Progress on dS and FRW holography

Gonzalo Torroba

SLAC, Stanford University

Based on works with X. Dong, B. Horn, S. Matsuura, E. Silverstein arXiv:1005.5403, 1108.5732, 1203.1680, and work in progress

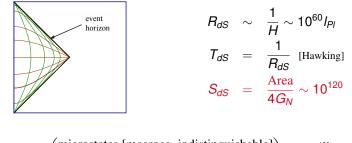
String Phenomenology Institute, CERN, July 2012

▲□▶▲□▶▲□▶▲□▶ □ のQ@

The measurement of the late time acceleration of the universe is one of the fundamental discoveries of the 20th century.

Deep implications of $\Lambda_{cc}>0$ and a de Sitter phase of exponential expansion.

★ Causal patch with event horizon:



 $QM: e^{S_{dS}} = \begin{pmatrix} \text{microstates [macrosc. indistinguishable]} \\ \text{that build our causal patch} \end{pmatrix} \sim e^{10^{120}}$?!

シック・ 川 ・ 川田・ 小田・ 小田・

Basic questions in theoretical cosmology:

What are the quantum mechanical degrees of freedom of our cosmology?

< □ > < 同 > < Ξ > < Ξ > < Ξ > < Ξ < </p>

- QM role of horizon and regions outside our causal patch?
- Can we understand quantum gravity on cosmological spacetimes?

Basic questions in theoretical cosmology:

- What are the quantum mechanical degrees of freedom of our cosmology?
- QM role of horizon and regions outside our causal patch?
- Can we understand quantum gravity on cosmological spacetimes?

Basic lesson from string theory: quantum gravity can be formulated as a holographic quantum mechanics theory. E.g. AdS/CFT.

→ Our goal: Formulate cosmology holographically.

In this talk we will present concrete steps towards a complete holographic duality for de Sitter and certain FRW cosmologies

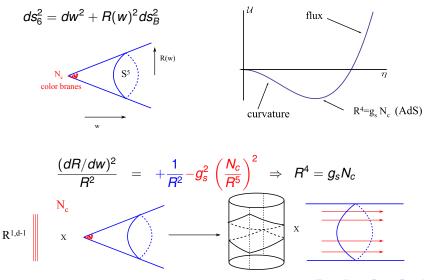
Road map for holographic formulation of cosmology:

▲□▶ ▲□▶ ▲□▶ ▲□▶ ▲□ ● のへぐ

- 1. Uplifting AdS/CFT to de Sitter
- 2. Holography for FRW cosmologies
- 3. RG flow in time-dependent QFT

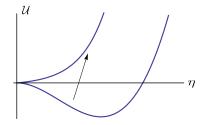
1. Uplifting AdS/CFT to de Sitter

Recall basics of AdS/CFT: Nc D3-branes in ten dimensions



▲□▶▲圖▶▲≣▶▲≣▶ ■ のなぐ

Cosmological sols from known AdS_d/CFT_{d-1} pairs by uplifting:



Need energy sources that compete w/ curvature, and such that they have a QFT interpretation in the d - 1 holographic dual.

Examples:

- cosmic strings in 4d
- magnetic monopole as Hopf fibration of S¹ over S²
- \Rightarrow magnetic flavors in dual QFT

Effective theory of de Sitter

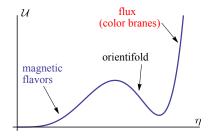
[Silverstein, ...]

String theory has three classes of energy sources:

- $\mathcal{U} \propto g_s^0$: curvature, "geometric" branes
- $\mathcal{U} \propto \pm g_s$: D-branes/orientifold planes
- $\mathcal{U} \propto g_s^2$: fluxes

A metastable de Sitter vacuum with only classical sources requires

◆□▶ ◆□▶ ◆□▶ ◆□▶ ● ● ● ●



Ten-dimensional perspective

• 1) Background geometry for color-branes:

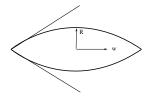
uplifted curvature + orientifold

$$ds^{2} = dw^{2} + R(w)^{2} ds^{2}_{B} \ \Rightarrow \ \frac{R'(w)^{2}}{R^{2}} \sim -\frac{1}{R^{2}} + \frac{g_{s}}{R^{n_{o}}}$$

with $n_O > 2$ related to the codimension of the orientifold

 $\rightsquigarrow R(w)$ first grows, reaches a maximum size, and then decreases

The cone of AdS/CFT has become a compact space!



◆□▶ ◆□▶ ◆□▶ ◆□▶ ● ● ● ●

Ten-dimensional perspective

• 1) Background geometry for color-branes:

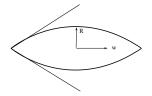
uplifted curvature + orientifold

$$ds^2 = dw^2 + R(w)^2 ds_B^2 \ \Rightarrow \ rac{R'(w)^2}{R^2} \sim -rac{1}{R^2} + rac{g_s}{R^{n_o}}$$

with $n_O > 2$ related to the codimension of the orientifold

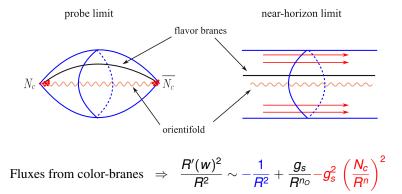
 $\rightsquigarrow R(w)$ first grows, reaches a maximum size, and then decreases

The cone of AdS/CFT has become a compact space!



◆□▶ ◆□▶ ◆□▶ ◆□▶ ● ● ● ●

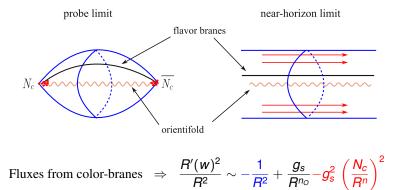
• 2) Charge conservation requires color branes at one tip and anti-branes at the other tip.



Fluxes dominate at small R ($2n > n_O$) preventing the singularities at the tips.

< □ > < 同 > < Ξ > < Ξ > < Ξ > < Ξ < </p>

Concrete example: dS₃ × Y₇ from uplifting D1-D5.
 Y₇: T⁴ fibration over S³.

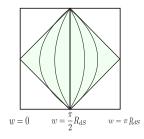


Fluxes dominate at small R ($2n > n_O$) preventing the singularities at the tips.

Concrete example: dS₃ × Y₇ from uplifting D1-D5.
 Y₇: T⁴ fibration over S³.

⇒ dS_d as two throat compactification. Microscopic d.o.f.: Two gauge theories with magnetic flavors, cutoff $1/R_{dS}$, and dynamical gravity in d - 1 dims. • Realizes the static patch of dS_d sliced by dS_{d-1}

$$ds_{dS_d}^2 = dw^2 + \sin^2(w/R_{dS}) ds_{dS_{d-1}}^2$$

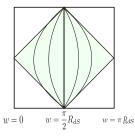


Two dS_{d-1} throats glued at $w = \frac{\pi}{2}R_{dS}$ \Rightarrow dS/dS correspondence of [Alishahiha, Karch, Silverstein, Tong]

▲□▶▲□▶▲□▶▲□▶ □ のQ@

• Realizes the static patch of dS_d sliced by dS_{d-1}

$$ds^2_{dS_d} = dw^2 + \sin^2(w/R_{dS}) ds^2_{dS_{d-1}}$$



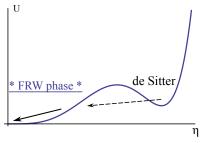
Two dS_{d-1} throats glued at $w = \frac{\pi}{2}R_{dS}$ \Rightarrow dS/dS correspondence of [Alishahiha, Karch, Silverstein, Tong]

Explanation of dS entropy in terms of gauge theory microstates:

$$S_{dS} = (M_{Pl}R_{dS})^{d-2} \sim N_{dof}$$
 of dual QFT

2. Holography for FRW cosmologies

Now focus on the FRW phase, where the 'uplifting' branes give the dominant contribution:



For concreteness, consider $AdS_5 \times S^5$ case.

→ (p,q) 7-branes play role of uplifting ingredient.

View S^5 as S_f^1 fibered over \mathbb{P}^2 . 7-branes wrap $AdS_5 \times S_f^1 \times \Sigma_2 (\subset \mathbb{P}^2)$

◆□▶ ◆□▶ ◆□▶ ◆□▶ ● ● ● ●

• (p,q) 7-branes compete with curvature: they are codimension 2 and have tension $T_7 \sim 1/g_s^2$.

E.g. 24 7-branes exactly cancel the curvature of \mathbb{P}^1 .

For n7-branes,

$$\Delta n \equiv n - n_*$$
 , $\mathcal{U}_R \propto rac{\Delta n}{R^2}$

< □ > < 同 > < Ξ > < Ξ > < Ξ > < Ξ < </p>

• (p,q) 7-branes compete with curvature: they are codimension 2 and have tension $T_7 \sim 1/g_s^2$.

E.g. 24 7-branes exactly cancel the curvature of \mathbb{P}^1 .

For n7-branes,

$$\Delta n \equiv n - n_*$$
, $\mathcal{U}_R \propto \frac{\Delta n}{R^2}$

♦ D3-(p,q)7 gauge theory contains electric and magnetic matter from (3-7) strings [Sen; Banks, Douglas, Seiberg; Argyres, Douglas; ...]

- Δn < 0: AdS/CFT sols. are known [Aharony, Fayazzudin, Maldacena; Polchinski, Silverstein]
- Δn ≥ 0: no static sol in the D3-(p,q)7 theory
 Dim O < unitarity bound; e.g. Dim(u) = ¹²/_{12-n} in SW

However: consistent time-dependent cosmologies for $\Delta n > 0!$

Cosmological 10d solution

Late time solution for $\Delta n > 0$, sourced by magnetic flavor branes (color flux subdominant), in string frame:

$$ds_s^2 = -dt_s^2 + \frac{t_s^2}{c^2}dH_4^2 + \frac{t_s^2}{c^2}dB_4^2 + dx_f^2$$
, $c^2 = \frac{7}{3}$

with B_4 is a compact 4-dim hyperbolic space.

- Internal space: S^5 + 7-branes $\Rightarrow \frac{t_s^2}{c^2} dB_4^2 + dx_f^2$
- 5d spacetime: open FRW (instead of AdS₅)

More general set of FRW solutions:

$$ds_{s}^{2} = -dt_{s}^{2} + \frac{t_{s}^{2}}{c^{2}}dH_{d-1}^{2} + \frac{t_{s}^{2}}{\hat{c}^{2}}dB_{2m}^{2} + dx_{f}^{2}$$

Goal: using this concrete solution, set up the holographic dictionary for FRW.

Warped solution

Basic requirement for holographic dual: warped region that redshifts energies. [Maldacena]

Focus on *d*-dim part and go to Einstein frame metric + change vars ...

$$\left[ds_d^2 = c^2 \left(T^{2/c} - w^2 \right)^{c-1} dw^2 + \left(1 - \frac{w^2}{T^{2/c}} \right)^{c-1} \left(-dT^2 + c^2 T^2 dH_{d-2}^2 \right) \right]$$

- (d 1) dual lives on $ds_{d-1}^2 = -dT^2 + c^2 T^2 dH_{d-2}^2$
- UV slice w = 0, two IR regions $w \to \pm T^{1/c}$:

$$E(w,T) = \left(1 - \frac{w^2}{T^{2/c}}\right)^{\frac{c-1}{2}} E_{pr} \ll M_{Pl}$$

< □ > < 同 > < Ξ > < Ξ > < Ξ > < Ξ < </p>

Dual QFT has time-dep couplings and nontrivial RG.

Properties of the holographic dual

• 1) (d-1)-dimensional gravity: $M_{Pl}^{d-1} \sim T$

At finite times, dual has propagating gravity, but as $T \to \infty$ gravity decouples! Suggests a precise QFT description of FRW physics at late times.

• 2) Field-theoretic degrees of freedom: $\Lambda_c \sim 1/T\,,\,N_{dof} \sim T^{d-2}$

→ System has finite cutoff and accumulates d.o.f. per lattice point. Agrees with parametric counting of magnetic flavors.

3) Correlation functions: semicl, ⟨𝒪(𝑥)𝒪(𝑥')⟩ ≈ exp [−𝔅(𝑥,𝑥')]
 Geodesics become shorter by moving along the radial direction.
 → Power-law correlator for KK modes in time-dependent QFT

3. RG flow in time dependent QFT

In our FRW construction, time-dependence allows for $n_{mag} - n_* > 0$ and leads to new scale invariant regimes.

 $[\text{Recall } n_{\text{mag}} - n_* > 0 \ \Rightarrow \Delta_{\mathcal{O}} < \Delta_{\text{unitary}} \text{ in static theory}]$

QFT question of more general relevance:

How do time-dependent couplings affect the IR dynamics of QFTs?

< □ > < 同 > < Ξ > < Ξ > < Ξ > < Ξ < </p>

3. RG flow in time dependent QFT

In our FRW construction, time-dependence allows for $n_{mag} - n_* > 0$ and leads to new scale invariant regimes.

 $[\text{Recall } n_{\text{mag}} - n_* > 0 \ \Rightarrow \Delta_{\mathcal{O}} < \Delta_{\text{unitary}} \text{ in static theory}]$

QFT question of more general relevance:

How do time-dependent couplings affect the IR dynamics of QFTs?

Tractable t-dep QFT: large-N CFT plus scalar field

•
$$L = L_{CFT} - \frac{1}{2} \left((\partial \phi)^2 + m^2 \phi^2 \right) + g(t) \phi \mathcal{O}_{CFT}$$

 $\Rightarrow L = L_{CFT} + \frac{g(t)^2}{2m^2} \mathcal{O}_{CFT}^2$ for $E \ll m$

Choose dim. of \mathcal{O} at g = 0 fixed pt. by $\Delta_{\mathcal{O}} = d/2 + \nu$, $(\nu > 1)$.

♦ Static case, $g = g_0$

- requiring that $g_0\phi \mathcal{O}$ be marginal in the IR sets

$$\Delta_{\phi} = d - \left(\frac{d}{2} + \nu\right) = \frac{d}{2} - \nu \Rightarrow \text{ violates unitarity!}$$

< □ > < 同 > < 三 > < 三 > < 三 > < ○ < ○ </p>

- instead, g_0 becomes irrelevant, and in the IR we have original CFT plus free (decoupled) scalar ϕ

Choose dim. of \mathcal{O} at g = 0 fixed pt. by $\Delta_{\mathcal{O}} = d/2 + \nu$, $(\nu > 1)$.

♦ Static case, $g = g_0$

- requiring that $g_0\phi \mathcal{O}$ be marginal in the IR sets

$$\Delta_{\phi} = d - \left(\frac{d}{2} + \nu\right) = \frac{d}{2} - \nu \Rightarrow \text{ violates unitarity!}$$

- instead, g_0 becomes irrelevant, and in the IR we have original CFT plus free (decoupled) scalar ϕ

\blacklozenge t-dep case, $g ightarrow g_0 t^{lpha}$

- no effect for $E \gg \partial g/g$
- for $\alpha > \nu$, g(t) dominates the IR, reversing direction of flow
- two-point functions can be calculated exactly at large-N

$$\langle \phi(\mathbf{x})\phi(\mathbf{x}') \rangle \rightarrow rac{\mathrm{const}}{t^{lpha}|\mathbf{x}-\mathbf{x}'|^{d-2
u}t'^{lpha}} \ \Rightarrow \ \Delta_{I\!R}(\phi) = rac{d}{2} - \nu + lpha$$

4. Conclusions and future directions

- We have constructed a holographic duality for de Sitter and certain FRW cosmologies.
- Holographic description in terms of two coupled QFTs with magnetic flavors. Explains entropy of cosmological sols.
- Time-dependent couplings can strongly affect the IR physics and induce novel scale invariant regimes
- Central role played by magnetic flavors. More general relevance for cosmological sols?
- Develop further the holographic description, with time-dependent and running couplings.

(ロ) (同) (三) (三) (三) (○) (○)

• Construct realistic dS4 solutions.