

The Experimental Status of ϑ_{13} from the Point of View of the Reactor Neutrino Disappearance Experiments

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Daya Bay collaboration

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References



- Daya Bay
 - F.P. An et al., Daya Bay Coll., “A side-by-side comparison of Daya Bay anti-neutrino detectors”, *Nucl. Inst. and Meth. A* 685 (2012), pp. 78-97
 - F.P. An et al., Daya Bay Coll., “Observation of electron anti-neutrino disappearance at Daya Bay”, *Phys. Rev. Lett.* 108, 171803 (2012)
 - D. Dwyer, Talk at Neutrino 2012, Kyoto, June 4, 2012
- Double Chooz
 - Y. Abe et al., Double Chooz Collaboration , “Indication for the disappearance of reactor electron antineutrinos in the Double Chooz experiment.”, arXiv:1112.6353 [hep-ex], *Phys.Rev.Lett.* 108 (2012) 131801
 - M. Ishitsuka, Talk at Neutrino 2012, Kyoto, June 4, 2012
- RENO
 - J.K. Ahn et al., Reno Collaboration, “Observation of Reactor Electron Anti-Neutrino Disappearance in the RENO Experiment”, *Phys.Rev.Lett.* 108 (2012) 191802
 - Soo-Bong Kim, Talk at Neutrino 2012, Kyoto, June 4, 2012

Neutrino flavor eigenstates $|\nu_f\rangle$, $f = e/\mu/\tau$ produced in weak
 Interactions are different from mass eigenstates $|\nu_i\rangle$, $i = 1/2/3$
 → non-diagonal Unitary
 mixing matrix:

$$U_{fi} \equiv \langle \nu_f | \nu_i \rangle \Rightarrow |\nu_f\rangle = \sum_{i=1}^3 U_{fi}^* |\nu_i\rangle$$

Canonical representation of Pontecorvo-Magi-Nakagawa-Sakata mixing matrix
 is done by ordered product of 12, 13 and 23 rotations, one CP phase δ connected
 to the smallest mixing angle θ_{13} and two Majorana phases $\alpha_{1,2}$.

$$\begin{aligned} \theta_{23} &\cong 45^\circ \\ \begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} &= \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos(\theta_{23}) & \sin(\theta_{23}) \\ 0 & -\sin(\theta_{23}) & \cos(\theta_{23}) \end{pmatrix} \cdot \begin{pmatrix} \cos(\theta_{13}) & 0 & \sin(\theta_{13}) \cdot e^{-i\delta} \\ 0 & 1 & 0 \\ -\sin(\theta_{13}) \cdot e^{i\delta} & 0 & \cos(\theta_{13}) \end{pmatrix}. \end{aligned}$$

$$\begin{pmatrix} \cos(\theta_{12}) & \sin(\theta_{12}) & 0 \\ -\sin(\theta_{12}) & \cos(\theta_{12}) & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} e^{i\alpha_1/2} & 0 & 0 \\ 0 & e^{i\alpha_2/2} & 0 \\ 0 & 0 & 1 \end{pmatrix} \cdot \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

$$\theta_{12} \cong 34^\circ$$

Majorana phases
 α are
 irrelevant for
 oscillations.

If masses of mass eigenstates are different
then probabilities oscillate:

$$P_{\nu_f \rightarrow \nu_g}(L) = \delta_{fg} - 4 \sum_{i < j} \Re(U_{gi} U_{fi}^* U_{fj} U_{gj}^*) \sin^2 \left(\frac{m_j^2 - m_i^2}{4\hbar c} \frac{L}{E} \right) + 2 \sum_{i < j} \Im(U_{gi} U_{fi}^* U_{fj} U_{gj}^*) \sin \left(\frac{m_j^2 - m_i^2}{2\hbar c} \frac{L}{E} \right)$$

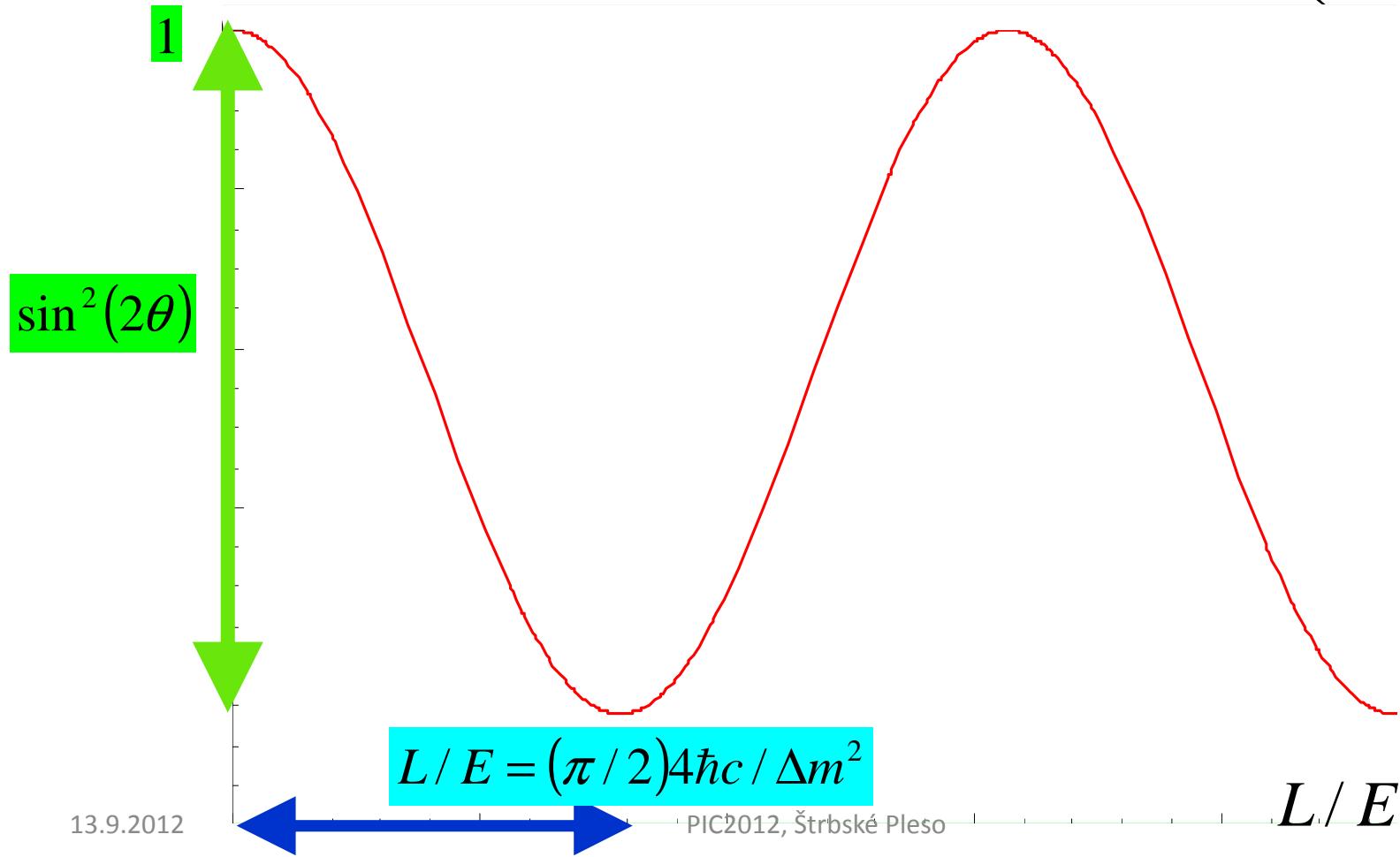
Last term is CP and T odd and it is ≠0 only if:

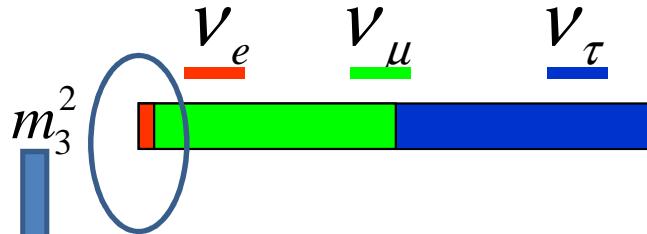
- $f \neq g$
- all three mixing angles $\neq 0$ and
- imaginary part of $\exp(i\delta) = \text{Sin}(\delta) \neq 0$

**Amplitude of oscillations = $\sin^2(2\theta)$,
oscillation length is inversely proportional to Δm^2**

$$P_{\nu \rightarrow \nu}(L/E)$$

$$P_{\nu \rightarrow \nu}(L/E) = 1 - \sin^2(2\theta) \sin^2\left(\frac{\Delta m^2}{4\hbar c} \frac{L}{E}\right)$$

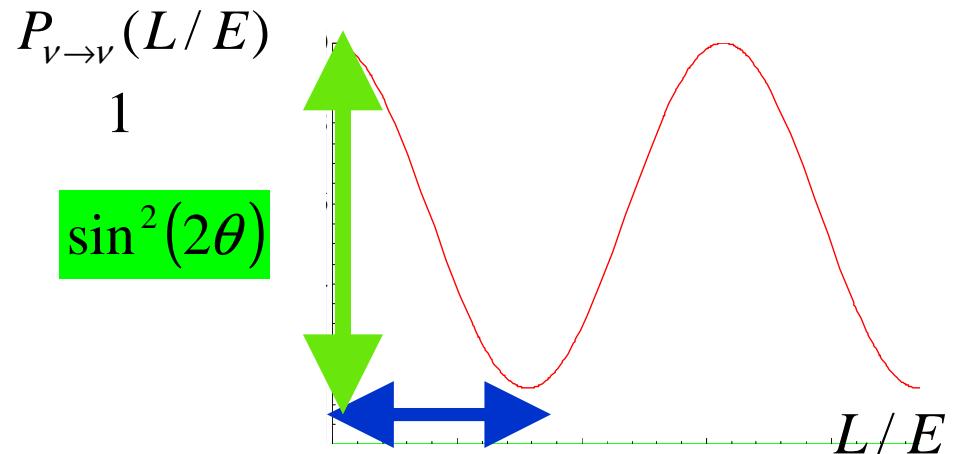
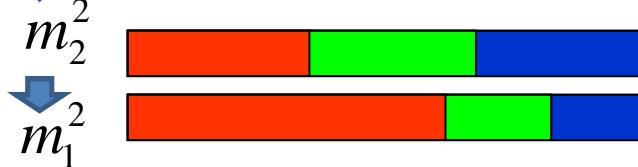




**Two Δm^2 differs app. by a factor of 30
→ two very different oscillation lengths**

$$|m_3^2 - m_1^2| \approx 2.3 \times 10^{-3} \text{ eV}^2$$

$$m_2^2 - m_1^2 \approx 7.6 \times 10^{-5} \text{ eV}^2$$



$$\begin{aligned} (L/E)_{1stMINIMUM} &= (\pi/2) 4\hbar c / \Delta m^2 \\ &\approx 0.5 \text{ km/MeV} = 500 \text{ km/GeV} \\ &\approx 15 \text{ km/MeV} = 15000 \text{ km/GeV} \end{aligned}$$

	0.5 km/MeV 500 km/GeV	15 km/MeV 15000 km/GeV
$\nu_e \rightarrow \nu_e$	Daya Bay Double Chooz/RENO	reactor KAMLAND Sun ν_e (+matter effect)
$\nu_\mu \rightarrow \nu_\mu$	atm. acc.	

Disappearance probability

$$P_{\nu_f \rightarrow \nu_f}(x) = 1 - \sum_{i < j} 4|U_{fi}|^2 |U_{fj}|^2 \sin^2 \left(1.267 \Delta m_{ij}^2 [eV^2] \frac{x[m]}{E[MeV]} \right)$$

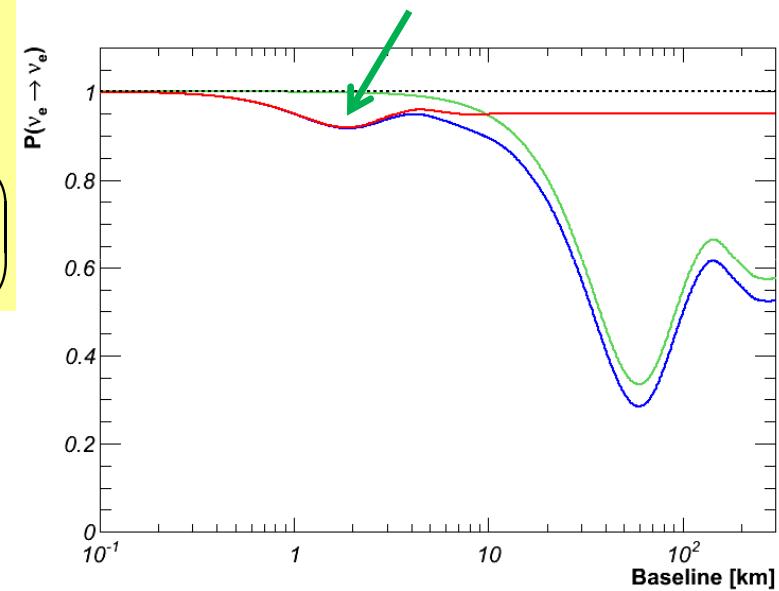
Disappearance probability for electron (anti)neutrinos:

$$P_{\nu_e \rightarrow \nu_e}(x) \xrightarrow{\Delta m_{31}^2 \approx \Delta m_{32}^2}$$

$$1 - \sin^2(2\theta_{13}) \sin^2 \left(1.267 \Delta m_{31}^2 [eV^2] \frac{x[m]}{E[MeV]} \right)$$

$$- \cos^4(\theta_{13}) \sin^2(2\theta_{12}) \sin^2 \left(1.267 \Delta m_{21}^2 [eV^2] \frac{x[m]}{E[MeV]} \right)$$

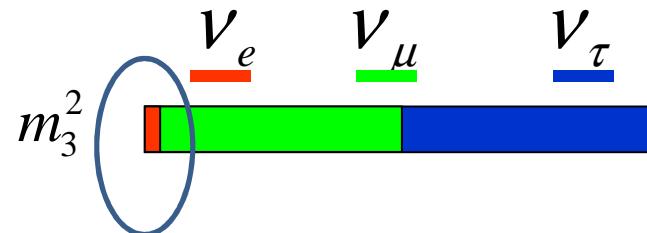
For E=4 MeV the first minimum is at ~2 km



$$P_{\nu_e \rightarrow \nu_e}(x) \approx 1 - \sin^2(2\theta_{13}) \sin^2 \left(1.267 \Delta m_{31}^2 [eV^2] \frac{x[m]}{E[MeV]} \right)$$

Mixing angle θ_{13}

$\sin^2(\theta_{13})$ is the fraction of electron neutrino in mass eigenstate m_3



Two ways to measure θ_{13}

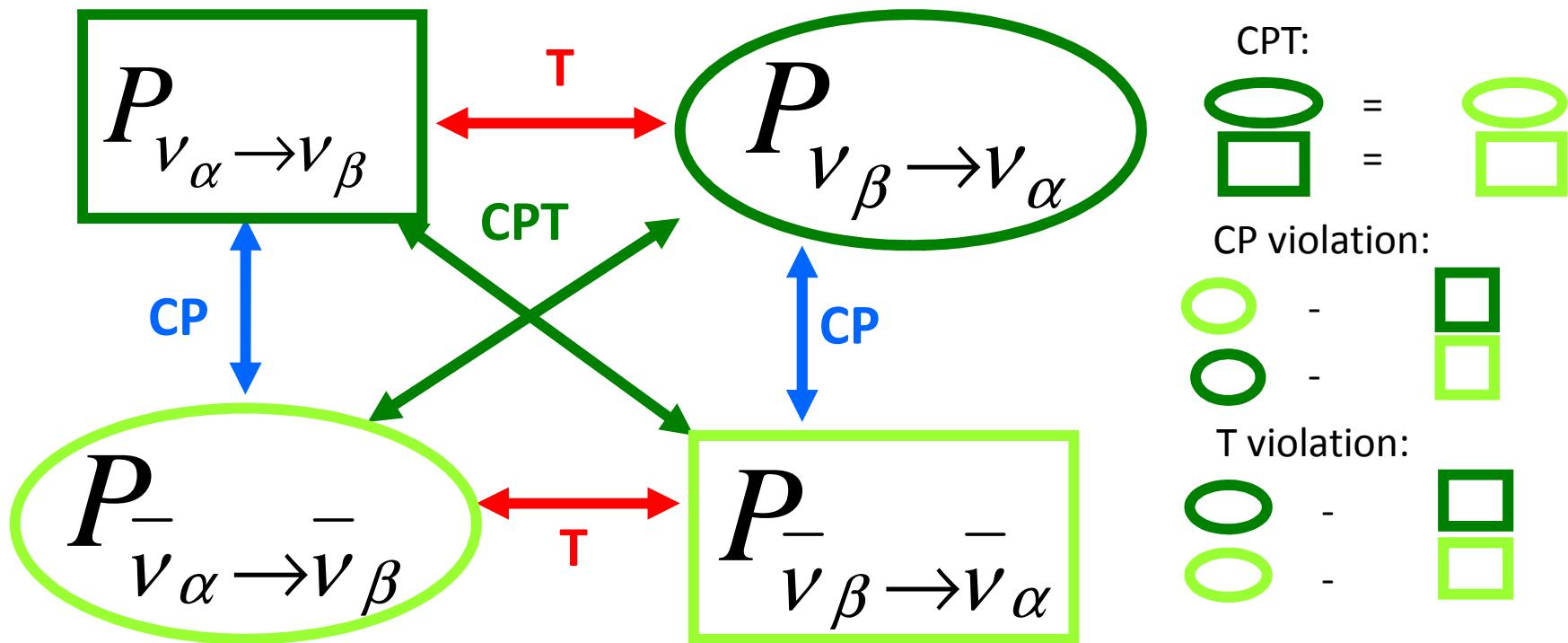
- To measure electron (anti)neutrino disappearance
- To measure electron (anti)neutrino appearance in muon (anti)neutrino beam

measurements has to be done at small values of $L/E \sim 0.5 \text{ km/MeV} = 500 \text{ km/GeV}$:

$$m_2^2$$
$$m_1^2$$

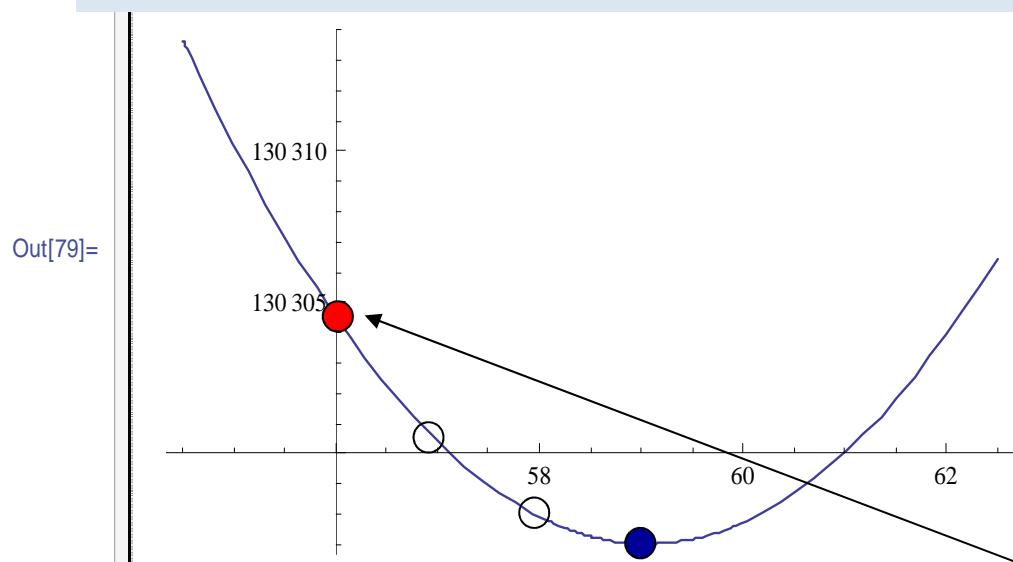


If $\theta_{13} \neq 0$ then CP and T violation in lepton sector could be investigated with neutrino oscillations

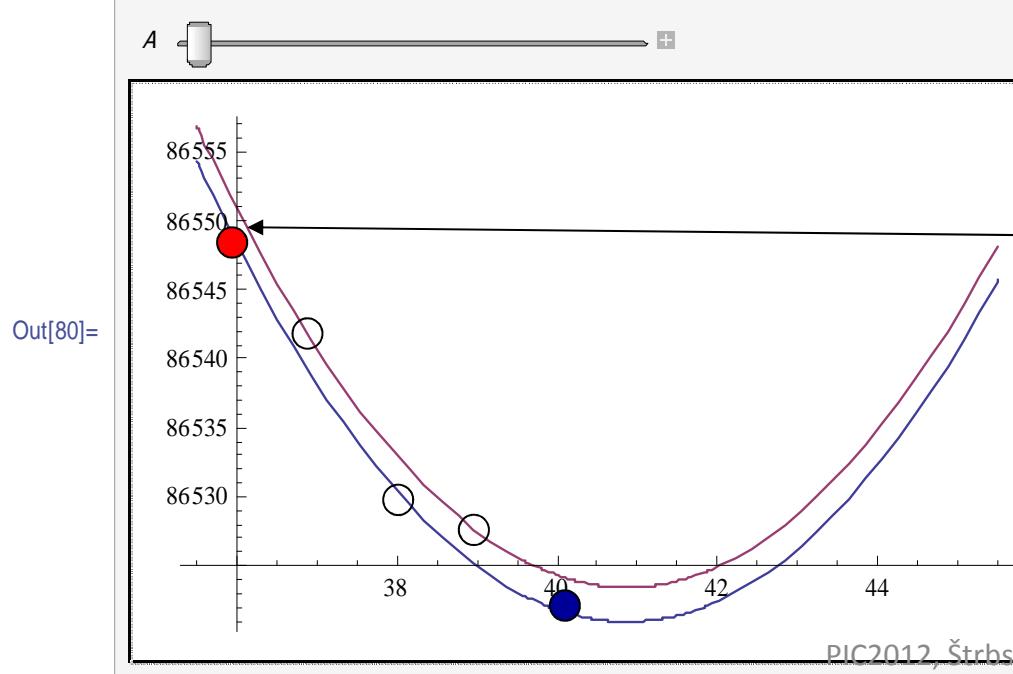


$$\begin{aligned}
 P_{\nu_\mu \rightarrow \nu_e} - P_{\bar{\nu}_\mu \rightarrow \bar{\nu}_e} = & -2 \sin(\delta) \cos(\theta_{13}) \sin(2\theta_{13}) \sin(2\theta_{12}) \sin(2\theta_{23}) \\
 & \times \sin\left(\frac{\Delta m_{12}^2}{4\hbar c} \frac{L}{E}\right) \sin\left(\frac{\Delta m_{31}^2}{4\hbar c} \frac{L}{E}\right) \sin\left(\frac{\Delta m_{23}^2}{4\hbar c} \frac{L}{E}\right)
 \end{aligned}$$

Nuclear reactors are powerful sources of electron antineutrinos

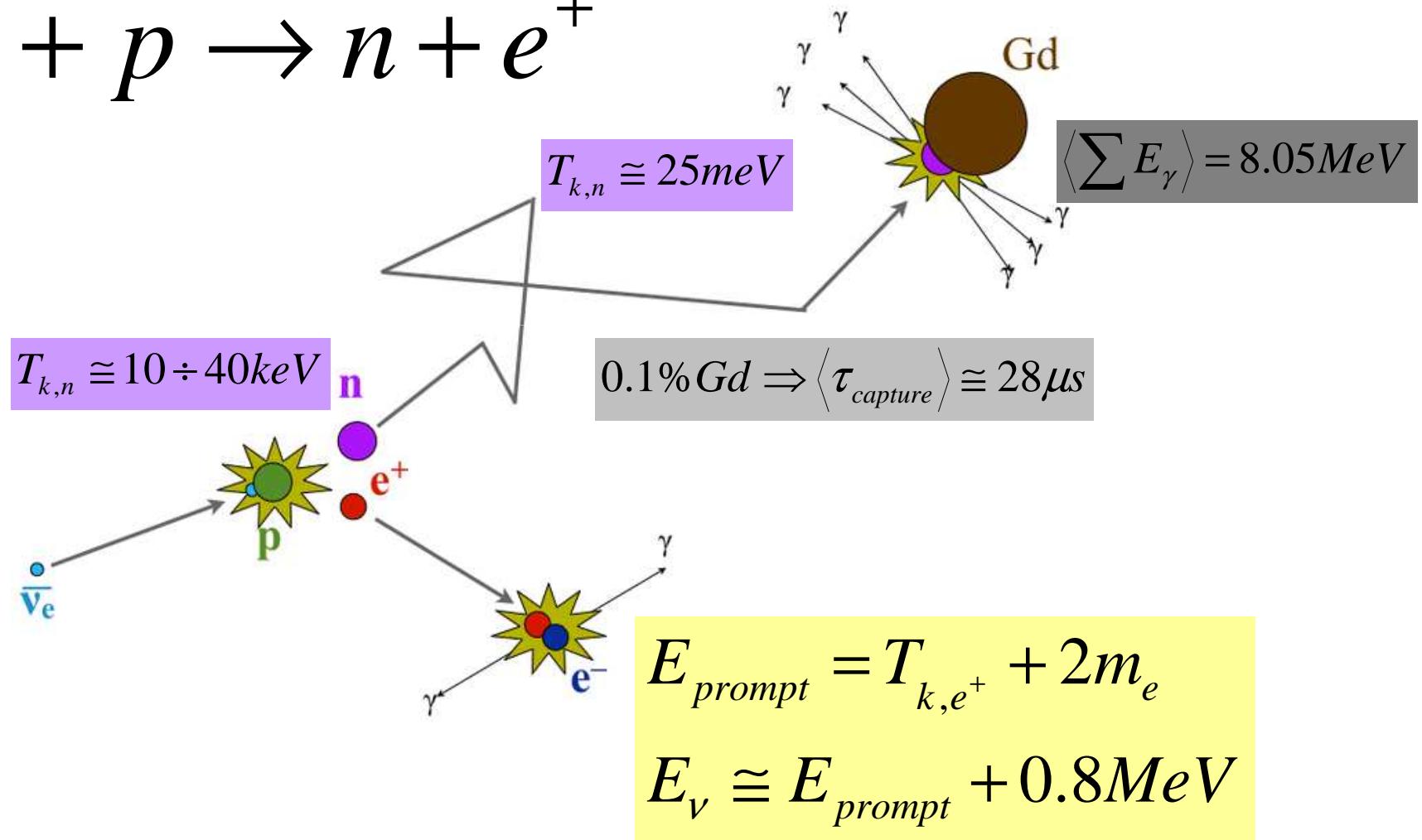
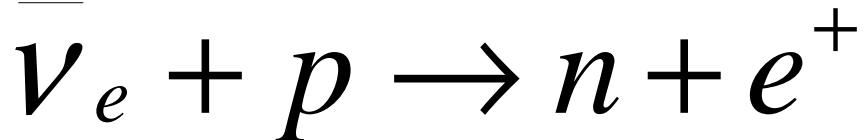


Fission products are neutron rich.
The valley of stability is reached by series
of beta- decays.
In average app. **6 electron antineutrinos**
are produced per fission.

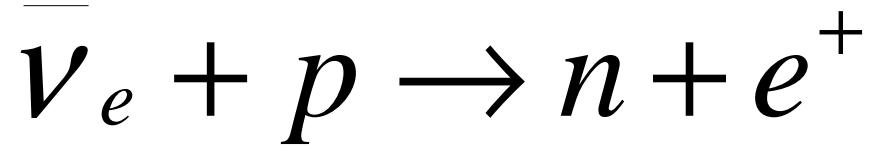


$$2 \cdot 10^{20} \bar{\nu}_e / s / GW_{th}$$

Detection of antineutrinos via Inverse Beta Decay (IBD). Coincidence of prompt signal from positron and delayed signal of neutron capture on Gd.



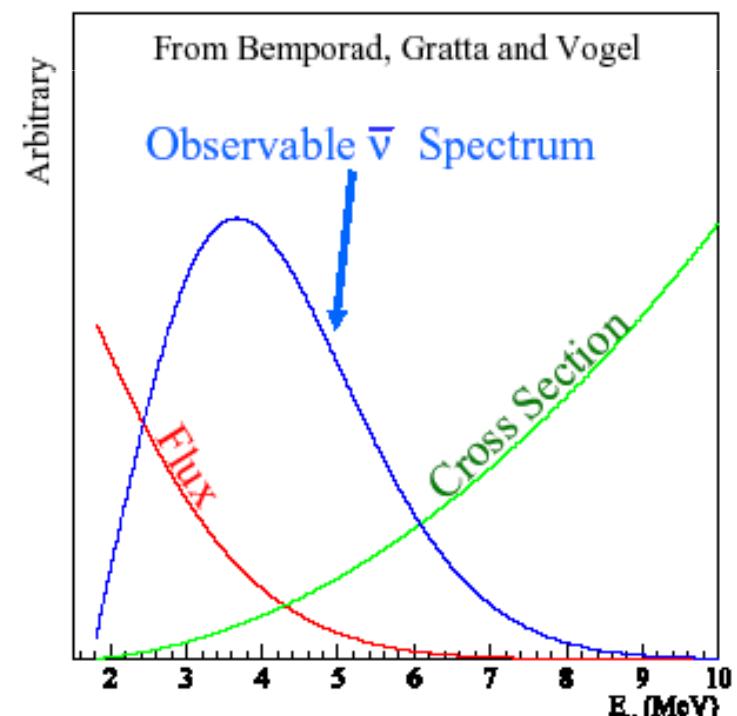
Detection of antineutrinos: Inverse Beta Decay (IBD)



$$E_{\nu,THR} = \frac{(m_n + m_e)^2 - m_p^2}{2m_p} = \frac{m_n + m_p + m_e}{2m_p} (m_n - m_p + m_e) \\ = 1.00096(m_n - m_p + m_e) = 1.83 \text{ MeV}$$

Only antineutrinos with energies larger than 1.8 MeV interact.

Detected energy spectrum is the product of reactor neutrino spectrum and IBD cross section and it reaches the maximum around 4 MeV → the first oscillation minimum is at 0.5 km/MeV → 2 km for 4 MeV



Neutron capture on Gadolinium

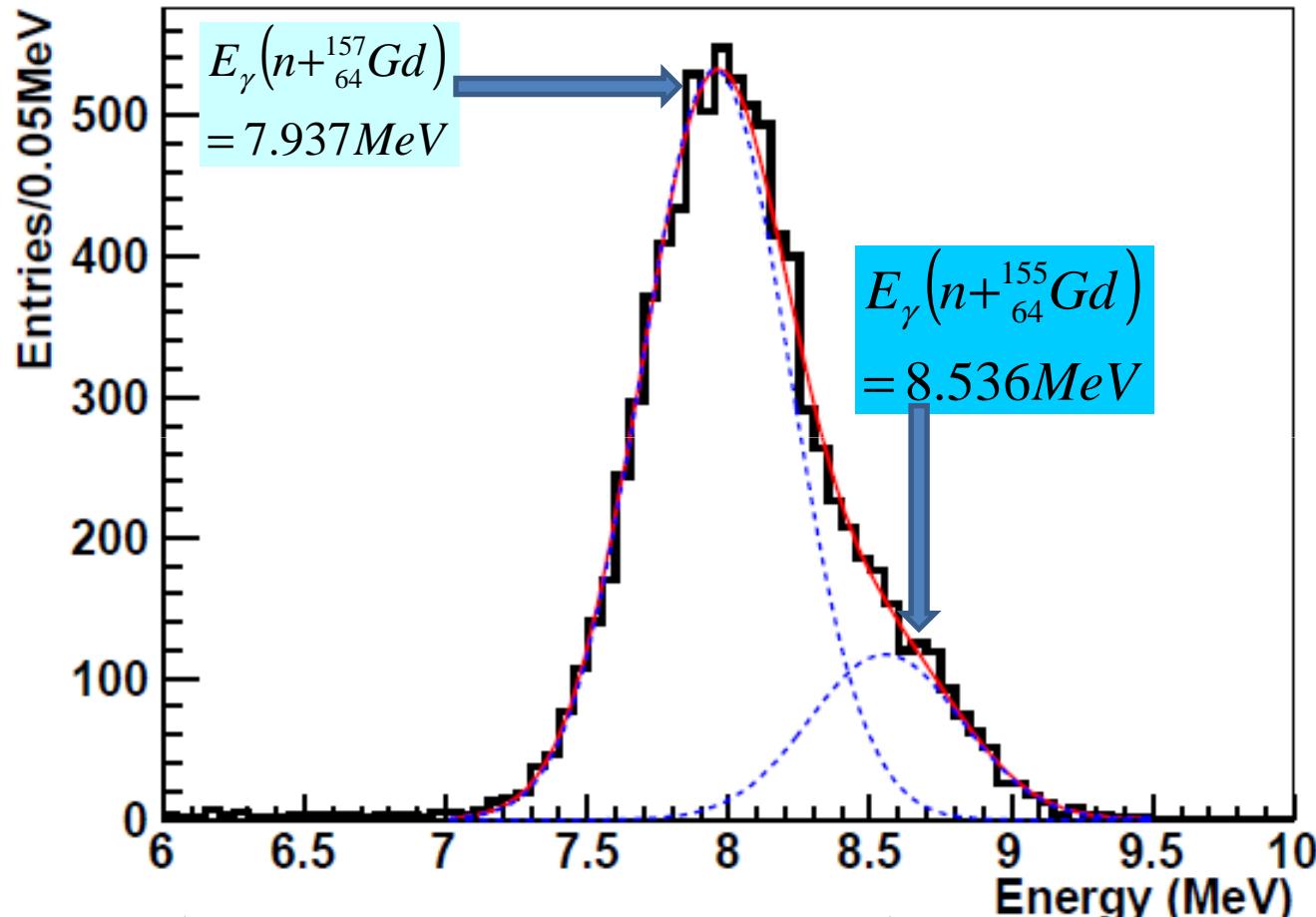
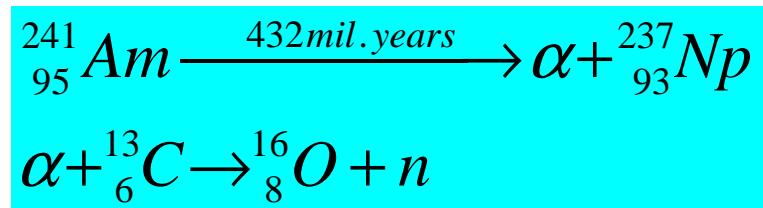


${}_{64}^A Gd$	$\sigma_{n + {}_{64}^A Gd \rightarrow {}_{64}^{A+1} Gd^*} [b]$	Abund. [%]	B[MeV/A]
$A = 152 \left(\xrightarrow{\alpha} {}_{62}^{148} Sm \right)$	735	0.20	8.233399
$A = 154$	85	2.18	8.224794
$A = 155$	60900	14.80	8.213248
$A = 156$	1.8	20.47	8.215320
$A = 157$	254000	15.65	8.203501
$A = 158$	2.2	24.84	8.201817
$A = 160 \left(\xrightarrow{2\beta} {}_{66}^{160} Dy \right)$	1.4	21.86	8.183010

$$\sum E_\gamma \left(n + {}_{64}^{155} Gd \right) = 8.536 MeV \quad \langle E_\gamma \rangle = 8.048 MeV$$

$$\sum E_\gamma \left(n + {}_{64}^{157} Gd \right) = 7.937 MeV$$

^{241}Am - ^{13}C Daya Bay



$$\langle E_\gamma \rangle = \frac{14.80\% \cdot 69000b \cdot 8.536\text{MeV} + 15.65\% \cdot 254000b \cdot 7.937\text{MeV}}{14.80\% \cdot 69000b + 15.65\% \cdot 254000b} = 8.048\text{MeV}$$

Three running experiments

Experiment	Power (GW)	Baseline(m) Near/Far	Detector(t) Near/Far	Overburden (MWE) Near/Far	Design sensitivity (90%CL)
Double Chooz	8.5	400/1050	~ /8.2 (8.2/8.2)	120/300	~ 0.03
Daya Bay	17.4	470, 576/1650	40, 20/60 (40, 40/80)	250, 265/860	~ 0.008
RENO	16.5	409/1444	16/16	120/450	~ 0.02

The experiments are constructed following the concept of two identical near/far detectors proposed by:

L.Mikaelyan and V.V.Sinev [Phys.Atom.Nucl.63:1002-1006,2000;
Yad.Fiz.63N6:1077-1081,2000]

Daya Bay, China

6 reactors:
17.4 GW total
(thermal) power

A total of eight functionally identical and moveable detectors in three detector halls.

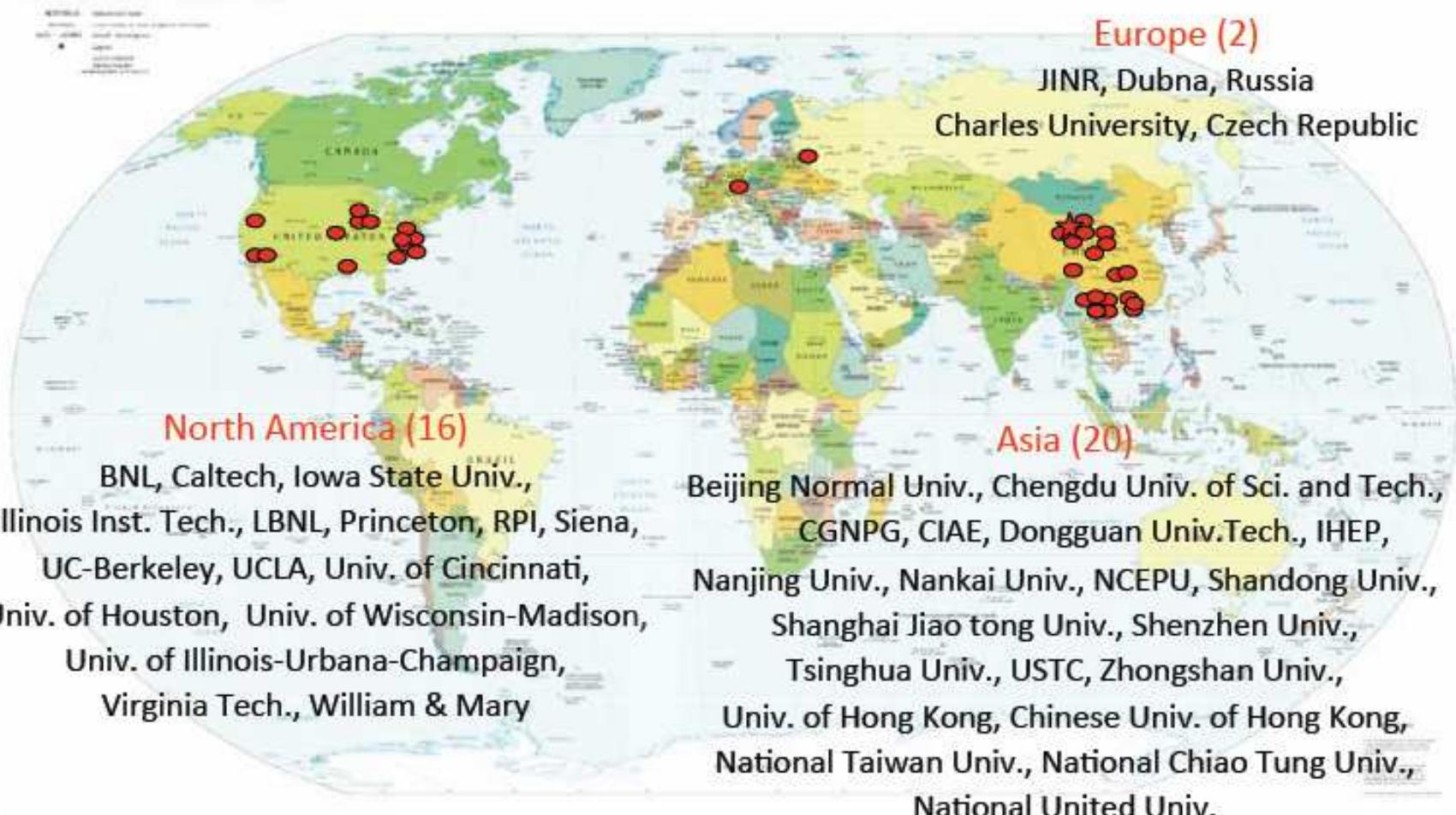
6 of the 8 detectors have been taking physics data since Dec. 2012

The remaining two detectors will be installed and commissioned this year.





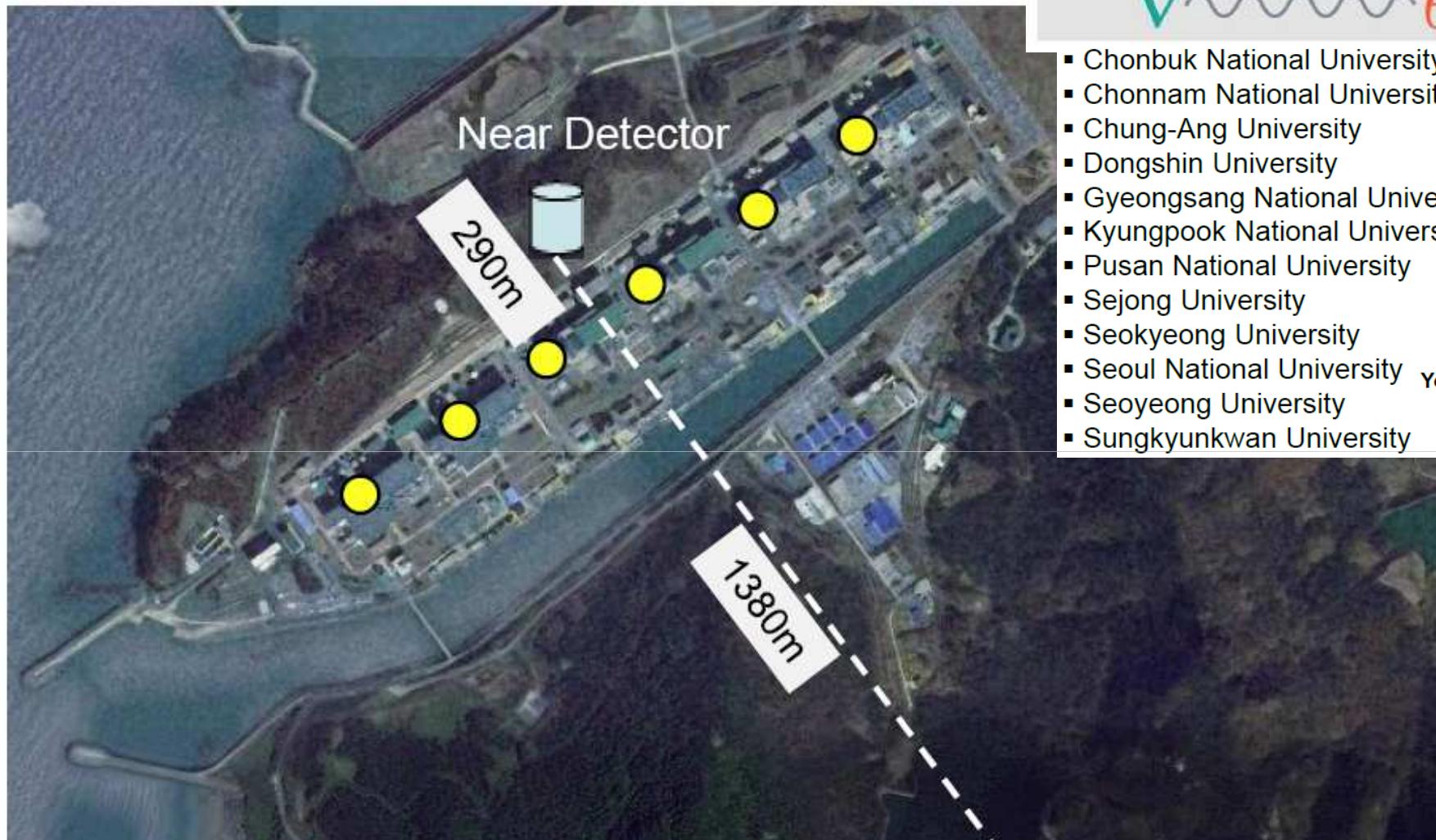
Political Map of the World, June 1999



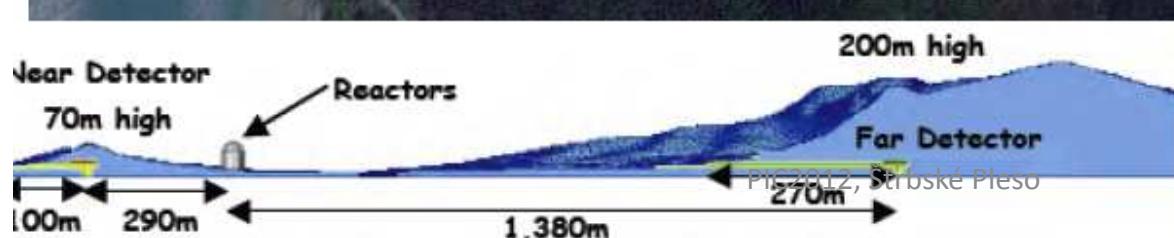
Double Chooz, France



RENO, South Korea



- Chonbuk National University
- Chonnam National University
- Chung-Ang University
- Dongshin University
- Gyeongsang National University
- Kyungpook National University
- Pusan National University
- Sejong University
- Seokyeong University
- Seoul National University YongGw
- Seoyeong University
- Sungkyunkwan University



Anti-neutrino detectors

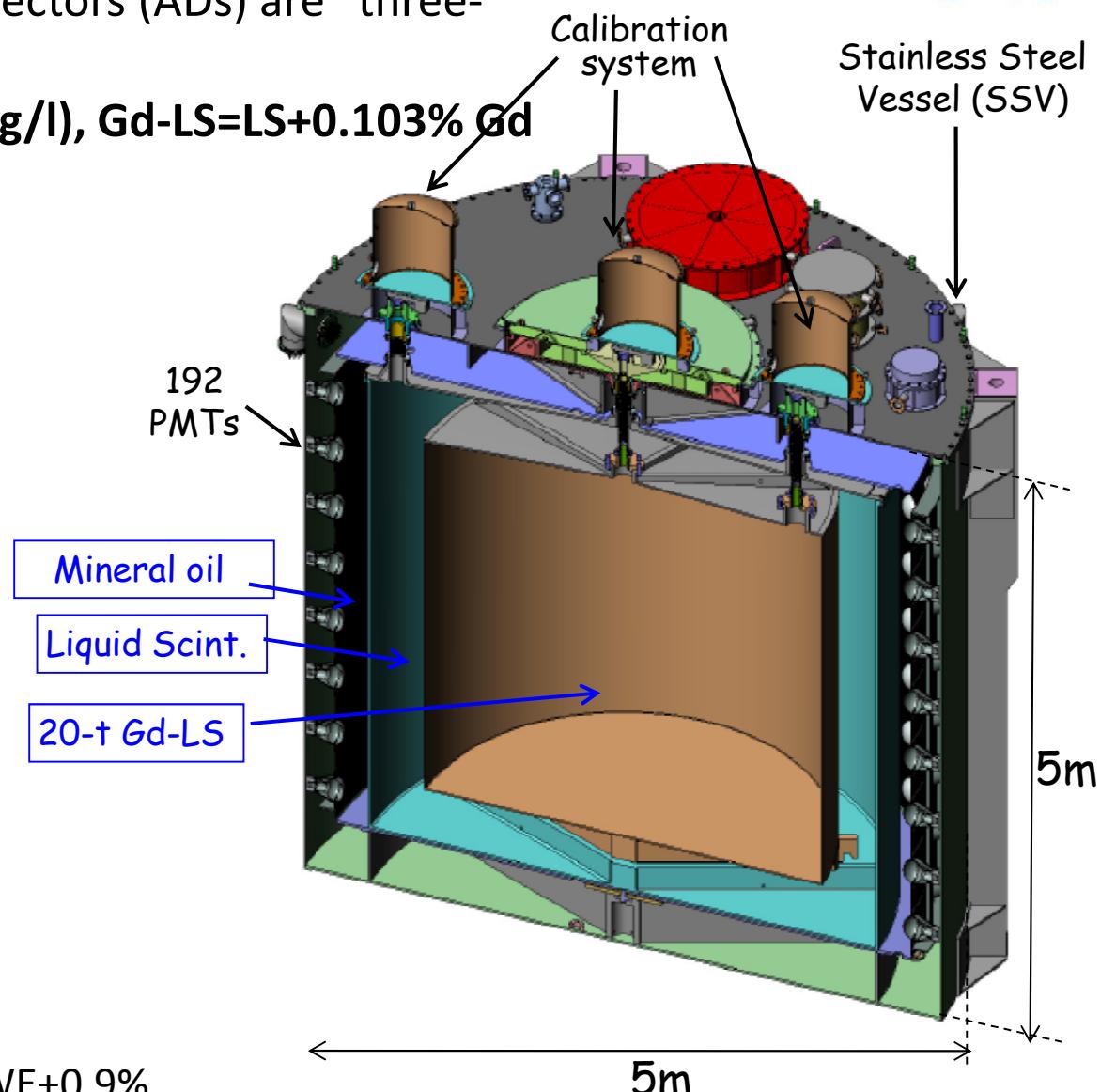


- The Daya Bay anti-neutrino detectors (ADs) are “three-zone” cylindrical modules.
- LS=LAB+PPO(3 g/l)+MSB(15 mg/l), Gd-LS=LS+0.103% Gd**

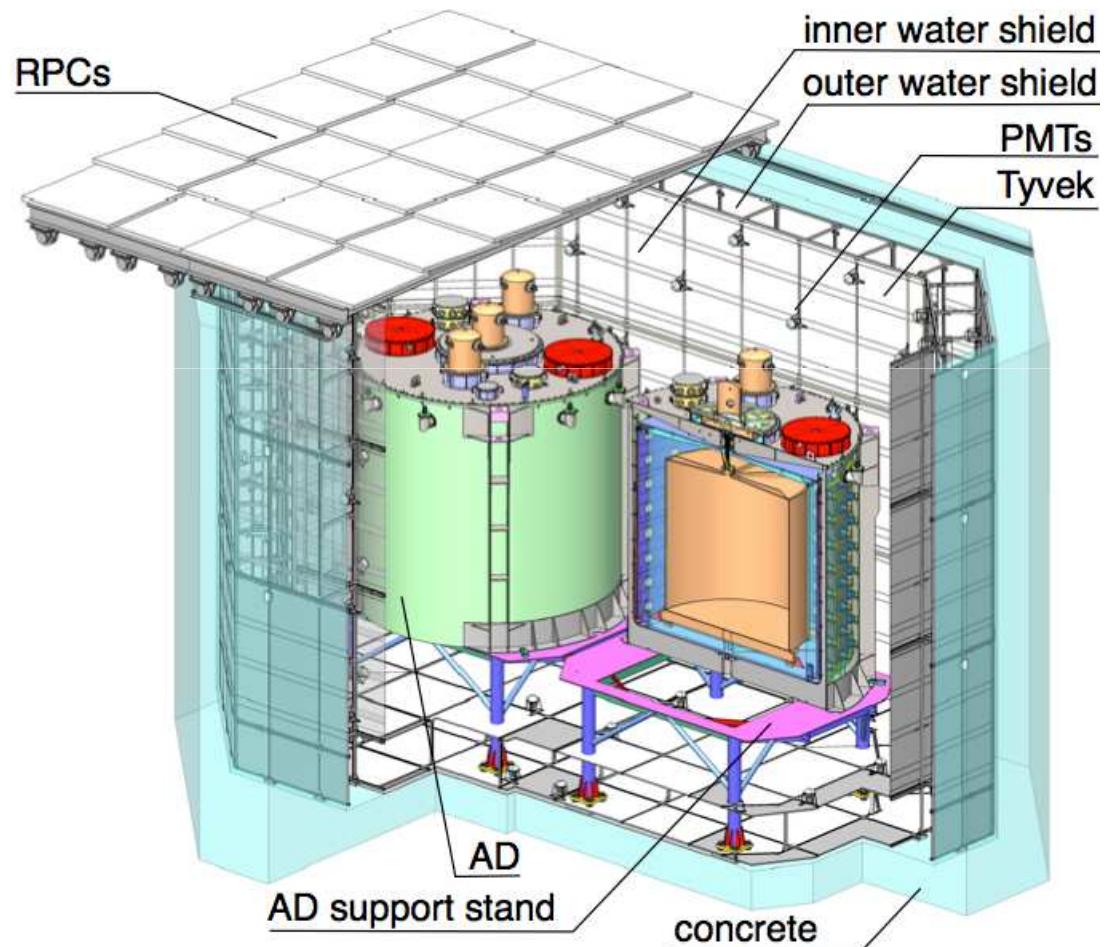
- Zones are separated by acrylic vessels:

Zone	Mass	Liquid	Purpose
Inner acrylic vessel	20 t	Gd-doped liquid scintillator	Anti-neutrino target
Outer acrylic vessel	20 t	Liquid scintillator	Gamma catcher (from target zone)
Stainless steel vessel	40 t	Mineral Oil	Radiation shielding

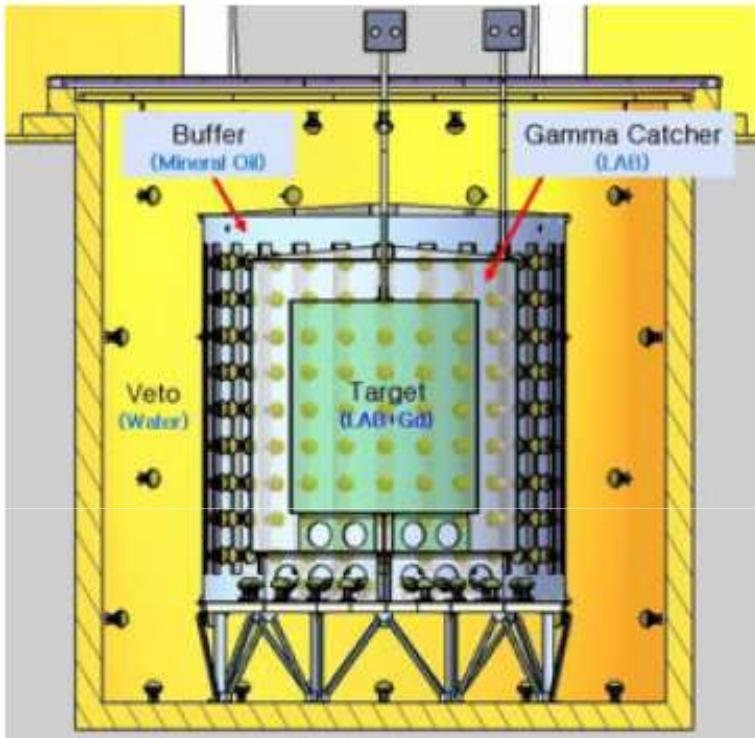
- Top and bottom reflectors are used to increase light yield
- Energy resolution: $\sigma_E/E = 7.5\%/\sqrt{E} + 0.9\%$



- Outer layer of water Čerenkov detector (on sides and bottom) is 1m thick, inner layer >1.5m. Water extends 2.5m above ADs
 - 288 8" PMTs in each near hall
 - 384 8" PMTs in Far Hall
- 4-layer RPC modules above pool
 - 54 modules in each near hall
 - 81 modules in Far Hall



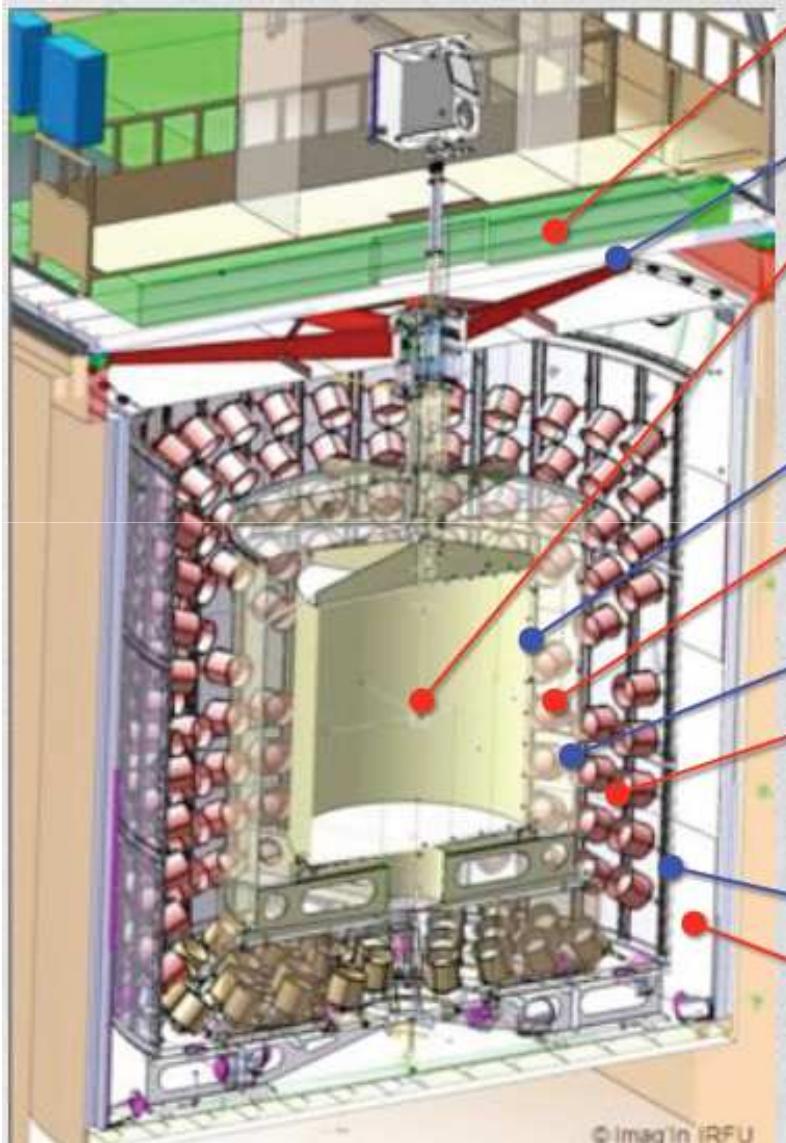
RENO Detector



- 354 ID +67 OD 10" PMTs
- Target : 16.5 ton Gd-LS, R=1.4m, H=3.2m
- Gamma Catcher : 30 ton LS, R=2.0m, H=4.4m
- Buffer : 65 ton mineral oil, R=2.7m, H=5.8m
- Veto : 350 ton water, R=4.2m, H=8.8m



Double Chooz detector



Outer Veto: Plastic scintillator strips
Identify cosmic μ

Steel shield (15cm thick)

v-target:
Gd loaded (1g/l) liquid scint. (10m^3)
Target of neutrino interaction
Neutrons captured on Gd

Acrylic vessel -----

γ -catcher: Liquid scintillator (22m^3)
Measure γ 's escaped from v-target

Acrylic vessel -----

Buffer:
Mineral oil (110m^3) & 390 10-inch PMT
Reduction of environmental γ 's

Steel tank -----

Inner Veto:
Liquid scintillator (90m^3) & 78 8-inch PMT
Identify cosmic μ & reduction neutrons

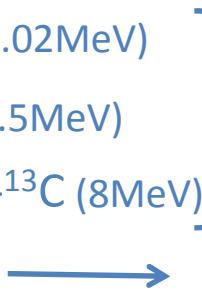
Daya Bay	Double Chooz	RENO
		Cs, 662 keV
Ge, 2x511 keV		Ge, 2x511 keV
Co60 2.5 MeV	Co60 2.5 MeV	Co60 2.5 MeV
	Cf252	Cf252
Am241-C13		
LED	LED	
		Laser

- Three sources + LED in each calibration unit, on a turn-table:

○ ^{68}Ge (1.02MeV)
 ○ ^{60}Co (2.5MeV)
 ○ $^{241}\text{Am}-^{13}\text{C}$ (8MeV)
 ○ LED

Energy calibration
 (linearity, detector
 response... etc)

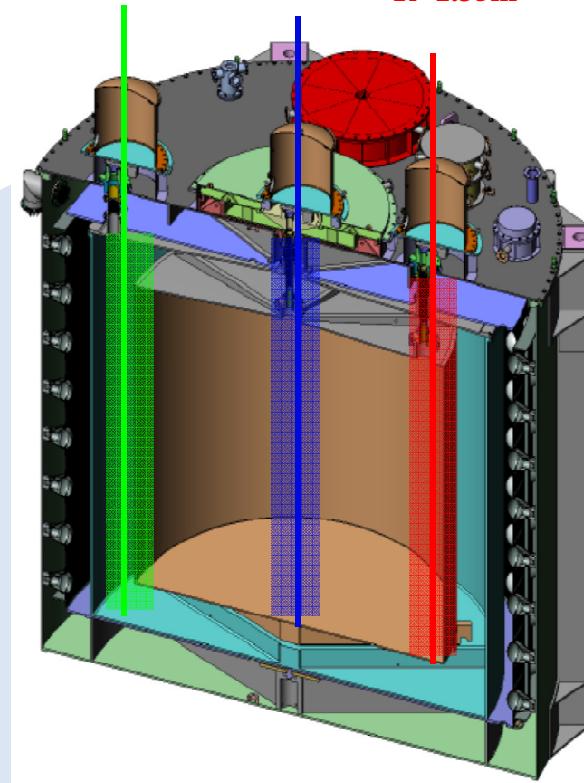
Timing, gain and
 relative QE



- Can also use spallation neutrons
 (uniformity, stability, calibration, ... etc).

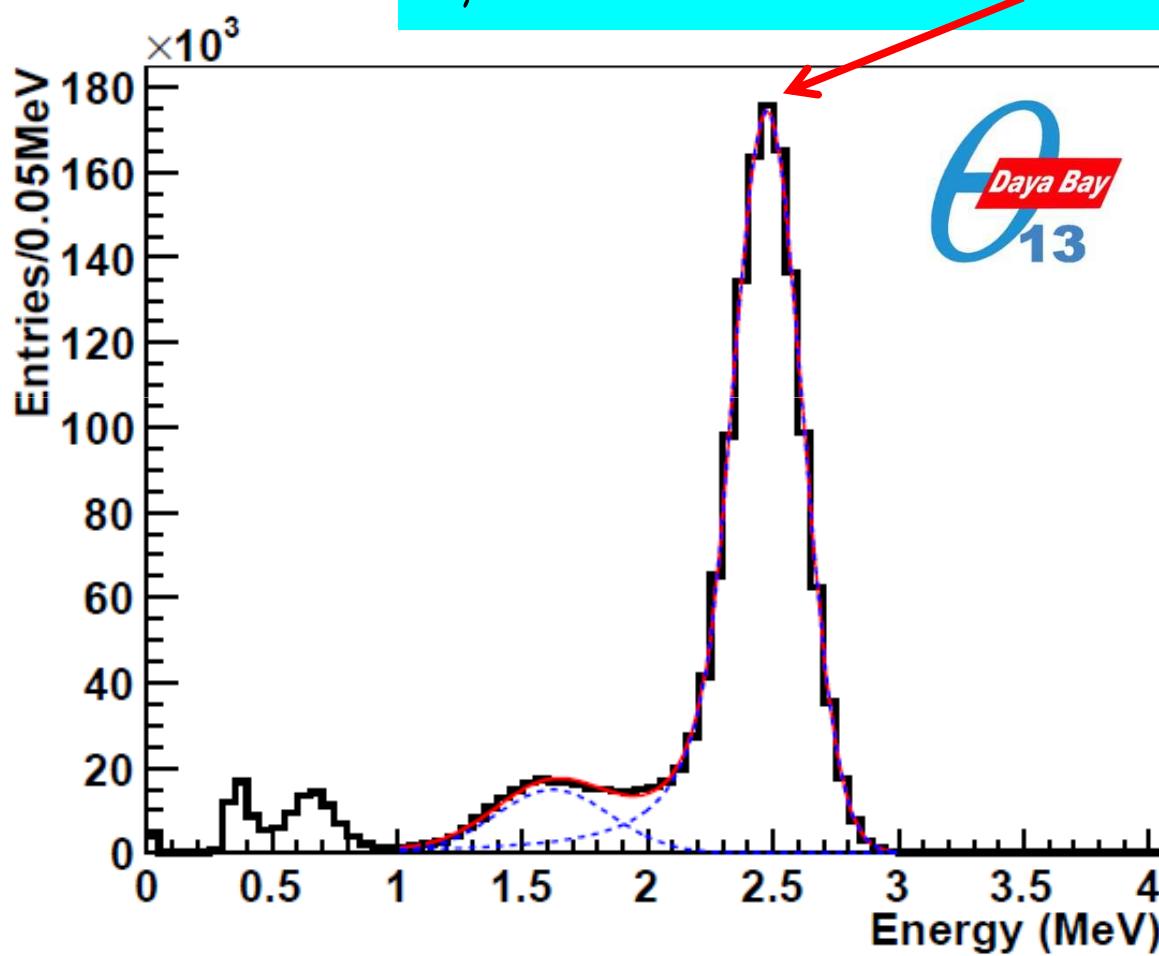
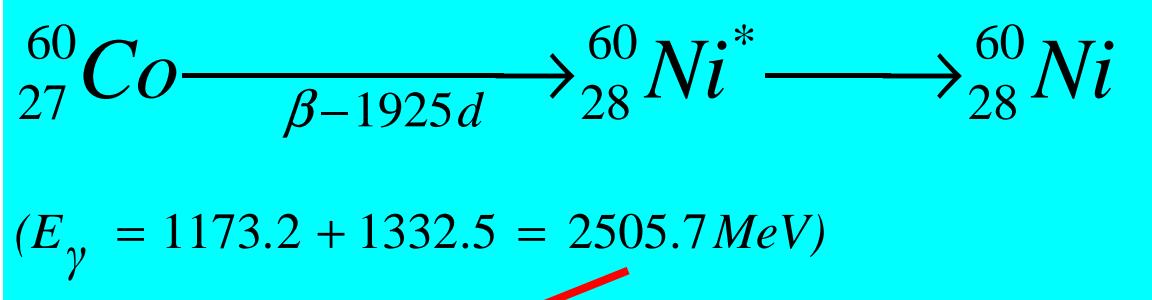
Detector calibration

R=1.775 m R=0 R=1.35m



Automated Calibration Units (Daya Bay)

Three calibration units per detector
 that deploy sources along z-axis



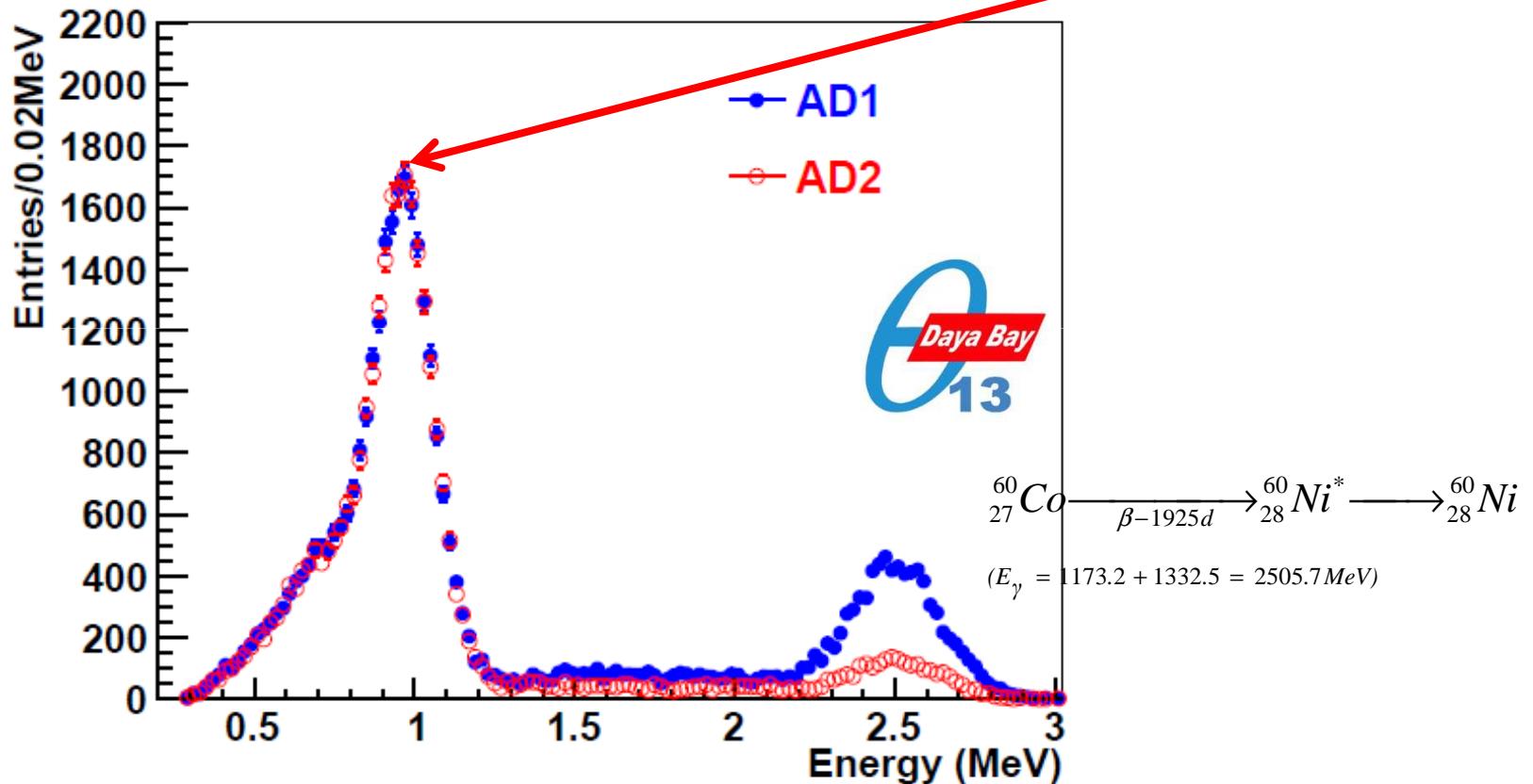
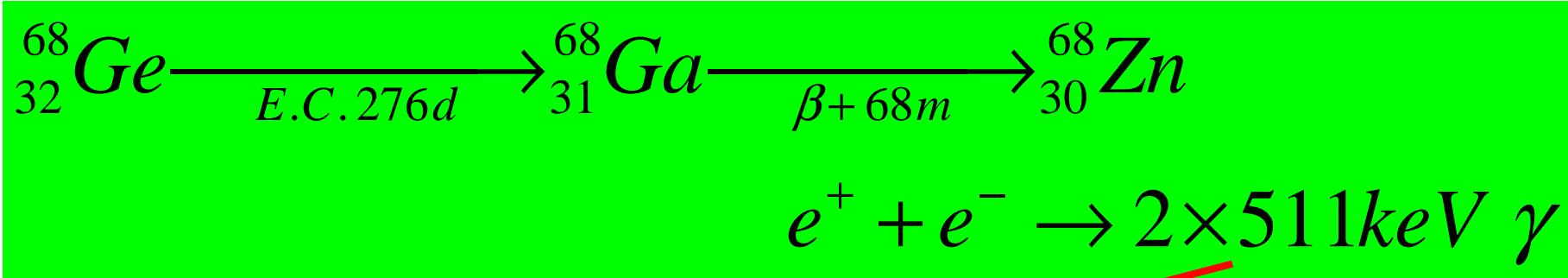
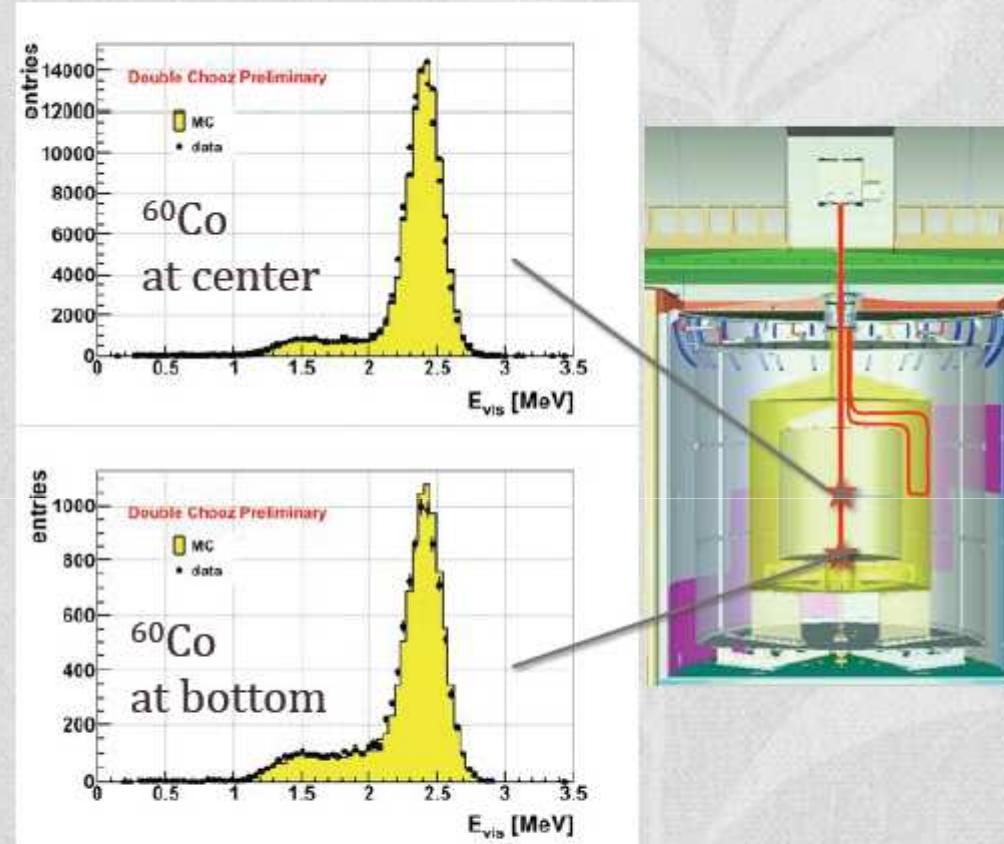


Figure 15: The energy spectrum of the ${}^{68}Ge$ source.
PIC2012, Štrbské Pleso

Calibration

Energy calibration

1. PMT and electronics gain non-linearity calibration
 - LED light injection system
2. Correction for position dependence & stability
 - Spallation neutron captures on H and Gd
3. Energy scale
 - Radioactive sources deployed into ν -target and γ -catcher



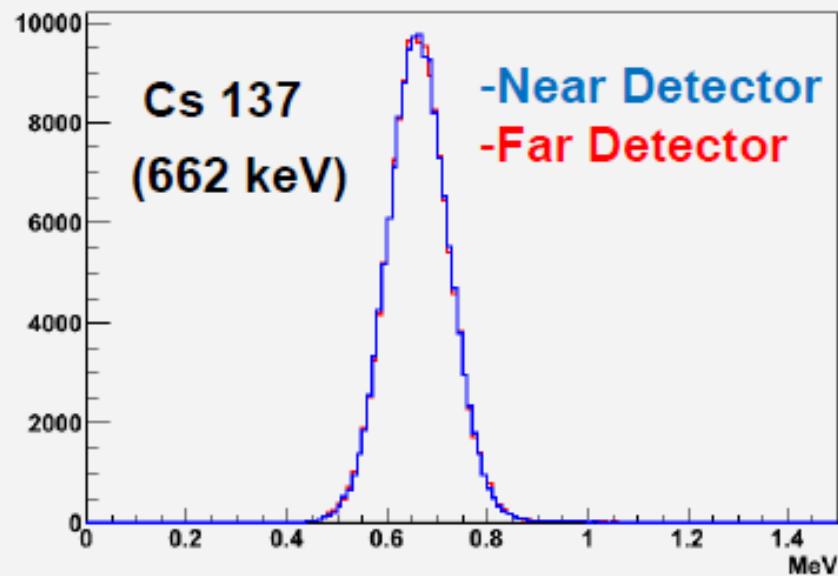
Neutron detection efficiency

Energy & time window, Gd fraction, spill in/out effects

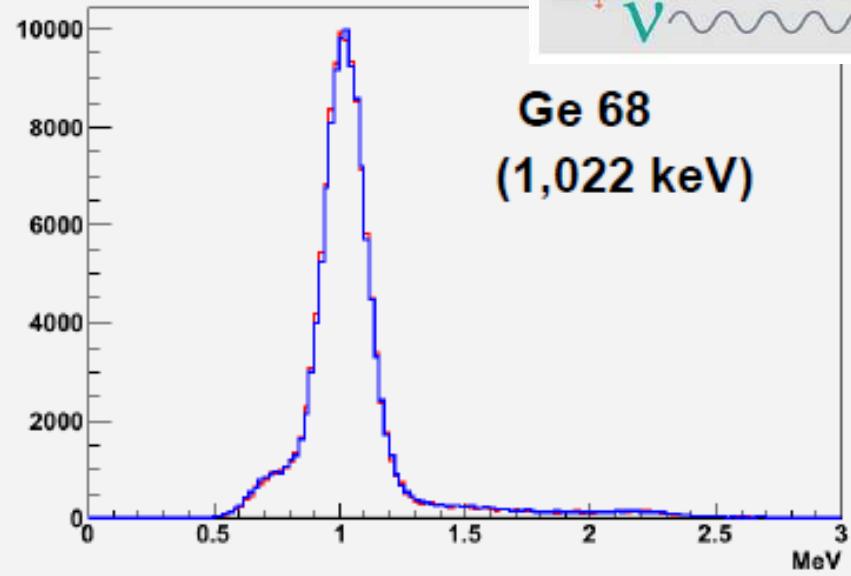
- ^{252}Cf source deployed into ν -target and γ -catcher



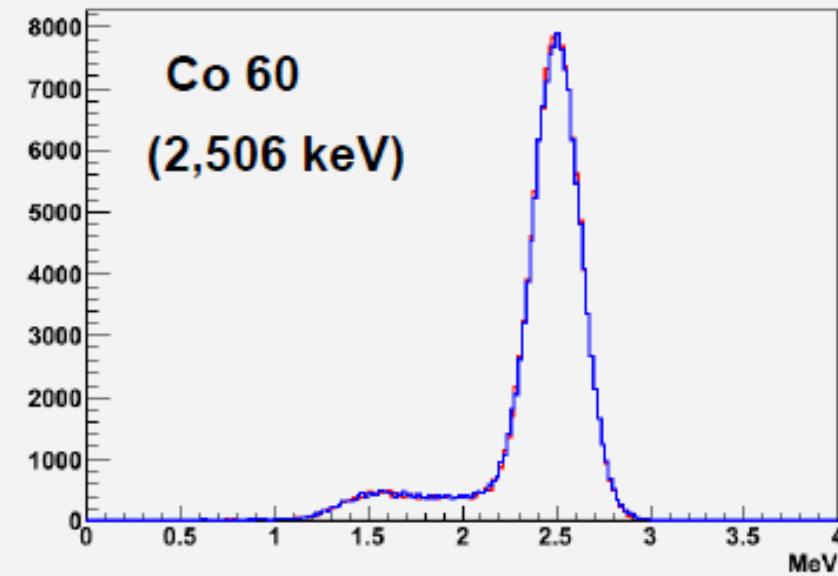
Energy Distribution(Cs)



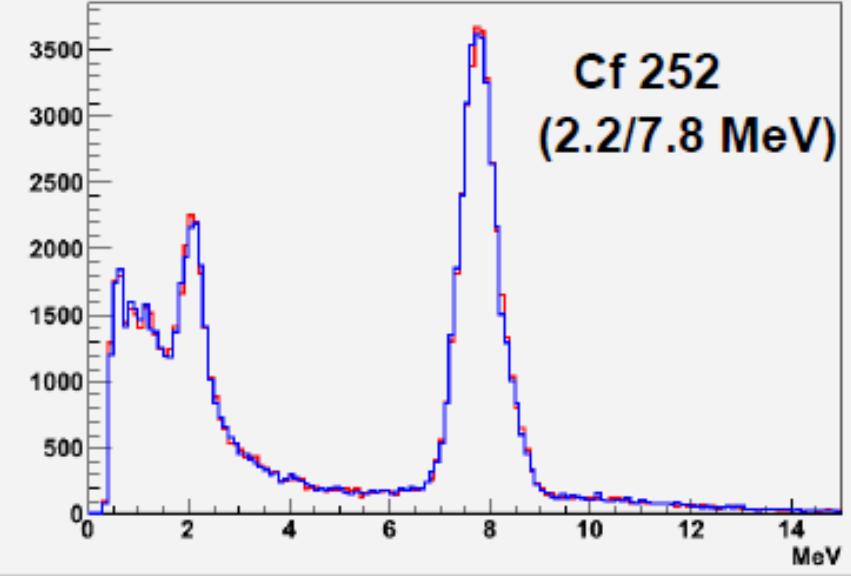
Energy Distribution(Ge)



Energy Distribution(Co)



Energy Distribution(Cf)



ENERGY RESOLUTION

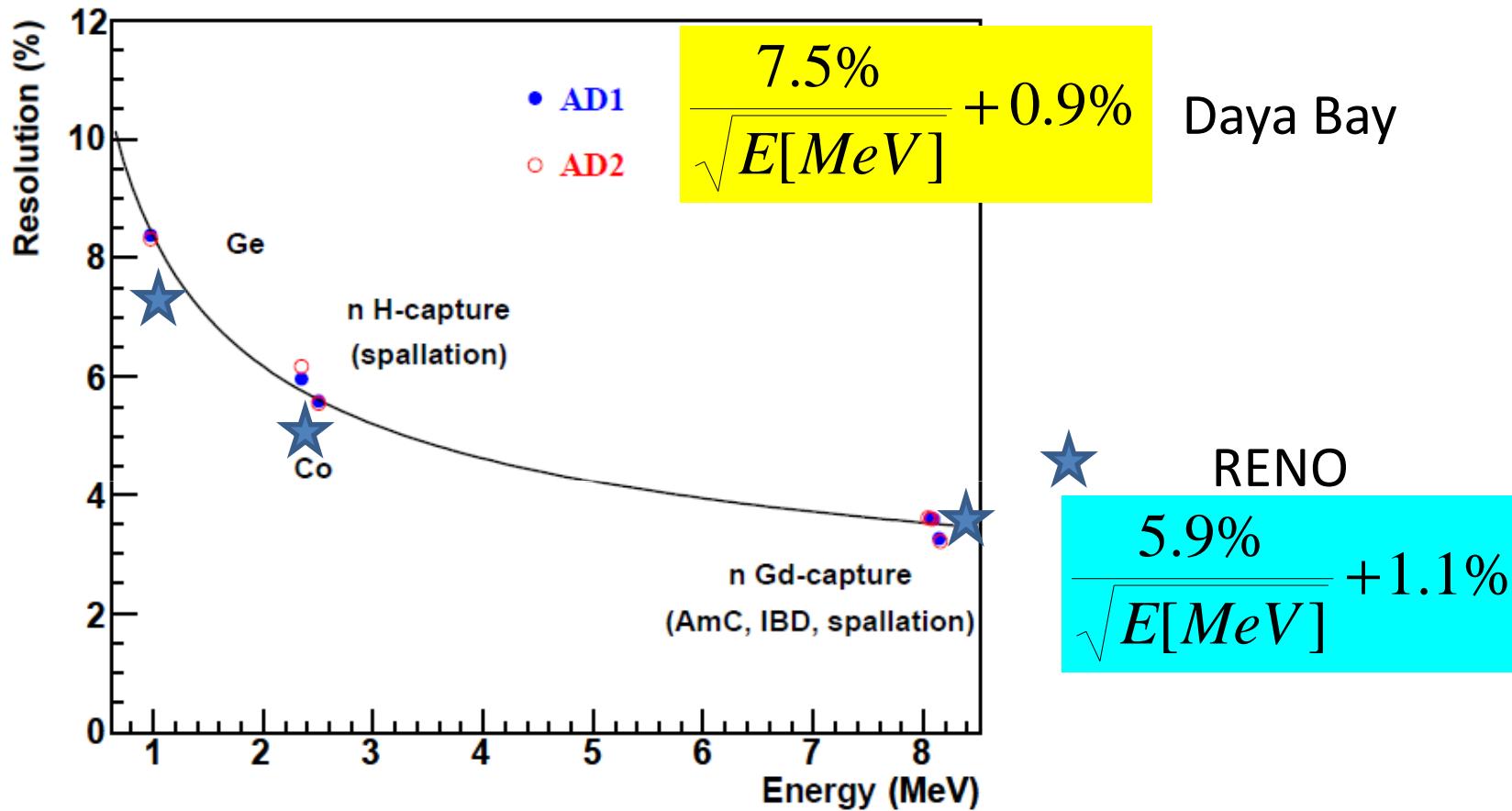
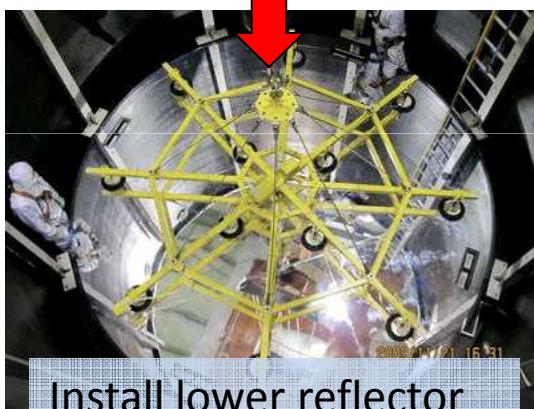


Figure 25: Resolution of reconstructed energy.

Assembly of Anti-neutrino detectors



Stainless Steel Vessel (SSV) in assembly pit



Install lower reflector



Install Acrylic Vessels

ADs are assembled in clean-room



Install PMT ladders



Install top reflector

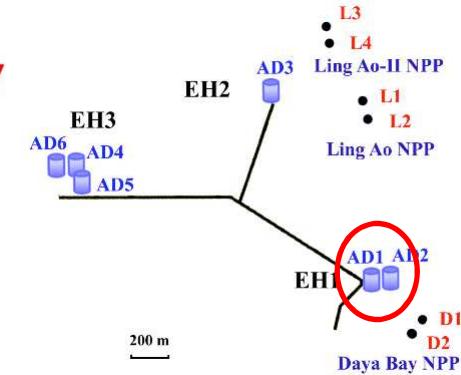


Close SSV lid

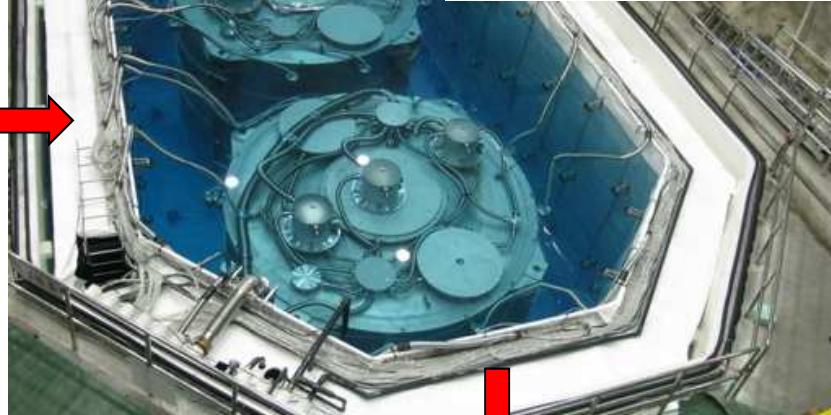


Install calibration units

Near Hall (EH1) Installation



Fill pool with purified water (~1 wk)



Install filled AD1 and AD2 in pool



Data taking started on 15 Aug 2011

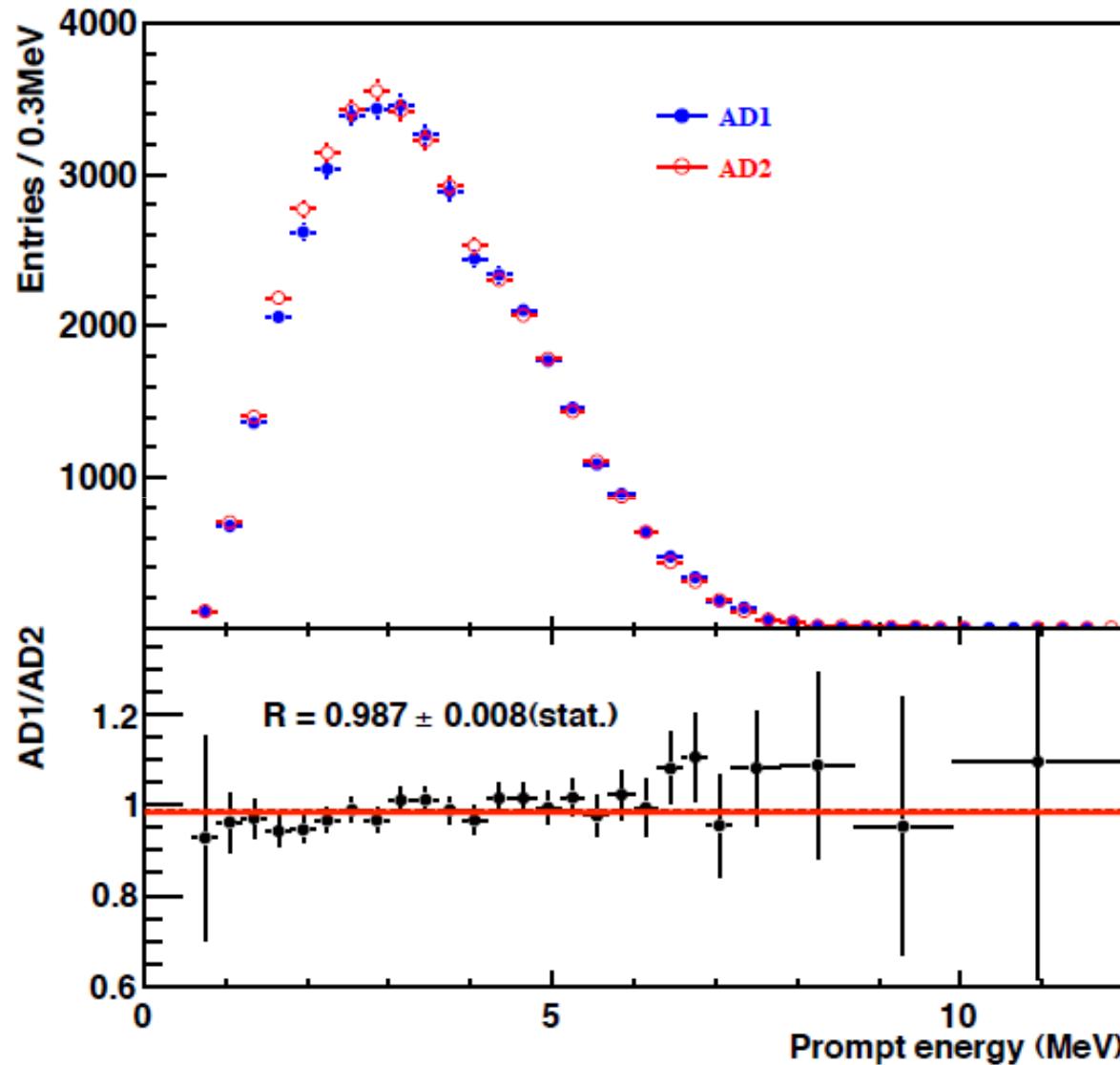


13.9.2012
Roll RPC over cover



PIC2012, Štroské Pleso
Place cover over pool

Detailed comparison of AD1 and AD2



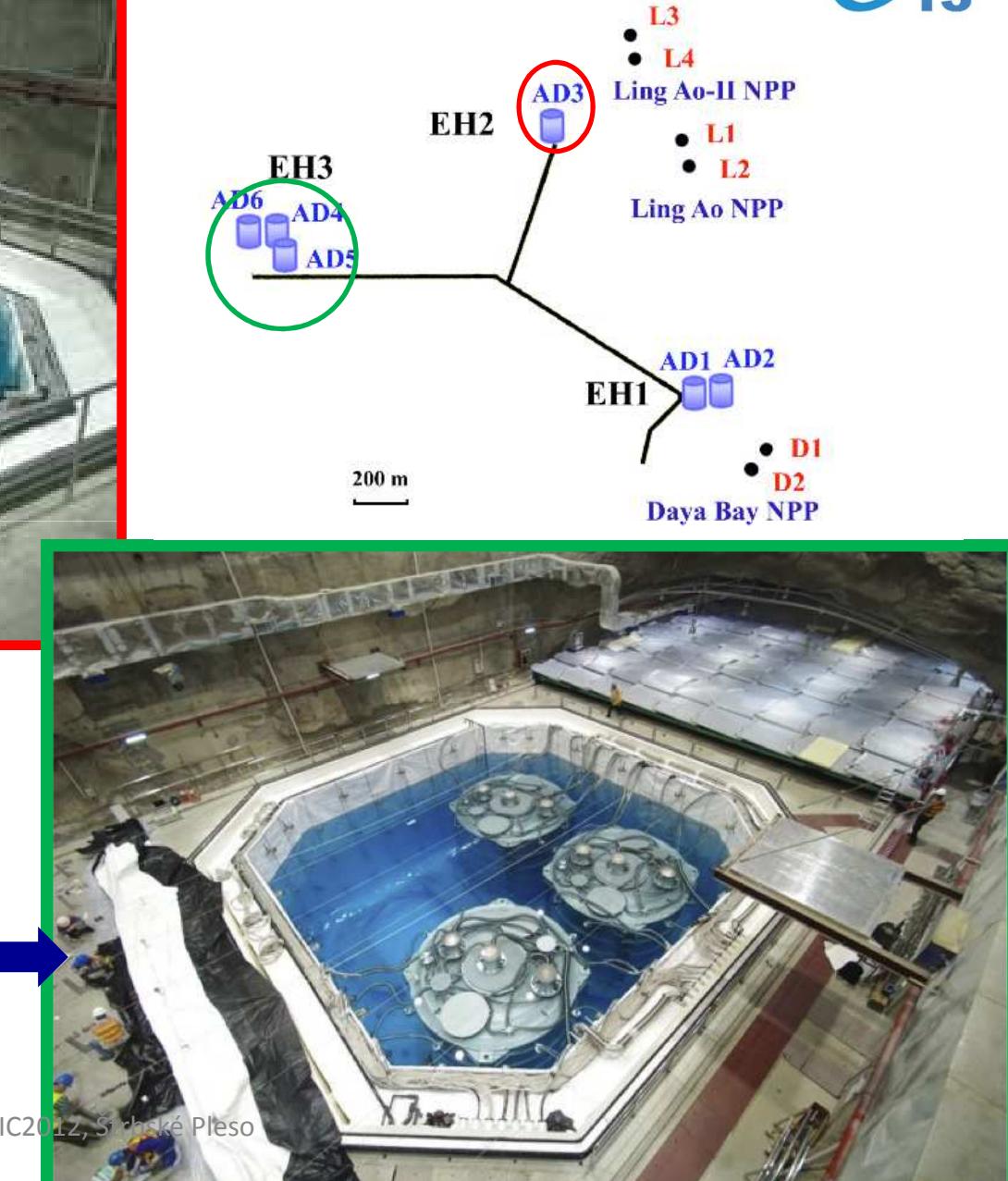
Ling Ao (EH2) and Far (EH3) Halls



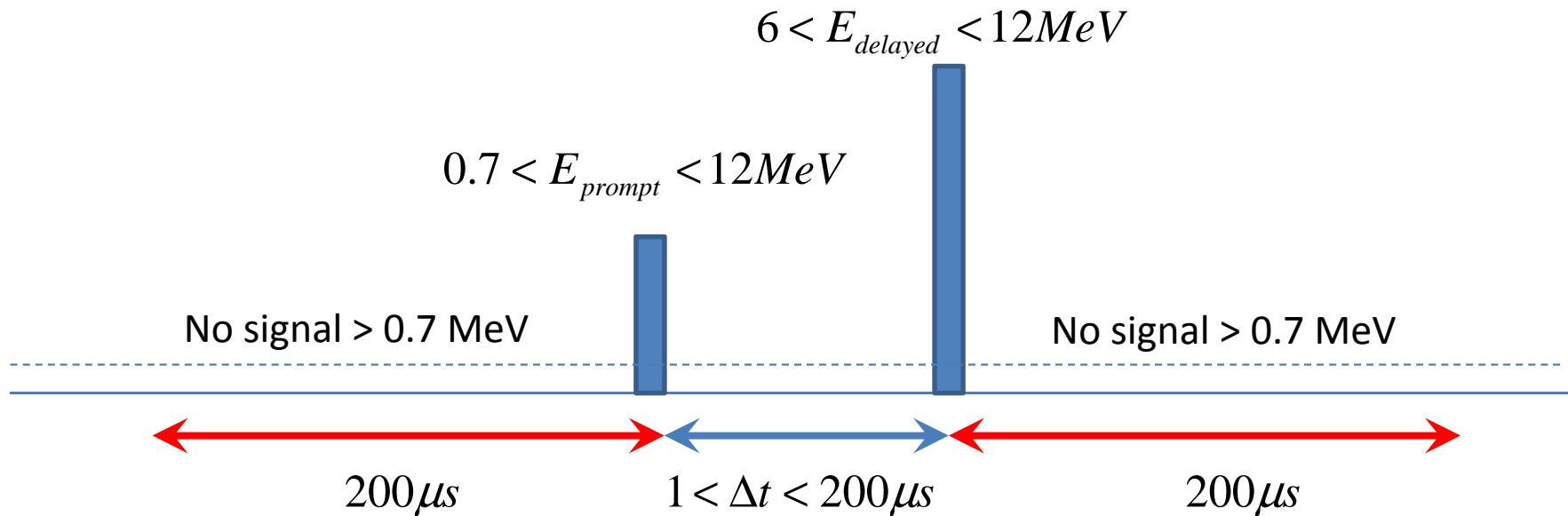
EH2 (Ling Ao Near Hall):
Began operation on 5 Nov 2011

EH3 (Far Hall):
Started data-taking on
24 Dec 2011

- Remaining two ADs will be installed in 2012

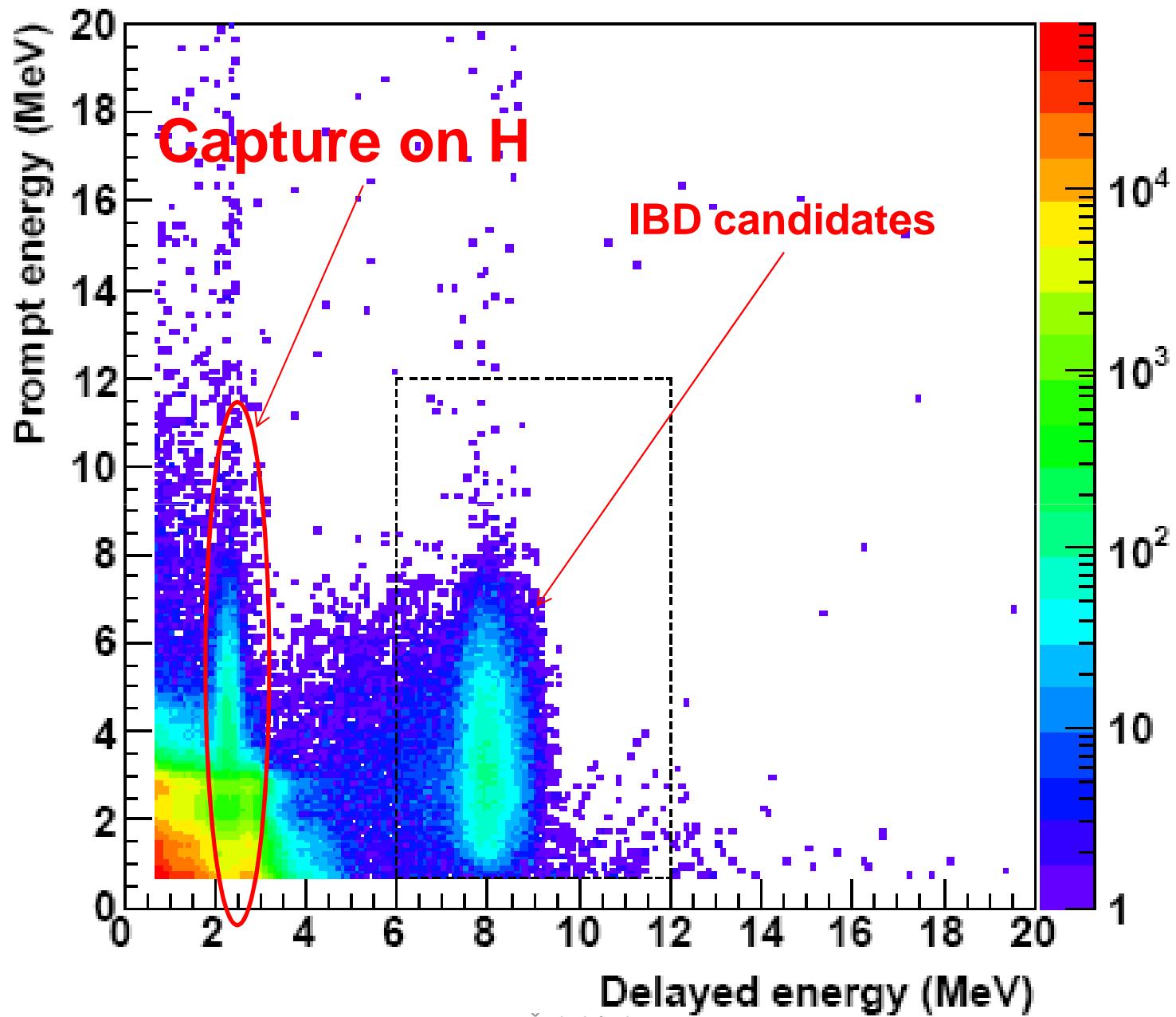


Inverse Beta Decay Selection



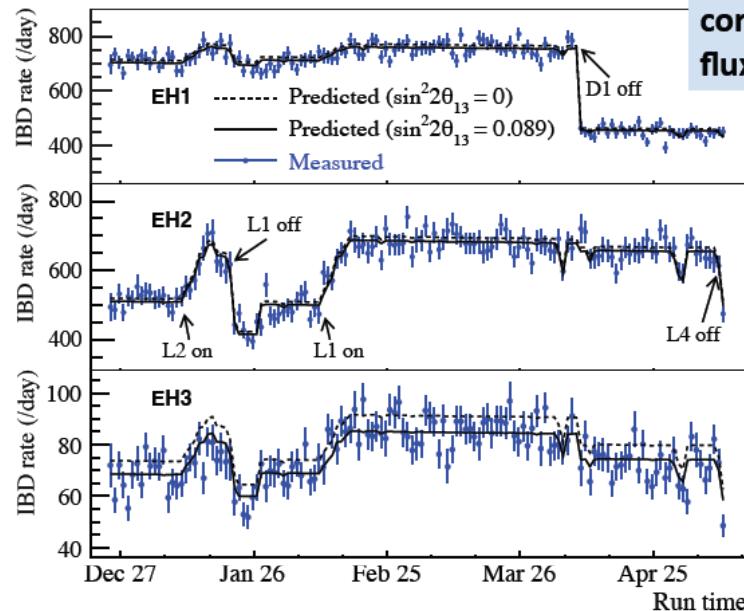
- **Prompt-delayed coincidence:**
 - Prompt positron: $0.7 \text{ MeV} < E_p < 12 \text{ MeV}$ (DYB, RENO), 12.2 MeV (Double Chooz)
 - Delayed neutron: $6.0 \text{ MeV} < E_d < 12 \text{ MeV}$ (DYB, RENO, Double Chooz)
 - Capture Time: $1 \mu s < \Delta t < 200 \mu s$ (DYB), $2 \mu s < \Delta t < 100 \mu s$ (Double Chooz, RENO)
- **Multiplicity:**
 - No signal $200 \mu s$ (Daya Bay), $100 \mu s$ (Double Chooz, RENO) around IBD

- Muon Veto
 - Pool muon: veto following **0.6 ms**
 - AD muon (> 20 MeV): veto following **1 ms**
 - AD shower muon (> 2.5 GeV): veto following **1 s**
that is > 5 T_{1/2} of Li9/He8 isotopes
 - Muon > 600 MeV veto 0.5 s (Double Chooz)
 - Muon > 1.5 GeV veto 0.01 s (RENO)



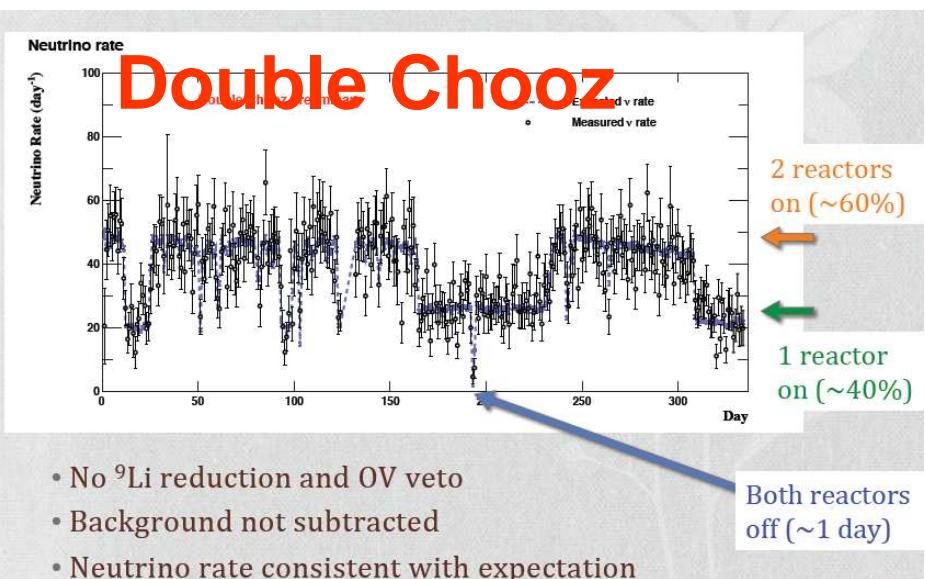
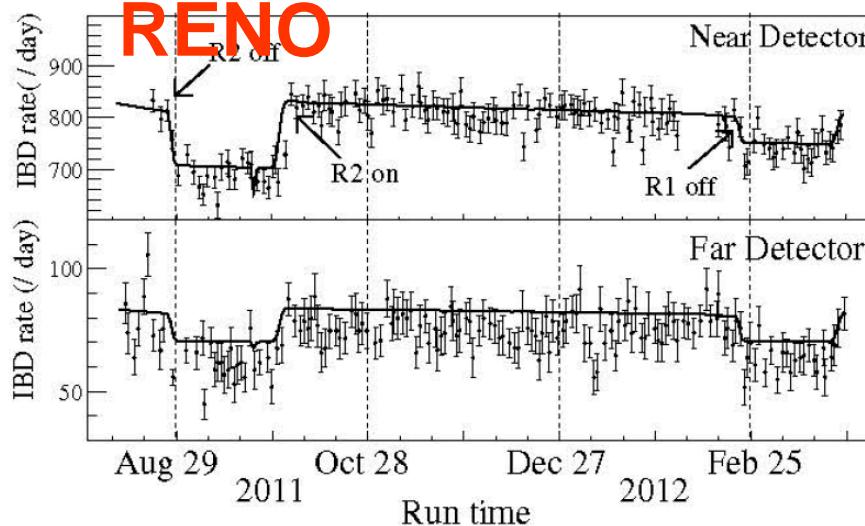


Antineutrino Rate vs. Time



Detected rate strongly correlated with reactor flux expectations.

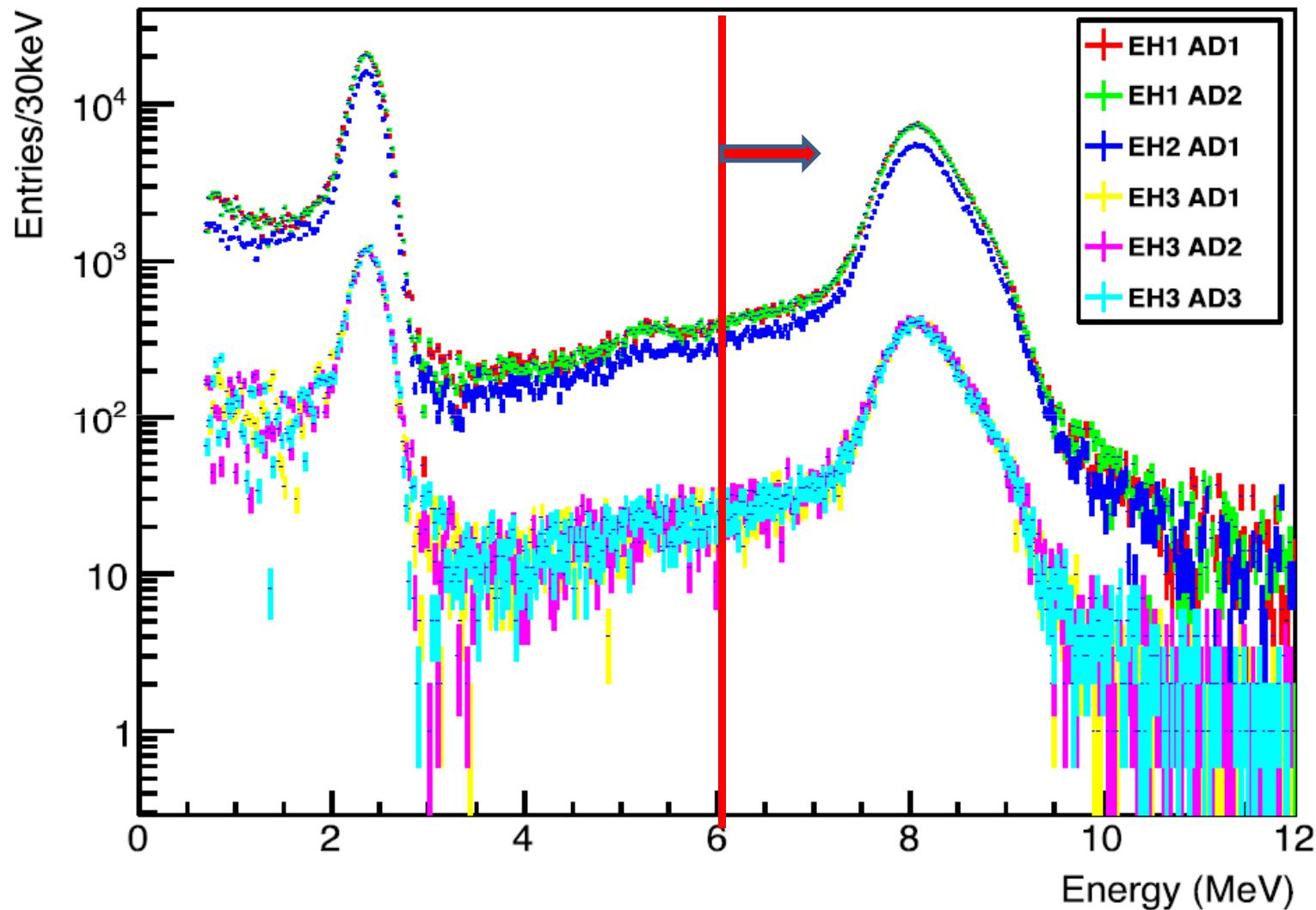
IBD candidates/day	Near	Far
Daya Bay	662+671(EH1) 613 (EH2)	78+75+77
Double Chooz		~40
RENO	779	73



Systematics

	Daya Bay		Reno		Double Chooz
	Corr.	Uncorr.	Corr.	Uncorr.	Corr/Uncorr.
Target proton	0.47%	0.03%	0.5%	0.1%	0.3%
Flasher cut	0.01%	0.01%	0.1%	0.01%	-
Delayed energy cut	0.6%	0.12%	0.5%	0.1%	0.7%
Prompt energy cut	0.1%	0.01%	0.1%	0.01%	-
Energy response	-	-	-	-	0.3%
Trigger efficiency					<0.1%
Multiplicity cut	0.02%	<0.01%	0.06%	0.04%	-
Capture time cut	0.12%	0.01%	0.5%	0.01%	0.5%
Gd capture ratio	0.8%	<0.1%	0.7%	0.1%	0.3%
Spill-in	1.5%	0.02%	1.0%	0.03%	0.3%
livetime	0.002%	<0.01%			-
Muon veto cut	-	-	0.06%	0.04%	-
Total	1.9%	<small>RNC2011 Štrbské Pleso</small>	1.5%	0.2%	1.0%

DELAYED ENERGY CUT



Gd content is monitored by measurement of the time of neutron capture on Gd

$$\langle \tau \rangle = \frac{1}{\langle v_n \cdot \sigma \rangle \cdot N_{Gd}}$$

Very simplified estimation using Maxwell Boltzmann distribution of neutron velocities, app. cross section and Gd concentration gives:

$$\frac{dP}{dv} = \sqrt{\frac{2}{\pi}} \left(\frac{m}{kT} \right)^3 v^2 e^{-\frac{mv^2}{2kT}} \Rightarrow m.p. \quad v = \sqrt{\frac{2kT}{m_n}} \xrightarrow{T=20^\circ C} 2200 \text{ m/s}$$

$$N_{Gd} = 0.103\% \cdot \rho \cdot \frac{N_A}{\langle A_{Gd} \rangle} = 0.00103 \cdot 0.86 \text{ g/cm}^3 \cdot \frac{6.022 / \text{mol}}{157.25 \text{ g/mol}} = 3.29 \cdot 10^{18} / \text{cm}^3$$

$$\langle \sigma \rangle = (0.148 \cdot 60900 + 0.1565 \cdot 254000) \cdot 10^{-24} \text{ cm}^2 = 4.876 \cdot 10^{-20} \text{ cm}^2$$

$$\langle \tau \rangle = \frac{1}{v \cdot \langle \sigma \rangle \cdot N_{Gd}} = 28.3 \mu\text{s}$$

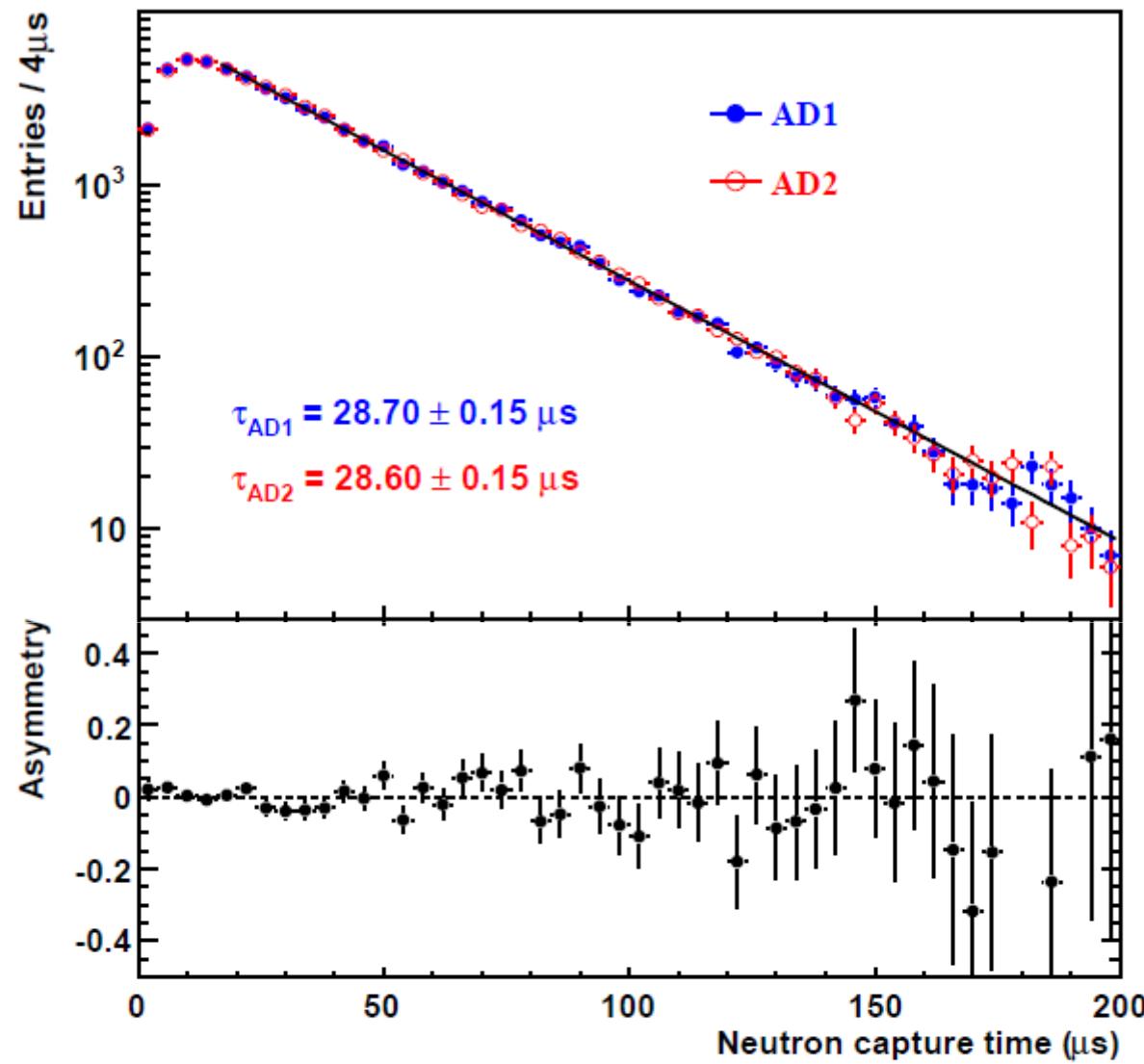


Figure 14: The neutron capture time on Gd from the Am- ^{13}C source at the detector center.

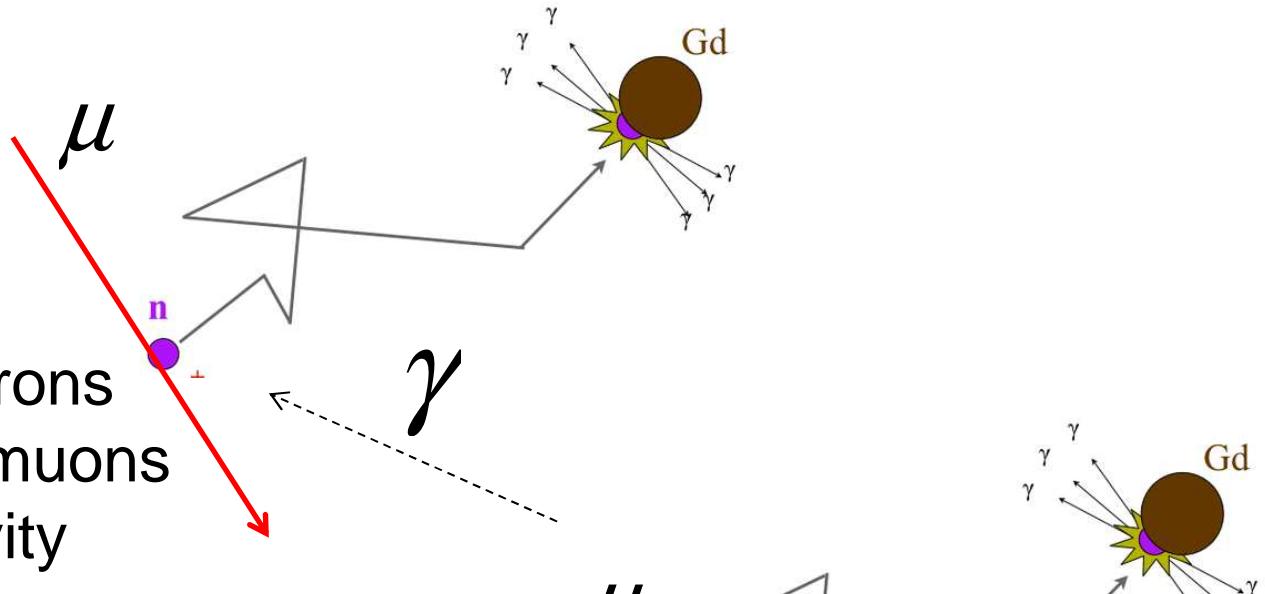
Reactor flux estimate

	Daya Bay		Reno		Double Chooz
	Corr.	Uncorr.	Corr.	Uncorr.	Corr./Uncorr.
Thermal power		0.5%		0.5%	0.5%
Fission fraction/Fuel composition		0.6%		0.7%	0.9%
Fission cross section /Bugey 4 measurement	3%		1.9%		1.4%
Reference spectra			0.5%		0.5%
IBD cross section			0.2%		0.2%
Energy per fission	0.2%		0.2%		0.2%
Baseline	0.02%		-		0.2%
Spent fuel		0.3%			
Total	3%	0.8%	2.0%	0.9%	1.8%

Backgrounds

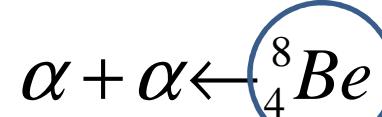
-Accidental

coincidences of neutrons produced by cosmic muons and natural radioactivity

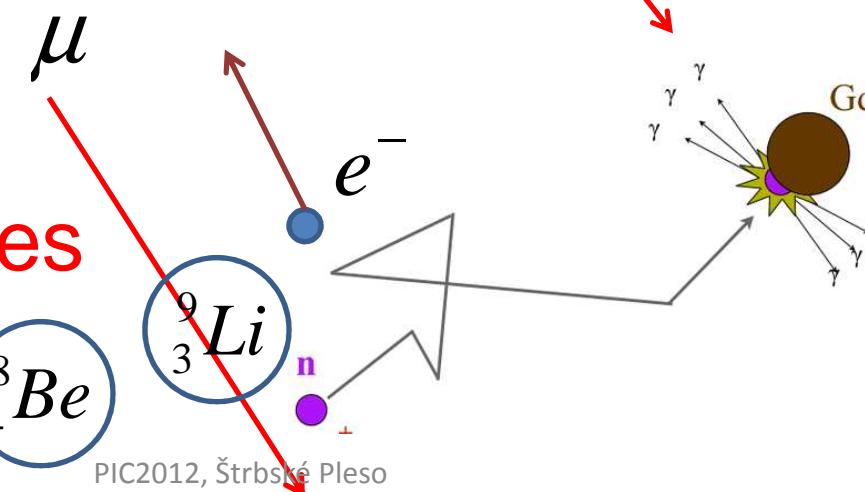


Fast neutrons

Li9 and He8 isotopes

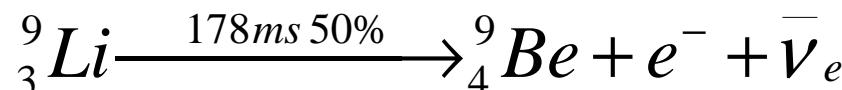
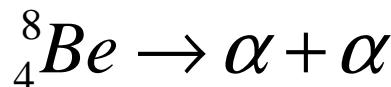
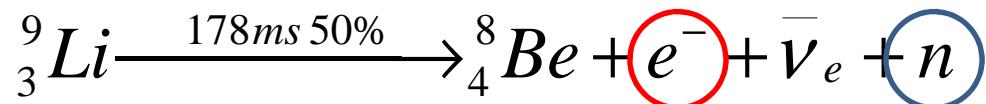
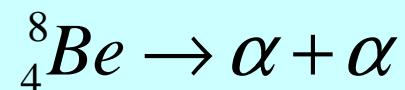
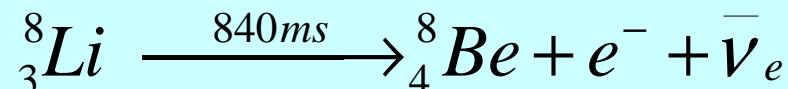
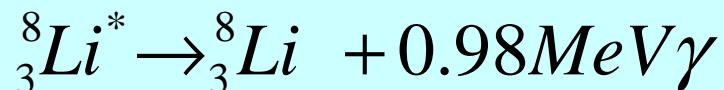
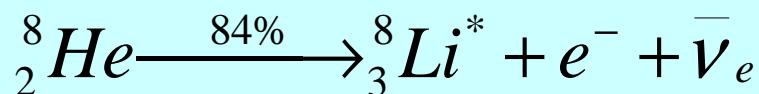
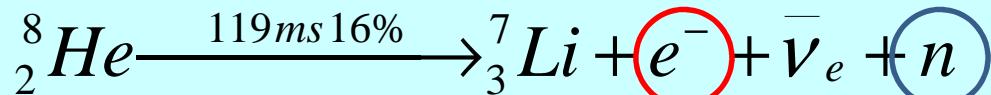


PIC2012, Štrbské Pleso



Li9 and He8 background

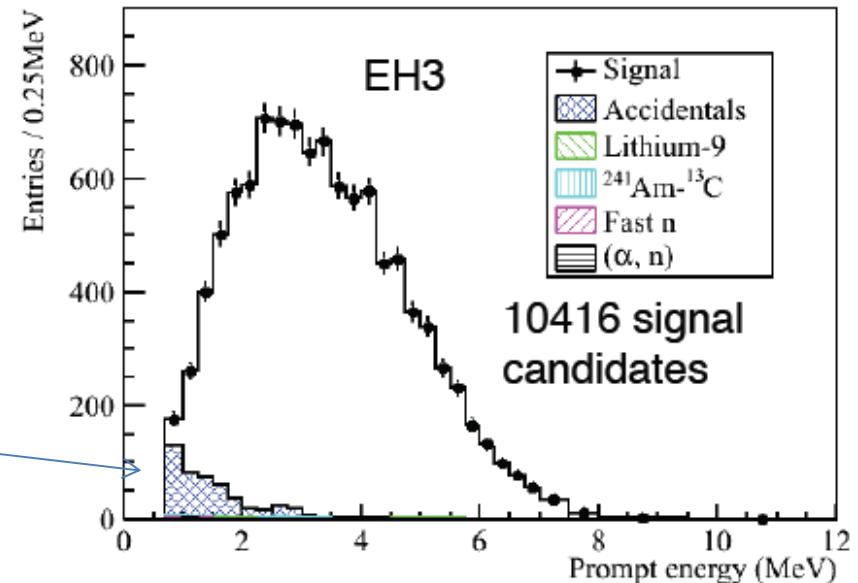
These isotopes are products of photonuclear interactions of cosmic muons on C



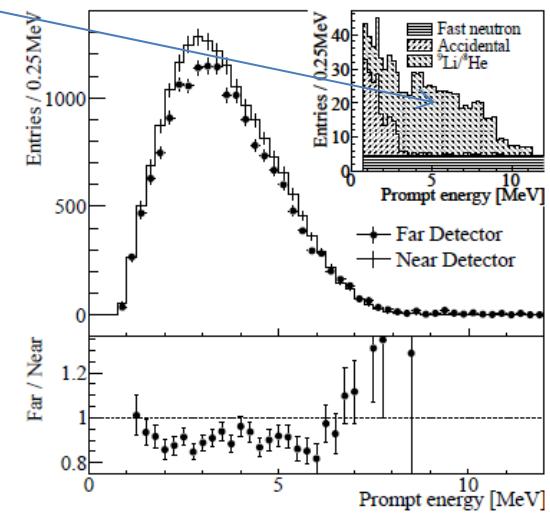
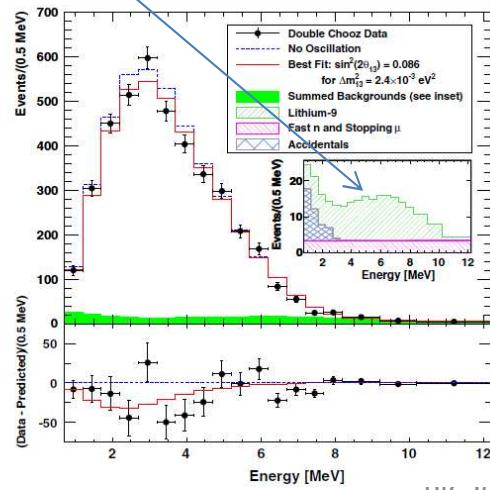
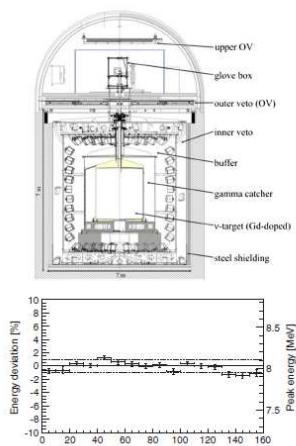
Backgrounds & uncertainties

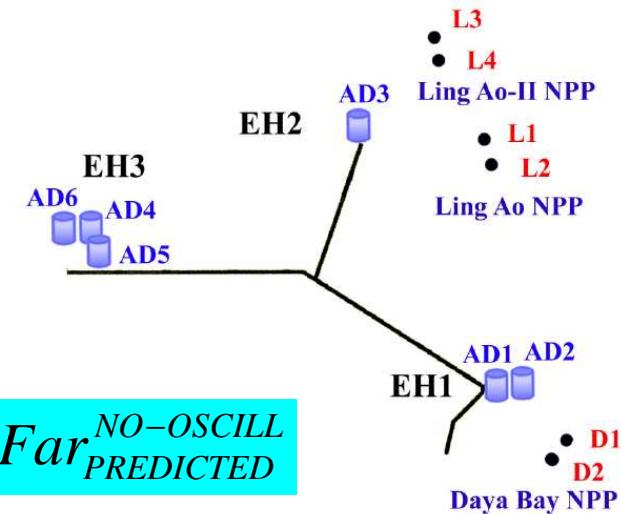
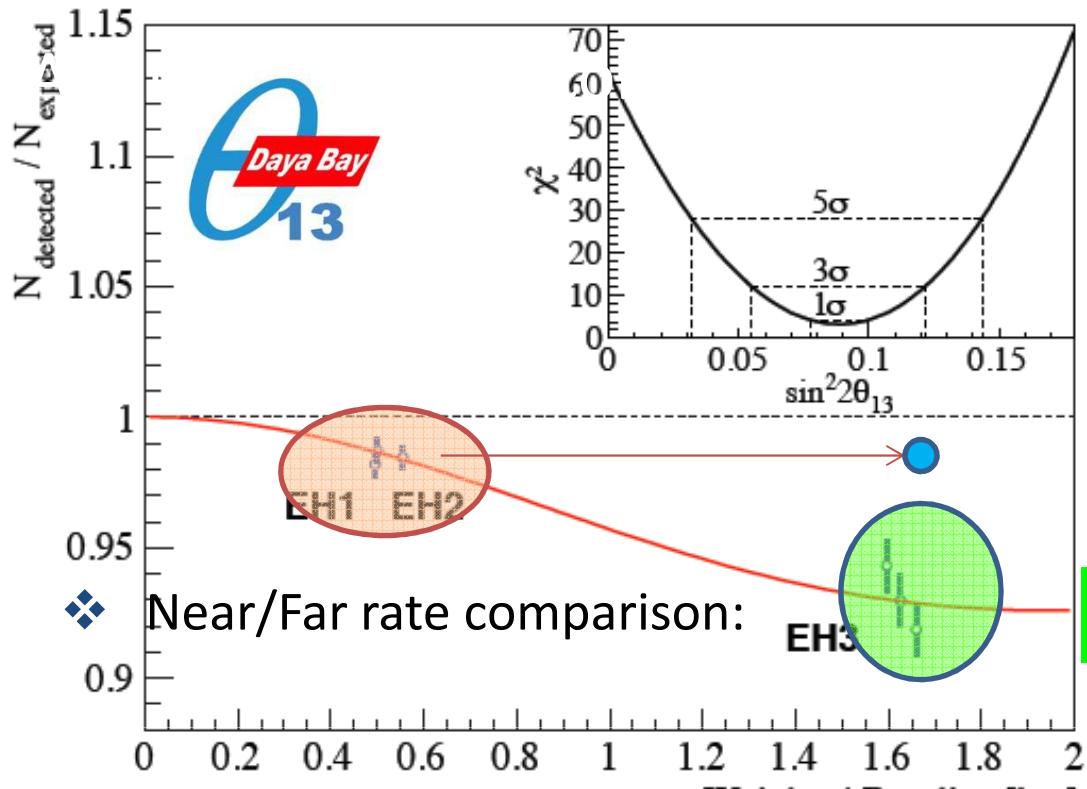
	Daya Bay		Reno		Double Chooz
	Near	Far	Near	Far	Far
Accidentals (B/S)	1.4%	4.0%	0.56%	0.93%	0.6%
Uncertainty($\Delta B/B$)	1.0%	1.4%	1.4%	4.4%	0.8%
Fast neutrons(B/S)	0.1%	0.06%	0.64%	1.3%	1.6%
Uncertainty($\Delta B/B$)	31%	40%	2.6%	6.2%	30%
$^8\text{He}/^9\text{Li}$ (B/S)	0.4%	0.3%	1.6%	3.6%	2.8%
Uncertainty ($\Delta B/B$)	52%	55%	48%	29%	50%
α -n(B/S)	0.01%	0.05%	-	-	-
Uncertainty($\Delta B/B$)	50%	50%	-	-	-
Am-C(B/S)	0.03%	0.3%	-	-	-
Uncertainty ($\Delta B/B$)	100%	100%	-	-	-
Total backgrounds(B/S)	1.9%	4.7%	2.8%	5.8%	5.0%
Total Uncertainties ($\Delta(B/S)$)	0.2%	0.35%	0.8%	1.1%	1.5%

Due to large overburden and strict showering muons veto cut Daya Bay Li/He background is suppressed and the background is dominated by accidental coincidences which are concentrated at low neutrino energies.



In Double Chooz and RENO the background is still dominated by decays of Li/He isotopes that spans the whole range of neutrino energies.





$$R = \frac{\text{Far}_{\text{measured}}}{\text{Far}_{\text{expected}}} = \frac{M_4 + M_5 + M_6}{\sum_{i=4}^6 (\alpha_i(M_1 + M_2) + \beta_i M_3)}$$

M_n : measured rates in each detector.

Weights α_i, β_i : determined from baselines and reactor fluxes,
no oscillations assumed.

$$R = 0.944 \pm 0.007 \text{ (stat)} \pm 0.003 \text{ (syst)}$$

➤ Unambiguous observation of antineutrino deficit at the far site!

Rate-only analysis



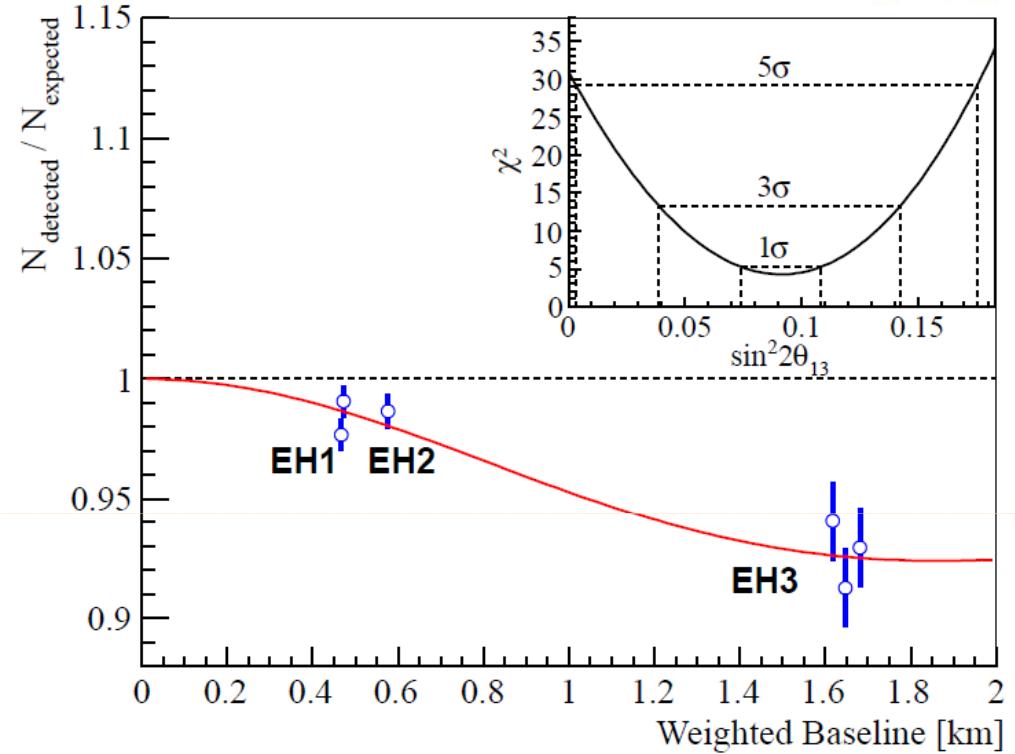
- ❖ Determine θ_{13} using measured rates in each detector:

Uses standard χ^2 approach ($\chi^2/\text{NDF}=4.26/4$)

$$\chi^2 = \sum_{d=1}^6 \frac{[M_d - T_d(1 + \varepsilon + \sum_r \omega_r^d \alpha_r + \varepsilon_d) + \eta_d]^2}{M_d + B_d} + \sum_r \frac{\alpha_r^2}{\sigma_r^2} + \sum_{d=1}^6 \left(\frac{\varepsilon_d^2}{\sigma_d^2} + \frac{\eta_d^2}{\sigma_B^2} \right),$$

[Absolute rate is not constrained.]

Consistent results obtained by independent analyses, different reactor flux models.



$$\sin^2 2 \theta_{13} = 0.089 \pm 0.010 \text{ (stat)} \pm 0.005 \text{ (syst)}$$

$\sin^2 2 \theta_{13} = 0$ excluded at 7.7σ

$$\theta_{13} \simeq 8.7^\circ$$

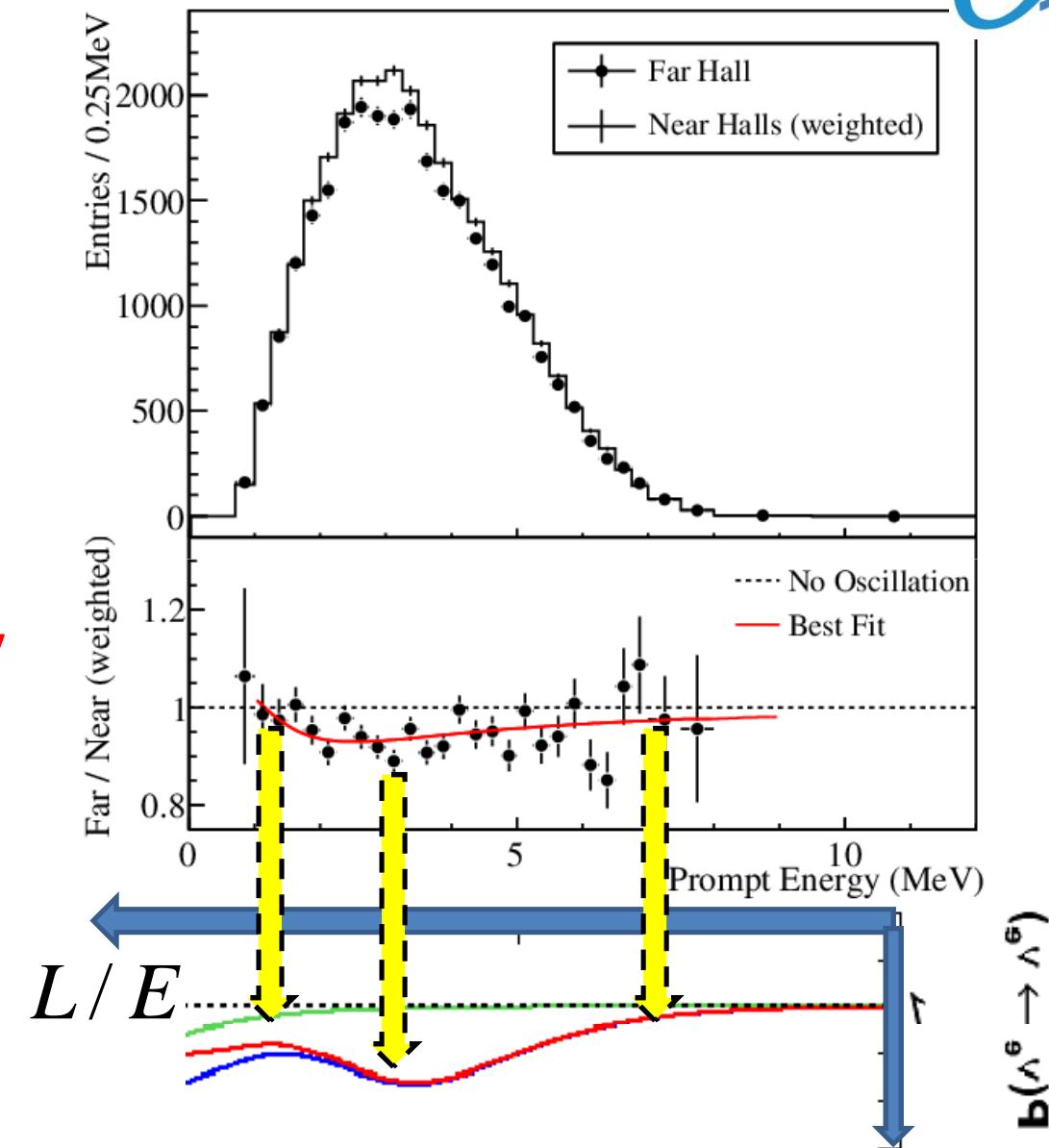
The smallest lepton mixing angle is comparable to largest (Cabibbo) quark mixing angle.

The disagreement of the spectra in far and near hall provides further evidence of neutrino oscillation.

The ratio of the spectra is consistent with the best-fit oscillation solution of

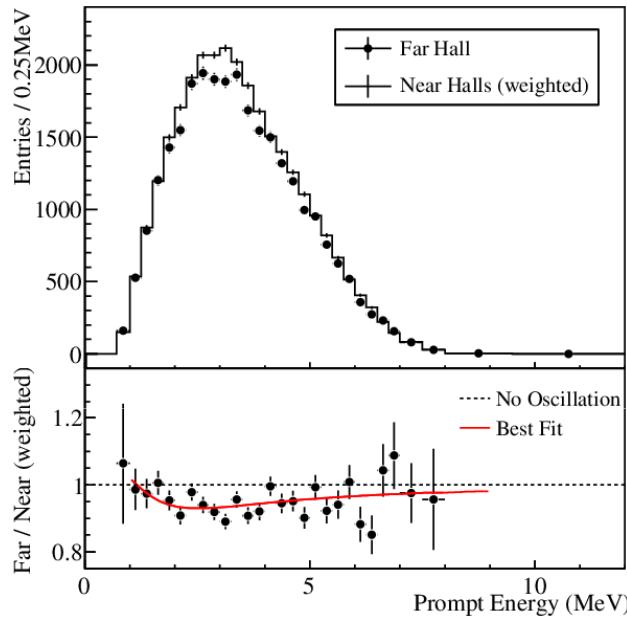
$\sin^2 2\theta_{13} = 0.089$
obtained from the rate-only analysis.

Currently the result is only from rate analysis!



Reactor neutrinos Disappearance: Latest results from Neutrino 2012

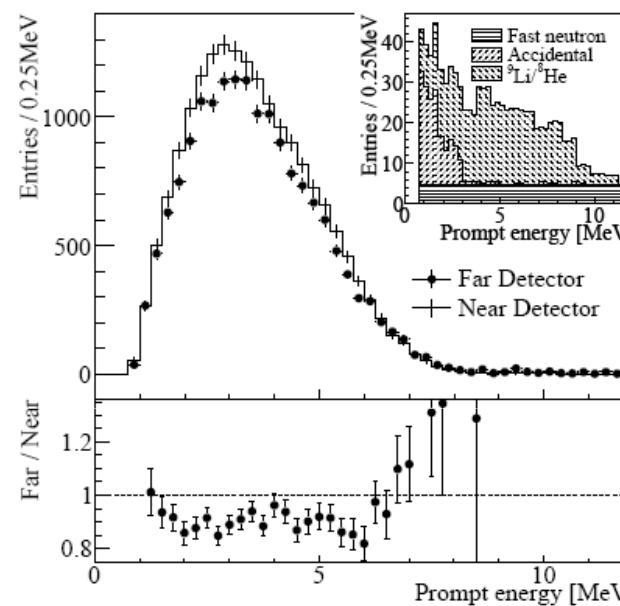
Daya Bay



$$R = 0.944 \pm 0.007 \pm 0.003$$

$\sin^2 2\theta_{13} = 0.089 \pm 0.010 \pm 0.005$
 7.7 σ for non-zero θ_{13}

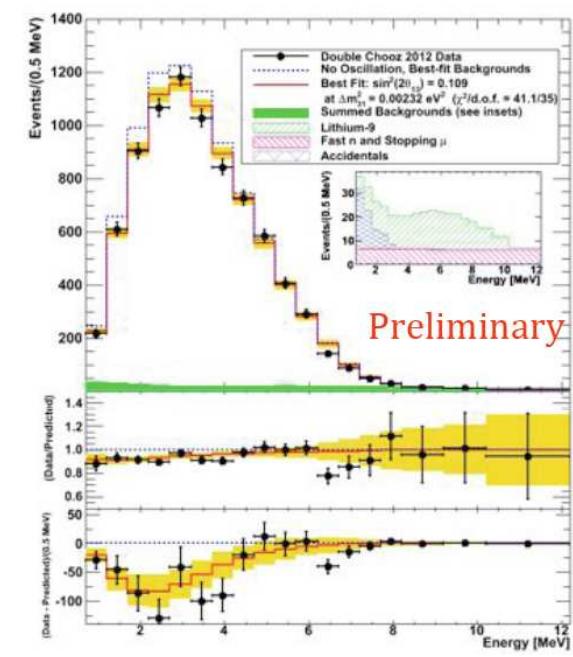
Reno



$$R = 0.920 \pm 0.009 \pm 0.014$$

$\sin^2 2\theta_{13} = 0.113 \pm 0.013 \pm 0.019$
 4. 9 σ for non-zero θ_{13}

Double Chooz



$$\sin^2 2\theta_{13} = 0.170 \pm 0.035 \pm 0.040$$

$\sin^2 2\theta_{13} = 0.109 \pm 0.030 \pm 0.025$
 2.9 σ for non-zero θ_{13}

CONCLUSIONS

-A non zero, surprisingly large value of the third mixing angle theta 13 has been measured in 2012. The result is important as it opens future searches for violation of CP in lepton sector.

-After 2011 hints for non zero value of theta13 from accelerator experiments, combined data and Double Chooz it is important that today we have:

- convincing results from Daya Bay with the significance of 7.7 sigma reported at Neutrino 2012 (the discovery of non zero value with significance exceeding 5 sigma was announced in March and published)

- observation paper published from RENO (significance close to 5 sigma announced in April and published) and

- latest results from Double Chooz (3 sigma significance reported at Neutrino2012 is now on arXiv).

-One can expect improvements in near future:

- reduction of statistical errors and systematic uncertainties with more data

- shape analyses

- completion of Daya Bay setup (2012)

- near detector at Double Chooz experiment (2013)

....

Already now the three experiments collected several hundred thousands antineutrino interactions and I am convinced that new, interesting analyses will be performed using such unique set of data.