

Heterotic Line Bundle Models in Four Dimensions



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model building: [arXiv:1106.4804](https://arxiv.org/abs/1106.4804), [arXiv:1202.1757](https://arxiv.org/abs/1202.1757) with Lara Anderson, James Gray and Eran Palti

moduli stabilization: [arXiv:0903.5088](https://arxiv.org/abs/0903.5088), [arXiv:0905.1748](https://arxiv.org/abs/0905.1748), [arXiv:1010.0255](https://arxiv.org/abs/1010.0255), [arXiv:1102.0011](https://arxiv.org/abs/1102.0011), [arXiv:1107.5076](https://arxiv.org/abs/1107.5076), with Lara Anderson, James Gray and Burt Ovrut

earlier work: [hep-th/0702210](https://arxiv.org/abs/hep-th/0702210), [arXiv:0805.2875](https://arxiv.org/abs/0805.2875), [arXiv:0911.1569](https://arxiv.org/abs/0911.1569), . . .
with Lara Anderson, James Gray, Yang-Hui He

Overview

- Introduction: Constructing Line Bundle Models
- Four-dimensional effective theory
- Phenomenological issues and examples
- Moduli stabilization
- Conclusion and outlook

Introduction

Data to define a heterotic line bundle model we need:

- A Calabi-Yau 3-fold X
- A line bundle sum $V = L_1 \oplus \cdots \oplus L_5$ on X ,
 $c_1(V) = 0$, so structure group is $S(U(1)^5)$.
- vanishing slopes $\mu(L_a) \equiv c_1(L_a) \wedge J^2 \stackrel{!}{=} 0$
- Anomaly: $c_2(TX) - c_2(V) - c_2(\tilde{V}) = [C]$
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standard-like model
(hopefully) with
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- Additional U(1)s are usually Green-Schwarz anomalous, hence, associated vector bosons are massive.
- Abelian bundles still carries many of the properties of the generic non-Abelian bundles.
- But we have to remember they are usually part of a moduli space of non-Abelian bundles.

Recall: 4d gauge group is $SU(5) \times S(U(1)^5) \cong SU(5) \times U(1)^4$

Label $S(U(1)^5)$ repr. by integer vectors $\mathbf{q} = (q_1, \dots, q_5)$ with identification $\mathbf{q} \sim \tilde{\mathbf{q}}$ iff $\mathbf{q} - \tilde{\mathbf{q}} \in \mathbb{Z}(1, 1, 1, 1, 1)$.

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for example: quintic $X \sim [\mathbb{P}^4 | 5]$ or bi-cubic $X \sim \left[\begin{array}{c|c} \mathbb{P}^2 & 3 \\ \hline \mathbb{P}^2 & 3 \end{array} \right]$

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Classification of freely-acting discrete symmetries

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Have scanned CICYs with freely-acting symmetries and $h^{1,1}(X) \leq 5$ (60 spaces) and line bundle sums $-9 + h^{1,1}(X) \leq k_a^i \leq 0 - h^{1,1}(X)$.

On CICYs with $h^{1,1}(X) = 4, 5$ we find
 ~ 2000 standard models.*

*standard model: SM gauge group times (anomalous) U(1)s, exact MSSM matter spectrum, one or more pairs of Higgs doublets, no exotics charged under standard model group.

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There will be many more!

(L. Anderson, A. Constantin, J. Gray, AL, E. Palti, in progress)

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matter

$$(\mathbf{10}^p) = (Q^p, u^p, e^p), \quad (\bar{\mathbf{5}}^p) = (d^p, L^p), \quad H, \quad \bar{H}, \quad S^\alpha \quad (p = 1, 2, 3)$$

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No exotics charged under SM group, no vector-like pairs etc.

moduli

$$\mathcal{S} = s + i\sigma, \quad T^i = t^i + 2i\chi^i, \quad Z, \quad \dots$$

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D-terms

$$D_a = k_a^i t_i - \sum_{\alpha} Q_a(S^{\alpha}) |S^{\alpha}|^2$$

4d U(1)s

bundle

all $S^{\alpha} = 0$:

$$U(1)^4$$

$$S(U(1)^5)$$

moduli

$$\mathcal{S} = s + i\sigma, \quad T^i = t^i + 2i\chi^i, \quad Z, \quad \dots$$

$$\delta_a \chi^i = -k_a^i \quad \delta_a \sigma = -k_a^i \beta_i$$

U(1) masses

$$M_{ab} = \mathbf{k}_a^T G \mathbf{k}_b$$

$$(\text{number of massless U(1)s}) = 4 - \text{rank}(k_a^i)$$

$$> 4 - (\text{number of } T^i)$$

D-terms

$$D_a = k_a^i t_i - \sum_{\alpha} Q_a(S^{\alpha}) |S^{\alpha}|^2$$

4d U(1)s

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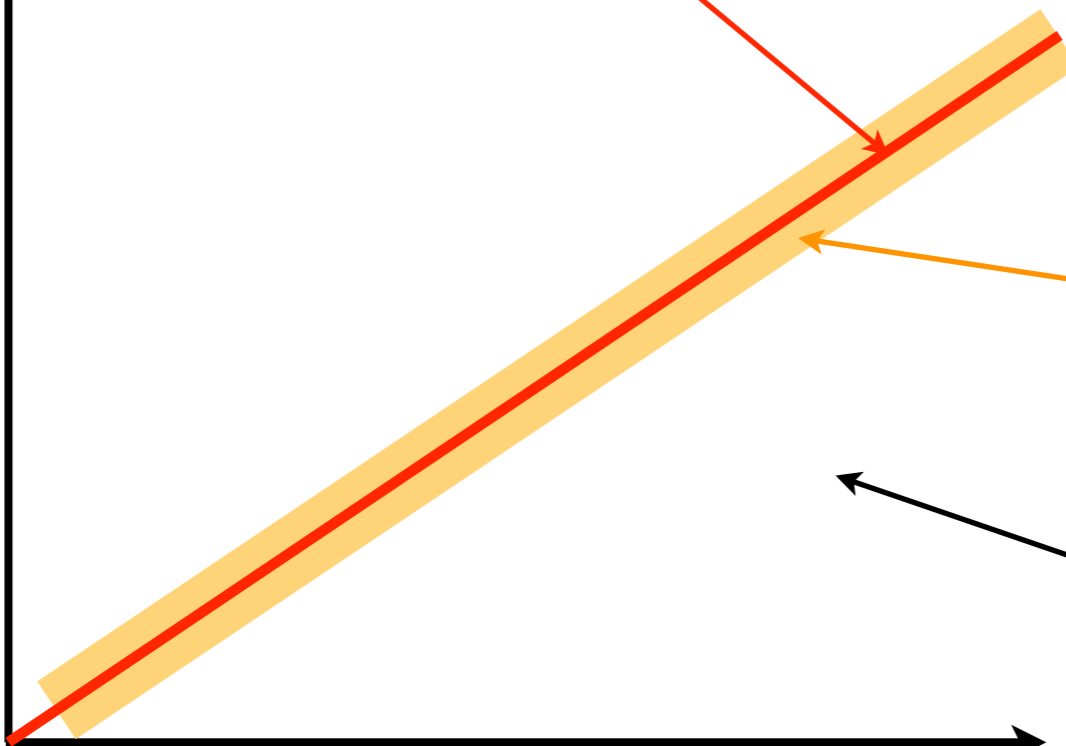
some $S^{\alpha} \neq 0$:

$$U(1)^{f-1}$$

$$S(U(n_1) \times \dots \times U(n_f))$$

t_2

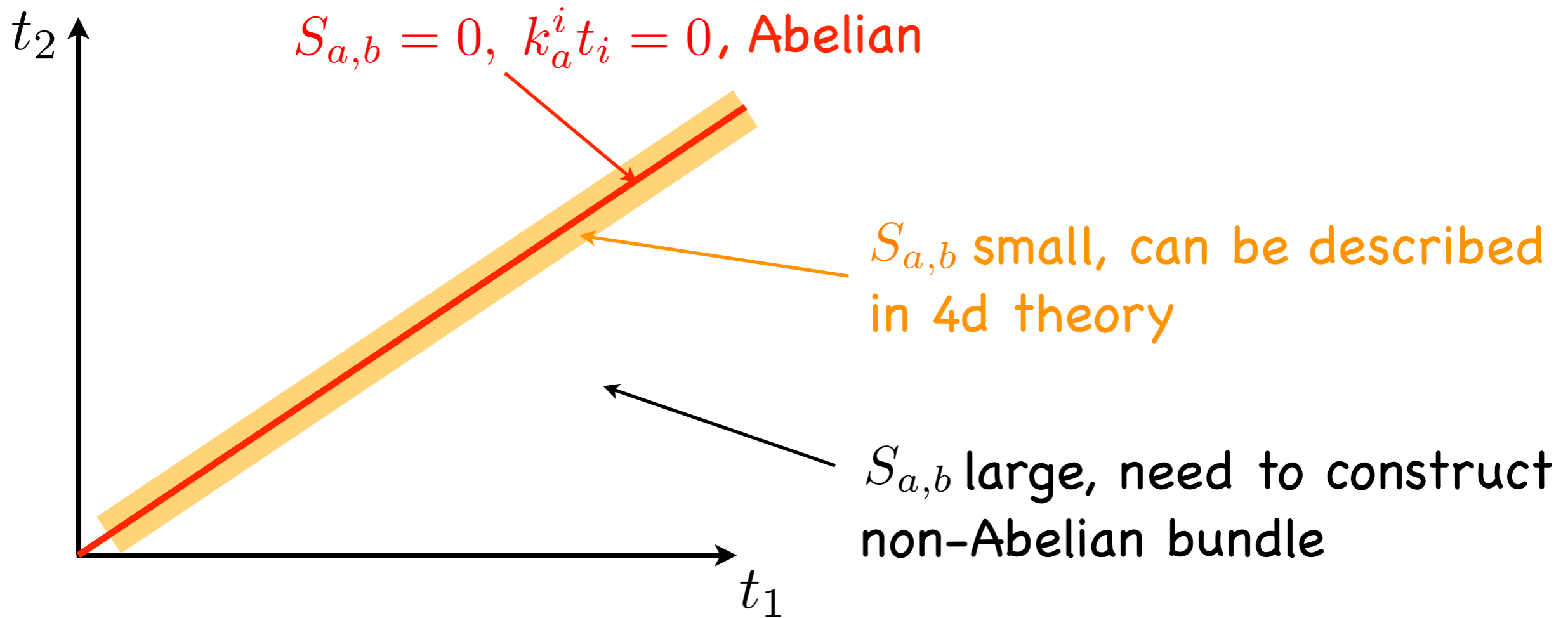
$S_{a,b} = 0, k_a^i t_i = 0, \text{Abelian}$



$S_{a,b}$ small, can be described
in 4d theory

$S_{a,b}$ large, need to construct
non-Abelian bundle

t_1



Example: $S_{4,5} \neq 0$

4d U(1)s: $U(1)^4 \rightarrow U(1)^3$

structure group: $S(U(1)^5) \rightarrow S(U(1)^3 \times U(2))$

bundle: $V = L_1 \oplus \cdots \oplus L_5 \rightarrow V = L_1 \oplus L_2 \oplus L_3 \oplus U$

$$0 \rightarrow L_4 \rightarrow U \rightarrow L_5 \rightarrow 0, \quad \text{Ext}^1(L_5, L_4) \cong H^1(X, L_4 \otimes L_5^*) \ni S_{4,5}$$

Constraints from $S(U(1)^5)$, e.g. superpotential:

$$W = \mu(S, Z) H \bar{H} + Y_{pq}^{(d)}(S, Z) H \bar{\mathbf{5}}^p \mathbf{10}^q + Y_{pq}^{(u)}(S, Z) \bar{H} \mathbf{10}^p \mathbf{10}^q$$

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$$\mu = \mu_0 + \mu_{1,\alpha} S^\alpha + \mu_{2,\alpha\beta} S^\alpha S^\beta + \dots$$

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Texture at each order in S^α

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$$Y = Y_0 + Y_{1,\alpha}S^\alpha + Y_{2,\alpha\beta}S^\alpha S^\beta + \dots$$

Texture at each order in S^α

How to get fermion masses?

top Yukawa: $\text{rk}(Y_0) > 0$ but only $\text{rk}(Y_0) = 0, 2, 3$ possible:

$$\bar{H}_{a,b}\mathbf{10}_a\mathbf{10}_b \text{ but } a \neq b \text{ since } \wedge^2 V = \sum_{a < b} L_a \otimes L_b$$

Two possible ways out:

- 1) Non-trivial equivariant blocks of line bundles \rightarrow analogue of F-theory monodromy.

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$$\wedge^2 U^* \neq 0 \Rightarrow \text{charge } -2\mathbf{e}_4$$

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How to generate other masses:

If some S^α small: Froggatt-Nielsen mechanism

If all S^α large: non-perturbative effects

... or combination of both

Note: $e^{-n_i T^i} \times (\text{matter})$ constrained by U(1)s as well

Phenomenological issues and example

CY data: ■ Cicy 6777, Symmetry 1

$$X = \begin{pmatrix} 1 & 1 & 0 & 0 \\ 0 & 0 & 0 & 2 \\ 0 & 0 & 2 & 0 \\ 2 & 0 & 0 & 0 \\ 1 & 1 & 1 & 1 \end{pmatrix}$$

$$\eta(X) = -64 \quad h^{1,1}(X) = 5 \quad h^{2,1}(X) = 37 \quad c_2(TX) = \{24, 24, 24, 24, 56\}$$

$$\kappa = 12 t_1 t_2 t_3 + 12 t_1 t_2 t_4 + 12 t_1 t_3 t_4 + 12 t_2 t_3 t_4 + 24 t_1 t_2 t_5 + 24 t_1 t_3 t_5 + 12 t_2 t_3 t_5 + 24 t_1 t_4 t_5 + 24 t_2 t_4 t_5 + 24 t_3 t_4 t_5 + 24 t_1 t_5^2 + 12 t_2 t_5^2 + 12 t_3 t_5^2 + 24 t_4 t_5^2 + 8 t_5^3$$

symmetry: 1 order: 2

Abelian: True block diagonal: True factors: {2}

$$\text{Action on coordinates: } \left\{ \begin{pmatrix} -1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & -1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & -1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & -1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 \end{pmatrix} \right\}$$

$$\text{Action on polynomials: } \left\{ \begin{pmatrix} -1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix} \right\}$$

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Phenomenological issues and example

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← volume

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Phenomenological issues and example

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$$X = \begin{pmatrix} 1 & 1 & 0 & 0 \\ 0 & 0 & 0 & 2 \\ 0 & 0 & 2 & 0 \\ 2 & 0 & 0 & 0 \\ 1 & 1 & 1 & 1 \end{pmatrix}$$

← definition of CY in $(\mathbb{P}^1)^{\times 4} \times \mathbb{P}^3$

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← \mathbb{Z}_2 symmetry

Action on polynomials: $\left\{ \begin{pmatrix} -1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix} \right\}$

bundle data:

■ Model number 13, Identifier {6777, 2, 1}

■ Basic properties

standard model? **True** massless U(1): **0** number of $5 \bar{5}$ pairs: **1** $c_2(V) = \{8, 6, 6, 14, 24\}$

$$V: (k_a^i) = \begin{pmatrix} 1 & 1 & 0 & 0 & -2 \\ 0 & -1 & 1 & -1 & 1 \\ 0 & -2 & 1 & 1 & 0 \\ -1 & 0 & 1 & 0 & 0 \\ 0 & 1 & -1 & 0 & 0 \end{pmatrix}$$

Cohomology of V:

L_2	=	{1, -1, -2, 0, 1}	$h[L_2]$	=	{0, 4, 0, 0}	$h[L_2, R]$	=	{{0, 0}, {2, 2}, {0, 0}, {0, 0}}
L_5	=	{-2, 1, 0, 0, 0}	$h[L_5]$	=	{0, 2, 0, 0}	$h[L_5, R]$	=	{{0, 0}, {1, 1}, {0, 0}, {0, 0}}
$L_2 \times L_4$	=	{1, -2, -1, 0, 1}	$h[L_2 \times L_4]$	=	{0, 4, 0, 0}	$h[L_2 \times L_4, R]$	=	{{0, 0}, {2, 2}, {0, 0}, {0, 0}}
$L_2 \times L_5$	=	{-1, 0, -2, 0, 1}	$h[L_2 \times L_5]$	=	{0, 1, 1, 0}	$h[L_2 \times L_5, R]$	=	{{0, 0}, {1, 0}, {1, 0}, {0, 0}}
$L_4 \times L_5$	=	{-2, 0, 1, 0, 0}	$h[L_4 \times L_5]$	=	{0, 2, 0, 0}	$h[L_4 \times L_5, R]$	=	{{0, 0}, {1, 1}, {0, 0}, {0, 0}}
$L_1 \times L_2^*$	=	{0, 1, 2, -1, -1}	$h[L_1 \times L_2^*]$	=	{0, 0, 4, 0}	$h[L_1 \times L_2^*, R]$	=	{{0, 0}, {0, 0}, {2, 2}, {0, 0}}
$L_1 \times L_5^*$	=	{3, -1, 0, -1, 0}	$h[L_1 \times L_5^*]$	=	{0, 0, 8, 0}	$h[L_1 \times L_5^*, R]$	=	{{0, 0}, {0, 0}, {4, 4}, {0, 0}}
$L_2 \times L_3^*$	=	{1, -2, -3, -1, 2}	$h[L_2 \times L_3^*]$	=	{0, 4, 0, 0}	$h[L_2 \times L_3^*, R]$	=	{{0, 0}, {2, 2}, {0, 0}, {0, 0}}
$L_2 \times L_4^*$	=	{1, 0, -3, 0, 1}	$h[L_2 \times L_4^*]$	=	{0, 12, 0, 0}	$h[L_2 \times L_4^*, R]$	=	{{0, 0}, {6, 6}, {0, 0}, {0, 0}}
$L_2 \times L_5^*$	=	{3, -2, -2, 0, 1}	$h[L_2 \times L_5^*]$	=	{0, 11, 3, 0}	$h[L_2 \times L_5^*, R]$	=	{{0, 0}, {-2 + r1, 13 - r1}, {-6 + r1, 9 - r1}, {0, 0}}
$L_4 \times L_5^*$	=	{2, -2, 1, 0, 0}	$h[L_4 \times L_5^*]$	=	{0, 6, 0, 0}	$h[L_4 \times L_5^*, R]$	=	{{0, 0}, {3, 3}, {0, 0}, {0, 0}}

Wilson line: **{0}, {1}** Equivariant structure: **{0}, {0}, {0}, {0}, {1}** Higgs pairs: **1**

Downstairs spectrum: **{2 10₂, 10₅, 2 $\bar{5}_{2,4}$, $\bar{5}_{4,5}$, H_{2,5}, $\bar{H}_{2,5}$, 2 S_{2,1}, 4 S_{5,1}, 2 S_{2,3}, 6 S_{2,4}, (13 - r1) S_{2,5}, (9 - r1) S_{5,2}, 3 S_{4,5}}}** Phys. Higgs: **{H_{2,5}, $\bar{H}_{2,5}$ }**

Legend: {r1 → rk₁[H₂[N* × I2 × I5*, 2], H₂[I2 × I5*, 2]]}

Transfer format: **{6, 1, 1, 4, 6, 5, 5, 9, 8, 10, 12, 18, 1, 7, 13}, {6, 6, -1, -1, -1, -1}**

rk(Y^(u)) = **{2, 3}** rk(Y^(d)) = **{0, 0}** dim. 4 operators absent: **{True, True}** dim. 5 operators absent: **{True, True}**

bundle data:

■ **Model number 13, Identifier {6777, 2, 1}**

■ **Basic properties**

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← integer matrix defining line bundle sum

Cohomology of V:

L_2	= {1, -1, -2, 0, 1}	$h[L_2]$	= {0, 4, 0, 0}	$h[L_2, R]$	= {{0, 0}, {2, 2}, {0, 0}, {0, 0}}
L_5	= {-2, 1, 0, 0, 0}	$h[L_5]$	= {0, 2, 0, 0}	$h[L_5, R]$	= {{0, 0}, {1, 1}, {0, 0}, {0, 0}}
$L_2 \times L_4$	= {1, -2, -1, 0, 1}	$h[L_2 \times L_4]$	= {0, 4, 0, 0}	$h[L_2 \times L_4, R]$	= {{0, 0}, {2, 2}, {0, 0}, {0, 0}}
$L_2 \times L_5$	= {-1, 0, -2, 0, 1}	$h[L_2 \times L_5]$	= {0, 1, 1, 0}	$h[L_2 \times L_5, R]$	= {{0, 0}, {1, 0}, {1, 0}, {0, 0}}
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$L_1 \times L_5^*$	= {3, -1, 0, -1, 0}	$h[L_1 \times L_5^*]$	= {0, 0, 8, 0}	$h[L_1 \times L_5^*, R]$	= {{0, 0}, {0, 0}, {4, 4}, {0, 0}}
$L_2 \times L_3^*$	= {1, -2, -3, -1, 2}	$h[L_2 \times L_3^*]$	= {0, 4, 0, 0}	$h[L_2 \times L_3^*, R]$	= {{0, 0}, {2, 2}, {0, 0}, {0, 0}}
$L_2 \times L_4^*$	= {1, 0, -3, 0, 1}	$h[L_2 \times L_4^*]$	= {0, 12, 0, 0}	$h[L_2 \times L_4^*, R]$	= {{0, 0}, {6, 6}, {0, 0}, {0, 0}}
$L_2 \times L_5^*$	= {3, -2, -2, 0, 1}	$h[L_2 \times L_5^*]$	= {0, 11, 3, 0}	$h[L_2 \times L_5^*, R]$	= {{0, 0}, {-2 + r1, 13 - r1}, {-6 + r1, 9 - r1}, {0, 0}}
$L_4 \times L_5^*$	= {2, -2, 1, 0, 0}	$h[L_4 \times L_5^*]$	= {0, 6, 0, 0}	$h[L_4 \times L_5^*, R]$	= {{0, 0}, {3, 3}, {0, 0}, {0, 0}}

Wilson line: **{{0}, {1}}** Equivariant structure: **{{0}, {0}, {0}, {0}, {1}}** Higgs pairs: **1**

Downstairs spectrum: **{2 10₂, 10₅, 2 $\bar{5}_{2,4}$, $\bar{5}_{4,5}$, H_{2,5}, $\bar{H}_{2,5}$, 2 S_{2,1}, 4 S_{5,1}, 2 S_{2,3}, 6 S_{2,4}, (13 - r1) S_{2,5}, (9 - r1) S_{5,2}, 3 S_{4,5}}}** Phys. Higgs: **{H_{2,5}, $\bar{H}_{2,5}}$ }**

Legend: {r1 → rk₁[H₂[N* × I₂ × I₅*, 2], H₂[I₂ × I₅*, 2]]}

Transfer format: **{{6, 1, 1, 4, 6, 5, 5, 9, 8, 10, 12, 18, 1, 7, 13}, {6, 6, -1, -1, -1, -1}}**

rk(Y^(u)) = **{2, 3}** rk(Y^(d)) = **{0, 0}** dim. 4 operators absent: **{True, True}** dim. 5 operators absent: **{True, True}**

bundle data:

■ **Model number 13, Identifier {6777, 2, 1}**

■ **Basic properties**

standard model? **True** massless U(1): **0** number of $5 \bar{5}$ pairs: **1** $c_2(V) = \{8, 6, 6, 14, 24\}$

$$V: (k_a^i) = \begin{pmatrix} 1 & 1 & 0 & 0 & -2 \\ 0 & -1 & 1 & -1 & 1 \\ 0 & -2 & 1 & 1 & 0 \\ -1 & 0 & 1 & 0 & 0 \\ 0 & 1 & -1 & 0 & 0 \end{pmatrix}$$

← integer matrix defining line bundle sum

Cohomology of V:

L_2	= {1, -1, -2, 0, 1}	$h[L_2]$	= {0, 4, 0, 0}	$h[L_2, R]$	= {{0, 0}, {2, 2}, {0, 0}, {0, 0}}
L_5	= {-2, 1, 0, 0, 0}	$h[L_5]$	= {0, 2, 0, 0}	$h[L_5, R]$	= {{0, 0}, {1, 1}, {0, 0}, {0, 0}}
$L_2 \times L_4$	= {1, -2, -1, 0, 1}	$h[L_2 \times L_4]$	= {0, 4, 0, 0}	$h[L_2 \times L_4, R]$	= {{0, 0}, {2, 2}, {0, 0}, {0, 0}}
$L_2 \times L_5$	= {-1, 0, -2, 0, 1}	$h[L_2 \times L_5]$	= {0, 1, 1, 0}	$h[L_2 \times L_5, R]$	= {{0, 0}, {1, 0}, {1, 0}, {0, 0}}
$L_4 \times L_5$	= {-2, 0, 1, 0, 0}	$h[L_4 \times L_5]$	= {0, 2, 0, 0}	$h[L_4 \times L_5, R]$	= {{0, 0}, {1, 1}, {0, 0}, {0, 0}}
$L_1 \times L_2^*$	= {0, 1, 2, -1, -1}	$h[L_1 \times L_2^*]$	= {0, 0, 4, 0}	$h[L_1 \times L_2^*, R]$	= {{0, 0}, {0, 0}, {2, 2}, {0, 0}}
$L_1 \times L_5^*$	= {3, -1, 0, -1, 0}	$h[L_1 \times L_5^*]$	= {0, 0, 8, 0}	$h[L_1 \times L_5^*, R]$	= {{0, 0}, {0, 0}, {4, 4}, {0, 0}}
$L_2 \times L_3^*$	= {1, -2, -3, -1, 2}	$h[L_2 \times L_3^*]$	= {0, 4, 0, 0}	$h[L_2 \times L_3^*, R]$	= {{0, 0}, {2, 2}, {0, 0}, {0, 0}}
$L_2 \times L_4^*$	= {1, 0, -3, 0, 1}	$h[L_2 \times L_4^*]$	= {0, 12, 0, 0}	$h[L_2 \times L_4^*, R]$	= {{0, 0}, {6, 6}, {0, 0}, {0, 0}}
$L_2 \times L_5^*$	= {3, -2, -2, 0, 1}	$h[L_2 \times L_5^*]$	= {0, 11, 3, 0}	$h[L_2 \times L_5^*, R]$	= {{0, 0}, {-2 + r1, 13 - r1}, {-6 + r1, 9 - r1}, {0, 0}}
$L_4 \times L_5^*$	= {2, -2, 1, 0, 0}	$h[L_4 \times L_5^*]$	= {0, 6, 0, 0}	$h[L_4 \times L_5^*, R]$	= {{0, 0}, {3, 3}, {0, 0}, {0, 0}}

Wilson line: **{{0}, {1}}** Equivariant structure: **{{0}, {0}, {0}, {0}, {1}}** Higgs pairs: **1**

Downstairs spectrum: **{2 10₂, 10₅, 2 $\bar{5}_{2,4}$, $\bar{5}_{4,5}$, H_{2,5}, $\bar{H}_{2,5}$, 2 S_{2,1}, 4 S_{5,1}, 2 S_{2,3}, 6 S_{2,4}, (13 - r1) S_{2,5}, (9 - r1) S_{5,2}, 3 S_{4,5}}}** Phys. Higgs: **{H_{2,5}, $\bar{H}_{2,5}$ }**

Legend: {r1 → rk₁[H₂[N* × I2 × I5*, 2], H₂[I2 × I5*, 2]]}

Transfer format: **{{6, 1, 1, 4, 6, 5, 5, 9, 8, 10, 12, 18, 1, 7, 13}, {6, 6, -1, -1, -1, -1}}**

rk(Y^(u)) = **{2, 3}** rk(Y^(d)) = **{0, 0}** dim. 4 operators absent: **{True, True}** dim. 5 operators absent: **{True, True}**

**spectrum: 10₂, 10₂, 10₅, $\bar{5}_{2,4}$, $\bar{5}_{2,4}$, $\bar{5}_{4,5}$, H_{2,5}, $\bar{H}_{2,5}$,
S_{5,1}, S_{2,3}, S_{2,4}, S_{2,5}, S_{5,2}, S_{4,5}**

allowed operators (1):

■ Operators

basic superpotential terms:

rank 2 for $\langle S^\alpha \rangle = 0$

$$\overline{H}10^p 10^q: Y^{(u)} = \begin{pmatrix} \{S_{5,2}, S_{2,5} S_{5,2}^2, S_{2,4} S_{4,5} S_{5,2}^2, S_{2,5}^2 S_{5,2}^3\} & \{S_{5,2}, S_{2,5} S_{5,2}^2, S_{2,4} S_{4,5} S_{5,2}^2, S_{2,5}^2 S_{5,2}^3\} & \{1, S_{2,5} S_{5,2}, S_{2,4} S_{4,5} S_{5,2}, S_{2,5}^2 S_{5,2}^2, S_{2,4} S_{2,5} S_{4,5} S_{5,2}^2\} \\ \{S_{5,2}, S_{2,5} S_{5,2}^2, S_{2,4} S_{4,5} S_{5,2}^2, S_{2,5}^2 S_{5,2}^3\} & \{S_{5,2}, S_{2,5} S_{5,2}^2, S_{2,4} S_{4,5} S_{5,2}^2, S_{2,5}^2 S_{5,2}^3\} & \{1, S_{2,5} S_{5,2}, S_{2,4} S_{4,5} S_{5,2}, S_{2,5}^2 S_{5,2}^2, S_{2,4} S_{2,5} S_{4,5} S_{5,2}^2\} \\ \{1, S_{2,5} S_{5,2}, S_{2,4} S_{4,5} S_{5,2}, S_{2,5}^2 S_{5,2}^2, S_{2,4} S_{2,5} S_{4,5} S_{5,2}^2\} & \{1, S_{2,5} S_{5,2}, S_{2,4} S_{4,5} S_{5,2}, S_{2,5}^2 S_{5,2}^2, S_{2,4} S_{2,5} S_{4,5} S_{5,2}^2\} & \{S_{2,5}, S_{2,4} S_{4,5}, S_{2,5}^2 S_{5,2}, S_{2,4} S_{2,5} S_{4,5} S_{5,2}, S_{2,4}^2 S_{4,5}^2 S_{5,2}, S_{2,5}^3 S_{5,2}^2\} \end{pmatrix}$$

$$H\overline{5}^p 10^q: Y^{(d)} = \begin{pmatrix} (0) & (0) & (0) \\ (0) & (0) & (0) \\ (0) & (0) & (0) \end{pmatrix} \leftarrow \text{rank 0}$$

$$H\overline{H}: \mu = \{1, S_{2,5} S_{5,2}, S_{2,4} S_{4,5} S_{5,2}, S_{2,5}^2 S_{5,2}^2, S_{2,4} S_{2,5} S_{4,5} S_{5,2}^2\} \leftarrow \mu\text{-term}$$

$$W_{\text{sing}} = \{S_{2,4} S_{4,5} S_{5,2}, S_{2,5}^2 S_{5,2}^2, S_{2,4} S_{2,5} S_{4,5} S_{5,2}^2\}$$

Majorana masses for RH neutrinos

R-parity violating terms in superpotential:

$$\overline{H}L^p: \rho = \begin{pmatrix} \{S_{2,4} S_{5,2}, S_{2,4} S_{2,5} S_{5,2}^2, S_{2,4}^2 S_{4,5} S_{5,2}^2\} \\ \{S_{2,4} S_{5,2}, S_{2,4} S_{2,5} S_{5,2}^2, S_{2,4}^2 S_{4,5} S_{5,2}^2\} \\ \{S_{2,4}, S_{2,4} S_{2,5} S_{5,2}, S_{2,4}^2 S_{4,5} S_{5,2}, S_{2,4} S_{2,5}^2 S_{5,2}^2\} \end{pmatrix}$$

$$10^p \overline{5}^q \overline{5}^r: \lambda = \{ \{ \{ \{ (0), (0), (0) \}, \{ (0), (0), (0) \}, \{ (0), (0), (0) \} \}, \{ \{ (0), (0), (0) \}, \{ (0), (0), (0) \}, \{ (0), (0), (0) \} \}, \{ \{ (0), (0), (0) \}, \{ (0), (0), (0) \}, \{ (0), (0), (0) \} \}, \{ \{ (0), (0), (0) \}, \{ (0), (0), (0) \}, \{ (0), (0), (0) \} \}, \{ \{ (0), (0), (0) \}, \{ (0), (0), (0) \}, \{ (0), (0), (0) \} \} \}$$

Dimension 5 operators in superpotential:

$$\overline{5}^p 10^q 10^r 10^s: \lambda' = \{ \{ \{ \{ \{ (0), (0), (0) \}, \{ (0), (0), (0) \}, \{ (0), (0), (0) \} \}, \{ \{ (0), (0), (0) \}, \{ (0), (0), (0) \}, \{ (0), (0), (0) \} \}, \{ \{ (0), (0), (0) \}, \{ (0), (0), (0) \}, \{ (0), (0), (0) \} \}, \{ \{ (0), (0), (0) \}, \{ (0), (0), (0) \}, \{ (0), (0), (0) \} \}, \{ \{ (0), (0), (0) \}, \{ (0), (0), (0) \}, \{ (0), (0), (0) \} \} \}, \{ \{ \{ (0), (0), (0) \}, \{ (0), (0), (0) \}, \{ (0), (0), (0) \} \}, \{ \{ (0), (0), (0) \}, \{ (0), (0), (0) \}, \{ (0), (0), (0) \} \}, \{ \{ (0), (0), (0) \}, \{ (0), (0), (0) \}, \{ (0), (0), (0) \} \}, \{ \{ (0), (0), (0) \}, \{ (0), (0), (0) \}, \{ (0), (0), (0) \} \}, \{ \{ (0), (0), (0) \}, \{ (0), (0), (0) \}, \{ (0), (0), (0) \} \} \}$$

D-terms:

$$FI\text{-terms: } k^i_{aK_i} = \begin{pmatrix} -4 t_1 t_2 - 4 t_1 t_3 + 4 t_2 t_4 + 4 t_3 t_4 - 8 t_1 t_5 + 8 t_4 t_5 \\ 4 t_1 t_3 + 8 t_2 t_3 - 4 t_1 t_4 + 4 t_2 t_4 + 8 t_3 t_4 - 8 t_1 t_5 + 8 t_2 t_5 + 12 t_3 t_5 + 4 t_5^2 \\ -4 t_2 t_4 - 4 t_3 t_4 + 8 t_1 t_5 + 4 t_2 t_5 + 4 t_3 t_5 + 8 t_5^2 \\ 4 t_1 t_2 - 4 t_1 t_3 + 4 t_2 t_4 - 4 t_3 t_4 + 4 t_2 t_5 - 4 t_3 t_5 \\ 4 t_1 t_3 - 8 t_2 t_3 + 4 t_1 t_4 - 8 t_2 t_4 - 4 t_3 t_4 + 8 t_1 t_5 - 16 t_2 t_5 - 12 t_3 t_5 - 8 t_4 t_5 - 12 t_5^2 \end{pmatrix}$$

$$\text{singlet D-terms: } q_{\alpha a} S^\alpha \overline{S}^{\beta} = \begin{pmatrix} -S_{2,1} S_{2,1}^\dagger - S_{5,1} S_{5,1}^\dagger \\ S_{2,1} S_{2,1}^\dagger + S_{2,3} S_{2,3}^\dagger + S_{2,4} S_{2,4}^\dagger + S_{2,5} S_{2,5}^\dagger - S_{5,2} S_{5,2}^\dagger \\ -S_{2,3} S_{2,3}^\dagger \\ -S_{2,4} S_{2,4}^\dagger + S_{4,5} S_{4,5}^\dagger \\ -S_{2,5} S_{2,5}^\dagger - S_{4,5} S_{4,5}^\dagger + S_{5,1} S_{5,1}^\dagger + S_{5,2} S_{5,2}^\dagger \end{pmatrix}$$

allowed operators (2):

Kinetic terms:

$$\text{GM term: } \tilde{\mu} = \{S_{2,5}^\dagger S_{5,2}^\dagger, S_{2,5} S_{5,1} S_{2,1}^\dagger, S_{2,1} S_{5,2} S_{5,1}^\dagger, S_{2,4} S_{4,5} S_{2,5}^\dagger, S_{5,1} S_{2,1}^\dagger S_{5,2}^\dagger, S_{2,1} S_{2,5}^\dagger S_{5,1}^\dagger, S_{2,5} S_{2,4}^\dagger S_{4,5}^\dagger, S_{2,4}^\dagger S_{4,5}^\dagger S_{5,2}^\dagger\}$$

$$\begin{aligned} \bar{5}^p \bar{5}^{\bar{q}\dagger}: K^{(\bar{5})} = & \left\{ \left\{ 1, S_{2,5} S_{5,2}, S_{2,5}^\dagger S_{5,2}^\dagger, S_{2,5} S_{5,1} S_{2,1}^\dagger, S_{2,1} S_{2,5}^\dagger S_{5,1}^\dagger, S_{2,1} S_{5,2} S_{5,1}^\dagger, S_{5,1} S_{2,1}^\dagger S_{5,2}^\dagger, S_{2,4} S_{4,5} S_{2,5}^\dagger, S_{2,5} S_{2,4}^\dagger S_{4,5}^\dagger, S_{2,4} S_{4,5} S_{5,2}, S_{2,4}^\dagger S_{4,5}^\dagger S_{5,2}^\dagger \right\}, \right. \\ & \left. \left\{ 1, S_{2,5} S_{5,2}, S_{2,5}^\dagger S_{5,2}^\dagger, S_{2,5} S_{5,1} S_{2,1}^\dagger, S_{2,1} S_{2,5}^\dagger S_{5,1}^\dagger, S_{2,1} S_{5,2} S_{5,1}^\dagger, S_{5,1} S_{2,1}^\dagger S_{5,2}^\dagger, S_{2,4} S_{4,5} S_{2,5}^\dagger, S_{2,5} S_{2,4}^\dagger S_{4,5}^\dagger, S_{2,4} S_{4,5} S_{5,2}, S_{2,4}^\dagger S_{4,5}^\dagger S_{5,2}^\dagger \right\}, \right. \\ & \left. \left\{ S_{2,5}, S_{5,2}, S_{2,1} S_{5,1}, S_{2,4} S_{4,5}, S_{2,5}^2 S_{5,2}, S_{2,5}^\dagger (S_{5,2}^\dagger)^2 \right\}, \right. \\ & \left. \left\{ 1, S_{2,5} S_{5,2}, S_{2,5}^\dagger S_{5,2}^\dagger, S_{2,5} S_{5,1} S_{2,1}^\dagger, S_{2,1} S_{2,5}^\dagger S_{5,1}^\dagger, S_{2,1} S_{5,2} S_{5,1}^\dagger, S_{5,1} S_{2,1}^\dagger S_{5,2}^\dagger, S_{2,4} S_{4,5} S_{2,5}^\dagger, S_{2,5} S_{2,4}^\dagger S_{4,5}^\dagger, S_{2,4} S_{4,5} S_{5,2}, S_{2,4}^\dagger S_{4,5}^\dagger S_{5,2}^\dagger \right\}, \right. \\ & \left. \left\{ S_{2,5}, S_{5,2}, S_{5,1} S_{2,1}, S_{2,4} S_{4,5}, (S_{2,5}^\dagger)^2 S_{5,2}, S_{2,5} S_{5,2}^2 \right\}, \left\{ S_{2,5}, S_{5,2}, S_{5,1} S_{2,1}, S_{2,4} S_{4,5}, (S_{2,5}^\dagger)^2 S_{5,2}, S_{2,5} S_{5,2}^2 \right\}, \right. \\ & \left. \left\{ 1, S_{2,5} S_{5,2}, S_{2,5}^\dagger S_{5,2}^\dagger, S_{2,5} S_{5,1} S_{2,1}^\dagger, S_{2,1} S_{2,5}^\dagger S_{5,1}^\dagger, S_{2,1} S_{5,2} S_{5,1}^\dagger, S_{5,1} S_{2,1}^\dagger S_{5,2}^\dagger, S_{2,4} S_{4,5} S_{2,5}^\dagger, S_{2,5} S_{2,4}^\dagger S_{4,5}^\dagger, S_{2,4} S_{4,5} S_{5,2}, S_{2,4}^\dagger S_{4,5}^\dagger S_{5,2}^\dagger \right\} \right\} \end{aligned}$$

$$\begin{aligned} 10^p 10^{\bar{q}\dagger}: K^{(10)} = & \left\{ \left\{ 1, S_{2,5} S_{5,2}, S_{2,5}^\dagger S_{5,2}^\dagger, S_{2,5} S_{5,1} S_{2,1}^\dagger, S_{2,1} S_{2,5}^\dagger S_{5,1}^\dagger, S_{2,1} S_{5,2} S_{5,1}^\dagger, S_{5,1} S_{2,1}^\dagger S_{5,2}^\dagger, S_{2,4} S_{4,5} S_{2,5}^\dagger, S_{2,5} S_{2,4}^\dagger S_{4,5}^\dagger, S_{2,4} S_{4,5} S_{5,2}, S_{2,4}^\dagger S_{4,5}^\dagger S_{5,2}^\dagger \right\}, \right. \\ & \left. \left\{ 1, S_{2,5} S_{5,2}, S_{2,5}^\dagger S_{5,2}^\dagger, S_{2,5} S_{5,1} S_{2,1}^\dagger, S_{2,1} S_{2,5}^\dagger S_{5,1}^\dagger, S_{2,1} S_{5,2} S_{5,1}^\dagger, S_{5,1} S_{2,1}^\dagger S_{5,2}^\dagger, S_{2,4} S_{4,5} S_{2,5}^\dagger, S_{2,5} S_{2,4}^\dagger S_{4,5}^\dagger, S_{2,4} S_{4,5} S_{5,2}, S_{2,4}^\dagger S_{4,5}^\dagger S_{5,2}^\dagger \right\}, \right. \\ & \left. \left\{ S_{2,5}, S_{5,2}, S_{5,1} S_{2,1}, S_{2,4} S_{4,5}, (S_{2,5}^\dagger)^2 S_{5,2}, S_{2,5} S_{5,2}^2 \right\}, \left\{ S_{2,5}, S_{5,2}, S_{5,1} S_{2,1}, S_{2,4} S_{4,5}, (S_{2,5}^\dagger)^2 S_{5,2}, S_{2,5} S_{5,2}^2 \right\}, \right. \\ & \left. \left\{ S_{2,5}, S_{5,2}, S_{2,1} S_{5,1}, S_{2,4} S_{4,5}, S_{2,5}^2 S_{5,2}, S_{2,5}^\dagger (S_{5,2}^\dagger)^2 \right\}, \left\{ 1, S_{2,5} S_{5,2}, S_{2,5}^\dagger S_{5,2}^\dagger, S_{2,5} S_{5,1} S_{2,1}^\dagger, S_{2,1} S_{2,5}^\dagger S_{5,1}^\dagger, S_{2,1} S_{5,2} S_{5,1}^\dagger, S_{5,1} S_{2,1}^\dagger S_{5,2}^\dagger, S_{2,4} S_{4,5} S_{2,5}^\dagger, S_{2,5} S_{2,4}^\dagger S_{4,5}^\dagger, S_{2,4} S_{4,5} S_{5,2}, S_{2,4}^\dagger S_{4,5}^\dagger S_{5,2}^\dagger \right\} \right\} \end{aligned}$$

$$L^p H^\dagger: \hat{\rho} = \begin{pmatrix} \{S_{2,4} S_{2,5}, S_{2,4} S_{5,2}, S_{2,4} S_{5,1} S_{2,1}^\dagger, S_{2,5} S_{5,2} S_{4,5}^\dagger\} \\ \{S_{2,4} S_{2,5}, S_{2,4} S_{5,2}, S_{2,4} S_{5,1} S_{2,1}^\dagger, S_{2,5} S_{5,2} S_{4,5}^\dagger\} \\ \{S_{2,4}, S_{2,5} S_{4,5}, S_{2,1} S_{4,5} S_{5,1}^\dagger, S_{2,4} S_{2,5} S_{5,2}, S_{2,4} S_{2,5} S_{5,2}^\dagger\} \end{pmatrix}$$

Overview of phenomenological properties:

standard models	no massless $U(1)$	1 Higgs pair	2 Higgs pairs	3 Higgs pairs	$\text{rk}(Y^{(u)}) > 0$	no proton decay, $\lambda = \lambda' = 0$	1 Higgs, $\text{rk}(Y^{(u)}) > 0$, $\lambda = \lambda' = 0$, $U(1)$ s massive
407	237	262	77	63	45	198	13

Moduli stabilization

- dilaton \mathcal{S}
 - $h^{1,1}(X)$ Kahler moduli T^i
 - $h^{2,1}(X)$ complex structure moduli Z^a
 - $h^1(X, V \otimes V^*)$ bundle moduli \mathcal{S}^α
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Analogue of IIB stabilization cannot be realized easily:

- Small W_0 difficult to achieve since there are fewer fluxes.
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Usually $c_2(V) < c_2(TX)$, so there is room for a hidden bundle \tilde{V}

Stabilize moduli by hidden E_8 flux

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Probably need to combine D-terms, matter superpotential for S^α and non-perturbative effects to stabilize all moduli and break Susy.

(B. Dundee, S. Raby, A. Westphal 2010;
S. Parameswaran, I. Zavala, S. Ramos-Sanchez, 2011)

Conclusion and outlook

- We can “mass-produce” heterotic CY standard models. Many more than the current 2000 models will be found.
- The additional $U(1)$ symmetries of these models constrain the 4d theory significantly.
- Calculation of superpotential terms feasible but matter field Kahler metric remains a problem.
- Finding discrete symmetries... (M. Klaput, C. Mishra, AL, in progress)
- More detailed model building (proton stability, fermion masses) is now possible.
- Moduli stabilization via hidden bundle looks promising. Remaining moduli have to be stabilized non-perturbatively, with Susy broken \rightarrow soft breaking parameters.

Thanks