generations
- Example
 - Semileptonic B decay
- Some SUSY implications
- Conclusion

Preliminaries

$$\mathcal{G}_F = SU(3)_Q \times SU(3)_U \times SU(3)_D$$

$$\mathcal{A}_u \equiv (Y_u Y_u^\dagger)_{\text{tr}} \quad \mathcal{A}_d \equiv (Y_d Y_d^\dagger)_{\text{tr}} \quad \Rightarrow \quad (8, 1, 1)$$
$$J \equiv i[\mathcal{A}_u, \mathcal{A}_d]$$

* OG, L. Mannelli and G. Perez, PLB **693**, 301 (2010) [arXiv:1002.0778]; JHEP **1010**, 046 (2010) [arXiv:1003.3869]

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$$\Rightarrow (8, 1, 1)$$

$SU(2)_L$ -invariant operator:

$$\mathcal{O}_L = \left[(X_L)^{ij} \bar{Q}_i \gamma^\mu Q_j \right] L_\mu$$

$$L_\mu = \sum_q \bar{q} \gamma_\mu q, \quad \sum_\ell \bar{\ell} \gamma_\mu \ell, \quad H^\dagger D_\mu H, \quad \dots$$

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Universality of CPV in Two Generations

$$X_L^d = \begin{pmatrix} a & c + id \\ c - id & -a \end{pmatrix}$$

CP violation in $s \rightarrow d$



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$$V^{\text{CKM}} = \begin{pmatrix} \cos \theta_C & \sin \theta_C \\ -\sin \theta_C & \cos \theta_C \end{pmatrix}$$

$$X_L^u = \begin{pmatrix} r & s + id \\ s - id & -r \end{pmatrix}$$

CP violation in $c \rightarrow u$



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CP violation in $s \rightarrow d$

$$V^{\text{CKM}} = \begin{pmatrix} \cos \theta_c & \sin \theta_c \\ -\sin \theta_c & \cos \theta_c \end{pmatrix}$$

$$d \propto X_L^{\text{CPV}}$$

$$X_L^u = \begin{pmatrix} r & s + id \\ s - id & -r \end{pmatrix}$$

CP violation in $c \rightarrow u$

Implication of the 3rd Generation

- Decompose X_L in terms of $SU(2)_Q$:

$$X_L = \begin{pmatrix} * & * & * \\ * & * & * \\ * & * & * \end{pmatrix}$$

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Diagram illustrating the decomposition of X_L into $SU(2)_Q$ representations:

- The top-left 2×2 block is labeled "adjoint" (red arrow).
- The top-right 2×1 column is labeled "Doublet - x_L " (blue arrow).
- The bottom-left 2×2 block is labeled "singlet" (green arrow).
- The bottom-right 1×1 element is labeled "singlet" (green arrow).

- $\Delta c, \Delta s=1$ processes are generated at order x_L^2 independently.
- Our argument applies to both contributions.

Example – Hadronic K and D Decays

- The bound on left-left (LL) operators from ϵ'/ϵ has been calculated in *.
- The weakest bound: $|\text{Im}(C_2^{(0)})| \lesssim 4.5 \times 10^{-5} \left(\frac{\Lambda_{\text{NP}}}{350 \text{ GeV}} \right)^2$
for $Q_2^{(0)} = (\bar{d}_\alpha s_\beta)_{V-A} \sum_q (\bar{q}_\beta q_\alpha)_{V-A}$
- The contribution to Δa_{CP} is:
$$\Delta a_{CP}^{\text{NP}} \approx 8.9 \sum_i \text{Im}(C_i^{\text{NP}}) \text{Im}(\Delta R_i^{\text{NP}})$$
- Hence for $\Delta R_i^{\text{NP}} \sim 1$: $\Delta a_{CP}^{\text{NP}} \lesssim 4 \times 10^{-4}$
- The contributions to Δa_{CP} are negligible!

* G. Isidori, J. F. Kamenik, Z. Ligeti and G. Perez, [arXiv:1111.4987]

Example – Semileptonic K and D Decays

- Analyze the CPV decay $K_L \rightarrow \pi^0 e^+ e^-$:

$$\frac{C_{sd}^{\ell_{R/L}}}{\Lambda_{\text{NP}}^2} (\bar{s}d)_{V-A} (\bar{\ell}\ell)_{V\pm A} \quad \Rightarrow \quad |\text{Im } C_{sd}^{e_{R/L}}| < 5.5 \times 10^{-4} \left(\frac{\Lambda_{\text{NP}}}{1 \text{ TeV}} \right)^2$$

$$\text{Br}(K_L \rightarrow \pi^0 e^+ e^-) < 2.8 \times 10^{-10} \quad (90\% \text{ CL})$$

- Apply the bound to the D CP asymmetry:

$$a_e^D \equiv \frac{\text{Br}(D^+ \rightarrow \pi^+ e^+ e^-) - \text{Br}(D^- \rightarrow \pi^- e^+ e^-)}{\text{Br}(D^+ \rightarrow \pi^+ e^+ e^-) + \text{Br}(D^- \rightarrow \pi^- e^+ e^-)}$$

Example – Semileptonic K and D Decays

- Assume that CPV is dominated by NP and decay rate is dominated by long distance SM:

$$|a_e^D| \lesssim \frac{2 \int d\rho |\text{Im}(\mathcal{A}_{\text{NP}})| |\mathcal{A}_{\text{SM}}|}{\int d\rho |\mathcal{A}_{\text{SM}}|^2} \lesssim 2 \sqrt{\frac{\int d\rho |\mathcal{A}_{\text{NP}}|^2}{\int d\rho |\mathcal{A}_{\text{SM}}|^2}}$$

- Use the experimental bound for the BR:

$$|a_e^D| \lesssim \left(\frac{1 \text{ TeV}}{\Lambda_{\text{NP}}} \right)^2 \frac{0.1 |\text{Im} C_{sd}^{e_{R/L}}|}{\sqrt{\text{Br}(D^+ \rightarrow \pi^+ e^+ e^-)}} \lesssim 0.02$$

- Robust bound on asymmetry from LL ops.

CPV Universality in 3rd Generation Decays

- Neglect the masses of the first 2 generations
 - CKM phase and 1-2 rotation become unphysical
 - Remove the 1-3 entry by additional $SU(2)$ rotation
- CKM is reduced to a rotation of $\theta \simeq \sqrt{\theta_{13}^2 + \theta_{23}^2}$

$$\mathcal{A}_d = \frac{y_b^2}{3} \begin{pmatrix} -1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & 2 \end{pmatrix} \quad \mathcal{A}_u = y_t^2 \begin{pmatrix} \spadesuit & 0 & 0 \\ 0 & \spadesuit & \spadesuit \\ 0 & \spadesuit & \spadesuit \end{pmatrix}$$

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protected by $U(1)_Q$

- Define “active”/”sterile” light flavors

CPV Universality in 3rd Generation Decays

- The only measure of CPV – $\text{Tr}(X_L \cdot J)$
 - ➔ $\text{Im}(X_L^u)_{a3} = \text{Im}(X_L^d)_{a3}$
- Active states coincide with 2nd generation quarks up to $\mathcal{O}(\lambda_C)$.

CPV Universality in 3rd Generation Decays

- The only measure of CPV – $\text{Tr}(X_L \cdot J)$

$$\longrightarrow \text{Im}(X_L^u)_{a3} = \text{Im}(X_L^d)_{a3}$$

- Active states coincide with 2nd generation quarks up to $\mathcal{O}(\lambda_C)$.

- Corrections – decompose X_L in terms of $SU(2)_{3a}$:

- Doublet constrained by $2 \rightarrow 1$ CPV transitions

$$X_L = \begin{pmatrix} \left[\begin{matrix} * \end{matrix} \right] & \left[\begin{matrix} * & * \end{matrix} \right] \\ \left[\begin{matrix} * \\ * \end{matrix} \right] & \left[\begin{matrix} * & * \\ * & * \end{matrix} \right] \end{pmatrix}$$

← Doublet
← singlet
← adjoint

- $m_c^2/m_t^2, m_s^2/m_b^2$

Example – Semileptonic B Decays

- **Observable – $B \rightarrow X_s \ell^+ \ell^-$ at $q^2 \equiv (p_{\ell^+} + p_{\ell^-})^2 \in [1, 6] \text{ GeV}^2$**

$$\frac{C_{bs}^H}{\Lambda_{\text{NP}}^2} (\bar{b}s)_{V-A} (H^\dagger D H) \quad \text{Br}(B \rightarrow X_s \ell^+ \ell^-)_{\text{low}} = (1.60 \pm 0.50) \times 10^{-6}$$

$$\longrightarrow |\text{Im}(C_{bs}^H)| < 8.7 \times 10^{-3} \left(\frac{\Lambda_{\text{NP}}}{1 \text{ TeV}} \right)^2$$


- **LHC sensitivity with 100 fb^{-1} at 14 TeV:**

$$|C_{tc(u)}^H| \lesssim 0.2 (\Lambda_{\text{NP}}/\text{TeV})^2$$

- Any CPV signal in $t \rightarrow cZ$ would have to be from $SU(3)_U$ breaking NP.

SUSY Implication

- The only source of $SU(3)_Q$ breaking is \tilde{m}_Q^2 , which is approximately $SU(2)_L$ -invariant.
- Define $\delta_{LL}^{ij} \equiv (\tilde{m}_Q^2)^{ij} / \bar{m}_Q^2$ with $\bar{m}_Q \equiv (m_{\tilde{Q}_1} + m_{\tilde{Q}_2})/2$
- Bound from ϵ'/ϵ : $\text{Im } \delta_{LL}^{12} \leq 0.5$ for $\bar{m}_{\tilde{Q}} = m_{\tilde{g}} = 500$ GeV

 $\delta_Q^{12} \equiv \frac{m_{\tilde{Q}_2} - m_{\tilde{Q}_1}}{m_{\tilde{Q}_2} + m_{\tilde{Q}_1}} \leq 0.25 \left(\frac{500 \text{ GeV}}{\bar{m}_{\tilde{Q}}} \right)$

- Not the strongest constraint *but* could have been derived more than 20 years ago!

Degeneracy of Squarks

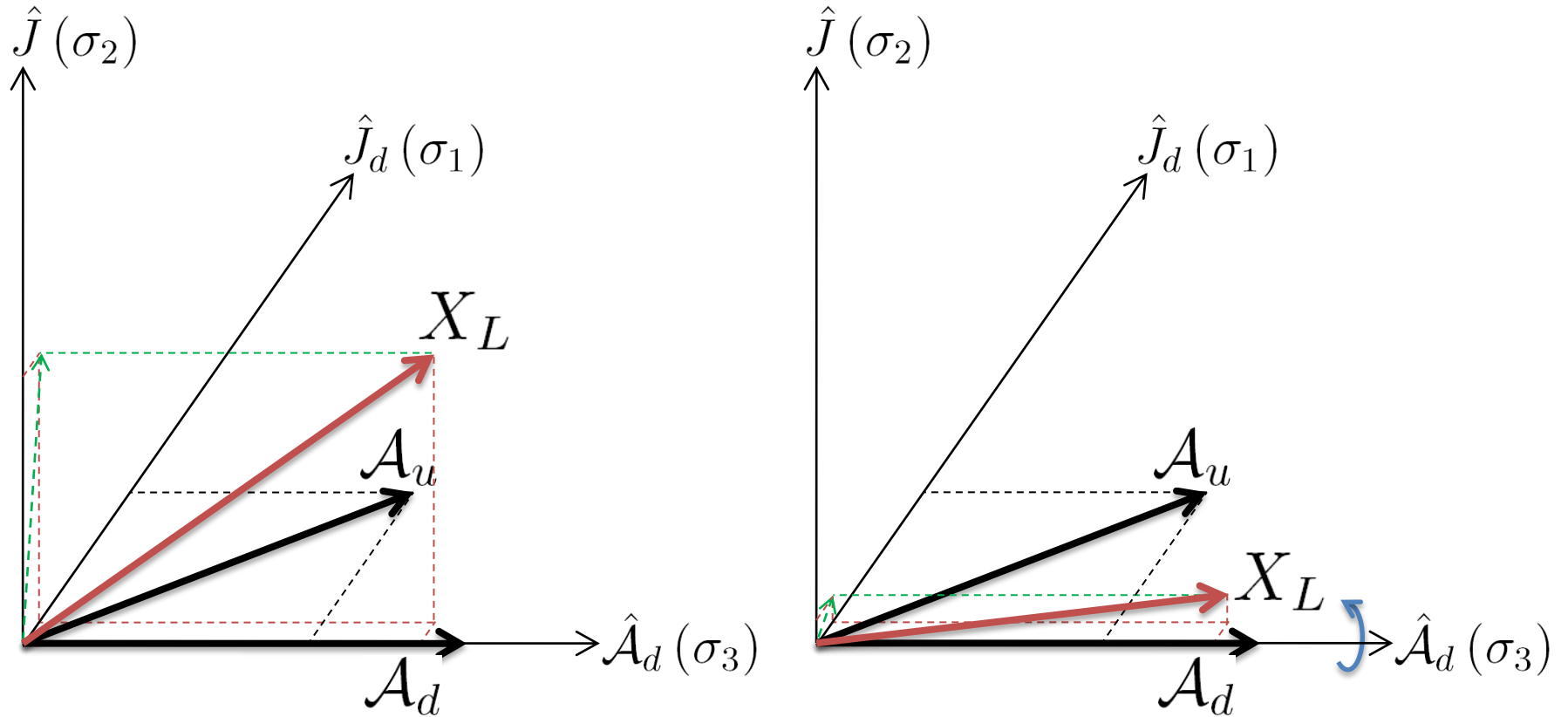
- Bounds on alignment models from K and D mixing and CPV*:

$$\begin{aligned} \mathcal{I}m(z_1^K) &= -\Lambda_{12}^2 \sin \alpha \sin 2\gamma, & \frac{m_{\tilde{Q}_2} - m_{\tilde{Q}_1}}{m_{\tilde{Q}_1} + m_{\tilde{Q}_2}} &\leq \begin{cases} 0.034 & \sin 2\gamma = 1 \\ 0.27 & \sin 2\gamma = 0 \end{cases} \\ \mathcal{I}m(z_1^D) &= -\Lambda_{12}^2 \sin(\alpha - 2\theta_c) \sin 2\gamma \end{aligned}$$

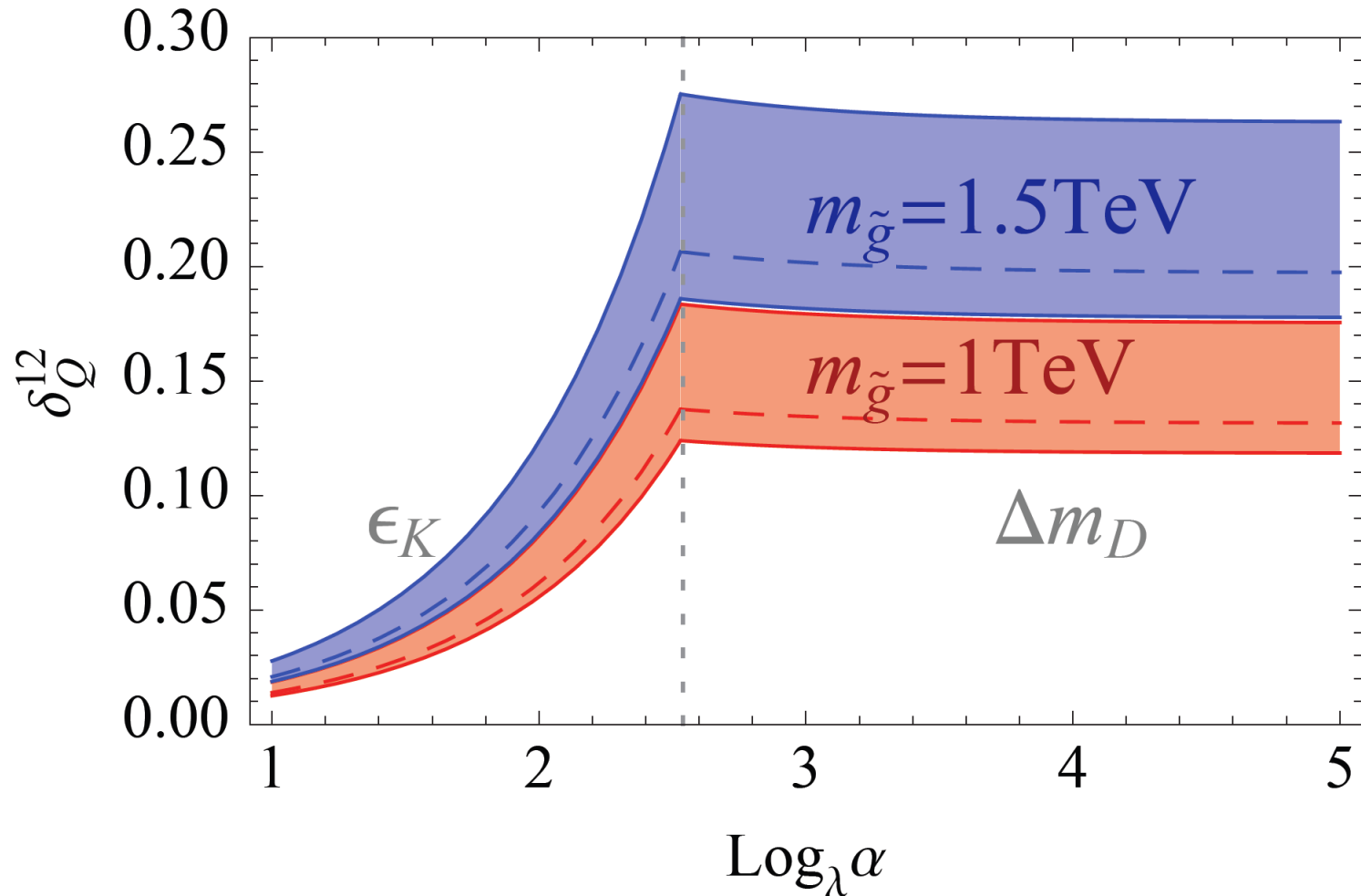
- The maximal phase case does not correspond to an alignment model.
- Alignment makes both real and imaginary parts small.

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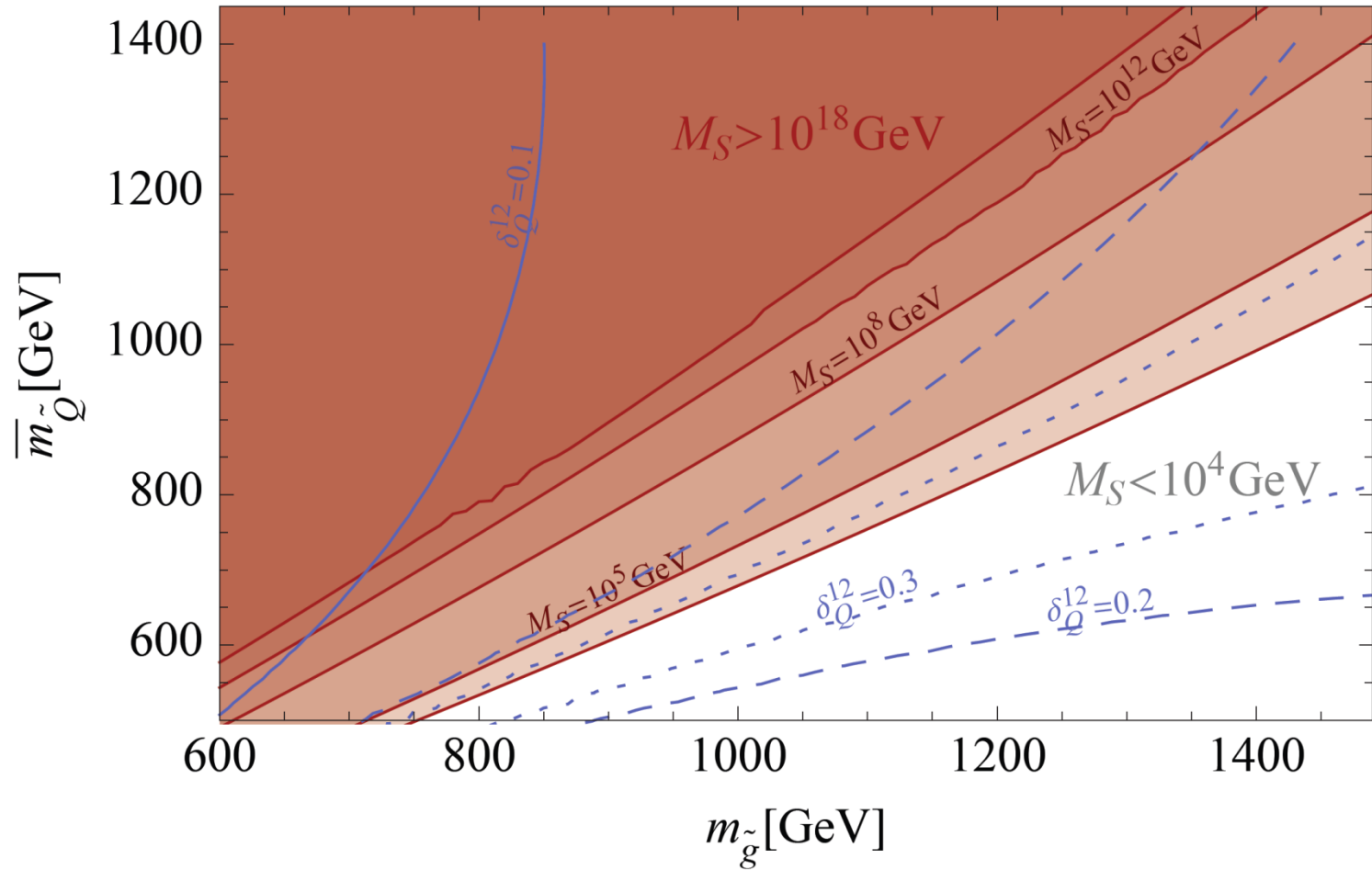


Degeneracy of Squarks

- No strong degeneracy required!
- Ex.: $m_{\tilde{g}}=1.3$ TeV, $m_{\tilde{Q}_1}=550$ GeV, $m_{\tilde{Q}_2}=950$ GeV
- This can be generated by*:
 - Anarchy at the SUSY breaking mediation scale
 - SUSY renormalization group flow to the TeV scale
 - Can lead to modest level of degeneracy

* Y. Nir and G. Raz, PRD **66**, 035007 (2002) [hep-ph/0206064]

Degeneracy of Squarks



Conclusion

- NP sources breaking $SU(3)_Q$ flavor symmetry induce universal CPV in $\Delta F=1$ processes.
- The *strongest* bound between up and down sectors applies.
 - Two generations: ϵ'/ϵ prohibits $SU(3)_Q$ -breaking explanations to Δa_{CP} measured by LHCb.
 - Three generations.
- SUSY alignment models can be consistent with data even with a low SUSY breaking mediation scale.

Conclusion

 Goodbye
blue stain!

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