

# Superconducting magnets for

# Accelerators

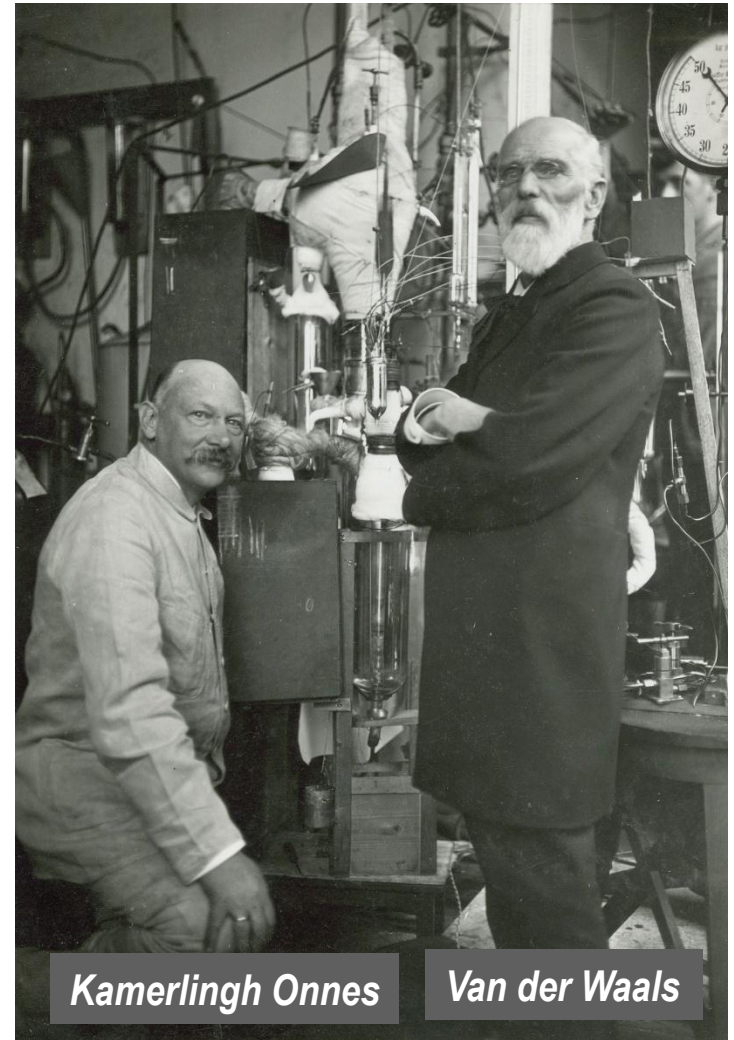
Martin N Wilson (Rutherford Lab → Oxford Instruments → CERN → OI → cons)

101 years of  
superconductivity

1908: Heike Kamerlingh  
Onnes liquefies helium

1911: HKO finds  
superconductivity

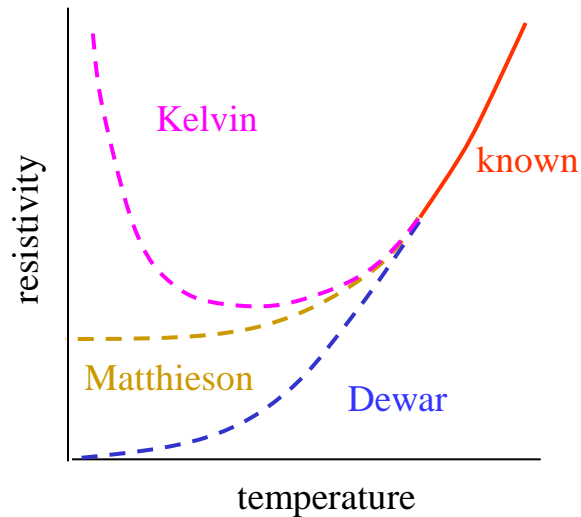
new science enabled by  
technology development



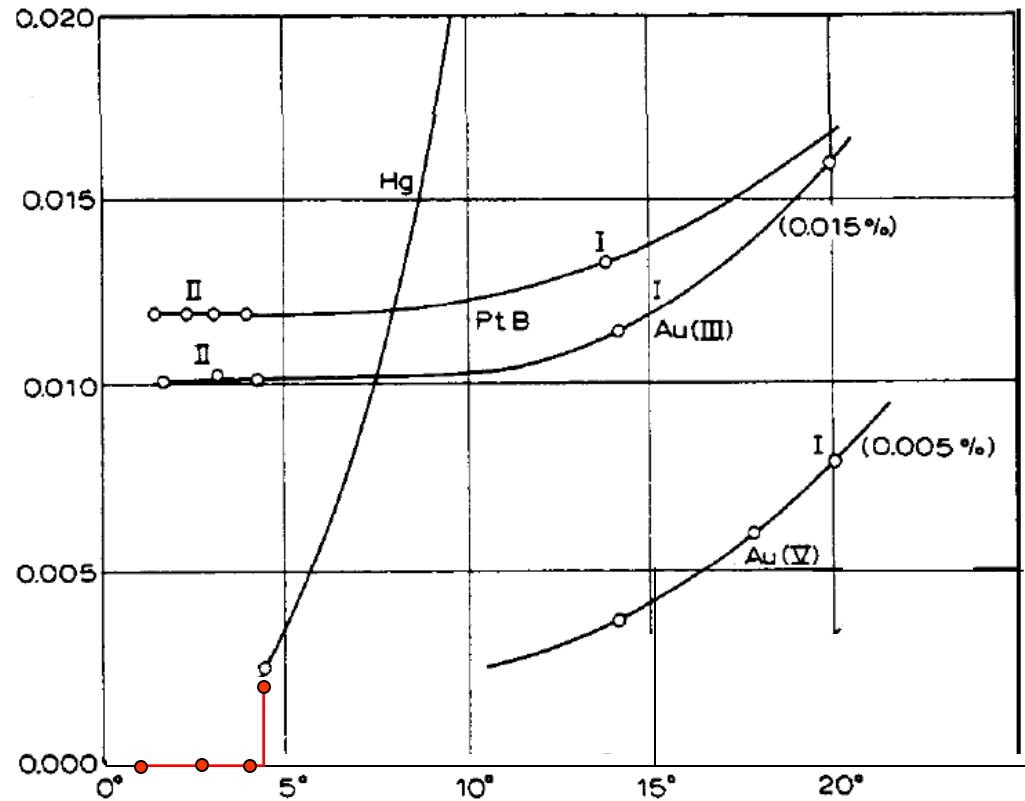
Kamerlingh Onnes

Van der Waals

# Resistivity at low temperature



- very different predictions of what might happen



- need high purity to test different theories
- can make Hg very pure by multiple distillation
- *but nobody expected this!*

*'Thus the mercury at 4.2 has entered a new state which ..... can be called the state of superconductivity'* (HKO Nobel lecture)

# How low is the resistance?

coil of lead wire, with shorted terminals

impose magnetic field when warm

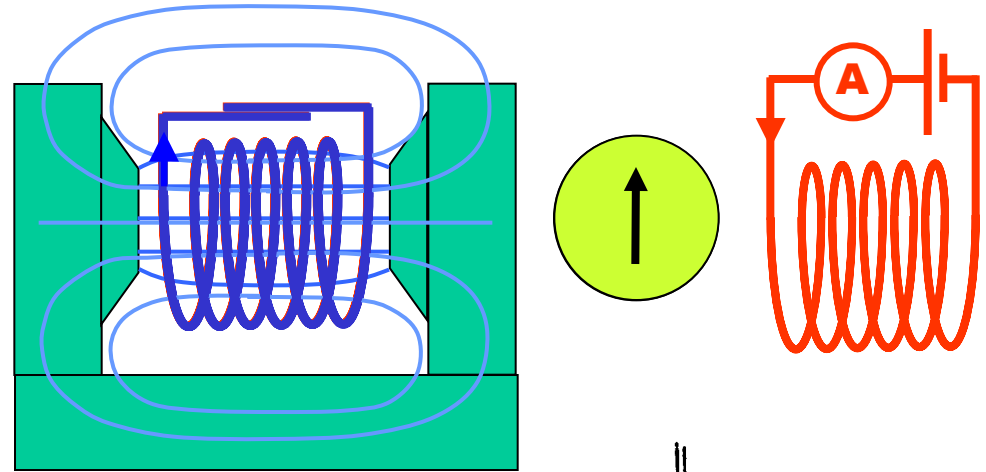
cool the coil

remove field - induces current

measure field from current

back off with a resistive coil

no change for hours - *persistent current*

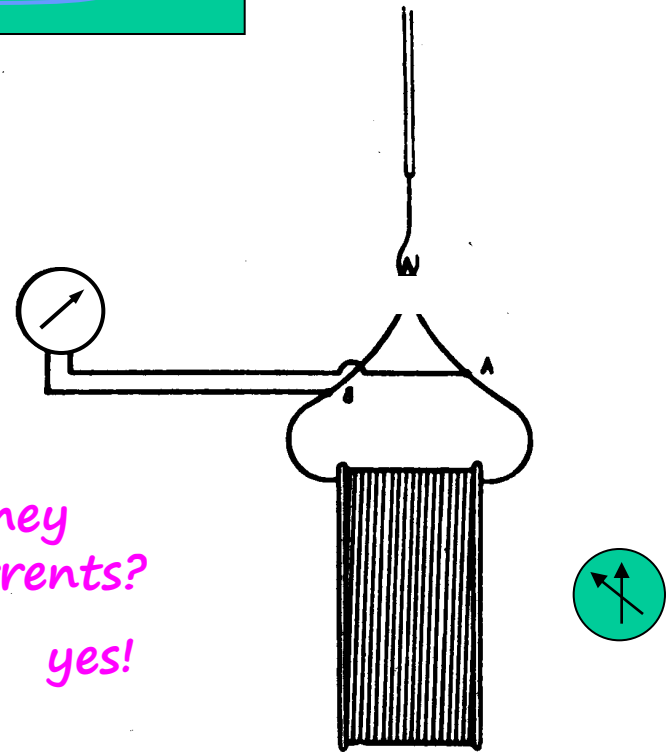


'It is uncanny to see ..... You can feel, almost tangibly how the ring of electrons in the wire turns around, around, around - slowly and almost without friction'

P Ehrenfest

but are they really currents?

yes!



# Magnets

' ..... bearing on the problem of producing intense magnetic field ..... a great number of Ampere windings can be located in a very small space without ..... heat being developed.....'

*Communication from the Physical Laboratory University of Leiden Sept 1913*

' ..... 100,000 Gauss could then be obtained by a coil of say 30 centimetres in diameter and the cooling with helium would require a plant which could be realized in Leiden with a relatively modest support.....'

*Third International Congress of Refrigeration, Chicago Sept 1913*

' ..... In field above this threshold value, a relatively large magnetic resistance arises at once.....'

'Thus an unexpected difficulty ..... faced us. The discovery of the strange property which causes this made up for the difficulties involved.'

*Nobel Prize Acceptance Lecture, Stockholm Dec 1913*

# *Superconductivity for accelerators - why bother?*

## *Abolish Ohm's Law*

- no power consumption  
(although do need refrigeration power)
- high current density  $\Rightarrow$  compact windings, high gradients
- ampere turns are cheap, so don't need iron  
(although often use it for shielding)

## *Consequences*

- lower power bills
- higher magnetic fields mean reduced bend radius
  - $\Rightarrow$  smaller rings
  - $\Rightarrow$  reduced capital cost
  - $\Rightarrow$  new technical possibilities  
(eg muon collider)
- higher quadrupole gradients
  - $\Rightarrow$  higher luminosity



# Plan of the Lectures

## 1 Introduction to Superconductors

- critical field, temperature & current
- superconductors for magnets
- manufacture of superconducting wires
- high temperature superconductors HTS

## 2 Magnets, 'Training' & Fine Filaments

- coil shapes for solenoids, dipoles & quadrupoles
- engineering current density & load lines
- degradation & training minimum quench energy
- critical state model & fine filaments

## 3 Magnetization, Cables & AC losses

- filamentary superconductors and magnetization
- coupling between filaments  $\Rightarrow$  magnetization
- why cables, coupling in cables
- AC losses in changing fields

## 4 Quenching and Cryogenics

- the quench process
- resistance growth, current decay, temperature rise
- quench protection schemes
- cryogenic fluids, refrigeration, cryostat design

## 5 Practical Matters

- LHC quench protection
- current leads
- accelerator magnet manufacture
- some superconducting accelerators

### Tutorial 1: Fine Filaments

- how filament size affects magnetization

### Tutorial 2: Quenching

- current decay and temperature rise

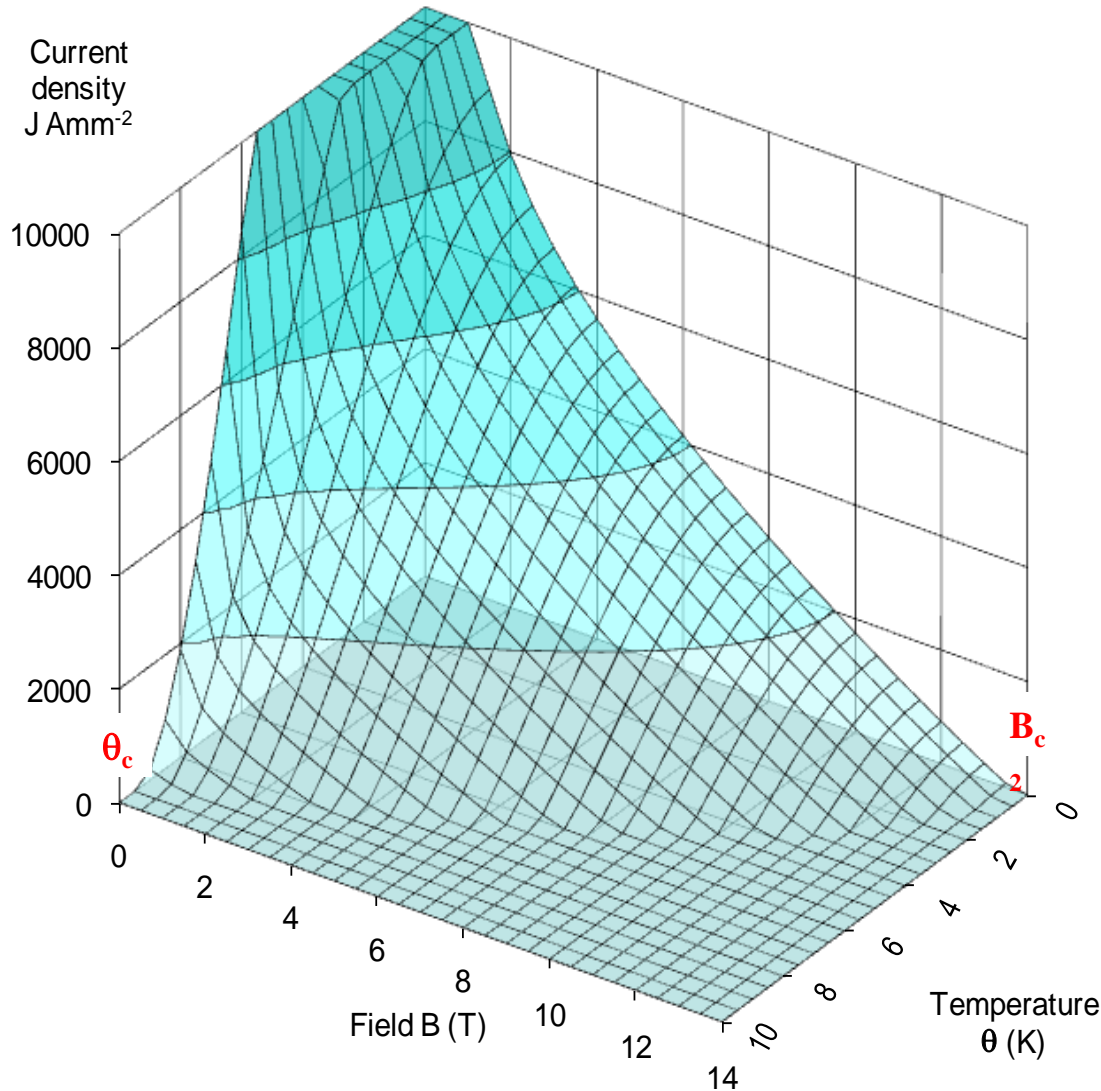


*get a feel for the numbers,  
bring a calculator*



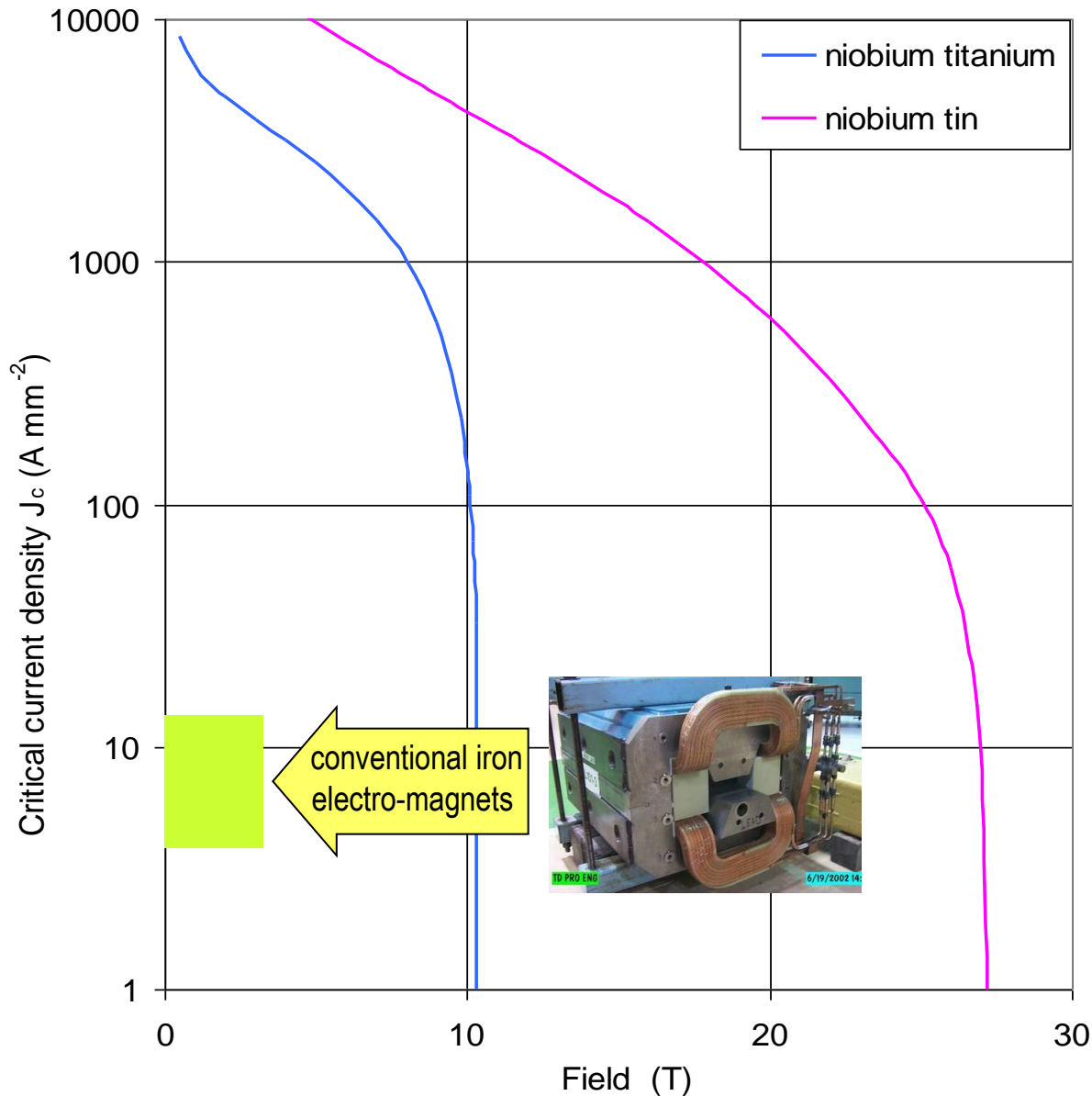
# The critical surface of niobium titanium

- **NbTi** is the standard commercial ‘work horse’ of the superconducting magnet business



- **critical surface** is the boundary between superconductivity and normal resistivity in  $J, B, \theta$  space
- superconductivity prevails everywhere below the surface, resistance everywhere above it
- upper critical field  $B_{c2}$  (at zero temperature and current)
- critical temperature  $\theta_c$  (at zero field and current)
- $B_{c2}$  and  $\theta_c$  are characteristic of the alloy composition
- critical current density  $J_c$  depends on processing

# The critical line at 4.2K

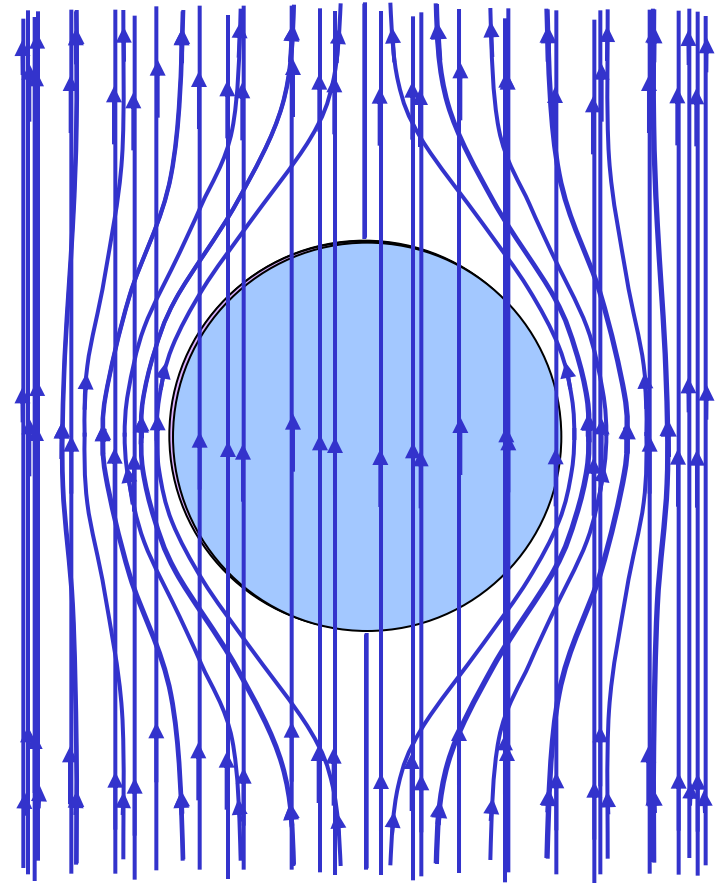


- magnets usually work in boiling liquid helium, so the critical surface is often represented by a curve of current versus field at 4.2K
- niobium tin Nb<sub>3</sub>Sn has a much higher performance than NbTi
- **but** Nb<sub>3</sub>Sn is a brittle intermetallic compound with poor mechanical properties
- both the field and current density of both superconductors are way above the capability of conventional electromagnets



# Two kinds of superconductor: type 1

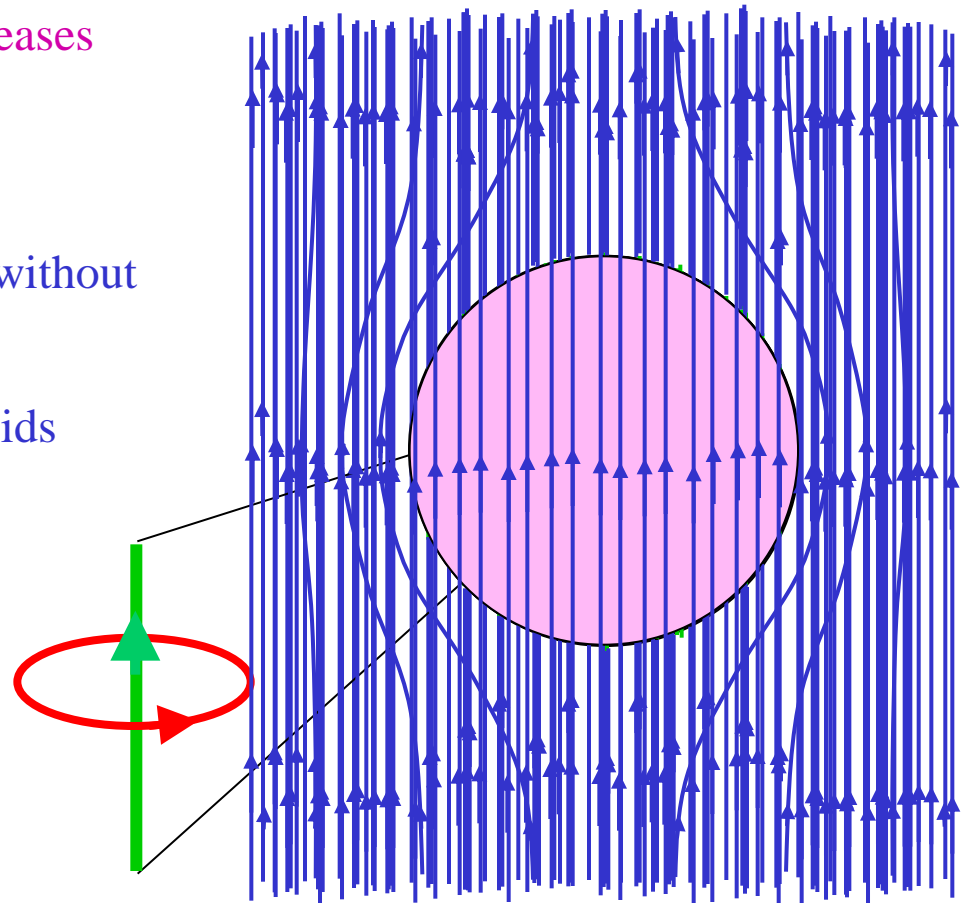
- the materials first discovered by Kammerlingh Onnes in 1911 - soft metals like lead, tin mercury
- sphere of metal at room temperature
- **apply magnetic field**
- **reduce the temperature - resistance decreases**
- **reduce the temperature some more - resistance decreases some more**
- at the critical temperature  $\theta_c$  the field is pushed out - the **Meissner effect** - superconductivity!
- increase the field - field is kept out
- increase the field some more - superconductivity is extinguished and the field jumps in
- thermodynamic critical field  $B_c$  is trade off between **reducing** energy via condensation to superconductivity and **increasing** energy by pushing out field  $\sim 0.1T$



*useless for magnets!*

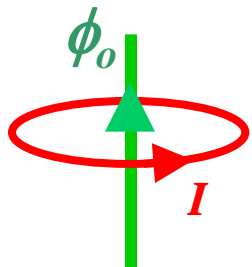
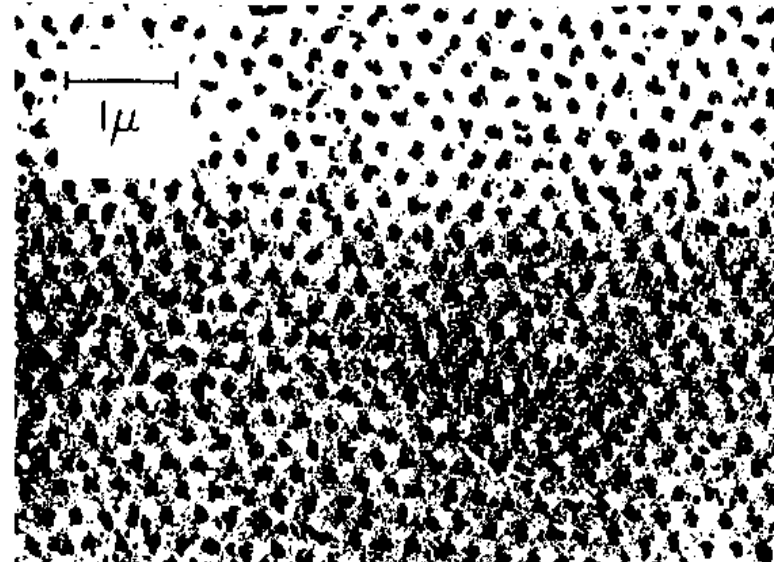
# Two kinds of superconductor: type 2

- apply magnetic field
- reduce the temperature - resistance decreases
- at the critical temperature  $\theta_c$  the field is pushed out
- increase the field - field jumps back in without quenching superconductivity
- it does so in the form of quantized fluxoids
- lower critical field  $B_{c1}$
- supercurrents encircle the resistive core of the fluxoid thereby screening field from the bulk material
- higher field  $\Rightarrow$  closer vortex spacing
- superconductivity is extinguished at the (much higher) upper critical field  $B_{c2}$



# Critical field: type 2 superconductors

- Meissner effect is not total, the magnetic field actually penetrates a small distance  $\lambda$  the **London Penetration Depth**.
- another characteristic distance is the **coherence length**  $\xi$  - the minimum distance over which the electronic state can change from superconducting to normal
- theory of Ginsburg, Landau, Abrikosov and Gorkov **GLAG** defines the ratio  $\kappa = \lambda / \xi$
- if  $\kappa > 1/\sqrt{2}$  material is **Type 2**
- magnetic field penetrates as discrete **fluxoids**



a single fluxoid encloses flux

$$\phi_0 = \frac{h}{2e} = 2 \times 10^{-15} \text{ Webers}$$

where  $h$  = Planck's constant,  
 $e$  = electronic charge

**upper critical field**

$$B_{c2} = \sqrt{2} \kappa B_c$$

in the 'dirty limit'  $\kappa \approx 2.4 \times 10^6 \gamma^{\frac{1}{2}} \rho_n$

where  $\rho_n$  is the normal state resistivity  
**- best superconductors are best resistors!**

thus the upper critical field

$$B_{c2} = 3.1 \times 10^3 \gamma \rho_n \theta_c$$

for NbTi:

$$\gamma \sim 900 \text{ J m}^{-3} \text{ K}^{-2}$$

$$\rho_n \sim 65 \times 10^{-8} \text{ W m}$$

$$\theta_c = 9.3 \text{ K}$$

hence  $B_{c2} \sim 18.5 \text{ T}$

Sommerfeld coefficient of electronic specific heat  $C_e = \gamma \theta$

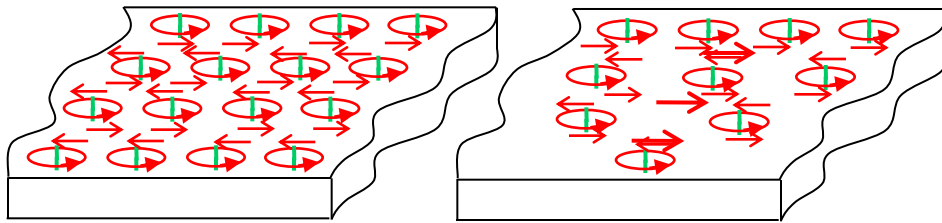
# Critical current density: type 2 superconductors

- fluxoids consist of resistive cores with supercurrents circulating round them.

spacing between the fluxoids is:-

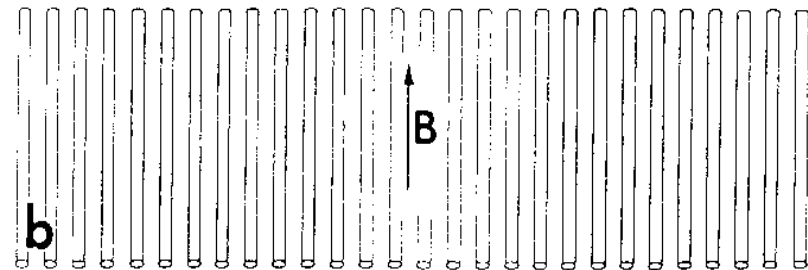
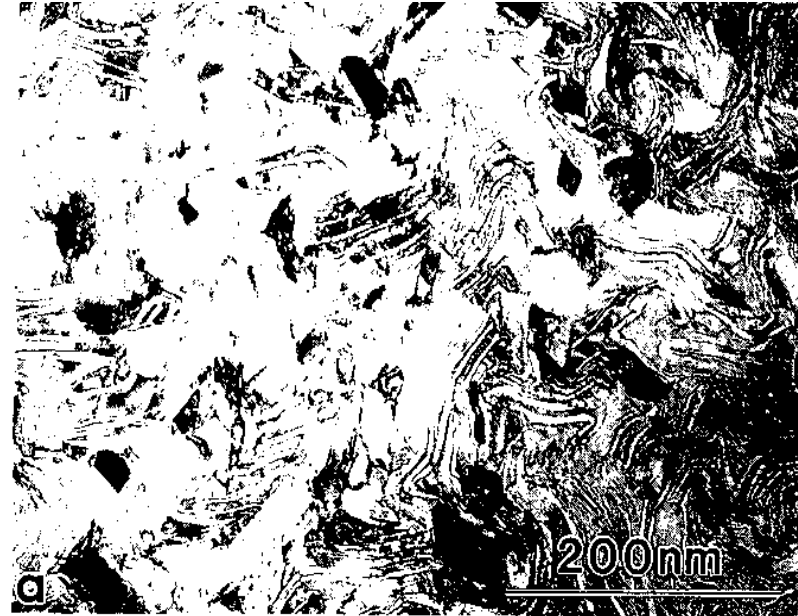
$$d = \left\{ \frac{2 \phi_o}{\sqrt{3} B} \right\}^{\frac{1}{2}} = 22nm \quad \text{at } 5T$$

- each fluxoid carries one unit of flux, so density of fluxoids = average field  
uniform density  $\Rightarrow$  uniform field  
 $\Rightarrow$  zero  $J$  (because  $Curl B = \mu_o J$ )
- to get a current density we must produce a **gradient** in the density of fluxoids



- fluxoids like to distribute uniformly
- so we must impose a gradient by inhomogeneities in the material, eg dislocations or precipitates

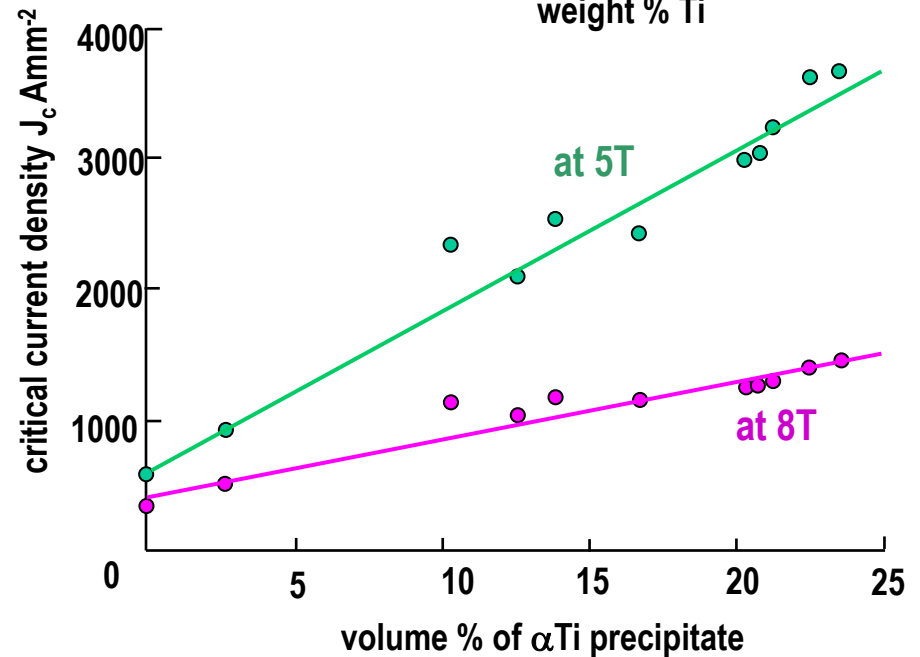
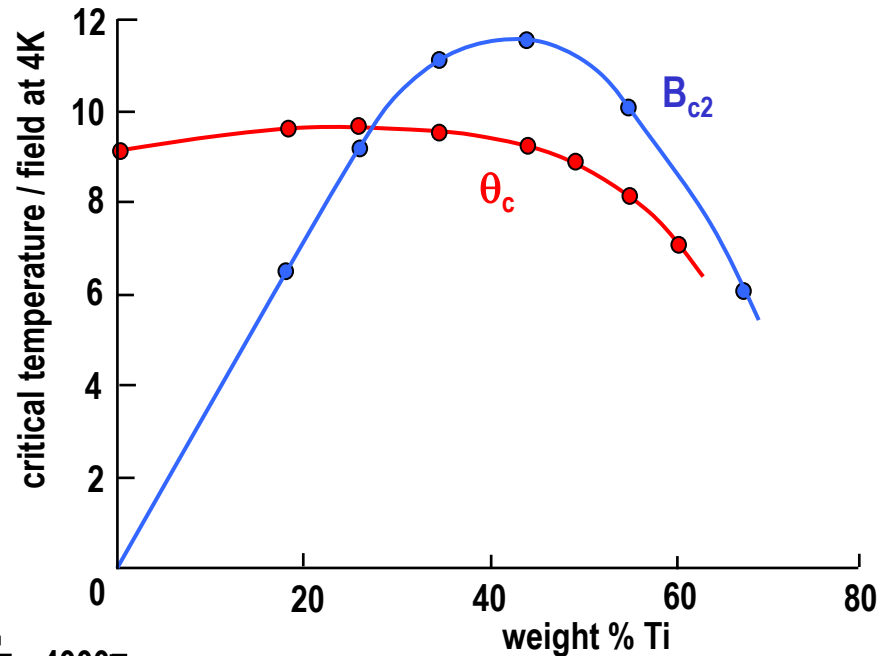
precipitates of  $\alpha$  Ti in Nb Ti



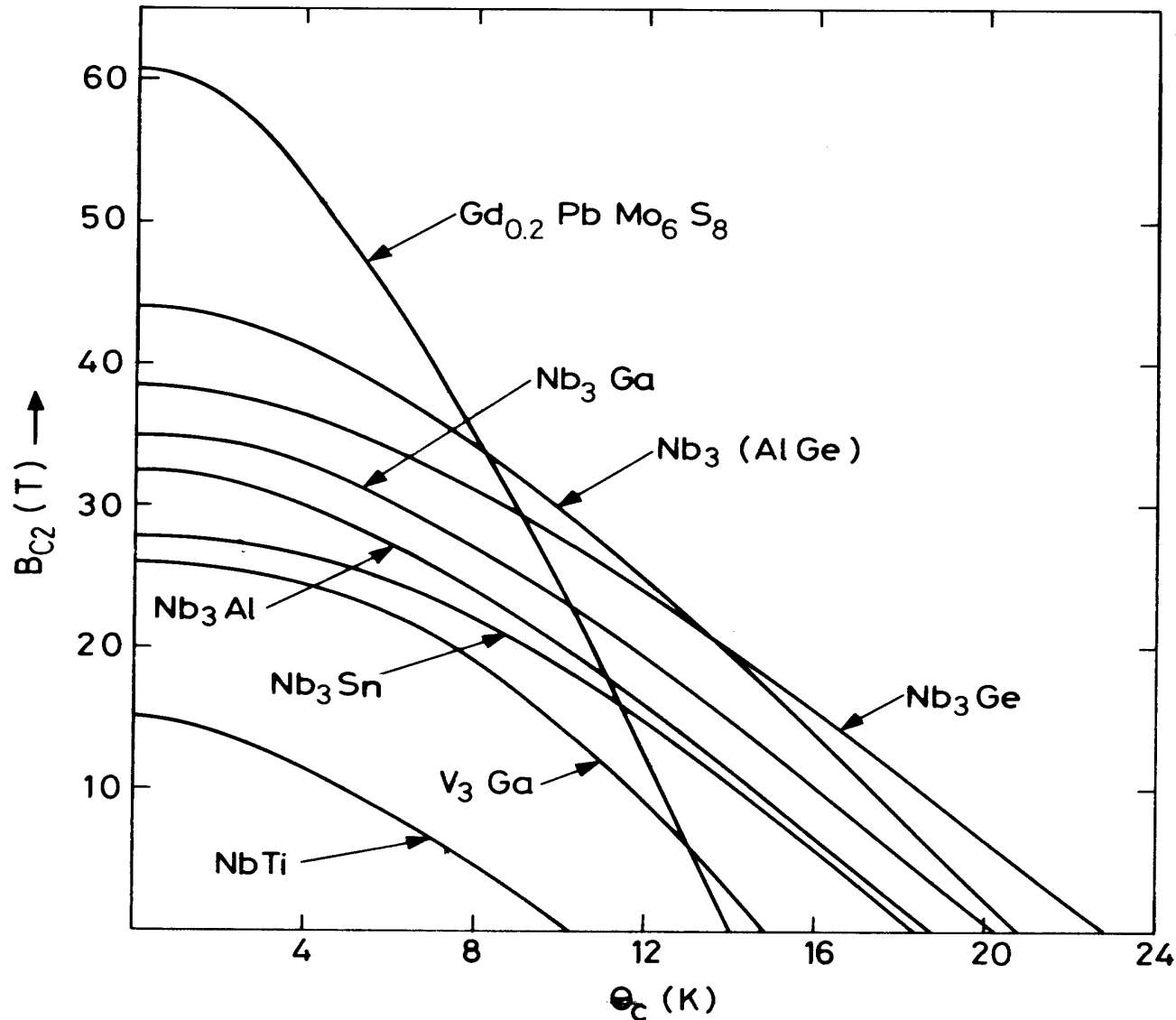
fluxoid lattice at 5T on the same scale

# Critical properties

- **Critical temperature  $\theta_c$** : choose the right material to have a large energy gap or 'depairing energy'  
*property of the material*
- **Upper Critical field  $B_{c2}$** : choose a Type 2 superconductor with a high critical temperature and a high normal state resistivity  
*property of the material*
- **Critical current density  $J_c$** : mess up the microstructure by cold working and precipitation heat treatments  
*hard work by the producer*



# Critical field & temperature of metallic superconductors

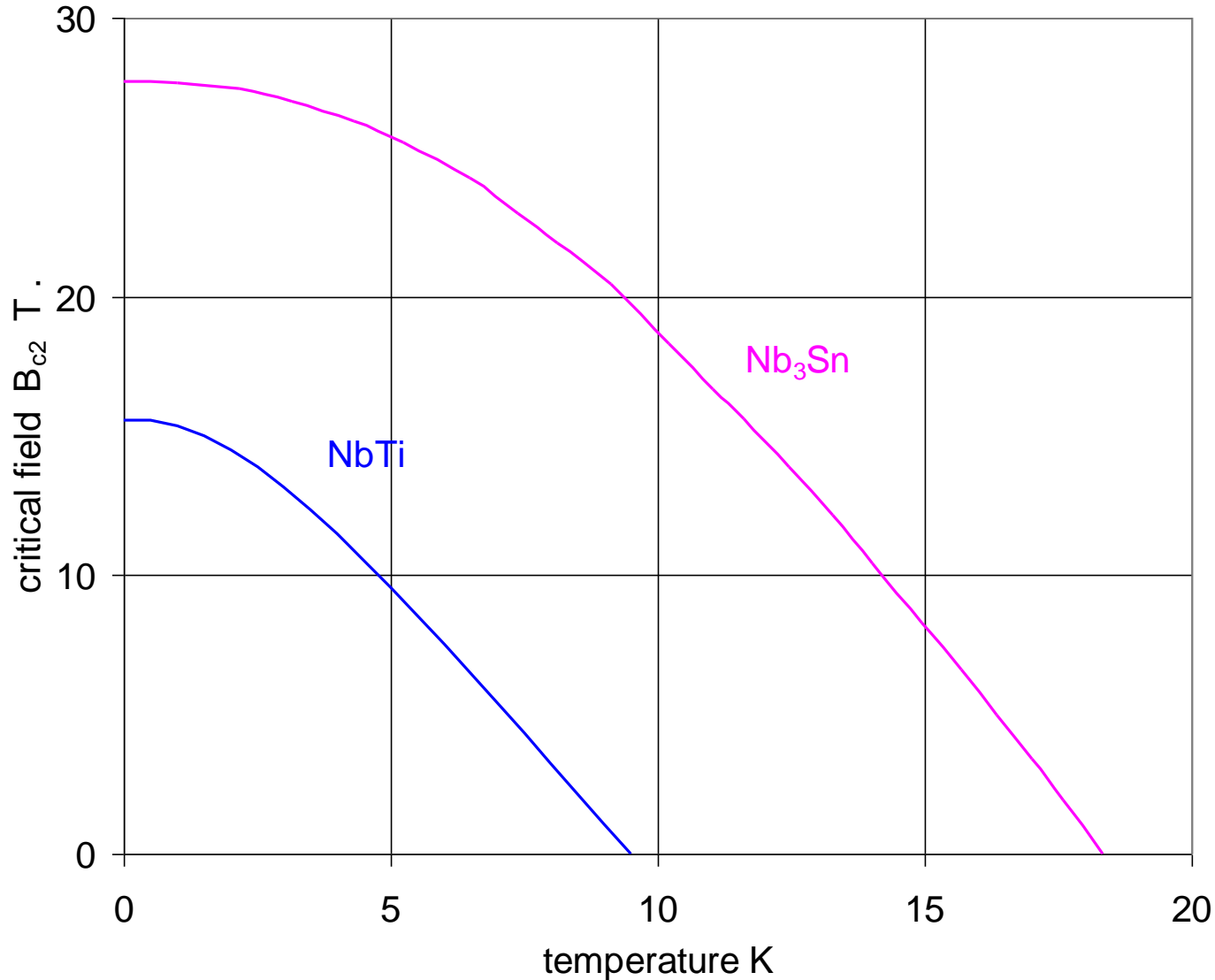


Note: of all the metallic superconductors, only NbTi is ductile.

All the rest are brittle intermetallic compounds



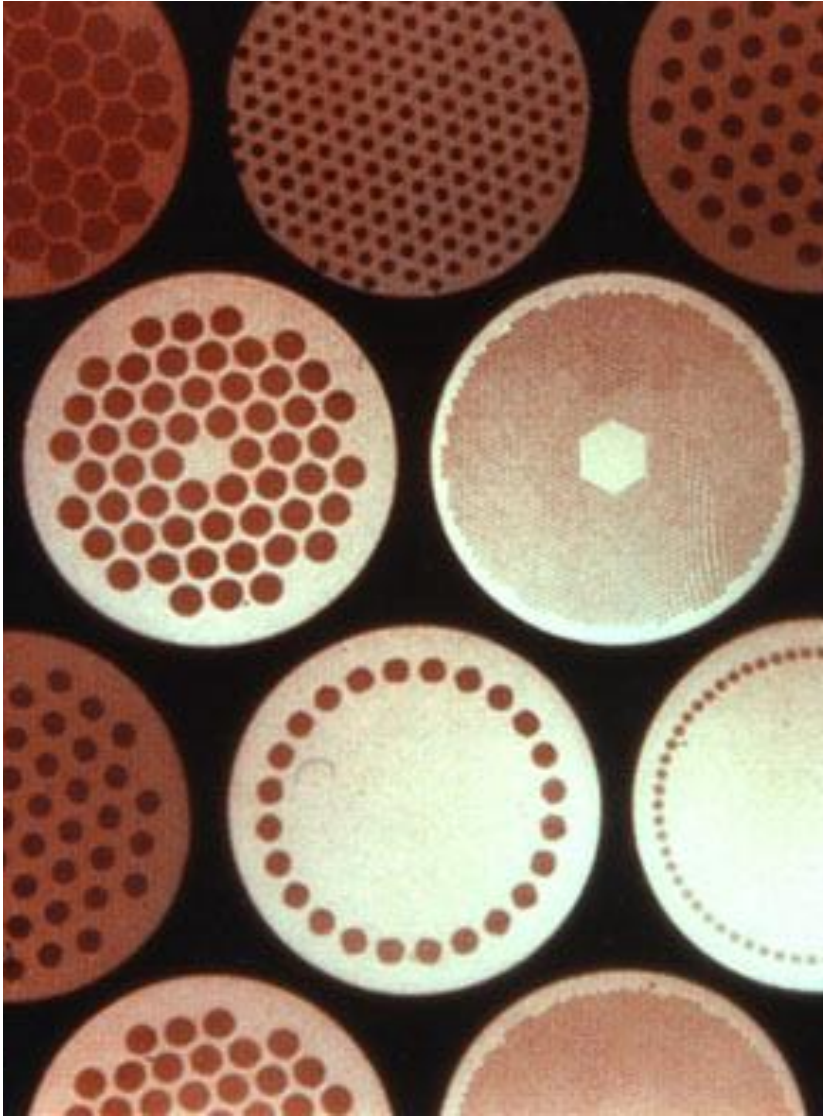
# Critical field & temperature of metallic superconductors



To date, all superconducting accelerators have used NbTi.

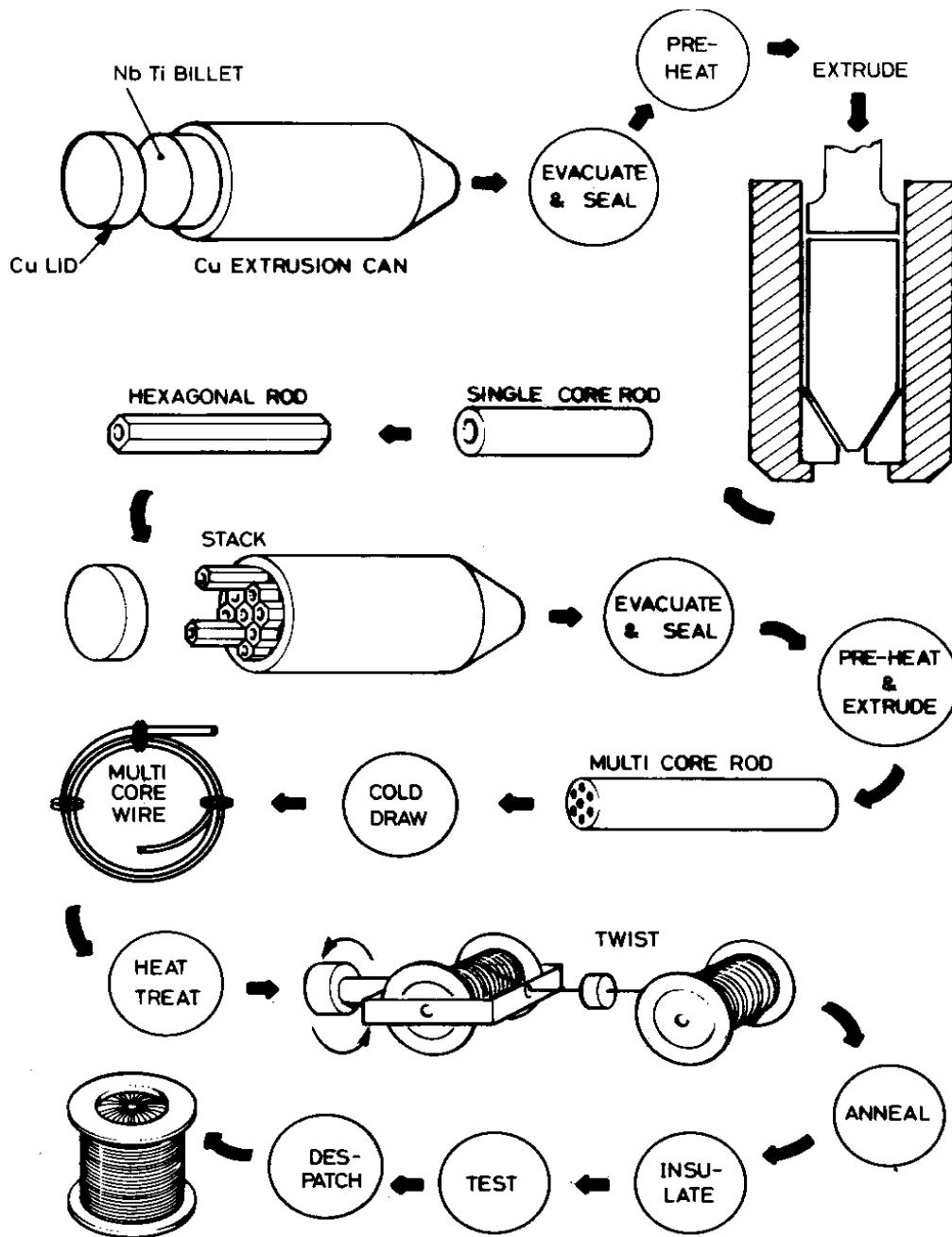
Of the intermetallics, only Nb<sub>3</sub>Sn has found significant use in magnets

# Practical wires for magnets



- some 40 years after its development, NbTi is still the most popular magnet conductor, with Nb<sub>3</sub>Sn being used for special high field magnets and HTS for some developmental prototypes.
- for reasons that will be described later, superconducting materials are always used in combination with a good normal conductor such as copper
- to ensure intimate mixing between the two, the superconductor is made in the form of fine filaments embedded in a matrix of copper
- typical dimensions are:
  - wire diameter = 0.3 - 1.0mm
  - filament diameter = 10 - 60μm
- for electromagnetic reasons, the composite wires are twisted so that the filaments look like a rope (see Lecture 3)

# NbTi manufacture



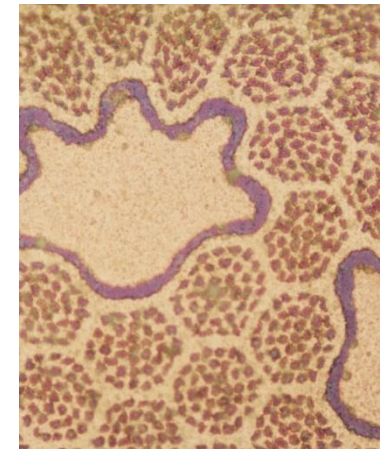
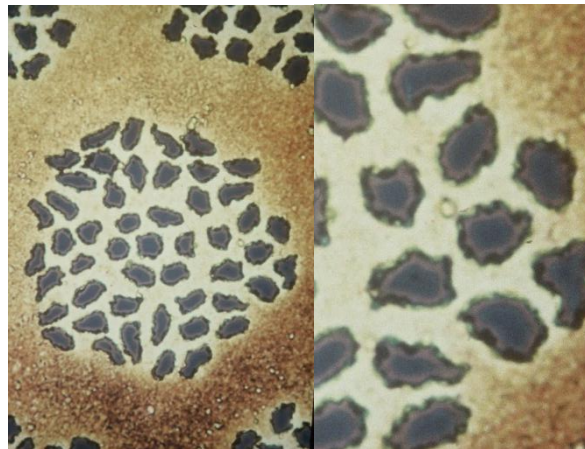
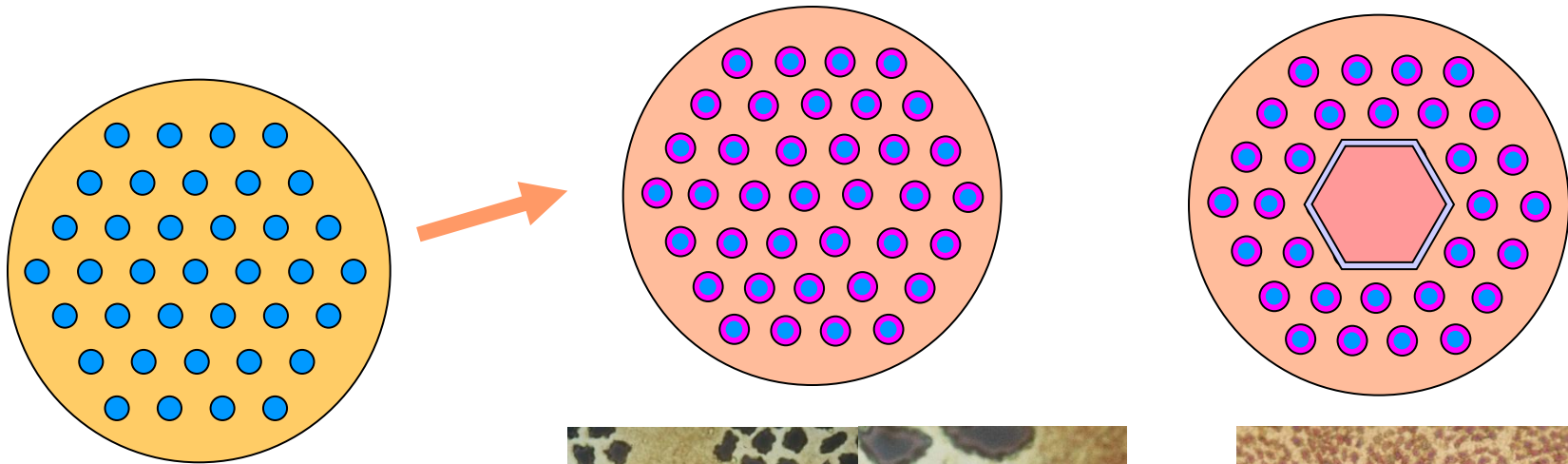
- vacuum melting of NbTi billets
- hot extrusion of the copper NbTi composite
- sequence of cold drawing and intermediate heat treatments to precipitate  $\alpha$ Ti phases as flux pinning centres
- for very fine filaments, must avoid the formation of brittle CuTi intermetallic compounds during heat treatment
  - usually done by enclosing the NbTi in a thin Nb shell
- twisting to avoid coupling - see lecture 2

# Filamentary $Nb_3Sn$ wire via the bronze route

$Nb_3Sn$  is a brittle material and cannot be drawn down. Instead must draw down pure niobium in a matrix of bronze (copper tin)

At final size the wire is heated ( $\sim 700^\circ C$  for some days) tin diffuses through the Cu and reacts with the Nb to form  $Nb_3Sn$

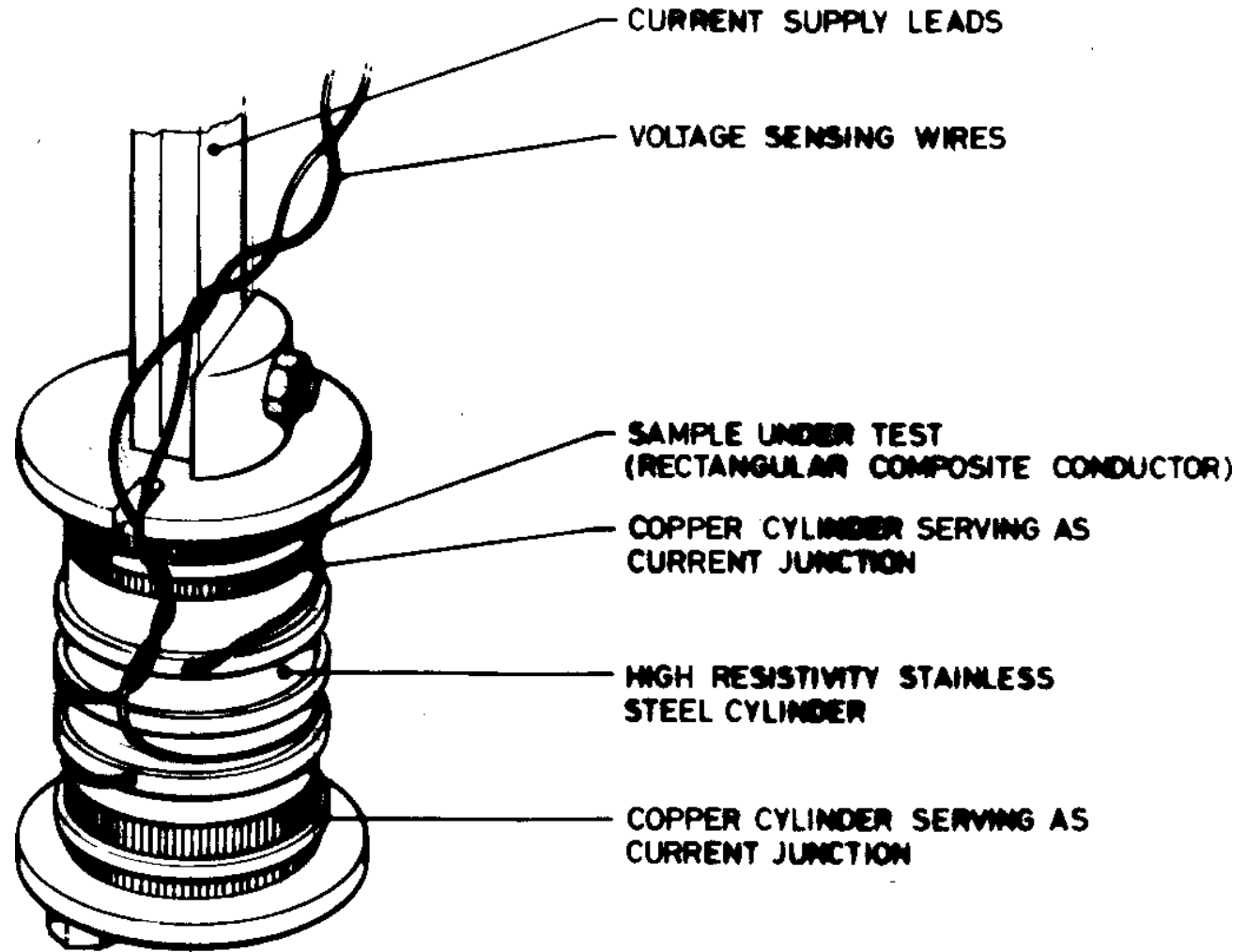
The remaining copper still contains  $\sim 3wt\%$  tin and has a high resistivity  $\sim 6 \times 10^{-8} \Omega m$ . So include 'islands' of pure copper surrounded by a diffusion barrier



- *BUT maximum ductile bronze is  $\sim 13wt\%$  tin,*
- *reaction slows at  $\sim 3wt\%$*
- *so low engineering  $J_c$*

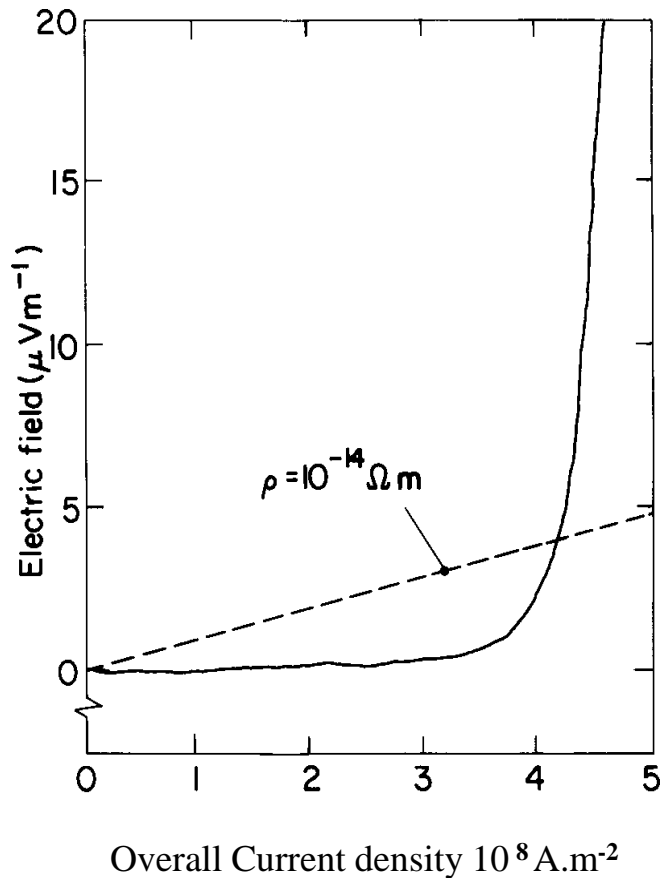
# Measurement of critical current

- this sample holder is placed in the bore of a superconducting solenoid, usually in liquid helium boiling at 4.2K
- at each field level the current is slowly increased and voltage across the test section is measured



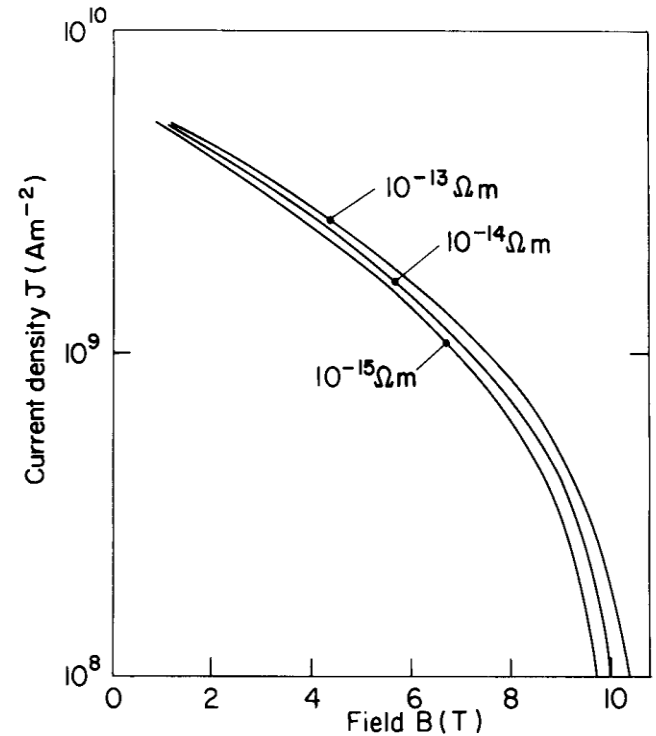
# Resistive transition 1

When measured sensitively, the boundary between superconducting and resistive states is not sharp, but slightly blurred.



If we measure  $J_c$  with voltage taps across the sample, we see that the voltage rises gradually.

To define  $J_c$ , we must therefore define a measurement sensitivity in terms of electric field or effective resistivity.



Commonly used definitions are  $\rho = 10^{-14} \Omega\text{m}$  or  $E = 1 \mu\text{V}\cdot\text{m}^{-1}$

Critical current defined at this level is about what you would expect the conductor in a resin impregnated solenoid to achieve. At higher resistivity, self heating starts to raise the internal temperature and reduce the critical current



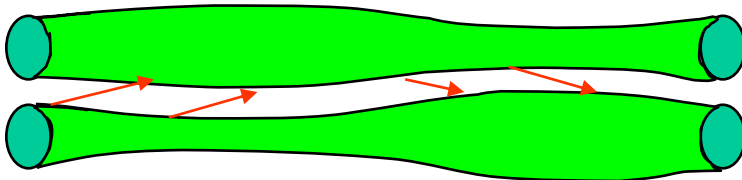
# Resistive transition 2

It has been found empirically that the resistive transition may be represented by a power law

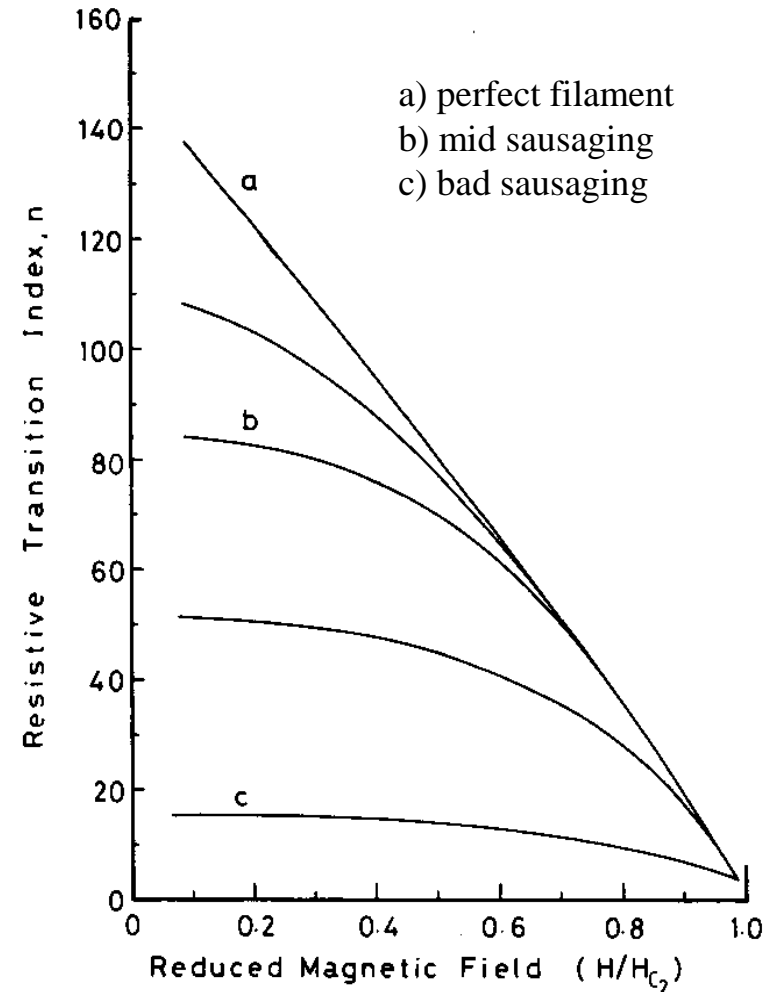
$$\rho(J) = \rho_o \left\{ \frac{J}{J_o} \right\}^n$$

where n is called the resistive transition index.

- the effect is partly within the filaments (flux flow) and partly between the filaments
- 'sausaging of the filaments, forces current to cross the copper matrix as critical current is approached.



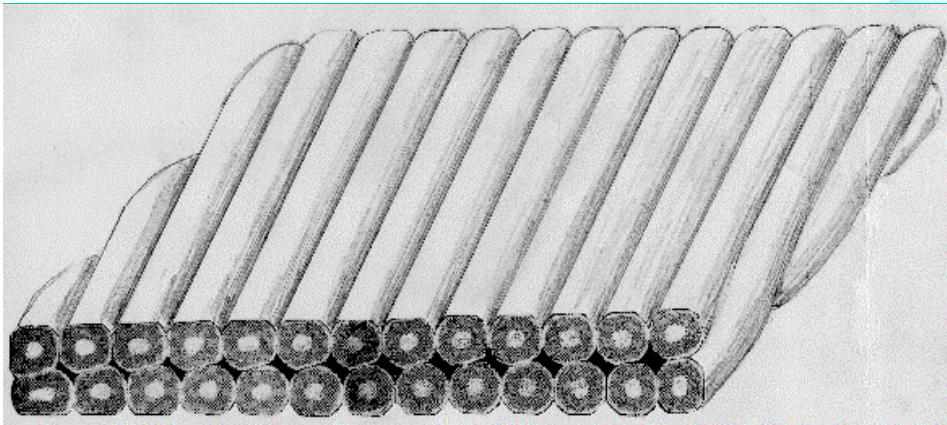
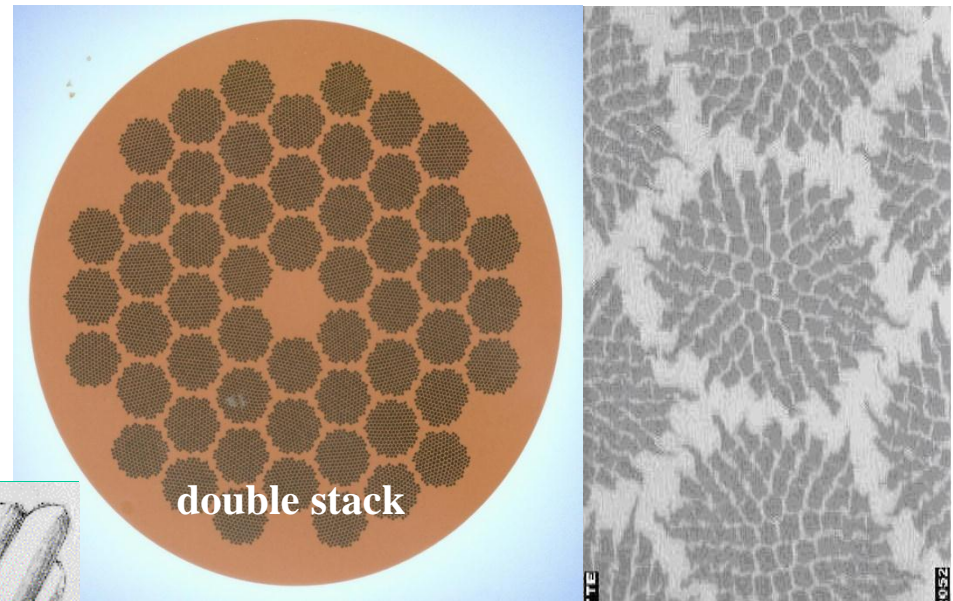
- resistive transition can be the main source of decay in persistent magnets
- 'n' is often taken as a measure of quality - look for  $n > 50$
- HTS conductors so far have low  $n \sim 5 - 10$



# Conductors for accelerator magnets



- to date, all superconducting accelerators have used NbTi superconductor.
- to control field errors and ac losses, the filaments must be  $< 10\mu\text{m}$  diameter (lectures 2 & 3)

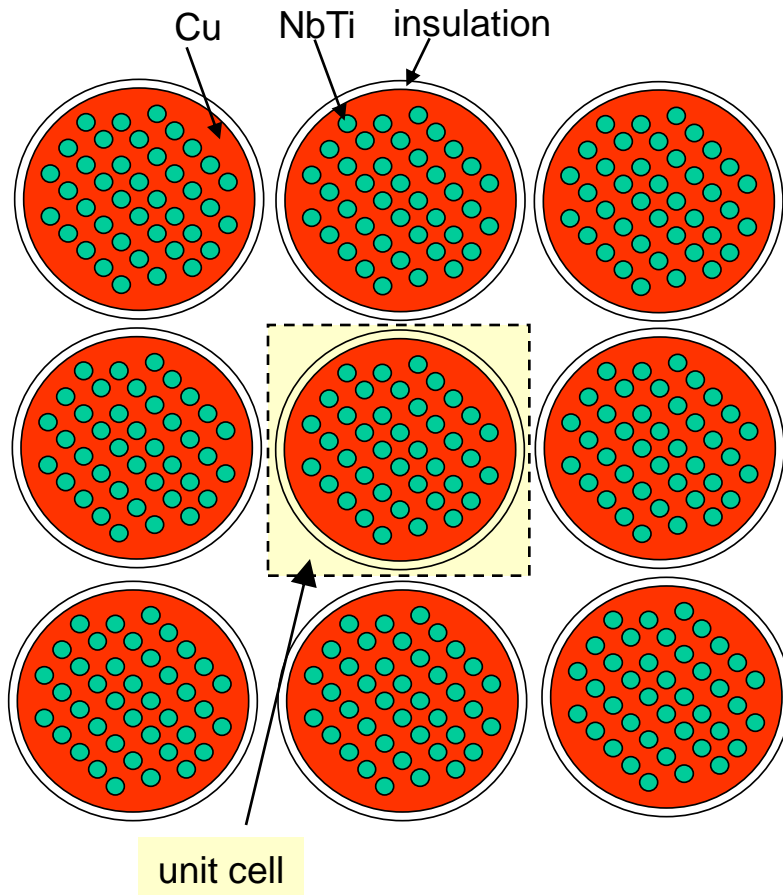


- to get the necessary high operating currents, many wires must be cabled together (lecture 3)

# Engineering current density

In designing a magnet, what really matters is the overall 'engineering' current density  $J_{eng}$

$$J_{eng} = \frac{\text{current}}{\text{unit cell area}} = J_{sup} \times \lambda_{sup}$$



fill factor within the wire  $\lambda_{wire} = \frac{1}{(1 + mat)}$

where  $mat$  = matrix : superconductor ratio

typically:

for NbTi  $mat = 1.5$  to  $3.0$  ie  $\lambda_{sup} = 0.4$  to  $0.25$

for Nb<sub>3</sub>Sn  $mat \sim 3.0$  ie  $\lambda_{sup} \sim 0.25$

for B2212  $mat = 3.0$  to  $4.0$  ie  $\lambda_{sup} = 0.25$  to  $0.2$

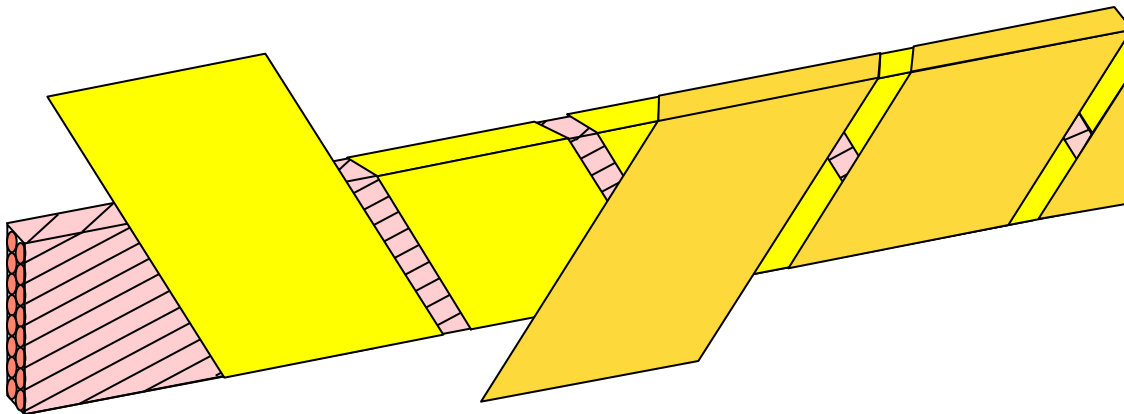
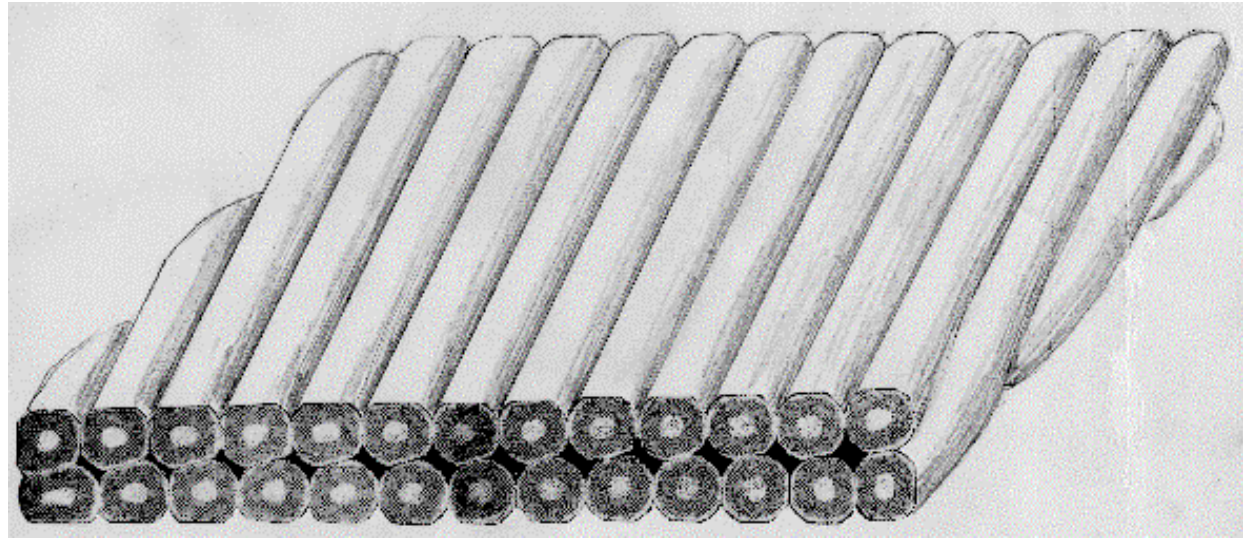
$\lambda_{winding}$  takes account of space occupied by insulation, cooling channels, mechanical reinforcement etc and is typically 0.7 to 0.8

So typically  $J_{eng}$  is only 15% to 30% of  $J_{supercon}$



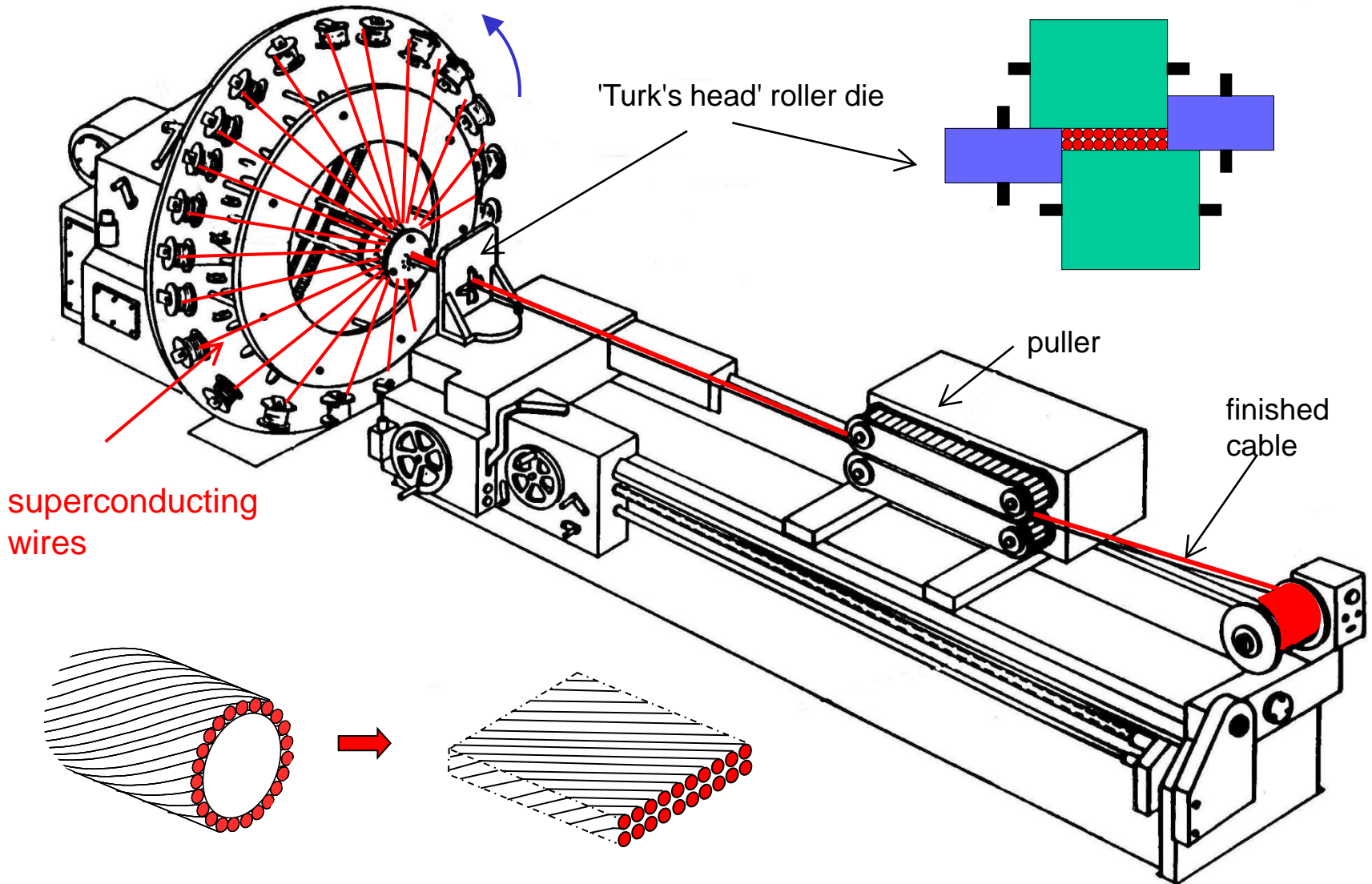
# Rutherford cable

- for high current applications, such as accelerators, we need many wires in parallel
- the most popular way of doing this is the Rutherford cable (see lecture 3)

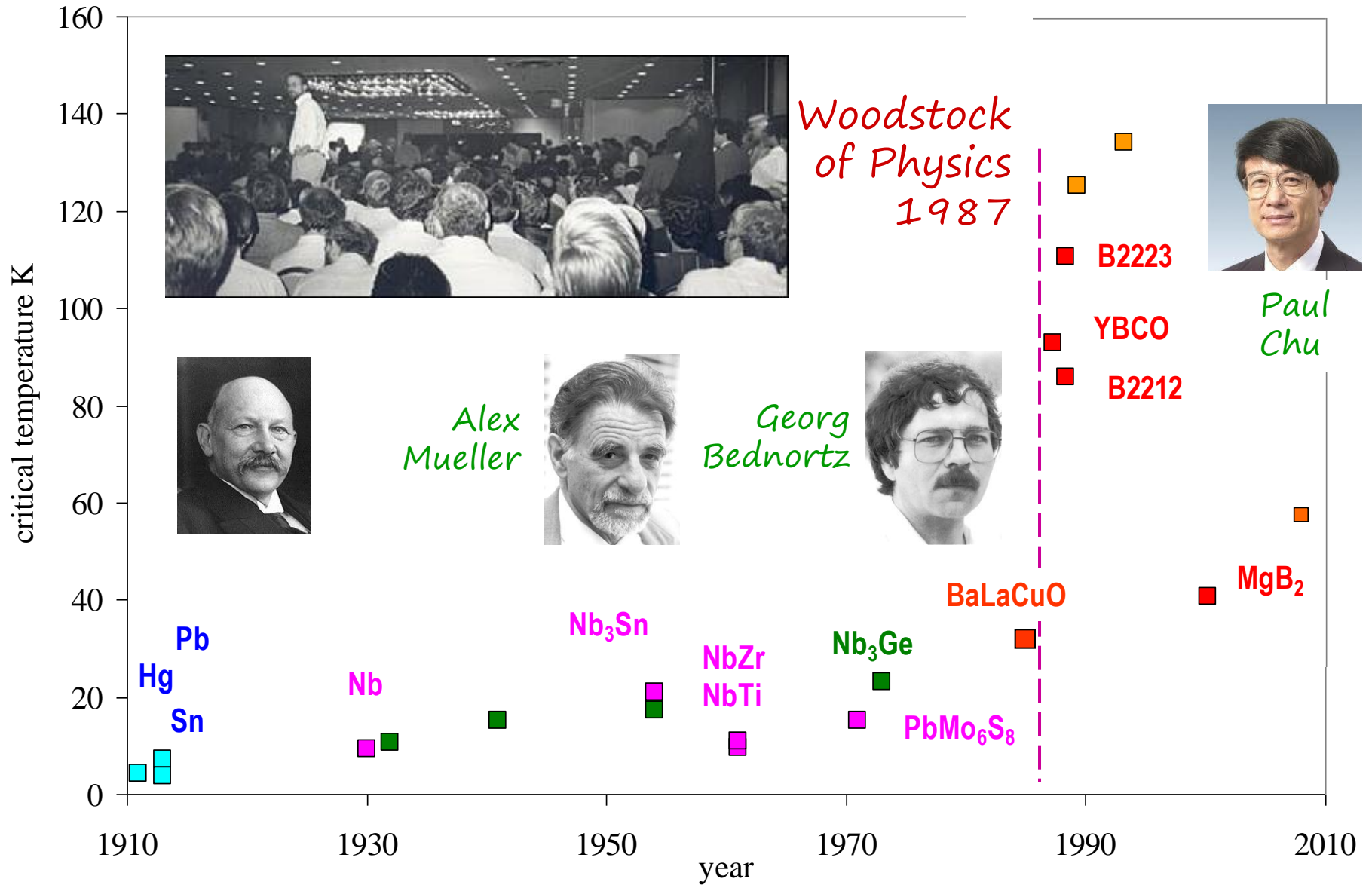


- Rutherford cable is usually insulated by wrapping it with Kapton tape

# Manufacture of Rutherford cable

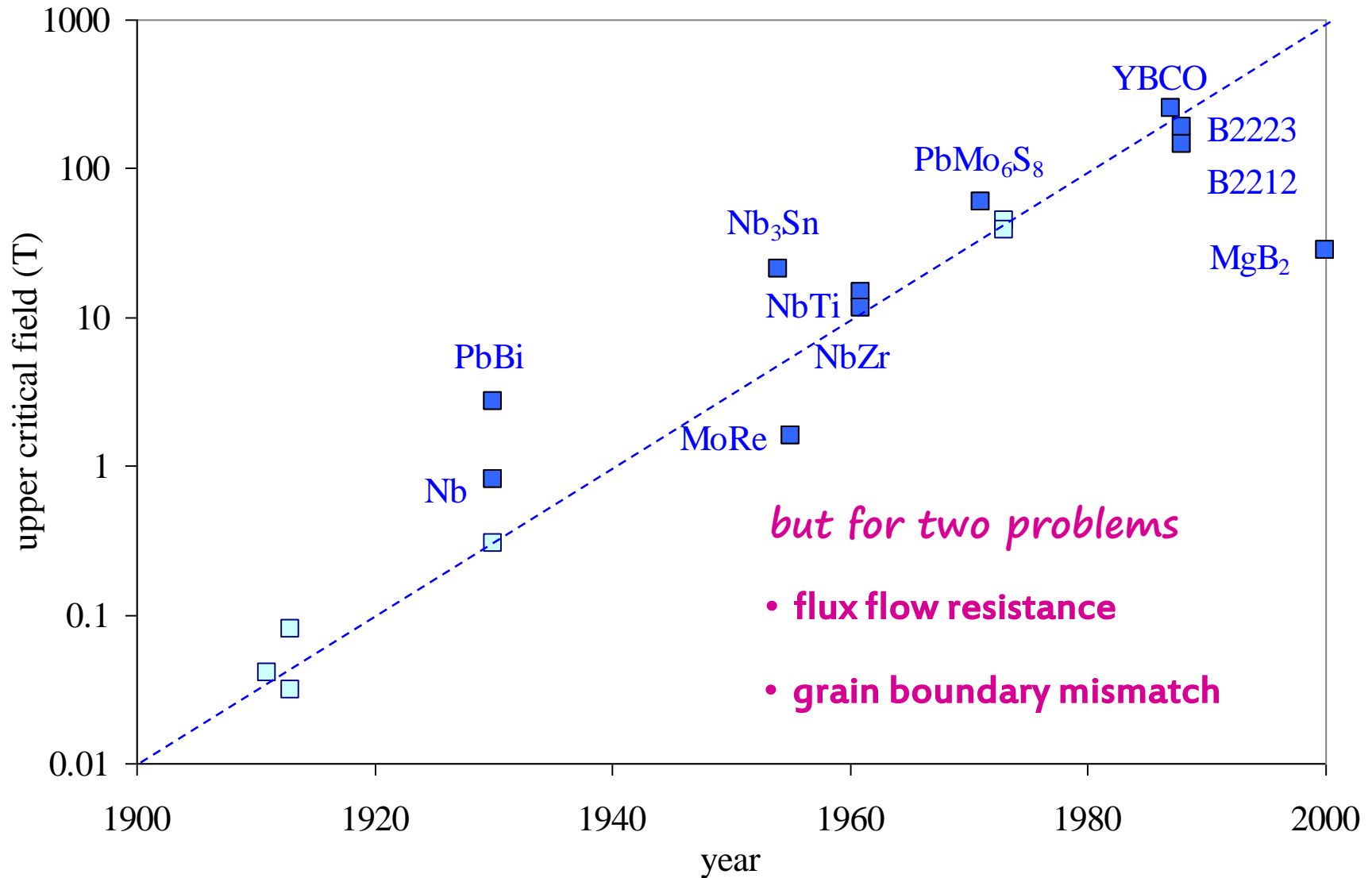


# A century of critical temperatures

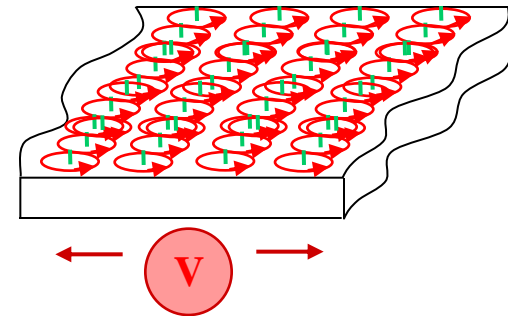
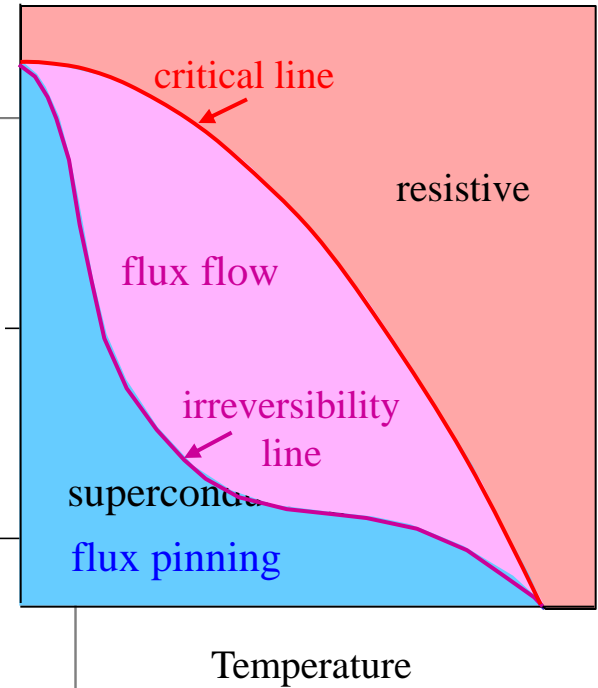
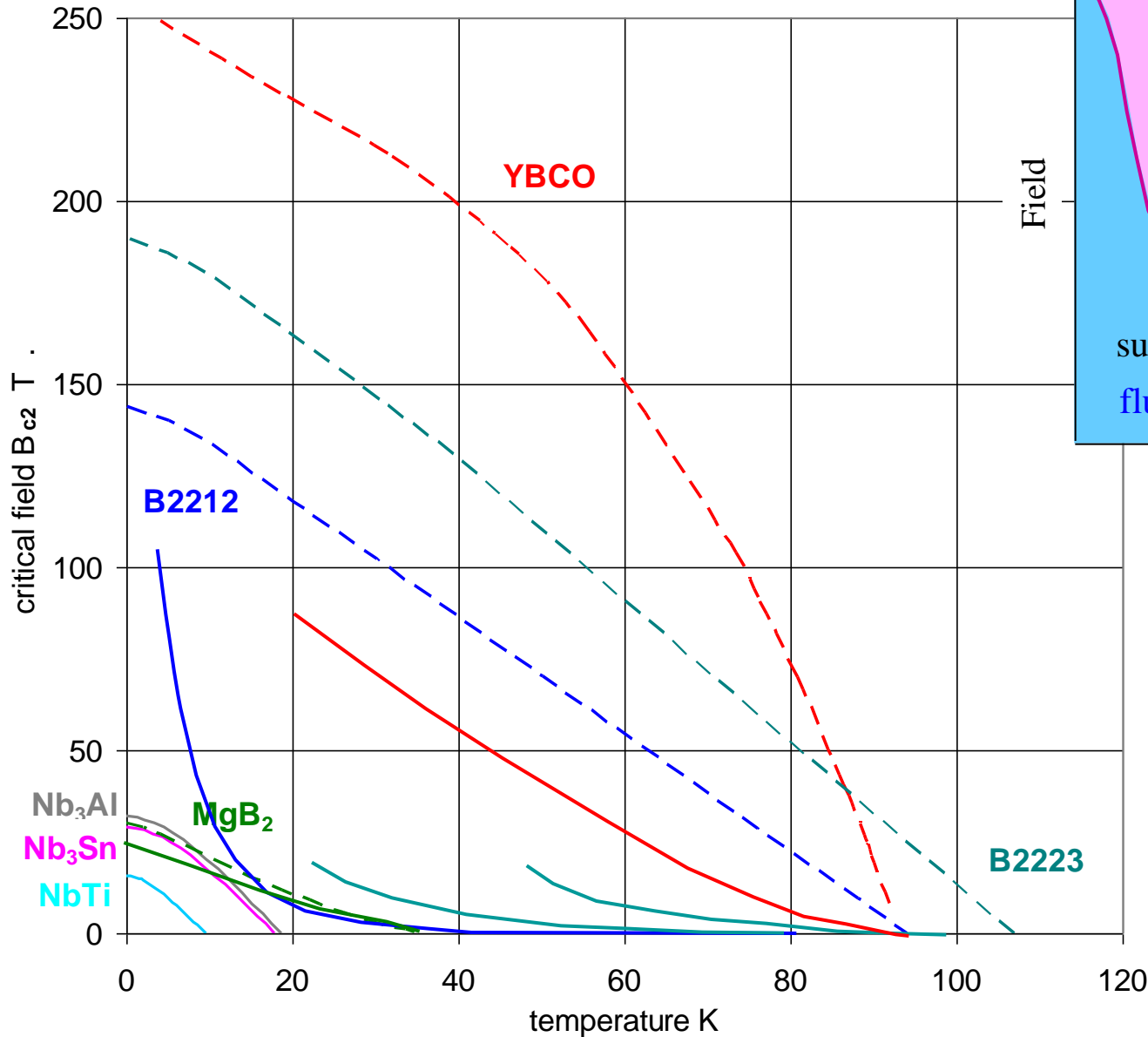




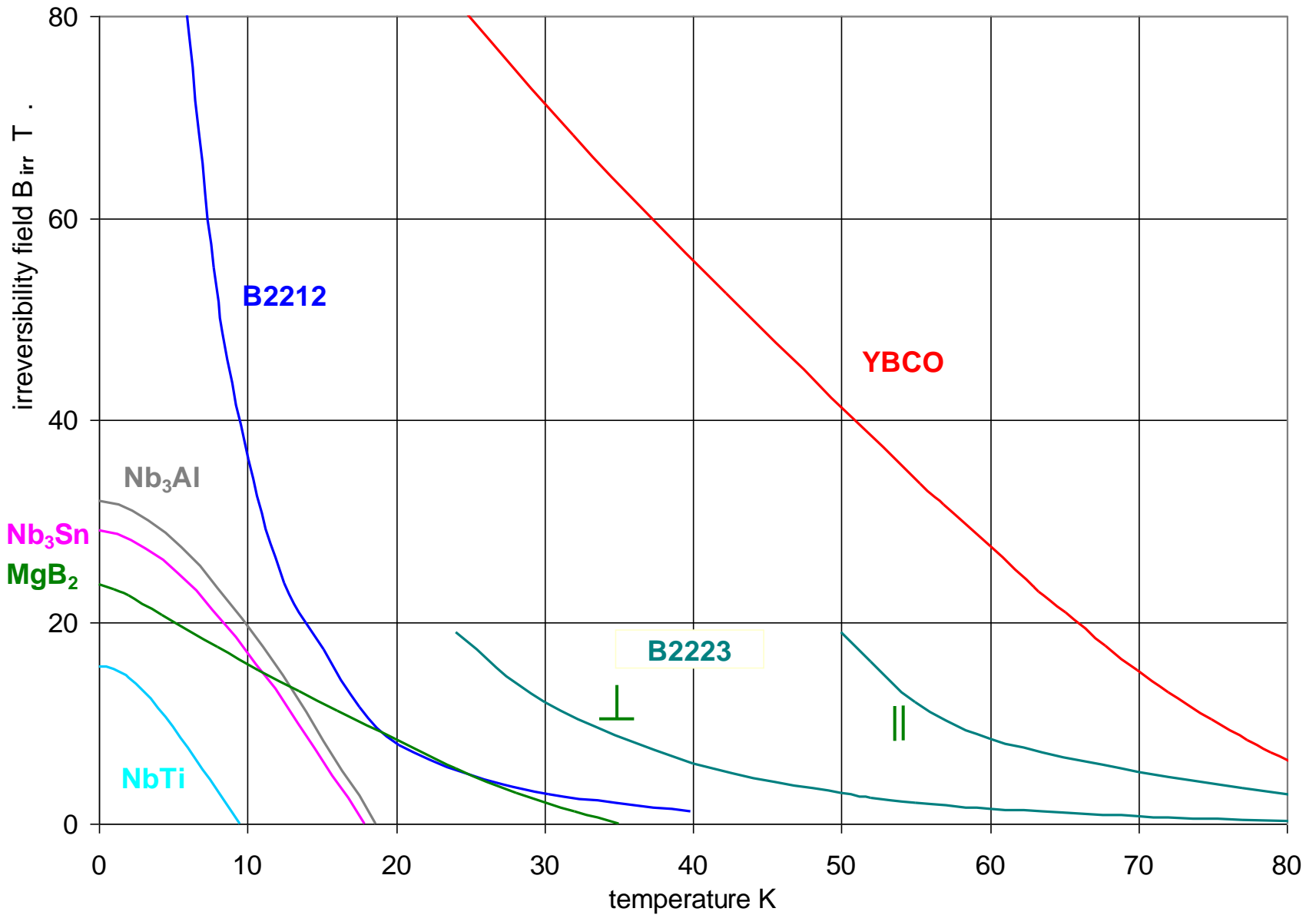
# Wonderful materials for magnets



# 1) Flux flow resistance

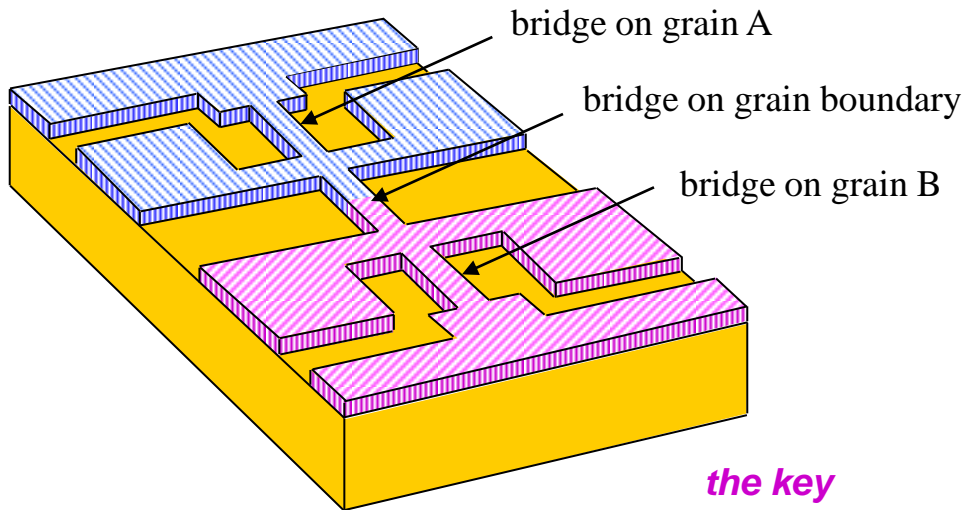
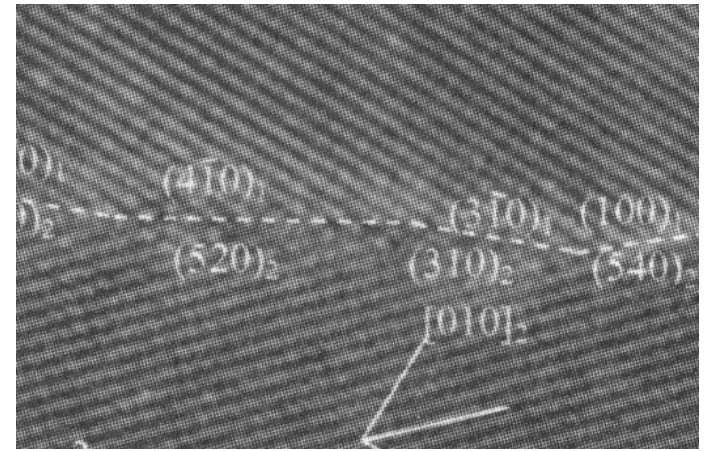


# Accessible fields for magnets

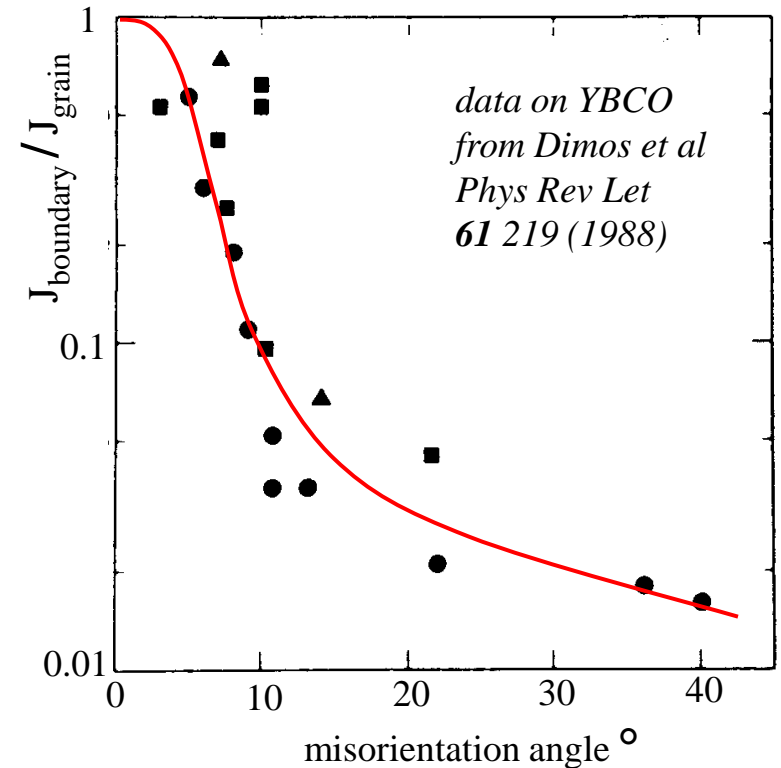


## 2) Grain boundary mismatch

- crystal planes in grains point in different directions
- critical currents are high within the grains
- $J_c$  across the grain boundary depends on the misorientation angle
- For good  $J_c$  must align the grains to within a few degrees



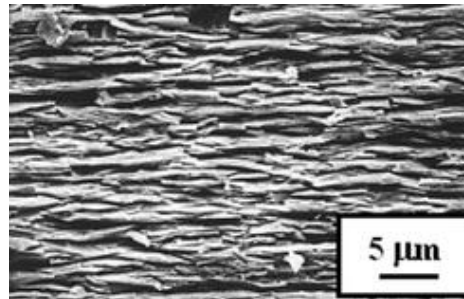
the key measurement of Dimos et al



# Practical HTS conductors

## B2212 & B2223

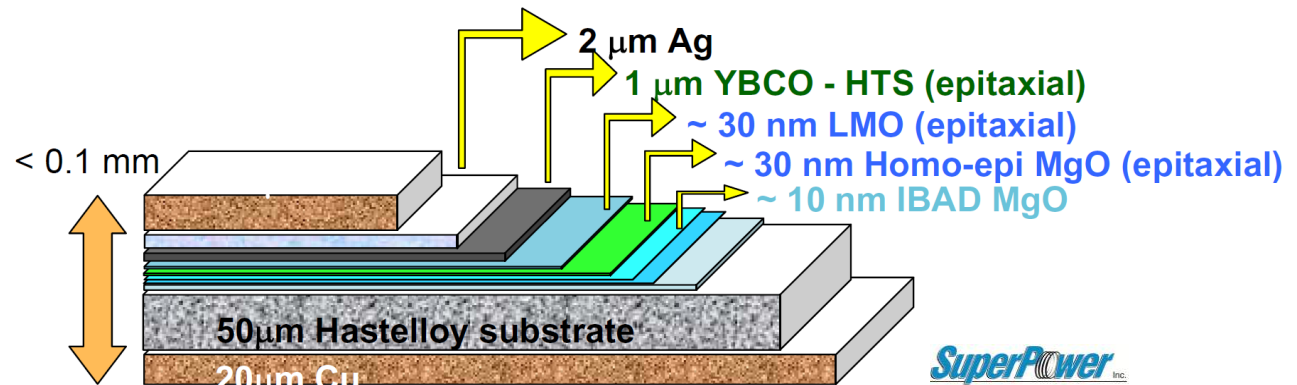
- grains tend to align when processed with silver
- but low irreversibility field



- OK in high field at low temperature  
- *high field inserts*
- OK in low field at high temperature  
- *power transmission cables*

## YBCO

- best irreversibility field
- very sensitive to grain boundary misalignment
- grains do not line up naturally
- deposit YBCO film on aligned substrate



- OK in high field **and** at high temperature

# Lecture 1: concluding remarks

- superconductors allow us to build magnets which burn no power (except refrigeration)
- ampere turns are cheap, so don't need iron
  - ⇒ fields higher than iron saturation (but still use iron for shielding)
- performance of all superconductors described by the critical surface in  $B J \theta$  space,
- three kinds of superconductor
  - **type 1**: low temperature, unsuitable for high field
  - **type 2**: low temperature, good for high field - but must create flux pinning to get current density
  - **HTS**: high temperature, high field - but current density is still a problem
- NbTi is the most common commercial superconductor - standard production process
- Nb<sub>3</sub>Sn has higher critical field & temperature - specialized commercial production
- BSCO high temperature **or** high field, but not both - prototype commercial production
- YBCO high temperature **and** high field, but must align the grains - research production
- measure  $I_c$  to check specification, the index  $n$  indicates quality
- for accelerators, so far it's only been NbTi, usually in Rutherford cables