

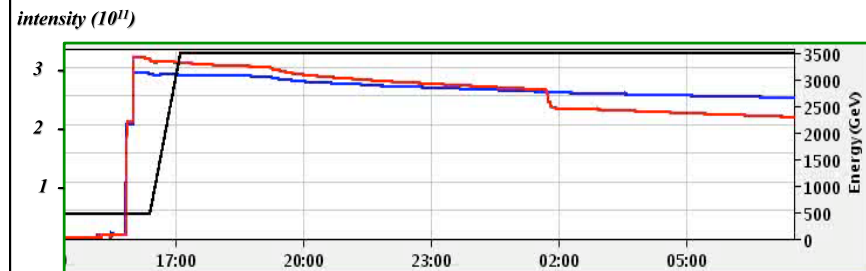
Introduction to Transverse Beam Dynamics

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IP5 *The Ideal World*
I.) Magnetic Fields and Particle Trajectories

Luminosity Run of a typical storage ring:

LHC Storage Ring: Protons accelerated and stored for 12 hours
 distance of particles travelling at about $v \approx c$
 $L = 10^{10} - 10^{11}$ km
 ... several times Sun - Pluto and back



- guide the particles on a well defined orbit („design orbit“)
- focus the particles to keep each single particle trajectory within the vacuum chamber of the storage ring, i.e. close to the design orbit.

I.) Introduction and Basic Ideas

„ ... in the end and after all it should be a kind of circular machine“
 → need transverse deflecting force

Lorentz force $\vec{F} = q * (\vec{E} + \vec{v} \times \vec{B})$

typical velocity in high energy machines: $v \approx c \approx 3 * 10^8 \text{ m/s}$

Example:

$$B = 1 \text{ T} \rightarrow F = q * 3 * 10^8 \frac{\text{m}}{\text{s}} * 1 \frac{\text{Vs}}{\text{m}^2}$$

$$F = q * 300 \frac{\text{MV}}{\text{m}}$$

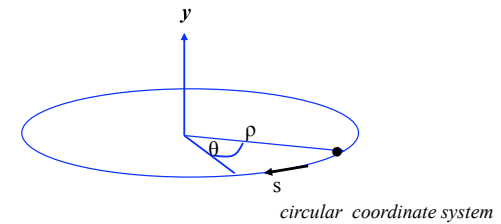
equivalent el. field ... E

technical limit for el. field

$$E \leq 1 \frac{\text{MV}}{\text{m}}$$

old greek dictum of wisdom:
 if you are clever, you use magnetic fields in an accelerator wherever it is possible.

The ideal circular orbit



condition for circular orbit:

Lorentz force $F_L = e v B$

centrifugal force $F_{centr} = \frac{\gamma m_0 v^2}{\rho}$

$$\frac{\gamma m_0 v^2}{\rho} = e v B$$

$$\frac{p}{e} = B \rho$$

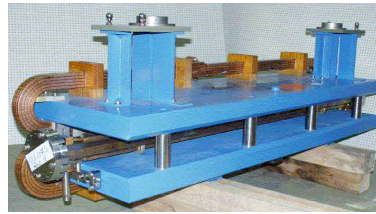
$B \rho =$ "beam rigidity"

1.) The Magnetic Guide Field

Dipole Magnets:

define the ideal orbit
homogeneous field created
by two flat pole shoes

$$B = \frac{\mu_0 n I}{h}$$



Normalise magnetic field to momentum:

$$\frac{p}{e} = B \rho \quad \longrightarrow \quad \frac{1}{\rho} = \frac{e B}{p}$$

convenient units:

$$B = [T] = \left[\frac{Vs}{m^2} \right] \quad p = \left[\frac{GeV}{c} \right]$$

Example LHC:

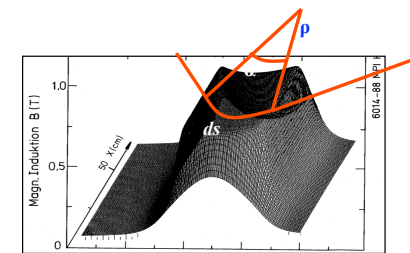
$$B = 8.3 T$$

$$p = 7000 \frac{GeV}{c}$$

$$\frac{1}{\rho} = e \frac{8.3 \frac{Vs}{m^2}}{7000 * 10^9 \frac{eV}{c}} = \frac{8.3 s * 3 * 10^8 \frac{m}{s}}{7000 * 10^9 m^2}$$

$$\frac{1}{\rho} = 0.333 \frac{8.3}{7000} 1/m$$

The Magnetic Guide Field



field map of a storage ring dipole magnet

$$\rho = 2.53 \text{ km} \quad \longrightarrow \quad 2\pi\rho = 17.6 \text{ km} \approx 66\%$$

$$B \approx 1 \dots 8 \text{ T}$$

rule of thumb: $\frac{1}{\rho} \approx 0.3 \frac{B [T]}{p [GeV/c]}$

„normalised bending strength“

2.) Quadrupole Magnets:

required: **focusing forces** to keep trajectories in vicinity of the ideal orbit

linear increasing Lorentz force

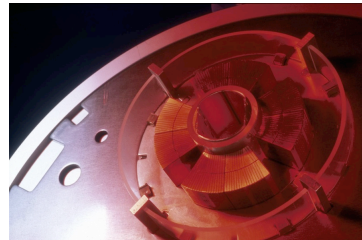
linear increasing magnetic field

$$B_y = g x \quad B_x = g y$$

normalised quadrupole field:

gradient of a quadrupole magnet: $g = \frac{2\mu_0 n I}{r^2}$

→ $k = \frac{g}{p/e}$



LHC main quadrupole magnet

$$g \approx 25 \dots 220 \text{ T/m}$$

simple rule:

$$k = 0.3 \frac{g(\text{T/m})}{p(\text{GeV}/c)}$$

what about the vertical plane:
... Maxwell

$$\vec{\nabla} \times \vec{B} = \cancel{\vec{J}} + \cancel{\frac{\partial \vec{E}}{\partial t}} = 0$$

$$\Rightarrow \frac{\partial B_y}{\partial x} = \frac{\partial B_x}{\partial y}$$

3.) The equation of motion:

Linear approximation:

* ideal particle → design orbit

* any other particle → coordinates x, y *small quantities*
 $x, y \ll \rho$

→ magnetic guide field: only linear terms in x & y of B
have to be taken into account

Taylor Expansion of the B field:

$$B_y(x) = B_{y0} + \frac{dB_y}{dx} x + \frac{1}{2!} \frac{d^2 B_y}{dx^2} x^2 + \frac{1}{3!} \frac{d^3 B_y}{dx^3} x^3 + \dots$$

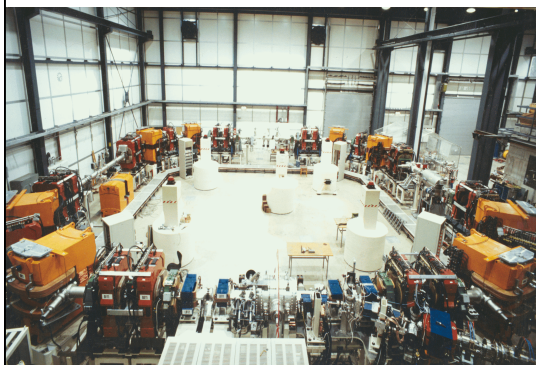
| normalise to momentum
 $p/e = B\rho$

$$\frac{B(x)}{p/e} = \frac{B_0}{B_0 \rho} + \frac{g^* x}{p/e} + \frac{1}{2!} \frac{eg'}{p/e} + \frac{1}{3!} \frac{eg''}{p/e} + \dots$$

The Equation of Motion:

$$\frac{B(x)}{p/e} = \frac{1}{\rho} + kx + \frac{1}{2!} m x^2 + \frac{1}{3!} n x^3 + \dots$$

only terms linear in x, y taken into account *dipole fields*
quadrupole fields



Separate Function Machines:

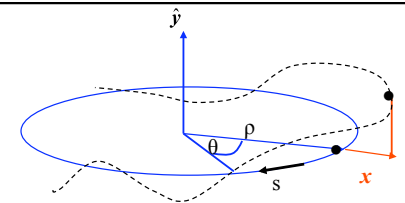
Split the magnets and optimise them according to their job:

bending, focusing etc

Example:
heavy ion storage ring TSR

* man sieht nur
dipole und quads → linear

Equation of Motion:



Consider local segment of a particle trajectory
... and remember the old days:
(Goldstein page 27)

radial acceleration:

$$a_r = \frac{d^2 \rho}{dt^2} - \rho \left(\frac{d\theta}{dt} \right)^2$$

Ideal orbit: $\rho = \text{const.}, \frac{d\rho}{dt} = 0$

Force: $F = m\rho \left(\frac{d\theta}{dt} \right)^2 = m\rho\omega^2$

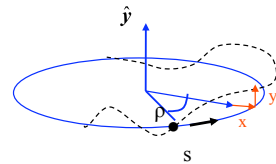
$$F = mv^2 / \rho$$

general trajectory: $\rho \rightarrow \rho + x$

$$F = m \frac{d^2}{dt^2} (x + \rho) - \frac{mv^2}{x + \rho} = e B_y v$$

$$F = m \frac{d^2}{dt^2}(x + \rho) - \frac{mv^2}{x + \rho} = e B_y v$$

① ②



① $\frac{d^2}{dt^2}(x + \rho) = \frac{d^2}{dt^2}x$... as $\rho = \text{const}$

② remember: $x \approx \text{mm}$, $\rho \approx \text{m}$... \rightarrow develop for small x

$$\frac{1}{x + \rho} \approx \frac{1}{\rho} \left(1 - \frac{x}{\rho}\right)$$

Taylor Expansion

$$f(x) = f(x_0) + \frac{(x - x_0)}{1!} f'(x_0) + \frac{(x - x_0)^2}{2!} f''(x_0) + \dots$$

$$m \frac{d^2 x}{dt^2} - \frac{mv^2}{\rho} \left(1 - \frac{x}{\rho}\right) = e B_y v$$

guide field in linear approx.

$$B_y = B_0 + x \frac{\partial B_y}{\partial x} \qquad m \frac{d^2 x}{dt^2} - \frac{mv^2}{\rho} \left(1 - \frac{x}{\rho}\right) = ev \left\{ B_0 + x \frac{\partial B_y}{\partial x} \right\} \quad | : m$$

$$\frac{d^2 x}{dt^2} - \frac{v^2}{\rho} \left(1 - \frac{x}{\rho}\right) = \frac{ev B_0}{m} + \frac{ev x g}{m}$$

independent variable: $t \rightarrow s$

$$\frac{dx}{dt} = \frac{dx}{ds} \frac{ds}{dt}$$

$$\frac{d^2 x}{dt^2} = \frac{d}{dt} \left(\frac{dx}{ds} \frac{ds}{dt} \right) = \frac{d}{ds} \left(\frac{dx}{ds} \frac{ds}{dt} \right) \frac{ds}{dt}$$

$$\frac{d^2 x}{dt^2} = x'' v^2 + \cancel{\frac{dx}{ds} \frac{dv}{ds}} v$$

$$x'' v^2 - \frac{v^2}{\rho} \left(1 - \frac{x}{\rho}\right) = \frac{ev B_0}{m} + \frac{ev x g}{m} \quad | : v^2$$

$$x'' - \frac{1}{\rho} \left(1 - \frac{x}{\rho}\right) = \frac{e B_0}{mv} + \frac{e x g}{mv}$$

$$mv = p$$

$$x'' - \frac{1}{\rho} + \frac{x}{\rho^2} = \frac{B_0}{p/e} + \frac{x g}{p/e}$$

normalize to momentum of particle

$$x'' - \cancel{\frac{1}{\rho}} + \frac{x}{\rho^2} = -\cancel{\frac{1}{\rho}} + k x$$

$$\frac{B_0}{p/e} = -\frac{1}{\rho}$$

$$\frac{g}{p/e} = k$$

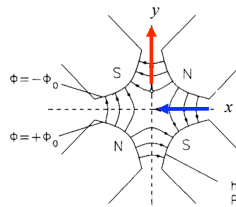
$$x'' + x \left(\frac{1}{\rho^2} - k\right) = 0$$

* Equation for the vertical motion:

$$\frac{1}{\rho^2} = 0 \quad \text{no dipoles ... in general ...}$$

$$k \leftrightarrow -k \quad \text{quadrupole field changes sign}$$

$$y'' + k y = 0$$



Remarks:

$$* \quad x'' + \left(\frac{1}{\rho^2} - k\right) \cdot x = 0$$

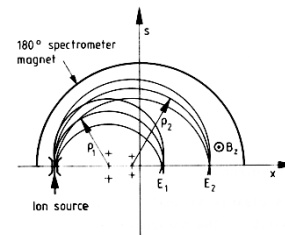
... there seems to be a focusing even without a quadrupole gradient

„weak focusing of dipole magnets“

$$k = 0 \quad \Rightarrow \quad x'' = -\frac{1}{\rho^2} x$$

even without quadrupoles there is a refocusing force (i.e. focusing) in the bending plane of the dipole magnets

... in large machines it is weak. (!)



Mass spectrometer: particles are separated according to their energy and focused due to the $1/\rho$ effect of the dipole

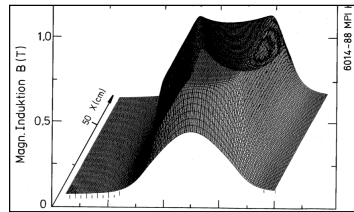
Hard Edge Model:

$$x'' + \left(\frac{1}{\rho^2} - k\right) \cdot x = 0$$

$$x''(s) + \left\{ \frac{1}{\rho^2(s)} - k(s) \right\} \cdot x(s) = 0$$

this equation is not really correct !!!

*bending and focusing forces
– even inside a magnet –
depend on the position “s”*



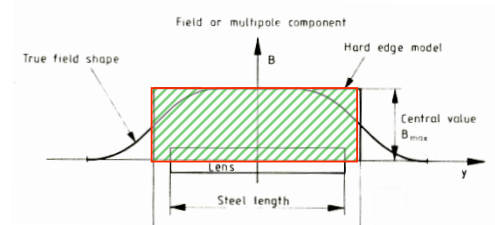
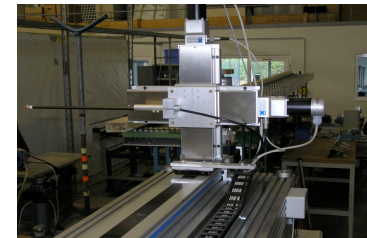
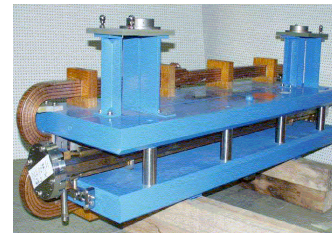
*magnetic field of a storage ring dipole
k = const within a magnet*

*But still: Inside a magnet the focusing
properties are constant !*

$$\frac{1}{\rho} = \text{const}$$

$$k = \text{const}$$

„effective magnetic length“



$$B * l_{eff} = \int_0^{l_{mag}} B ds$$

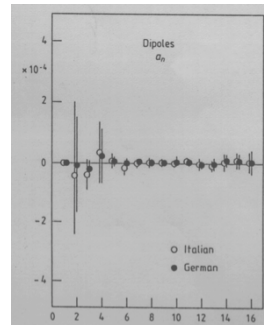
*** **Multipoles**

Taylor Expansion of the B field:

$$B_y(x) = B_{y0} + \frac{dB_y}{dx} x + \frac{1}{2!} \frac{d^2 B_y}{dx^2} x^2 + \frac{1}{3!} \frac{d^3 B_y}{dx^3} x^3 + \dots \quad : B_{y0}$$

divide by the main field
to get the relative error
contribution

→ definition of multipole coefficients.



Multipole contributions to the HERA s.c. dipole field

4.) Solution of Trajectory Equations

$$\left. \begin{array}{l} \text{Define ... hor. plane: } K = 1/\rho^2 - k \\ \text{... vert. Plane: } K = k \end{array} \right\} x'' + K x = 0$$

Differential Equation of harmonic oscillator ... with spring constant K

Ansatz: $x(s) = a_1 \cdot \cos(\omega s) + a_2 \cdot \sin(\omega s)$

general solution: linear combination of two independent solutions

$$x'(s) = -a_1 \omega \sin(\omega s) + a_2 \omega \cos(\omega s)$$

$$x''(s) = -a_1 \omega^2 \cos(\omega s) - a_2 \omega^2 \sin(\omega s) = -\omega^2 x(s) \quad \rightarrow \quad \omega = \sqrt{K}$$

general solution:

$$x(s) = a_1 \cos(\sqrt{K}s) + a_2 \sin(\sqrt{K}s)$$

determine a_1, a_2 by boundary conditions:

$$s = 0 \quad \rightarrow \quad \begin{cases} x(0) = x_0 & , & a_1 = x_0 \\ x'(0) = x'_0 & , & a_2 = \frac{x'_0}{\sqrt{|K|}} \end{cases}$$

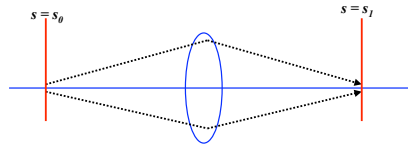
Hor. Focusing Quadrupole $K > 0$:

$$x(s) = x_0 \cdot \cos(\sqrt{|K|}s) + x'_0 \cdot \frac{1}{\sqrt{|K|}} \sin(\sqrt{|K|}s)$$

$$x'(s) = -x_0 \cdot \sqrt{|K|} \cdot \sin(\sqrt{|K|}s) + x'_0 \cdot \cos(\sqrt{|K|}s)$$

For convenience expressed in matrix formalism:

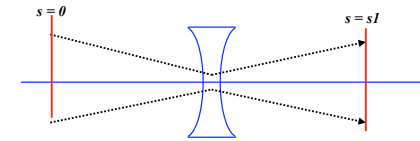
$$\begin{pmatrix} x \\ x' \end{pmatrix}_{s_1} = M_{foc} * \begin{pmatrix} x \\ x' \end{pmatrix}_{s_0}$$



$$M_{foc} = \begin{pmatrix} \cos(\sqrt{|K|}l) & \frac{1}{\sqrt{|K|}} \sin(\sqrt{|K|}l) \\ -\sqrt{|K|} \sin(\sqrt{|K|}l) & \cos(\sqrt{|K|}l) \end{pmatrix}$$

hor. defocusing quadrupole:

$$x'' - K x = 0$$



Remember from school:

$$f(s) = \cosh(s) \quad , \quad f'(s) = \sinh(s)$$

Ansatz: $x(s) = a_1 \cdot \cosh(\omega s) + a_2 \cdot \sinh(\omega s)$

$$M_{defoc} = \begin{pmatrix} \cosh \sqrt{|K|}l & \frac{1}{\sqrt{|K|}} \sinh \sqrt{|K|}l \\ \sqrt{|K|} \sinh \sqrt{|K|}l & \cosh \sqrt{|K|}l \end{pmatrix}$$

drift space:

$$K = 0$$

$$M_{drift} = \begin{pmatrix} 1 & l \\ 0 & 1 \end{pmatrix}$$

! with the assumptions made, the motion in the horizontal and vertical planes are independent „ ... the particle motion in x & y is uncoupled“

Thin Lens Approximation:

matrix of a quadrupole lens

$$M = \begin{pmatrix} \cos\sqrt{|k|}l & \frac{1}{\sqrt{|k|}}\sin\sqrt{|k|}l \\ -\sqrt{|k|}\sin\sqrt{|k|}l & \cos\sqrt{|k|}l \end{pmatrix}$$

in many practical cases we have the situation:

$$f = \frac{1}{kl_q} \gg l_q \quad \dots \text{focal length of the lens is much bigger than the length of the magnet}$$

limes: $l_q \rightarrow 0$ while keeping $kl_q = \text{const}$

$$M_x = \begin{pmatrix} 1 & 0 \\ -\frac{1}{f} & 1 \end{pmatrix}$$

$$M_y = \begin{pmatrix} 1 & 0 \\ \frac{1}{f} & 1 \end{pmatrix}$$

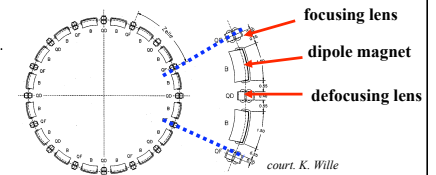
... useful for fast (and in large machines still quite accurate) „back on the envelope calculations“ ... and for the guided studies !

Transformation through a system of lattice elements

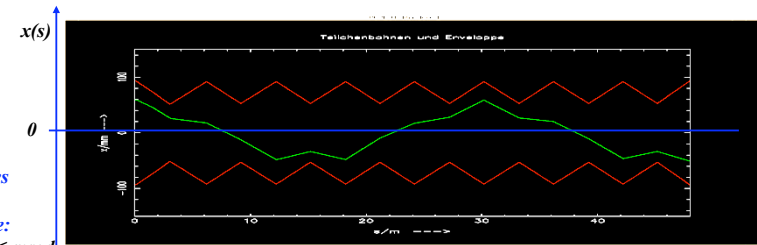
combine the single element solutions by multiplication of the matrices

$$M_{total} = M_{QF} * M_D * M_{QD} * M_{Bend} * M_D * \dots$$

$$\begin{pmatrix} x \\ x' \end{pmatrix}_{s_2} = M(s_2, s_1) * \begin{pmatrix} x \\ x' \end{pmatrix}_{s_1}$$



in each accelerator element the particle trajectory corresponds to the movement of a harmonic oscillator „



typical values in a strong foc. machine:

$$x \approx \text{mm}, x' \leq \text{mrad}$$

5.) Orbit & Tune:

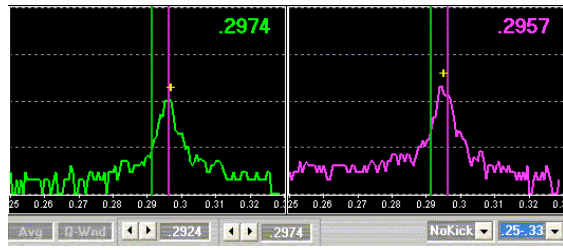
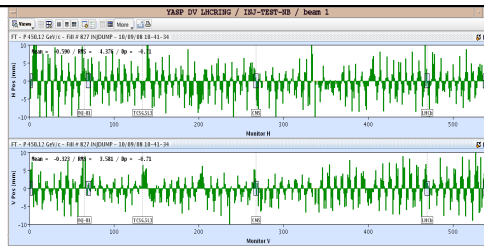
Tune: number of oscillations per turn

64.31
59.32

Relevant for beam stability:
non integer part

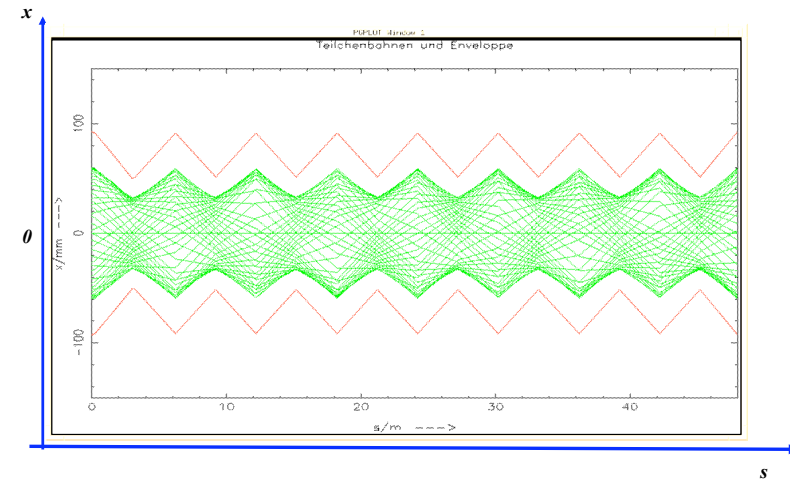
LHC revolution frequency: 11.3 kHz

$$0.31 * 11.3 = 3.5 \text{ kHz}$$



Question: what will happen, if the particle performs a second turn ?

... or a third one or ... 10^{10} turns



Résumé:

beam rigidity: $B \cdot \rho = \frac{p}{q}$

bending strength of a dipole: $\frac{1}{\rho} [m^{-1}] = \frac{0.2998 \cdot B_0(T)}{p(\text{GeV}/c)}$

focusing strength of a quadrupole: $k [m^{-2}] = \frac{0.2998 \cdot g}{p(\text{GeV}/c)}$

focal length of a quadrupole: $f = \frac{1}{k \cdot l_q}$

equation of motion: $x'' + Kx = \frac{1}{\rho} \frac{\Delta p}{p}$

matrix of a foc. quadrupole: $x_{s2} = M \cdot x_{s1}$

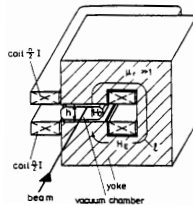
$$M = \begin{pmatrix} \cos\sqrt{|K|}l & \frac{1}{\sqrt{|K|}}\sin\sqrt{|K|}l \\ -\sqrt{|K|}\sin\sqrt{|K|}l & \cos\sqrt{|K|}l \end{pmatrix}, \quad M = \begin{pmatrix} 1 & 0 \\ \frac{1}{f} & 1 \end{pmatrix}$$

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- 2.) **Klaus Wille:** *Physics of Particle Accelerators and Synchrotron Radiation Facilities*, Teubner, Stuttgart 1992
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- 5.) **M.S. Livingston, J.P. Blewett:** *Particle Accelerators*, Mc Graw-Hill, New York, 1962
- 6.) *The CERN Accelerator School (CAS) Proceedings*
- 7.) **Frank Hinterberger:** *Physik der Teilchenbeschleuniger*, Springer Verlag 1997
- 8.) **Mathew Sands:** *The Physics of e+ e- Storage Rings*, SLAC report 121, 1970
- 9.) **D. Edwards, M. Syphers:** *An Introduction to the Physics of Particle Accelerators*, SSC Lab 1990

Dipole Magnets:

homogeneous field created by two flat pole shoes



Field Calculation:

3rd Maxwell equation for a static field: $\nabla \times \vec{H} = \vec{j}$

according to Stokes theorem:

$$\int_S (\nabla \times \vec{H}) \cdot \vec{n} \, da = \oint \vec{H} \cdot d\vec{l} = \int_S \vec{j} \cdot \vec{n} \, da = N \cdot I$$

$$\oint \vec{H} \cdot d\vec{l} = H_0 \cdot h + H_{Fe} \cdot l_{Fe}$$

in matter we get with $\mu_r \approx 1000$

$$\oint \vec{H} \cdot d\vec{l} = H_0 \cdot h + \frac{H_0}{\mu_r} \cdot l_{Fe} \approx H_0 \cdot h$$

Magnetic field of a dipole magnet:

$$H_0 = \frac{B_0}{\mu_0} \rightarrow B_0 = \frac{\mu_0 N I}{h}$$

$h = \text{gap height}$

Quadrupole Magnets:

Calculation of the Quadrupole Field:

$$\oint H ds = N \cdot I$$

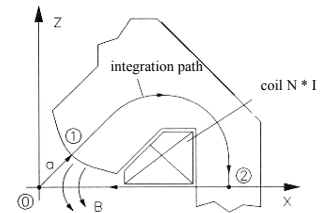
$$\oint H ds = \int_0^1 H_0 ds + \int_1^2 H_{Fe} ds + \int_2^0 H ds = N \cdot I$$

$H_{Fe} = H_0 / \mu_{Fe}$ $H \perp ds$
 $\mu_{Fe} \approx 1000$

now we know that $H = B / \mu_0$
and we require $B(r) = -g \cdot r$

$$\int_0^1 H_0 ds = \int_0^a \frac{B_0}{\mu_0} dr = \int_0^a \frac{g \cdot r}{\mu_0} dr = N \cdot I$$

gradient of a quadrupole field: $g = \frac{2\mu_0 \cdot N \cdot I}{r^2}$



cour. K. Wille