Measurement of longitudinal Parameters

The longitudinal dynamics is described by the longitudinal emittance as given by:

- Spread of the bunches ***l*** in time, length ***or*** rf-phase.
- Momentum spread ***δ = Δp/p***, or energy spread ***ΔW/W***

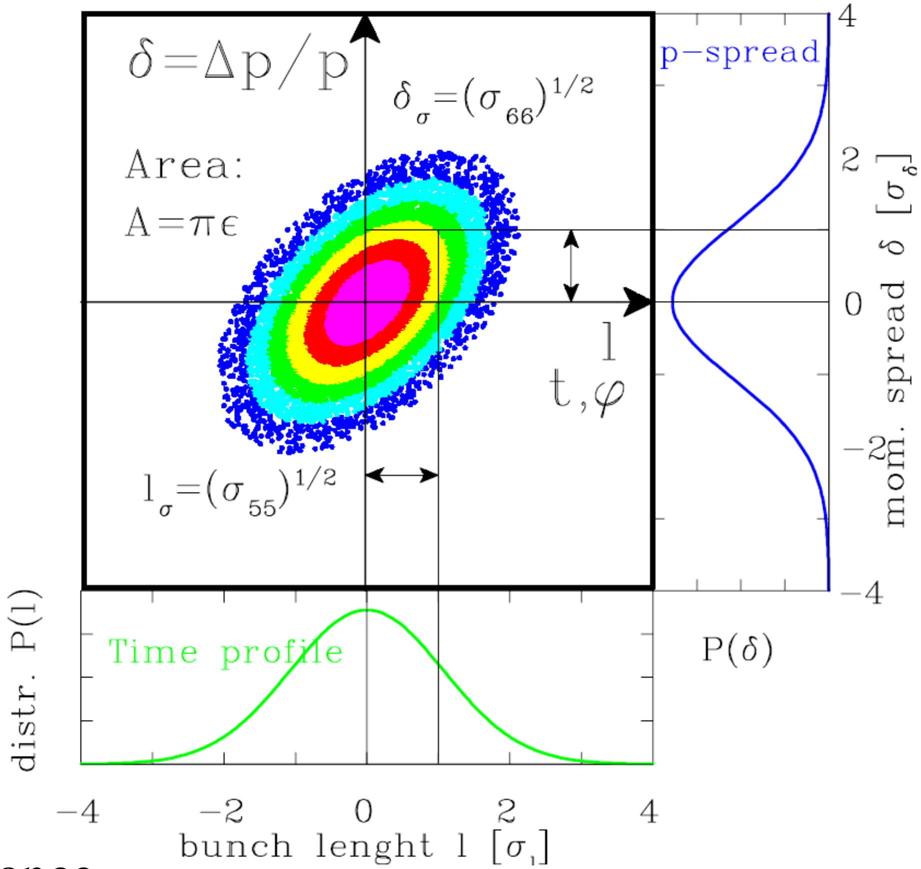
$$\Rightarrow \varepsilon_{long} = \frac{1}{\pi}\int_A dl \cdot d\delta$$

or with density function ***ρ(l,δ)***

$$\Rightarrow \varepsilon_{long} = \frac{1}{\pi}\int \rho(l,\delta)\, dl \cdot d\delta$$

The normalized value is preserved:

$$\varepsilon_{long}^{norm} = \beta\gamma \cdot \varepsilon_{long}$$



Discussed devices:

- Pick-ups for bunch length and emittance.
- Other techniques: Special detectors (low $\boldsymbol{E_{kin}}$ protons), streak cameras (e^-)

The Bunch Position measured by a Pick-Up

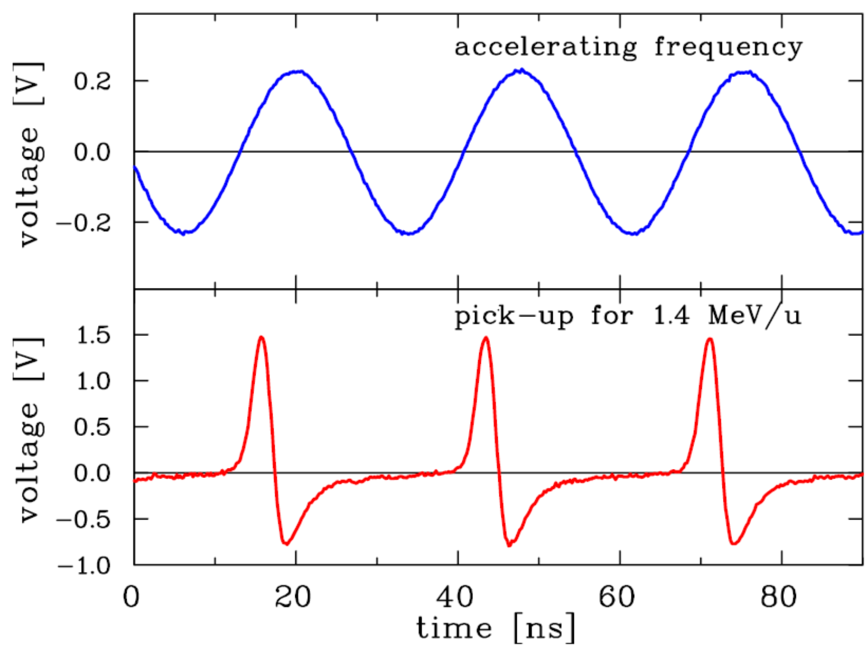
The ***bunch position*** is given relative to the accelerating rf.

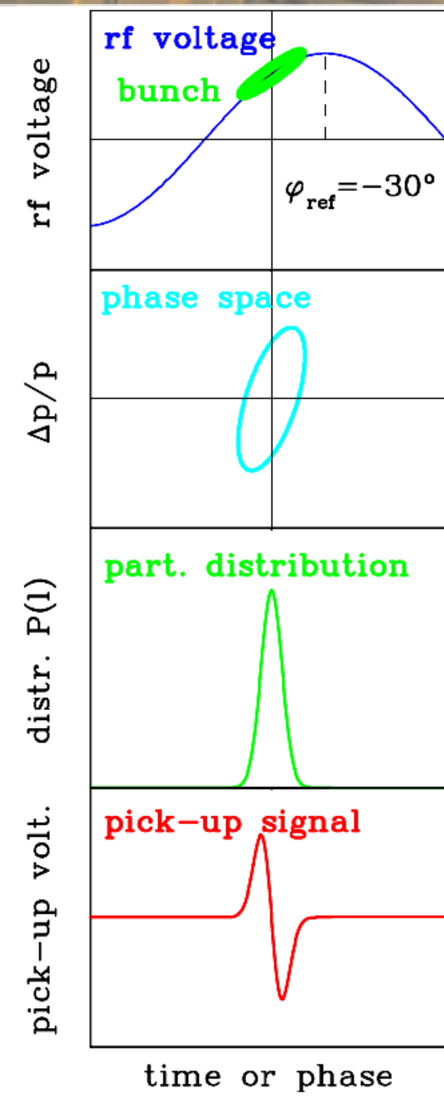
e.g. φ_{ref}=-30^{o} inside a rf cavity

must be well aligned for optimal acceleration

Transverse correspondence: Beam position

Example: Pick-up signal and 36 MHz rf at GSI-LINAC:





Outline:

- Definition of longitudinal phase space
- Proton LINAC: Determination of mean energy

  used for alignment of cavities phase and amplitude
- Determination of longitudinal emittance
- Bunch length measurement for non-relativistic beams
- Bunch length measurement for relativistic beams
- Summary

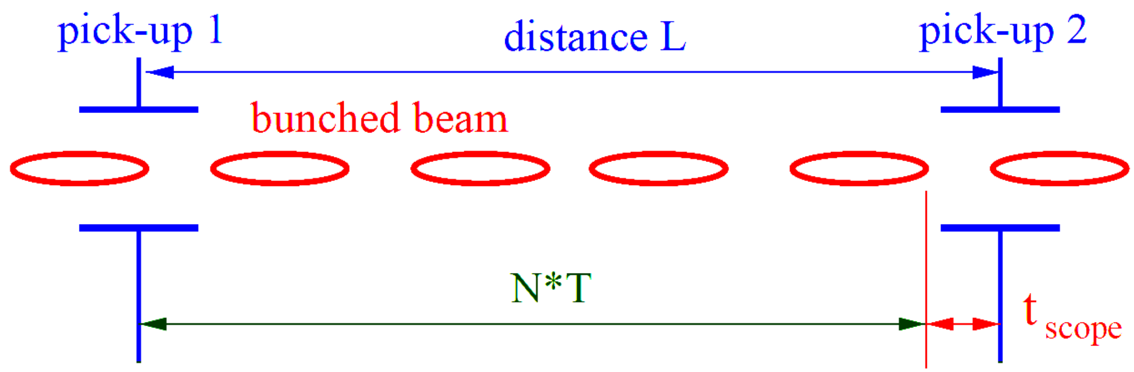
Determination of non-relativistic mean Energy using Pick-Ups

The energy delivered by a LINAC is sensitive to the mechanics, rf-phase and amplitude.

For non-relativistic energies at proton LINACs time-of-flight (TOF) with two pick-ups is used:

$$\beta c = \frac{L}{NT + t_{\mathrm{scope}}}$$

$\rightarrow$ the velocity $\boldsymbol{\beta}$ is measured.



Example: Time-of-flight signal from two pick-ups at 1.4 MeV/u:

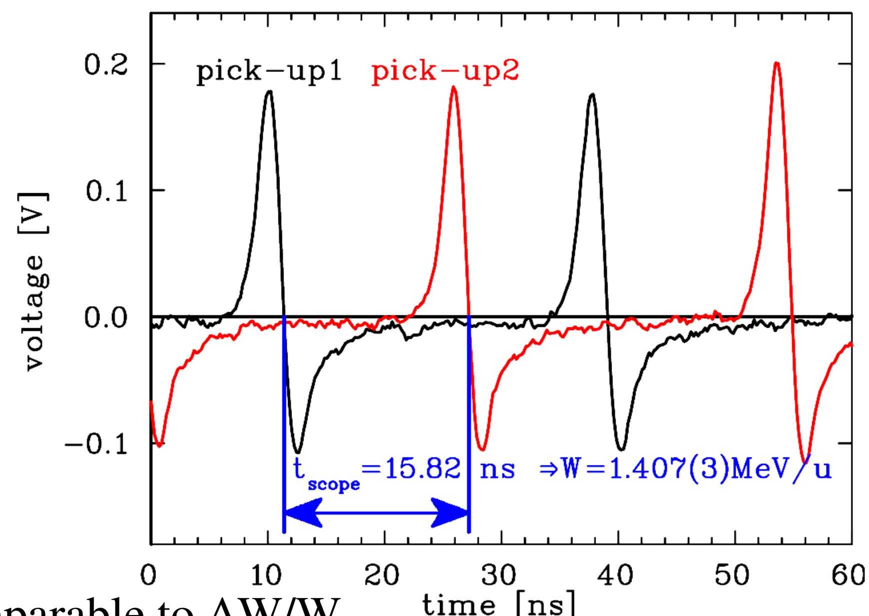
The reading is t_{scope}= 15.82(5)ns with

f_{rf} = 36.136MHz $\Leftrightarrow$ T = 27.673ns

L = 1.629 m and N = 3

$\Rightarrow \boldsymbol{\beta} = 0.05497(7)$

$\Rightarrow$ W=1.407(3) MeV/u



The accuracy is typically 0.1 % i.e. comparable to ΔW/W

Precision of TOF Measurement for non-relativistic Energy

The precision of TOF is given by the accuracy in time and distance reading:

$$\frac{\Delta\beta}{\beta} = \sqrt{\left(\frac{\Delta L}{L}\right)^2 + \left(\frac{\Delta t}{NT + t_{\text{scope}}}\right)^2}$$

Accuracy of scope reading ***Δt*** ≈100 ps, uncertainty in distance **ΔL** ≈ 1 mm.

Example: GSI-LINAC: $\boldsymbol{L}$ = 3.25 m and $\boldsymbol{f_{rf}}$ = 36 MHz:

location (GSI-slang)		RFQ	IH1	IH2	AL4
energy W	[MeV/u]	0.12	0.75	1.4	11.4
velocity β	%	1.6	4.0	5.5	15.5
total TOF	[ns]	677	271	197	70
bunch spacing $\beta c/f_{rf}$	[cm]	13	33	45	129
Number of bunches N		25	9	7	2
resolution $\Delta W/W$	%	0.07	0.10	0.12	0.22

- The accuracy is typically 0.1 % (same order of magnitude as ***ΔW/W***)
- The length has to be matched to the velocity
- Due to the distance of ≈ 3 m, different solutions for the # of bunches ***N*** are possible

→ A third pick-up has to be installed closed by, to get an unique solution.

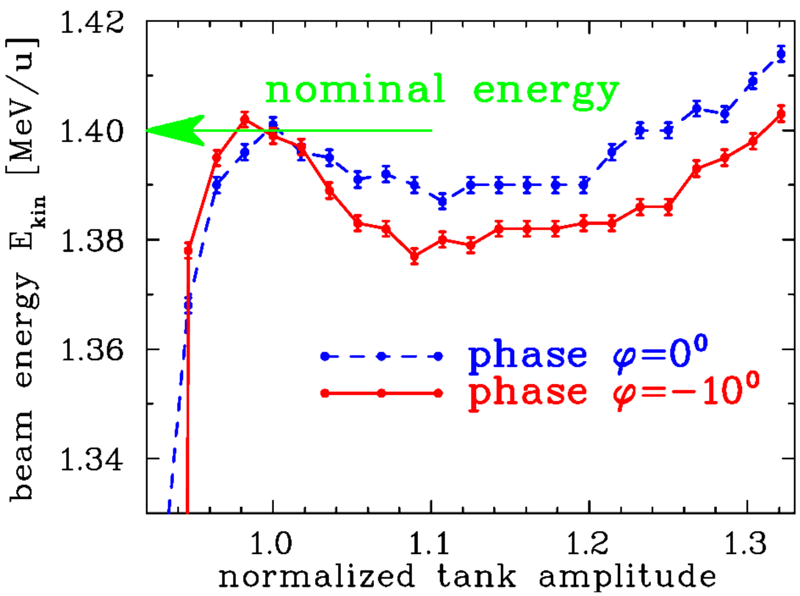
Cavity Alignment using a TOF Measurement

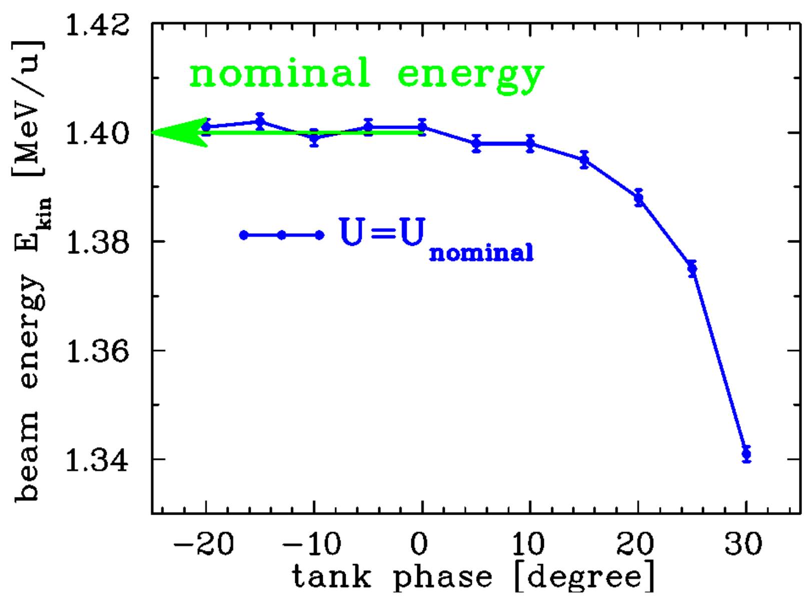
The mean energy is important for the matching between LINAC module.

It depends on phase and amplitude of the rf wave inside the cavities.

Example: Energy at GSI LINAC (nominal energy 1.400 MeV/u):

(distance between pick-ups: $\boldsymbol{L}$ = 1.97 m $\Rightarrow$ $\boldsymbol{N}$ = 4 bunches)





- ***Proton LINACs:*** Amplitude and phase should be carefully aligned by precise TOF
- ***Electron LINACs***: Due to relativistic velocity, TOF is not applicable.

Outline:

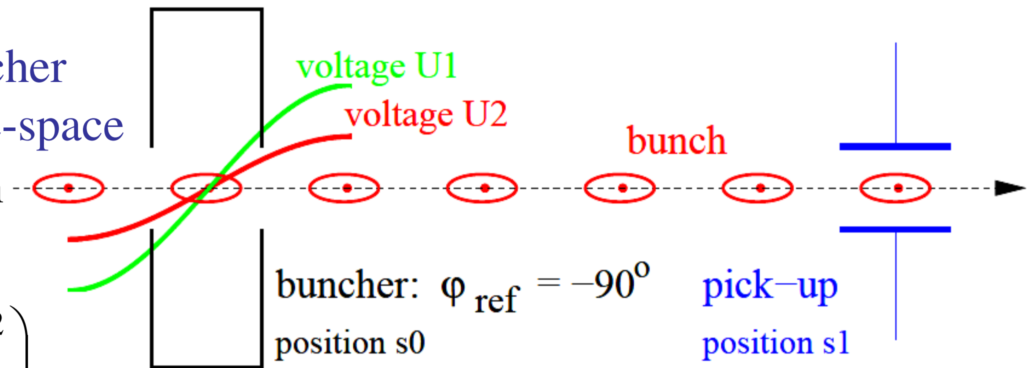
- Definition of longitudinal phase space
- Proton LINAC: Determination of mean energy
  used for alignment of cavities phase and amplitude
- **Determination of longitudinal emittance**
  LINAC: variation of bunch length
  Synchrotron: Topographic reconstruction
- **Bunch length measurement for non-relativistic beams**
- **Bunch length measurement for relativistic beams**
- **Summary**

Longitudinal focusing:

Variation of the bunch shape by a rf-buncher

$\rightarrow$ components 5 and 6 from 6-dim phase-space

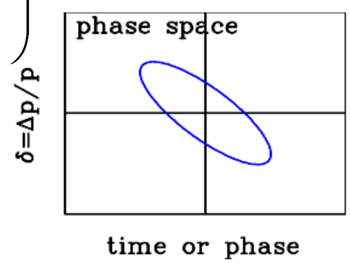
Transversal corres.: Quadrupole variation



- Transfer matrix of buncher & drift:

$$R_{buncher} = \begin{pmatrix} 1 & 0 \\ -1/f & 1 \end{pmatrix}, \; R_{drift} = \begin{pmatrix} 1 & L/\gamma^2 \\ 0 & 1 \end{pmatrix}$$

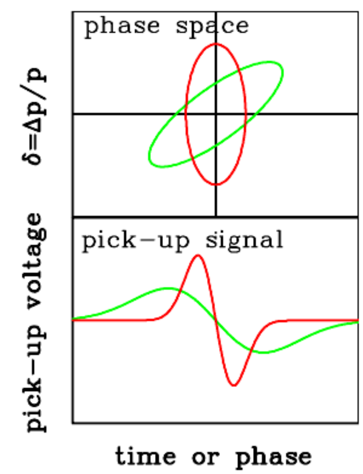
with focal length: $\quad 1/f = \dfrac{2\pi f_{rf}}{Apv^2} \cdot U$



- Variation of buncher amplitude ***U***

  $\Rightarrow$ different bunch width at s_1:

  beam matrix $\Delta t^2_{max} = \sigma_{55}(1, f)$



- System of redundant linear equations for $\sigma_{ij}(0)$:

$$\begin{aligned}
\sigma_{55}(1,f_1) &= R_{55}^2(f_1)\cdot\sigma_{55}(0) + 2R_{55}(f_1)R_{56}(f_1)\cdot\sigma_{56}(0) + R_{56}^2(f_1)\cdot\sigma_{66}(0) && \text{focusing } f_1 \\
&\vdots \\
\sigma_{55}(n,f_n) &= R_{55}^2(f_n)\cdot\sigma_{55}(0) + 2R_{55}(f_n)R_{56}(f_n)\cdot\sigma_{56}(0) + R_{56}^2(f_n)\cdot\sigma_{66}(0) && \text{focusing } f_n
\end{aligned}$$

6-dim Phase Space for Accelerators

The particle trajectory is described with the 6-dim vector $\vec{x}^{\,t} = (x, x', y, y', l, \delta)$

For linear beam behavior the 6x6 transport matrix R is used:

The transformation from location $\boldsymbol{s_0}$ to $\boldsymbol{s_1}$ is:

$$\vec{x}(s_1) = R \cdot \vec{x}(s_0)$$

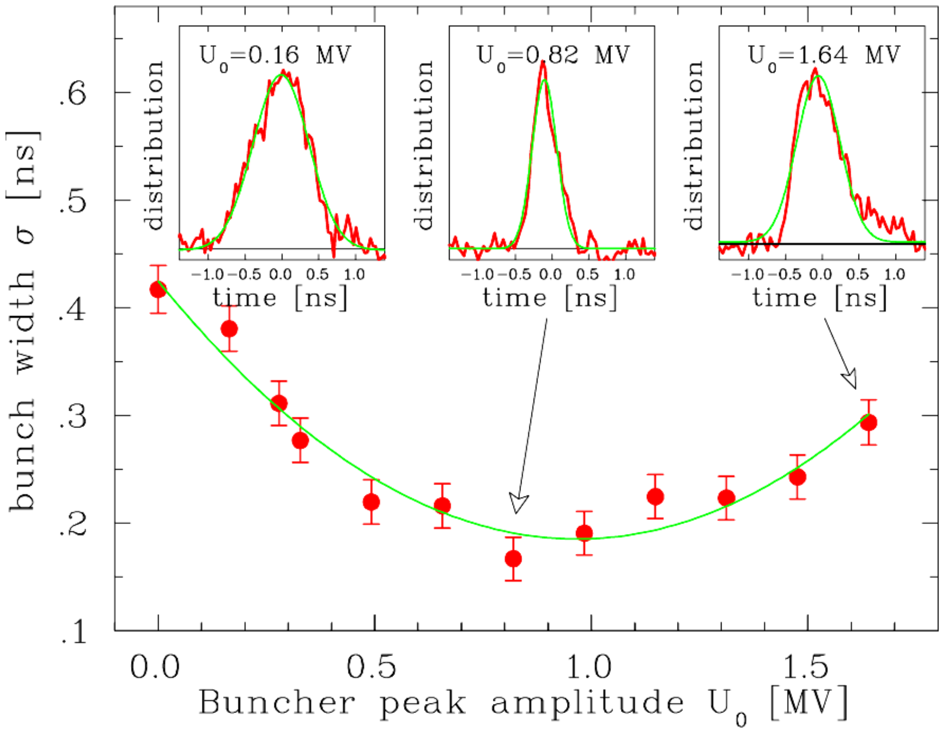
$$= \begin{pmatrix} R_{11} & R_{12} & R_{13} & R_{14} & R_{15} & R_{16} \\ R_{21} & R_{22} & \dots & =0\ \dots & \dots & =0 \dots \\ R_{31} & \dots & R_{33} & R_{34} & \dots & \dots \\ R_{41} & =0\ \dots & R_{43} & R_{44} & \dots & =0 \dots \\ R_{51} & \dots & \dots & \dots & R_{55} & R_{56} \\ R_{61} & =0 \dots & \dots & =0 \dots & R_{65} & R_{66} \end{pmatrix} \cdot \begin{pmatrix} x \\ x' \\ y \\ y' \\ l \\ \delta \end{pmatrix}$$

R separates in 3 matrices only if the horizontal, vertical and longitudinal planes do **not** couple,

e.g. no dispersion $\boldsymbol{D=-R_{16}=0}$

Example GSI LINAC:

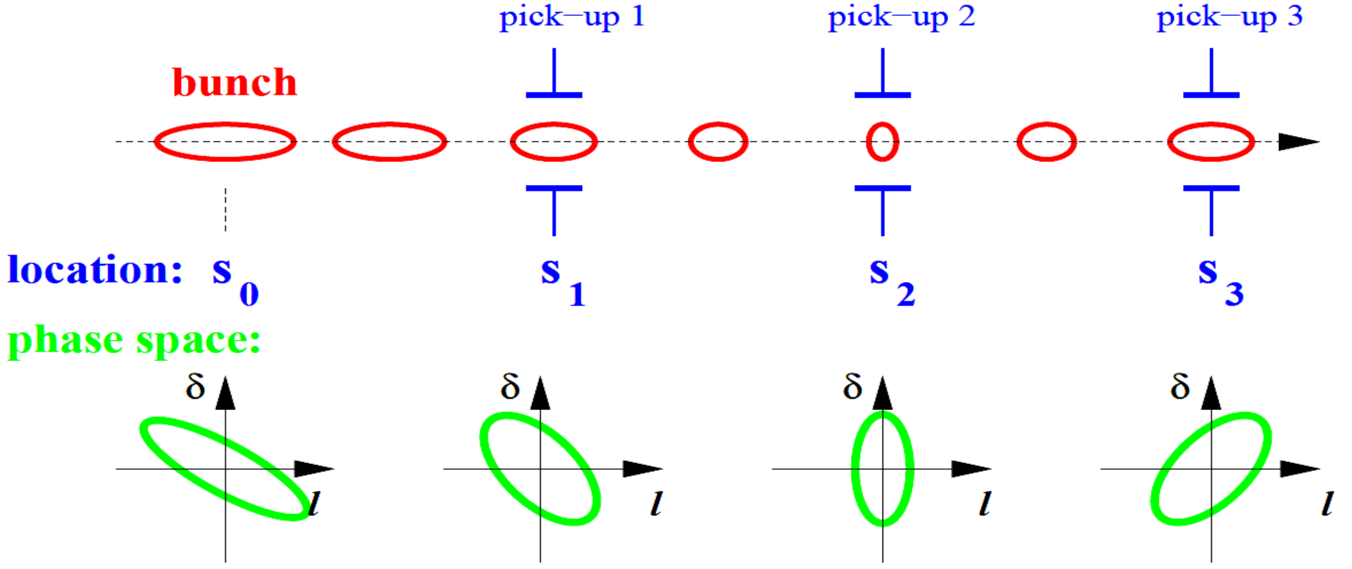
The voltage at the single gap -resonator is varied for 11.4 MeV/u Ni^{14+} beam, 31 m drift:



- The structure of short bunches can be determined with special monitor
- This example: The resolution is better than 50 ps or 2^o for 108 MHz
- Typical bunch length at proton LINACs: 30 to 200 ps

Longitudinal Emittance within a Transfer-Line

As for the 'three grid' method, the emittance can be determined in a transfer line.



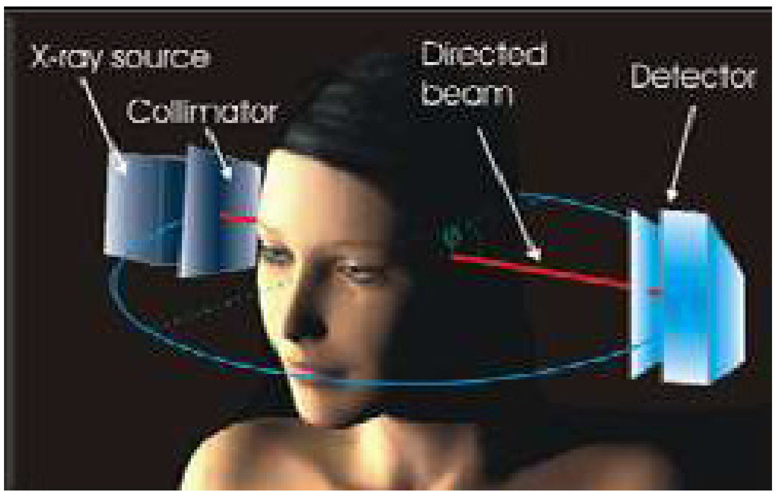
The system of redundant linear equations with the transfer matrix **R(i)** to location s_i:

$$\begin{aligned}
\sigma_{55}(1) &= R_{55}^2(1)\cdot\sigma_{55}(0) + 2R_{55}(1)R_{56}(1)\cdot\sigma_{56}(0) + R_{56}^2(1)\cdot\sigma_{66}(0) \qquad \mathbf{R}(1): s_0 \rightarrow s_1 \\
&\vdots \\
\sigma_{55}(n) &= R_{55}^2(n)\cdot\sigma_{55}(0) + 2R_{55}(n)R_{56}(n)\cdot\sigma_{56}(0) + R_{56}^2(n)\cdot\sigma_{66}(0) \qquad \mathbf{R}(n): s_0 \rightarrow s_n
\end{aligned}$$

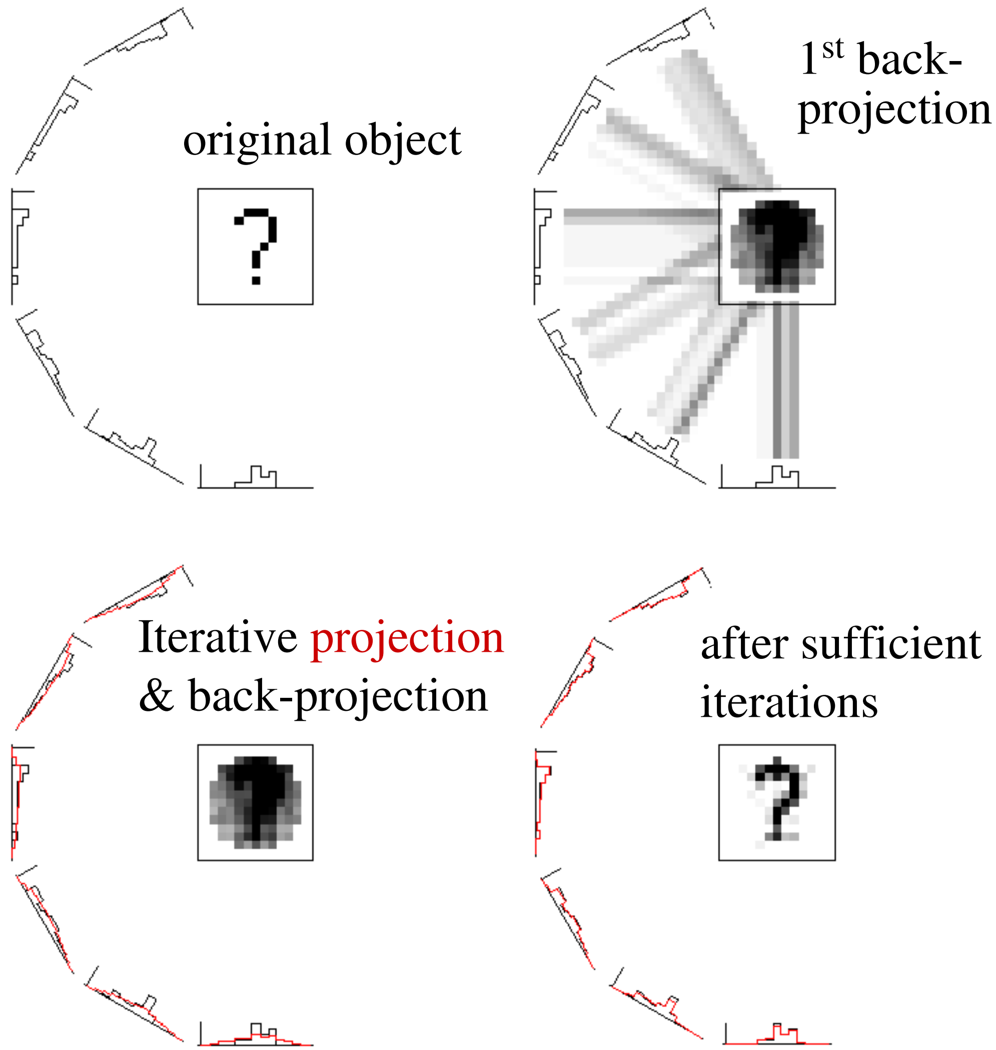
Assumptions:
- Bunches much longer than pick-up or relativistic E -field: $E_\perp >> E_{||}$
- Gaussian distribution without space-charge effects.

Tomography is medical image method

Tomography:
2-dim reconstruction of
sufficient 1-dim projections



Algebraic back projection:
Iterative process by redistributing
the 2-dim image and considering the
differences to the previous iteration step.



Longitudinal Emittance using tomographic Reconstruction

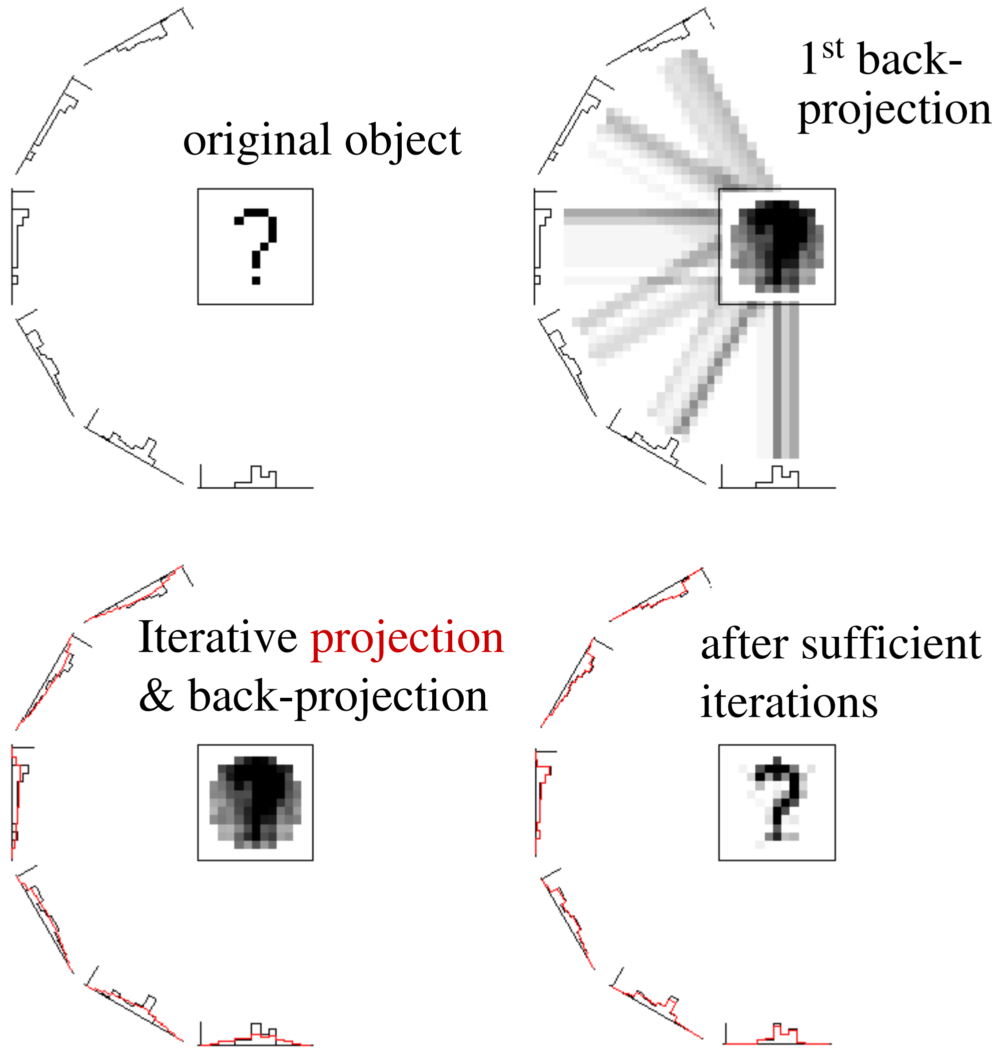
Tomography is medical image method

Tomography:
2-dim reconstruction of sufficient 1-dim projections

Application at accelerators:
Longitudinal emittance evolution in synchrotrons.

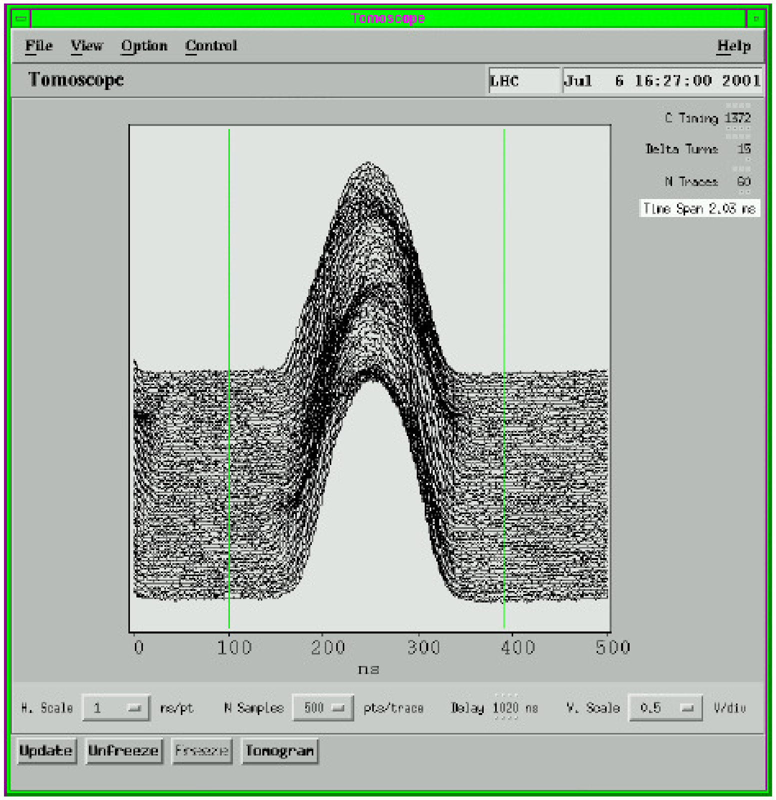
Bunch observation:
Each revolution,
the bunch shape changes a bit
due to synchrotron oscillations.
Fulfilled condition: $f_{synch} \ll f_{ref}$.

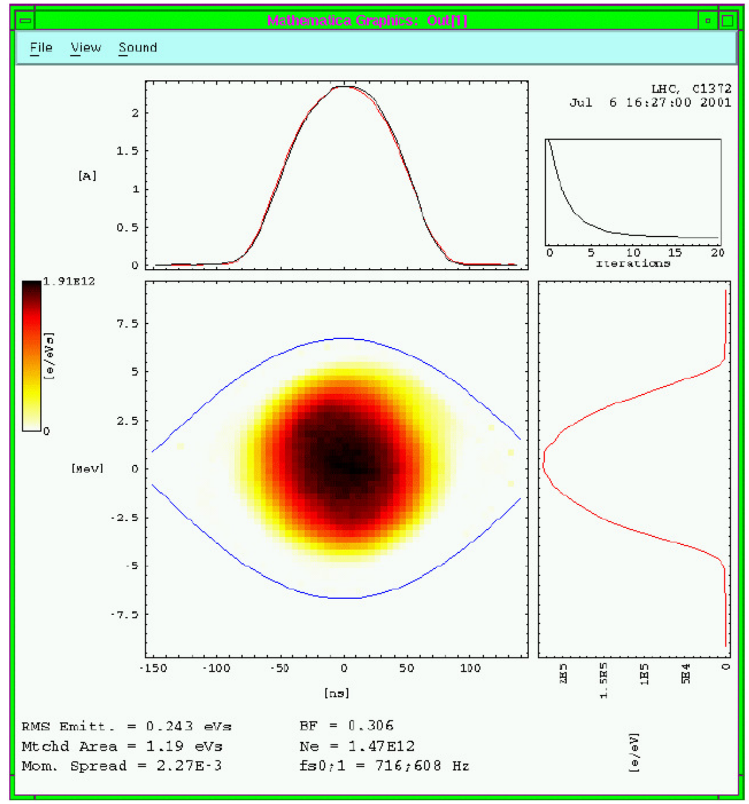
Algebraic back projection:
Iterative process by redistributing the 2-dim image and considering the differences to the previous iteration step.



Results of tomographic Reconstruction at a Synchrotron I

Bunches from 500 turns at the CERN PS and the phase space for the first time slice, measured with a wall current monitor:

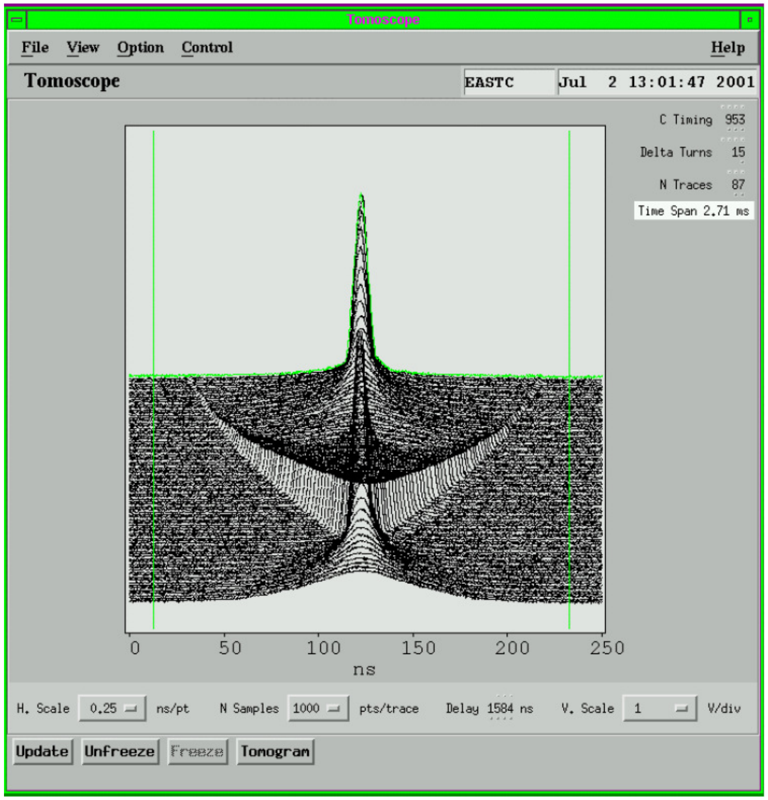


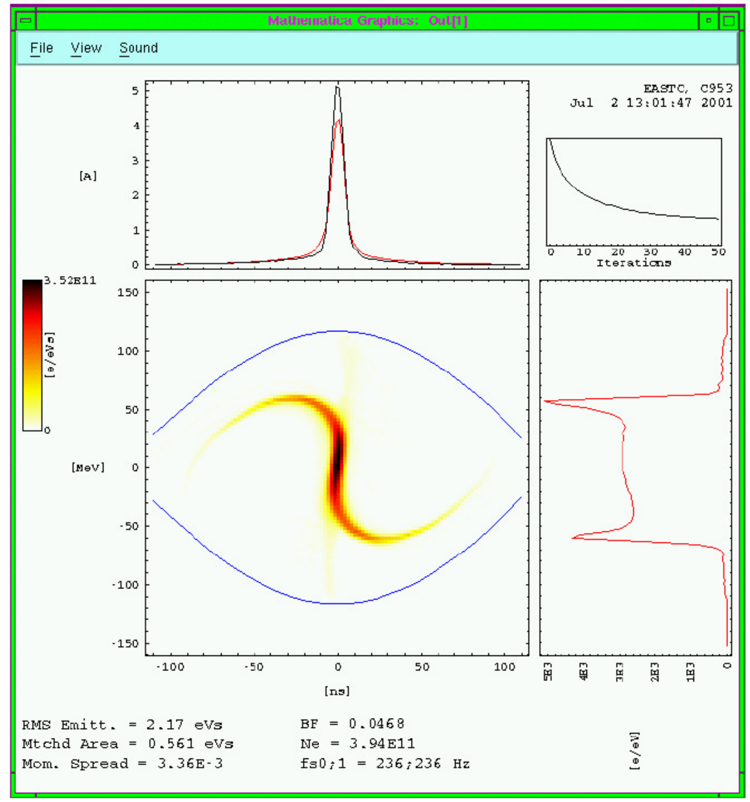


Typical bucket filling. Important knowledge for bunch 'gymnastics'.

Results of tomographic Reconstruction at a Synchrotron II

Bunches from 500 turns at the CERN PS and the phase space for the first time slice, measured with a wall current monitor:





Mismatched bunch shown oscillations and filamentation due to ‘bunch-rotation’.

Resistive Wall Current Monitor

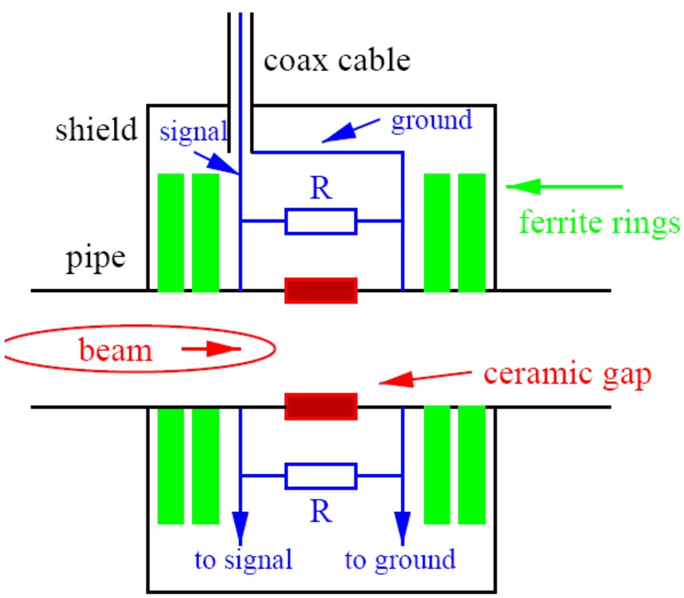
Broadband observation of bunches can be performed with a resistive Wall Current Monitor

Principle:
- Ceramic gap bridged with n =10…100 resistors of R = 10…100 Ω
- Measurement of voltage drop for $R_{tot} = R/n$ = 1…10 Ω
- Ferrite rings with high L → forces low frequency components through resistors

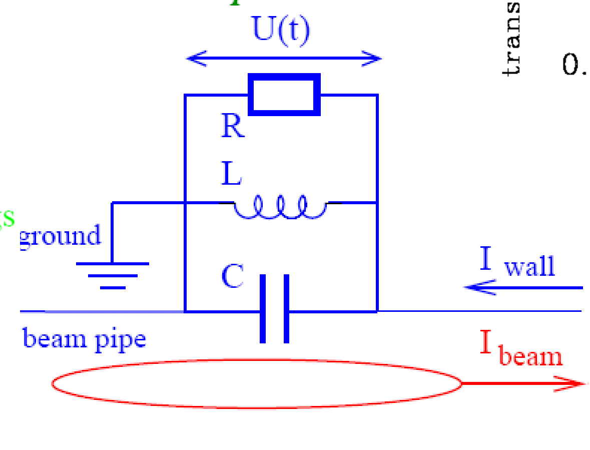
Bandwidth: typically $f_{low} = R_{tot}/(2\pi L) \approx$ **10 kHz**

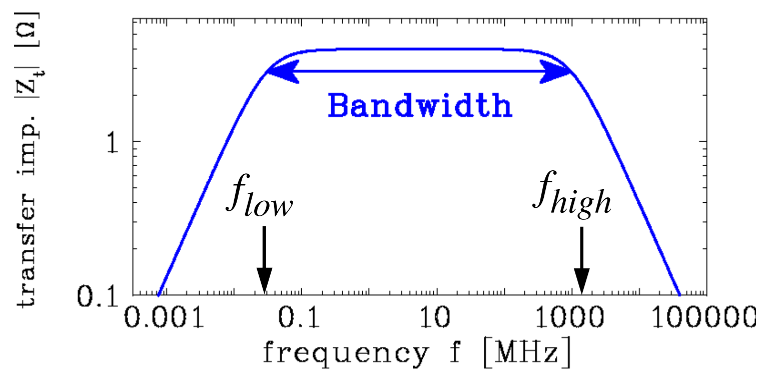
$f_{high} = 1/(2\pi R_{tot} C) \approx$ **1 GHz**

Application: Broadband bunch observation.



WCM equivalent circuit



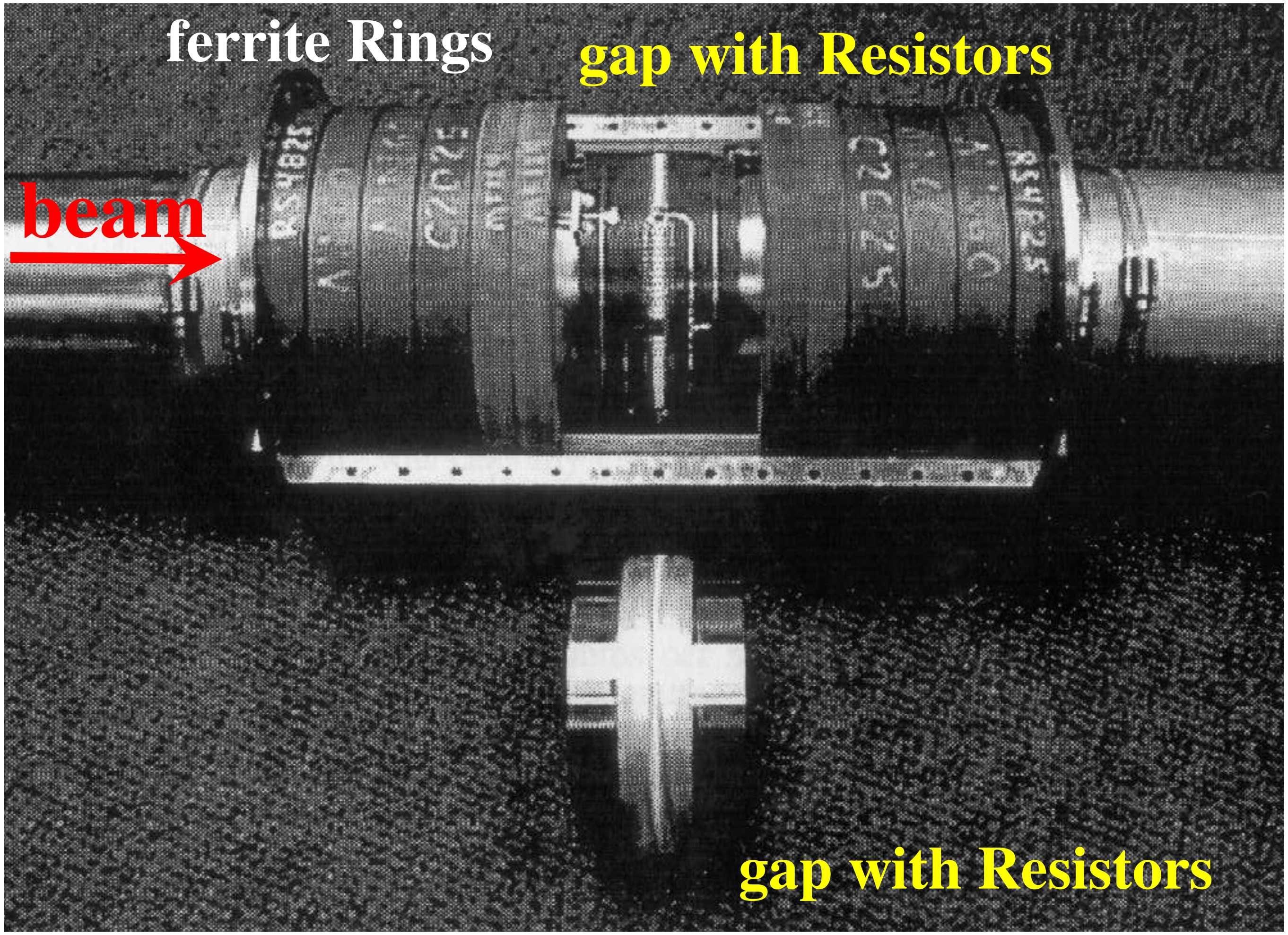


$$\frac{1}{Z_t} = \frac{1}{R_{tot}} + \frac{1}{i\omega L} + i\omega C$$

Within bandwidth: $Z_t \cong R_{tot}$

Resistive Wall Current Monitor

Example: Realization at Fermi-Laboratory



Outline:

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  used for alignment of cavities phase and amplitude
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  determination of particle arrival
- **Bunch length measurement for relativistic beams**
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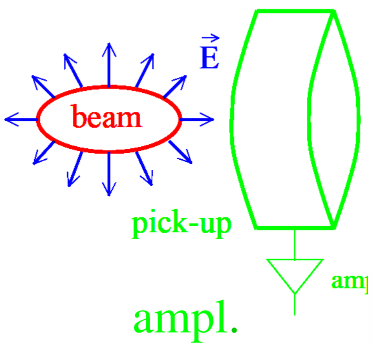
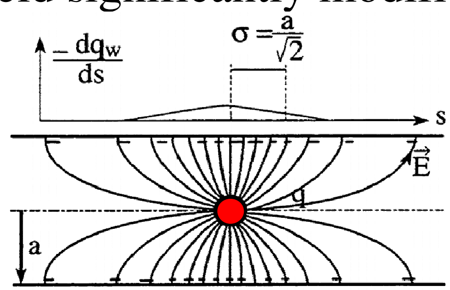
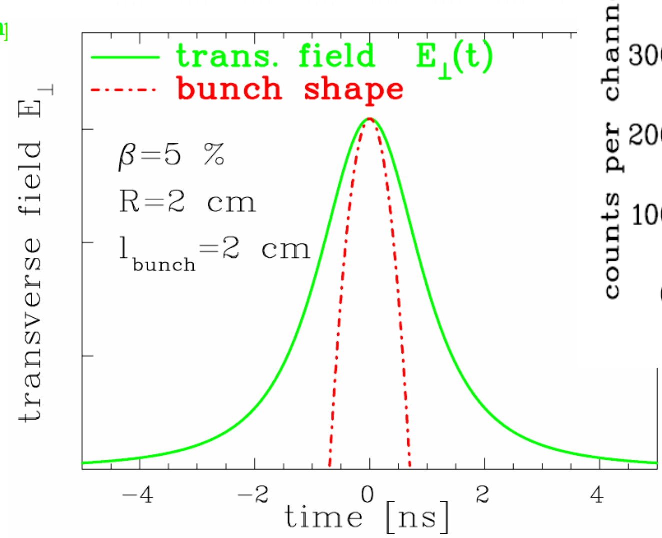
Pick-ups are used for:
- precise for bunch-center relative to rf
- course image of bunch shape

But:

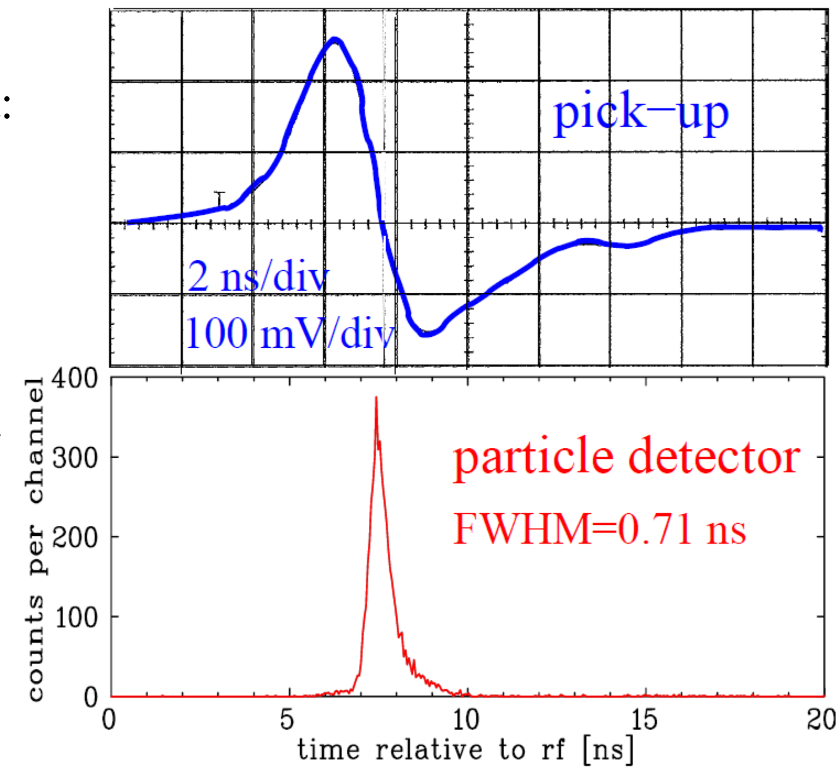
For β << 1 → long. E-field significantly modified:







Example: Comparison pick-up – particle counter:
Ar^{1+} with 1.4 MeV/u (β = 5.5%)



⇒ the pick-up signal is insensitive to bunch 'fine-structure'

Low Velocity Effect: General Consideration

Lorentz transformation of single point-like charge:

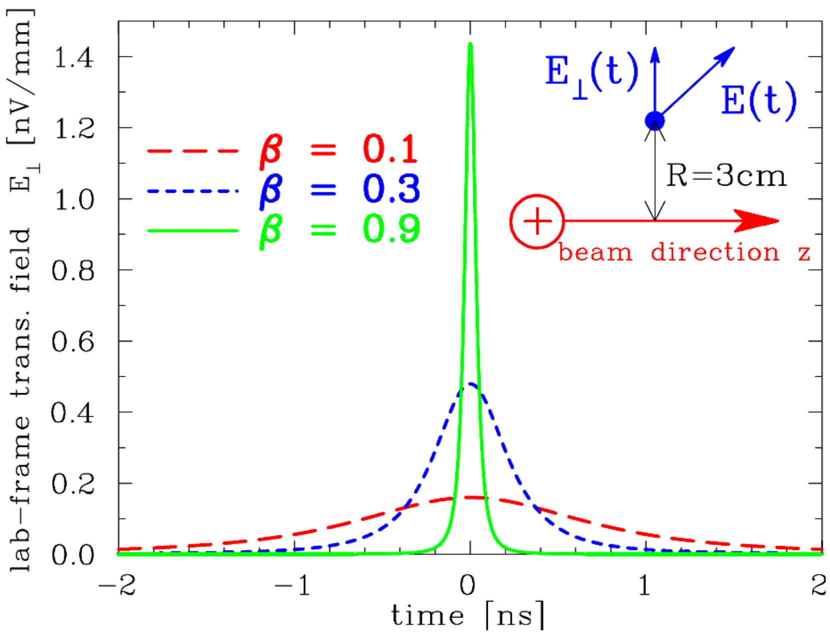
Lorentz boost ***and*** transformation of time: $E_\perp(t) = \gamma \cdot E'_\perp(t')$ and $t \rightarrow t'$

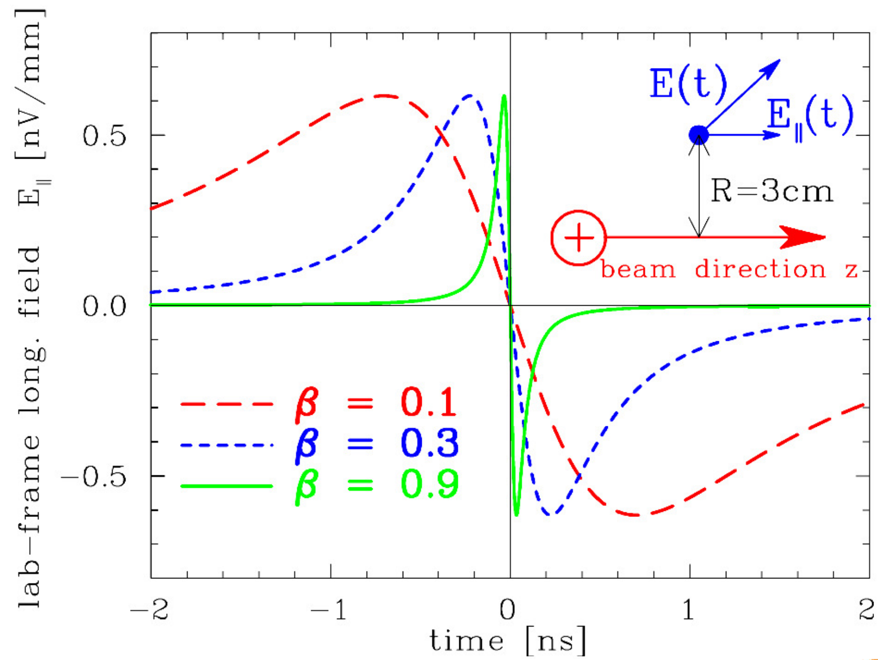
Trans. $\boldsymbol{E}_\perp$ lab.-frame of a point charge:

$$E_\perp(t) = \frac{e}{4\pi\varepsilon_0} \cdot \frac{\gamma R}{\left[R^2 + (\gamma\beta ct)^2\right]^{3/2}}$$

Long. $\boldsymbol{E}_\parallel$ lab.-frame of a point charge:

$$E_\parallel(t) = -\frac{e}{4\pi\varepsilon_0} \cdot \frac{\gamma\beta ct}{\left[R^2 + (\gamma\beta ct)^2\right]^{3/2}}$$





Broadband coaxial Faraday Cups for Bunch Structure

The bunch structure can be observed with cups, having a bandwidth up to several GHz.

Bandwidth and rise time: BW [GHz] = $\mathbf{0.3/t_{rise}}$[ns]

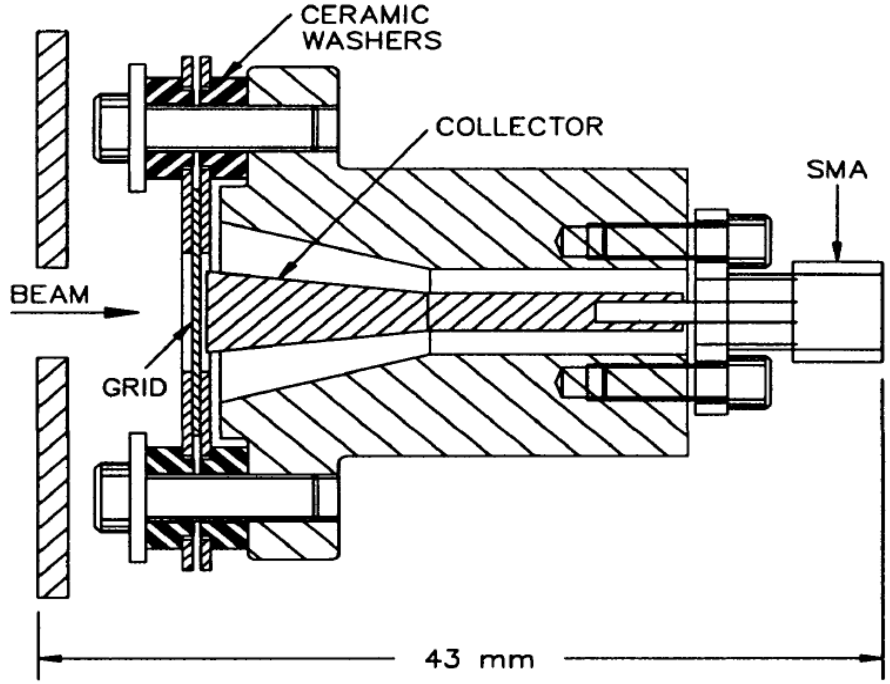
Impedance of a

coaxial transmission line:

$$Z_0 = \frac{Z_c}{2\pi} \cdot \ln \frac{r_{outer}}{r_{inner}}$$

with $$Z_c = \sqrt{\frac{\mu_0 \mu_r}{\varepsilon_0 \varepsilon_r}}$$

$\rightarrow$ impedance matching to prevent for reflections



Voltage reflection: $$\rho_V = \frac{Z - Z_0}{Z + Z_0}$$

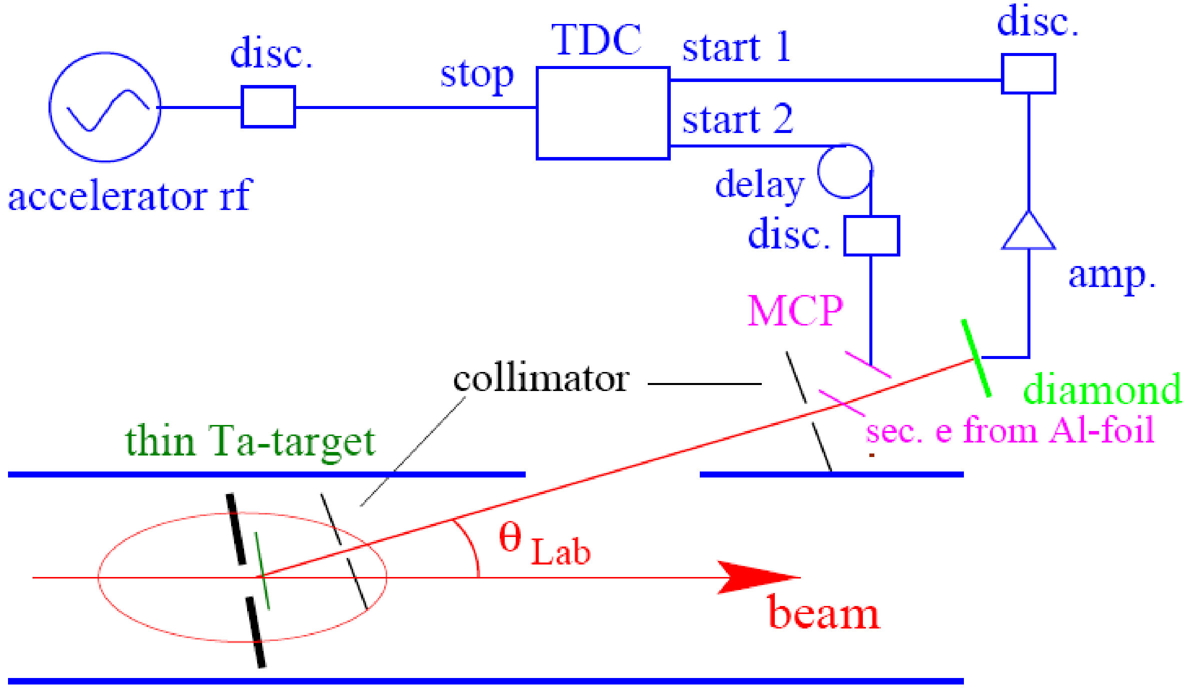
Voltage Standing Wave Ratio: $$\mathrm{VSWR} = \frac{Z}{Z_0} = \frac{1+\rho_V}{1-\rho_V}$$

$\boldsymbol{Z = Z_0}$: no reflection. $\boldsymbol{Z = 0} \Rightarrow \boldsymbol{\rho_V} = -1$: short circuit. $\boldsymbol{Z} = \infty \Rightarrow \boldsymbol{\rho_V} = 1$: open circuit.

Realization of a Broadband coaxial Faraday Cup

Time-of-Flight using Particle Detectors

The time of arrival of the particle is determined relative to the accelerating rf:



Realization at GSI-LINAC: Less than one particle per bunch due to ***single*** particle counting:

→ **Foil (130 nm):** attenuation $\approx 10^{-9}$ by Rutherford scat. ⇒finite solid angle $\Delta\Omega_{lab} = 2.5 \cdot 10^{-4}$

→ **Stop-detectors:** Fast detector with 1 ns pulse width (diamond)

→ **TDC:** Time relative to rf, resolution less than 25 ps (corresponding to 0.3^o in phase)

→ **Start-detector**: 2nd thin Al foil (50 nm) for secondary e^- acc. toward an MCP +50 Ω anode

⇒ **Result:** determination of phase and energy of ***individual*** particles.

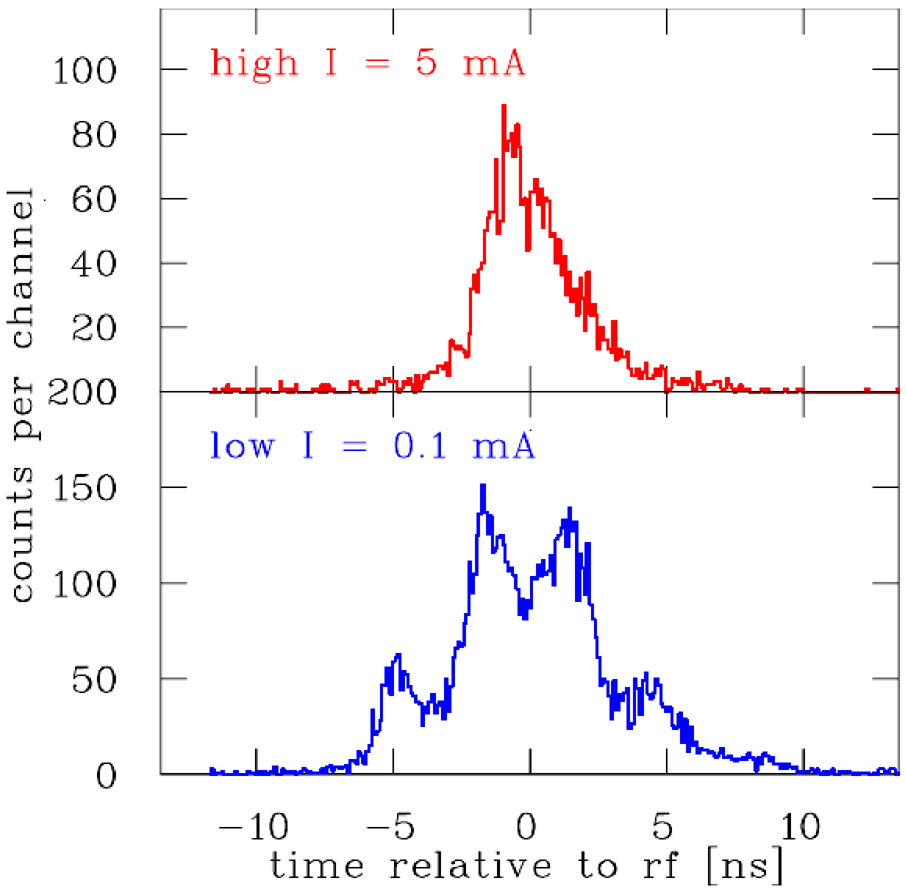
Result of Bunch Structure Determination at low E_{kin}

The bunch structure is dependent on the amplitude or phase setting

→ wrong bunching (RFQ), emittance blow-up, filamentation...

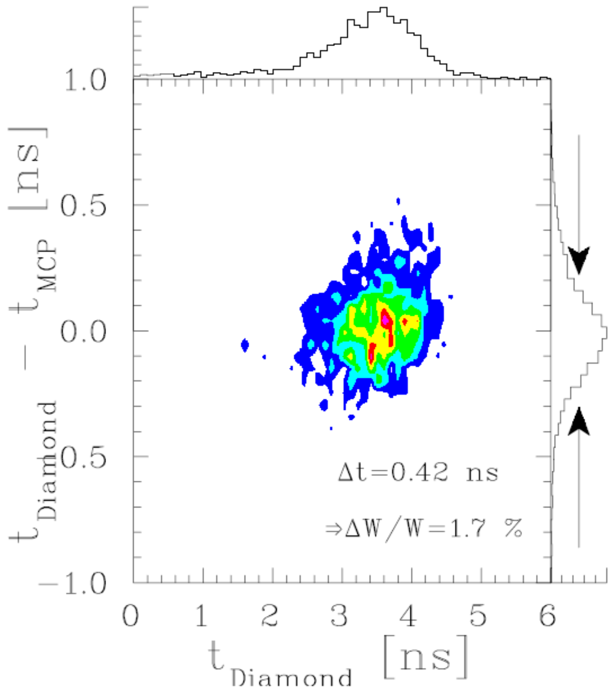
→ non-Gaussian distributions are possible

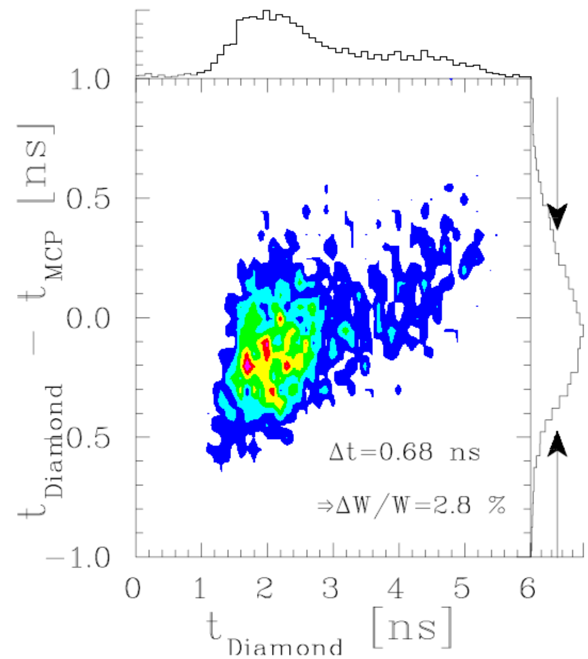
Example: The bunch shape at 120 keV/u from GSI-LINAC with different currents:



No 'standard' method for longitudinal emittance measurement is available!
Using two detectors in coincidence and a drift space in between,
the phase and the energy of a individual particle can be determined.
$\Rightarrow$ for many particles the longitudinal phase space can be spanned.

Example: GSI-LINAC at 1.4 MeV/u with ***low*** and ***high*** current Ar beam
The effect of the emittance blow-up due to the large space-charge is seen.





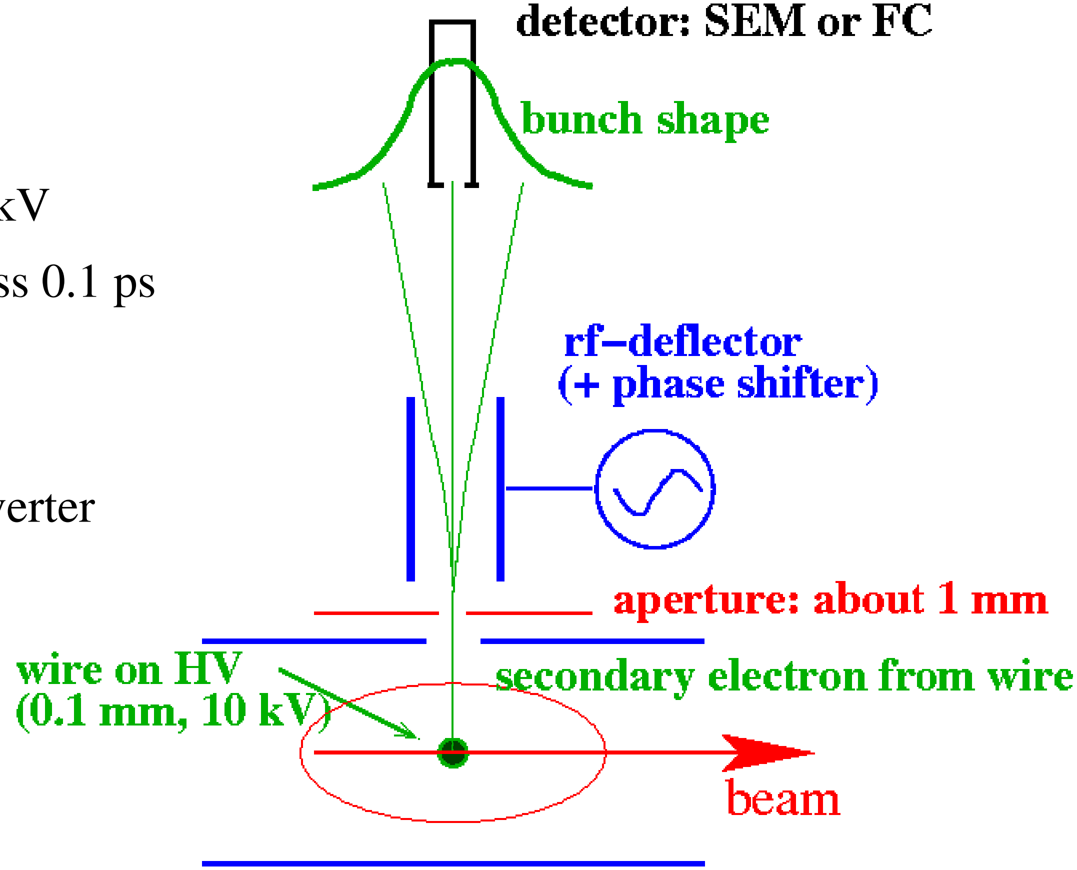
Secondary e^- liberated from a wire carrying the time information.

$\rightarrow$ Bunch Shape Monitor (BSM)

Working principle:

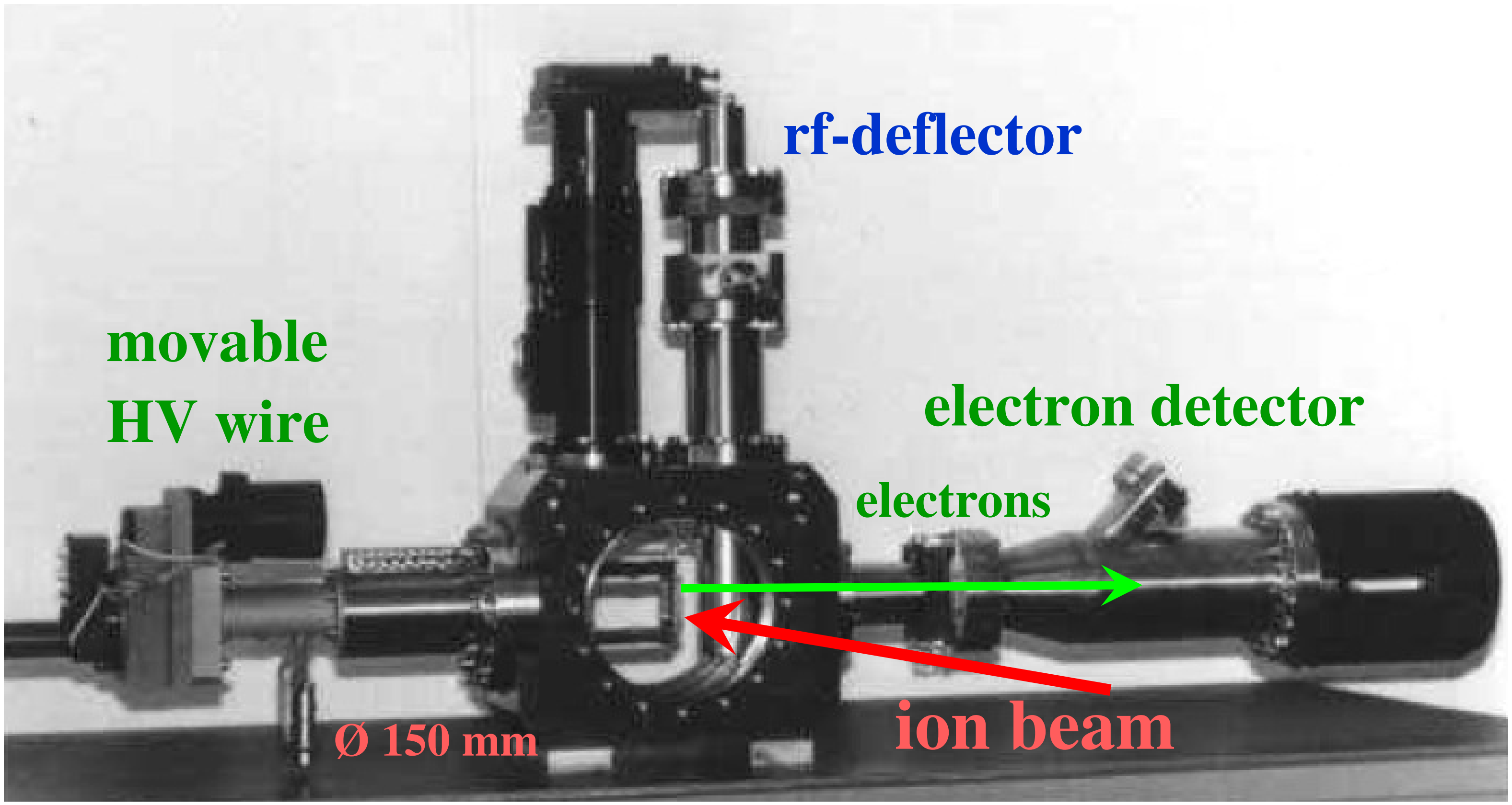
- insertion of a 0.1 mm wire at ≈ 10 kV
- emission of secondary e^- within less 0.1 ps
- secondary e^- are accelerated
- toward an rf-deflector
- rf-deflector as 'time-to-space' converter
- detector with a thin slit
- slow shift of the phase
- resolution $\approx 1^o < 10$ ps
- Measurements are comparable

  to that obtained with particle detectors.



SEM: secondary electron multiplier

Realization of Bunch Shape Monitor at CERN LINAC2



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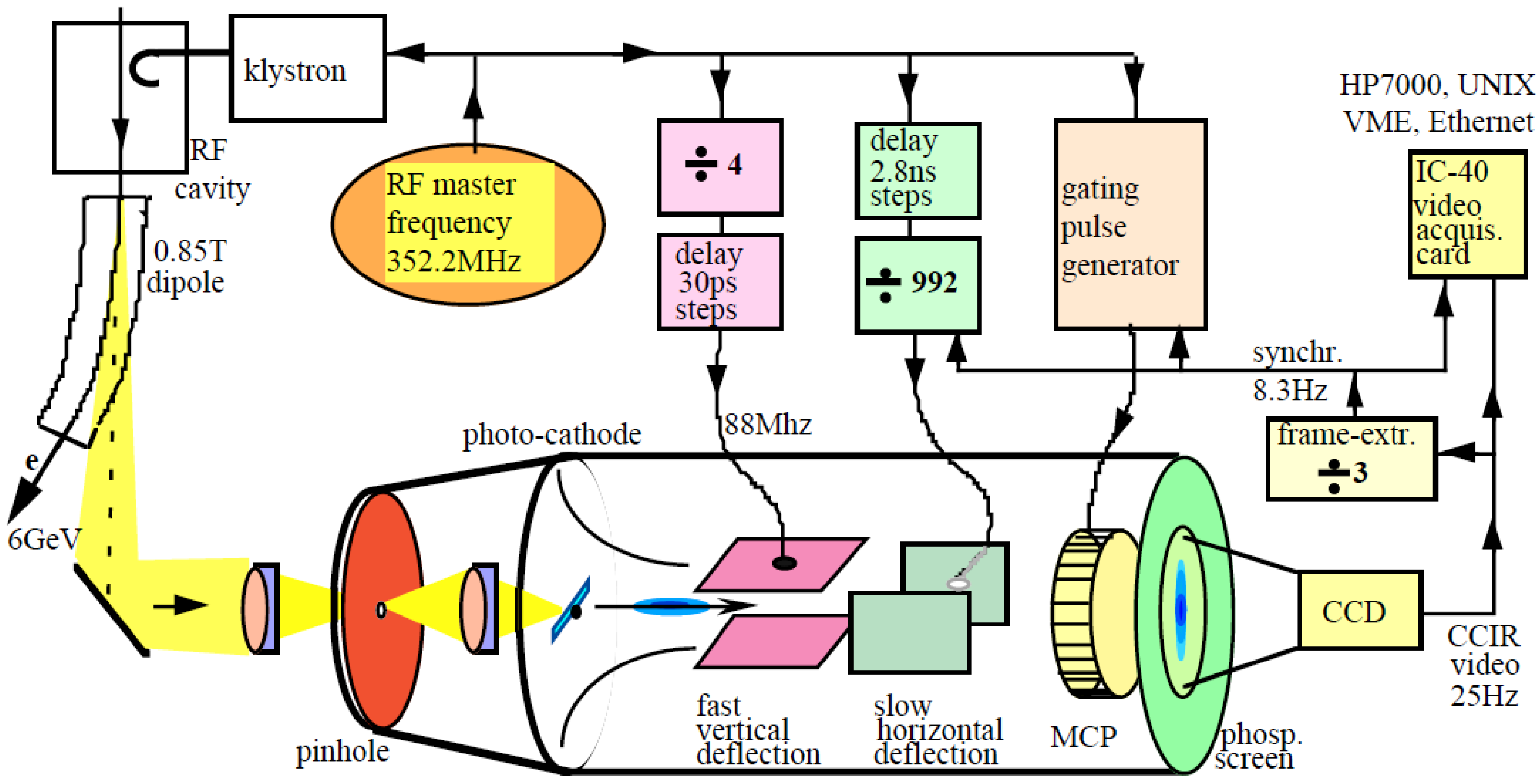
  synchrotron light monitor or electro-optical modulation
- **Summary**

Bunch Length Measurement for relativistic e^-

Electron bunches are too short (σ_t < 300 ps) to be covered by the bandwidth of pick-ups (f < 1 GHz $\Leftrightarrow$ t_{rise}>300 ps) for structure determination.

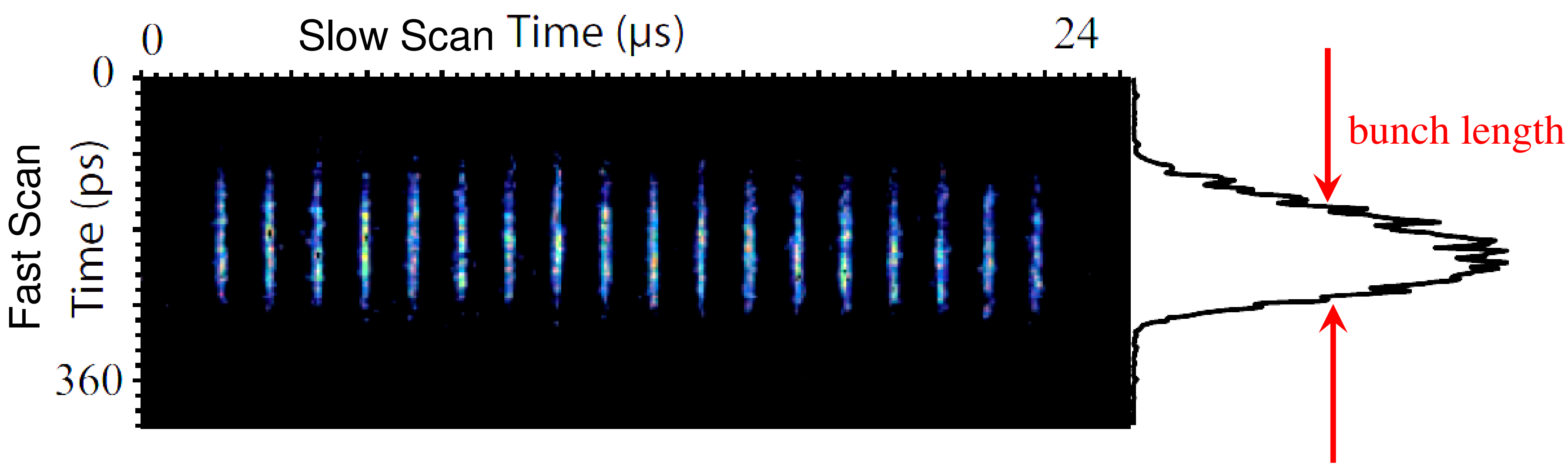
→ Time resolved observation of synchr. light with a streak camera: Resolution ≈ 1 ps.



Results of Bunch Length Measurement by a Streak Camera

The streak camera delivers a fast scan in vertical direction (here 360 ps full scale) and a slower scan in horizontal direction (24 μs).

Example: Bunch length at the synch. Light source SOLEIL for $\boldsymbol{U_{rf}}$ = 2 MV for slow direction 24 μs and scaling for fast scan 360 ps: measure $\boldsymbol{\sigma_t}$ = 35 ps.



The Artist View of a Streak Camera

FARADAY CUP 1998

Purpose. To recognize and encourage innovative achievements in the field of accelerator beam instrumentation.

Award. The Faraday Cup Award consists of a US$ 5000 prize and a certificate to be presented at the next Beam Instrumentation Workshop. Winners participating in the BIW will be given a $1000 travel allowance.

Eligibility. Nominations are open to contributors of all nations regardless of the geographical location at which the work was done.
The Award goes normally to one person, but may be shared by recipients having contributed to the same accomplishment. It will normally be awarded to scientists in the early stage of their career. Nominations of candidates shall remain active for 2 competitions.

Establishment and support. The Award was established in 1991 with the support of the Beam Instrumentation Workshop Organizing Committee.

Rules. The Faraday Cup shall be awarded for an outstanding contribution to the development of an innovative beam diagnostics instrument of proven workability. The Faraday Cup is only awarded for published contribution and delivered performance –as opposed to theoretical performance. Rules are available on request.

Award Committee. The Beam Instrumentation Workshop Organizing Committee.

Nominations. The nomination package shall include the name of the candidate, relevant publications, a statement outlining his/her personal contribution and that of others, two letters from co-workers familiar with the candidate and his contribution. Two master copies suitable for photocopying of this package must be submitted not later than the 15th of November 1997 to Steven Smith c/o BIW'98 Secretariat, SLAC, Stanford University, Stanford CA 94305-4085, U.S.A..

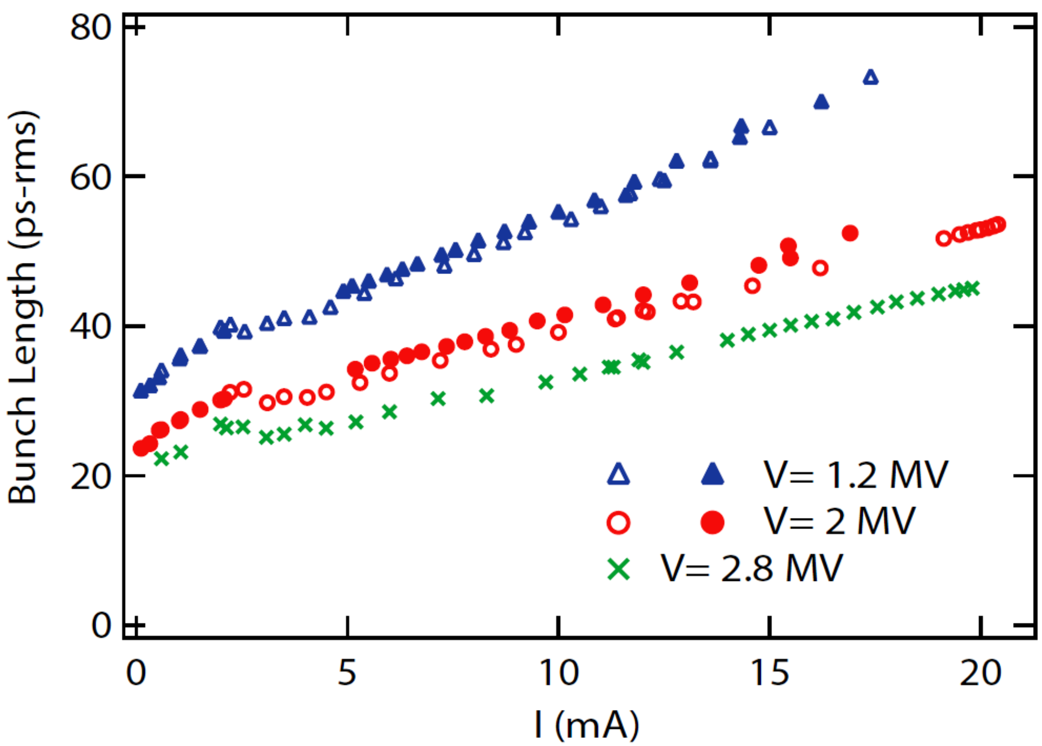
The Importance of Bunch Length by Streak Camera

Short bunches are desired by the synchrotron light users for time resolved spectroscopy.

The bunch focusing is changed by the rf-amplitude.

Example: Bunch length σ_t as a function of stored current

(space-charge de-focusing, impedance broadening) for different rf-amplitudes at SOLEIL:



For Free Electron Lasers → bunch length below 1 ps is used

→ below resolution of streak camera

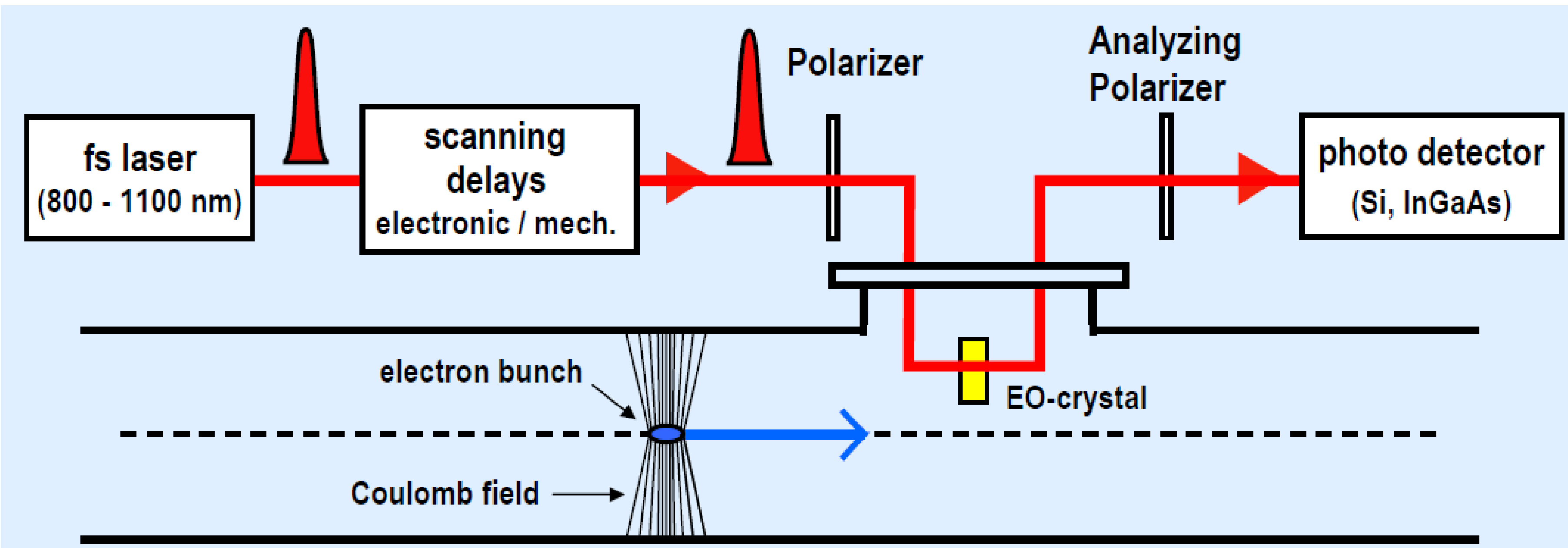
→ short laser pulses with $t \approx$ **10 fs and** electro-optical modulator

Electro optical modulator: birefringent, rotation angle depends on external electric field

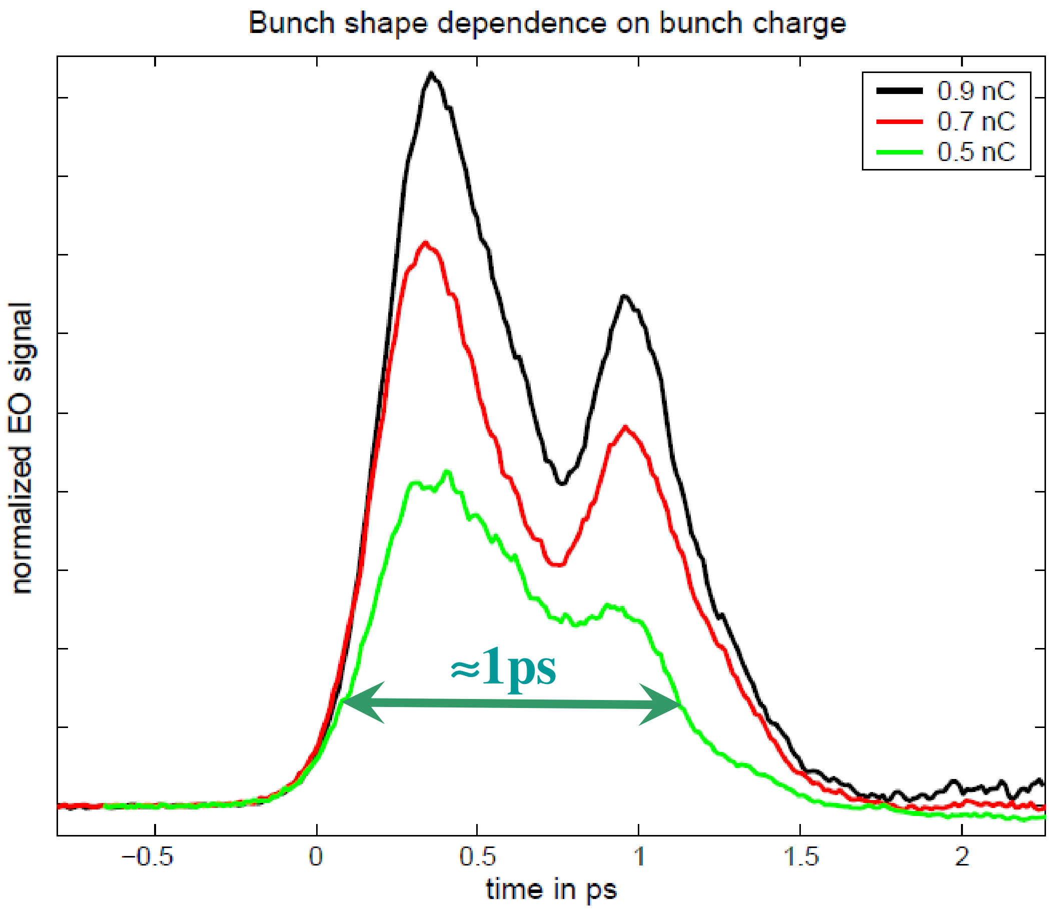
Relativistic electron bunches: transverse field $E_{\perp,\,lab} = \gamma E_{\perp,\,rest}$ carries the time information

Scanning of delay between bunch and laser → time profile after several pulses

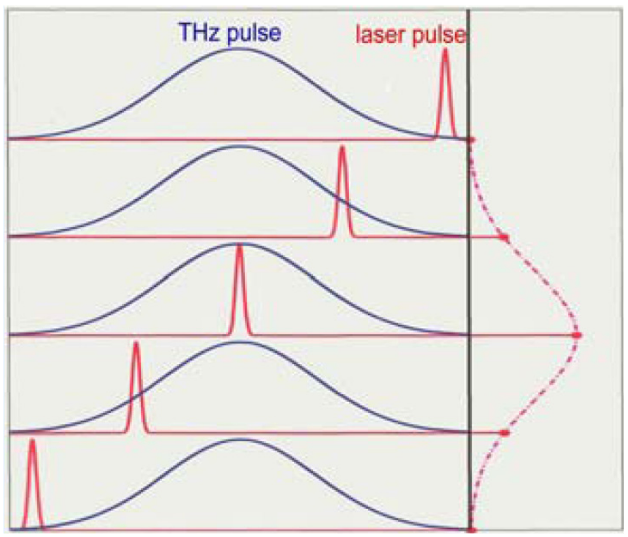
Additionally, single shot modifications successfully tested.



Example: Bunch length at FEL test facility FLASH



Scanning of the short laser pulse relative to bunch:

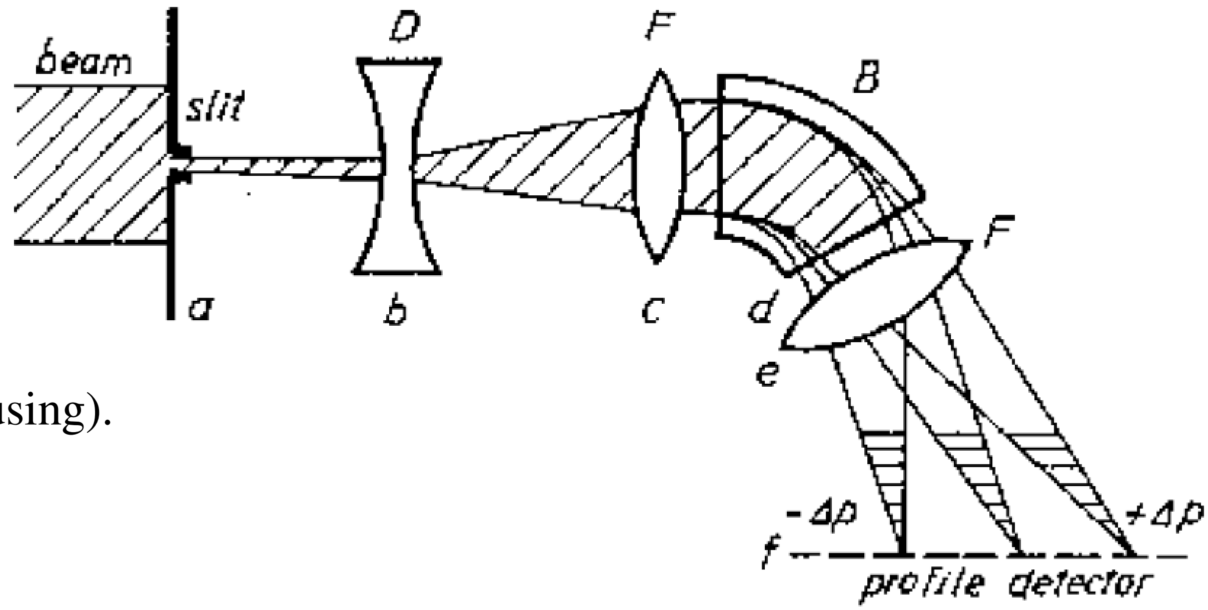


Results at FLASH, Hamburg, see B. Steffen et al., FEL Conf. Stanford, p. 549, 2005.

Measurement of Energy Spread by magnetic Spectrometer

The momentum $\boldsymbol{\delta = \Delta p/p}$ or energy spread $\boldsymbol{\Delta W/W}$ can be determined with a magnetic spectrometer:
→ Via dispersion, the momentum is shifted to a spatial distance.

The right optics has
to be chosen
to separate the transverse and
longitudinal parameters
(transverse point-to-point focusing).



Longitudinal ↔ transverse correspondences:

- position relative to rf ↔ transverse center-of-mass
- bunch structure in time ↔ transverse profile in space
- momentum or energy spread ↔ transverse divergence.

Determination uses:

Broadband pick-ups:
- position relative to rf, mean energy
- emittance at transfer lines or synchrotron via tomography
  assumption: bunches longer than pick-up.

Particle detectors:
- TOF or secondary e^- from wire
  → for non-relativistic proton beams
  reason: ***E***-field does not reflect bunch shape.

Streak cameras:
- time resolved monitoring of synchrotron radiation
  → for relativistic e^--beams, t_{bunch} < 1ns
  reason: too short bunches for rf electronics.

Laser scanning:
- Electro-optical modulation of short laser pulse
  → highest time resolution.