Electroweak and Hints of New Phenomena at the Tevatron

Andrew Askew June 8, 2012





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- Specifically:
 - Dileptons (Z p_T and A_{FB})
 - Zγ and Search for GMSB in Zγ
 - Trileptons and multileptons (WZ,ZZ)



- Several recent EWK/NP searches (diboson, multilepton final states, analyses which are sensitive to the description from the SM, and physics beyond).
- High precision measurement of the W mass (which gives indirect sensitivity to the search for the Higgs).



Run II: The Fermilab Tevatron



20.00

5 35

65 95 125 155

185

Weekly Integrated Luminosity

Week #

(Week 1 starts 03/05/01)

215 245 275 305 335 365 395 425 455 485 515 545

- Bunch crossing 396 ns
- Peak Luminosity: 4.2x10³² cm⁻² sec⁻¹
- Run II total delivered: 12 fb⁻¹
- In this talk, 2.1-8.6 fb⁻¹

2000.00

0 00



The Experiments:



 No experimental talk is complete unless one shows our beloved detectors. I will come back to this too, when I talk about the W mass.



- In principle, a simple topology, but can yield insights into higher order effects from QCD and fundamental constants of the SM:
 - Z p_T distribution
 - A_{FB} , $sin^2\theta_W$





- Same 2.1 fb⁻¹ dataset as used for previously published rapidity distribution, and angular coefficients.
 - 66 < M_{ee} < 116 GeV</p>
 - Three categories: Both electrons |η|<1.1, One electron |η|<1.1 and one electron 1.2 <|η|<2.8, Both electrons 1.2 <|η|<2.8.





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PRD 84, 012007 (2011)



CDF also made a preliminary measurement of $sin^2\theta_w$ =0.2320 +/- 0.008^{+0.001}-0.0009 using 2.1 fb⁻¹



PRD 84, 012007 (2011)



 A quick comparison between the new DØ measurement and other measurements.



In diboson final states, the SM makes unambiguous statements about the couplings, so any deviation is a sign of new physics. Thus measuring these cross sections are both measurements and searches.





Zγ is also a final state that is sensitive to GMSB SUSY (as well as a heretofore unexploited state) A complement to the more typical channels (yy

+missing E_{T})





- In a similar vein, one can measure the other diboson cross sections, implicitly searching for new physics.
- These serve as both a test of the gauge structure of the SM, and an important legacy of the Tevatron.



arXiv:1201.5652, accepted by PRD



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In a different sense, one can choose the trilepton state (ee+l or track, μμ +l or track), and define control regions to ensure proper descriptions of the data. Then armed with that confidence, invade the areas where signals from SUSY are expected.



l with that	Definition of Control and Signal Regions				
invada tha	Region	$M_{\ell\ell} \ { m cut} \ ({ m Gev}/c^2)$	$(\not\!\!E_T) \operatorname{cut} (\operatorname{GeV})$	N_{jet} cut	
invaue the	Region0	$M_{\ell\ell} > 20$	$\not\!\!\!E_T < 10$	—	
signals	Region1	$76 < M_{\ell\ell} < 106$	$\not\!\!\!E_T > 15$	$N_{jet} \leq 1$	
	Region2	$76 < M_{\ell\ell} < 106$	$\not\!\!\!E_T > 15$	$N_{jet} \geq 2$	
are	Region3	$20 < M_{\ell\ell} < 76 \text{ or } M_{\ell\ell} > 106$	$\not\!\!E_T < 10$	$N_{jet} \leq 1$	
uic	Region4	$20 < M_{\ell\ell} < 76 \text{ or } M_{\ell\ell} > 106$	$\not\!\!\!E_T < 10$	$N_{jet} \geq 2$	
	Region5	$76 < M_{\ell\ell} < 106$	$\not\!\!\!E_T < 10$	$N_{jet} \leq 1$	
	Region6	$76 < M_{\ell\ell} < 106$	$\not\!\!\!E_T < 10$	$N_{jet} \geq 2$	
	Region7	$20 < M_{\ell\ell} < 76 \text{ or } M_{\ell\ell} > 106$	$\not\!\!\!E_T > 15$	$N_{jet} \leq 1$	
	Region8	$20 < M_{\ell\ell} < 76 \text{ or } M_{\ell\ell} > 106$	$\not\!\!\!E_T > 15$	$N_{jet} \geq 2$	
	Region9	$20 < M_{\ell\ell} < 76 \text{ or } M_{\ell\ell} > 106$	$\not\!\!\!E_T > 20$	$N_{jet} \leq 1$	
	Region10	$76 < M_{\ell\ell} < 106$	_	-	
Andrew Askew, EWK and Hints of NP at	Region 11	$M_{\ell\ell} > 20$	_	- 10	



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- In a different sense, one can choose the trilepton state (ee+l or track, µµ+l or track), and define control regions to ensure proper descriptions of the data.
- Then armed with that confidence, invade the areas where signals from SUSY are expected.
- Optimize final cuts for the desired brand of SUSY.

Optimization cuts			
$M_{\ell_1\ell_2}$	$> M_{\tilde{\chi}_{1}^{\pm}} - M_{\tilde{\chi}_{1}^{0}}$		
$M_{\ell_1\ell_3}$	$< 75 \ { m GeV}/c^2$		
$M_{\ell_2 \ell_3}$	$< 75 \ { m GeV}/c^2$		
₽ _T	$> 25 { m ~GeV}$		
$p_{T,2}$	(> 8 and < 36 - 65) GeV/c		
$p_{T,3}$	$> 8 ~{ m GeV}/c$		

Optimized Trilepton Yields for Benchmark						
Channel	SM background	SUSY signal	Observation			
ee+lepton	1.5 ± 0.4	8.0 ± 0.8	3			
ee+track	11.6 ± 1.7	7.6 ± 0.8	13			
$\mu\mu$ +lepton	0.5 ± 0.1	6.7 ± 0.7	0			
$\mu\mu + track$	3.6 ± 1.0	6.2 ± 0.6	3			



- At first explicitly a search for high mass resonances decaying to ZZ.
 - In addition, the four charged lepton channel actually appeared to hint at a bona fide excess.
 - Naturally, the thing to do would be to look in alternative ZZ channels: llvv, and lljj.





- Subsequent checks in the alternatives revealed no excesses.
 - Note on the plots displayed that the empty histogram would be representative of a similar signal to the implied excess.



PRD 85, 012008, (2012)

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PRD 85, 012008, (2012)

- Subsequent checks in the alternatives revealed no excesses.
 - Shown are the corresponding limit plots, nothing to see here.



W Mass Measurement

Precise measurements of M_{W} and M₊ constrain the SM Higgs mass. Prior to this year, the status looked something like the plot on the right (from the W mass point of view).





Methodology

This entire business sounds easy in general, but the level of detail required to reach the desired precision on the mass is extreme, and requires an attention to detail unlike most other hadron collider measurements.





Methodology



Both experiments will follow the same general strategy:

- Build up a model of the detector and the inactive material, verify lepton energy scales and recoil energy resolution.
- Then using this model, generate templates of the W mass lineshape and compare with data. Both experiments also do this procedure blinded. But methods for doing the above are specific to the strengths of the different detectors.



- CDF starts with its very precise tracking.
 - Verify tracking alignment with cosmics.
 - Momentum scale and linearity from J/ ψ and Ψ
 - Perform independent measurement of $Z \rightarrow \mu \mu$ mass as an important scale check.





Once the tracking scale is nailed, can then transfer this to the calorimeter through E/p measurement (energy loss, radiative component and width), and in bins of E_{τ} . Can then likewise check against description of the Z peak.



PRL 108, 151803 (2012)



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 Use M_T, E_T and missing E_T for both electron and muon channels, and then combine taking into account correlations.

W-boson mass (MeV)

 $\begin{array}{l} 80\,408\,\pm\,19_{stat.}\,\pm\,18_{syst.} \\ 80\,393\,\pm\,21_{stat.}\,\pm\,19_{syst.} \end{array}$

 $80\,431 \pm 25_{\rm stat.} \pm 22_{\rm syst.}$

 $80\,379 \pm 16_{\rm stat} \pm 16_{\rm syst}$

 $\begin{array}{l} 80\,348\,\pm\,18_{\rm stat.}\,\pm\,18_{\rm syst.} \\ 80\,406\,\pm\,22_{\rm stat.}\,\pm\,20_{\rm syst.} \end{array}$



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Distribution

 $m_T(e, \nu)$

 $m_T(\mu, \nu)$

 $p_T^{\ell}(e)$

 $p_T^{\nu}(e)$

 $p_T^{\ell}(\mu)$

 $p_T^{\nu}(\mu)$



Reach a final measurement of M_w=80387 +/- 19 MeV, using 2.2 fb⁻¹.



PRL 108, 151803 (2012)



W Mass: DØ



 As previously stated, the game is the same, but the method is different (solely studying W→ev here). Tracking is not as precise, and there is a lot of material to account for prior to the EM calorimetry.

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- Sampling calorimeter, very sensitive to the amount of material prior to calorimeter.
 Allow dead material and the energy scale of each layer to float, and minimize the global data/MC χ², using
 - $Z \rightarrow$ ee events split into categories based on rapidity.



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W Mass: DØ



 Cross check this, to make sure that the description is robust (done for 1 fb⁻¹ analysis and repeated for subsequent 4.3 fb⁻¹)



Energy scale and offset (with the corrected material description), verified for different ranges of instantaneous luminosity (potential systematic effects from pileup).





W Mass: DØ



 After taking into account correlations, even though all three variables are used, only p_T and M_T contribute to the final measured value value (when combined with earlier 1 fb⁻¹ result): M_W=80375+/-23 MeV



W Mass: DØ

	ΔM_W (MeV)		
Source	m_T	p_T^e	₿ _T
Electron energy calibration	16	17	16
Electron resolution model	2	2	3
Electron shower modeling	4	6	7
Electron energy loss model	4	4	4
Hadronic recoil model	5	6	14
Electron efficiencies	1	3	5
Backgrounds	2	2	2
Experimental subtotal	18	20	24
PDF	11	11	14
QED	7	7	9
Boson p_T	2	5	2
Production subtotal	13	14	17
Total	22	24	29

• After taking into account correlations, even though all three variables are used, only p_T and M_T contribute to the final measured value (when combined with earlier 1 fb⁻¹ result): $M_W = 80375 + 723$ MeV



W Mass: Tevatron

- Not yet the final word for W mass at the Tevatron, but here is the new combined value.
- Tevatron Run II results now dominate the precision on the world average.



arXiv: 1204.0042



- The search for new physics at the Tevatron continues in many directions!
- I've barely scratched the surface of all of these topics, and there were still more I couldn't cover because of time.
- Check out the references that I've strewn through this talk, as well as:
 - http://www-cdf.fnal.gov/physics/physics.html
 - http://www-do.fnal.gov/Run2Physics/WWW/ results.htm





W+jets bump



 Here are the dijet invariant mass distributions for W+jets.



W+jets bump



Here we are with background subtracted. Word is that CDF will have an additional update soon.