

# **The Entropic Approach to Understanding the Cosmological Constant**

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with Andrew Frey and Gil Holder

CERN, 8 August 2007

# Outline

- History of the cosmological constant problem
- The anthropic approach: successes and shortcomings
- A new approach: the causal entropic principle
- New results
- A connection to dark matter?

# History of the Cosmological Constant

1917: Einstein introduces  $\Lambda$  to obtain a static universe

$$\mathcal{L} = \frac{1}{16\pi G} (R - 2\Lambda)$$

1912–1924: Slipher measures redshifts of galaxies, suggesting expanding universe before Hubble (1929)

1923: Einstein writes to Weyl,

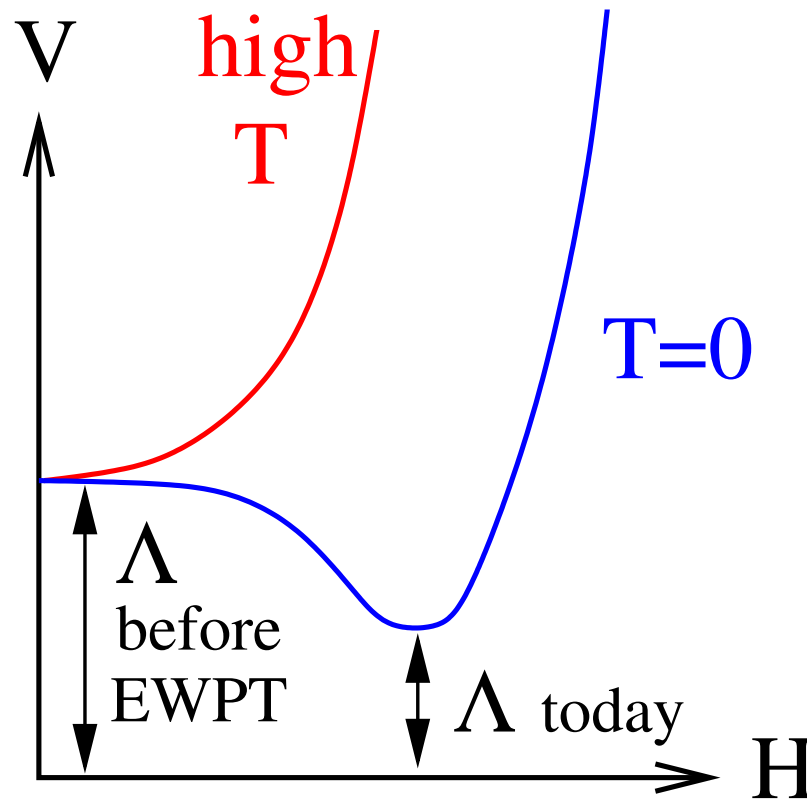
**“If there is no quasi-static world, then away with the cosmological term!”**

1967-1968:  $\Lambda$  invoked to get loitering universe, to explain excess of quasars at  $z = 1.95$

1967: Zeldovich proposes quantum fluctuations as origin of  $\Lambda$ , but they give too big result!

# History of the Cosmological Constant

1970's: spontaneous symmetry breaking in electroweak theory forces particle theorists to take vacuum energy more seriously



Early 1980's: first attempts to explain why  $\Lambda$  is small.  
A. Zee refers to the  $\Lambda$  "problem"

# History of the Cosmological Constant

1984: Hawking proposes probability distribution for  $\Lambda$  peaked at zero, from Euclidean quantum gravity,

$$P \sim \exp (M_p^2 / \Lambda)$$

1986-1988: Linde's eternal inflation provides universes with different values of  $\Lambda$

1987: Weinberg derives anthropic bound on  $\Lambda$  (but Hawking discussed idea at conference in 1981)

1988: Coleman's wormhole approach (> 500 citations)  $\Rightarrow$

1987-1988: Brown and Teitelboim suggest relaxation of  $\Lambda$  through nucleation of bubbles with successively smaller values of  $\Lambda$

1989: Weinberg's no-go theorem against smooth relaxation of  $\Lambda \rightarrow 0$  (in review article, 1300 citations)

# Coleman's wormhole mechanism

*S Coleman / Cosmological constant*

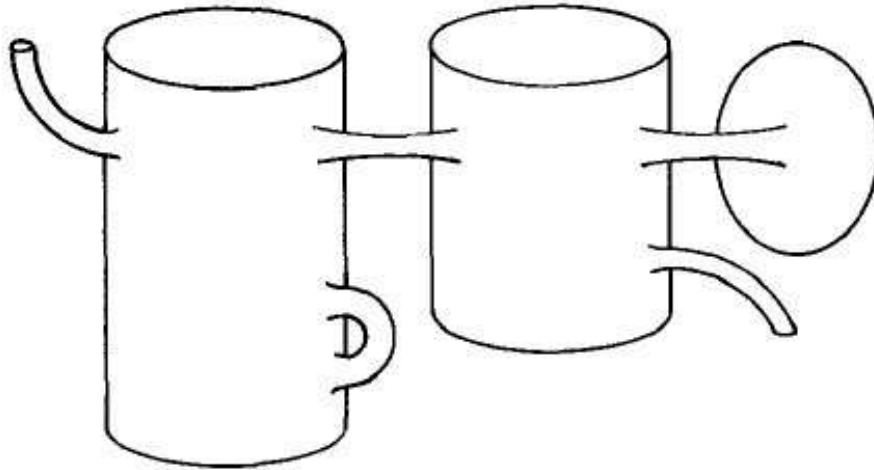


Fig 3 A manifold with wormholes and baby universes

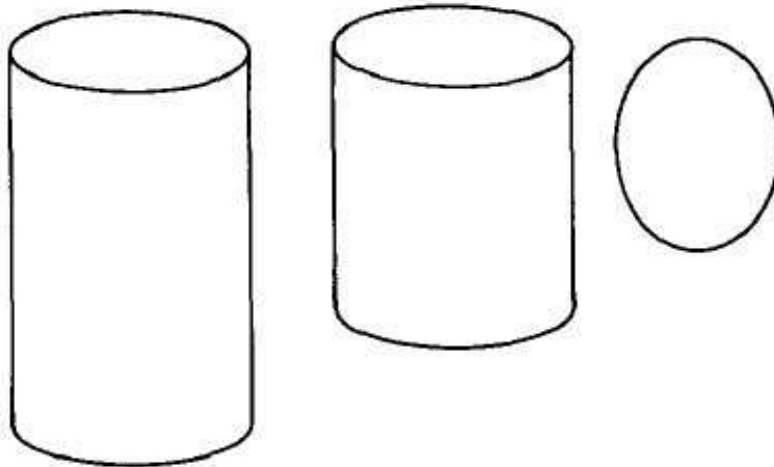


Fig 4 The manifold of fig 3, with wormholes and baby universes stripped away

Wormholes  $\rightarrow$   
superselection  
sectors with  
different values of  
 $\Lambda$

Probability  
distribution for  $\Lambda$ :

$$P \sim e^{e^{1/\Lambda}}$$

Other constants  
of nature also  
predicted, e.g.

$$\theta_{\text{QCD}} = \pi$$

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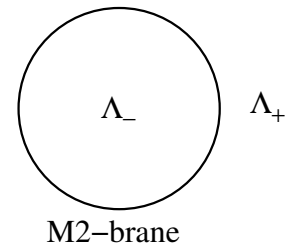
# History of the Cosmological Constant

1990-91: Efstathiou, Turner propose  $\Lambda$ CDM universe to make  $\Omega_{\text{CDM}} \sim 0.2$  consistent with  $\Omega = 1$ , as already indicated by CMB. Astronomers don't like it!

1998: High- $z$  type Ia supernovae establish  $\Lambda > 0$  ( $> 3000$  citations); also CMB and galaxy clusters  $\Rightarrow$

1999: Huterer and Turner coin the term "Dark Energy"

2000: Bousso and Polchinski realize the mechanism of Brown and Teitelboim within string theory

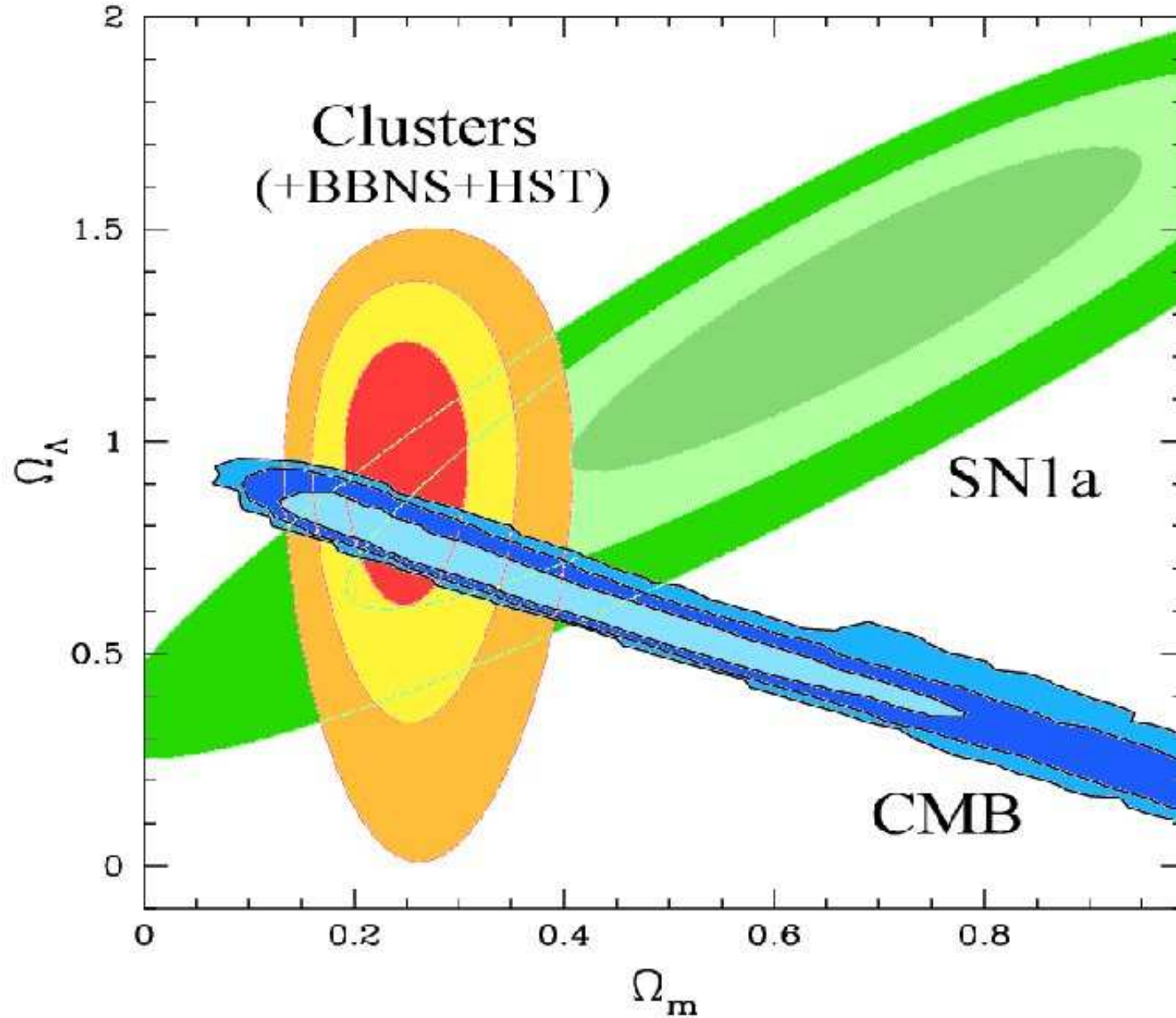


2000 – present: many authors try to use braneworlds to self-tune  $\Lambda \rightarrow 0$ . But this violates Weinberg's no-go theorem.



# Observational evidence for $\Lambda$

Turner and Huterer, 2007



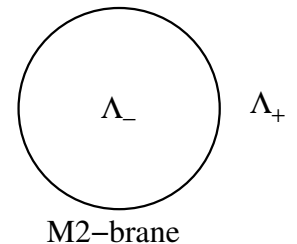
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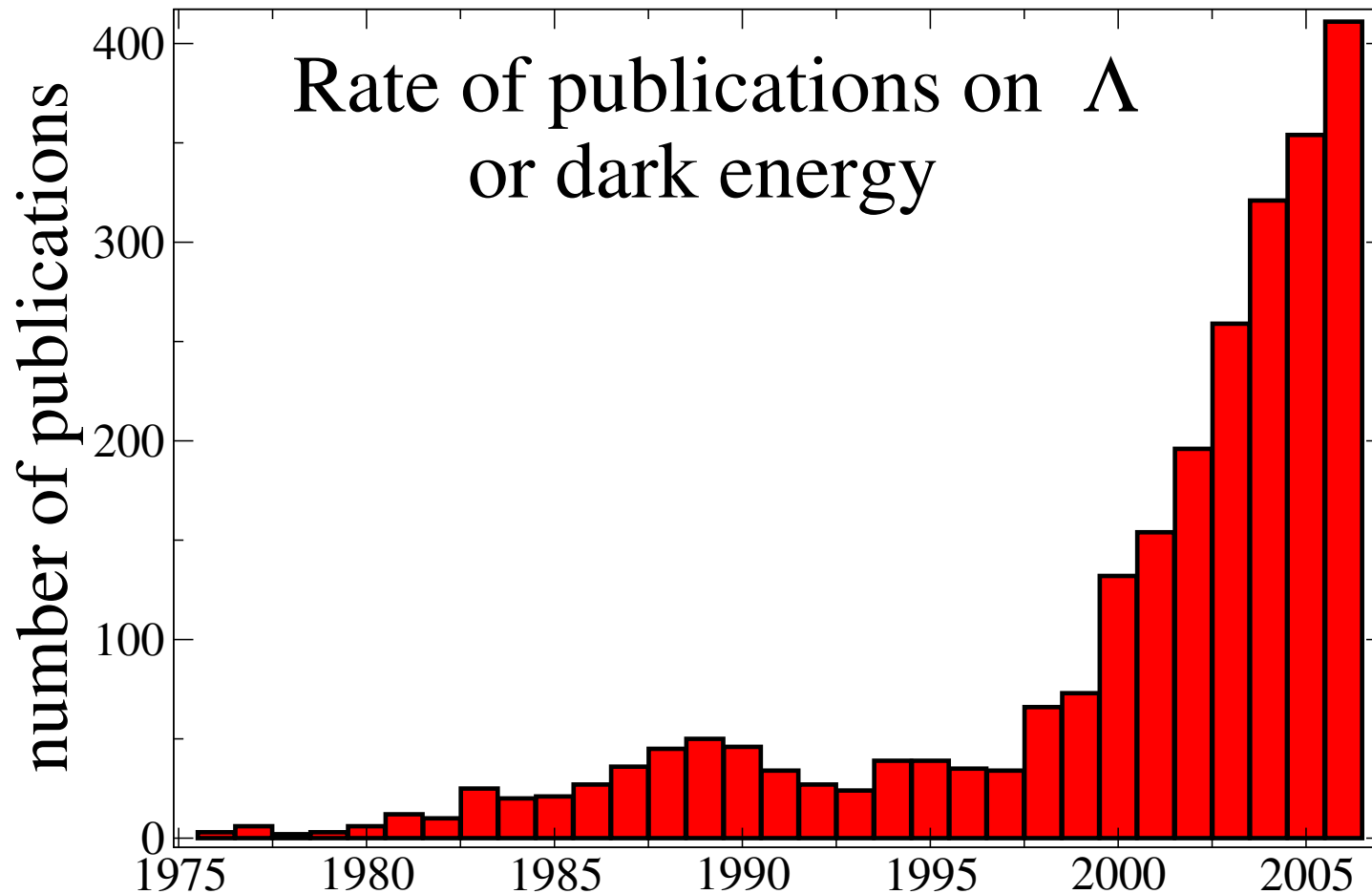
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# History of the Cosmological Constant

2003: Susskind notes that string theory landscape justifies existence of universes with many possible values of  $\Lambda$ .

Douglas counts  $\sim 10^{500}$  string theory vacua—can explain existence of universes with  $\Lambda \sim 10^{-123} M_p^2$

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# The lesson of history

What do we learn from these recollections?

- Much effort (and \$\$\$) is being spent on understanding  $\Lambda$
- Self-tuning where  $\Lambda \rightarrow 0$  continuously is ruled out
- The most promising ideas involve ensemble of universes with different values of  $\Lambda$ , and probability distribution  $P(\Lambda)$ .
- The probabilistic view cannot be divorced from anthropic considerations since

$$P_{\text{obs}}(\Lambda) = P_{\text{prior}}(\Lambda) \times (\text{Probability to have observers})$$

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# The Anthropic Approach

There is no agreed-upon explanation of the hierarchy

$$\frac{\rho_\Lambda}{M_p^4} \cong 10^{-123}$$

but one idea predicted such a value before its measurement: the anthropic principle (Weinberg, 1987). Too-fast expansion (or collapse) of universe prevents formation of structure at redshift  $z_c$  unless

$$-\rho_{m,0} < \rho_\Lambda < 3(1 + z_c)^3 \rho_{m,0} \sim 400 \rho_{m,0}$$

for  $z_c = 4$ . If prior probability distribution for  $\Lambda$  is flat, then probability of our universe is

$$P(\rho_\Lambda \lesssim \rho_{m,0}) \cong \frac{1}{400} = 0.25\%$$

Not as unlikely as  $\frac{1}{10^{123}}$ !

# Problems with Anthropic Approach

We now know of dwarf galaxies which formed at  $z = 10$ .  
This weakens anthropic constraint to

$$\rho_{\Lambda} < 5000 \rho_{m,0}$$

Our universe looks less likely.

What if other quantities are allowed to vary, like density contrast

$$Q = \frac{\delta\rho}{\rho} \text{ from inflation} \cong 2 \times 10^{-5} ?$$

Structure can form more easily, despite fast expansion.

Tegmark and Rees (1997): bound on  $\Lambda$  is further weakened,

$$\rho_{\Lambda} < 5000 \left( \frac{Q}{Q_{\text{obs}}} \right)^3 \rho_{m,0}$$

If  $Q \rightarrow 0.1$ , then [bound on  $\rho_{\Lambda}$ ]  $\rightarrow 10^{14} \rho_{m,0}$

# Problems with Anthropic Approach

To save anthropic explanation, need further assumptions, e.g., life requires galaxies of some minimum mass (heavy elements, needed for life, not retained if  $M_{\text{galaxy}} < 10^9 M_{\odot}$ ).

Such assumptions seem arbitrary and hard to justify. Perhaps life can exist without carbon or water? Perhaps “observers” completely unlike us can exist?



# Argument in favor of anthropics

If  $\Lambda_{\text{obs}}$  were 0, a dynamical mechanism for explaining it would be likely.

But why should a physics mechanism coincidentally predict  $\Lambda$  in the bizarre anthropic range?

That sounds more anthropocentric than admitting our existence is a data point!

Perhaps  $\Lambda = 0$  and dark energy is quintessence, but I think quintessence creates more problems than  $\Lambda$  by itself.

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# The Causal Entropic Principle

Idea suggested by

R. Bousso, R. Harnik, G. D. Kribs and G. Perez,  
“Predicting the Cosmological Constant from the Causal  
Entropic Principle,” arXiv:hep-th/0702115.

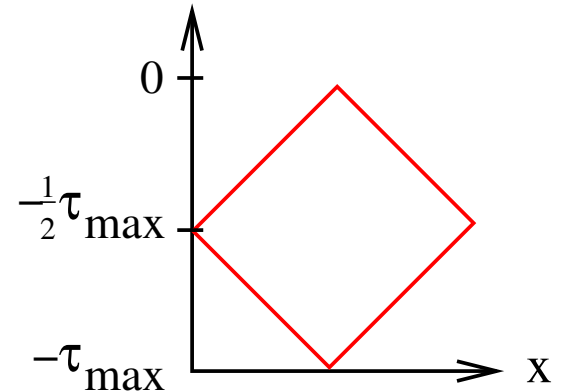
A replacement for the anthropic principle:

$$\left[ \text{Probability of finding observers} \right] \propto$$
$$\left[ \text{entropy produced in causally connected region of universe} \right]$$

- Simple, quantitative, makes minimal assumptions about nature of observers;
- Circumvents measure problem of eternal inflation

# The Causal Diamond

- Conformal time  $ds^2 = a^2(\tau)(d\tau^2 - d\vec{x}^2)$
- Start at some early time, e.g., reheating
- End at  $t = \infty$ , but  $\tau$  is finite, due to de Sitter horizon,



In Planck units,

$$\tau_{\max} = \int \frac{dt}{a(t)} \sim 2.8 t_{\Lambda}^{1/3} \equiv 2.8 \left( \frac{3}{\Lambda} \right)^{1/6}$$

assuming universe is first matter-, then  $\Lambda$ -dominated,

$$a(t) = \left[ t_{\Lambda} \sinh \left( \frac{3t}{2\Lambda} \right) \right]^{2/3} .$$

Comoving volume of diamond  $V_c = \frac{4\pi}{3} r^3(\tau)$  is finite at all times

# Entropic Principle

Amount of entropy created inside the diamond determines probability for evolving complexity, structure, observers:

$$\frac{dP}{d\Lambda} = \Delta S(\Lambda)$$

$\Delta S$  indicates free energy was available, needed for making measurements. Minimal assumption: observers must obey laws of thermodynamics!

Opposite situation: universe which remains in thermal equilibrium. No complexity can arise.

Same approach might be used to predict other variables besides  $\Lambda$ , if they vary over landscape.

# How to compute $\Delta S$

Product of two factors:

$$\Delta S = \int dt V_c(t) \frac{dS}{dV_c dt}$$

$V_c$  = comoving volume of causal diamond; determined by cosmology;

$\frac{dS}{dV_c dt}$  = entropy production rate per unit volume; depends on astrophysics. Sources of entropy:

Stars:  $\Delta S_* \cong 10^6$  per baryon

Active Galactic Nuclei:  $< \frac{1}{6} \Delta S_*$

Supernovae:  $< \frac{1}{10} \Delta S_*$

Galaxy cooling:  $< \frac{1}{100} \Delta S_*$

# Entropy production by stars

Starlight is the biggest source of entropy:

1 UV  $\gamma$  scattered by dust  $\implies > 100$  IR (20 meV)  $\gamma$ 's.

Rate of entropy production per comoving volume:

$$\frac{dS}{dV_c dt}(t) = \int_0^t dt' \frac{d^2 S}{dM_* dt}(t - t') \dot{\rho}_*(t')$$

$\frac{d^2 S}{dM_* dt}(t - t')$  = rate of entropy production per stellar mass, at time  $t$ , from stars born at time  $t'$ ,

$$\frac{d^2 S}{dM_* dt}(t - t') = \frac{1}{\langle M \rangle} \int_{0.08 M_\odot}^{M_{\max}(t-t')} dM \xi_{\text{IMF}}(M) \frac{d^2 s}{dN_* dt}$$

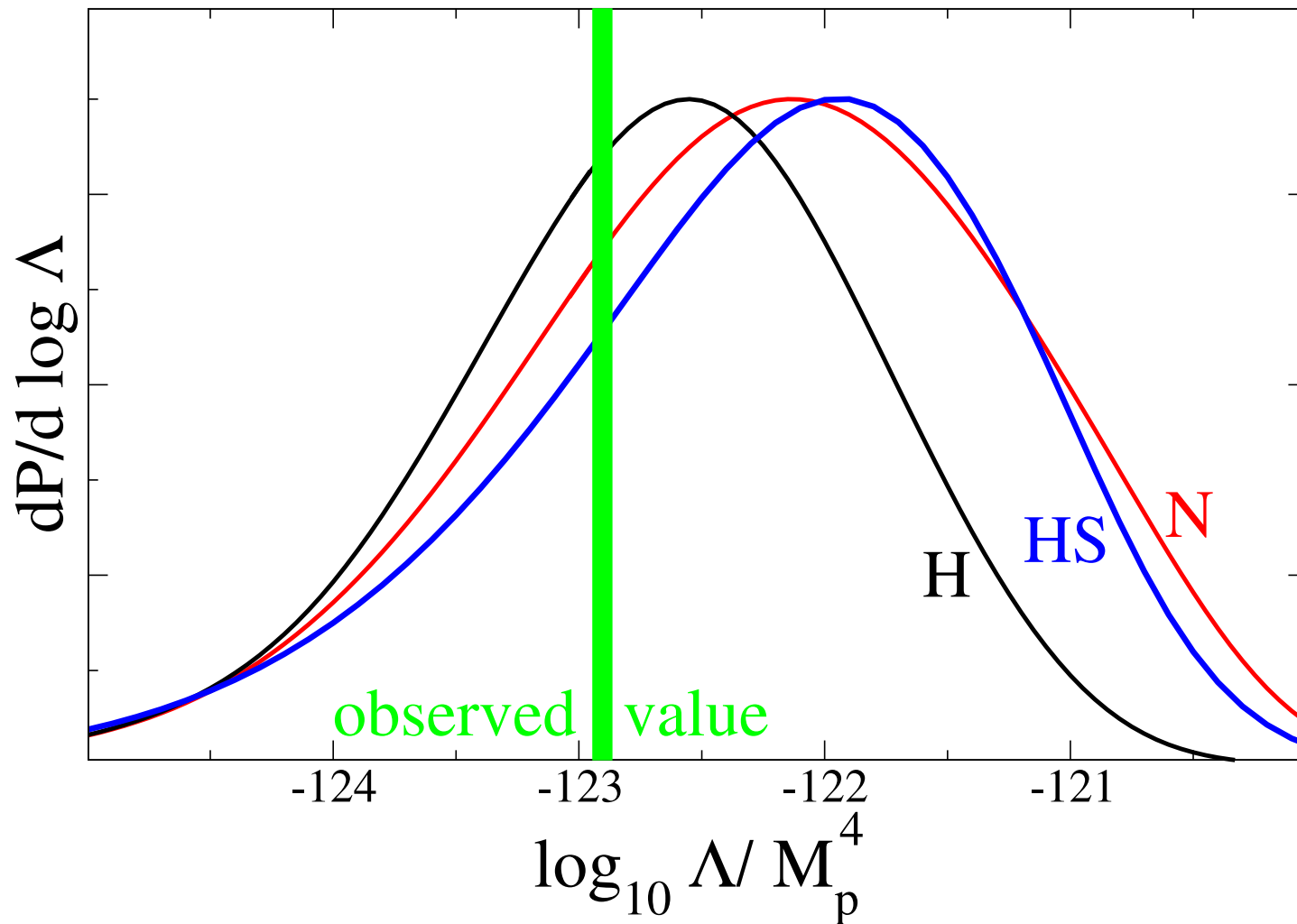
$\dot{\rho}_*$  = stellar mass formation rate

$\xi_{\text{IMF}}$  = “initial mass function”  $\sim M^{-2.35}$

$\frac{d^2 s}{dN_* dt}$  = rate per star =  $L_*/T_{\text{eff}} \sim (M/M_\odot)^{3.5} \times 10^{54} \text{ yr}^{-1}$

# Probability Distributions for $\Lambda$

For different published star formation rates  $\dot{\rho}_*$   
(Hopkins-Beacom, Nagamine *et al.*, Hernquist-Springel)



Probability to observe  $\Lambda = 10^{-123} M_p^4$  is  $O(1)$ !

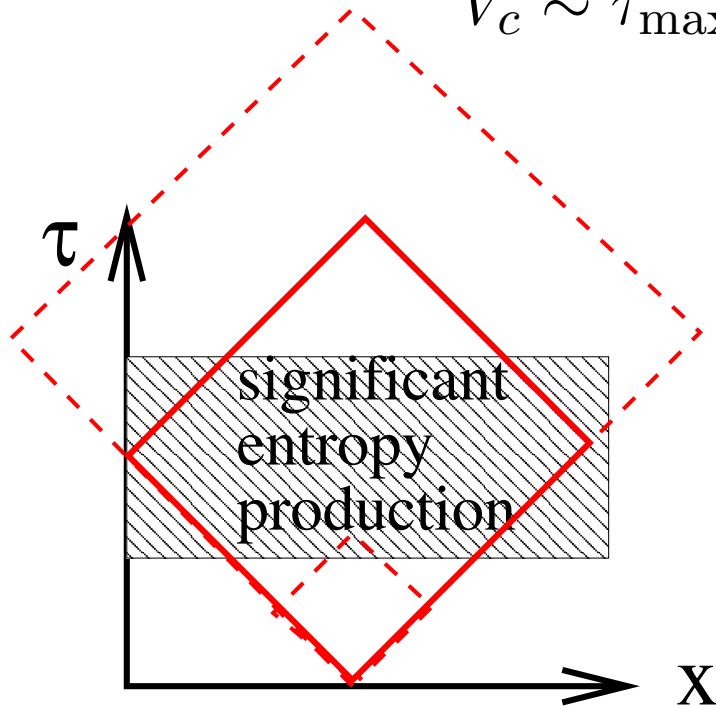


# Basic predictions are robust

The basic prediction remains unaffected by details of  $\dot{\rho}_*$ :  
 $dP/d \log \Lambda$  peaks near  $\Lambda_{\text{obs}}$ .

Easy to understand: Main  $\Lambda$ -dependence is through volume of causal diamond,

$$V_c \sim \tau_{\text{max}} \sim \frac{1}{\sqrt{\Lambda}}, \quad \int dt V_c \sim \frac{1}{\Lambda}$$



Preferred  $\Lambda$  has time of maximum diamond volume coinciding with time of maximum entropy production by stars.

# Contrast with anthropic prediction

Suppose  $\Lambda_a$  is anthropic cutoff on  $\Lambda$ :

$$\frac{dP}{d\Lambda} = \begin{cases} \text{const.}, & \Lambda < \Lambda_a \\ 0, & \Lambda > \Lambda_a \end{cases}$$

Then

$$\frac{dP}{d \log \Lambda} = \begin{cases} c\Lambda, & \Lambda < \Lambda_a \\ 0, & \Lambda > \Lambda_a \end{cases}$$

is sharply peaked at  $\Lambda_a$ .

Causal diamond volume's  $\Lambda$ -dependence changes this to

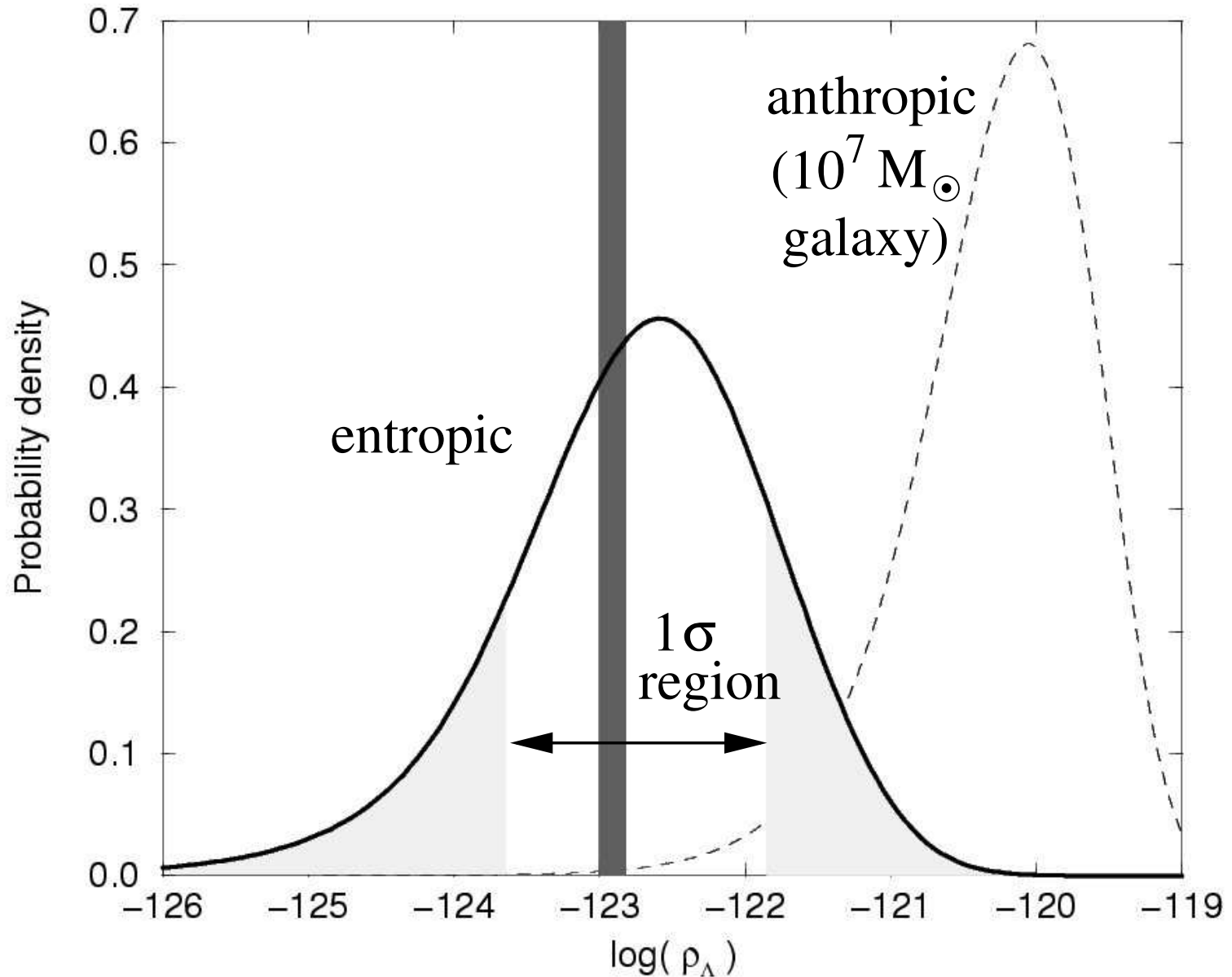
$$\frac{dP}{d\Lambda} = \begin{cases} \frac{c}{\Lambda}, & \Lambda < \Lambda_a \\ 0, & \Lambda > \Lambda_a \end{cases} \implies \frac{dP}{d \log \Lambda} = \begin{cases} c, & \Lambda < \Lambda_a \\ 0, & \Lambda > \Lambda_a \end{cases}$$

ignoring time-dependence of  $\frac{dS}{dV_c dt}$ . Including  $\frac{dS}{dV_c dt}$ ,  $c \rightarrow c(\Lambda)$ ,

$\frac{dP}{d \log \Lambda}$  is no longer peaked at  $\Lambda_a$

# Contrast with anthropic prediction

from Bousoo *et al.*



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# How can this idea be further tested?

Other quantities besides  $\Lambda$  should be scanned by the landscape. Can we predict their values? *E.g.*,

$Q$  (density contrast)

$k$  (spatial curvature)

$\Omega_b/\Omega_{\text{DM}}$  (baryon to dark matter ratio)

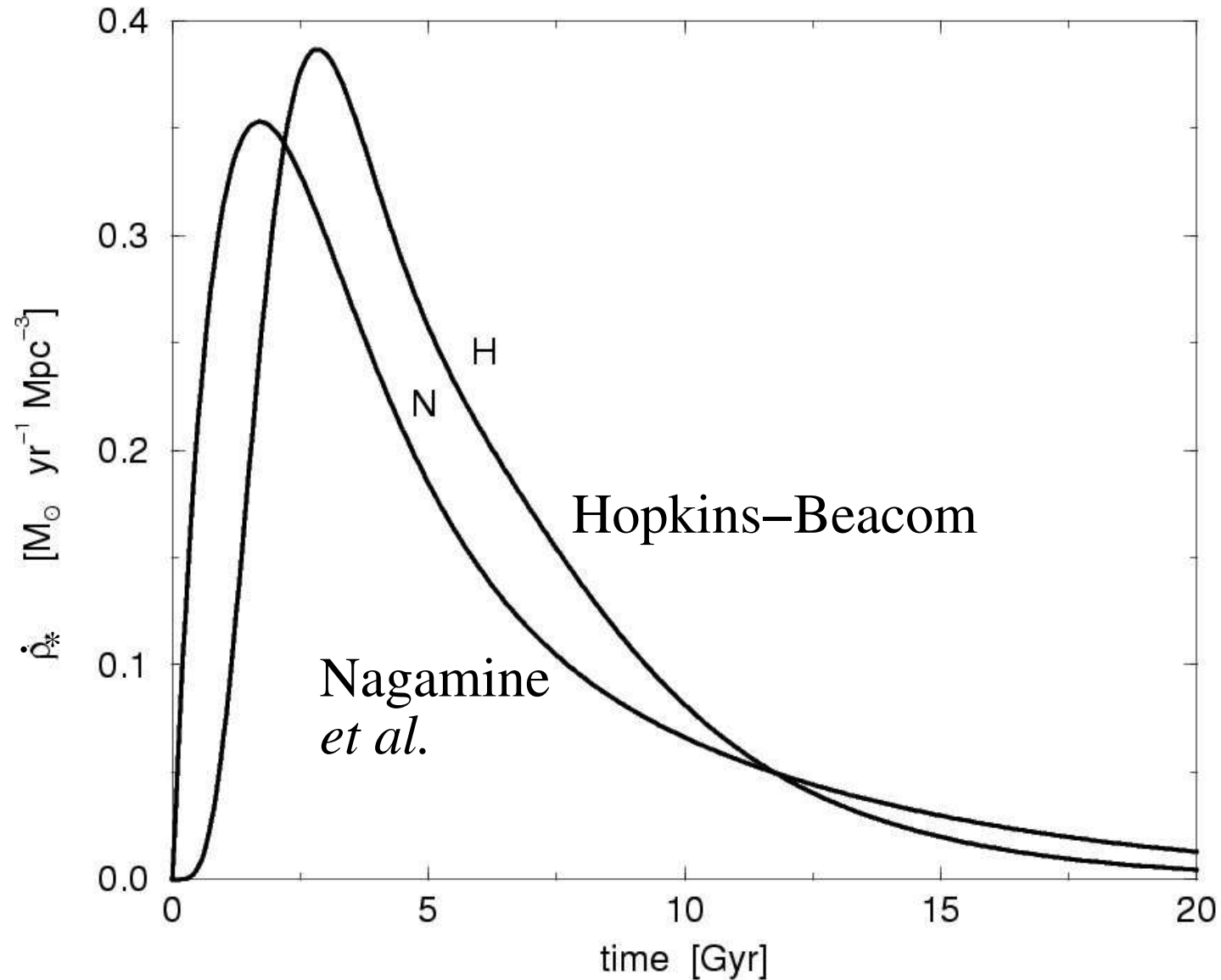
$n_m/n_\gamma$  (matter to photon ratio)

Start small: explore region of landscape containing universes like ours— $\Delta S$  dominated by starlight

Challenge: how does star formation rate  $\dot{\rho}_*$  depend on these quantities?

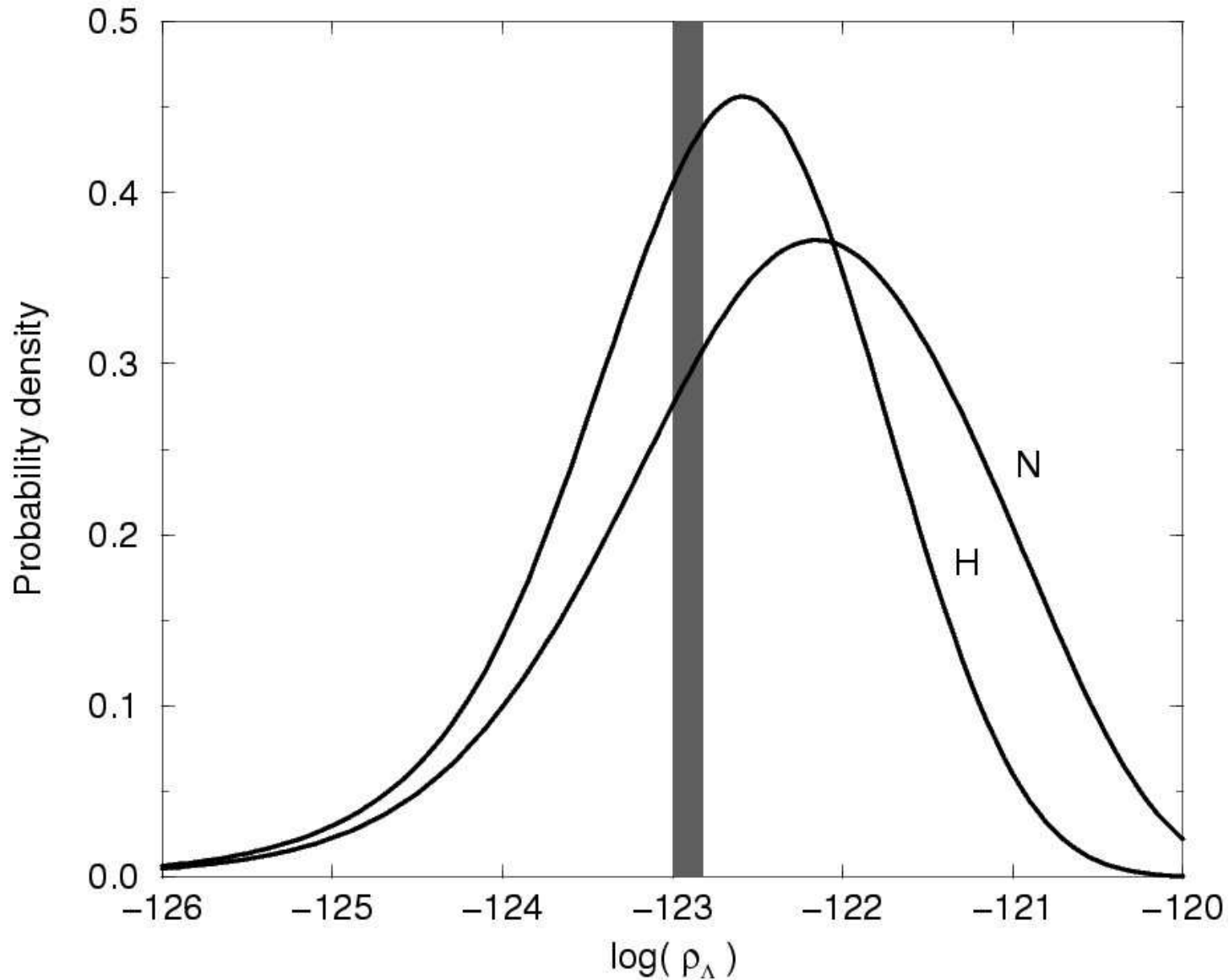
# Star Formation History, $\dot{\rho}_*$

Bouso *et al.* tried two published versions of  $\dot{\rho}_*$ , based on fits to data/simulations,



# Star Formation History, $\dot{\rho}_*$

leading to somewhat different predictions for  $dP/d\log\Lambda$ ,



# Star Formation History, $\dot{\rho}_*$

However, published SFR's are only valid for value of  $\Lambda$  in *our* universe. How to generalize to other values of  $\Lambda$ ?

Bousso *et al.* apply Press-Schechter formalism to  $\dot{\rho}_*$  as an overall multiplicative factor,

$$\dot{\rho}_* \rightarrow F(\Lambda)\dot{\rho}_*$$

$F$  quantifies how efficiently matter collapses as structure forms



# Press-Schechter Formalism

Using: statistics of Gaussian distribution of  $\delta\rho$  fluctuations,  
how fluctuations grow in spherical collapse model;  
Obtain: fraction of matter collapsed into DM halos  $\geq M$ ,

$$F(M, t) = \text{erfc} \left[ \frac{\delta_c}{\sqrt{2} \sigma(M, t)} \right]$$

$\sigma$  = r.m.s. fluctuations of mass  $M$

$\delta_c = 1.686$  from studying growth of perturbations

To model  $\Lambda$ -dependence of  $\dot{\rho}_*$ , Bousso *et al.* rescale

$$\dot{\rho}_* \rightarrow F(10^7 M_\odot, t_{\text{max}}) \dot{\rho}_*$$

$t_{\text{max}}$  = time when  $\dot{\rho}_*$  is maximized,  $10^7 M_\odot$  = typical  $M_{\text{galaxy}}$

But this is a crude approximation ...

# Hernquist-Springel SFR

Hernquist and Springel (2003) derived star formation rate from first principles + fits to simulations,

$$\dot{\rho}_* \propto \left[ \frac{\chi \bar{\chi}}{(\chi^m + \bar{\chi}^m)^{1/m}} \right]^{9/2\eta} \frac{1}{1 + \frac{5}{2} \sqrt{\pi} u \exp(u^2)}$$

where

$$\chi = \left( \frac{H(t)}{H_0} \right)^{2/3}, \quad u \propto \sigma_4(t)$$

and  $\sigma_4 =$  r.m.s. density fluctuation at  $M = M_4$  where  $T_{\text{virial}} = 10^4$  K.

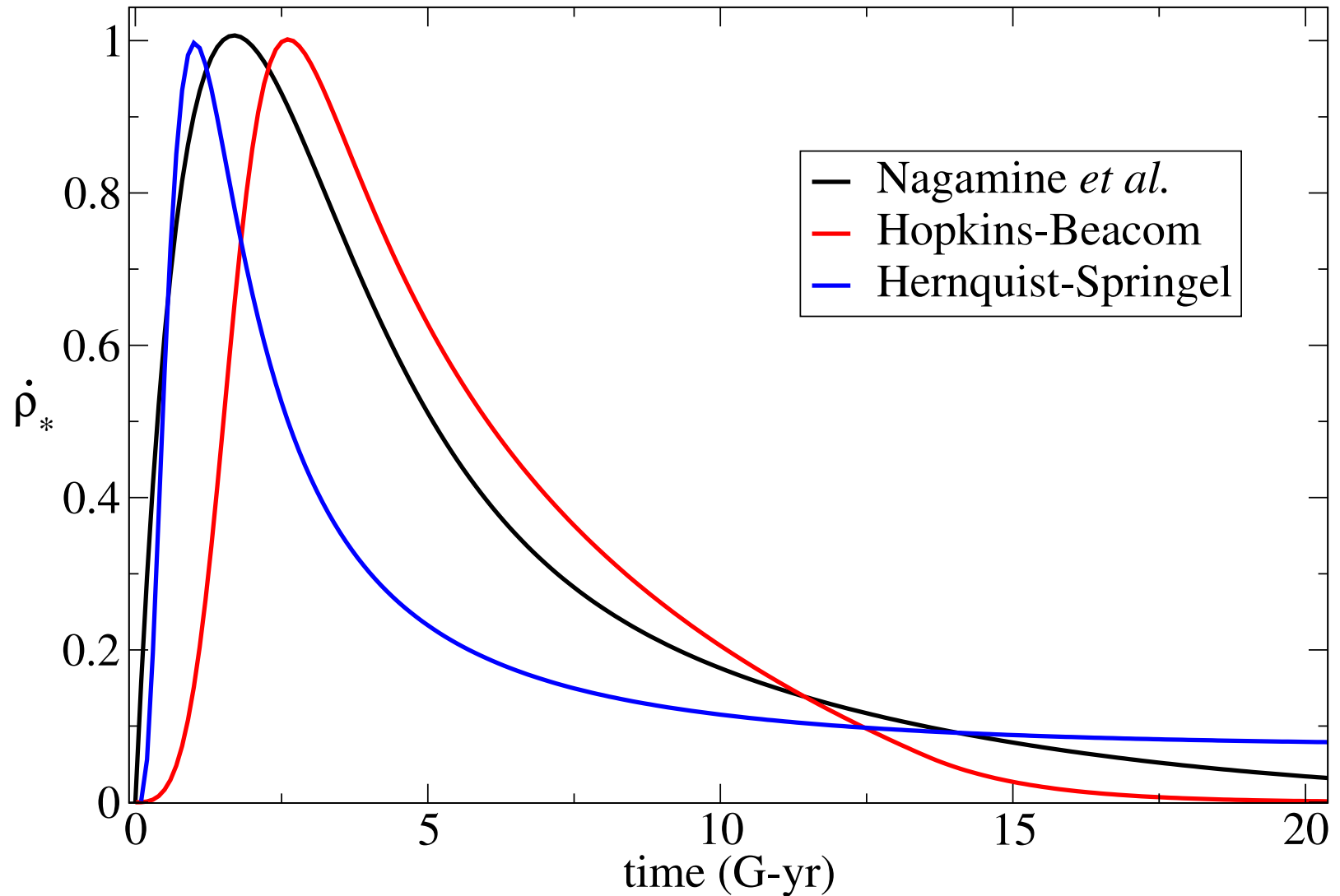
Tegmark, Aguirre, Rees, Wilczek (2005) show that

$$\sigma(M, t) \propto Q \rho_{\Lambda}^{-1/3} f(x(t)) g(M), \quad x(t) \equiv \frac{\rho_{\Lambda}}{\rho_m(t)}$$

Note  $\Lambda \sim Q^3$  to keep  $\sigma$  constant—large  $Q$  greatly weakens anthropic bound

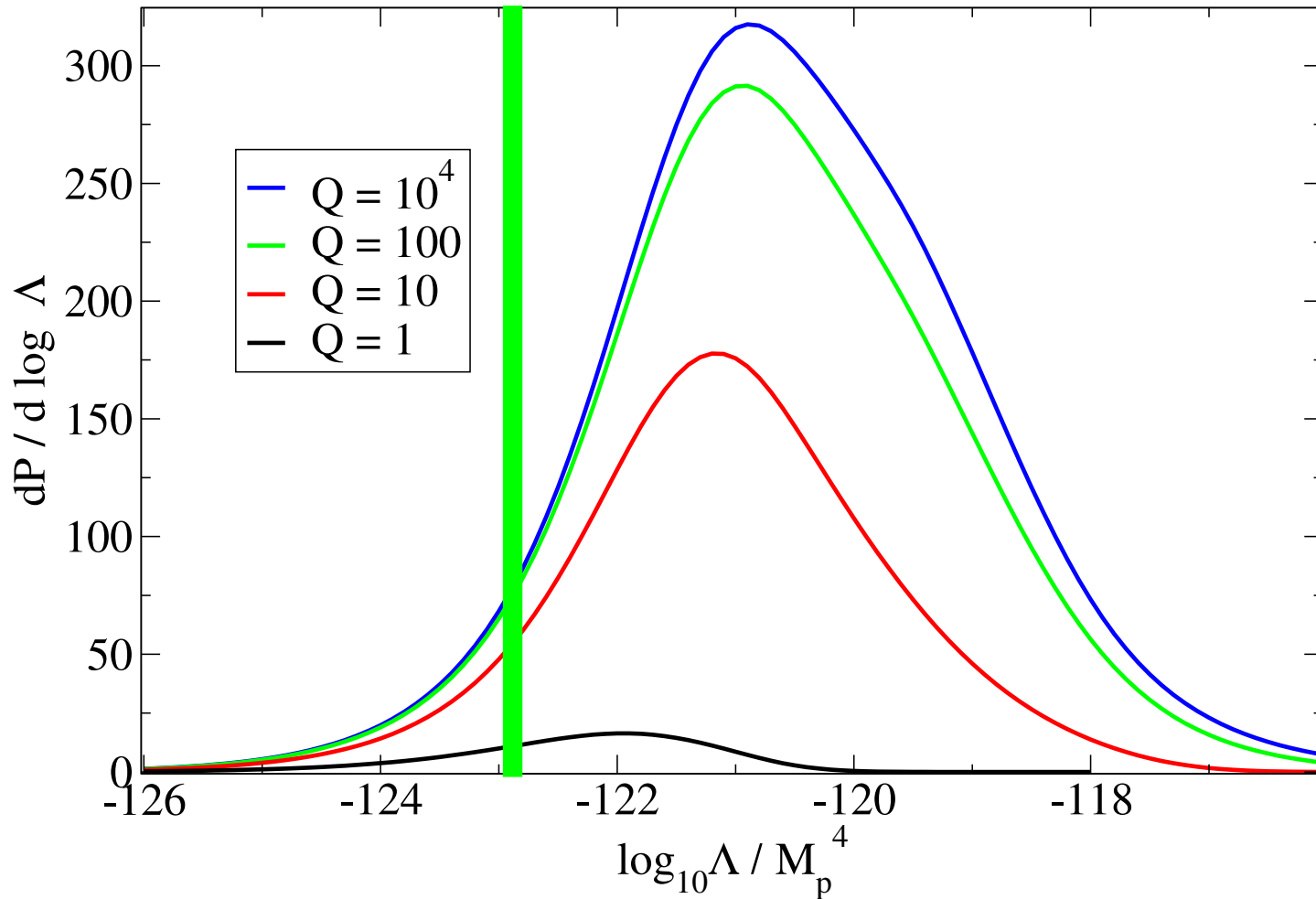
# Comparison of SFR's

The three different star formation rates considered:



# Dependence on Density Contrast

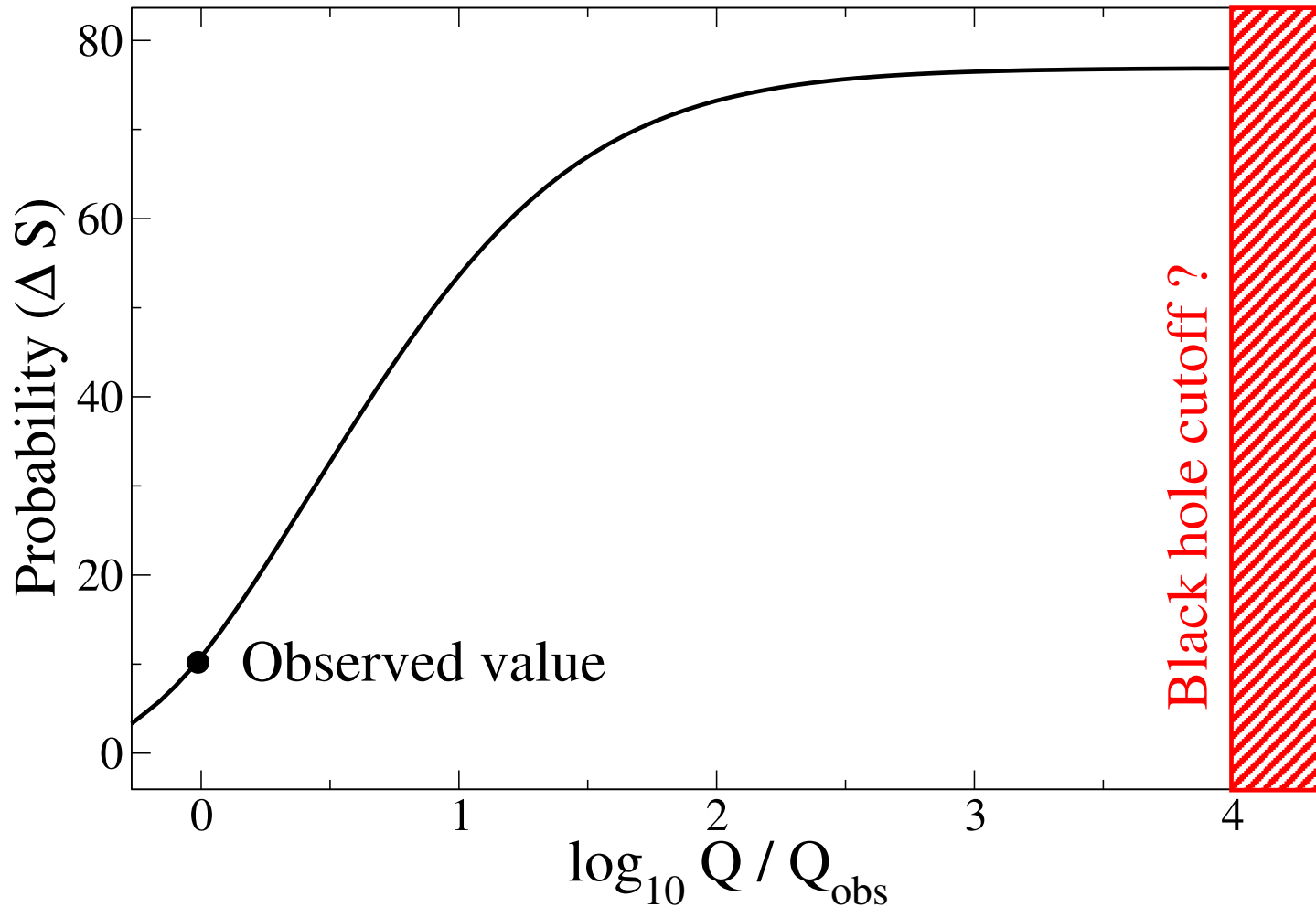
We can use the H-S SFR to see effect of varying  $Q$  (renormalized to  $Q = 1$  in our universe):



Most likely value of  $\Lambda$  is fairly insensitive to value of  $Q$ —unlike anthropic argument where  $\Lambda \sim Q^3$ .

# Probability distribution for $Q$

Keeping  $\Lambda = \Lambda_{\text{obs}}$ ,



Observed value is  $\sim 8$  times less likely than at high  $Q$

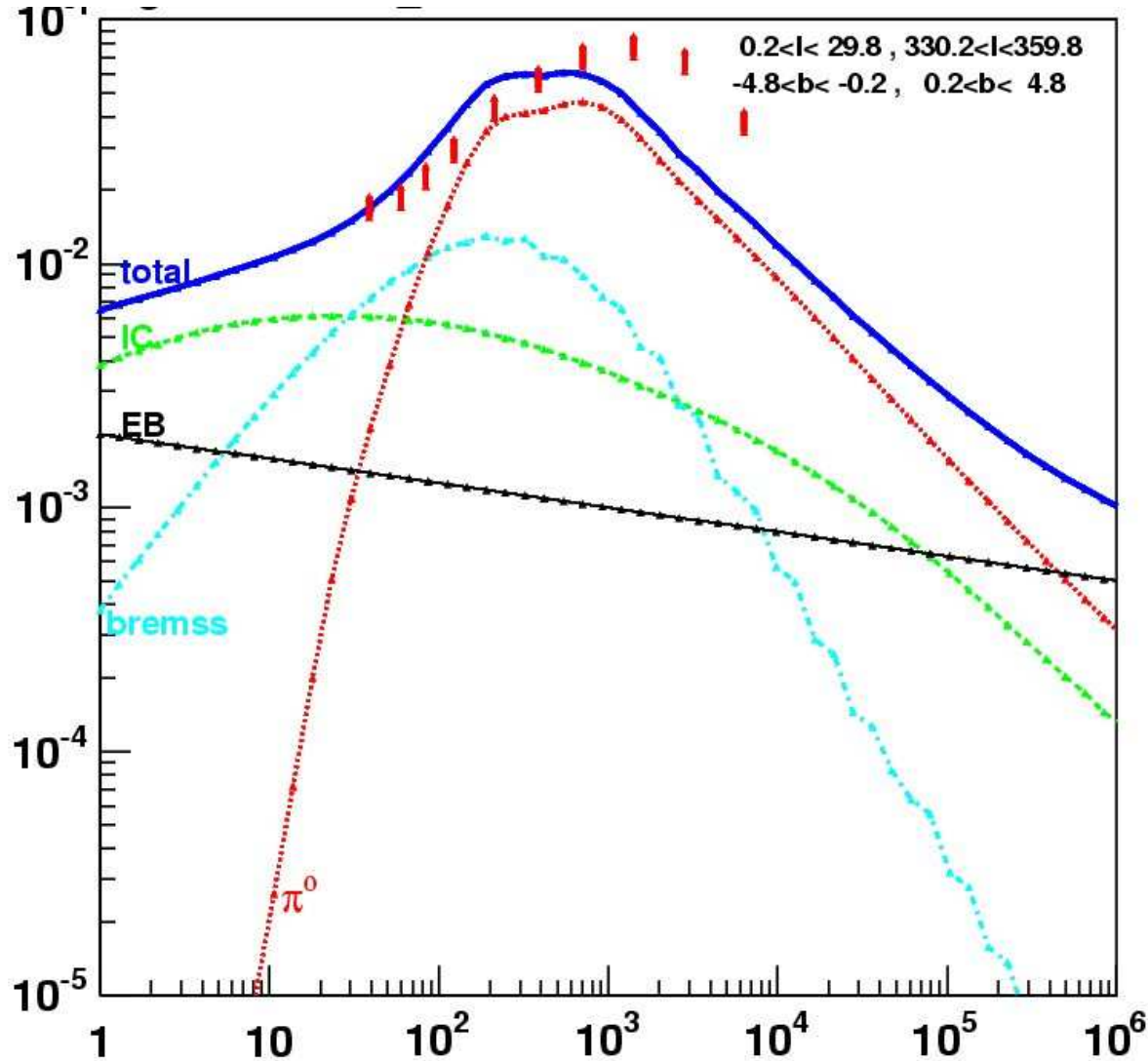
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Dark matter can also be a source of entropy if it decays

# Hints of decaying dark matter

- EGRET observes excess  $\gamma$  rays with  $E \sim 2 - 10$  GeV.

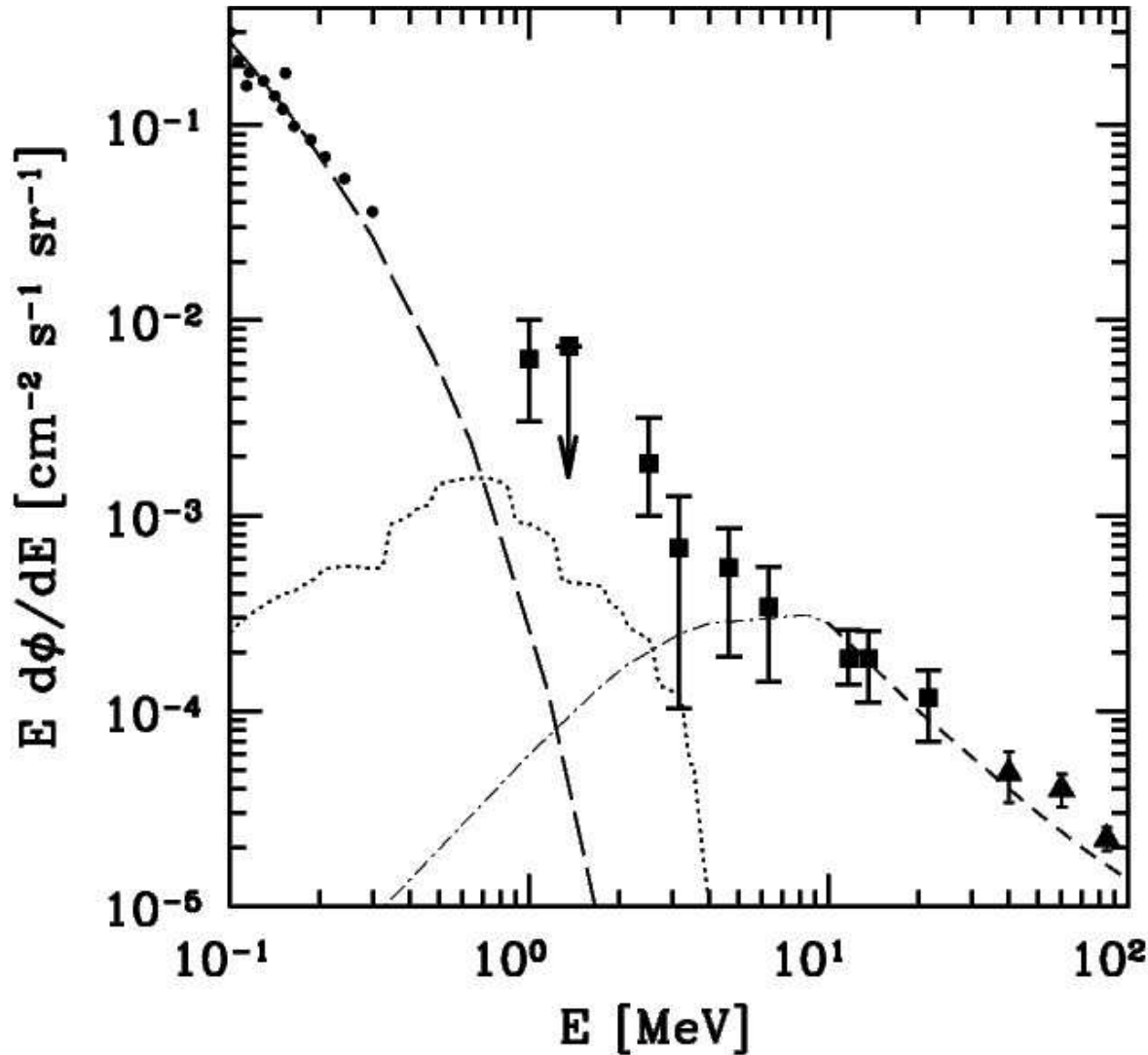


From de Boer *et al.*,  
astro-ph/0408272.

Buchmuller, Covi,  
Hamaguchi, Ibarra,  
Yanagida,  
hep-ph/0702184,  
interpret as due to  
gravitino DM decays

# Hints of decaying dark matter

- COMPTEL observes excess  $\gamma$  rays with  $E \sim 1 - 5$  MeV.



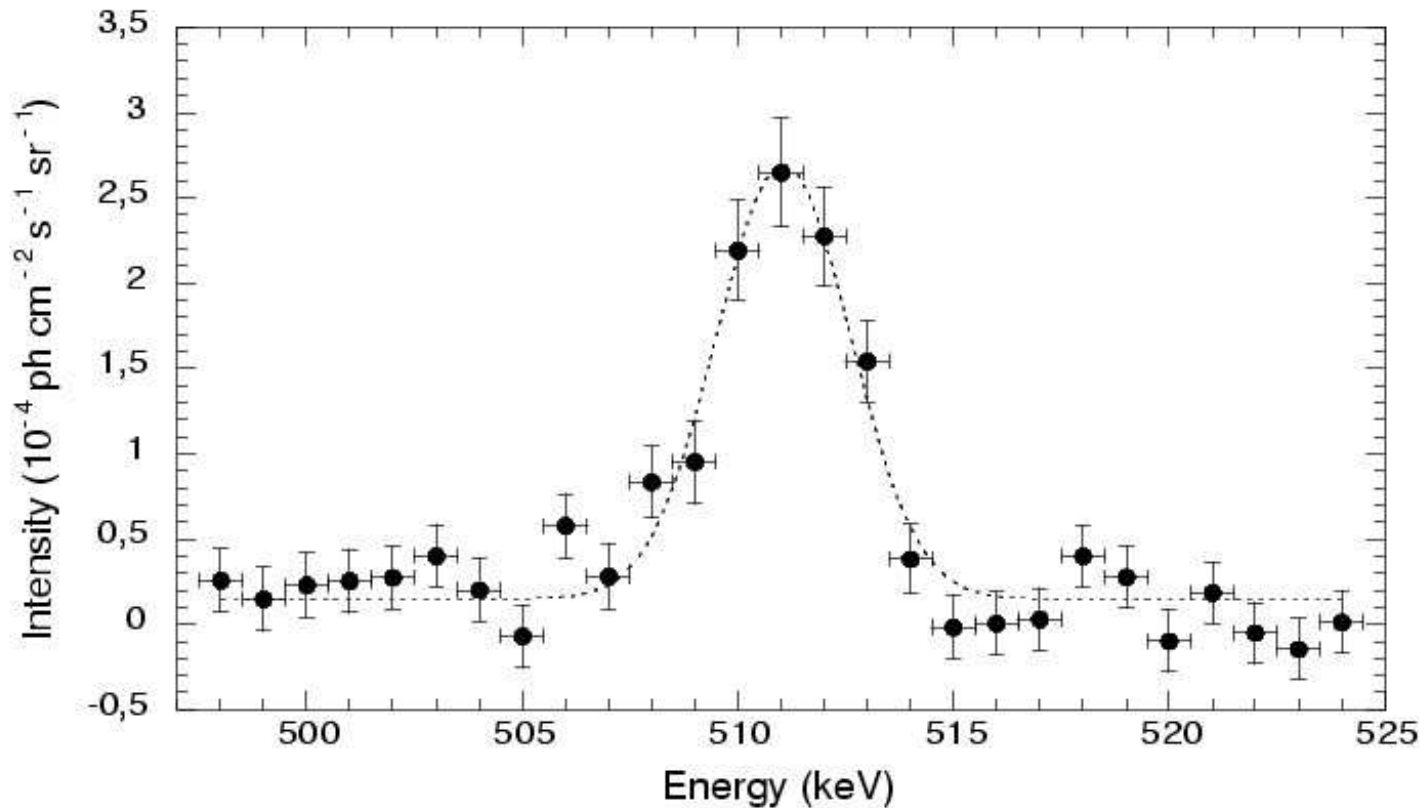
From Cembranos,  
Feng, Strigari,  
astro-ph/0704.1658

They consider KK  
and SUSY DM with  
small  $\Delta m$ ,  
 $\Delta m = M_{DM^*} - M_{DM}$



# Hints of decaying dark matter

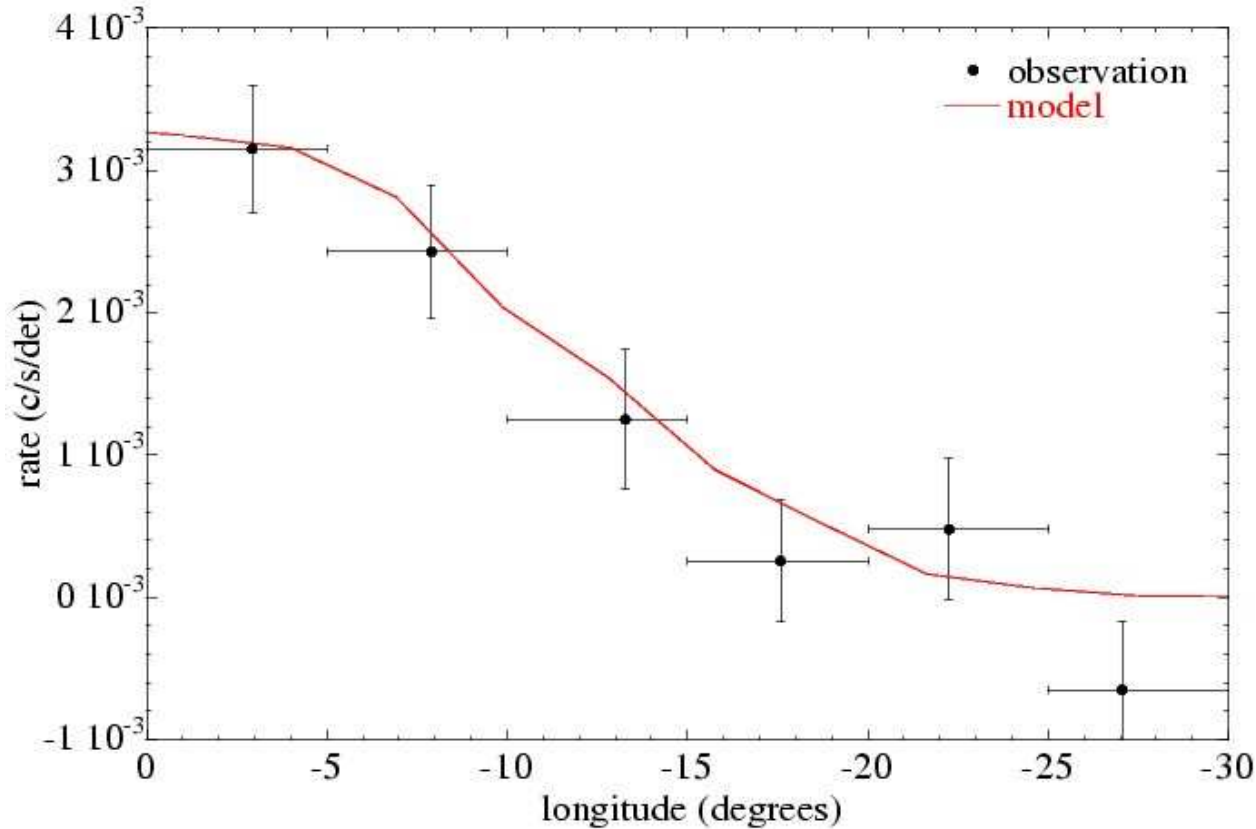
- INTEGRAL/SPI observes excess 511 keV  $\gamma$ 's from galactic center



**Fig. 3.** 511 keV flux spectrum obtained using a gaussian centred on the GC with a FWHM of  $10^\circ$ .

# Hints of decaying dark matter

- INTEGRAL/SPI observes excess 511 keV  $\gamma$ 's from galactic center



**Fig. 2.** Rate induced by galactic 511 keV photons as a function of longitude. The response to a gaussian source (FWHM =  $10^\circ$ ) is also shown for comparison.

# Hints of decaying dark matter

- INTEGRAL/SPI observes excess 511 keV  $\gamma$ 's from galactic center

$\sim$  MeV (or weak scale\*) DM decaying into  $e^+e^-$  has been suggested,

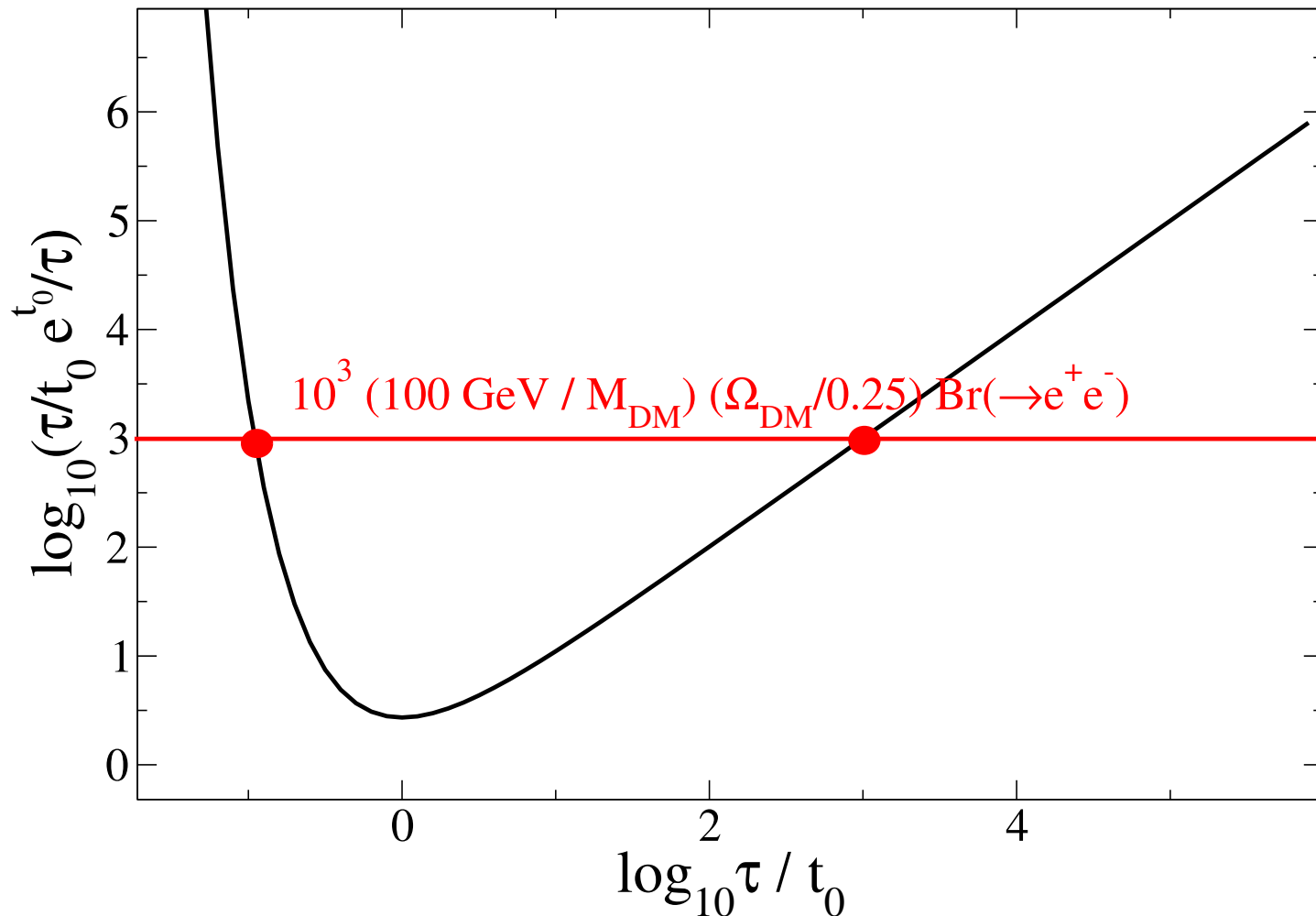
Picciotto, Pospelov	hep-ph/0402178:	sterile neutrinos
Hooper, Wang	hep-ph/0402220:	axinos
Kawasaki, Yanagida	hep-ph/0505167:	moduli
Kasuya, Takahashi	astro-ph/0508391:	Q-balls
Kasuya, Kawasaki	astro-ph/0602296:	moduli
*Pospelov, Ritz	hep-ph/0703128:	heavy excited WIMPs
Conlon, Quevedo	hep-ph/0705.3460:	string theory volume modulus

But why should DM be decaying only today?  $\Rightarrow$

# Lifetime needed for SPI anomaly

To match observed 511 keV line, need

$$\frac{\tau}{t_0} e^{t_0/\tau} \cong 10^3 \left( \frac{100 \text{ GeV}}{M_{DM}} \right) \left( \frac{\Omega_{DM}}{0.25} \right) \text{Br}(\text{DM} \rightarrow e^+ e^-)$$



# Hints of decaying dark matter

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*Pospelov, Ritz	hep-ph/0703128:	heavy excited WIMPs
Conlon, Quevedo	hep-ph/0705.3460:	string theory volume modulus

But why should DM be decaying only today?  
New coincidence problem

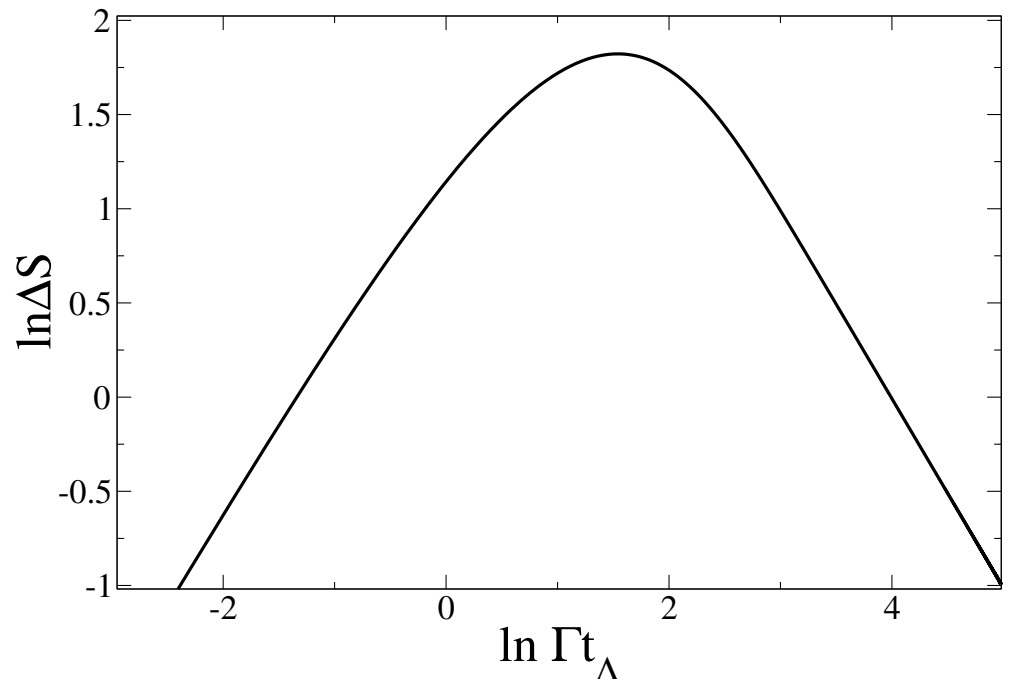
# Entropy from Decaying DM

Decaying DM could be major new source of entropy production in causal diamond. Simpler to compute than star contribution:

$$\frac{dS}{dV_c dt} = g_s \Gamma n e^{-\Gamma t}$$

Leads to entropy change  $\Delta S_d = t_\Lambda f(\Gamma t_\Lambda)$ , where  $t_\Lambda \equiv \sqrt{3/\Lambda}$

$f$  is maximized when  $1/\Gamma \cong t_\Lambda/4.7$ . Since  $t_\Lambda = 16.7$  Gyr, could explain why DM is decaying just today



# Decaying DM versus stars

How does  $\Delta S_d$  compare to  $\Delta S$  from stars? From decays,

$$\frac{dS_d}{dV_c dt} \cong 10^{66} g_s \left( \frac{\Gamma}{\text{Gyr}^{-1}} \right) \left( \frac{\text{eV}}{m} \right) \left( \frac{\rho_{DM}}{(10^{-3} \text{eV})^4} \right) \text{Mpc}^{-3} \text{y}^{-1}$$

where  $m$  = DM mass,  $g_s$  = number of decay products.

From stars, peak value is

$$\frac{dS_*}{dV_c dt} = 10^{63} \text{Mpc}^{-3} \text{y}^{-1}$$

$$\implies \Delta S_d > \Delta S_* \quad \text{if} \quad \frac{m}{g_s} \leq \text{keV}.$$

Need  $m \sim \text{MeV}$ ,  $g_s \sim 10^3$  to explain excess positrons. Such large  $g_s$  can come from synchrotron radiation by decay  $e^-$ 's.

If decay is GUT-suppressed, then  $\Gamma^{-1} \sim \frac{M_{\text{GUT}}^2}{m^3} \sim 10 \text{ Gyr}$

# Conclusions

- Causal entropic principle overcomes many deficiencies of anthropic principle
- Predicts  $\Lambda$  better than anthropic approach
- No problem from eternal inflation for defining measure on landscape
- No arbitrary assumptions about nature of observers
- Robust against variation of density contrast
- Can be applied to universes in the landscape which look quite different from ours
- Explanation for decaying dark matter?



# “cosmo-ph” archive

the cosmologically-oriented papers in astro-ph:

<http://www.physics.mcgill.ca/~jcline/astroph/>