

HIGGS HUNTING FROM THE BOTTOM UP

COLLABORATIVE SEARCH STRATEGIES AT THE LHC

LHC₂TSP: CERN - MARCH 2012

Jamison Galloway

based on work with A. Azatov and R. Contino

(arXiv:1202.3415 and in progress)



SAPIENZA
UNIVERSITÀ DI ROMA

HIGGS HUNTING FROM THE BOTTOM UP

COLLABORATIVE SEARCH STRATEGIES AT THE LHC*

LHC₂TSP: CERN - MARCH 2012

Jamison Galloway

based on work with A. Azatov and R. Contino

(arXiv:1202.3415 and in progress)



SAPIENZA
UNIVERSITÀ DI ROMA


*Or, a demonstration of principles described in yesterday's discussion of the Les Houches Recommendations

Entering a data-driven era...

- Focus on bridging any gaps between theory and experiment:
- Identify the essential theory that is useful for experimentalists, and...
- The essential experiment (statistics) for theorists

Entering a data-driven era...


- Focus on bridging any gaps between theory and experiment:
- Identify the essential theory that is useful for experimentalists, and...
- The essential experiment (statistics) for theorists



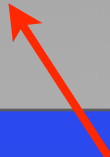
Establishes necessary
knowledge for how to
report data:
not just to make
interpretation easier,
but to make it
possible at all
(examples to come)

Entering a data-driven era...

- Focus on bridging any gaps between theory and experiment:
- Identify the essential theory that is useful for experimentalists, and...
- The essential experiment (statistics) for theorists



Establishes necessary knowledge for how to report data:
not just to make interpretation easier,
but to make it
possible at all
(examples to come)



Which models are ruled out, which are still in play?
Answering accurately requires an understanding of essential statistics...

OUTLINE

It is in exploring the common ground between theory that one can isolate the important points regarding efficient communication of results

*inasmuch as one is interested in constructing likelihoods and exclusions in *general* parameter spaces*

OUTLINE

It is in exploring the common ground between theory that one can isolate the **important points** regarding efficient communication of results

*inasmuch as one is interested in constructing likelihoods and exclusions in *general* parameter spaces*

Three such items by our count
= emphasis of this talk...

Getting at answers with an obvious example: The Higgs

The theory we know has to be augmented (unitarity, renorm'ability):

Getting at answers with an obvious example: The Higgs

The theory we **know** has to be augmented (unitarity, renorm'ability):

Three massive vectors, triplet of approximate $SU(2)$

$$U = \exp [2i\tau_a \pi_a(x)/v]$$

$$\mapsto LUR^\dagger$$

described at leading order:

$$\begin{aligned} \Delta\mathcal{L} = & \frac{v^2}{4} \text{tr} [(D_\mu U)^\dagger (D^\mu U)] \\ & - \frac{v}{\sqrt{2}} \psi_i^c U^\dagger \times \lambda_{ij} \psi_j + \text{h.c.} \end{aligned}$$

Getting at answers with an obvious example: The Higgs

The theory we **know** has to be **augmented** (unitarity, renorm'ability):

Three massive vectors, triplet of approximate $SU(2)$

$$U = \exp [2i\tau_a \pi_a(x)/v]$$
$$\mapsto LUR^\dagger$$

described at leading order:

$$\Delta\mathcal{L} = \frac{v^2}{4} \text{tr} [(D_\mu U)^\dagger (D^\mu U)] \times \left(1 + 2a \frac{h}{v} + b \frac{h^2}{v^2} + \dots \right)$$
$$- \frac{v}{\sqrt{2}} \psi_i^c U^\dagger \times \lambda_{ij} \psi_j + \text{h.c.} \times \left(1 + c \frac{h}{v} + \dots \right)$$

Assumption: the (custodial singlet) 'Higgs' might not be single-handedly responsible for unitarization, etc.

OTHER NEW PHYSICS enters at potentially low scales,

$$\text{e.g. } \Lambda_{\text{C.H.}} \sim 4\pi v / \sqrt{1 - a^2} \ll M_{\text{P}} .$$

Likewise "a" tells us about other Higgses in the spectrum

e.g. $a = 1 \Rightarrow H^\pm, A^0$ completely decoupled in MSSM

Getting at answers with an obvious example: The Higgs

The theory we **know** has to be **augmented** (unitarity, renorm'ability):

Three massive vectors, triplet of approximate $SU(2)$

$$U = \exp [2i\tau_a \pi_a(x)/v]$$

$$\mapsto LUR^\dagger$$

described at leading order:

$$\Delta\mathcal{L} = \frac{v^2}{4} \text{tr} [(D_\mu U)^\dagger (D^\mu U)] \times \left(1 + 2a \frac{h}{v} + b \frac{h^2}{v^2} + \dots \right) \\ - \frac{v}{\sqrt{2}} \psi_i^c U^\dagger \times \lambda_{ij} \psi_j + \text{h.c.} \times \left(1 + c \frac{h}{v} + \dots \right)$$

Focus on:

$$a \equiv \frac{g}{g_{\text{SM}}} \\ c \equiv \frac{y}{y_{\text{SM}}}$$

Getting at answers with an obvious example: The Higgs

The theory we **know** has to be **augmented** (unitarity, renorm'ability):

Three massive vectors, triplet of approximate $SU(2)$

$$U = \exp [2i\tau_a \pi_a(x)/v]$$
$$\mapsto LUR^\dagger$$

described at leading order:

$$\Delta\mathcal{L} = \frac{v^2}{4} \text{tr} [(D_\mu U)^\dagger (D^\mu U)] \times \left(1 + 2a \frac{h}{v} + b \frac{h^2}{v^2} + \dots \right)$$
$$- \frac{v}{\sqrt{2}} \psi_i^c U^\dagger \times \lambda_{ij} \psi_j + \text{h.c.} \times \left(1 + c \frac{h}{v} + \dots \right)$$

Focus on:

$$a \equiv \frac{g}{g_{\text{SM}}}$$
$$c \equiv \frac{y}{y_{\text{SM}}}$$

with theory predictions:

$$\star \text{ SM: } \begin{cases} a = b = 1 \\ c = 1 \end{cases}$$

$$\star \text{ "MCHM4": } \begin{cases} \xi \equiv v^2/f^2 \\ a = c = \sqrt{1-\xi} \end{cases}$$

$$\star \text{ "MCHM5": } \begin{cases} a = \sqrt{1-\xi} \\ c = \frac{1-2\xi}{\sqrt{1-\xi}} \end{cases}$$

Getting at answers with an obvious example: The Higgs

The theory we **know** has to be **augmented** (unitarity, renorm'ability):

Three massive vectors, triplet of approximate $SU(2)$

$$U = \exp [2i\tau_a \pi_a(x)/v]$$
$$\mapsto LUR^\dagger$$

described at leading order:

$$\Delta\mathcal{L} = \frac{v^2}{4} \text{tr} [(D_\mu U)^\dagger (D^\mu U)] \times \left(1 + 2a \frac{h}{v} + b \frac{h^2}{v^2} + \dots \right)$$
$$- \frac{v}{\sqrt{2}} \psi_i^c U^\dagger \times \lambda_{ij} \psi_j + \text{h.c.} \times \left(1 + c \frac{h}{v} + \dots \right)$$

Focus on:

with theory predictions:

$$a \equiv \frac{g}{g_{\text{SM}}}$$
$$c \equiv \frac{y}{y_{\text{SM}}}$$

A minimally-prejudiced question:

**HOW CAN WE
CONSTRAIN
THIS SPACE?**

$$\left\{ \begin{array}{l} \xi \equiv v^2 / f^2 \\ a = c = \sqrt{1 - \xi} \\ a = \sqrt{1 - \xi} \\ c = \frac{1 - 2\xi}{\sqrt{1 - \xi}} \end{array} \right.$$

Some “need to know” statistics

Given background, signal, and observed events: construct likelihood:

$$P(n|n_{\text{obs}}) = \frac{n^{n_{\text{obs}}} e^{-n}}{n_{\text{obs}}!} \times \pi(n)$$
$$\xrightarrow{\text{A.L.}} \exp\left[\frac{-(n - n_{\text{obs}})^2}{2n_{\text{obs}}}\right] \times \pi(n)$$

Some “need to know” statistics

Given **background**, **signal**, and observed events: construct likelihood:

$$P(n|n_{\text{obs}}) = \frac{n^{n_{\text{obs}}} e^{-n}}{n_{\text{obs}}!} \times \pi(n)$$
$$\xrightarrow{\text{A.L.}} \exp\left[\frac{-(n - n_{\text{obs}})^2}{2n_{\text{obs}}}\right] \times \pi(n)$$

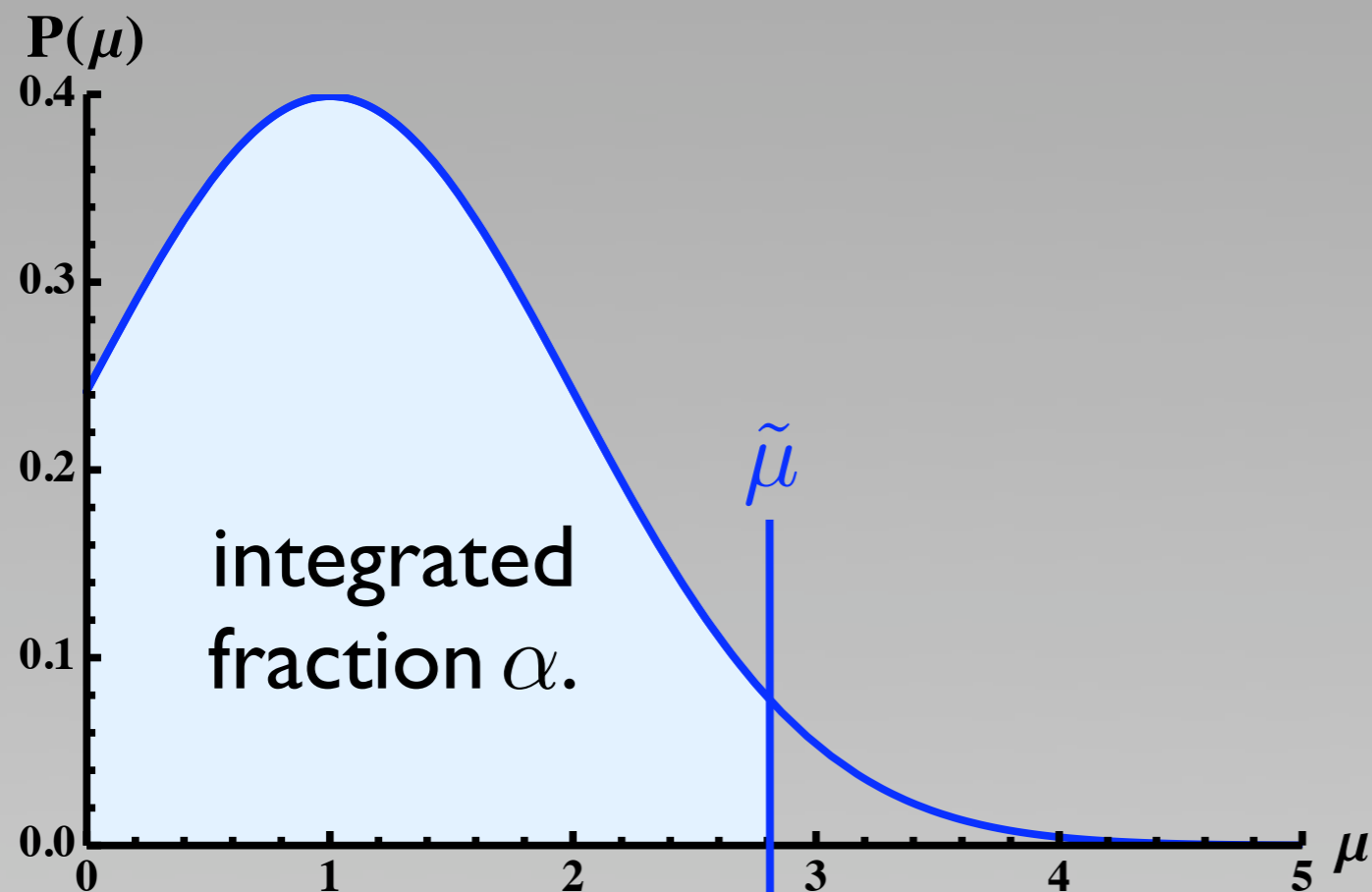
$$n = n_B + \mu n_S^{\text{SM}} \Rightarrow P(\mu) = \pi(\mu) \times \exp\left[\frac{(n_B + \mu n_S - n_{\text{obs}})^2}{2n_{\text{obs}}}\right]$$

Some “need to know” statistics

Given **background**, **signal**, and observed events: construct likelihood:

$$P(n|n_{\text{obs}}) = \frac{n^{n_{\text{obs}}} e^{-n}}{n_{\text{obs}}!} \times \pi(n)$$
$$\xrightarrow{\text{A.L.}} \exp\left[\frac{-(n - n_{\text{obs}})^2}{2n_{\text{obs}}}\right] \times \pi(n)$$

$$n = n_B + \mu n_S^{\text{SM}} \Rightarrow P(\mu) = \pi(\mu) \times \exp\left[\frac{(n_B + \mu n_S - n_{\text{obs}})^2}{2n_{\text{obs}}}\right]$$

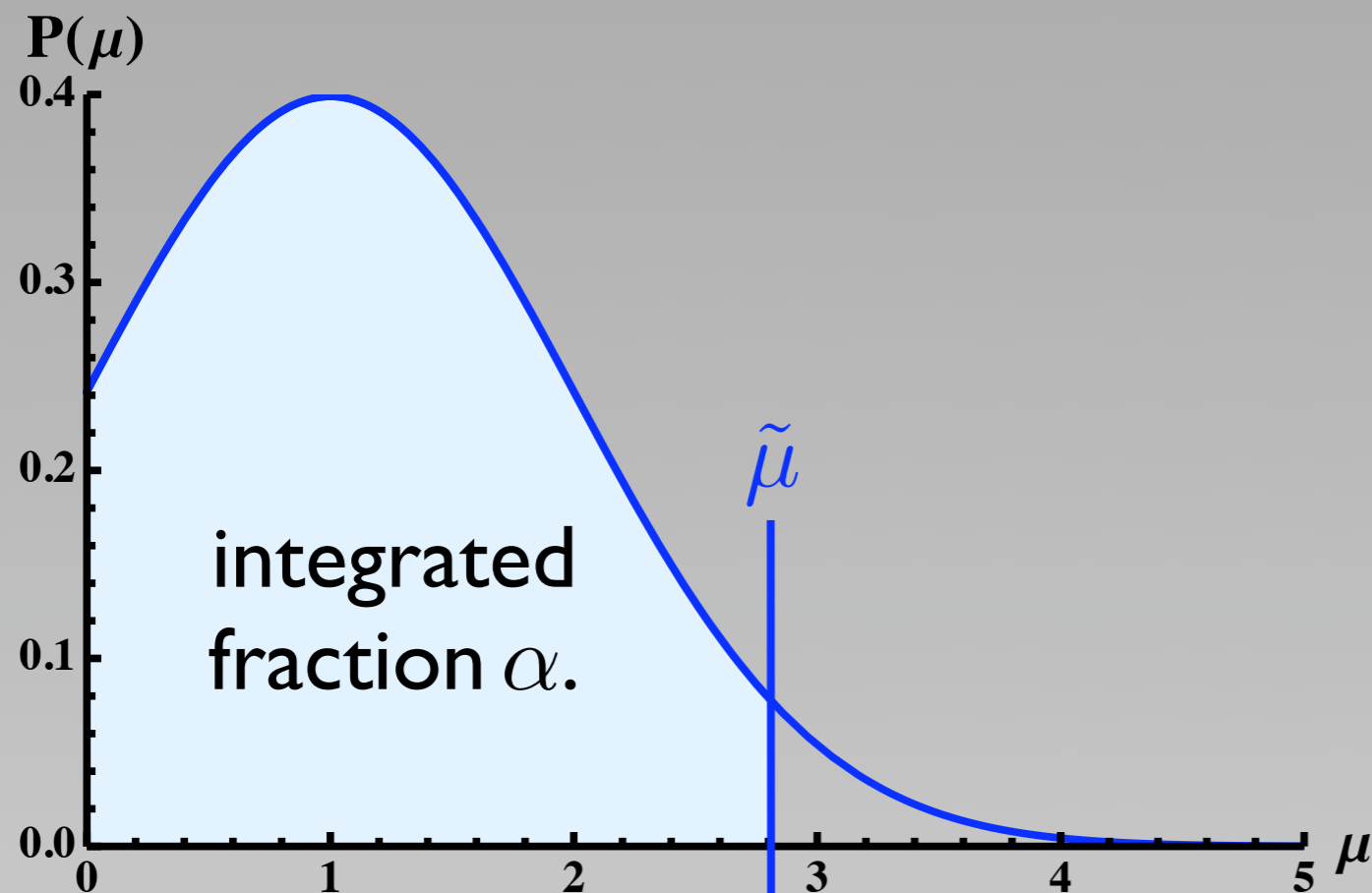


Some “need to know” statistics

Given **background**, **signal**, and observed events: construct likelihood:

$$P(n|n_{\text{obs}}) = \frac{n^{n_{\text{obs}}} e^{-n}}{n_{\text{obs}}!} \times \pi(n)$$
$$\xrightarrow{\text{A.L.}} \exp\left[\frac{-(n - n_{\text{obs}})^2}{2n_{\text{obs}}}\right] \times \pi(n)$$

$$n = n_B + \mu n_S^{\text{SM}} \Rightarrow P(\mu) = \pi(\mu) \times \exp\left[\frac{(n_B + \mu n_S - n_{\text{obs}})^2}{2n_{\text{obs}}}\right]$$



$\tilde{\mu}$: upper bound on signal strength modifier at CL = alpha.

Two versions:

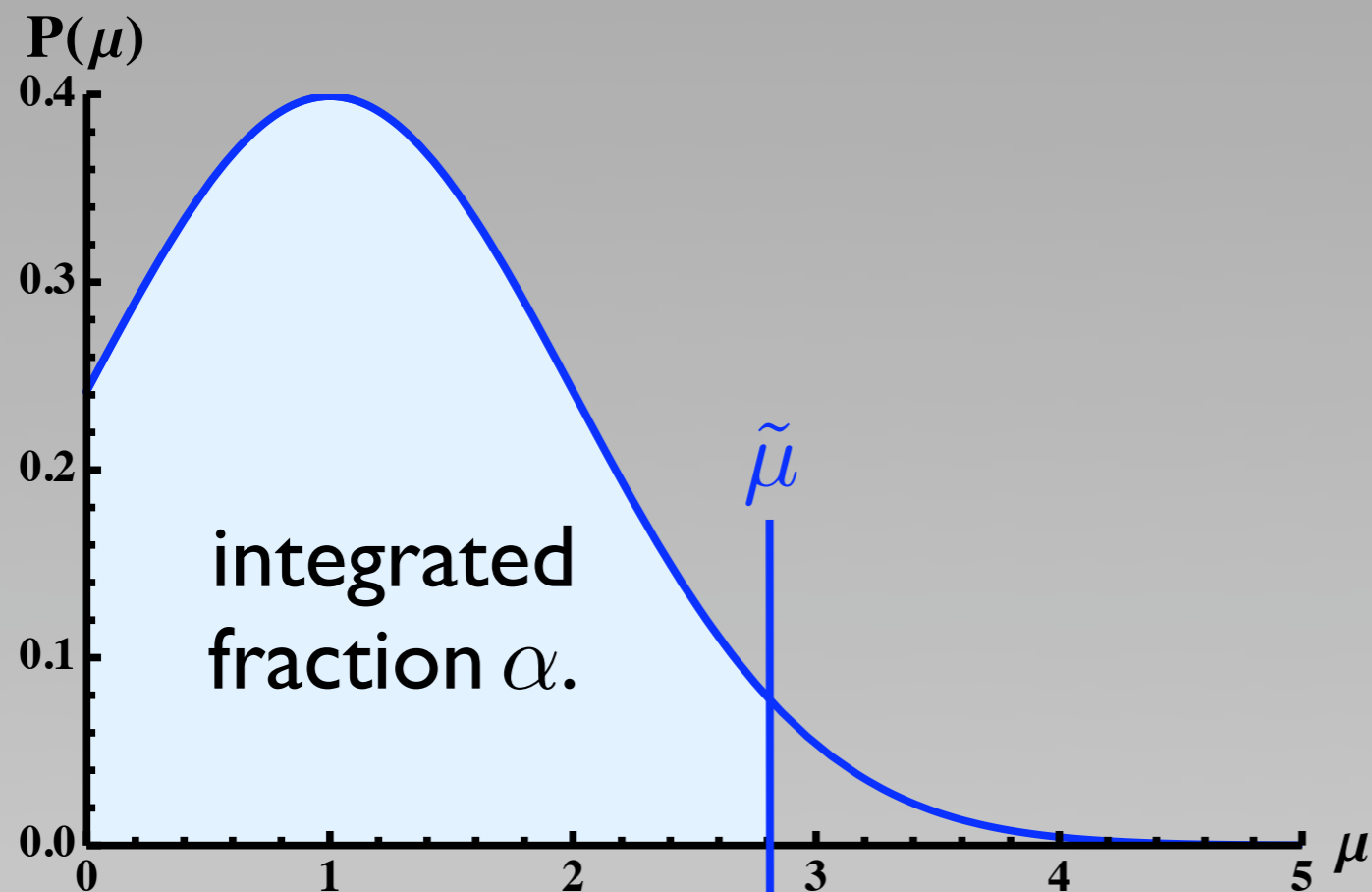
1. Expected (background only hypothesis)
2. Observed (compared to data)

Some “need to know” statistics

Given **background**, **signal**, and observed events: construct likelihood:

$$P(n|n_{\text{obs}}) = \frac{n^{n_{\text{obs}}} e^{-n}}{n_{\text{obs}}!} \times \pi(n)$$
$$\xrightarrow{\text{A.L.}} \exp\left[\frac{-(n - n_{\text{obs}})^2}{2n_{\text{obs}}}\right] \times \pi(n)$$

$$n = n_B + \mu n_S^{\text{SM}} \Rightarrow P(\mu) = \pi(\mu) \times \exp\left[\frac{(n_B + \mu n_S - n_{\text{obs}})^2}{2n_{\text{obs}}}\right]$$



$\tilde{\mu}$: upper bound on signal strength modifier at CL = alpha.

Two versions:

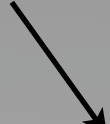
1. Expected (background only hypothesis)
2. Observed (compared to data)

A closer look at “signal strength modifier”

We want to compare number of *observed* signal events in SM units:

A closer look at “signal strength modifier”

We want to compare number of *observed* signal events in SM units:

$$n_S^{(i)} = \left(\int dt \mathcal{L} \right) \times \sum_p \sigma_p^{(i)} \times \zeta_{p,i} \times \text{BR}(h \rightarrow i)$$


$$\Rightarrow \mu = \frac{\sum_p \sigma_p^{(i)} \times \zeta_{p,i} \times \text{BR}(h \rightarrow i)}{\left[\sum_p \sigma_p^{(i)} \times \zeta_{p,i} \times \text{BR}(h \rightarrow i) \right]_{\text{SM}}}$$

A closer look at “signal strength modifier”

We want to compare number of *observed* signal events in SM units:

$$n_S^{(i)} = \left(\int dt \mathcal{L} \right) \times \sum_p \sigma_p^{(i)} \times \zeta_{p,i} \times \text{BR}(h \rightarrow i)$$
$$\Rightarrow \mu = \frac{\sum_p \sigma_p^{(i)} \times \zeta_{p,i} \times \text{BR}(h \rightarrow i)}{\left[\sum_p \sigma_p^{(i)} \times \zeta_{p,i} \times \text{BR}(h \rightarrow i) \right]_{\text{SM}}}$$

Efficiencies not always provided, so unknown to theorists

Best we can do: *assume* that $\zeta_{p,i} = \zeta_i \forall p$.[†]

[†] Safely justified for SM and SM-like ($a = c$), but not in general.

A closer look at “signal strength modifier”

We want to compare number of *observed* signal events in SM units:

$$n_S^{(i)} = \left(\int dt \mathcal{L} \right) \times \sum_p \sigma_p^{(i)} \times \zeta_{p,i} \times \text{BR}(h \rightarrow i)$$
$$\Rightarrow \mu = \frac{\sum_p \sigma_p^{(i)} \times \zeta_{p,i} \times \text{BR}(h \rightarrow i)}{\left[\sum_p \sigma_p^{(i)} \times \zeta_{p,i} \times \text{BR}(h \rightarrow i) \right]_{\text{SM}}}$$

Efficiencies not always provided, so unknown to theorists

Best we can do: *assume* that $\zeta_{p,i} = \zeta_i \forall p$.[†]

$$\mu \rightarrow \left. \frac{\sum_p \sigma_p^{(i)} \times \text{BR}(h \rightarrow i)}{\left[\sum_p \sigma_p^{(i)} \times \text{BR}(h \rightarrow i) \right]_{\text{SM}}} \right\} \begin{array}{l} \text{This can be related purely to theory, but} \\ \text{it's only approximate} \end{array}$$

EFFICIENCIES NEEDED

[†] Safely justified for SM and SM-like ($a = c$), but not in general.

A closer look at “signal strength modifier”

*****IMPORTANT POINT NUMBER ONE*****

Tracing back from number of events to underlying theory
REQUIRES knowledge of efficiencies
when considering generalized spaces.

Moving on: Comparison to Likelihood

From here just map theory parameters to μ and compare to $P(\mu)$...

Moving on: Comparison to Likelihood

From here just map theory parameters to μ and compare to $P(\mu)$...

NOT YET AVAILABLE *

Moving on: Comparison to Likelihood

From here just map theory parameters to μ and compare to $P(\mu)$...

NOT YET AVAILABLE *

* but let's not despair
(although it would be nice to have)

Moving on: Comparison to Likelihood

*****IMPORTANTish POINT NUMBER TWO*****

Constructing exclusions for underlying theory
WOULD BE HELPED BY knowledge of likelihoods directly.

Moving on: Comparison to *RECONSTRUCTED* Likelihood

(Three variables, only two constraints: we need to be slightly clever)

Assume asymptotic limit, i.e. Poisson \longrightarrow Gaussian:

$$P(n_B + \mu n_S | n_{\text{obs}}) = \pi(\mu) \times \exp \left[\frac{-(n_B + \mu n_S - n_{\text{obs}})^2}{2n_{\text{obs}}} \right]$$

Moving on: Comparison to *RECONSTRUCTED* Likelihood

(Three variables, only two constraints: we need to be slightly clever)

Assume asymptotic limit, i.e. Poisson \longrightarrow Gaussian:

$$P(n_B + \mu n_S | n_{\text{obs}}) = \pi(\mu) \times \exp \left[\frac{-(n_B + \mu n_S - n_{\text{obs}})^2}{2n_{\text{obs}}} \right]$$

$$\Rightarrow \tilde{\mu}_{\text{exp}}^{95\%} = 1.96 \times \frac{\sqrt{n_B}}{n_S}$$

Moving on: Comparison to *RECONSTRUCTED* Likelihood

(Three variables, only two constraints: we need to be slightly clever)

Assume asymptotic limit, i.e. Poisson \longrightarrow Gaussian:

$$P(n_B + \mu n_S | n_{\text{obs}}) = \pi(\mu) \times \exp \left[\frac{-(n_B + \mu n_S - n_{\text{obs}})^2}{2n_{\text{obs}}} \right]$$

$$\Rightarrow \tilde{\mu}_{\text{exp}}^{95\%} = 1.96 \times \frac{\sqrt{n_B}}{n_S}$$

For observed exclusion, use a simple rewriting:

$$P(n_B + \mu n_S | n_{\text{obs}}) = \pi(\mu) \times \exp \left[-\frac{1}{2} \left(\mu \frac{n_S}{\sqrt{n_B}} \frac{\sqrt{n_B}}{\sqrt{n_{\text{obs}}}} + \delta \right)^2 \right]; \quad \delta \equiv \frac{n_B - n_{\text{obs}}}{\sqrt{n_{\text{obs}}}}$$

Moving on: Comparison to *RECONSTRUCTED* Likelihood

(Three variables, only two constraints: we need to be slightly clever)

Assume asymptotic limit, i.e. Poisson \longrightarrow Gaussian:

$$P(n_B + \mu n_S | n_{\text{obs}}) = \pi(\mu) \times \exp \left[\frac{-(n_B + \mu n_S - n_{\text{obs}})^2}{2n_{\text{obs}}} \right]$$

$$\Rightarrow \tilde{\mu}_{\text{exp}}^{95\%} = 1.96 \times \frac{\sqrt{n_B}}{n_S}$$

For observed exclusion, use a simple rewriting:

$$P(n_B + \mu n_S | n_{\text{obs}}) = \pi(\mu) \times \exp \left[-\frac{1}{2} \left(\mu \frac{n_S}{\sqrt{n_B}} \frac{\sqrt{n_B}}{\sqrt{n_{\text{obs}}}} + \delta \right)^2 \right]; \quad \delta \equiv \frac{n_B - n_{\text{obs}}}{\sqrt{n_{\text{obs}}}}$$

Now make the assumption $\frac{n_{\text{obs}} - n_B}{n_{\text{obs}}} \ll 1$

Moving on: Comparison to *RECONSTRUCTED* Likelihood

(Three variables, only two constraints: we need to be slightly clever)

Assume asymptotic limit, i.e. Poisson \longrightarrow Gaussian:

$$P(n_B + \mu n_S | n_{\text{obs}}) = \pi(\mu) \times \exp \left[\frac{-(n_B + \mu n_S - n_{\text{obs}})^2}{2n_{\text{obs}}} \right]$$

$$\Rightarrow \tilde{\mu}_{\text{exp}}^{95\%} = 1.96 \times \frac{\sqrt{n_B}}{n_S}$$

For observed exclusion, use a simple rewriting:

$$P(n_B + \mu n_S | n_{\text{obs}}) = \pi(\mu) \times \exp \left[-\frac{1}{2} \left(\mu \frac{n_S}{\sqrt{n_B}} \frac{\sqrt{n_B}}{\sqrt{n_{\text{obs}}}} + \delta \right)^2 \right]; \quad \delta \equiv \frac{n_B - n_{\text{obs}}}{\sqrt{n_{\text{obs}}}}$$

Now make the assumption $\frac{n_{\text{obs}} - n_B}{n_{\text{obs}}} \ll 1$

Moving on: Comparison to *RECONSTRUCTED* Likelihood

(Three variables, only two constraints: we need to be slightly clever)

Assume asymptotic limit, i.e. Poisson \longrightarrow Gaussian:

$$P(n_B + \mu n_S | n_{\text{obs}}) = \pi(\mu) \times \exp \left[\frac{-(n_B + \mu n_S - n_{\text{obs}})^2}{2n_{\text{obs}}} \right]$$

$$\Rightarrow \tilde{\mu}_{\text{exp}}^{95\%} = 1.96 \times \frac{\sqrt{n_B}}{n_S}$$

For observed exclusion, use a simple rewriting:

$$P(n_B + \mu n_S | n_{\text{obs}}) = \pi(\mu) \times \exp \left[-\frac{1}{2} \left(\mu \frac{n_S}{\sqrt{n_B}} \frac{\sqrt{n_B}}{\sqrt{n_{\text{obs}}}} + \delta \right)^2 \right]; \quad \delta \equiv \frac{n_B - n_{\text{obs}}}{\sqrt{n_{\text{obs}}}}$$

Now make the assumption $\frac{n_{\text{obs}} - n_B}{n_{\text{obs}}} \ll 1$

Moving on: Comparison to *RECONSTRUCTED* Likelihood

$$P(\mu) = N \times \exp \left[-\frac{1}{2} \left(\frac{1.96 \times \mu}{\tilde{\mu}_{\text{exp}}^{(95\%)}} + \delta \right)^2 \right]$$

Moving on: Comparison to *RECONSTRUCTED* Likelihood

$$P(\mu) = N \times \exp \left[-\frac{1}{2} \left(\frac{1.96 \times \mu}{\tilde{\mu}_{\text{exp}}^{(95\%)}} + \delta \right)^2 \right]$$

Solve for remaining parameter using observed exclusion limit:

$$0.95 = \int_0^{\tilde{\mu}_{\text{obs}}^{(95\%)}} d\mu P(\mu)$$

Moving on: Comparison to *RECONSTRUCTED* Likelihood

$$P(\mu) = N \times \exp \left[-\frac{1}{2} \left(\frac{1.96 \times \mu}{\tilde{\mu}_{\text{exp}}^{(95\%)}} + \delta \right)^2 \right]$$

Solve for remaining parameter using observed exclusion limit:

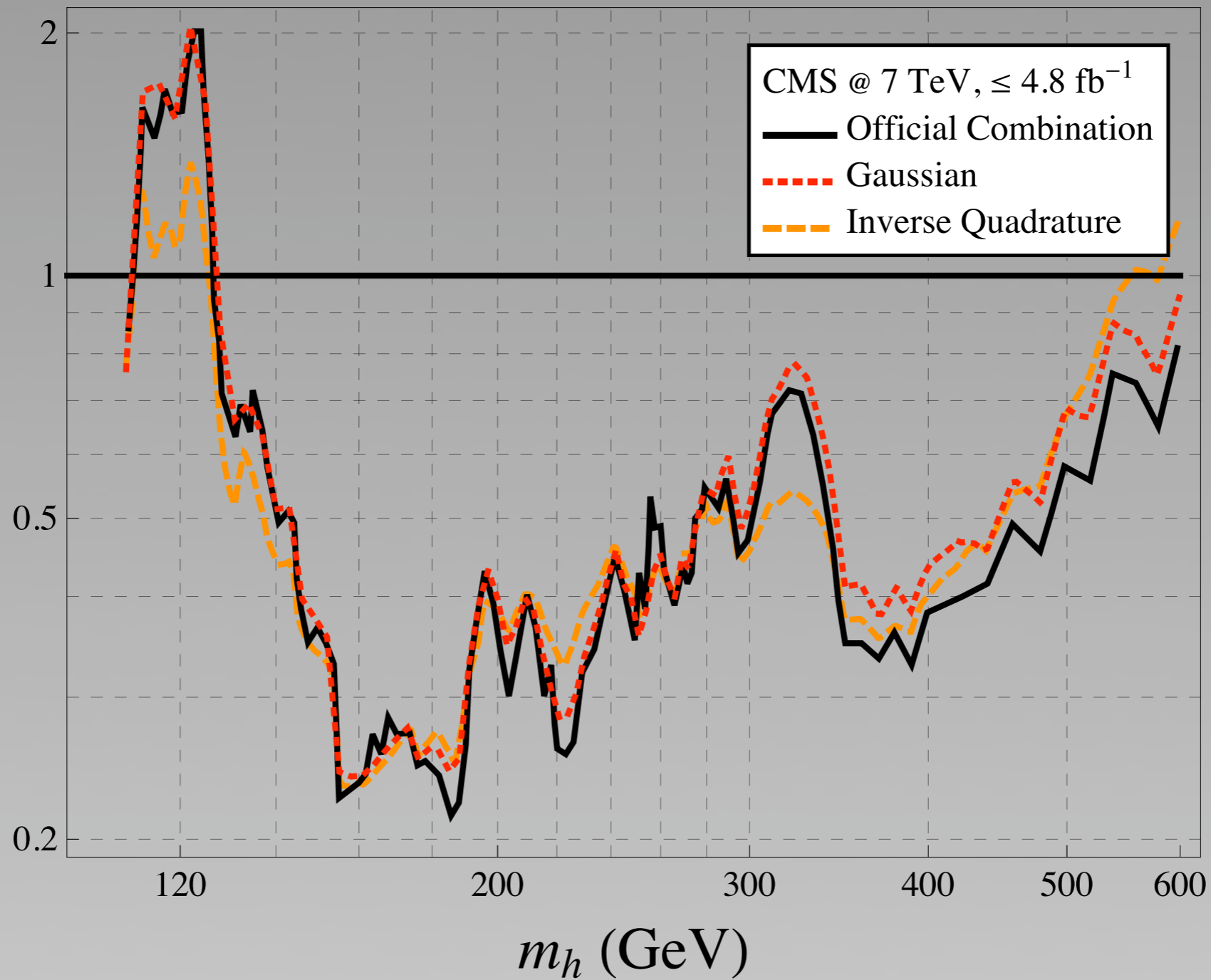
$$0.95 = \int_0^{\tilde{\mu}_{\text{obs}}^{(95\%)}} d\mu P(\mu)$$

RECAP:

- o Expected exclusion tells us about s/b
- o Observed tells us delta, completes determination of (AL) likelihood
- o Good news: can be done over whole mass range, not just at 'peaks' where information on best fit is available

How well does this method do?

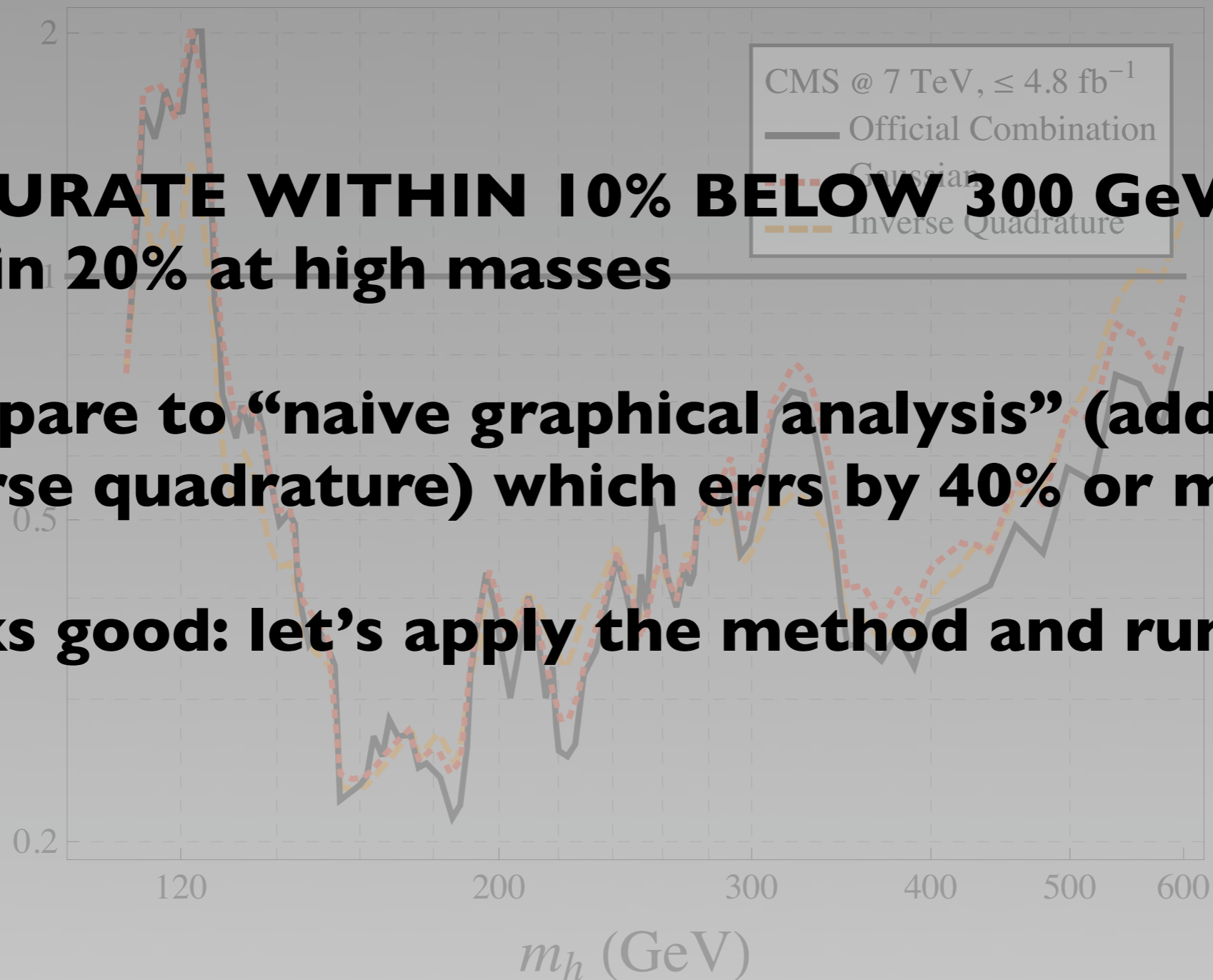
One possible check: the total combination



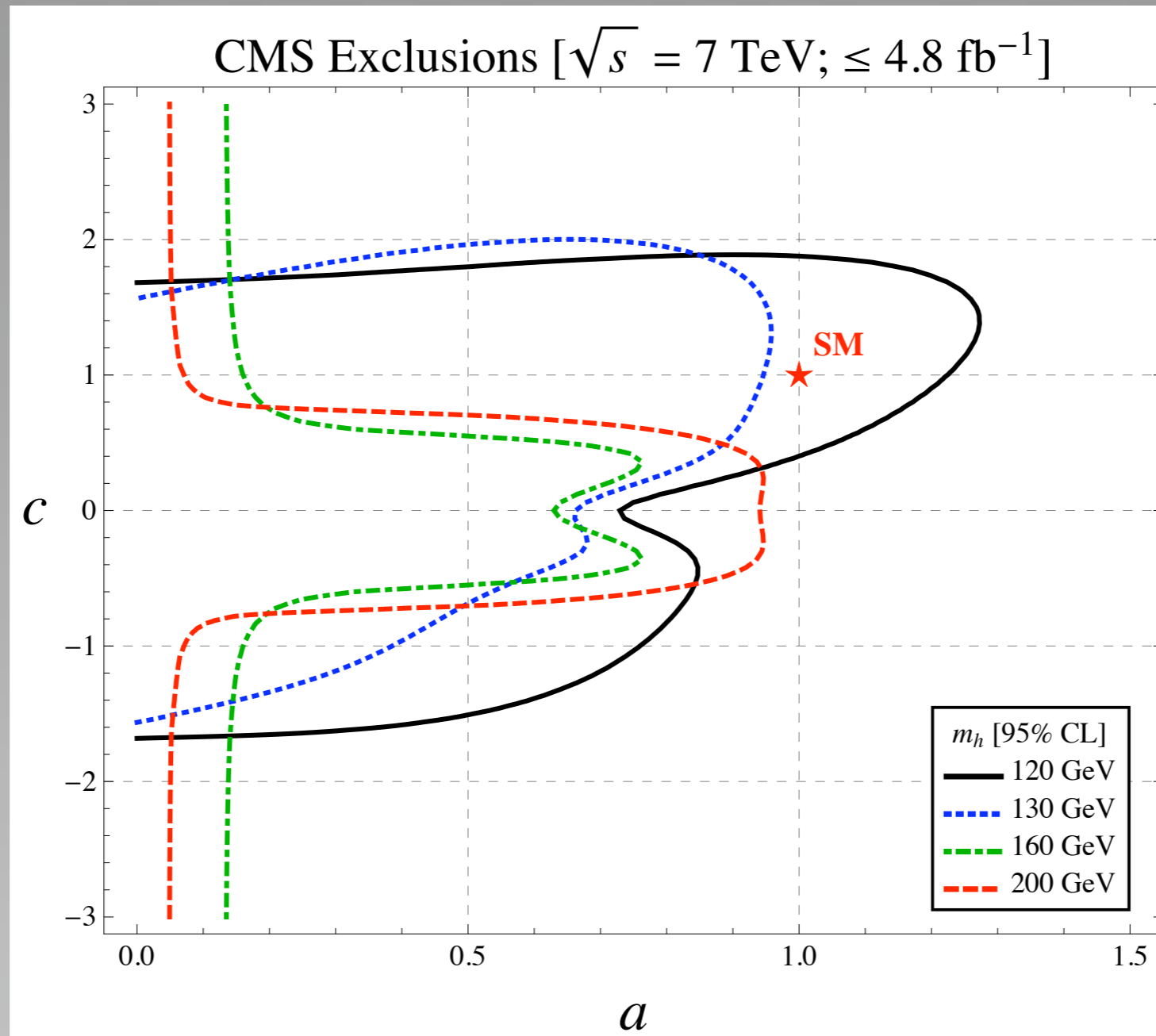
How well does this method do?

One possible check: the total combination

- o **ACCURATE WITHIN 10% BELOW 300 GeV; within 20% at high masses**
- o **Compare to “naive graphical analysis” (adding in inverse quadrature) which errs by 40% or more**
- o **Looks good: let’s apply the method and run with it**



Status report for unpopular mass points



Status report for the Higgs at 125(?)(!)

Five channels for a light Higgs:

1. WW
2. $\gamma\gamma$
3. ZZ
4. $\tau\tau$
5. bb

Status report for the Higgs at 125(?)(!)

Five channels for a light Higgs:

1. WW 2. $\gamma\gamma$ 3. ZZ 4. $\tau\tau$ 5. bb

- 1, 2. Zero Jet, same/opposite flavor lepton } ← Inclusive
3, 4. One Jet, same/opposite flavor lepton }
5. Two Jets ← VBF

Status report for the Higgs at 125(?)(!)

Five channels for a light Higgs:

1. WW 2. $\gamma\gamma$ 3. ZZ 4. $\tau\tau$ 5. bb

1. Both in barrel, $\min(R_9) > 0.94$
2. Both in barrel, $\min(R_9) < 0.94$
3. \geq One in endcap, $\min(R_9) > 0.94$
4. \geq One in endcap, $\min(R_9) < 0.94$
5. Dijet tag
- ← Inclusive
- ← VBF

Photon candidates with high values of R_9 are mostly unconverted and have less background than those with lower values. Photon candidates in the barrel have less background than those in the endcap. For this reason it has been found useful to divide photon candidates into four categories and apply a different selection in each category, using more stringent requirements in categories with higher background and worse resolution.

Status report for the Higgs at 125(?)(!)

Five channels for a light Higgs:

1. WW
2. $\gamma\gamma$
3. ZZ
4. $\tau\tau$
5. bb

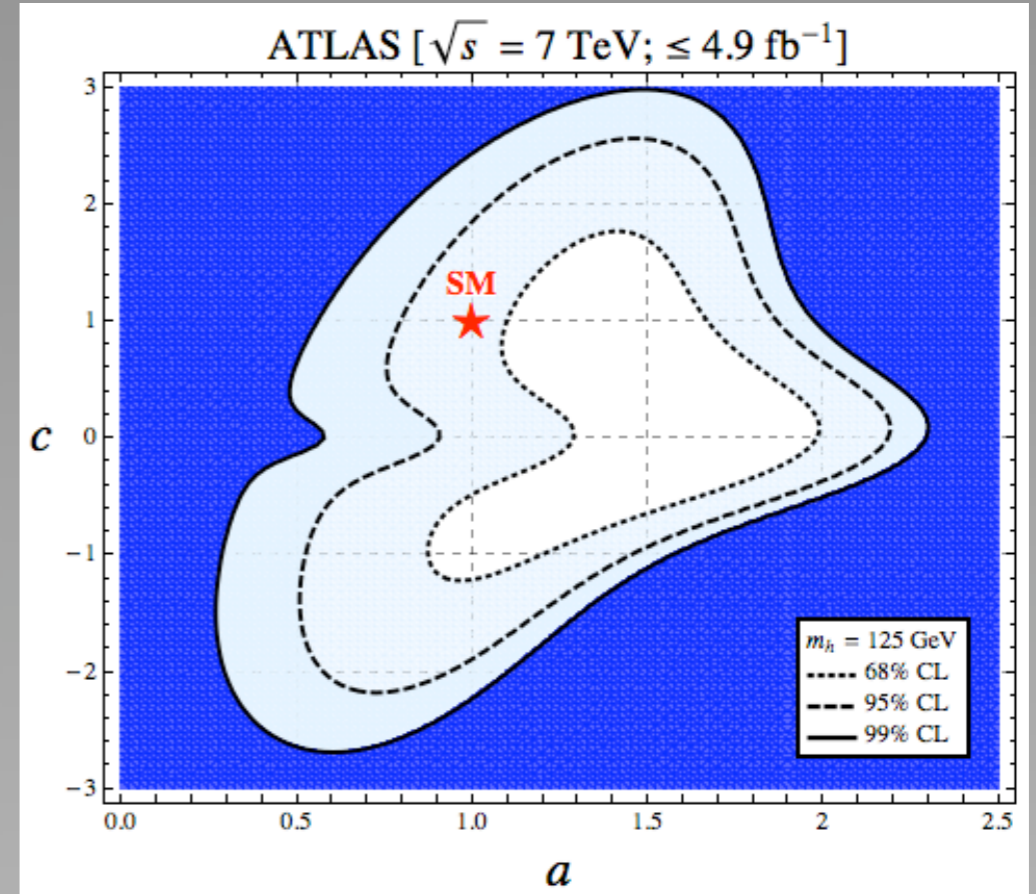
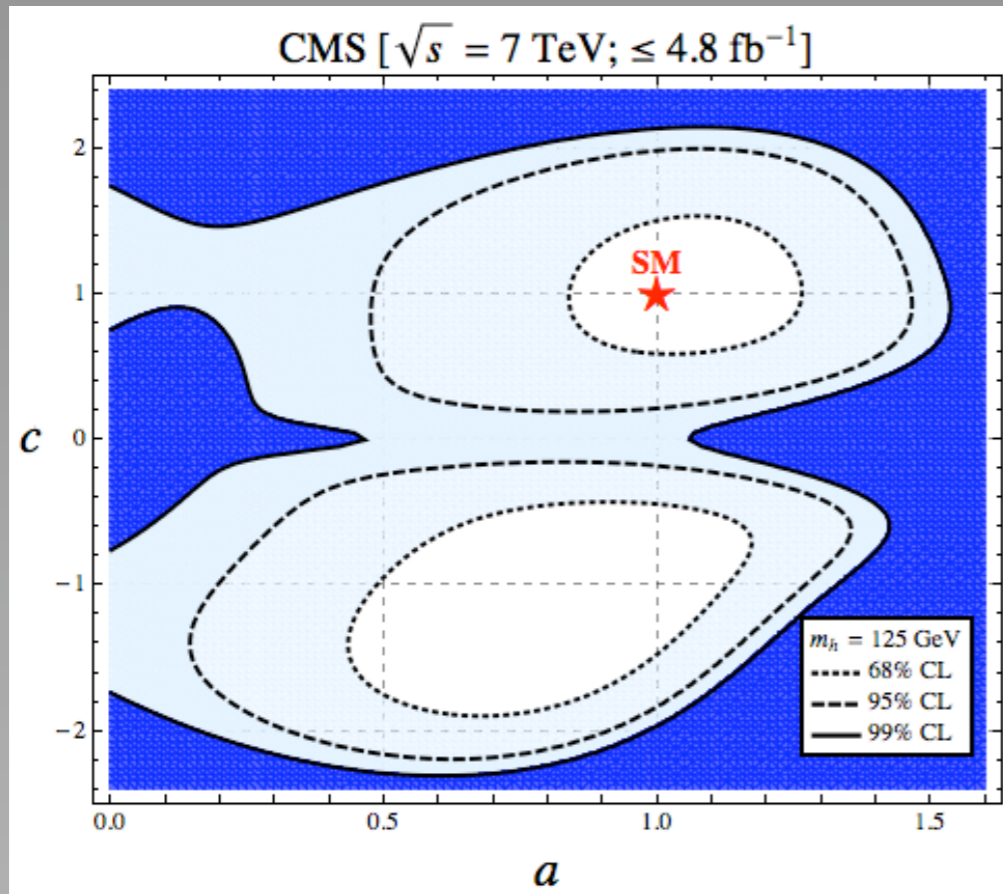
Inclusive

VBF + GF + “Boosted”

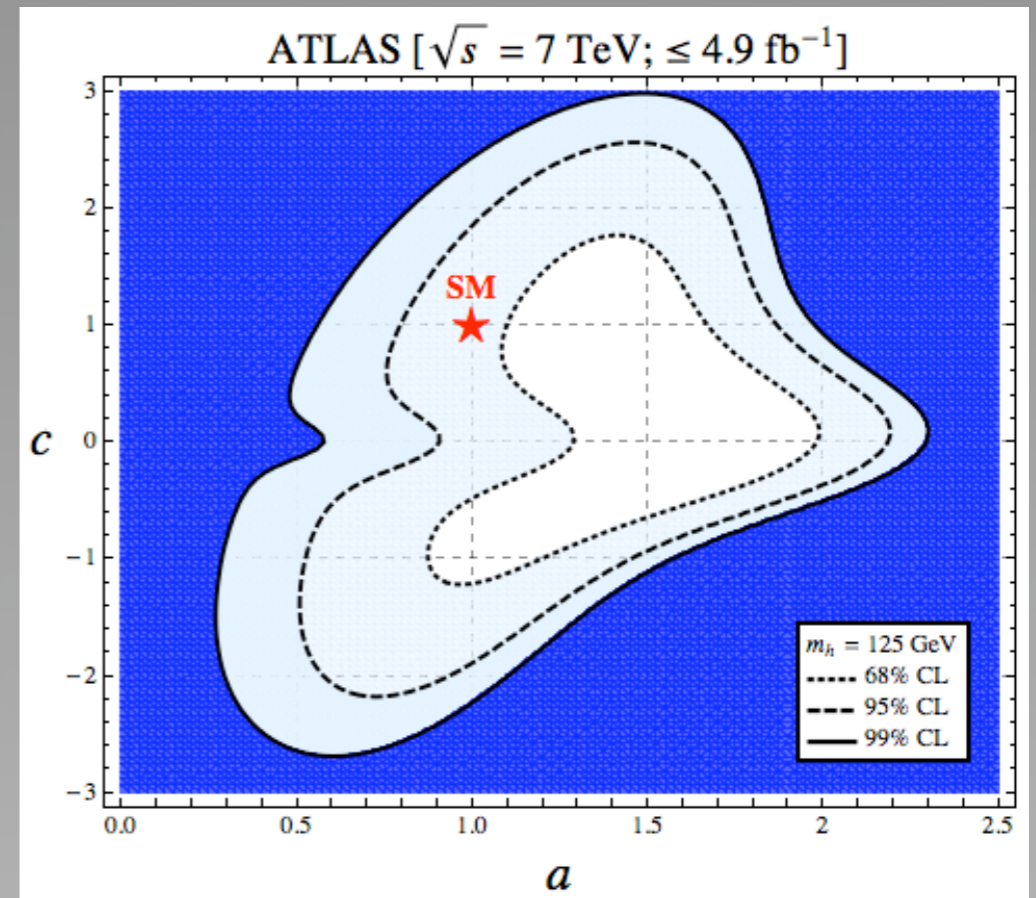
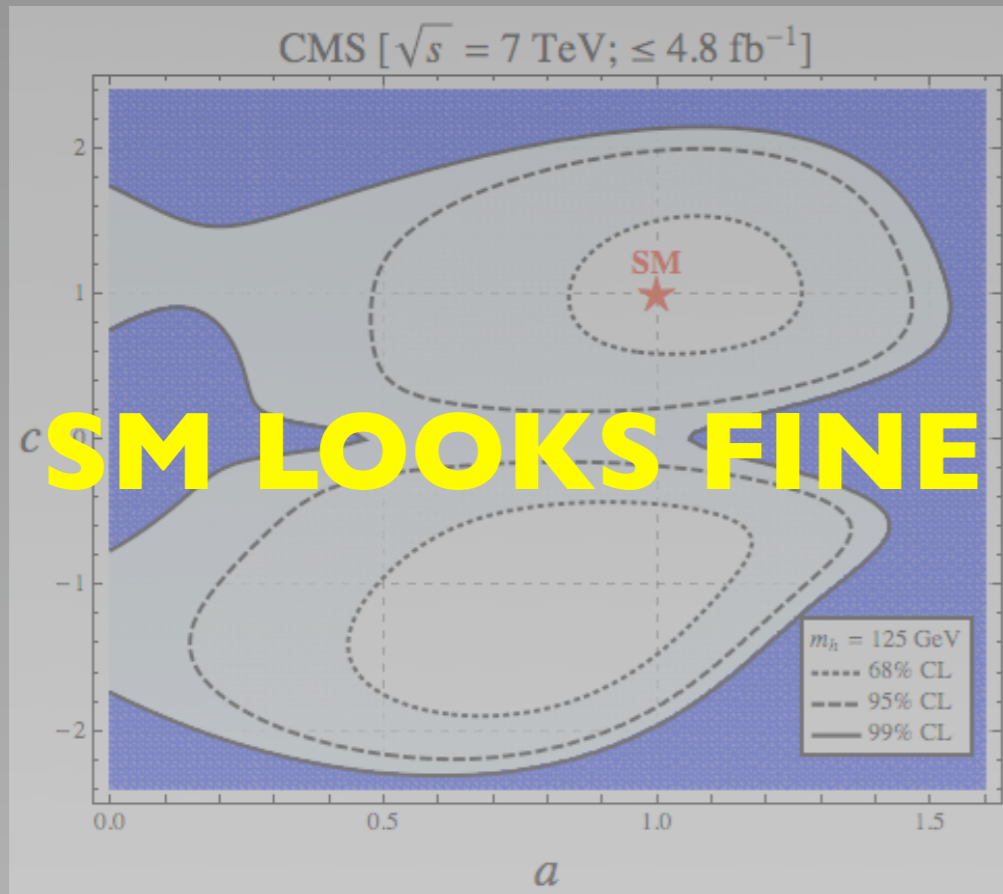
(combined limit given; event numbers for one mass)

Associated Production

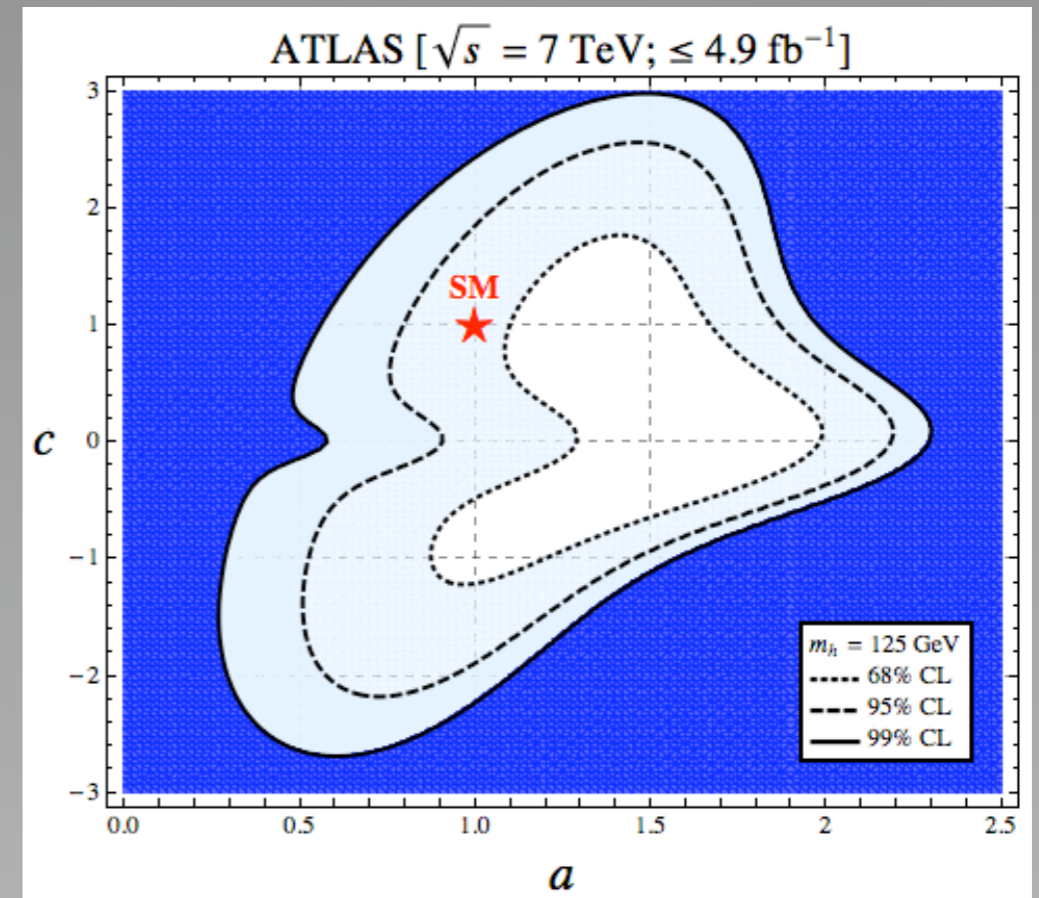
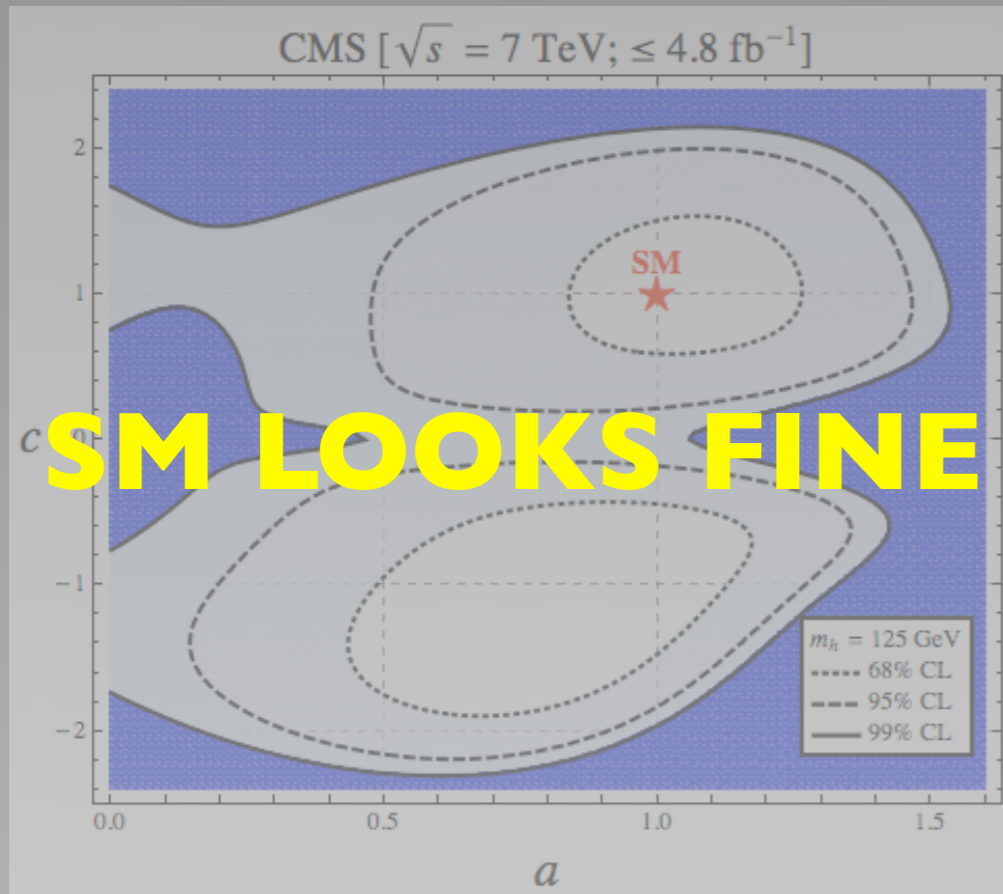
Status report for the Higgs at 125(?)(!)



Status report for the Higgs at 125(?)(!)

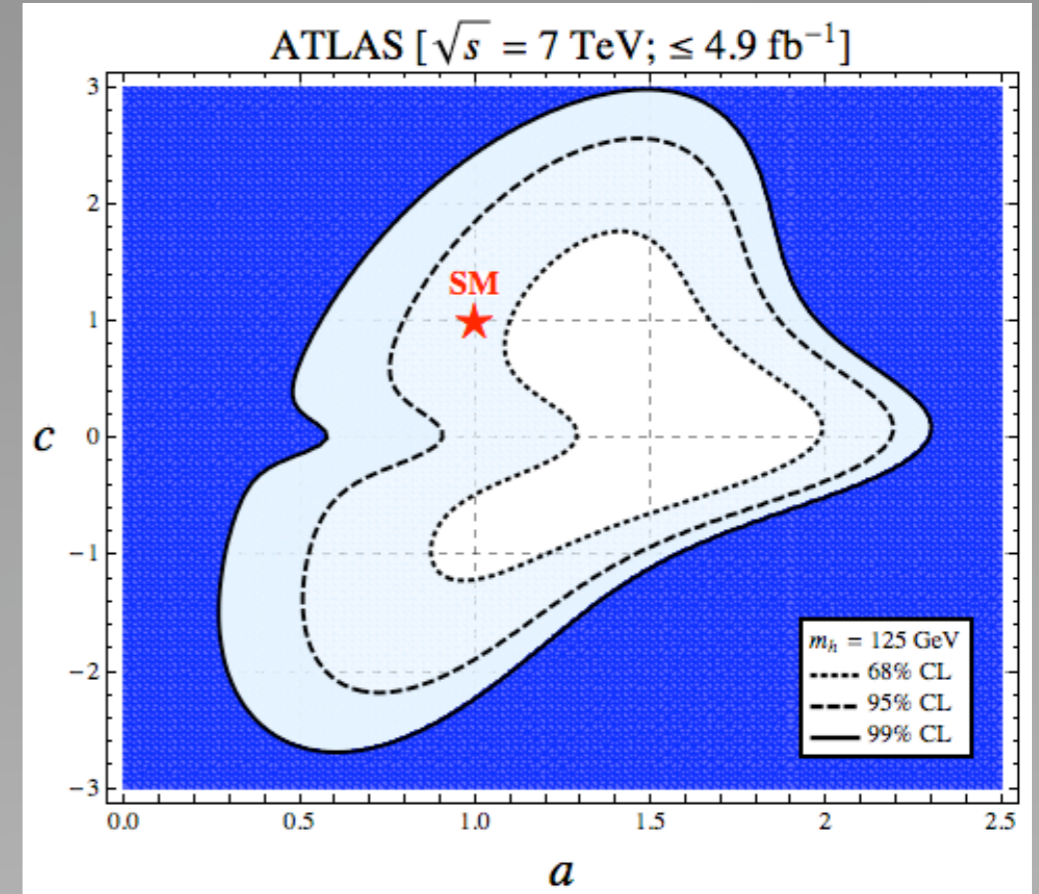
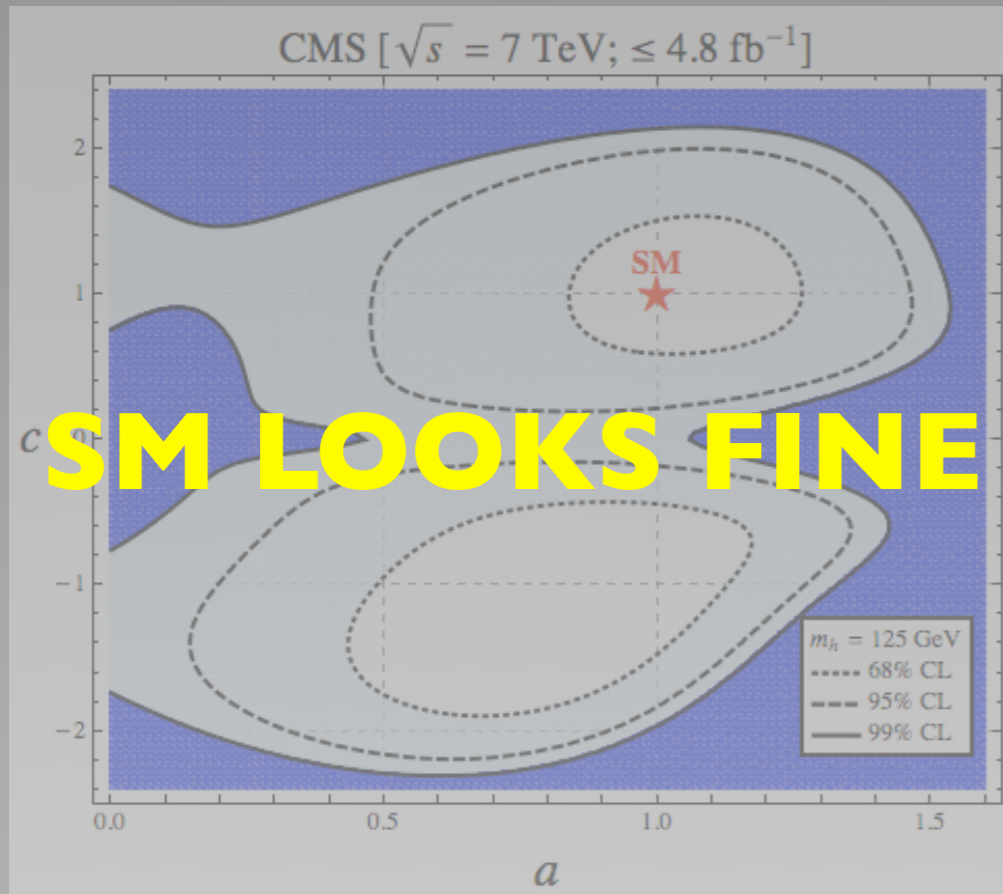


Status report for the Higgs at 125(?)(!)

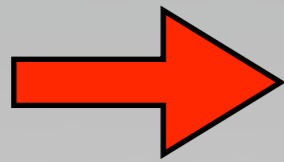


ATLAS seems to disfavor the SM:
how should we take this?

Status report for the Higgs at 125(?)(!)



ATLAS seems to disfavor the SM:
how should we take this?

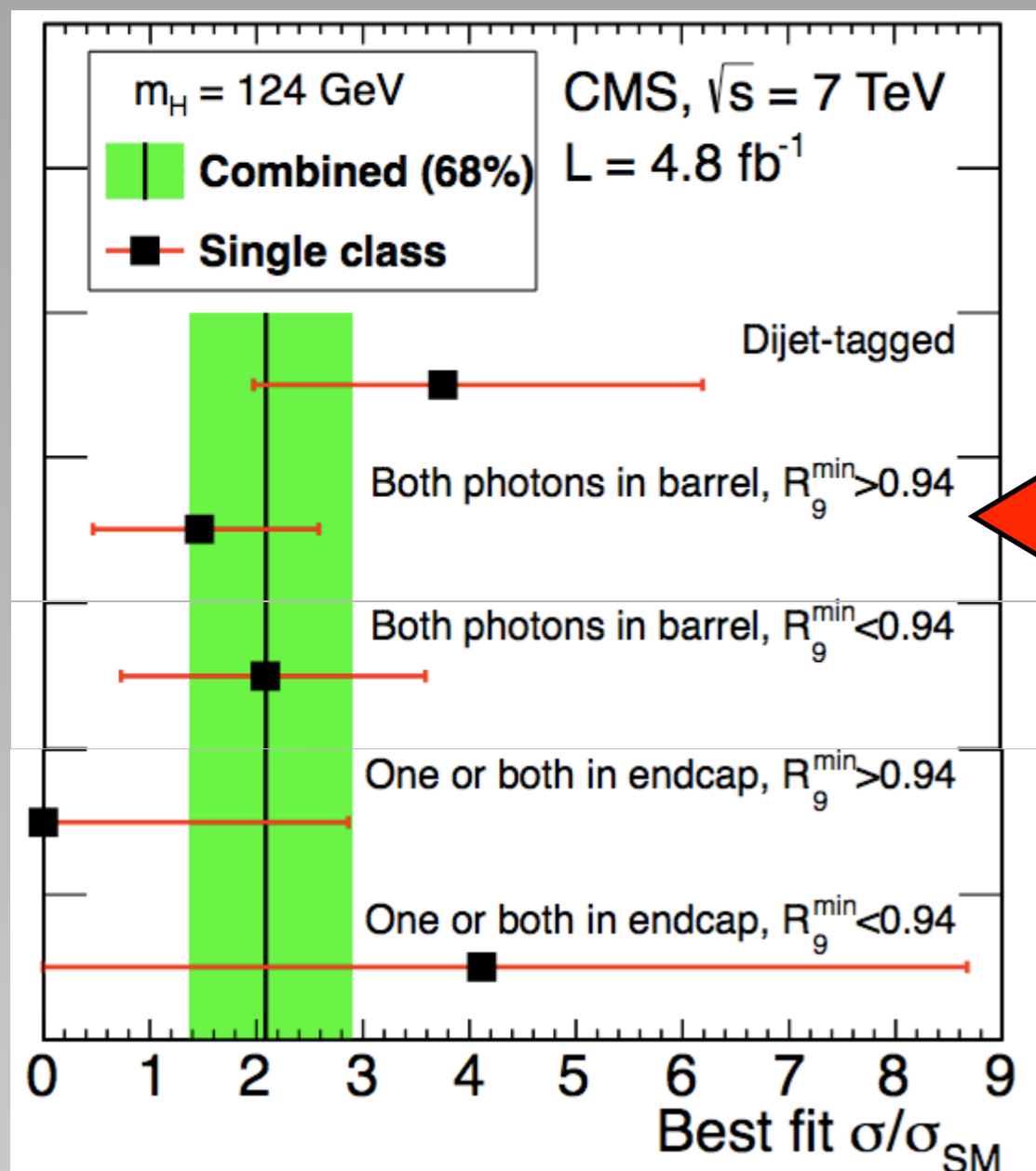


NOT VERY SERIOUSLY
stay tuned...

Final Point: The Need for Exclusive Searching and Reporting

About the displayed CMS results:

- All WW subchannels treated individually
- Others (except bb) treated inclusively
- Can do better for gamma gamma exactly at peak

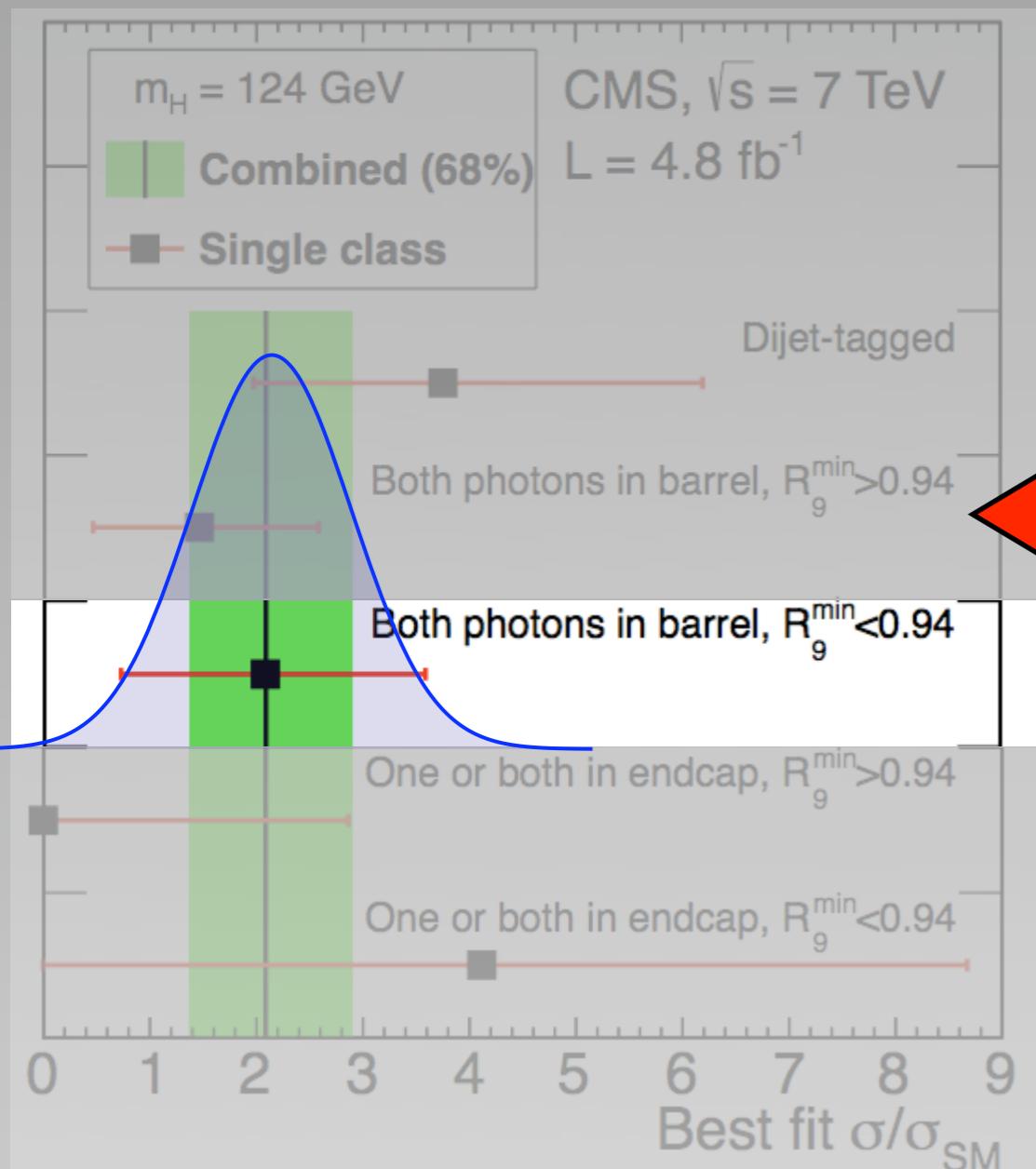


Different method:
Fit each band with
appropriate distribution
(approx. Gaussian)

Final Point: The Need for Exclusive Searching and Reporting

About the displayed CMS results:

- All WW subchannels treated individually
- Others (except bb) treated inclusively
- Can do better for gamma gamma exactly at peak

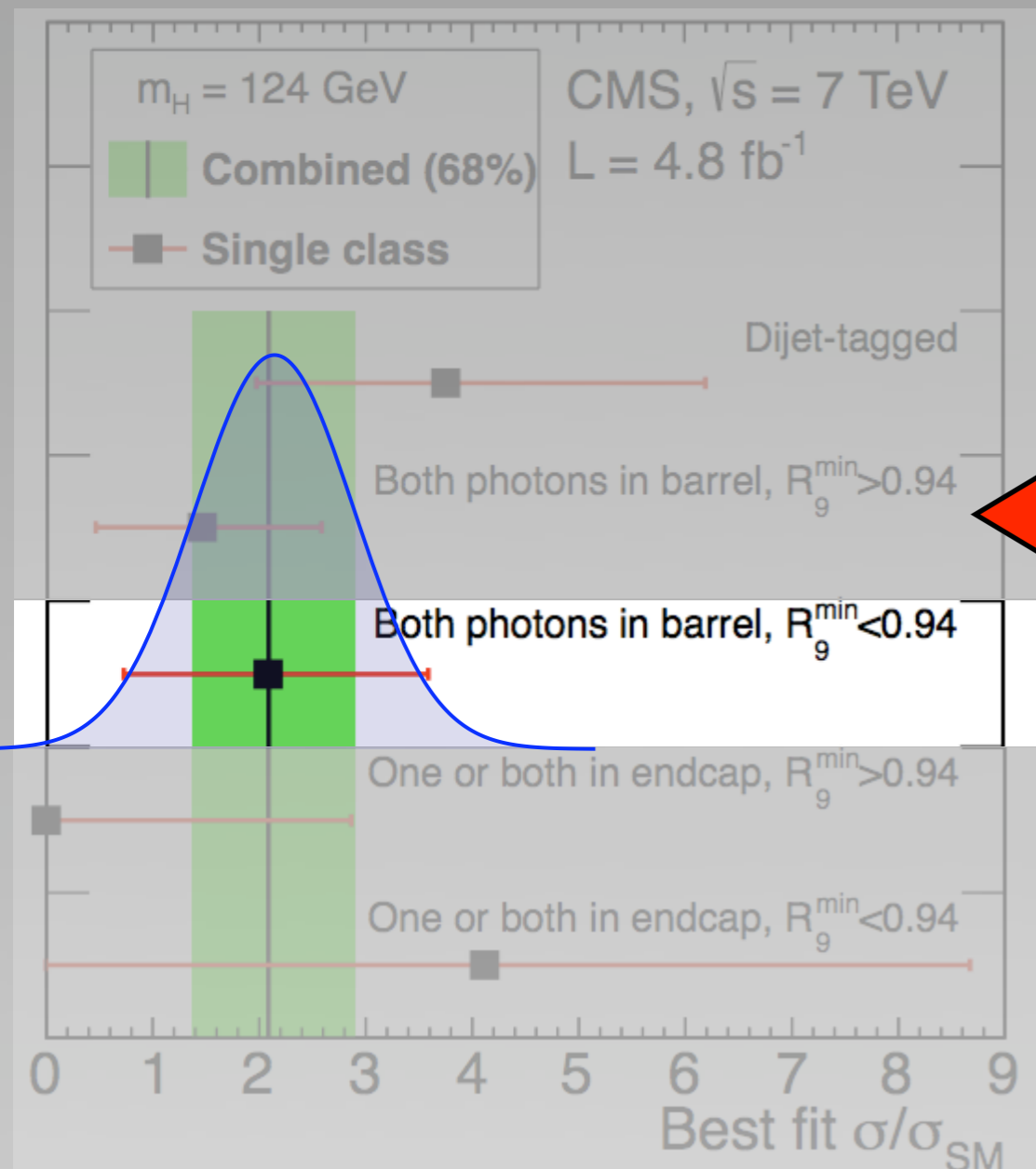


e.g. ...

Final Point: The **Need** for **Exclusive** Searching and Reporting

About the displayed CMS results:

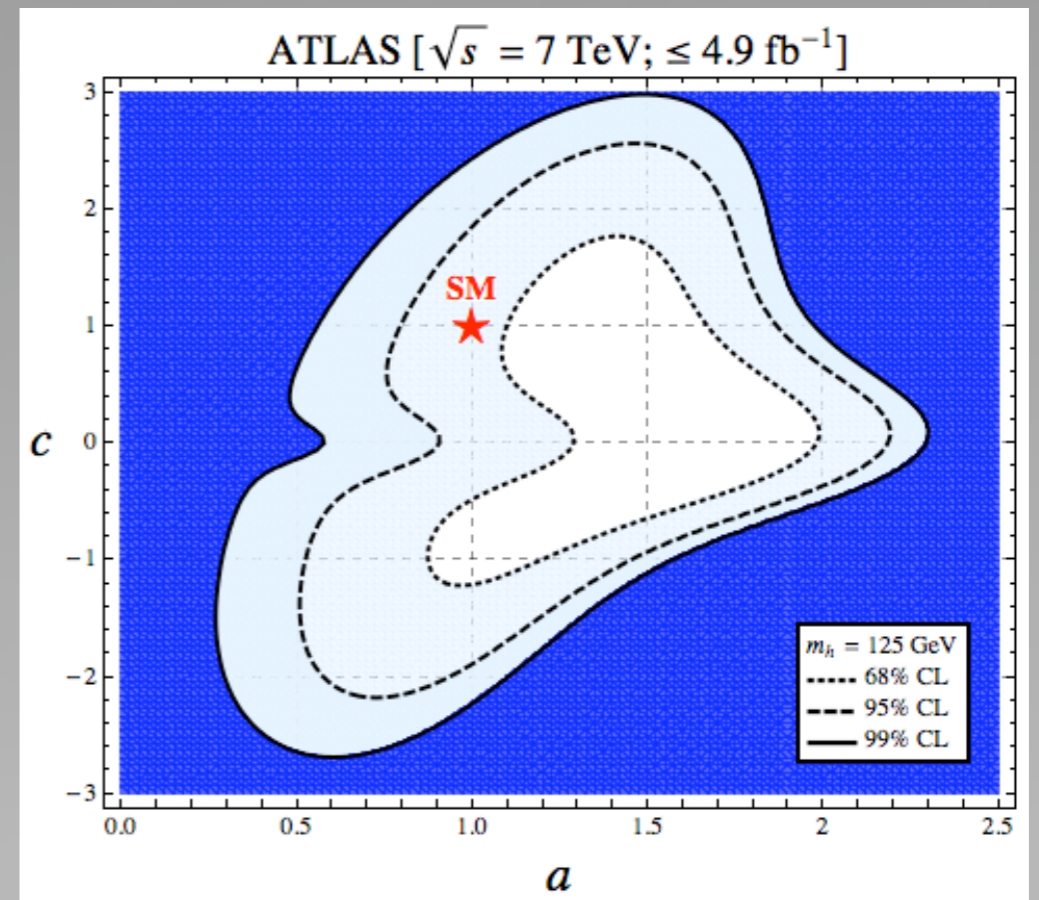
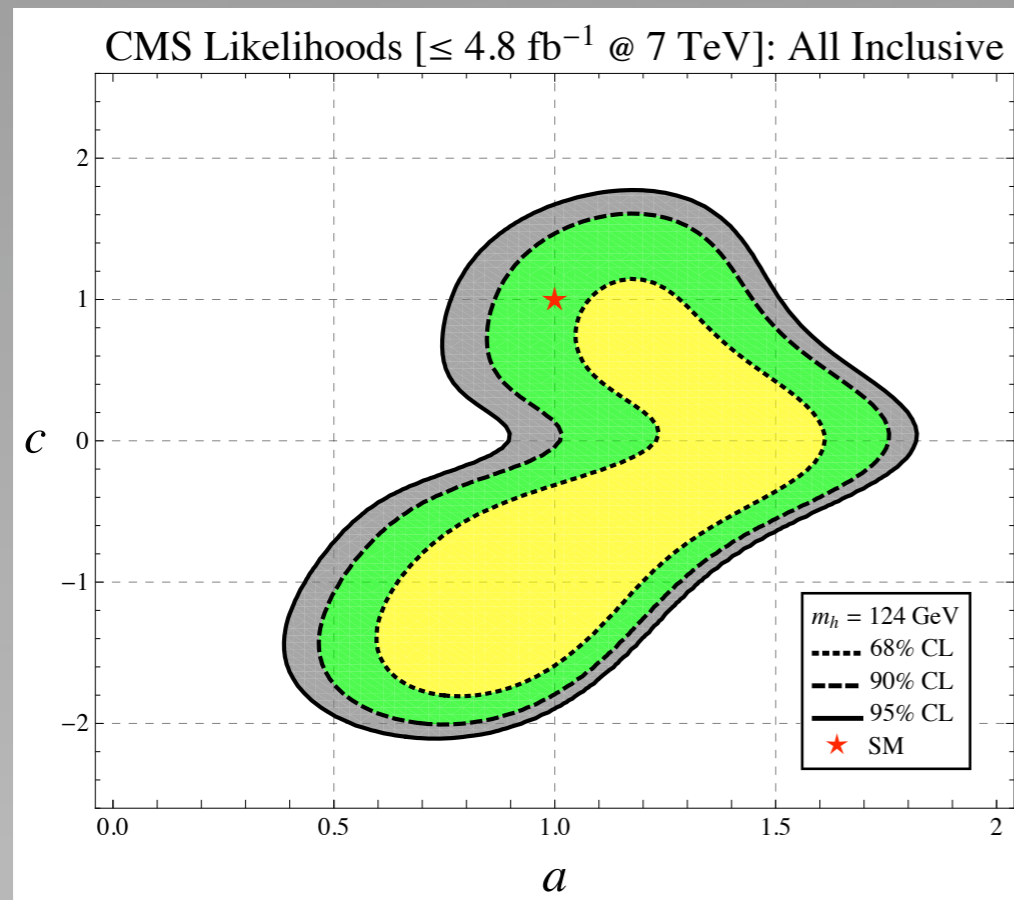
- All WW subchannels treated individually
- Others (except bb) treated inclusively
- Can do better for gamma gamma exactly at peak



Total likelihood given by product of all

Final Point: The Need for Exclusive Searching and Reporting

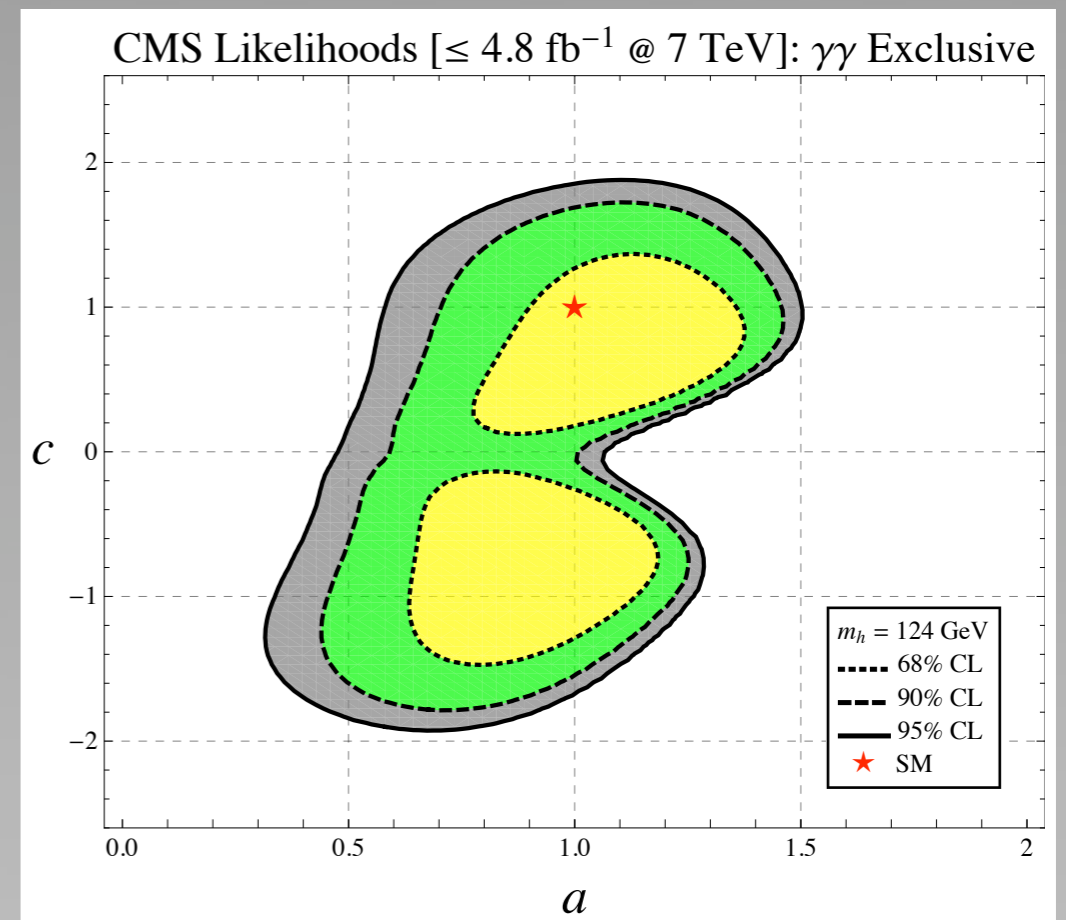
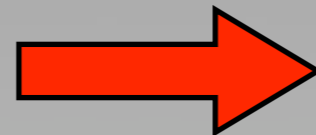
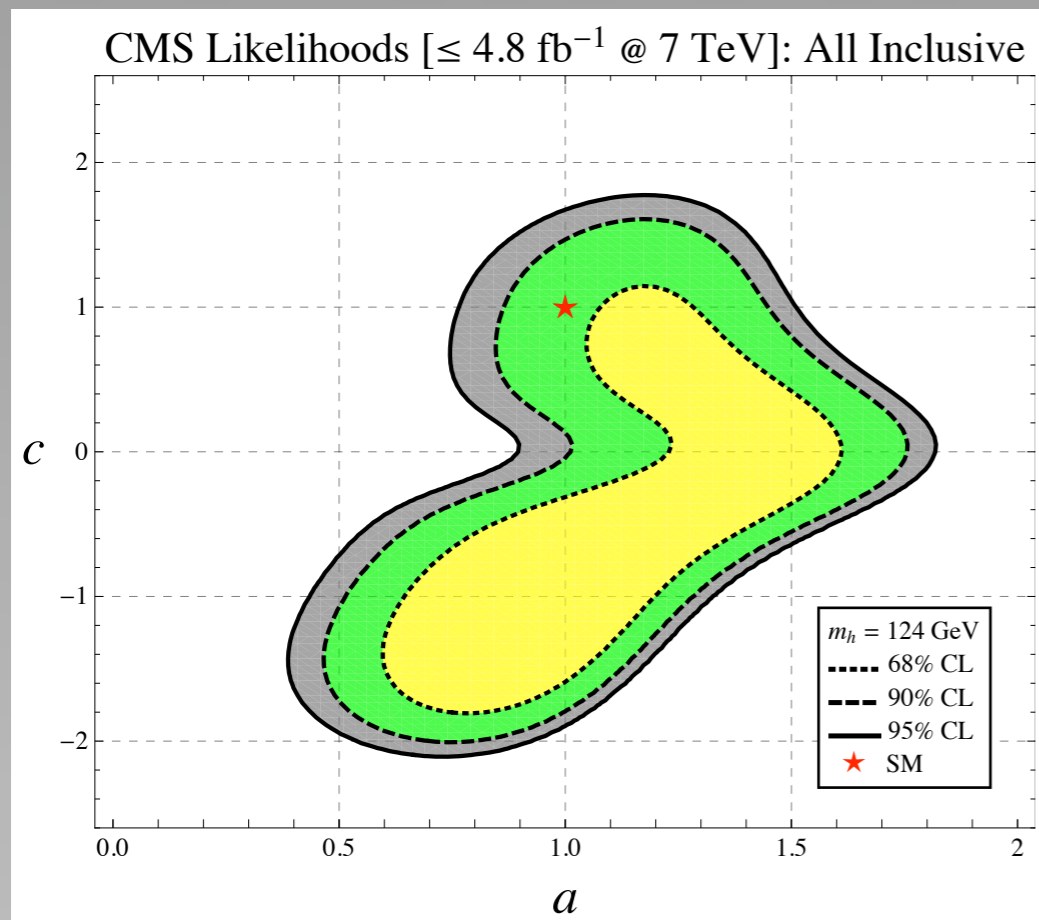
Side-by-side comparison of INCLUSIVE results:



(There **are** real differences, but we see a distinctive -- qualitative -- similarity here)

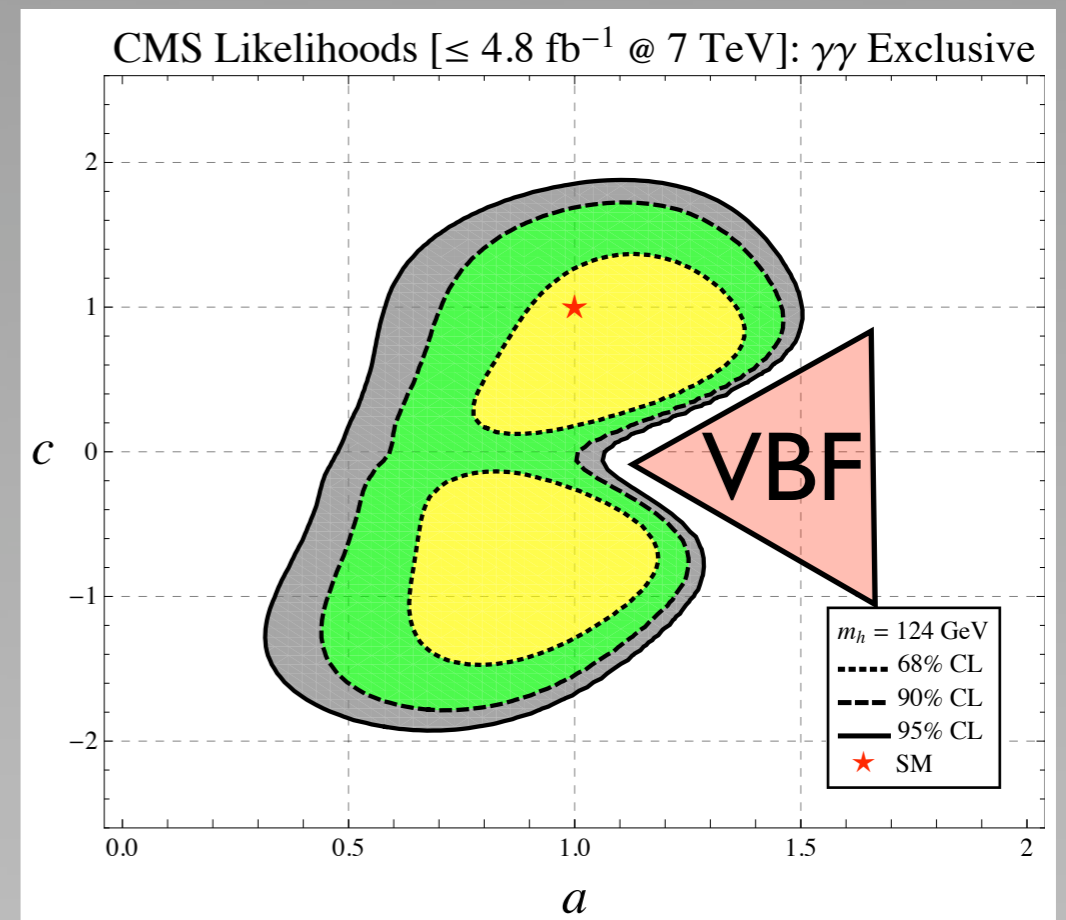
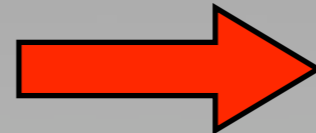
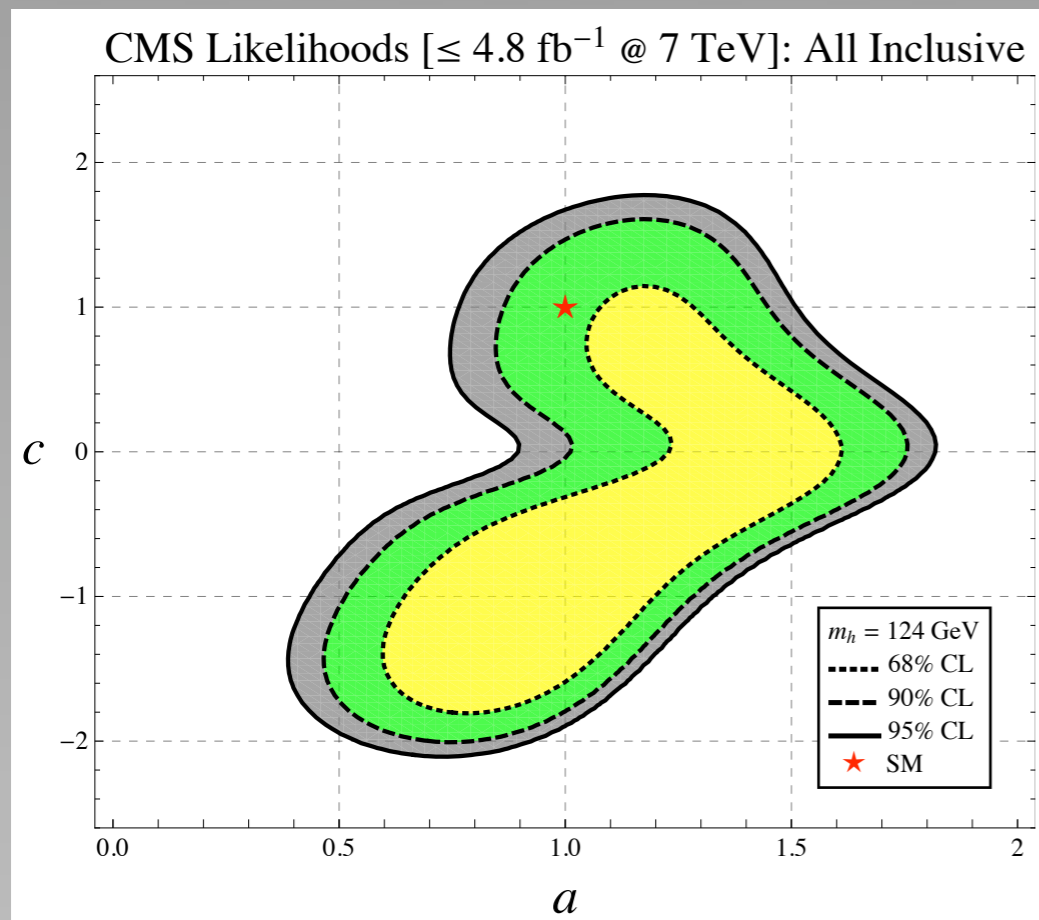
Final Point: The Need for Exclusive Searching and Reporting

Now treat **gamma gamma** subchannels:



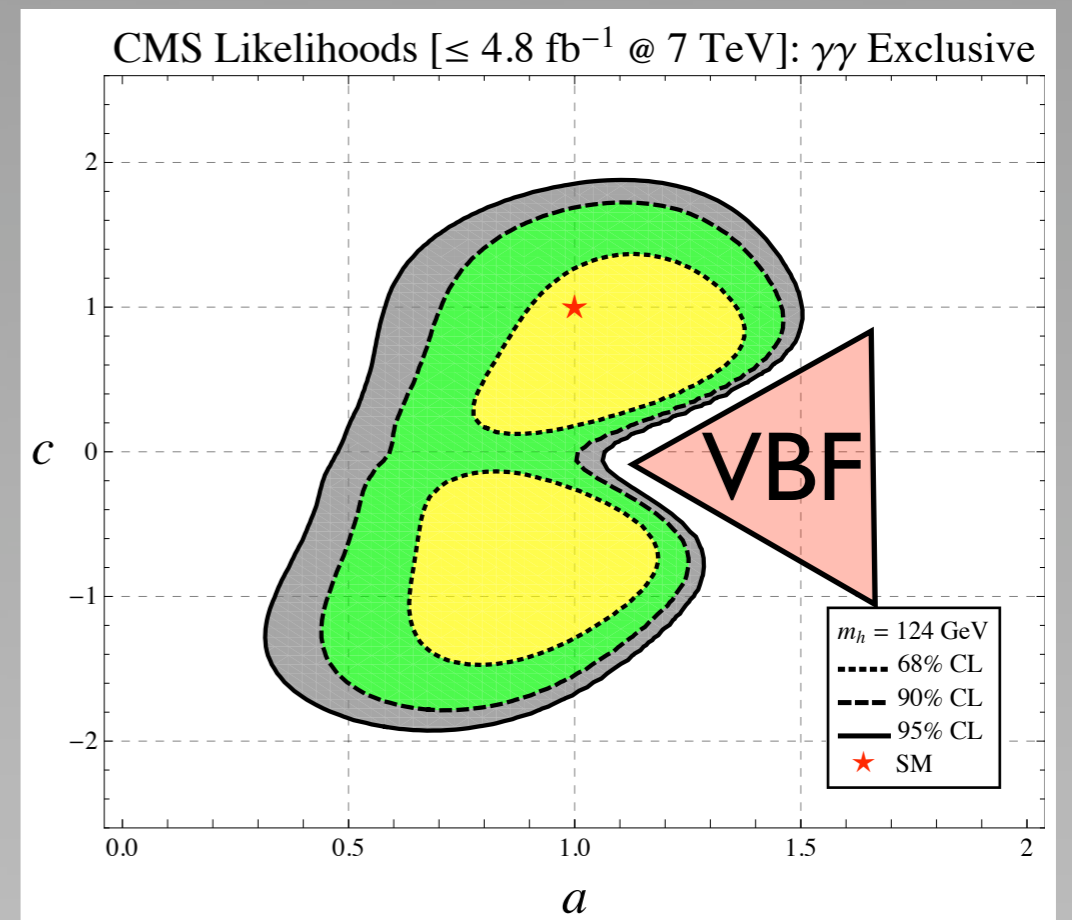
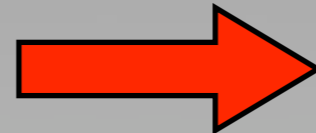
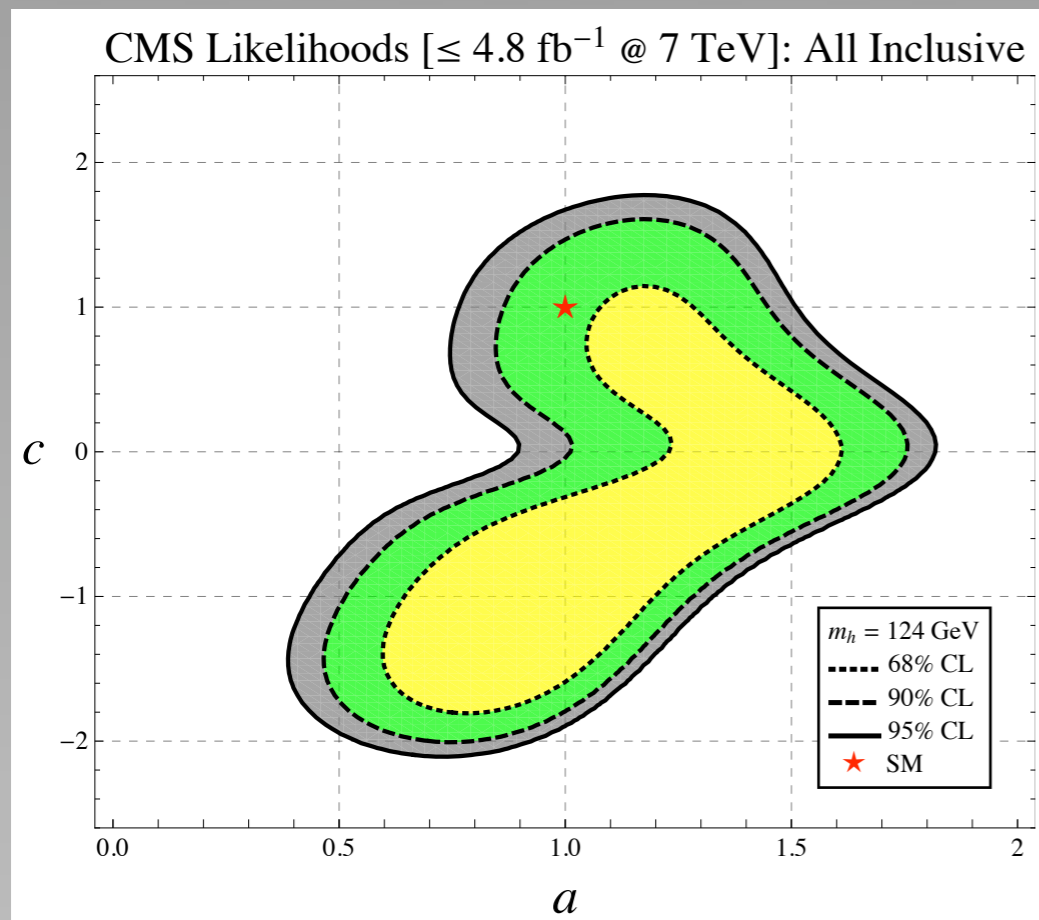
Final Point: The Need for Exclusive Searching and Reporting

Now treat **gamma gamma** subchannels:

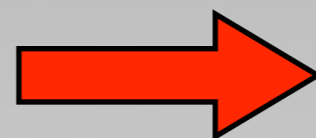


Final Point: The Need for Exclusive Searching and Reporting

Now treat **gamma gamma** subchannels:



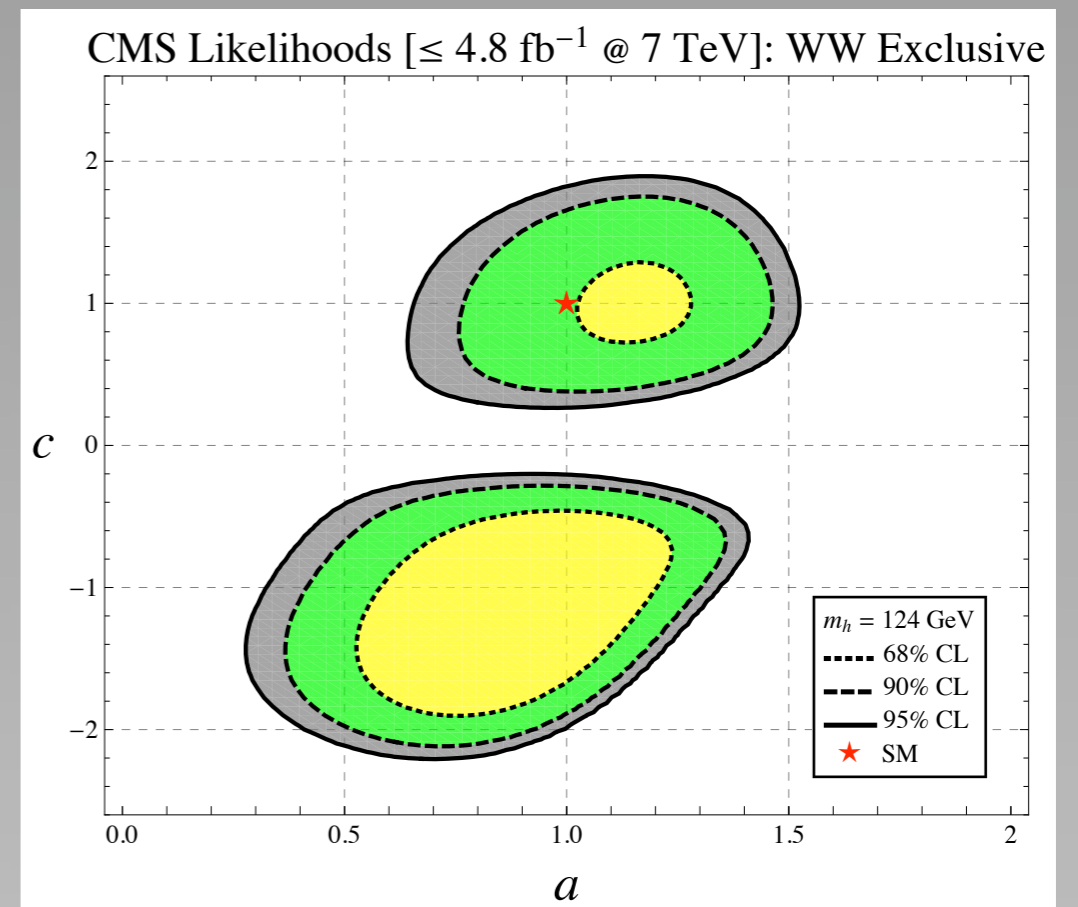
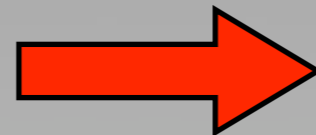
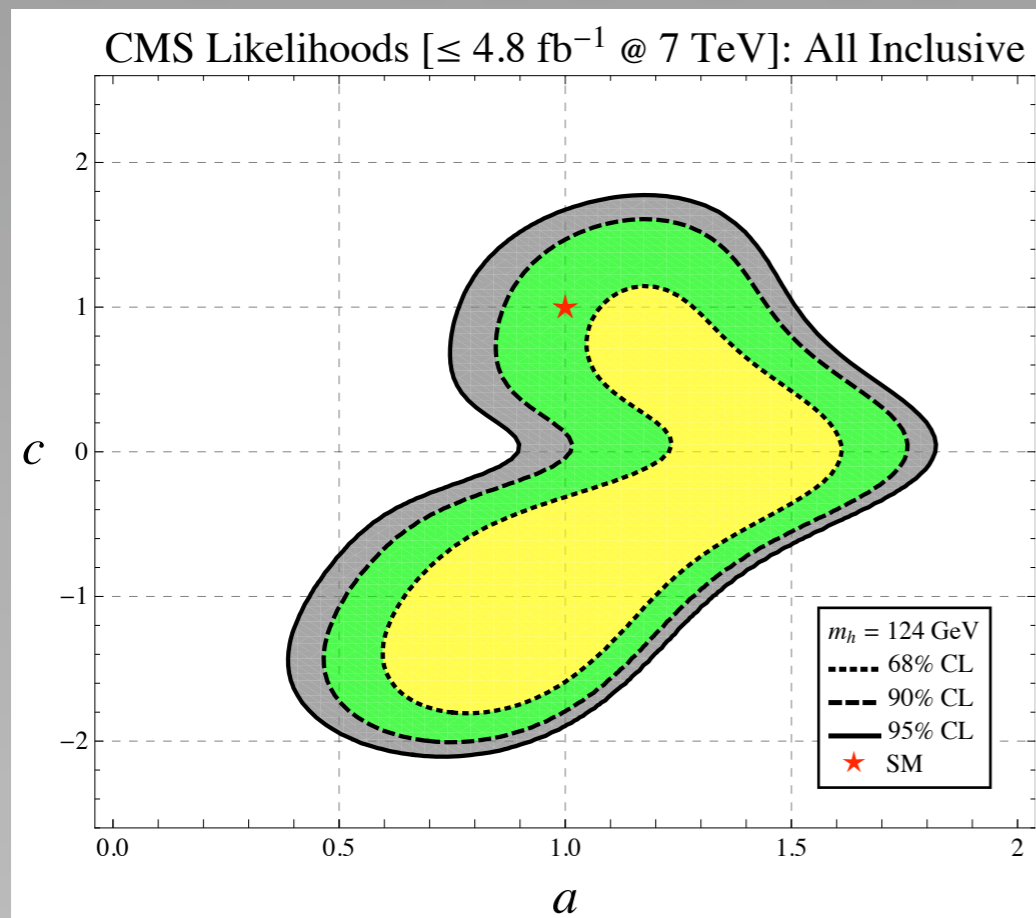
near $c = 0$ line, $R \sim a^2$



Excess in dijet
fit with gauge coupling

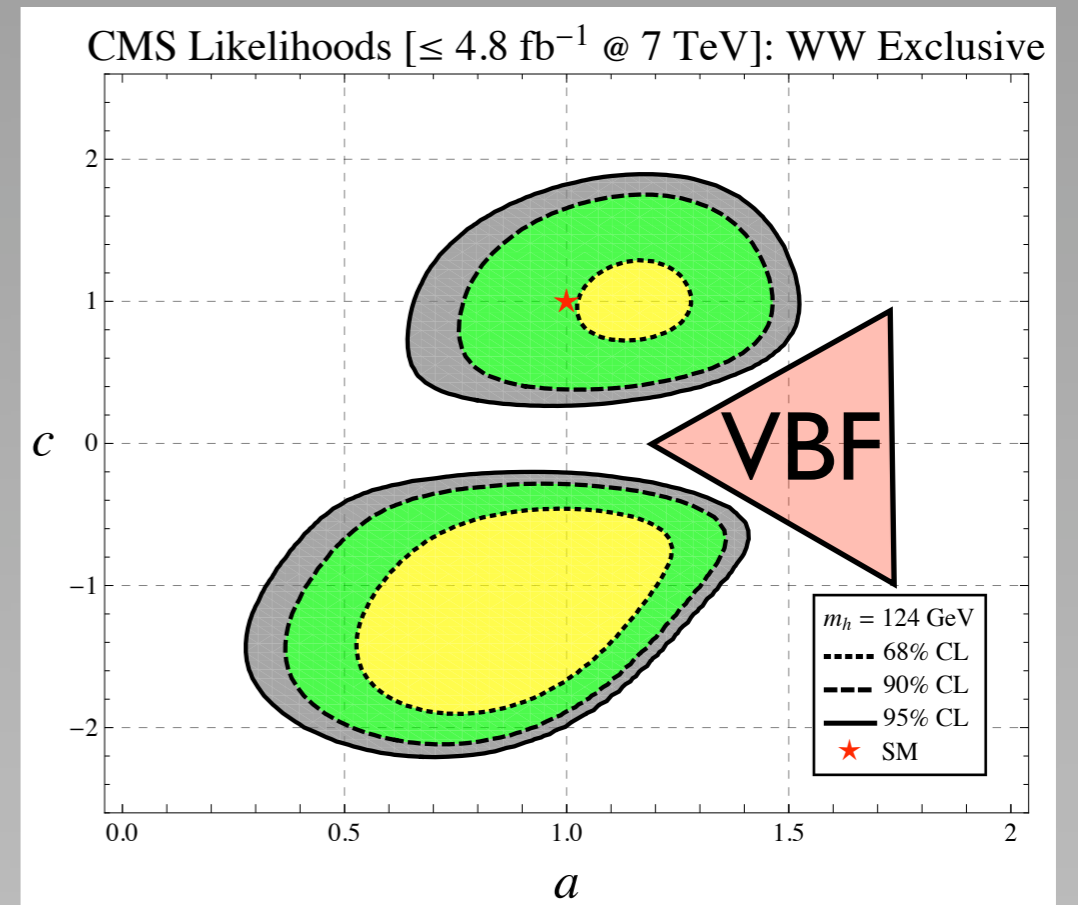
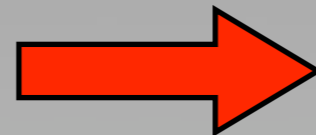
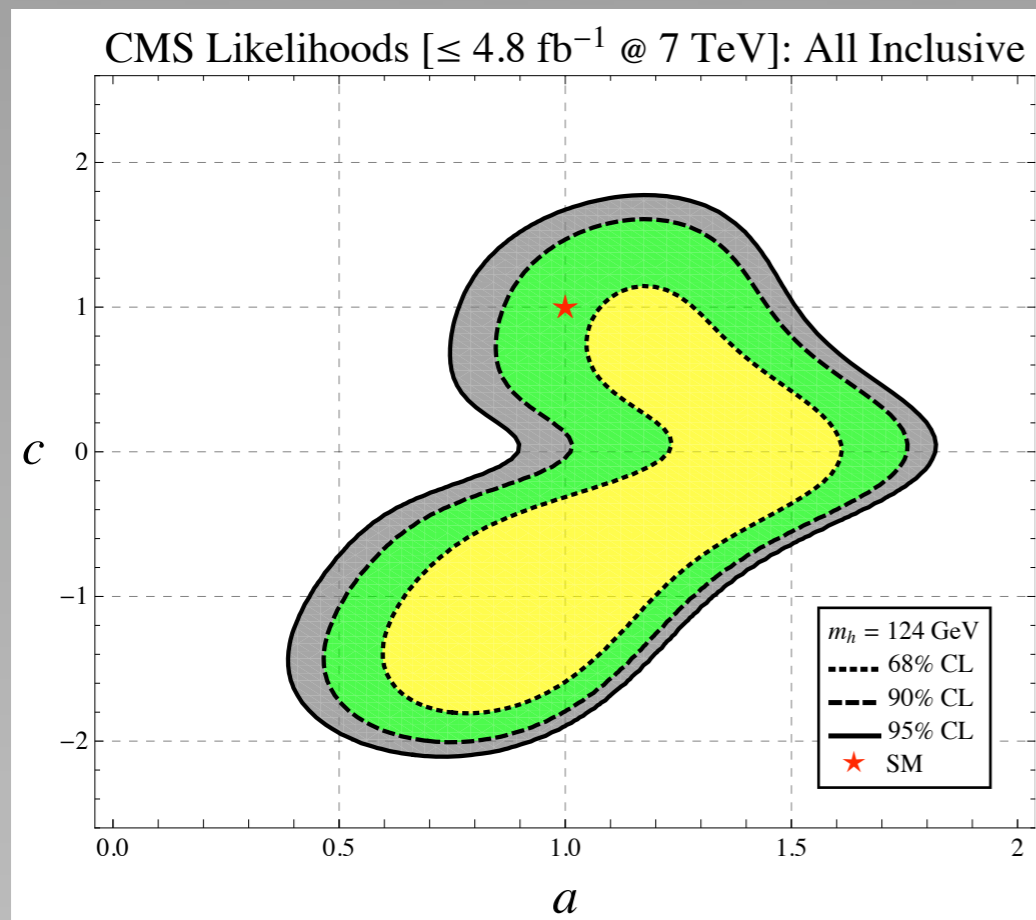
Final Point: The Need for Exclusive Searching and Reporting

WW subchannels:



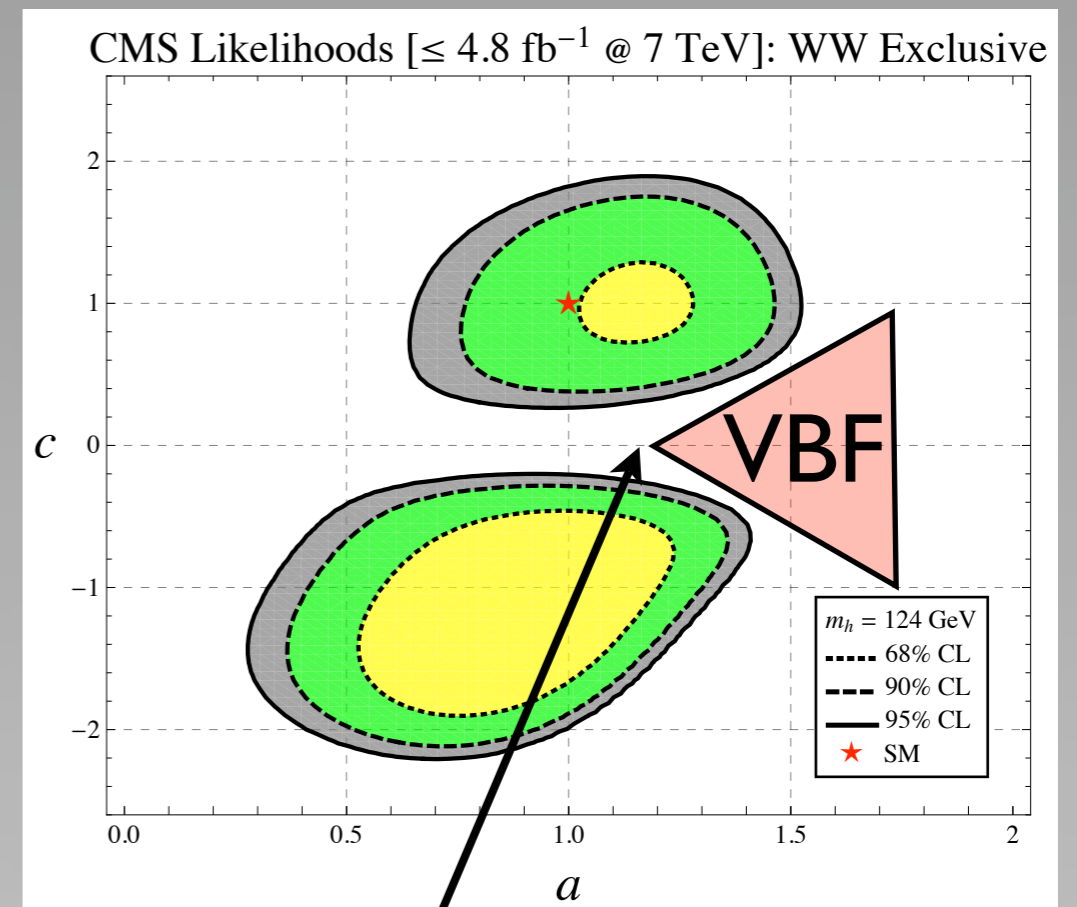
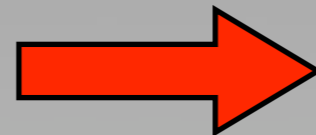
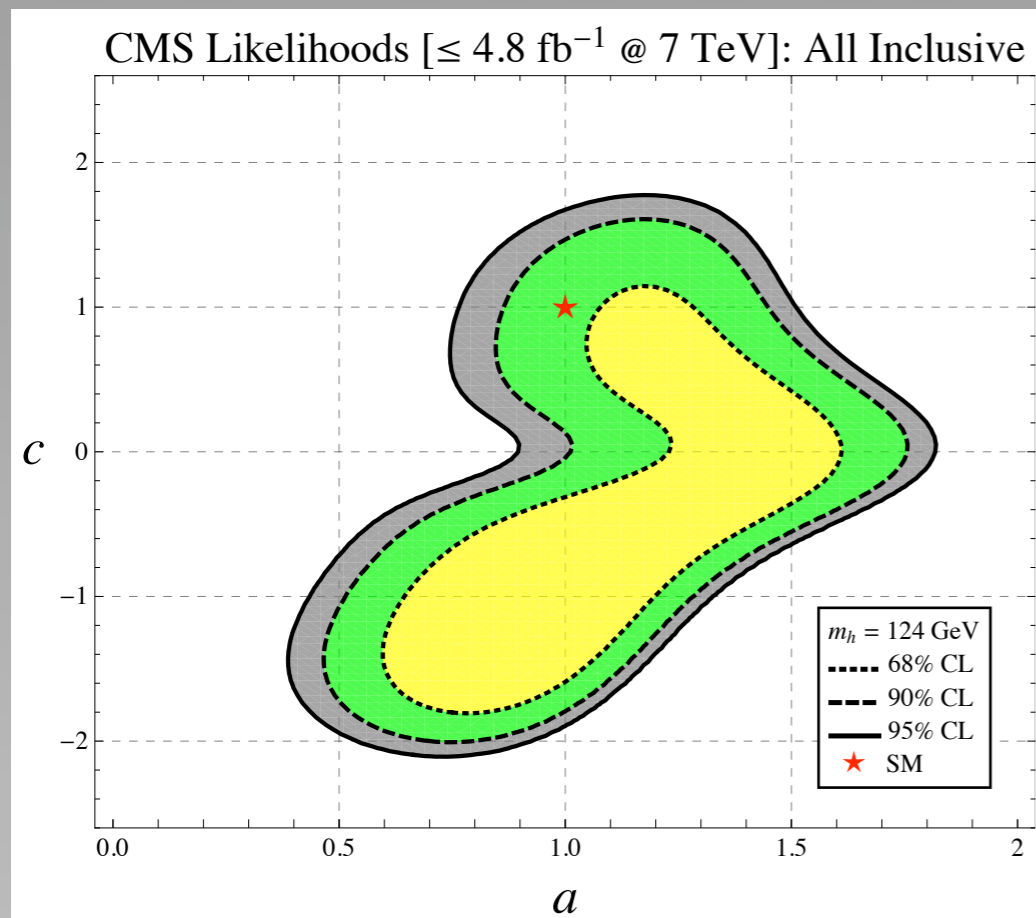
Final Point: The Need for Exclusive Searching and Reporting

WW subchannels:



Final Point: The Need for Exclusive Searching and Reporting

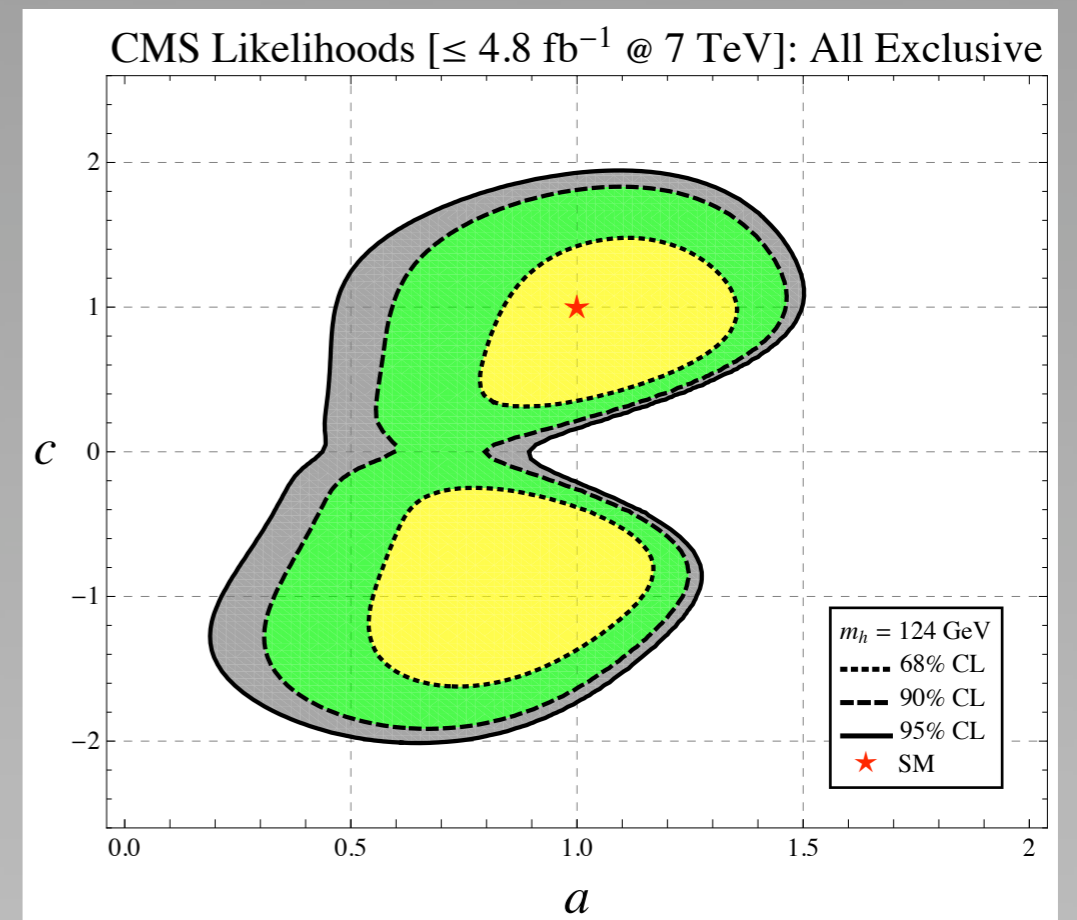
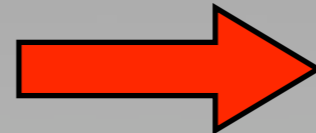
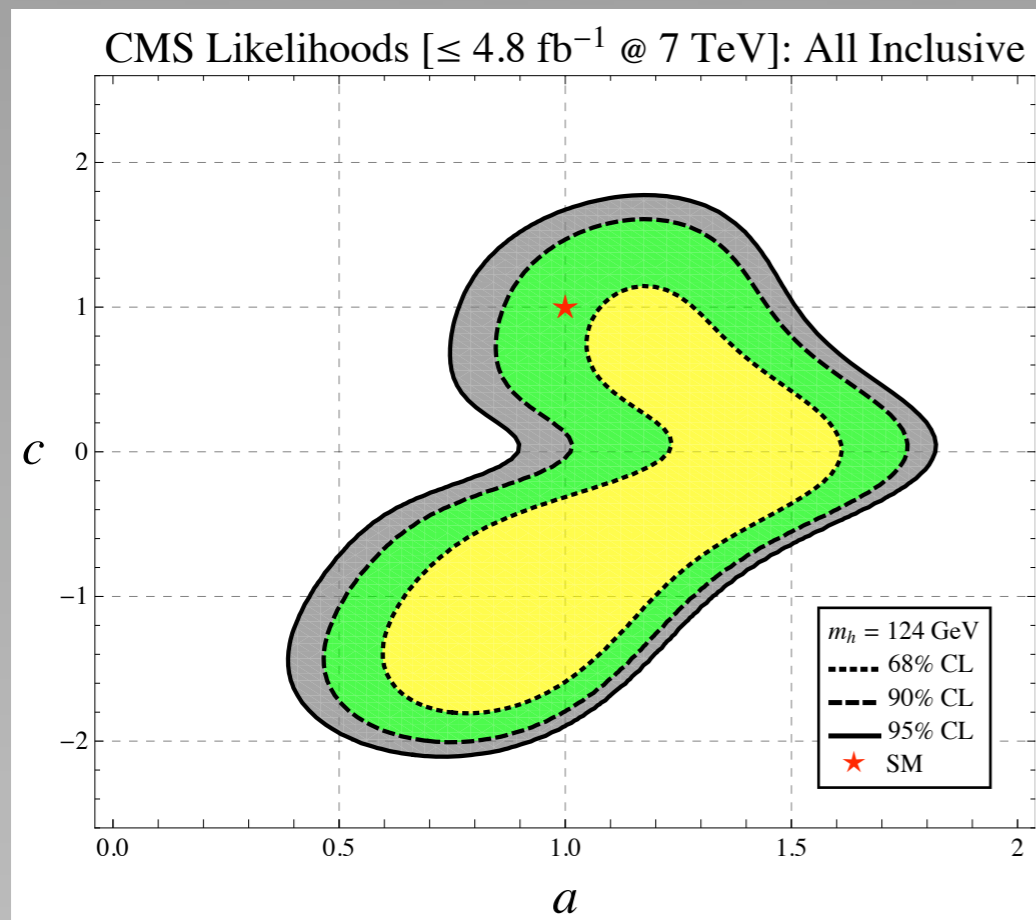
WW subchannels:



Note VBF cuts deeper in this case:
signal deficit in this subchannel
BG ~ 11 , obs. ~ 8

Final Point: The Need for Exclusive Searching and Reporting

ALL subchannels:



Final Point: The **Need** for **Exclusive** Searching and Reporting

*****IMPORTANT POINT NUMBER THREE*****

Tracing back from events to underlying theory **REQUIRES**
separate presentation of limits from each subchannel.

To Conclude

THREE IMPORTANT POINTS

- I. Cut efficiencies truly *needed* for constraining generic spaces
- II. Direct likelihoods would be *nice* to have
- III. Exclusions in generic spaces *need* exclusive searches

To Conclude

THREE IMPORTANT POINTS

- I. Cut efficiencies truly *needed* for constraining generic spaces
- II. Direct likelihoods would be *nice* to have
- III. Exclusions in generic spaces *need* exclusive searches

And some 'incidentals':

- I. SM Higgs looking very good at 125 GeV
- II. Other masses still in play, but very non-SM couplings
- III. Time will tell us more, but we can already *tell ourselves* a LOT:
Well-tested techniques in place to explore
the parameter space of your choice...

Choose a space, any space...

Choose a space, any space...

Which has Nature chosen?

We're well on our way
to an answer...