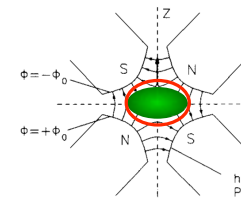
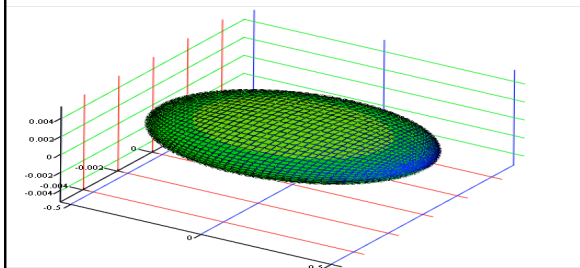


# Introduction to Transverse Beam Optics

Bernhard Holzer,  
CERN-LHC

## II.) The Ideal World: Particle Trajectories & Beams



Bunch in a storage ring

### Reminder of Part I

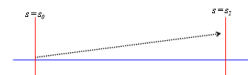
#### Equation of Motion:

$$x'' + Kx = 0 \quad K = 1/\rho^2 - k \dots \text{hor. plane:}$$

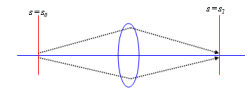
$$K = k \quad \dots \text{vert. Plane:}$$

#### Solution of Trajectory Equations

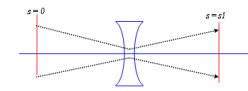
$$\begin{pmatrix} x \\ x' \end{pmatrix}_{s1} = M^* \begin{pmatrix} x \\ x' \end{pmatrix}_{s0}$$



$$M_{drift} = \begin{pmatrix} 1 & l \\ 0 & 1 \end{pmatrix}$$



$$M_{foc} = \begin{pmatrix} \cos(\sqrt{|K|}l) & \frac{1}{\sqrt{|K|}} \sin(\sqrt{|K|}l) \\ -\sqrt{|K|} \sin(\sqrt{|K|}l) & \cos(\sqrt{|K|}l) \end{pmatrix}$$



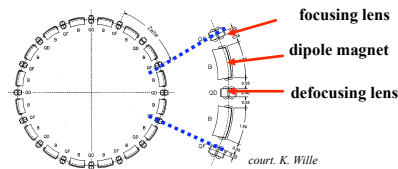
$$M_{defoc} = \begin{pmatrix} \cosh(\sqrt{|K|}l) & \frac{1}{\sqrt{|K|}} \sinh(\sqrt{|K|}l) \\ \sqrt{|K|} \sinh(\sqrt{|K|}l) & \cosh(\sqrt{|K|}l) \end{pmatrix}$$

**Transformation through a system of lattice elements**

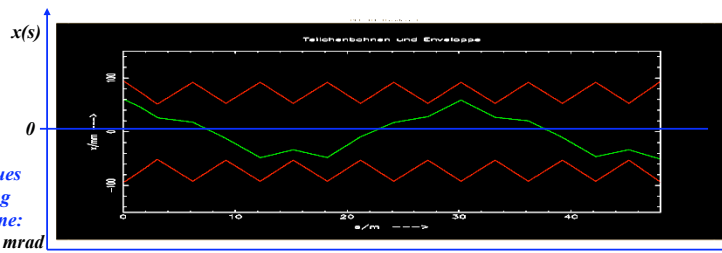
combine the single element solutions by multiplication of the matrices

$$M_{total} = M_{QF} * M_D * M_{QD} * M_{Bend} * M_D * \dots$$

$$\begin{pmatrix} x \\ x' \end{pmatrix}_{s2} = M(s_2, s_1) * \begin{pmatrix} x \\ x' \end{pmatrix}_{s1}$$



in each accelerator element the particle trajectory corresponds to the movement of a harmonic oscillator „

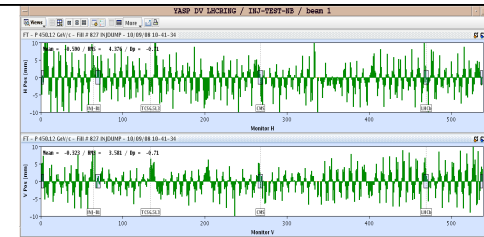


**5.) Orbit & Tune:**

Tune: number of oscillations per turn

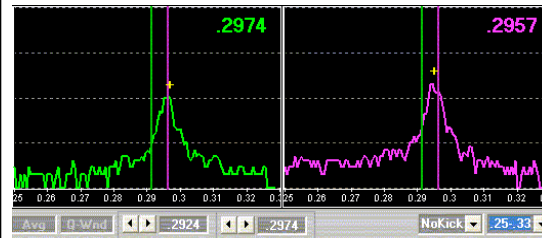
64.31  
59.32

Relevant for beam stability:  
non integer part



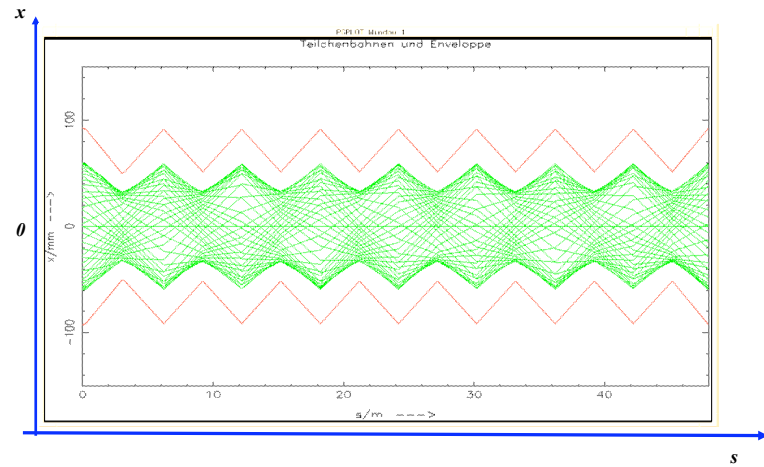
LHC revolution frequency: 11.3 kHz

$$0.31 * 11.3 = 3.5 kHz$$



**Question: what will happen, if the particle performs a second turn ?**

... or a third one or ...  $10^{10}$  turns

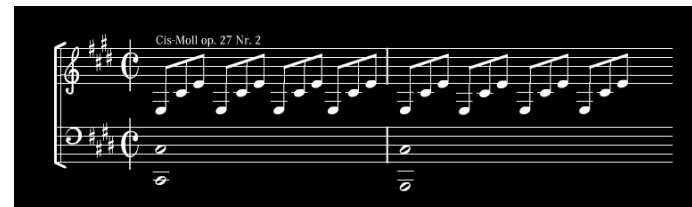


19th century:

Ludwig van Beethoven: „Mondschein Sonate“



Sonate Nr. 14 in cis-Moll (op. 27/II, 1801)



**Astronomer Hill:**

*differential equation for motions with periodic focusing properties  
„Hill's equation“*



*Example: particle motion with periodic coefficient*

*equation of motion:*  $x''(s) - k(s)x(s) = 0$

*restoring force  $\neq$  const,  
 $k(s)$  = depending on the position  $s$   
 $k(s+L) = k(s)$ , periodic function*

*we expect a kind of quasi harmonic oscillation: amplitude & phase will depend on the position  $s$  in the ring.*

**7.) The Beta Function**

**General solution of Hill's equation:**

(i)  $x(s) = \sqrt{\epsilon} \sqrt{\beta(s)} \cdot \cos(\psi(s) + \phi)$

$\epsilon, \Phi$  = integration constants determined by initial conditions

$\beta(s)$  periodic function given by focusing properties of the lattice  $\leftrightarrow$  quadrupoles

$$\beta(s + L) = \beta(s)$$

Inserting (i) into the equation of motion ...

$$\psi(s) = \int_0^s \frac{ds}{\beta(s)}$$

$\Psi(s)$  = „phase advance“ of the oscillation between point „0“ and „s“ in the lattice.

For one complete revolution: number of oscillations per turn „Tune“

$$Q_y = \frac{1}{2\pi} \oint \frac{ds}{\beta(s)}$$

### 8.) Beam Emittance and Phase Space Ellipse:

general solution of Hill equation

$$\left\{ \begin{array}{l} (1) \quad x(s) = \sqrt{\varepsilon} * \sqrt{\beta(s)} * \cos(\psi(s) + \phi) \\ (2) \quad x'(s) = -\frac{\sqrt{\varepsilon}}{\sqrt{\beta(s)}} * \{ \alpha(s) * \cos(\psi(s) + \phi) + \sin(\psi(s) + \phi) \} \end{array} \right.$$

from (1) we get

$$\cos(\psi(s) + \phi) = \frac{x(s)}{\sqrt{\varepsilon} * \sqrt{\beta(s)}}$$

Insert into (2) and solve for  $\varepsilon$

$$\alpha(s) = -\frac{1}{2} \beta'(s)$$

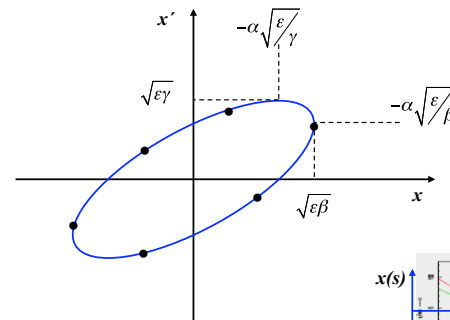
$$\gamma(s) = \frac{1 + \alpha(s)^2}{\beta(s)}$$

$$\varepsilon = \gamma(s) * x^2(s) + 2\alpha(s)x(s)x'(s) + \beta(s)x'(s)^2$$

- \*  $\varepsilon$  is a constant of the motion ... it is independent of „s“
- \* parametric representation of an ellipse in the  $x \ x'$  space
- \* shape and orientation of ellipse are given by  $\alpha, \beta, \gamma$

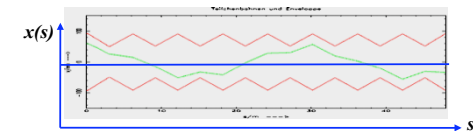
### Beam Emittance and Phase Space Ellipse

$$\varepsilon = \gamma(s) x^2(s) + 2\alpha(s)x(s)x'(s) + \beta(s) x'^2(s)$$



**Liouville: in reasonable storage rings area in phase space is constant.**

$$A = \pi * \varepsilon = \text{const}$$



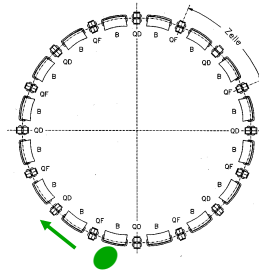
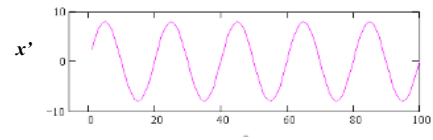
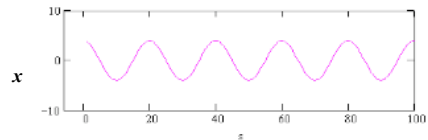
$\varepsilon$  beam emittance = **woozilycity** of the particle ensemble, **intrinsic beam parameter**, cannot be changed by the foc. properties.

**Scientifiquely speaking: area covered in transverse  $x, x'$  phase space ... and it is constant !!!**

### Particle Tracking in a Storage Ring

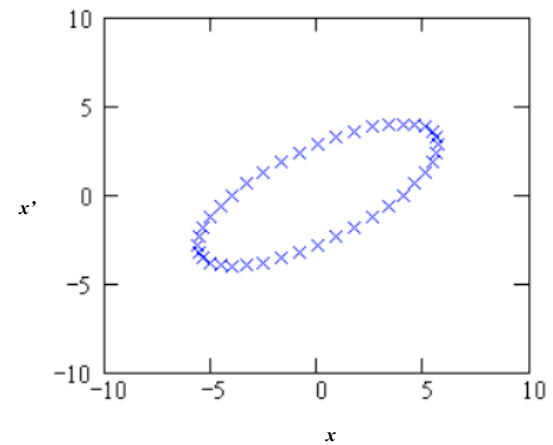
Calculate  $x$ ,  $x'$  for each linear accelerator element according to matrix formalism

plot  $x$ ,  $x'$  as a function of „ $s$ “



... and now the ellipse:

note for each turn  $x$ ,  $x'$  at a given position „ $s_1$ “ and plot in the phase space diagram



### Emittance & The Phase Space Ellipse

particle trajectory:  $x(s) = \sqrt{\varepsilon} \sqrt{\beta(s)} \cos\{\psi(s) + \phi\}$

max. Amplitude:  $\hat{x}(s) = \sqrt{\varepsilon\beta}$   $\rightarrow$   $x'$  at that position ...?

... put  $\hat{x}(s)$  into  $\varepsilon = \gamma(s) \cdot x^2(s) + 2\alpha(s)x(s)x'(s) + \beta(s)x'(s)^2$  and solve for  $x'$

$$\varepsilon = \gamma \cdot \varepsilon\beta + 2\alpha\sqrt{\varepsilon\beta} \cdot x' + \beta x'^2$$

$\rightarrow$   $x' = -\alpha \cdot \sqrt{\varepsilon / \beta}$

\* A high  $\beta$ -function means a large beam size and a small beam divergence. !  
... et vice versa !!!

\* In the middle of a quadrupole  $\beta$  is maximum,  $\alpha = \text{zero}$  }  $x' = 0$  ... and the ellipse is flat

### Phase Space Ellipse

$$\varepsilon = \gamma(s) x^2(s) + 2\alpha(s)x(s)x'(s) + \beta(s)x'^2(s)$$

$$\alpha(s) = -\frac{1}{2}\beta'(s)$$

$$\gamma(s) = \frac{1+\alpha(s)^2}{\beta(s)}$$

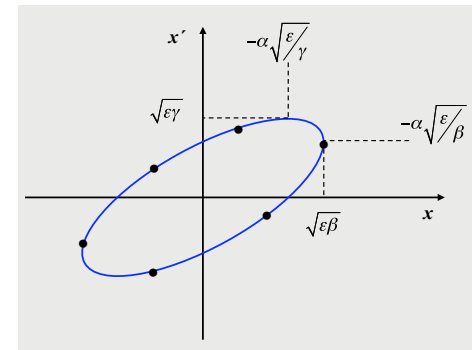
$\rightarrow$   $\varepsilon = \frac{x^2}{\beta} + \frac{\alpha^2 x^2}{\beta} + 2\alpha \cdot x x' + \beta \cdot x'^2$

... solve for  $x'$   $x'_{1,2} = \frac{-\alpha \cdot x \pm \sqrt{\varepsilon\beta - x^2}}{\beta}$

... and determine  $\hat{x}'$  via:  $\frac{dx'}{dx} = 0$

$\rightarrow$   $\hat{x}' = \sqrt{\varepsilon\gamma}$

$\rightarrow$   $\hat{x} = \pm\alpha \sqrt{\varepsilon/\gamma}$

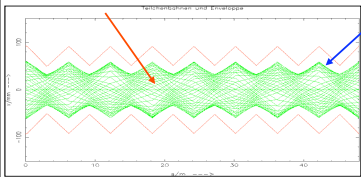


shape and orientation of the phase space ellipse depend on the Twiss parameters  $\beta$   $\alpha$   $\gamma$

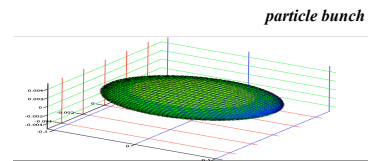
### Emittance of the Particle Ensemble:

$$x(s) = \sqrt{\varepsilon \beta(s)} \cdot \cos(\Psi(s) + \phi)$$

$$\dot{x}(s) = \sqrt{\varepsilon} \sqrt{\beta'(s)}$$



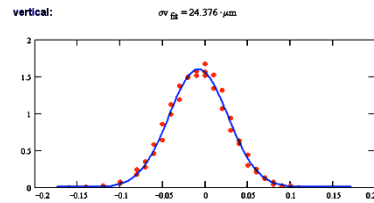
single particle trajectories,  $N \approx 10^{11}$  per bunch



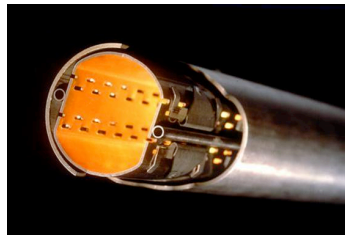
particle bunch

Gauß  
Particle Distribution:  $\rho(x) = \frac{N \cdot e}{\sqrt{2\pi}\sigma_x} \cdot e^{-\frac{1}{2}\frac{x^2}{\sigma_x^2}}$

particle at distance  $1\sigma$  from centre  $\leftrightarrow$  68.3 % of all beam particles

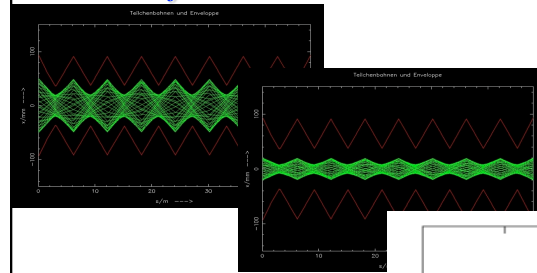


LHC:  $\sigma = \sqrt{\varepsilon \cdot \beta} = \sqrt{5 \cdot 10^{-10} \text{ m} \cdot 180 \text{ m}} = 0.3 \text{ mm}$



aperture requirements:  $r_0 \geq 10 \cdot \sigma$

### Emittance of the Particle Ensemble:

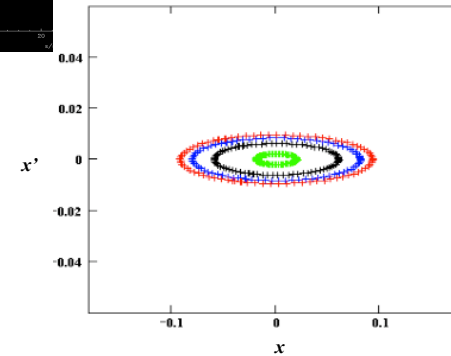


Example: LHC  
beam parameters in the arc

$$\beta(x) \approx 180 \text{ m}$$

$$\varepsilon \approx 5 \cdot 10^{-10} \text{ rad} \cdot \text{m} \quad (\leftrightarrow 1\sigma)$$

$$\sigma = \sqrt{\varepsilon \beta} \approx 0.3 \text{ mm}$$





**9.) Transfer Matrix M** ... yes we had the topic already

general solution of Hill's equation

$$\begin{cases} x(s) = \sqrt{\varepsilon} \sqrt{\beta(s)} \cos\{\psi(s) + \phi\} \\ x'(s) = \frac{-\sqrt{\varepsilon}}{\sqrt{\beta(s)}} [\alpha(s) \cos\{\psi(s) + \phi\} + \sin\{\psi(s) + \phi\}] \end{cases}$$

remember the trigonometrical gymnastics:  $\sin(a + b) = \dots$  etc

$$\begin{aligned} x(s) &= \sqrt{\varepsilon} \sqrt{\beta_s} (\cos\psi_s \cos\phi - \sin\psi_s \sin\phi) \\ x'(s) &= \frac{-\sqrt{\varepsilon}}{\sqrt{\beta_s}} [\alpha_s \cos\psi_s \cos\phi - \alpha_s \sin\psi_s \sin\phi + \sin\psi_s \cos\phi + \cos\psi_s \sin\phi] \end{aligned}$$

starting at point  $s(0) = s_0$ , where we put  $\Psi(0) = 0$

$$\left. \begin{aligned} \cos\phi &= \frac{x_0}{\sqrt{\varepsilon\beta_0}} \\ \sin\phi &= -\frac{1}{\sqrt{\varepsilon}} (x'_0 \sqrt{\beta_0} + \frac{\alpha_0 x_0}{\sqrt{\beta_0}}) \end{aligned} \right\} \text{inserting above ...}$$

$$\begin{aligned} x(s) &= \sqrt{\frac{\beta_s}{\beta_0}} \{ \cos\psi_s + \alpha_0 \sin\psi_s \} x_0 + \sqrt{\beta_s \beta_0} \sin\psi_s x'_0 \\ x'(s) &= \frac{1}{\sqrt{\beta_s \beta_0}} \{ (\alpha_0 - \alpha_s) \cos\psi_s - (1 + \alpha_0 \alpha_s) \sin\psi_s \} x_0 + \sqrt{\frac{\beta_0}{\beta_s}} \{ \cos\psi_s - \alpha_s \sin\psi_s \} x'_0 \end{aligned}$$

which can be expressed ... for convenience ... in matrix form  $\begin{pmatrix} x \\ x' \end{pmatrix}_s = M \begin{pmatrix} x \\ x' \end{pmatrix}_0$

$$M = \begin{pmatrix} \sqrt{\frac{\beta_s}{\beta_0}} (\cos\psi_s + \alpha_0 \sin\psi_s) & \sqrt{\beta_s \beta_0} \sin\psi_s \\ \frac{(\alpha_0 - \alpha_s) \cos\psi_s - (1 + \alpha_0 \alpha_s) \sin\psi_s}{\sqrt{\beta_s \beta_0}} & \sqrt{\frac{\beta_0}{\beta_s}} (\cos\psi_s - \alpha_s \sin\psi_s) \end{pmatrix}$$

- \* we can calculate the single particle trajectories between two locations in the ring, if we know the  $\alpha \beta \gamma$  at these positions.
- \* and nothing but the  $\alpha \beta \gamma$  at these positions.

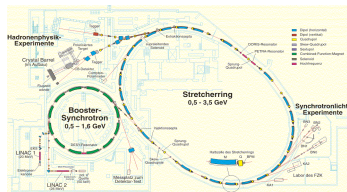
\* ... !

\* Äquivalenz der Matrizen

### 10.) Periodic Lattices

transfer matrix for particle trajectories as a function of the lattice parameters

$$M = \begin{pmatrix} \sqrt{\frac{\beta_s}{\beta_0}} (\cos\psi_s + \alpha_0 \sin\psi_s) & \sqrt{\beta_s \beta_0} \sin\psi_s \\ \frac{(\alpha_0 - \alpha_s) \cos\psi_s - (1 + \alpha_0 \alpha_s) \sin\psi_s}{\sqrt{\beta_s \beta_0}} & \sqrt{\frac{\beta_0}{\beta_s}} (\cos\psi_s - \alpha_s \sin\psi_s) \end{pmatrix}$$



Delta Electron Storage Ring

„This rather formidable looking matrix simplifies considerably if we consider one complete turn ...“

$$M(s) = \begin{pmatrix} \cos\psi_{turn} + \alpha_s \sin\psi_{turn} & \beta_s \sin\psi_{turn} \\ -\gamma_s \sin\psi_{turn} & \cos\psi_{turn} - \alpha_s \sin\psi_{turn} \end{pmatrix}$$

$$\psi_{turn} = \int_s^{s+L} \frac{ds}{\beta(s)} \quad \psi_{turn} = \text{phase advance per period}$$

Tune: Phase advance per turn in units of  $2\pi$

$$Q = \frac{1}{2\pi} * \oint \frac{ds}{\beta(s)}$$

### Stability Criterion:

Question: what will happen, if we do not make too many mistakes and your particle performs one complete turn ?



Matrix for 1 turn:

$$M = \begin{pmatrix} \cos\psi_{turn} + \alpha_s \sin\psi_{turn} & \beta_s \sin\psi_{turn} \\ -\gamma_s \sin\psi_{turn} & \cos\psi_{turn} - \alpha_s \sin\psi_{turn} \end{pmatrix} = \underbrace{\cos\psi}_{\mathbf{I}} \cdot \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} + \underbrace{\sin\psi}_{\mathbf{J}} \begin{pmatrix} \alpha & \beta \\ -\gamma & -\alpha \end{pmatrix}$$

Matrix for N turns:

$$M^N = (\mathbf{I} \cos\psi + \mathbf{J} \sin\psi)^N = \mathbf{I} \cos N\psi + \mathbf{J} \sin N\psi$$

The motion for N turns remains bounded, if the elements of  $M^N$  remain bounded

$$\psi = \text{real} \quad \Leftrightarrow \quad |\cos\psi| < 1 \quad \Leftrightarrow \quad |\text{Trace}(M)| < 2$$

stability criterion .... proof for the disbelieving colleagues !!

$$\text{Matrix for 1 turn: } M = \begin{pmatrix} \cos\psi_{turn} + \alpha_s \sin\psi_{turn} & \beta_s \sin\psi_{turn} \\ -\gamma_s \sin\psi_{turn} & \cos\psi_{turn} - \alpha_s \sin\psi_{turn} \end{pmatrix} = \underbrace{\cos\psi}_{\mathbf{I}} \cdot \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} + \sin\psi \underbrace{\begin{pmatrix} \alpha & \beta \\ -\gamma & -\alpha \end{pmatrix}}_{\mathbf{J}}$$

Matrix for 2 turns:

$$M^2 = (\mathbf{I} * \cos\psi_1 + \mathbf{J} * \sin\psi_1) * (\mathbf{I} * \cos\psi_2 + \mathbf{J} * \sin\psi_2) \\ = \mathbf{I}^2 * \cos\psi_1 \cos\psi_2 + \mathbf{I}\mathbf{J} * \cos\psi_1 \sin\psi_2 + \mathbf{J}\mathbf{I} * \sin\psi_1 \cos\psi_2 + \mathbf{J}^2 \sin\psi_1 \sin\psi_2$$

now ...

$$\mathbf{I}^2 = \mathbf{I}$$

$$\left. \begin{aligned} \mathbf{I} * \mathbf{J} &= \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} * \begin{pmatrix} \alpha & \beta \\ -\gamma & -\alpha \end{pmatrix} = \begin{pmatrix} \alpha & \beta \\ -\gamma & -\alpha \end{pmatrix} \\ \mathbf{J} * \mathbf{I} &= \begin{pmatrix} \alpha & \beta \\ -\gamma & -\alpha \end{pmatrix} * \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} = \begin{pmatrix} \alpha & \beta \\ -\gamma & -\alpha \end{pmatrix} \end{aligned} \right\} \mathbf{I} * \mathbf{J} = \mathbf{J} * \mathbf{I}$$

$$\mathbf{J}^2 = \begin{pmatrix} \alpha & \beta \\ -\gamma & -\alpha \end{pmatrix} * \begin{pmatrix} \alpha & \beta \\ -\gamma & -\alpha \end{pmatrix} = \begin{pmatrix} \alpha^2 - \gamma\beta & \alpha\beta - \beta\alpha \\ -\gamma\alpha + \alpha\gamma & \alpha^2 - \gamma\beta \end{pmatrix} = \begin{pmatrix} -1 & 0 \\ 0 & -1 \end{pmatrix} = -\mathbf{I}$$

$$M^2 = \mathbf{I} * \cos(\psi_1 + \psi_2) + \mathbf{J} * \sin(\psi_1 + \psi_2)$$

$$M^2 = \mathbf{I} * \cos(2\psi) + \mathbf{J} * \sin(2\psi)$$

## 11.) Transformation of $\alpha, \beta, \gamma$

consider two positions in the storage ring:  $s_0, s$

$$\begin{pmatrix} x \\ x' \end{pmatrix}_s = M \cdot \begin{pmatrix} x \\ x' \end{pmatrix}_{s_0}$$

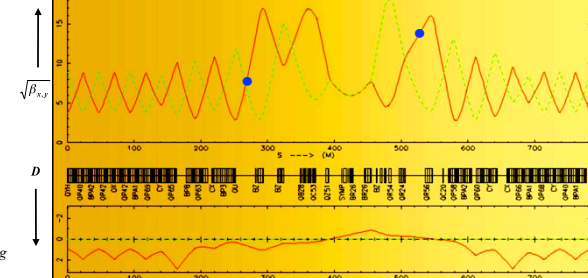
where ...  $M = M_{QF} \cdot M_{QD} \cdot M_B \cdot M_{Drift} \cdot M_{QF} \cdot \dots$

$$M = \begin{pmatrix} C & S \\ C' & S' \end{pmatrix}$$

since  $\varepsilon = \text{const}$  (Liouville):

$$\varepsilon = \beta x'^2 + 2\alpha x x' + \gamma x^2$$

$$\varepsilon = \beta_0 x_0'^2 + 2\alpha_0 x_0 x_0' + \gamma_0 x_0^2$$



express  $x_0, x'_0$  as a function of  $x, x'$ .

$$\begin{pmatrix} x \\ x' \end{pmatrix}_s = M \cdot \begin{pmatrix} x \\ x' \end{pmatrix}_{s_0}$$

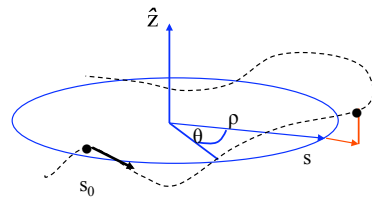
... remember  $W = CS - SC' - I$

$$\begin{pmatrix} x \\ x' \end{pmatrix}_0 = M^{-1} \cdot \begin{pmatrix} x \\ x' \end{pmatrix}_s$$

$$M^{-1} = \begin{pmatrix} S' & -S \\ -C' & C \end{pmatrix}$$



$$\begin{aligned} x_0 &= S'x - Sx' \\ x'_0 &= -C'x + Cx' \end{aligned}$$



inserting into  $\epsilon$

$$\epsilon = \beta x'^2 + 2\alpha xx' + \gamma x^2$$

$$\epsilon = \beta_0(Cx' - C'x)^2 + 2\alpha_0(S'x - Sx')(Cx' - C'x) + \gamma_0(S'x - Sx')^2$$

sort via  $x, x'$  and compare the coefficients to get ....

$$\beta(s) = C^2\beta_0 - 2SC\alpha_0 + S^2\gamma_0$$

$$\alpha(s) = -CC'\beta_0 + (SC' + S'C)\alpha_0 - SS'\gamma_0$$

$$\gamma(s) = C'^2\beta_0 - 2S'C'\alpha_0 + S'^2\gamma_0$$

in matrix notation:

$$\begin{pmatrix} \beta \\ \alpha \\ \gamma \end{pmatrix}_s = \begin{pmatrix} C^2 & -2SC & S^2 \\ -CC' & SC' + CS' & -SS' \\ C'^2 & -2S'C' & S'^2 \end{pmatrix} \begin{pmatrix} \beta_0 \\ \alpha_0 \\ \gamma_0 \end{pmatrix}$$



- 1.) *this expression is important*
- 2.) *given the twiss parameters  $\alpha, \beta, \gamma$  at any point in the lattice we can transform them and calculate their values at any other point in the ring.*
- 3.) *the transfer matrix is given by the focusing properties of the lattice elements, the elements of  $M$  are just those that we used to calculate single particle trajectories.*
- 4.) *go back to point 1.)*

### Résumé:

*equation of motion:*  $x''(s) + K(s)x(s) = 0$  ,  $K = 1/\rho^2 - k$

*general solution of Hill's equation:*  $x(s) = \sqrt{\varepsilon} \sqrt{\beta(s)} \cdot \cos(\psi(s) + \phi)$

*phase advance & tune:*  $\psi_{12}(s) = \int_{s_1}^{s_2} \frac{1}{\beta(s)} ds$  ,  $Q(s) = \frac{1}{2\pi} \oint \frac{1}{\beta(s)} ds$

*emittance:*  $\varepsilon = \gamma(s) x^2(s) + 2\alpha(s)x(s)x'(s) + \beta(s) x'^2(s)$

*transfer matrix from  $s_1 \rightarrow s_2$ :* 
$$M_{1 \rightarrow 2} = \begin{pmatrix} \sqrt{\frac{\beta_2}{\beta_1}} (\cos\psi_{21} + \alpha_1 \sin\psi_{21}) & \sqrt{\beta_2 \beta_1} \sin\psi_{21} \\ \frac{(\alpha_1 - \alpha_2) \cos\psi_{21} - (1 + \alpha_1 \alpha_2) \sin\psi_{21}}{\sqrt{\beta_2 \beta_1}} & \sqrt{\frac{\beta_1}{\beta_2}} (\cos\psi_{21} - \alpha_2 \sin\psi_{21}) \end{pmatrix}$$

*matrix for 1 turn:* 
$$M(s) = \begin{pmatrix} \cos\psi_{turn} + \alpha_s \sin\psi_{turn} & \beta_s \sin\psi_{turn} \\ -\gamma_s \sin\psi_{turn} & \cos\psi_{turn} - \alpha_s \sin\psi_{turn} \end{pmatrix}$$

*stability criterion:*  $|\text{Trace}(M)| < 2$