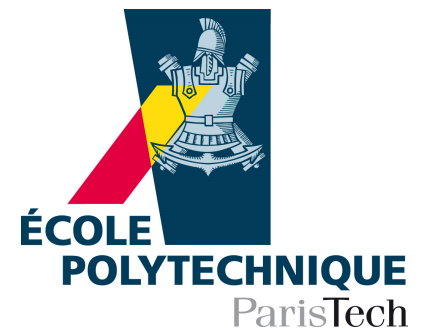
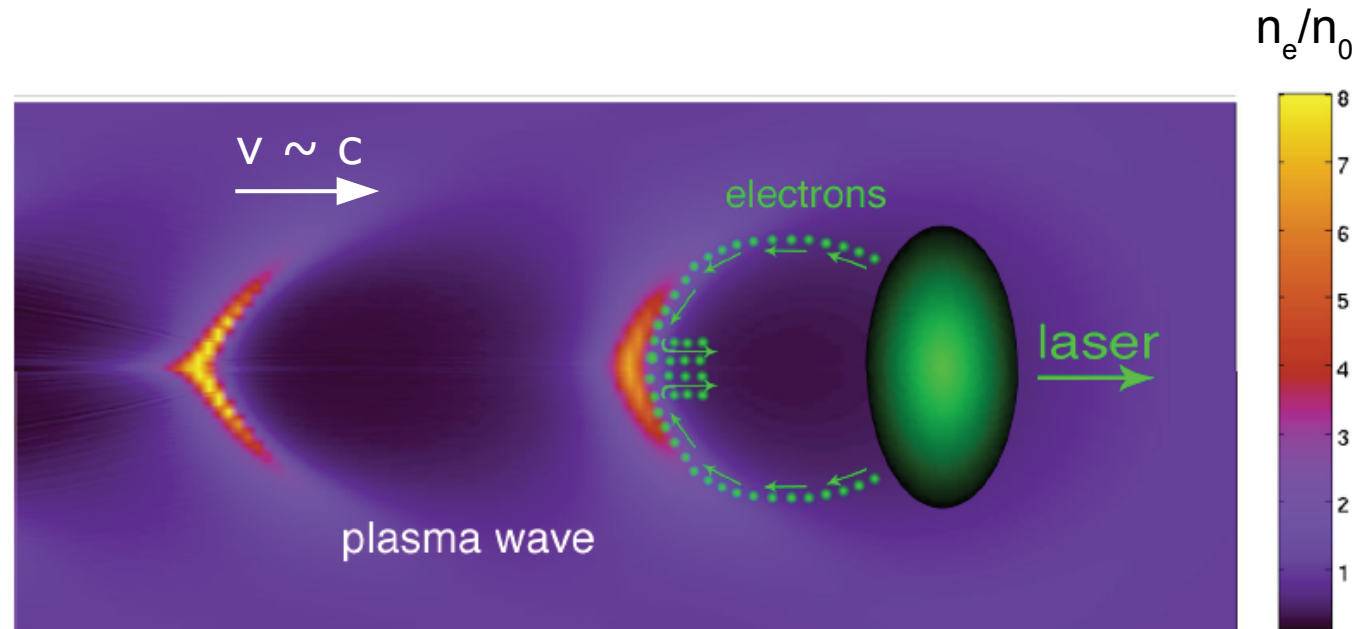
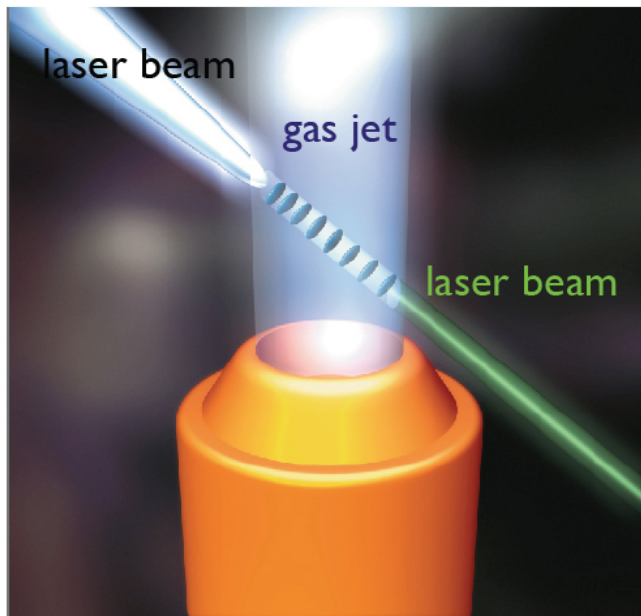

Emittance measurement in laser plasma accelerators

C. Thaury, S. Corde, K. Ta Phuoc, X. Davoine, R. Lehe
A. Rousse and V. Malka

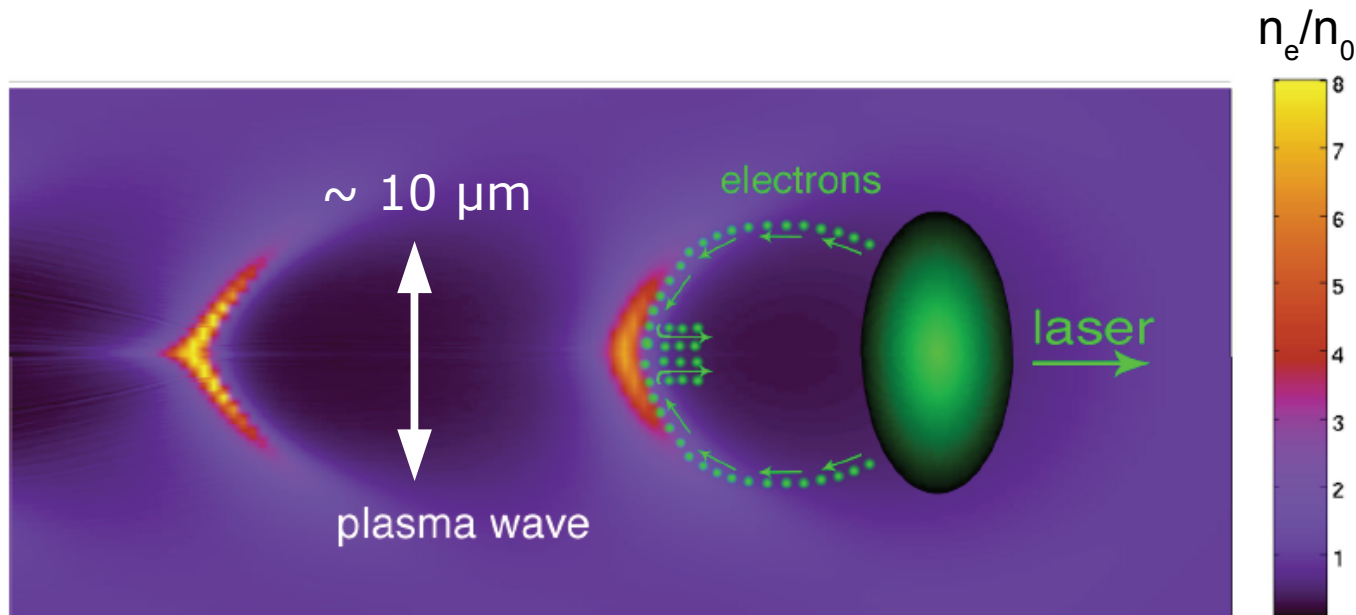


Laser-plasma acceleration



- The rising edge of the laser ionizes the gas and creates a plasma.
- The laser pulse triggers electric fields ~ 100 GV/m in its wake.
- Trapped electrons are accelerated up to ~ 100 MeV – 1 GeV within few mm.

Laser-plasma acceleration

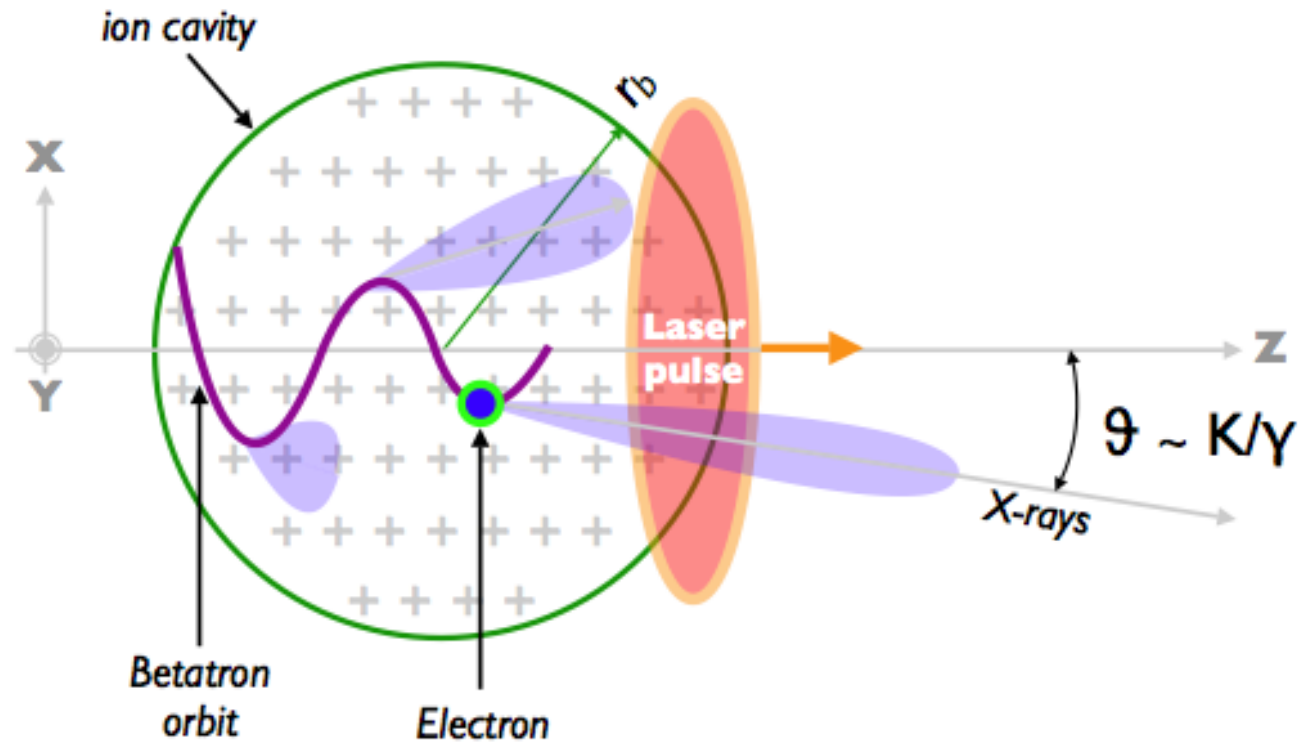


Main properties :

- Up to 1 GeV electron energy.
- ~ 10 -100 pC charge, fs duration, $\Delta E/E \sim 1$ -10%.
- Transverse size $\sim \mu\text{m}$, divergence $\sim \text{mrad}$
➔ $\epsilon_N \sim \text{p.m.m.mrad}$.

Because of the large divergence, conventional emittance measurement techniques are not suitable.

Betatron emission in laser accelerator



Longitudinal force Transverse force

$$\frac{d\vec{p}}{dt} = \vec{F}_{\parallel} + \vec{F}_{\perp}$$

$$\vec{F}_{\perp} = -\frac{1}{2} m \alpha \omega_p^2 r \vec{e}_r$$



**Harmonic oscillator
in the transverse
direction**

Betatron emission in laser accelerator

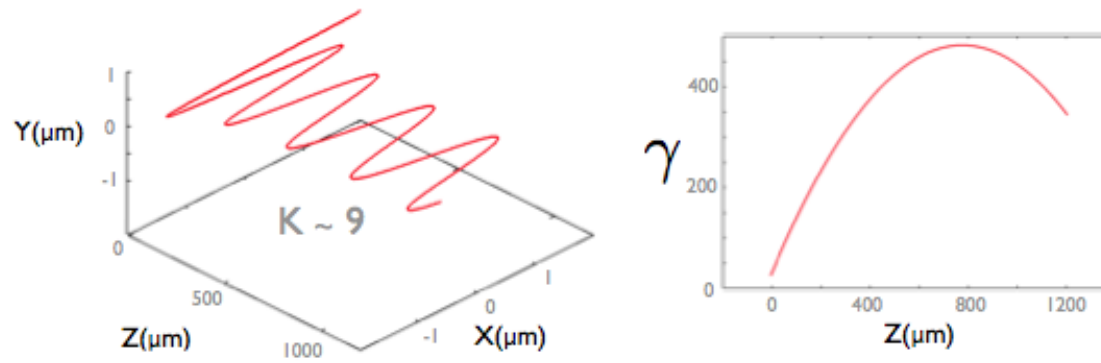
Betatron oscillations for an adiabatic acceleration

$$\hat{x}(t) = \sqrt{\hat{J}} \sin \left[\int_0^t \omega_\beta(t') dt' + \phi \right]$$
$$\hat{p}_x(t) = \sqrt{\hat{J}} \cos \left[\int_0^t \omega_\beta(t') dt' + \phi \right]$$

where

$$\hat{x} = \left(\frac{\alpha\gamma}{8} \right)^{1/4} k_p x$$
$$\hat{p}_x = \frac{p_x}{(2\alpha\gamma)^{1/4} mc}$$
$$\hat{J} = \frac{k_p J}{mc} = \hat{x}^2 + \hat{p}_x^2$$

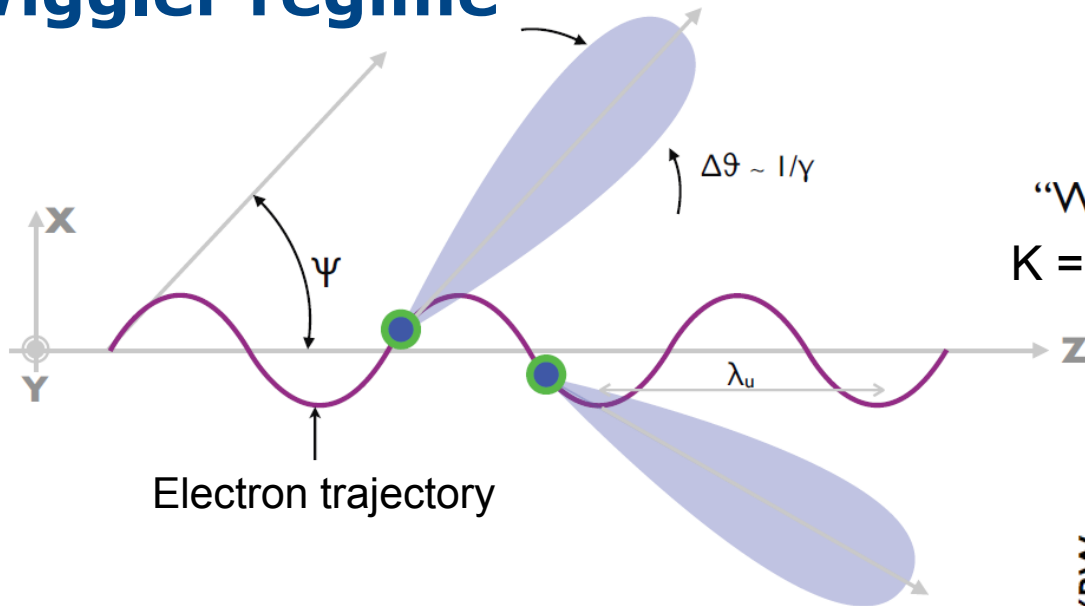
➔ Sinusoidal oscillations with a time varying frequency $\omega_\beta = \sqrt{\alpha\omega_p^2/2\gamma(t)}$



Plasma cavity = undulator

Betatron emission in laser accelerator

Wiggler regime



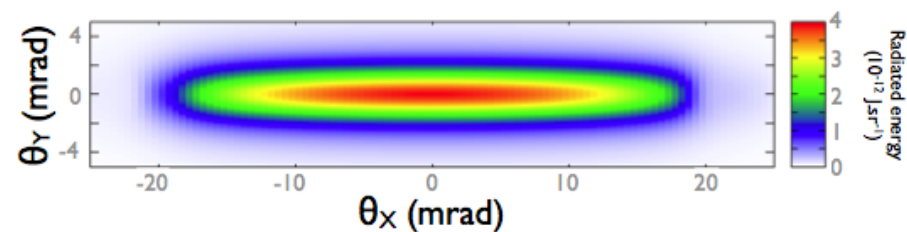
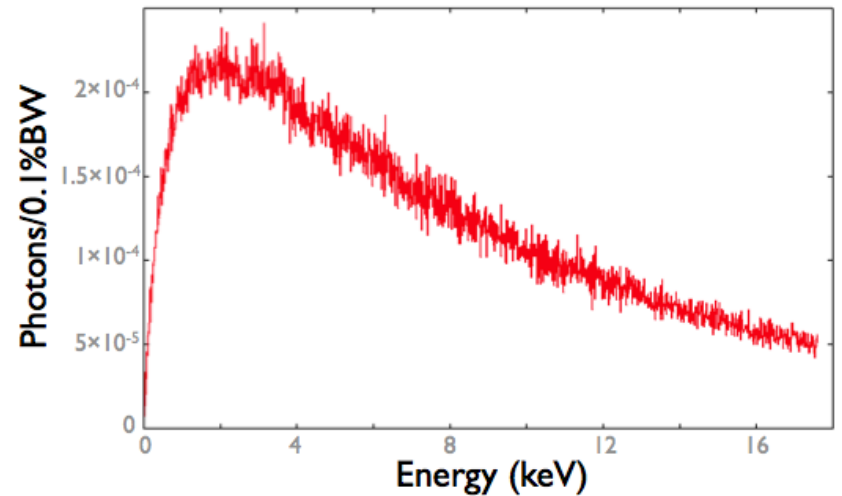
X-ray properties

$$E_c = \frac{3\hbar\omega_p^2}{4c} \gamma^2 \alpha r_\beta$$

$$\theta = \frac{p_\beta}{\gamma mc} = \frac{k_p}{\sqrt{2}} \gamma^{-1/2} \alpha^{1/2} r_\beta$$

$$r_\beta = \sqrt{\hat{J}} \left(\frac{\alpha\gamma}{8} \right)^{-1/4} k_p^{-1}$$

$$p_\beta = \sqrt{\hat{J}} (2\alpha\gamma)^{1/4} mc$$



Principle of emittance measurement

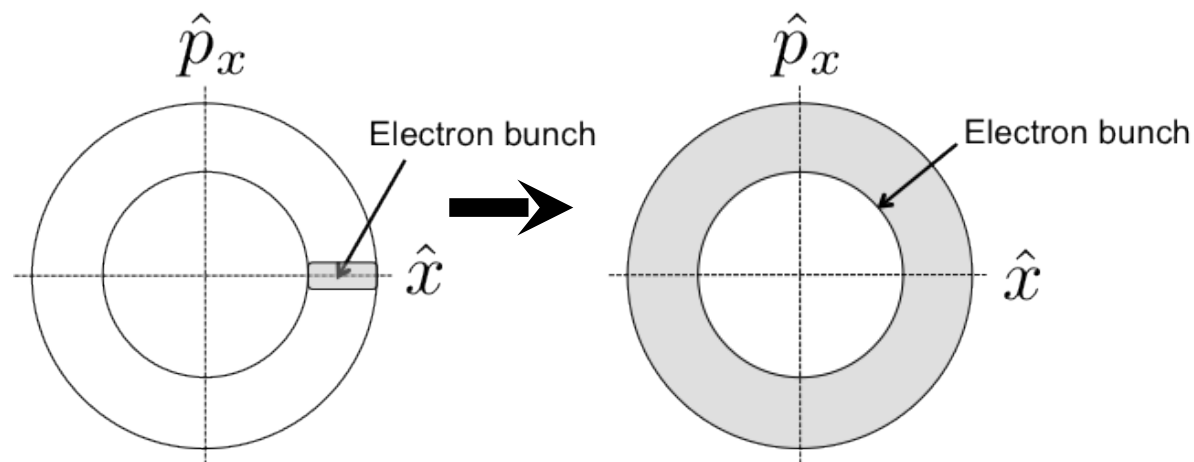
Three variables (γ , α , ϵ_N) \Rightarrow three measurements

- Electron spectrum $\Rightarrow \gamma$
- X-ray angular profile $\Rightarrow p_\beta / (\gamma m c) \propto \epsilon_N^{1/2} \alpha^{1/4} \gamma^{-3/4}$
- X-ray spectrum $\Rightarrow \gamma^2 \alpha r_\beta \propto \epsilon_N^{1/2} \gamma^{7/4} \alpha^{3/4}$
- X-ray source size $\Rightarrow r_\beta \propto \epsilon_N^{1/2} (\alpha \gamma)^{-1/4}$

Principle of emittance measurement

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- **X-ray source size** $\Rightarrow r_\beta \propto \epsilon_N^{1/2} (\alpha \gamma)^{-1/4}$
- Symmetrized emittance = non-coherent upper limit of the normalized emittance.



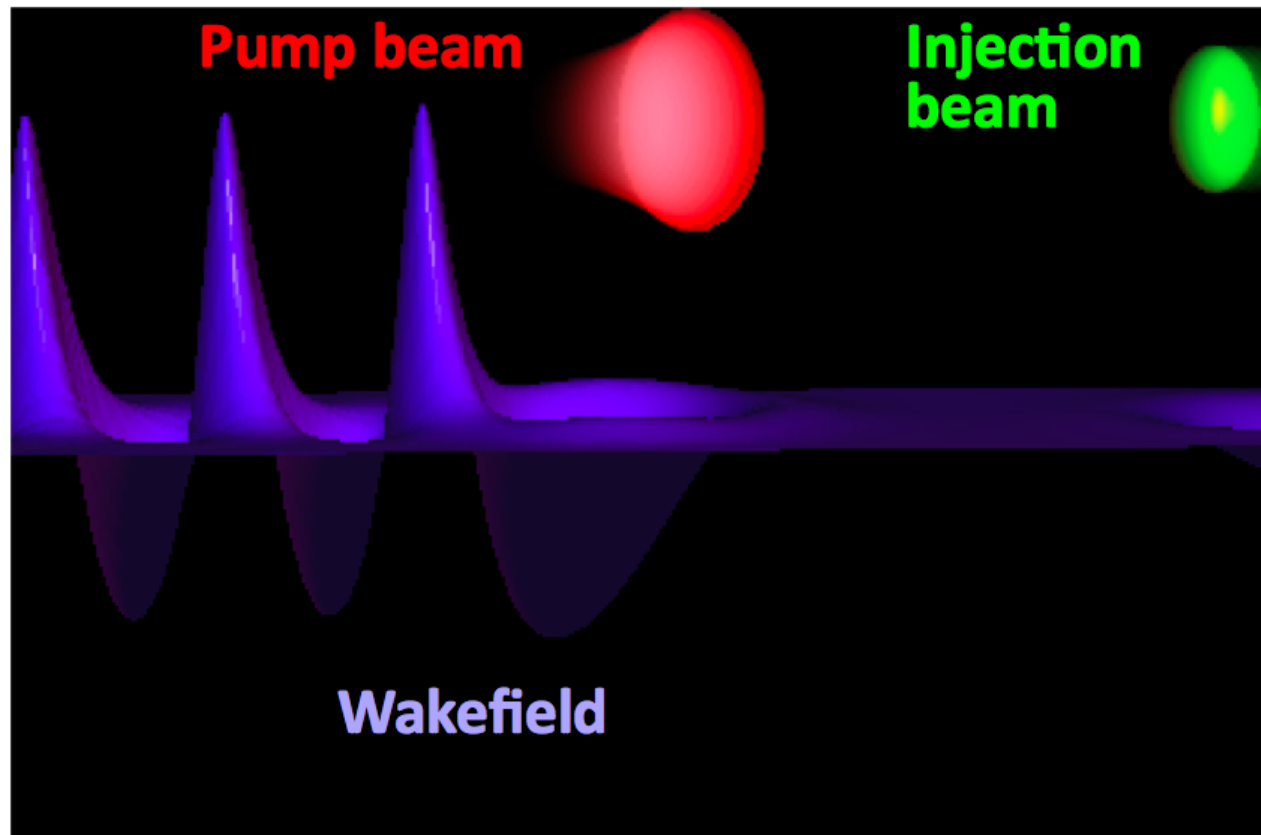
Principle of emittance measurement

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- **Electron spectrum** $\Rightarrow \gamma$
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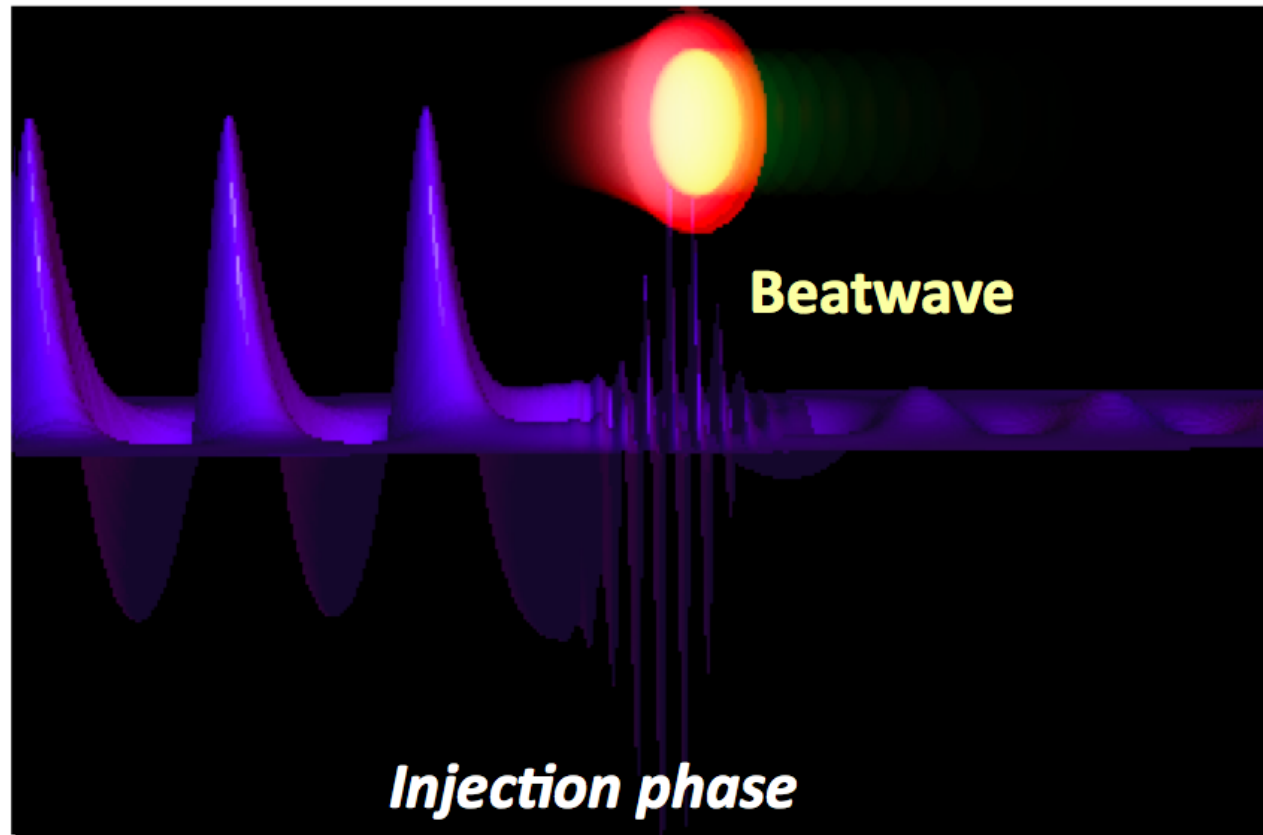
A precise measure requires a good stability and a control over the acceleration

Colliding injection



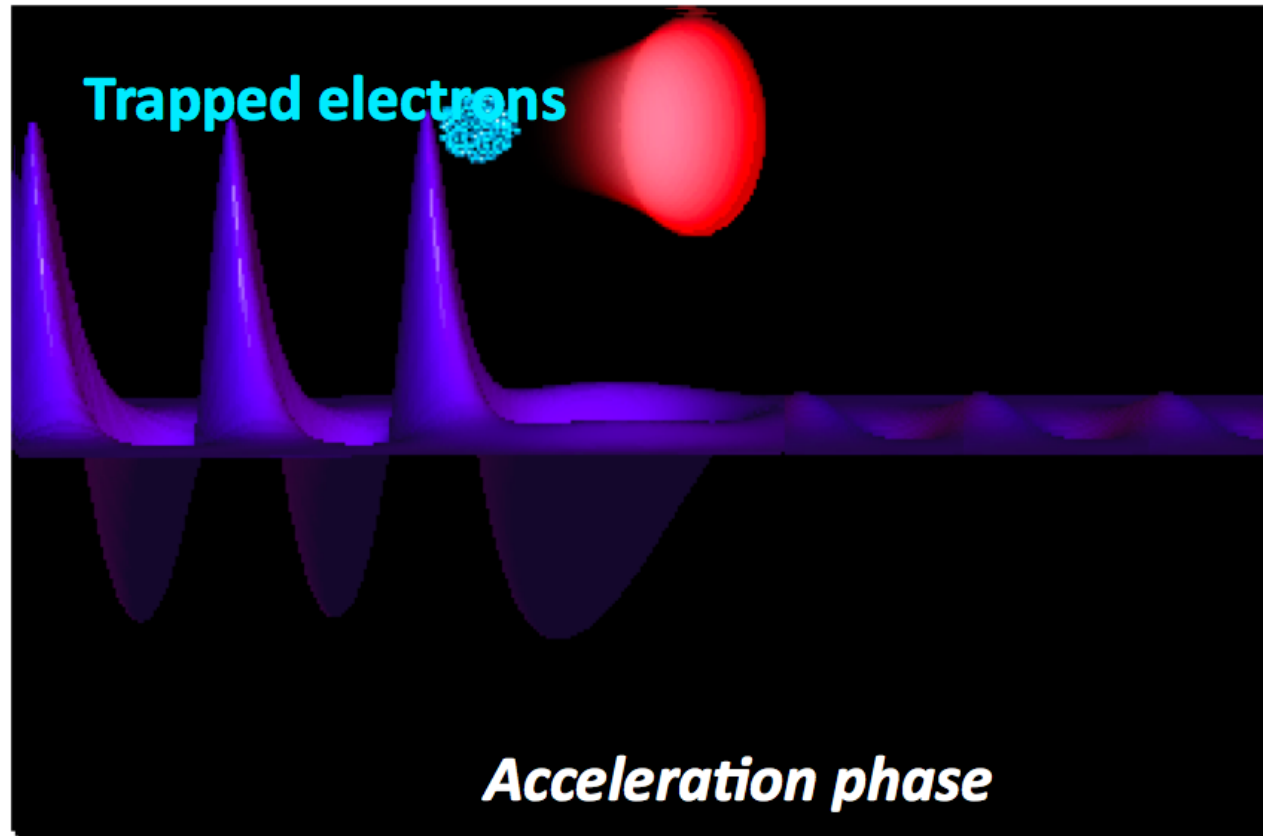
- Pump beam ➡ accelerating structure.
- Injection beam ➡ local injection.

Colliding injection



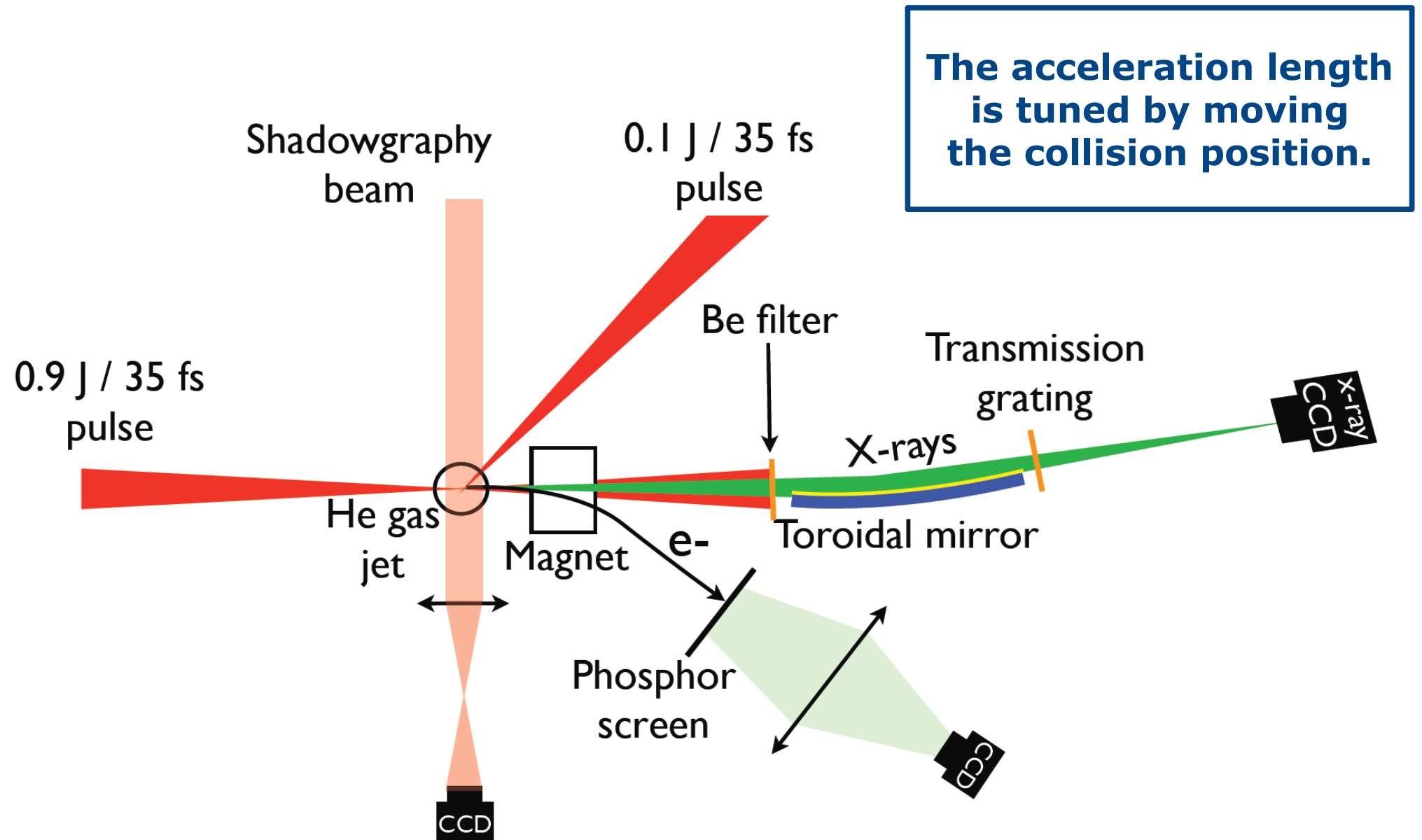
- Pump beam ➡ accelerating structure.
- Injection beam ➡ local injection.
- During the collision, some electrons are heated by the beat-wave ponderomotive force

Colliding injection

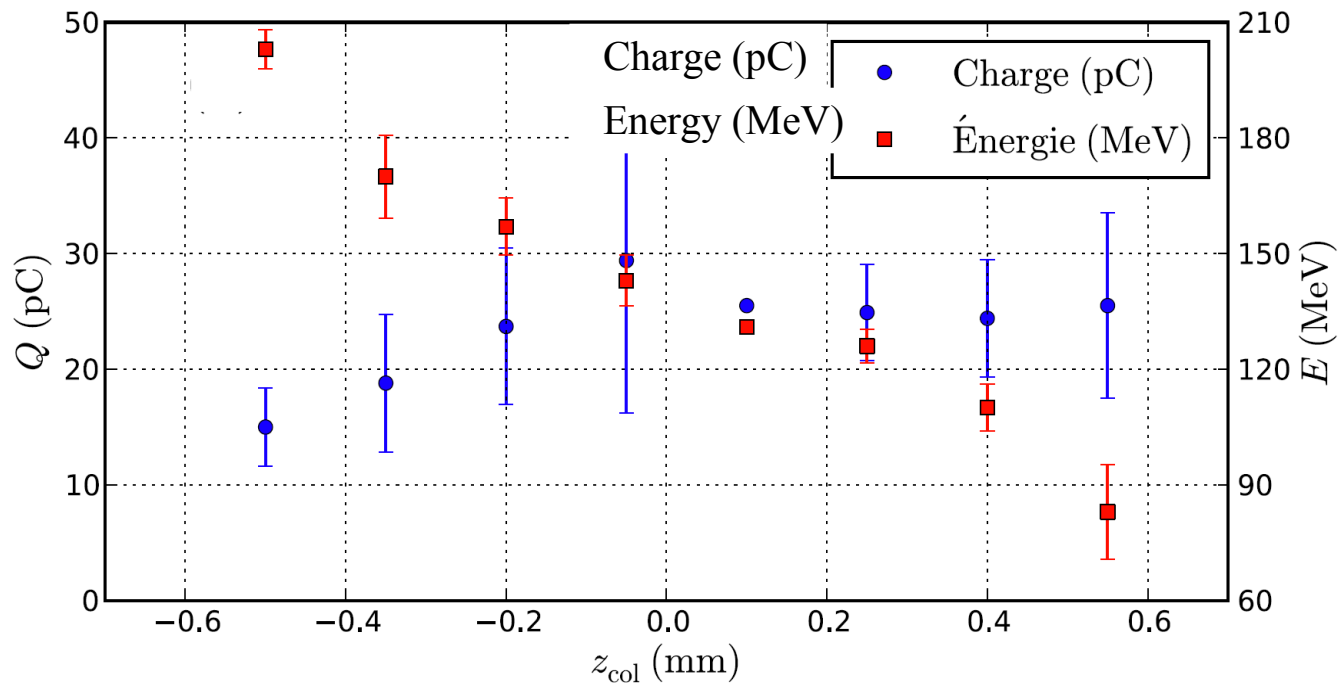


- Pump beam ➡ accelerating structure.
- Injection beam ➡ local injection.
- During the collision, some electrons are heated by the beat-wave ponderomotive force ➡ they gain enough energy to be trapped.

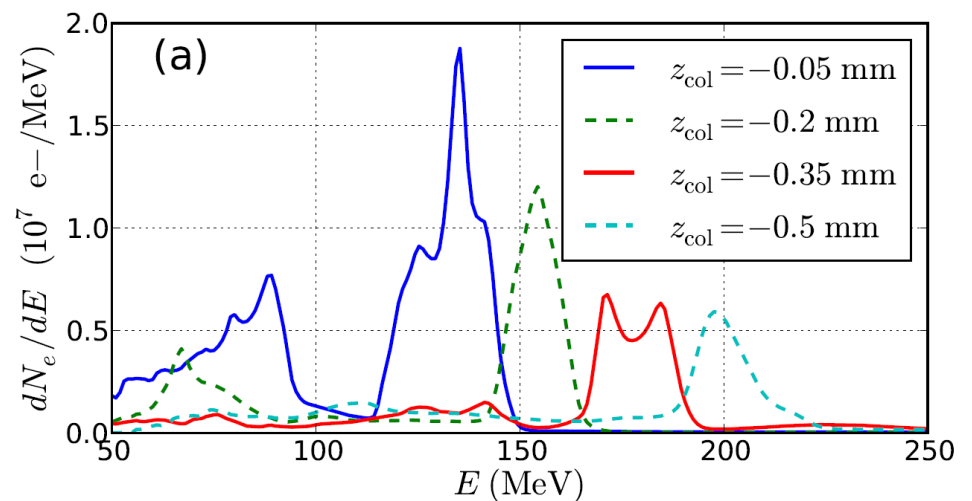
Experimental setup



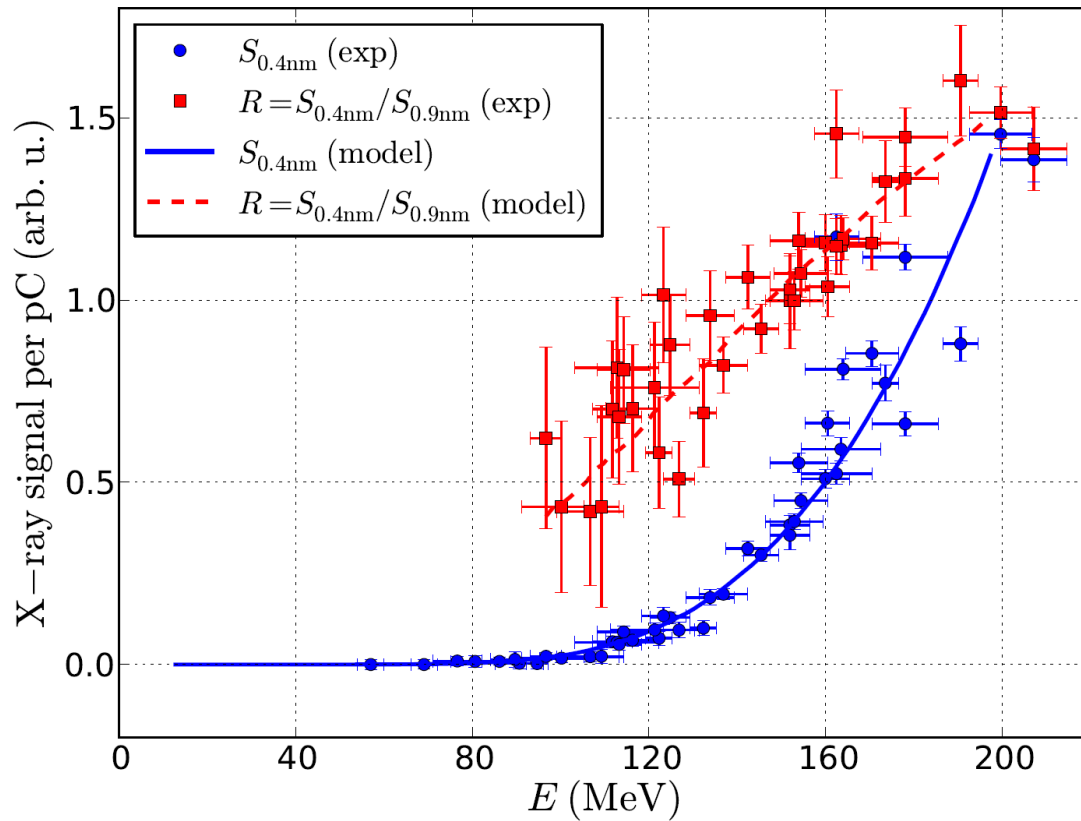
First measurement : electron energy



- Charge ~ 25 pC
- Energy $\in [80, 210]$ MeV

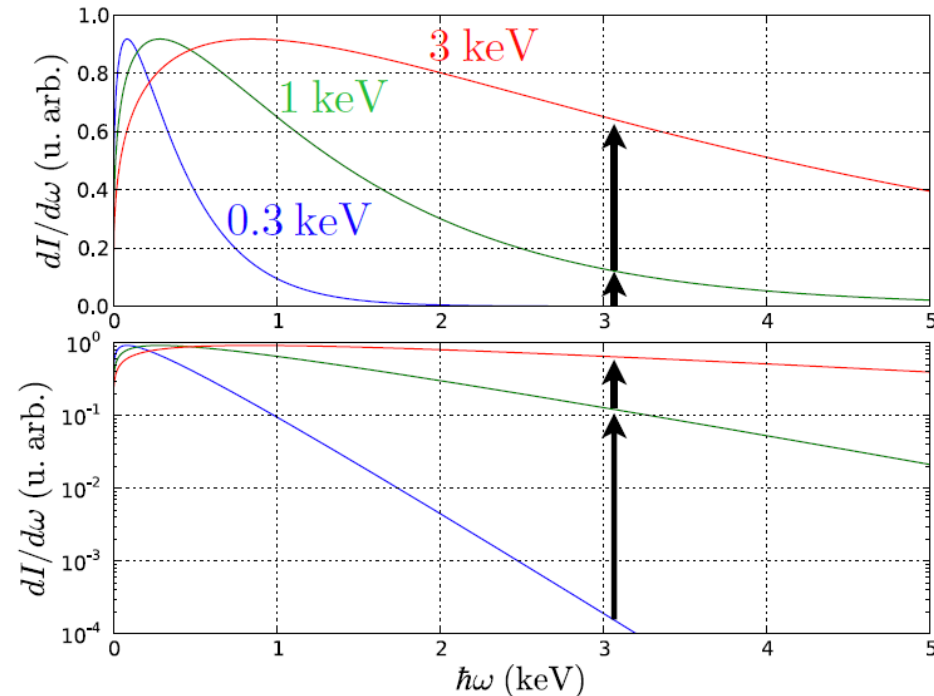


Second measurement : X-ray critical energy



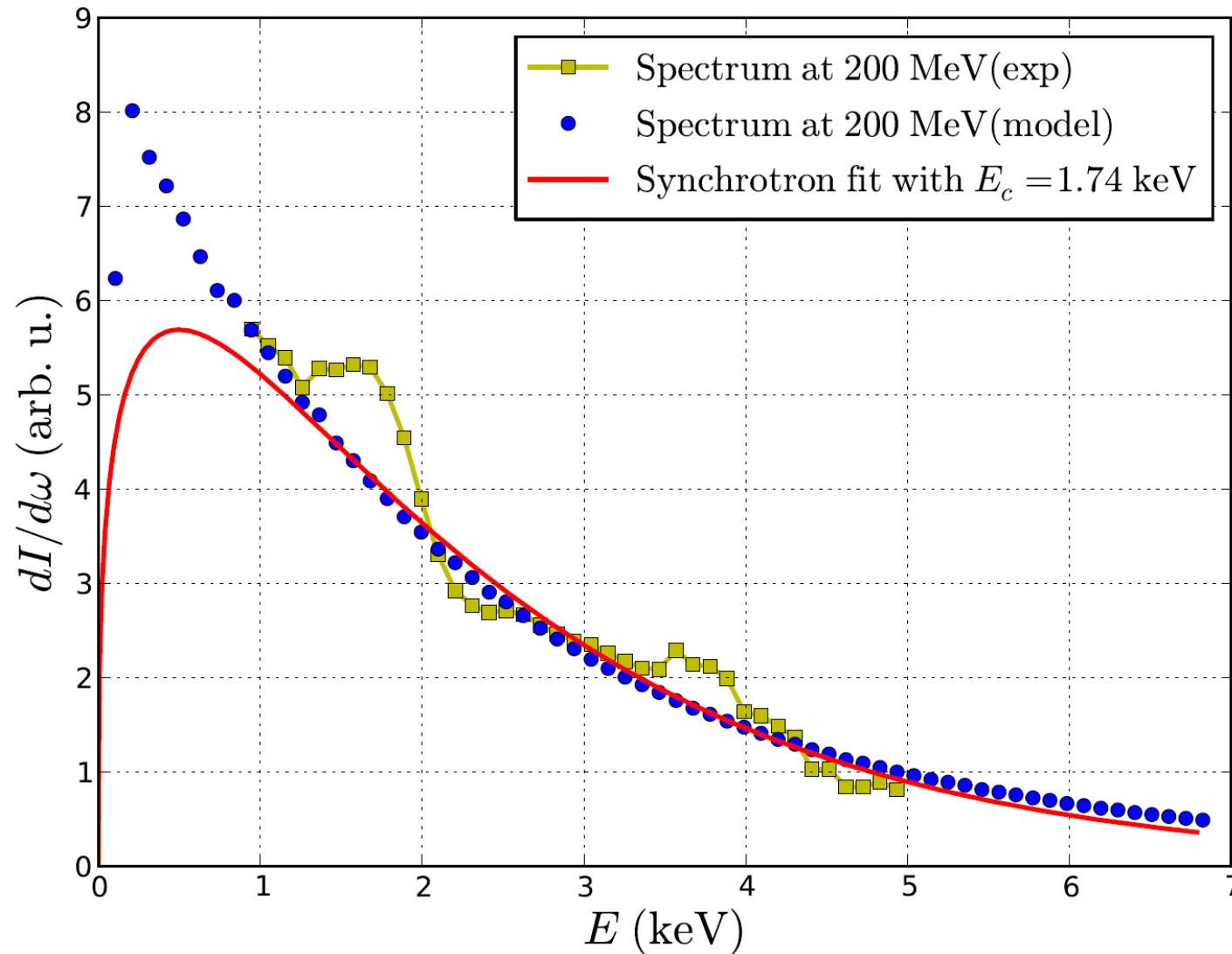
$$E_C \propto \epsilon_N^{1/2} \gamma^{7/4} \alpha^{3/4}$$

Critical energy increases with γ .



Good agreement between model and experiment
↓
validation of the method

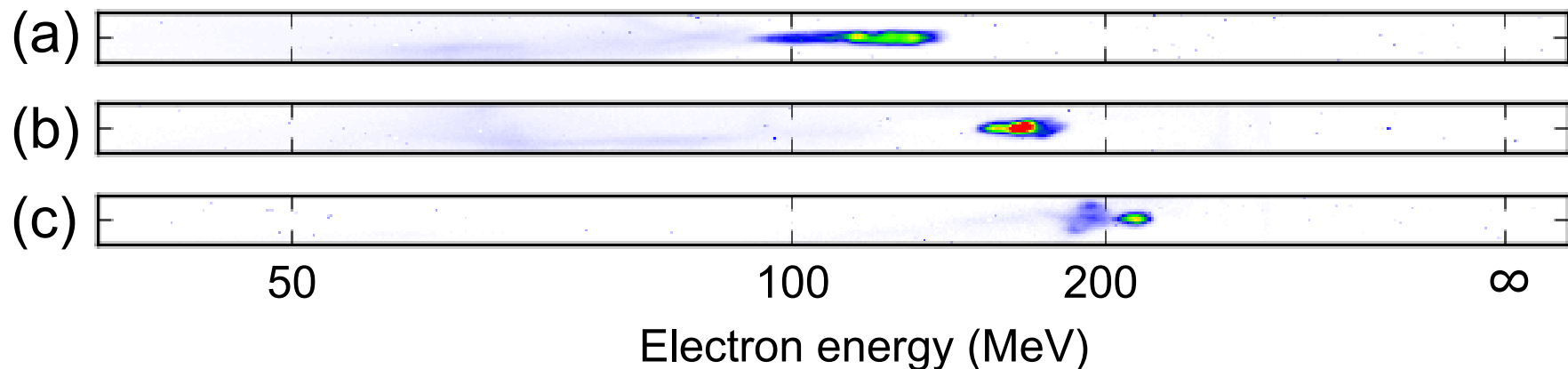
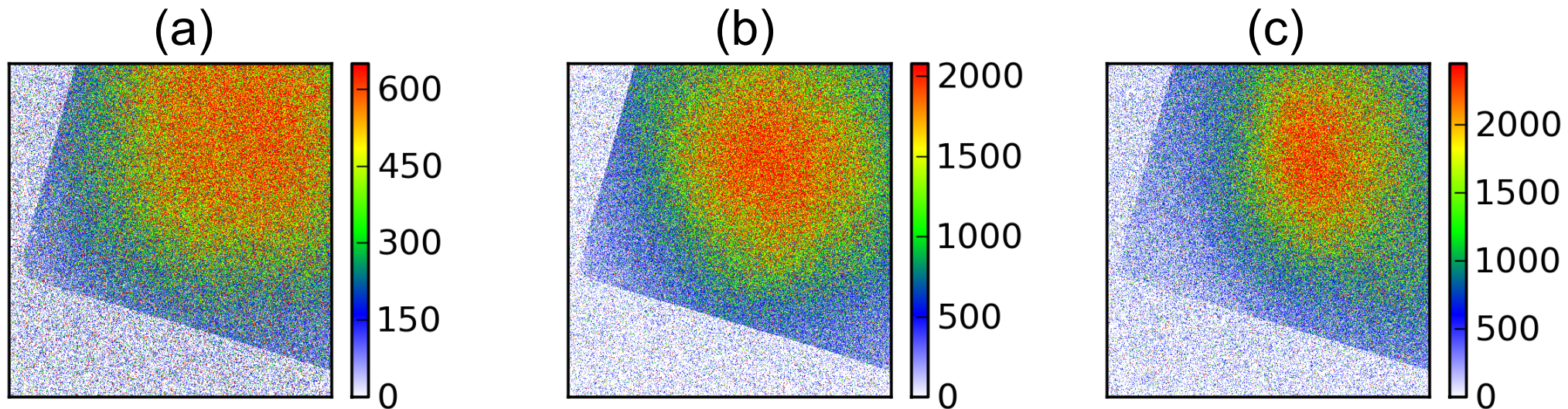
Second measurement : X-ray critical energy



$$E_c = 1.7 \pm 0.5 \text{ KeV} \Rightarrow \varepsilon_N = (0.53 \pm 0.36) \alpha^{-3/2} \text{ n.mm.mrad}$$

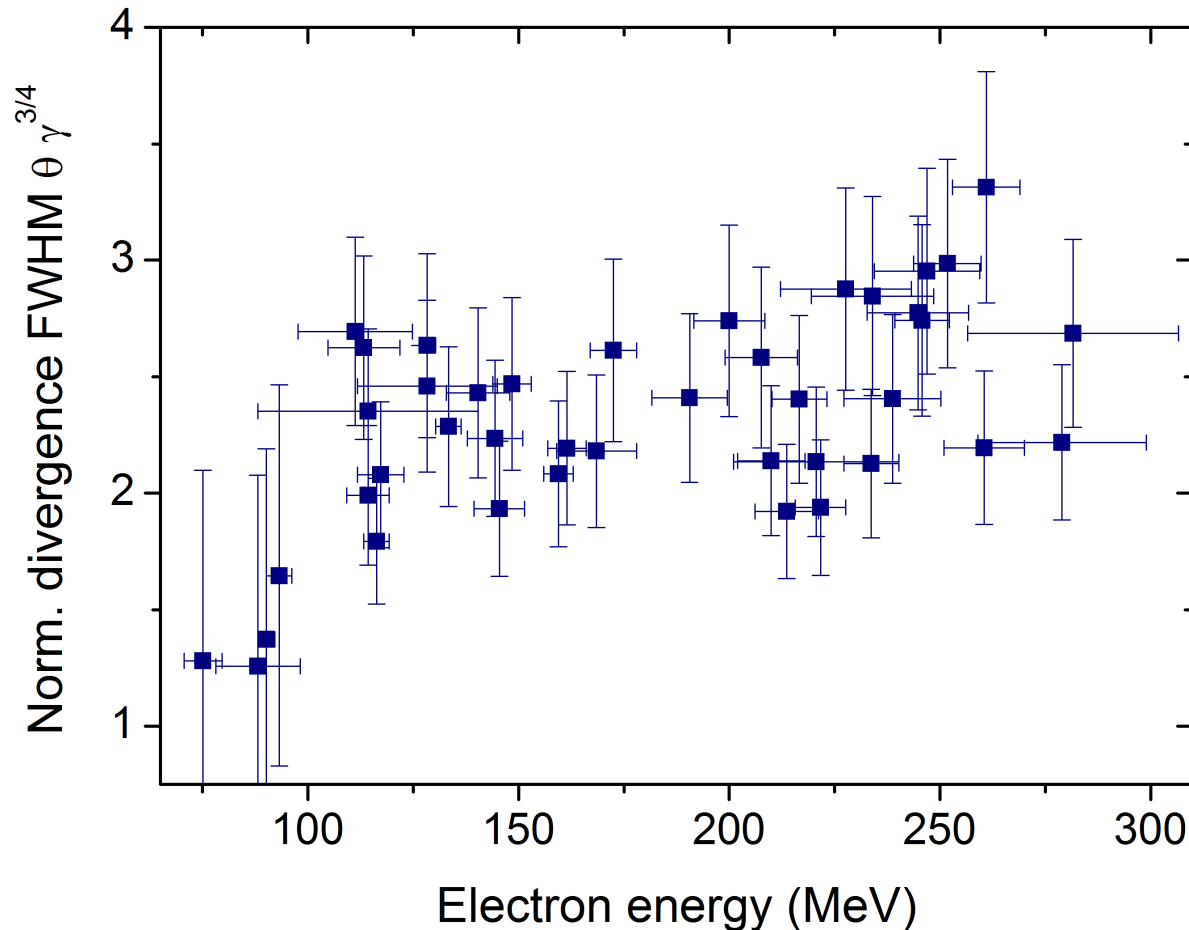
Third measurement : X-ray divergence

$\theta \propto \epsilon_N^{1/2} \alpha^{1/4} \gamma^{-3/4}$ \Rightarrow divergence decreases when γ increases



Third measurement : X-ray divergence

$$\theta \propto \epsilon_N^{1/2} \alpha^{1/4} \gamma^{-3/4}$$



The normalized divergence $\theta \gamma^{3/4}$ increases

Influence of the angular momentum

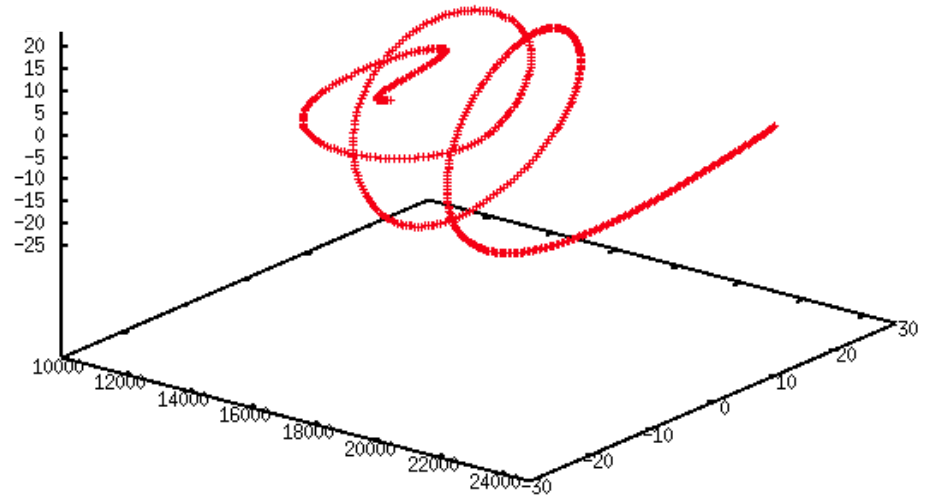
$$\hat{x}(t) = \sqrt{\hat{J}_x} \sin \left[\int_0^t \omega_\beta(t') dt' + \phi_x \right]$$

$$\hat{p}_x(t) = \sqrt{\hat{J}_x} \cos \left[\int_0^t \omega_\beta(t') dt' + \phi_x \right]$$

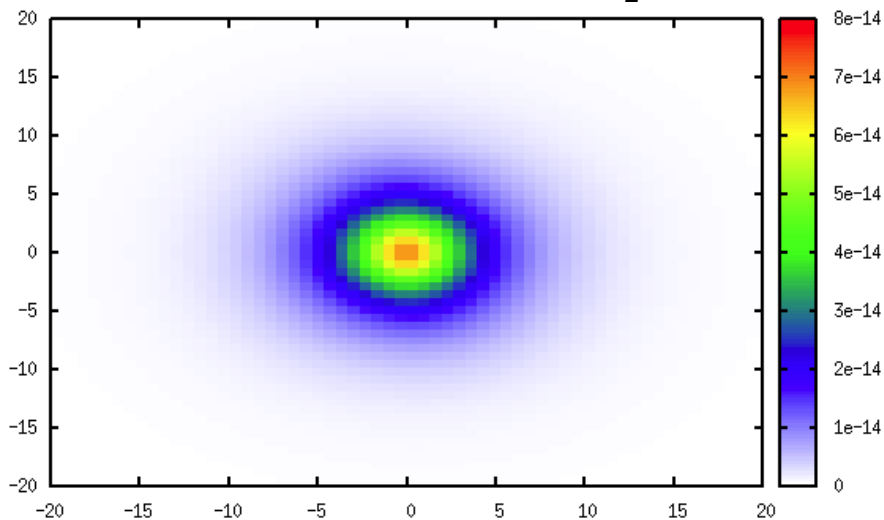
$$\hat{y}(t) = \sqrt{\hat{J}_y} \sin \left[\int_0^t \omega_\beta(t') dt' + \phi_y \right]$$

$$\hat{p}_y(t) = \sqrt{\hat{J}_y} \cos \left[\int_0^t \omega_\beta(t') dt' + \phi_y \right]$$

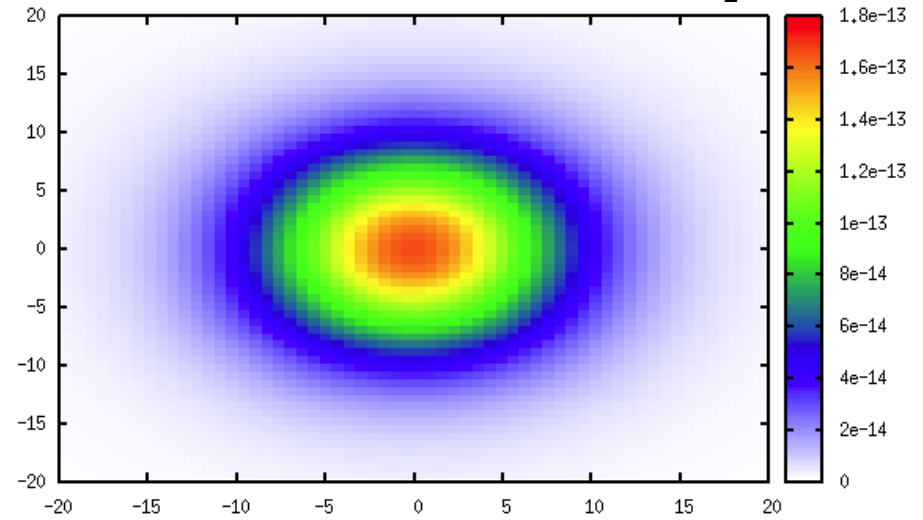
$\phi_x \neq \phi_y \rightarrow$ Non-zero angular momentum L_z



Simul. divergence for $L_z=0$



Simul. divergence for max. L_z

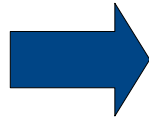


Angular momentum growth

$L_z(0)=0$ and non-perfectly symmetrical fields

$$F_x = -\alpha m \omega_p^2 x/2$$

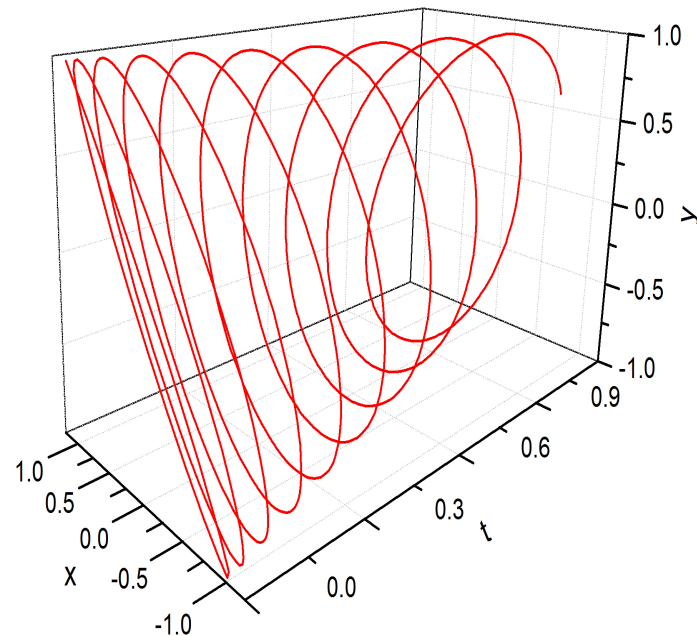
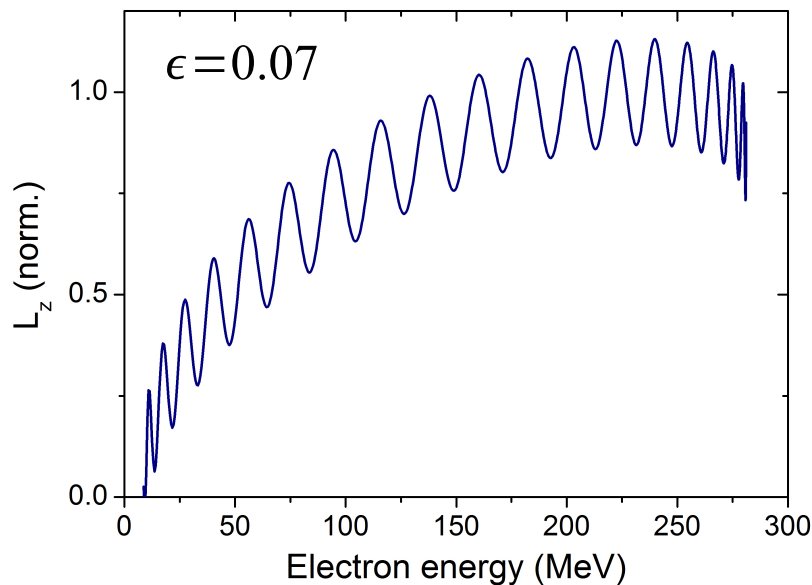
$$F_y = -\alpha(1+\epsilon) m \omega_p^2 x/2$$



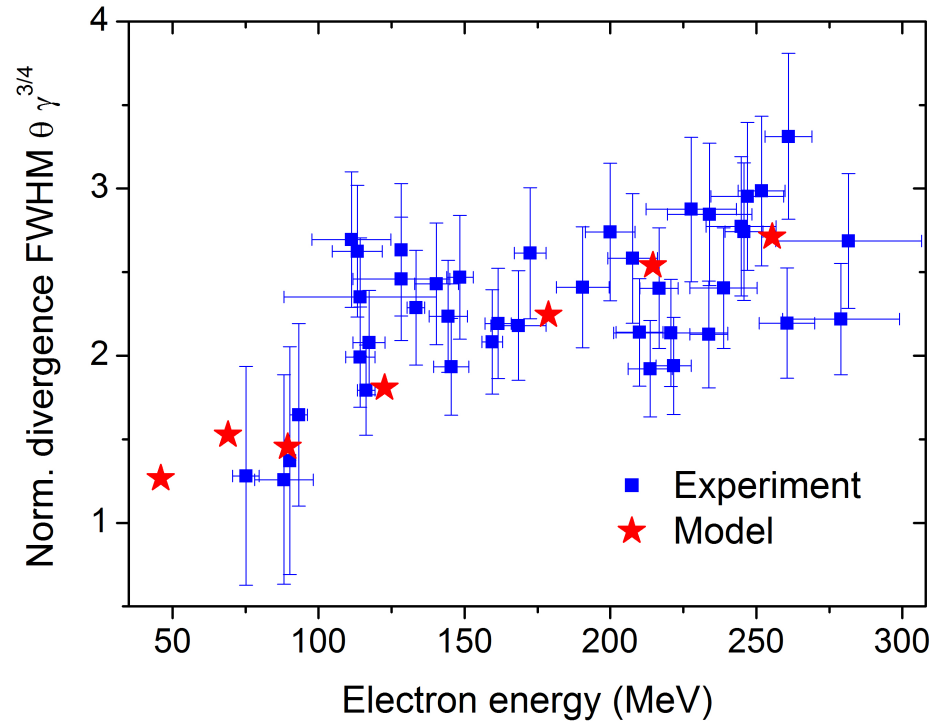
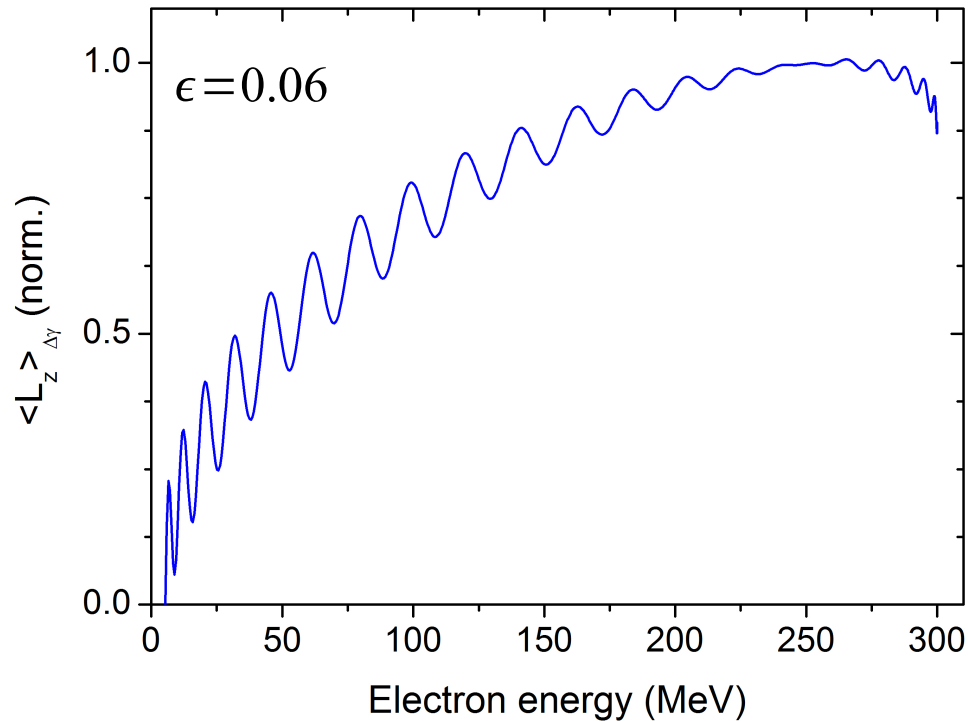
$$\phi_x = \int_0^t \omega_p \sqrt{\alpha/2 \gamma(t)} dt$$

$$\phi_y = \int_0^t \omega_p \sqrt{\alpha(1+\epsilon)/2 \gamma(t)} dt$$

$$L_z = \sqrt{J_x J_y} \left[(1-\epsilon) \sin(\phi_x - \phi_y) - 2\epsilon \sin(\phi_x + \phi_y) \right]$$



Angular momentum growth



A 4th measurement would be required to get L_z

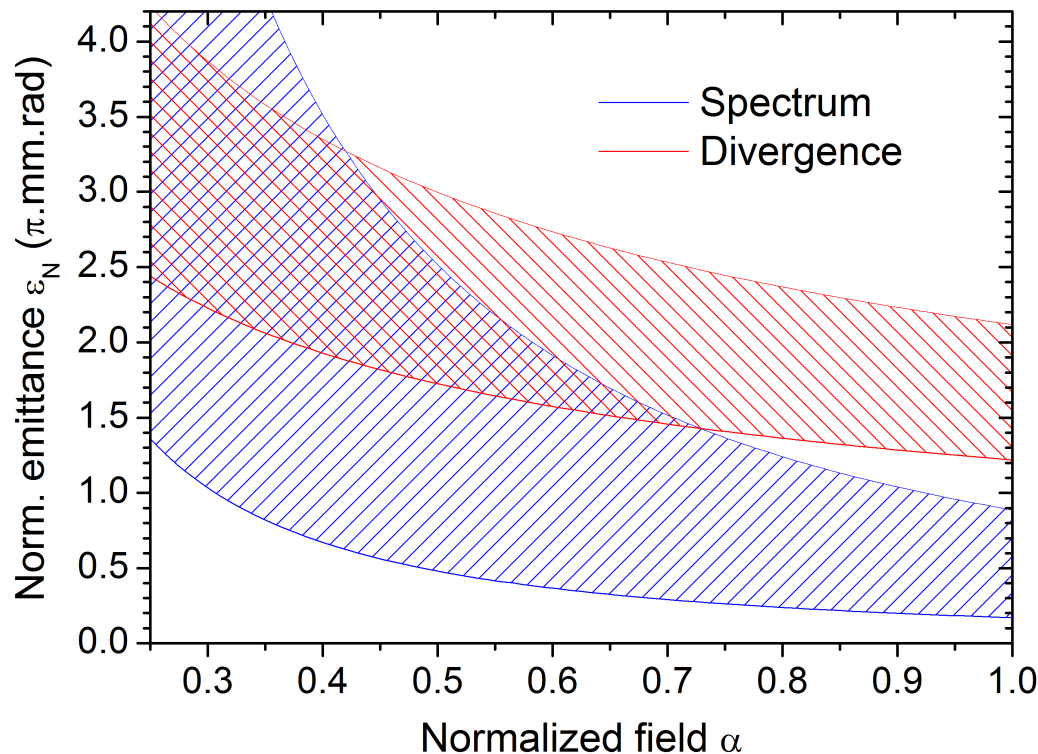
➡ assume that $\gamma^{3/4} \theta^{\text{FWHM}}$ is max. for $E = 260$ MeV

For $E = 260$ MeV, $\gamma^{3/4} \theta^{\text{FWHM}} = 1900 \pm 300$ mrad

➡ $\epsilon_N = (1.7 \pm 0.4) \alpha^{-1/2}$ n.mm.mrad

An estimate of the emittance

- Spectrum $\Rightarrow \varepsilon_N = (0.53 \pm 0.36) \alpha^{-3/2}$ π .mm.mrad
- Divergence $\Rightarrow \varepsilon_N = (1.7 \pm 0.4) \alpha^{-1/2}$ π .mm.mrad
- Theory $\Rightarrow 0.25 \lesssim \alpha \leq 1$



1.4π .mm.mrad $\lesssim \varepsilon_N \lesssim 4.2 \pi$.mm.mrad
Most probable : $\alpha = 0.32$ and $\varepsilon_N \approx 3 \pi$.mm.mrad

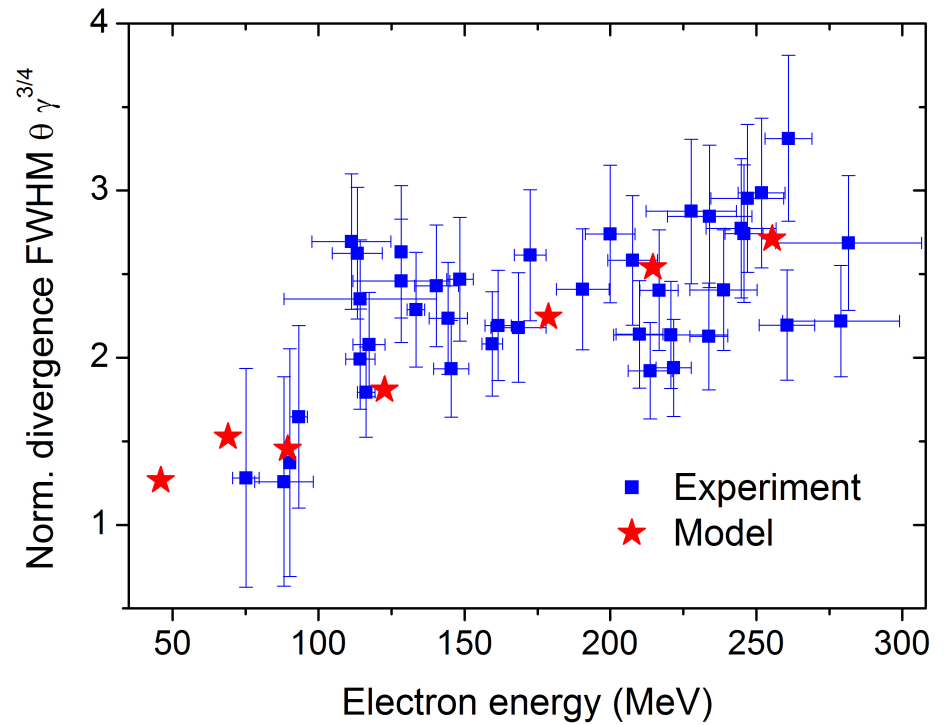
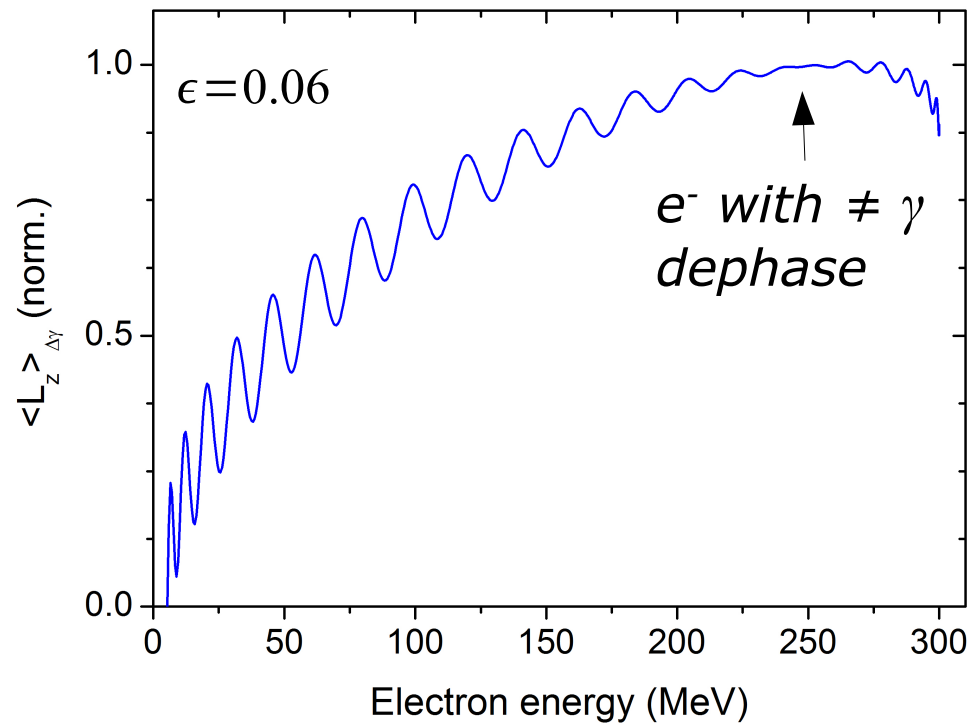
Summary

- Proof of principle experiment of betatron based emittance measurement.
- $\varepsilon_N < 4$ p.mm.mrad.
- Method works also for higher electron energy.
- Potentially single shot.

Outlook

- Reduce the error bars.
- Perform a 4th measurement to get L_z .
- Theoretical study of angular momentum growth.

Angular momentum growth

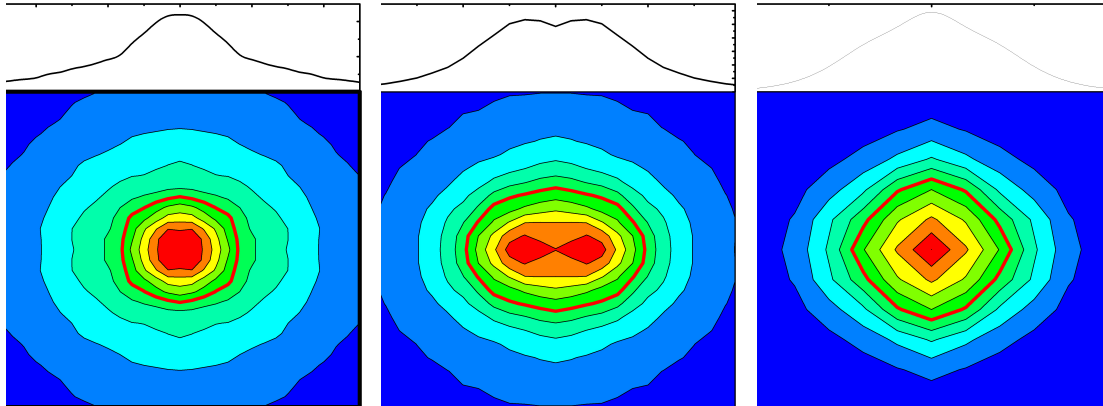


Model $\gamma^{3/4} \theta$

50 MeV

70 MeV

200 MeV



Experimental $\gamma^{3/4} \theta$

120 MeV

230 MeV

