Discovery potential of a light H+ at the LHC in the single top mode



October 10

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arXiv:1207.4071 - to appear in JHEP

Motivation

• Should we look for a charged Higgs in single top production?

$\sigma(pp \to t(+\bar{t}) + X \approx \frac{1}{3}\,\sigma(pp \to t\bar{t} + X)$

• Discover particles or exclude models



picture stolen from Y. Yamamoto

Multi-Higgs?

and Outline

- The models (2HDM no CP, with CP, with extensions?)
- Current bounds
- The analysis
- Conclusions

The softly broken Z_2 symmetric 2HDM potential

$$V(\Phi_{1}, \Phi_{2}) = m_{1}^{2} \Phi_{1}^{\dagger} \Phi_{1} + m_{2}^{2} \Phi_{2}^{\dagger} \Phi_{2} - (m_{12}^{2} \Phi_{1}^{\dagger} \Phi_{2} + \text{h.c}) + \frac{1}{2} \lambda_{1} (\Phi_{1}^{\dagger} \Phi_{1})^{2} + \frac{1}{2} \lambda_{2} (\Phi_{2}^{\dagger} \Phi_{2})^{2} + \lambda_{3} (\Phi_{1}^{\dagger} \Phi_{1}) (\Phi_{2}^{\dagger} \Phi_{2}) + \lambda_{4} (\Phi_{1}^{\dagger} \Phi_{2}) (\Phi_{2}^{\dagger} \Phi_{1}) + \frac{1}{2} \lambda_{5} [(\Phi_{1}^{\dagger} \Phi_{2})^{2} + \text{h.c.}]$$

 $\phi_1 \rightarrow \phi_1 \quad \phi_2 \rightarrow -\phi_2$

- m_{12}^2 and λ_5 real, vacuum configuration (CP-conserving) $\langle \phi_1 \rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v_1 \end{pmatrix}; \ \langle \phi_2 \rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v_2 \end{pmatrix}$

7 free parameters + M_W : m_h , m_H , m_A , $m_{H^{\pm}}$, $\tan\beta$, $\sin\alpha M^2 = \frac{m_{12}^2}{\sin\beta\cos\beta}$

- m_{12}^2 and λ_5 complex, vacuum configuration (explicit CP-violating)

$$\left\langle \phi_{1} \right\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v_{1} \end{pmatrix}; \quad \left\langle \phi_{2} \right\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v_{2} \end{pmatrix}$$

$$\operatorname{Im}[m_{12}^2] = v_1 v_2 \operatorname{Im}[\lambda_5]$$

I. Ginzburg, M. Krawczyk and P. Osland, hep-ph/ 0211371.

8 free parameters + M_W : $m_{H^{\pm}}$, $\tan\beta$, others...

Common features



Extending the Z₂ symmetry to the fermions – 4 independent Yukawa Lagrangians

			•••	• •
	Ι	II	III	IV
up	Φ_2	Φ_2	Φ_2	Φ_2
down	Φ_2	Φ_1	Φ_1	Φ_2
lepton	Φ_2	Φ_1	Φ_2	Φ_1
			•••	• •• •

III = I' = Y = Flipped IV = II' = X = Leptonic



an eta

- Production process is the same.
- Fermionic decay widths are the same.

Extensions that include 2HDM as a sub-model – example: Model X, in M. Aoki, S. Kanemura and O. Seto PRL102 (2009) 051805

2HDM Lagrangian (CP conserving to CP-violating potential)

• <u>couplings that involve gauge bosons</u>

$$sin(\beta - \alpha) \implies C = c_{\beta}R_{11} + s_{\beta}R_{12}$$

• couplings that involve fermions

 $\sin \alpha \quad \tan \beta$

	c_1c_2	$s_1 c_2$	s_2
R =	$-(c_1s_2s_3+s_1c_3)$	$c_1c_3 - s_1s_2s_3$	$c_{2}s_{3}$
	$\langle -c_1s_2c_3+s_1s_3$	$-(c_1s_3+s_1s_2c_3)$	$c_2 c_3$ /

		—		
	Type I	Type II	Lepton	Flipped
			Specific	
Up	$\frac{R_{12}}{s_{\beta}} - ic_{\beta} \frac{R_{13}}{s_{\beta}}$	$\frac{R_{12}}{s_{\beta}} - ic_{\beta} \frac{R_{13}}{s_{\beta}}$	$\frac{R_{12}}{s_{\beta}} - ic_{\beta} \frac{R_{13}}{s_{\beta}}$	$\frac{R_{12}}{s_{\beta}} - ic_{\beta}\frac{R_{13}}{s_{\beta}}$
Down	$\frac{R_{12}}{s_{\beta}} + ic_{\beta} \frac{R_{13}}{s_{\beta}}$	$rac{R_{11}}{c_{eta}} - i s_{eta} rac{R_{13}}{c_{eta}}$	$\frac{R_{12}}{s_{\beta}} + ic_{\beta} \frac{R_{13}}{s_{\beta}}$	$\frac{R_{11}}{c_{\beta}} - is_{\beta} \frac{R_{13}}{c_{\beta}}$
Leptons	$\frac{R_{12}}{s_{\beta}} + ic_{\beta} \frac{R_{13}}{s_{\beta}}$	$rac{R_{11}}{c_{eta}} - i s_{eta} rac{R_{13}}{c_{eta}}$	$rac{R_{11}}{c_{eta}} - is_{eta} rac{R_{13}}{c_{eta}}$	$\frac{R_{12}}{s_{\beta}} + ic_{\beta} \frac{R_{13}}{s_{\beta}}$

Model	$g_{ar{u}dH^+}$	$g_{l\bar{\nu}H^+}$
Ι	$\frac{ig}{\sqrt{2}M_W}V_{ud}\left[-m_d/\tan\beta P_R + m_u/\tan\beta P_L\right]$	$\frac{ig}{\sqrt{2}M_W}\left[-m_l/\tan\beta P_R\right]$
II	$\frac{ig}{\sqrt{2}M_W}V_{ud}\left[m_d\tan\beta P_R + m_u/\tan\beta P_L\right]$	$\frac{i g}{\sqrt{2} M_W} \left[m_l \tan \beta P_R \right]$
Υ	$\frac{ig}{\sqrt{2}M_W}V_{ud}\left[m_d\tan\beta P_R+m_u/\tan\beta P_L\right]$	$\frac{ig}{\sqrt{2}M_W}\left[-m_l/\tan\beta P_R\right]$
X	$\frac{ig}{\sqrt{2}M_W} V_{ud} \left[-m_d / \tan\beta P_R + m_u / \tan\beta P_L \right]$	$\frac{i\tilde{g}}{\sqrt{2}M_W}\left[m_l\tan\beta P_R\right]$



top and charged Higgs Branching Ratios in models I and X

light particles are involved.

Constraints - How light is light?

• <u>DIRECT BOUNDS</u> LEP bound for a charged Higgs $e^+e^- \rightarrow H^+H^-$

$$BR(H^+ \to \tau^+ \nu) + BR(H^+ \to c\bar{s}) = 1 \qquad m_{H^{\pm}} > 79.3 \, GeV$$
$$BR(H^+ \to \tau^+ \nu) \approx 1 \qquad m_{H^{\pm}} > 89.6 \, GeV \qquad \text{(Model X)}$$

• **INDIRECT BOUNDS** B factories





Can we probe "oval" region? (Models I and X)



Only Higgs pair production can increase with tanß.

For resonant production



So far

A class of models (2HDM and 2HDM-like) can be studied with single top production.

Constraints imply (at least) charged Higgs above 80 GeV (LEP), and tanβ above 1 (B-physics and R_b - charged Higgs mass below top mass).

Yukawa couplings fall with tan²β (type I,X). Other processes (double charged Higgs production) do not have that limitation but a) cross sections are smaller b) strong dependence on the remaining model parameters (resonant production).

After ttbar, searches based on single top are then the best way to maximize the exclusion in the (mass, tanβ) plane.

> Could be important in "the MSSM" if $|\Delta b|$ is large! Oscar talk this morning.

Single top@14TeV (7 and 8 TeV bad)

<u>Signal</u> - single top (after cuts signal is t-channel contribution)



Charged Higgs decays to tau nu and only leptonic final states are considered.

<u>Background (irreducible)</u> - single top

Background (reducible) - ttbar, W+nj, Wc+nj, Wbb+nj (n<4)

"leptonic transverse mass"

$$M_T^{l\nu} = \sqrt{2p_{Tl}\not{p}_T - 2(p_{xl}\not{p}_x + p_{yl}\not{p}_y)},$$

 $30 \,{
m GeV} < M_T^{l\nu} < 60 \,{
m GeV}$ Masses: 90 to 130 GeV $30 \,{
m GeV} < M_T^{l\nu} < 60 \,{
m GeV}$ or $M_T^{l\nu} > 85 \,{
m GeV}$ above 130 GeV



Background

		Events (1 fb^{-1})	Events (1 fb^{-1})	
POWHEG	Process	before cuts	after cuts	Efficiency $(\%)$
(AcerMC)	Single top $(t$ -channel)	246600	18.5	0.0075
4	Single top $(s$ -channel)	10650	0	0
	Single top (tW)	66000	0.7	0.0010
ΡΥΤΗΤΑ	$t\bar{t}$ (semileptonic)	371133	0	0
	$t\bar{t}$ (leptonic)	88940	1.8	0.0020
(top decays)	$t\bar{t}$ (hadronic)	387169	0	0
	W+0j	4.3×10^{7}	0	0
	W+1j	8.8×10^{6}	0.6	$7.1 imes 10^{-6}$
	W+2j	2.8×10^{6}	3.9	0.00014
AlpGen <	W+3j	1.2×10^{6}	0	0
	Wc+0j	7.4×10^5	0	0
(MCFM)	Wc+1j	3.2×10^5	3.2	0.0010
	Wc+2j	2.0×10^{5}	1.0	0.0011
	Wc+3j	8.9×10^4	0	0
	Wbb+0j	6638	0	0
(nadronisation)	Wbb+1j	6582	0.5	0.007
+	Wbb+2j	3746	0	0
DELPHES	Wbb+3j	2896	0	0

(detector effects)

charged Higgs mass 120 GeV







B($t \rightarrow bH^{+}$



Type I 2HDM-like models



<u>Theoretical bounds</u> to the rescue (in the exact Z_2 model, m_{12} =0). Taking the lightest neutral Higgs 125 GeV

$$0.18\lesssim ext{tan}\,eta\lesssim 5.59$$

B. Gorczyca, M. Krawczyk, arXiv: 1112.5086

Similar trend for type I models. Branching ratios are slightly smaller than the type X ones.





Insert result



Then (14 TeV prediction) and now (7 TeV result)



Conclusions

Single top channel looks good.

The factor of ~3 in production is turned into a factor between 2 and 3 after the analysis.

ATLAS and CMS can improve the factor.

Light charged Higgs in Type-II-like models will be excluded (charged Higgs discovered) by the end of the 14 TeV run.

In Type-I (X) models, tanβ will be probed in the range 1 to 15 (my guess), depending on the charged Higgs mass, by the end of the 13/14 TeV run. the end

a) 1 electron P_T > 30 GeV and $|\eta| < 2.5$ or 1 muon P_T > 20 GeV and $|\eta| < 2.5$

- b) Veto events with two or more leptons with $P_T > 10 \text{ GeV}$
- c) Exclude events with leptons with $P_T > 55 \text{ GeV}$
- d) Veto events with missing E_T below 50 GeV
- e) Exactly one b-tagged jet with 20 GeV < P_T < 75 GeV

f) top quark invariant mass above 280 GeV $a = \left(\frac{p_{zl}}{p_l}\right)^2 - 1$,

$$p\!\!\!/_z = \frac{-b\pm\sqrt{b^2-4ac}}{2a}$$

g) "leptonic transverse mass"

$$M_T^{l\nu} = \sqrt{2p_{Tl}\not\!p_T - 2(p_{xl}\not\!p_x + p_{yl}\not\!p_y)},$$

 $30 \,\mathrm{GeV} < M_T^{l
u} < 60 \,\mathrm{GeV}$ charged Higgs masses 90 to 130 GeV $30 \,\mathrm{GeV} < M_T^{l
u} < 60 \,\mathrm{GeV}$ or $M_T^{l
u} > 85 \,\mathrm{GeV}$ above 130 GeV

 $b = 2\left(\frac{p_{xl}\not{p}_x + p_{yl}\not{p}_y}{p_l} + \frac{m_{H^+}^2}{2p_l}\right)\frac{p_{zl}}{p_l},$

h) One and only one non-b jet, P_T above 30 GeV

i) Accept events where jets have $|\eta| > 2.5$

Parametrisation (8) W. Khater and P. Osland, Nucl. Phys. B 661, 209 (2003).

2 charged, H^{\pm} , and 3 neutral, h_1 , h_2 and h_3 3 masses

$$\stackrel{\bullet}{\longrightarrow} \begin{pmatrix} h_1 \\ h_2 \\ h_3 \end{pmatrix} = R \begin{pmatrix} \eta_1 \\ \eta_2 \\ \eta_3 \end{pmatrix} \quad R \mathcal{M}^2 R^T = \operatorname{diag} \left(m_1^2, m_2^2, m_3^2 \right)$$

$$R = \begin{pmatrix} c_1 c_2 & s_1 c_2 & s_2 \\ -(c_1 s_2 s_3 + s_1 c_3) & c_1 c_3 - s_1 s_2 s_3 & c_2 s_3 \\ -c_1 s_2 c_3 + s_1 s_3 & -(c_1 s_3 + s_1 s_2 c_3) & c_2 c_3 \end{pmatrix}$$
3 angles

 \implies Re $[m_{12}^2]$ soft breaking term

 \rightarrow tan β ratio of vacuum expectation values

$$\implies m_3^2 = \frac{m_1^2 R_{13} (R_{12} \tan \beta - R_{11}) + m_2^2 R_{23} (R_{22} \tan \beta - R_{21})}{R_{33} (R_{31} - R_{32} \tan \beta)}$$

Experimental

• **INDIRECT BOUNDS** B factories

Models I and X



 $\tan\beta > 1$

 $m_{H^{\pm}} = 100 \, GeV$

Combined "model independent" results

