



Andreas Crivellin

**Albert Einstein Center for Fundamental Physics
ITP Bern**

**Tauonic-B decays
and 2HDMs**

Outline:

- n Tauonic B decays

- n $B \rightarrow TV$

- n $B \rightarrow DTV$

- n $B \rightarrow D^*TV$

- n Two Higgs Doublet Models

- n Type II

- n Type III

- n 2HDM III and the MSSM

- n Conclusions

Tauonic B decays

- n Tree-level decays in the SM via W-boson
- n Sensitive to a charged Higgs due to the heavy tau lepton in the final state.

Observable	SM	Experiment	Significance
$\text{Br}[B \rightarrow \tau \nu]$	$(0.719^{+0.115}_{-0.076}) \times 10^{-4}$	$(1.15 \pm 0.23) \times 10^{-4}$	1.6σ
$\text{Br}[B \rightarrow D \tau \nu] / \text{Br}[B \rightarrow D \nu]$	0.297 ± 0.017	0.440 ± 0.072	2.0σ
$\text{Br}[B \rightarrow D^* \tau \nu] / \text{Br}[B \rightarrow D^* \nu]$	0.252 ± 0.003	0.332 ± 0.030	2.7σ

➡ All three observables are above the SM prediction

B → TV

$$Br[B \rightarrow \tau \nu] = \frac{G_F^2 |V_{ub}|^2}{8\pi} m_\tau^2 f_B^2 m_B \left(1 - \frac{m_\tau^2}{m_B^2} \right) \tau_B \left| 1 + \frac{m_B^2}{m_b m_t} \frac{C_R^{ub} - C_L^{ub}}{C_{SM}^{ub}} \right|^2$$

n V_{ub} can be determined from

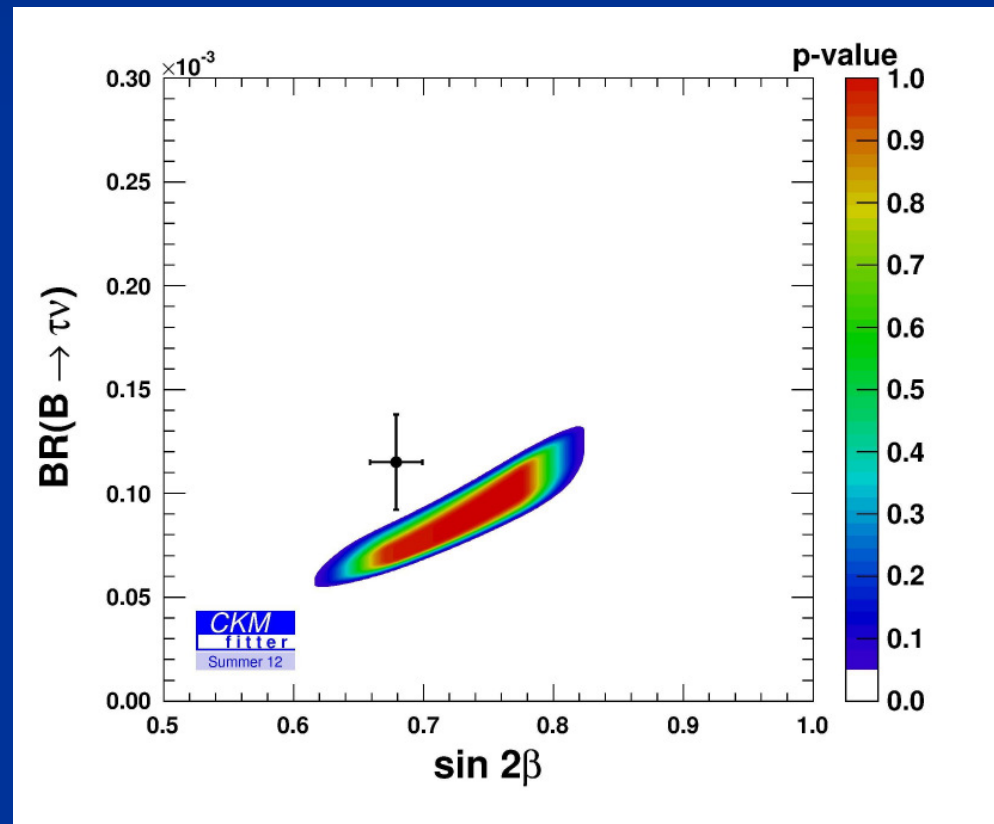
n $B \rightarrow \pi l \nu$

n inclusive decay

n Global fit to the CKM matrix

Different determinations do not agree

➔ V_{ub} problem

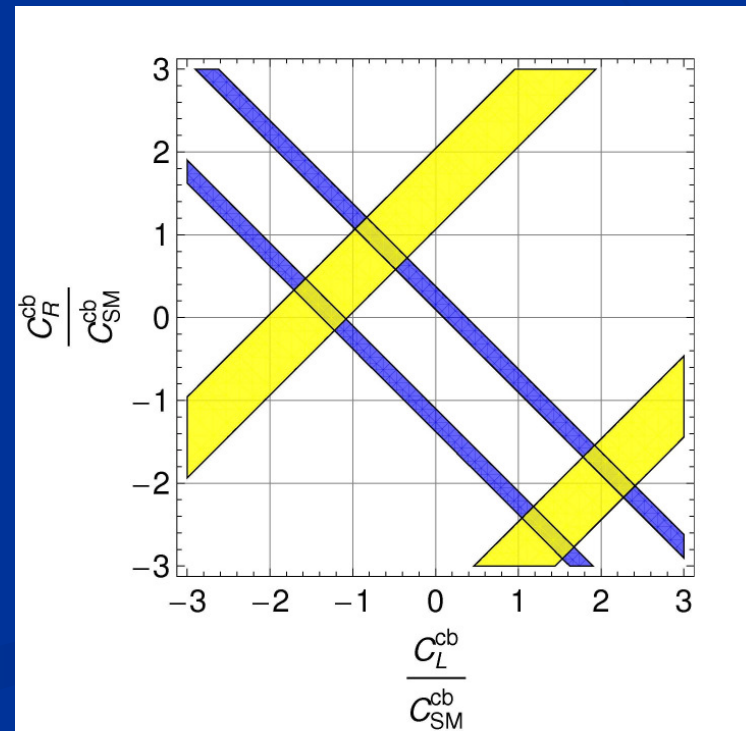


$B \rightarrow D^{(*)} \tau \nu$

$$R(D) = \frac{\text{Br}[B \rightarrow D \tau \nu]}{\text{Br}[B \rightarrow D \nu]} = R_{SM}(D) \left(1 + 1.5 \text{Re} \left[\frac{C_R^{cb} + C_L^{cb}}{C_{SM}^{cb}} \right] + 1.0 \left| \frac{C_R^{cb} + C_L^{cb}}{C_{SM}^{cb}} \right|^2 \right)$$

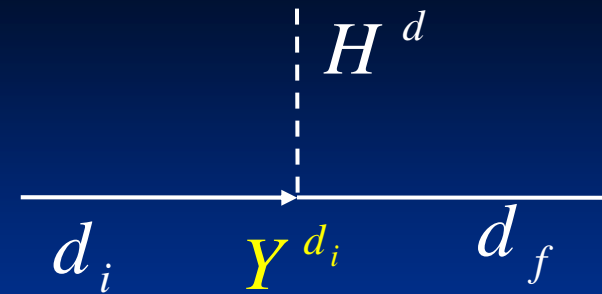
$$R(D^*) = \frac{\text{Br}[B \rightarrow D^* \tau \nu]}{\text{Br}[B \rightarrow D^* \nu]} = R_{SM}(D^*) \left(1 + 0.12 \text{Re} \left[\frac{C_R^{cb} - C_L^{cb}}{C_{SM}^{cb}} \right] + 0.05 \left| \frac{C_R^{cb} - C_L^{cb}}{C_{SM}^{cb}} \right|^2 \right)$$

- § Form factors uncertainties drop out to a large extent in the ratios $R(D)$ and $R(D^*)$.
- § $R(D^*)$ less sensitive to NP
- § C_R cannot explain $R(D)$ and $R(D^*)$ simultaneously but C_L can.



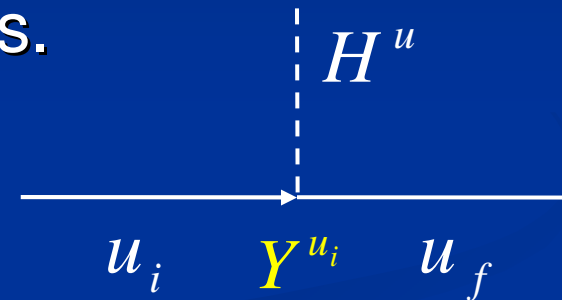
2HDM of type II

- One Higgs doublet couples only to down-quarks the other Higgs doublet only to up-quarks.



- 2 additional free parameters:

$\tan(\beta)=v_u/v_d$ and the heavy Higgs mass M_H



- Neutral Higgs-quark couplings are flavour-conserving.
- Charged Higgs contribution to $b \rightarrow s\gamma$ requires $m_H > 380 \text{ GeV}$ T. Hermann, M. Misiak and M. Steinhauser 1208.2788

Tauonic B decays in the 2HDM II

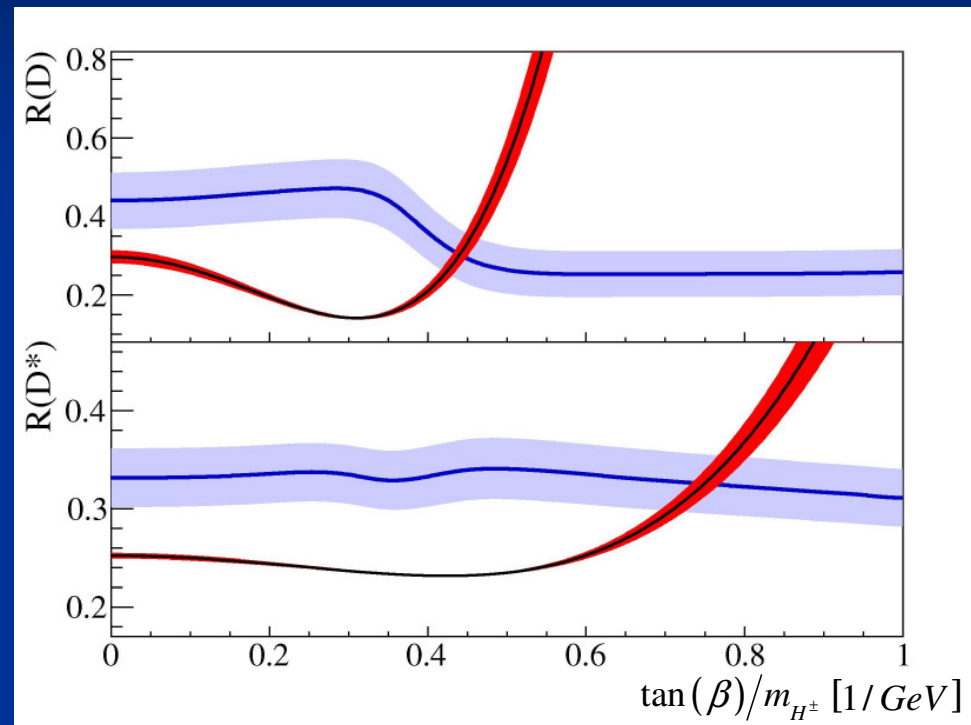
$$C_R^{qb} = \frac{-1}{m_{H^\pm}^2} V_{qb} \frac{m_b m_\tau}{v^2} \tan^2(\beta)$$

$$C_L^{qb} \approx 0$$

- n Contribution to $B \rightarrow \tau \nu$ necessarily destructive.
- n $\tan(\beta)/m_{H^\pm}$ needed for $R(D^*)$ too large.
- n Cannot explain $B \rightarrow D^{(*)} \tau \nu$ and $B \rightarrow D \tau \nu$ simultaneously.

BaBar collaboration 1205.5442

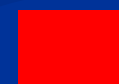
➡ Disfavored by current data



arXiv:1205.5442



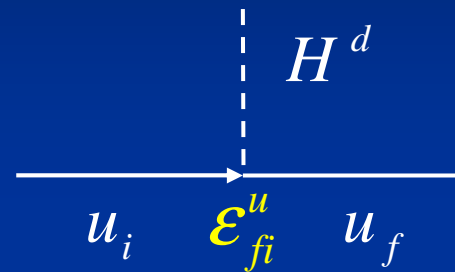
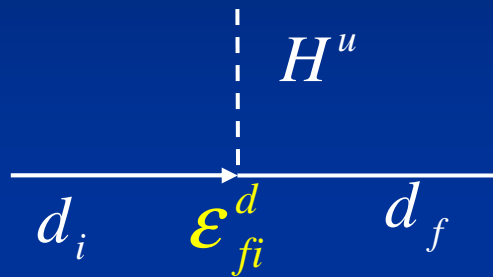
measurement



theory prediction

2HDM of type III

- Both Higgs doublets couple simultaneously to up and down quarks \longrightarrow Flavour-changing neutral Higgs couplings



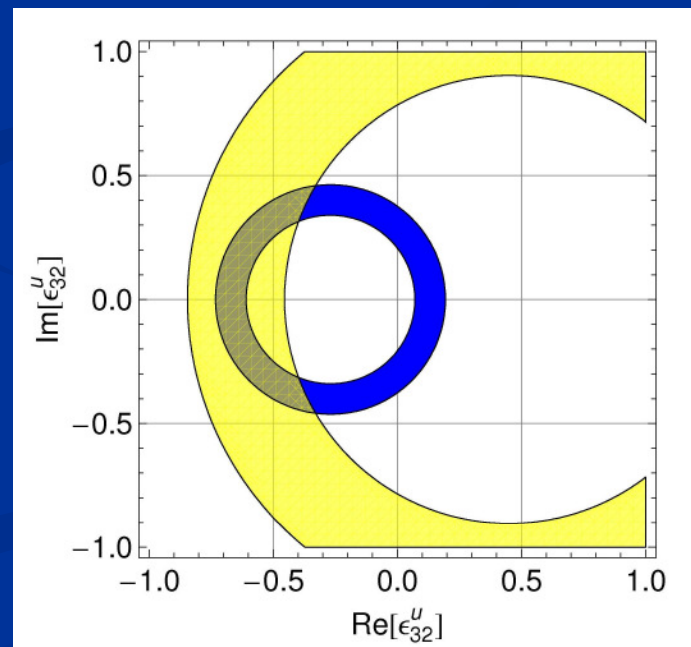
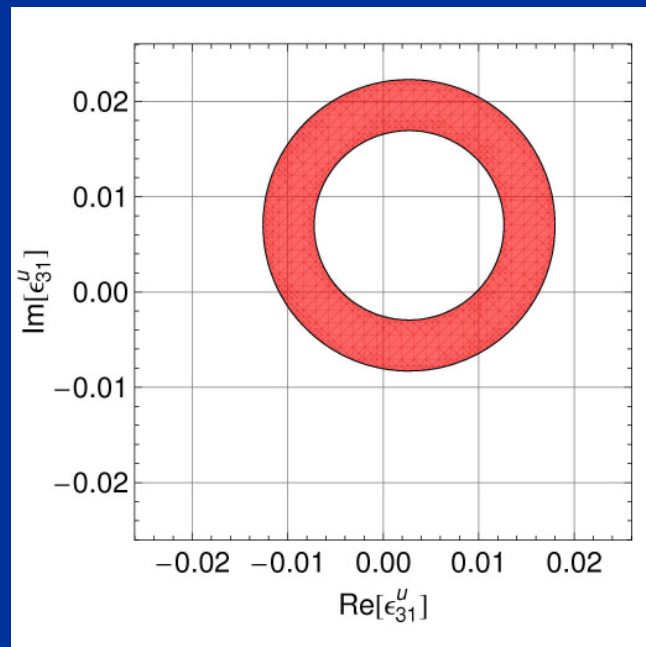
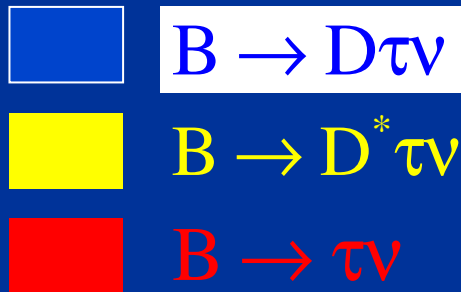
- All flavour-changing elements \mathcal{E}_{fi}^d and $\mathcal{E}_{12,21}^u$ are constrained from FCNCs processes.
- Also $\mathcal{E}_{13,23}^u$ constrained from charged Higgs diagrams, but $\mathcal{E}_{31,32}^u$ is unconstrained.
- Contribution to tauonic B decays $C_L^{qb} \approx \frac{\mathcal{E}_{3q}^{u*}}{m_{H^\pm}^2} \frac{m_\tau}{v} \tan(\beta)$
 $\mathcal{E}_{31,32}^u$ unconstrained from FCNCs
- 2HDM III with MFV cannot explain $B \rightarrow \tau \nu$, $B \rightarrow D^{(*)} \tau \nu$ and $B \rightarrow D \tau \nu$

2HDM of type III with flavour-violation in the up-sector

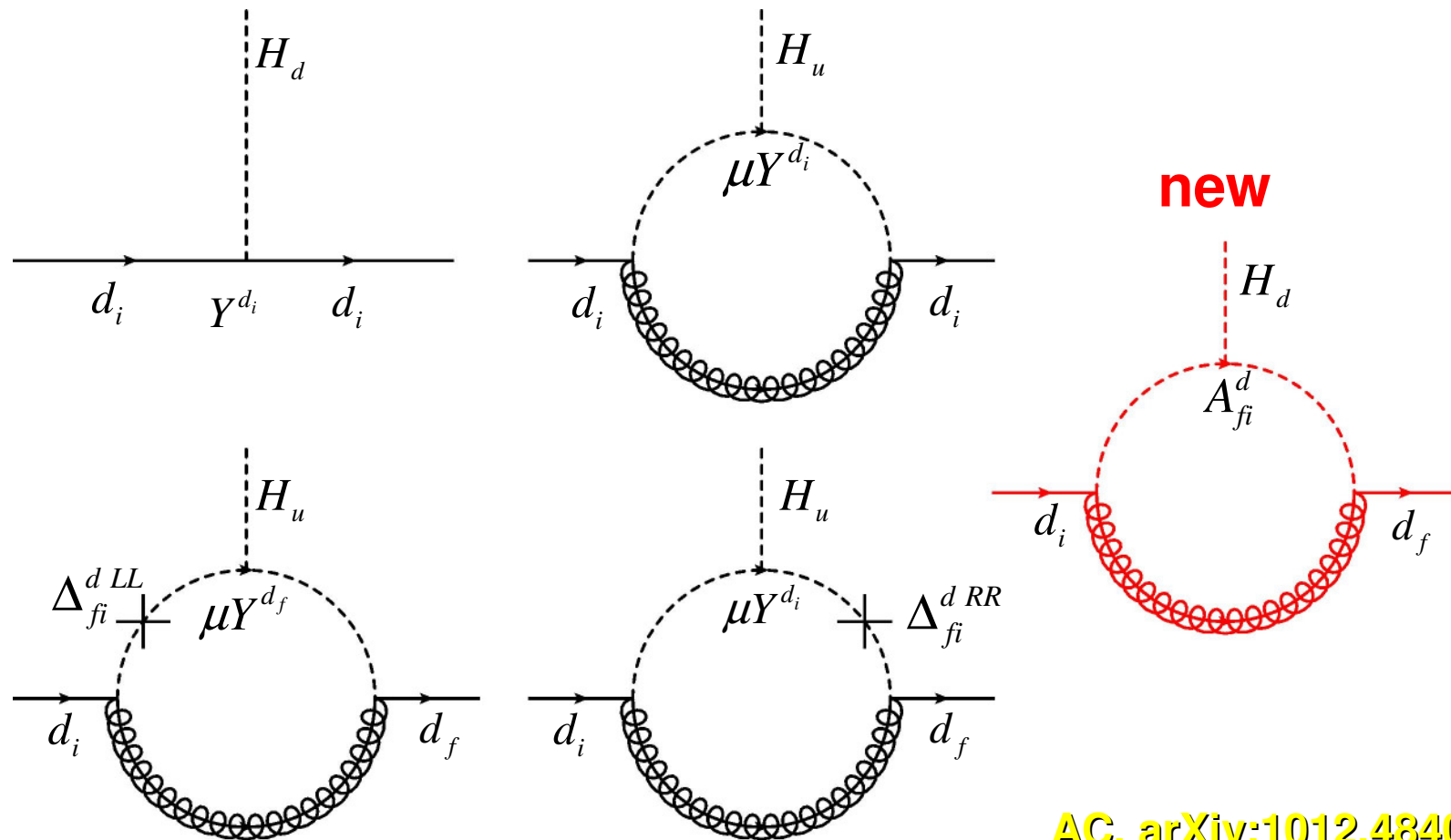
A.C., C. Greub, A. Kokulu 1208.2788

- Constructive contribution to $B \rightarrow \tau \nu$ using \mathcal{E}_{31}^u is possible.
- $B \rightarrow D^{(*)} \tau \nu$ and $B \rightarrow D \tau \nu$ can be explained simultaneously using \mathcal{E}_{32}^u . \longrightarrow Check model via $H^0, A^0 \rightarrow \bar{t} c$

Allowed regions from:



Effective Higgs-quark-quark vertices in the MSSM



MSSM and the decoupling limit

- n Effective Higgs vertices generate important threshold corrections to the relation between Yukawa couplings and quark masses.


➔ **Resummation of all ($\tan(\beta)$) enhanced contribution necessary.**

L. Hall, R. Rattazzi, U. Sarid hep-ph/9306309,
A.C., L. Hofer, J. Rosiek, arXiv:1103.4272

- n MSSM corrections are too small to generate $\epsilon_{32,31}^u$ needed for $B \rightarrow \tau \nu$ and $B \rightarrow D^{(*)} \tau \nu$.
- n NLO calculation in preparation including analytic results and $\tan(\beta)$ resummation in the generic MSSM.

Δ_b at order α_s^2 A.C., Christoph Greub arXiv:1012.xxxx

Conclusions

- n First hints for violation of lepton universality tauonic B-decays.
- n 2HDM II disfavored by current data.
- n 2HDM of type III with flavour-violation in the up-sector can explain the $B \rightarrow \tau \nu$, $B \rightarrow D^* \tau \nu$ and $B \rightarrow D \tau \nu$.
- n 2HDM III is the decoupling limit of the MSSM  NLO matching is important.

Effects of a right-handed W-coupling on V_{ub}

