

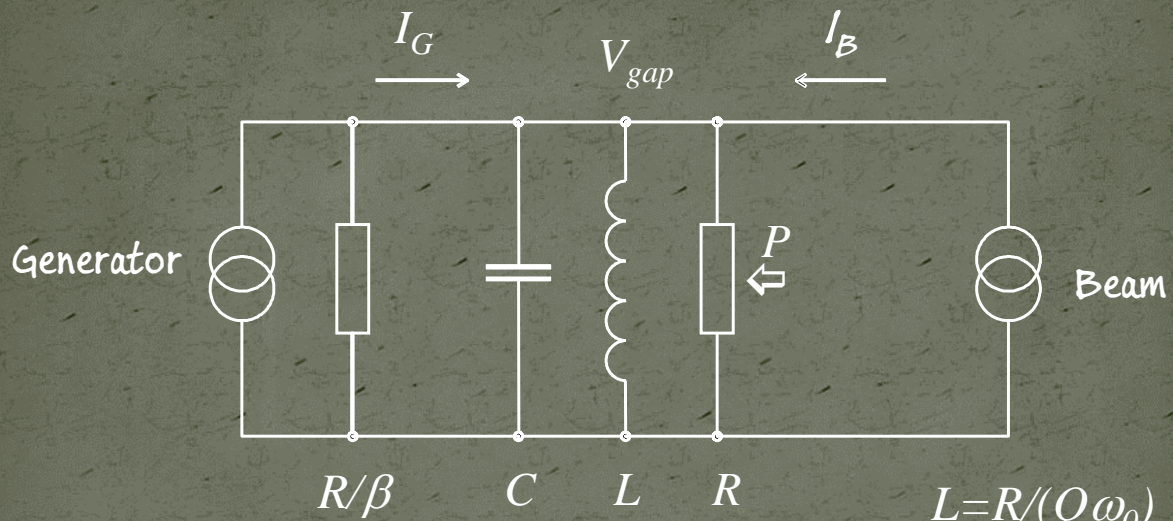
RF Systems II

Erk Jensen, CERN BE-RF

Characterizing a cavity

Cavity resonator – equivalent circuit

Simplification: single mode



β : coupling factor

$\underbrace{\hspace{10em}}$
Cavity

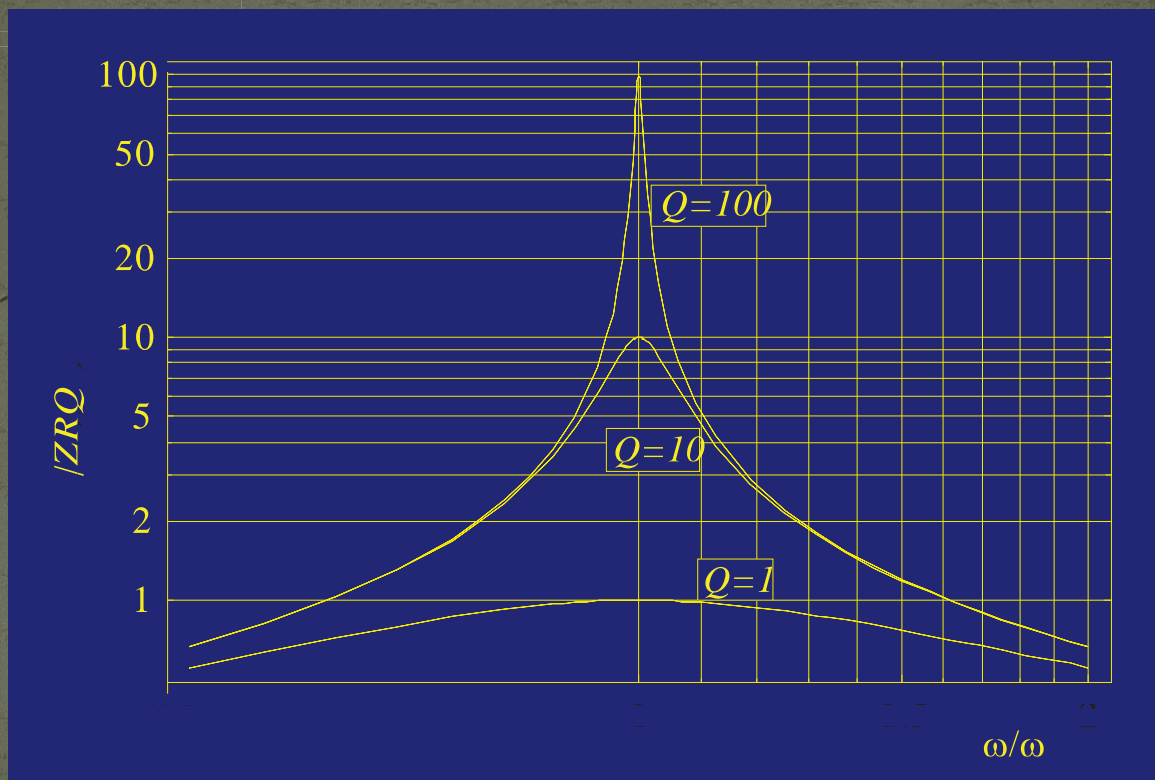
$$C=Q/(R\omega_0)$$

R : Shunt impedance

$$\sqrt{L/C}: R\text{-upon-}Q$$

We have used this before when explaining the “fast feedback”

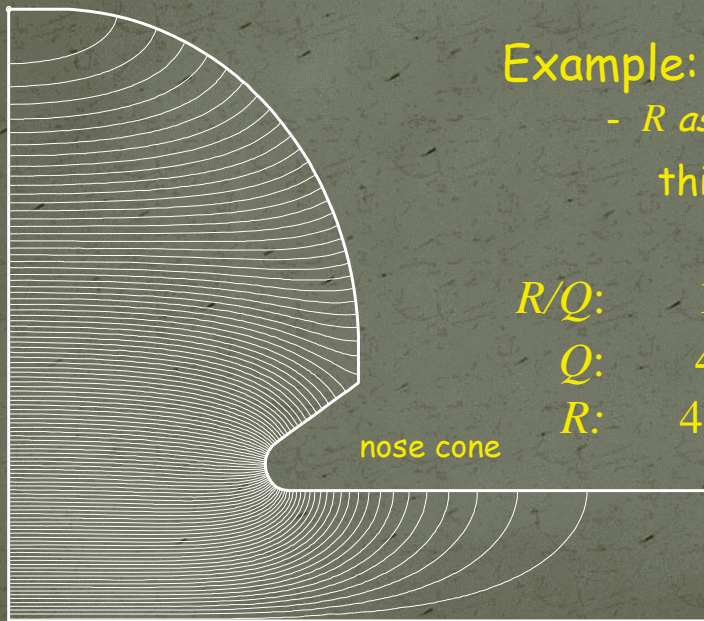
Resonance



Reentrant cavity

Nose cones increase transit time factor, round outer shape minimizes losses.

Nose cone example Freq = 500.003



Example: KEK photon factory 500 MHz
- R as good as it gets -

	this cavity	optimized pillbox
R/Q :	111 Ω	107.5 Ω
Q :	44270	41630
R :	4.9 M Ω	4.47 M Ω

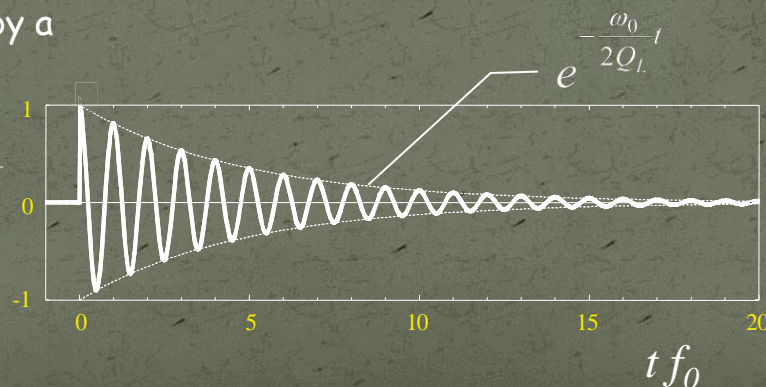
Loss factor

$$k_{loss} = \frac{\omega_0 R}{2 Q} = \frac{V_{gap}^2}{4 W} = \frac{1}{2 C}$$

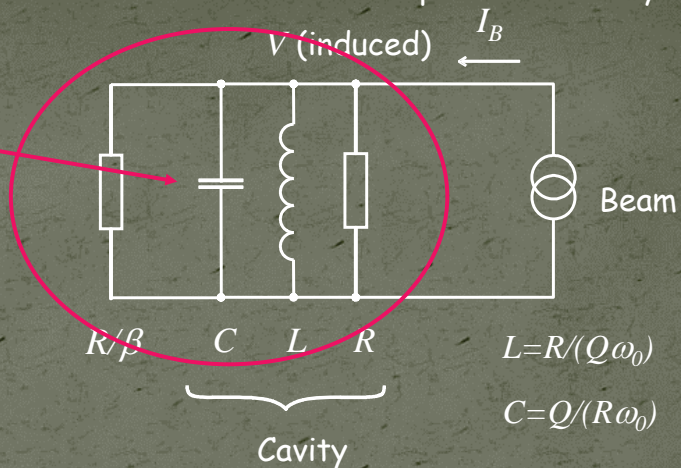
Energy deposited by a single charge q : $k_{loss} q^2$

Voltage induced by a single charge q :

$$\frac{V_{gap}}{2 k_{loss} q}$$



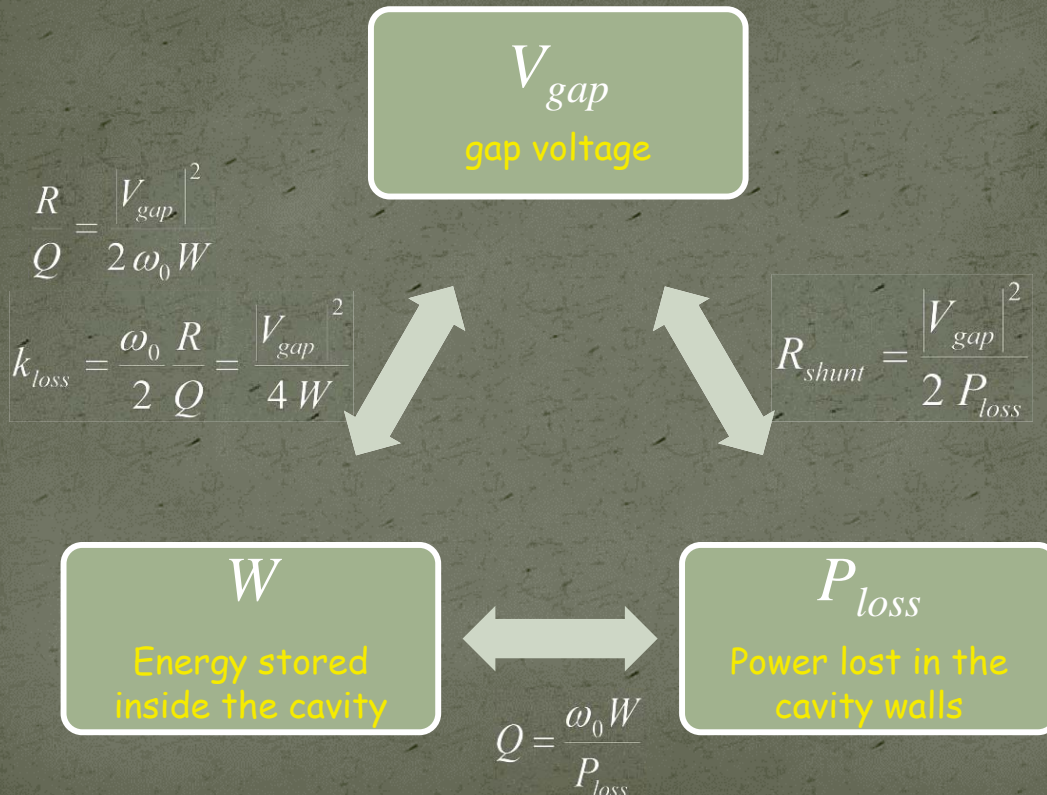
Impedance seen by the beam



$$L = R / (Q \omega_0)$$

$$C = Q / (R \omega_0)$$

Summary: relations V_{gap} , W , P_{loss}



Let's talk about RF → beam efficiency!

- With zero beam current, RF power fed into the cavity excites a gap voltage, but it will be entirely lost in the cavity walls; this is characterized by the shunt impedance R :

$$|V_{acc}| = \frac{1}{2} \left(\sqrt{(4P)R} \right)$$

- A non-zero beam current induces a voltage reducing the gap voltage*); this is known as **beam loading** and normally considered a disadvantage.

$$|V_{acc}| = \frac{1}{2} \left(\sqrt{(4P + I_{beam}^2 R)R} - I_{beam} R \right)$$

- But: if we define the RF to beam efficiency as “increase of beam power” divided by “RF input power”, we find that large efficiency can be obtained only with large beam loading (at the expense of reduced accelerating voltage).
- Example: CLIC drive beam accelerated with 98% RF to beam efficiency.

*) for an accelerated beam! For a decelerated beam the voltage is increased

Cavity parameters

- Resonance frequency

$$\omega_0 = \frac{1}{\sqrt{L \cdot C}}$$

- Transit time factor

field varies while particle is traversing the gap

$$\int E_z e^{j \frac{\omega}{c} z} dz$$

$$\int E_z dz$$

Circuit definition

$$|V_{gap}|^2 = 2 R_{shunt} P_{loss}$$

- Shunt impedance
gap voltage – power relation

Linac definition

$$|V_{gap}|^2 = R_{shunt} P_{loss}$$

- Q factor

$$\omega_0 W = Q P_{loss}$$

- R/Q
independent of losses – only geometry!

$$\frac{R}{Q} = \frac{|V_{gap}|^2}{2 \omega_0 W} = \sqrt{\frac{L}{C}}$$

$$\frac{R}{Q} = \frac{|V_{gap}|^2}{\omega_0 W}$$

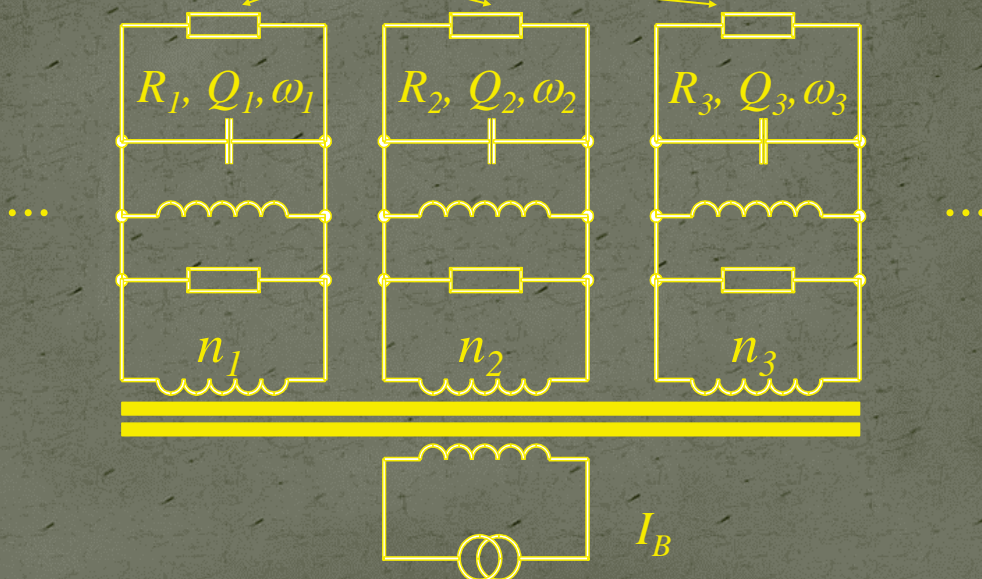
- loss factor

$$k_{loss} = \frac{\omega_0 R}{2 Q} = \frac{|V_{gap}|^2}{4 W}$$

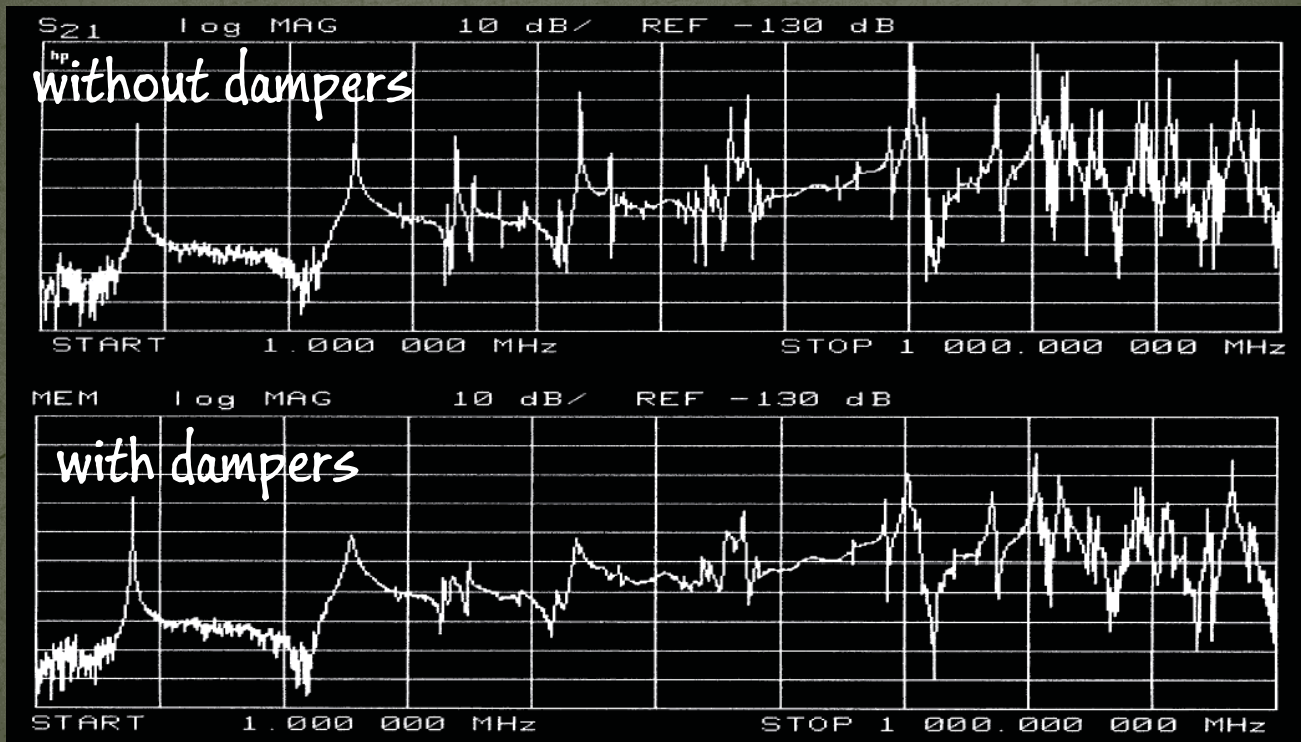
$$k_{loss} = \frac{\omega_0 R}{4 Q} = \frac{|V_{gap}|^2}{4 W}$$

Higher order modes (HOM's)

external dampers

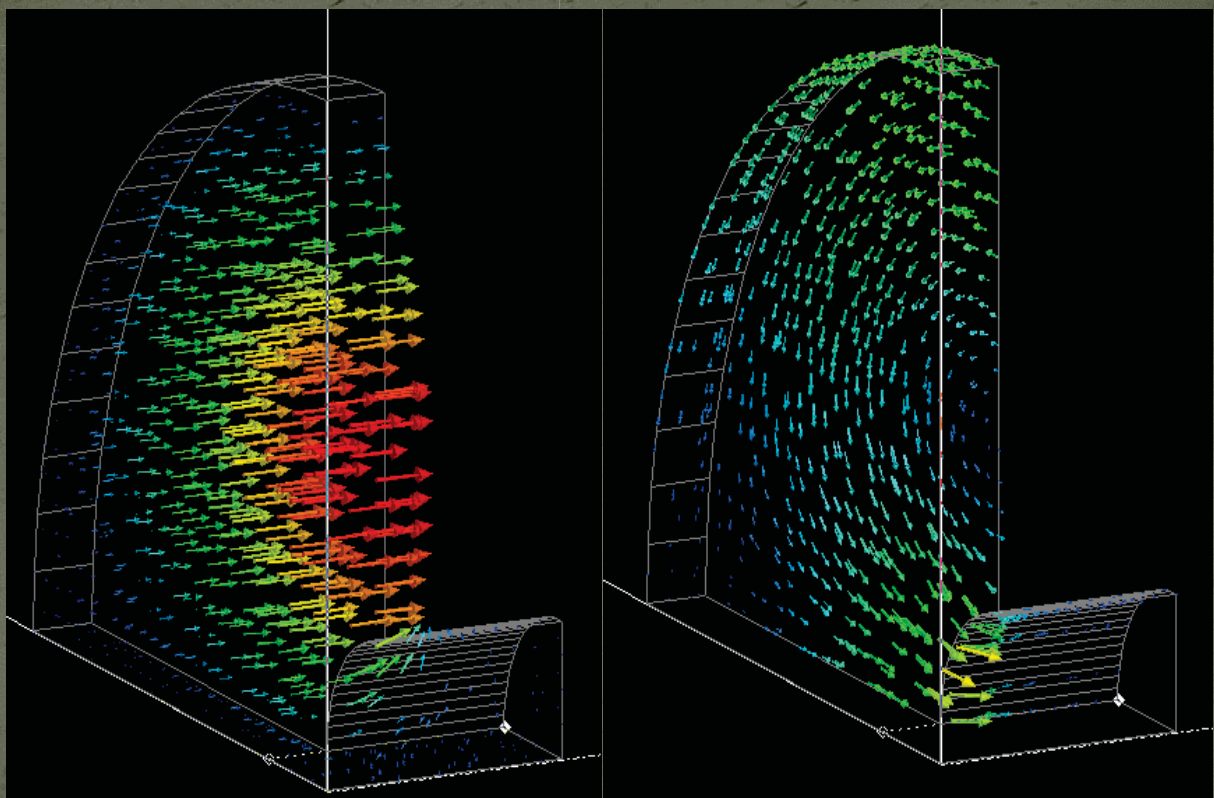


HOM (measured spectrum)



Pillbox: Dipole mode

(TM_{110}) (only 1/8 shown)



electric field (@ 0°)

magnetic field (@ 90°)

Panofsky-Wenzel theorem

For particles moving virtually at $v = c$, the integrated transverse force (kick) can be determined from the transverse variation of the integrated longitudinal force!

$$j \frac{\omega}{c} \vec{F}_{\perp} = \nabla_{\perp} F_{\parallel}$$

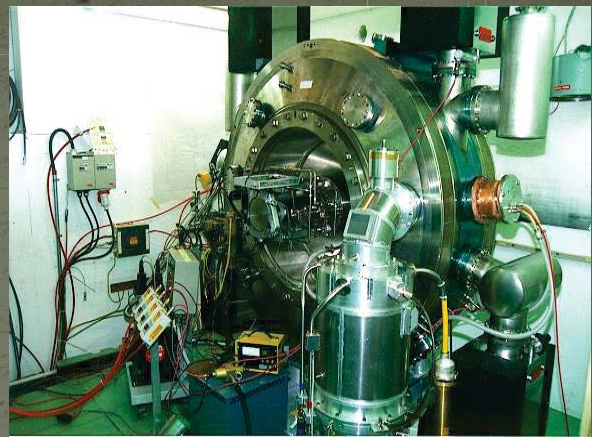
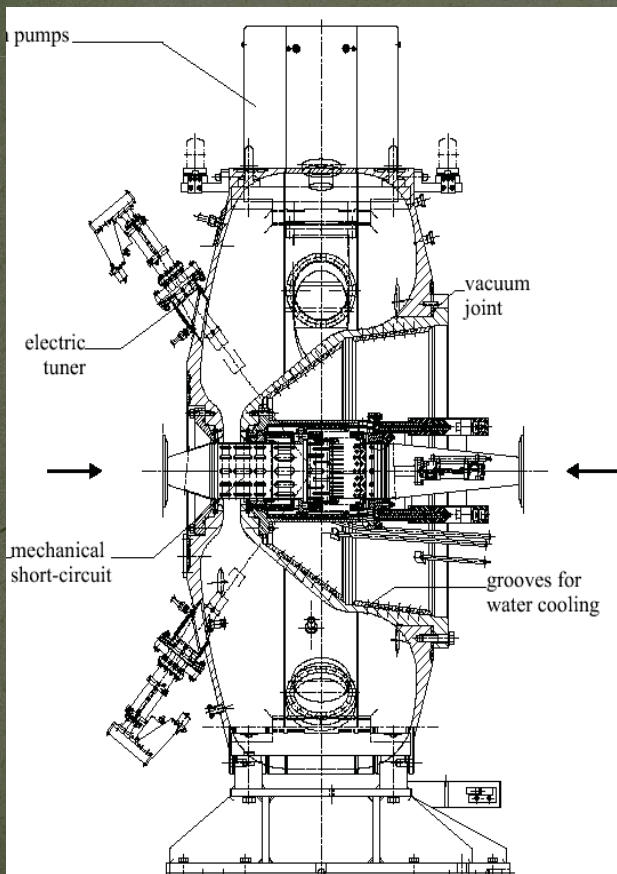
Pure TE modes: No net transverse force !

Transverse modes are characterized by

- the transverse impedance in ω -domain
- the transverse loss factor (kick factor) in t -domain !

W.K.H. Panofsky, W.A. Wenzel: "Some Considerations Concerning the Transverse Deflection of Charged Particles in Radio-Frequency Fields", RSI 27, 1957]

CERN/PS 80 MHz cavity (for LHC)



inductive (loop) coupling, low self-inductance

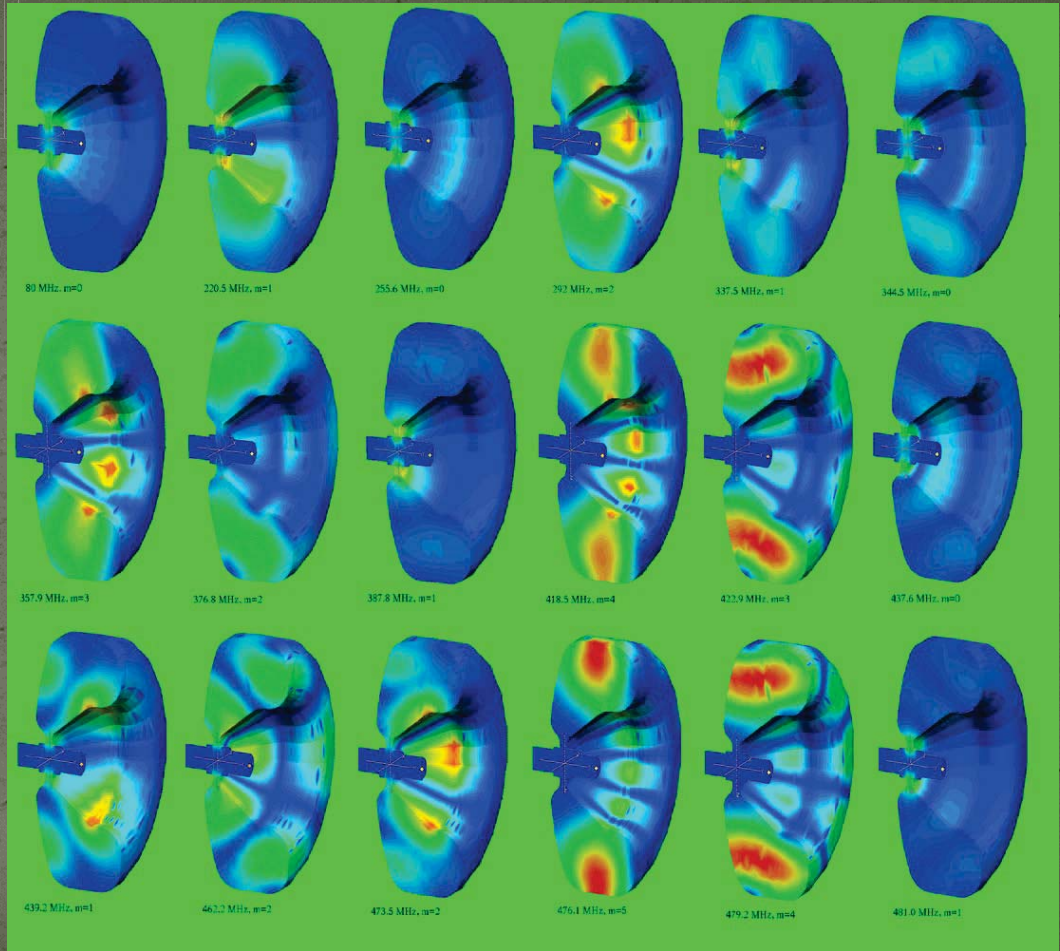
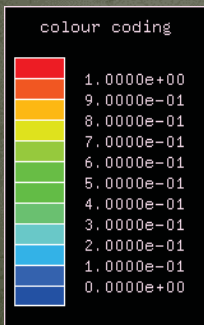
HOM's

Example shown:

80 MHz cavity PS for LHC.

Color-coded:

$$\vec{E}$$



More examples of cavities

PS 19 MHz cavity (prototype, photo: 1966)

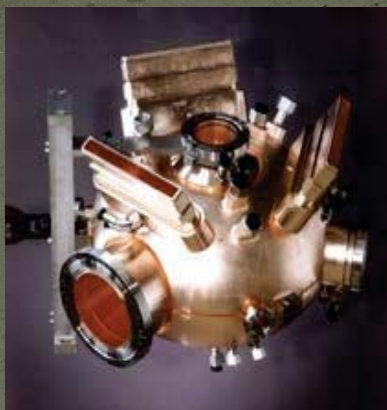


2nd Nov, 2012

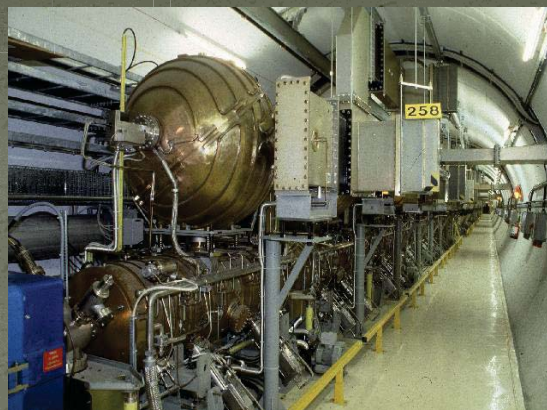
CAS Granada - EJ: RF Systems II

17

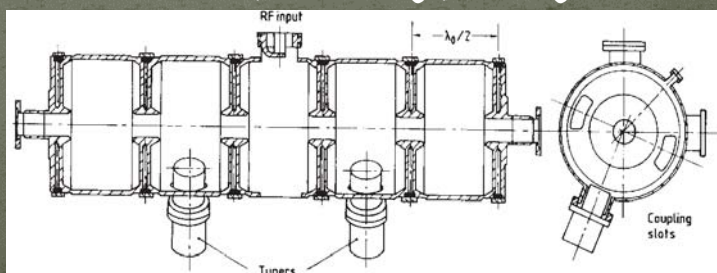
Examples of cavities



PEP II cavity
476 MHz, single cell,
1 MV gap with 150 kW,
strong HOM damping,



LEP normal-conducting Cu RF cavities,
352 MHz. 5 cell standing wave +
spherical cavity for energy storage, 3 MV

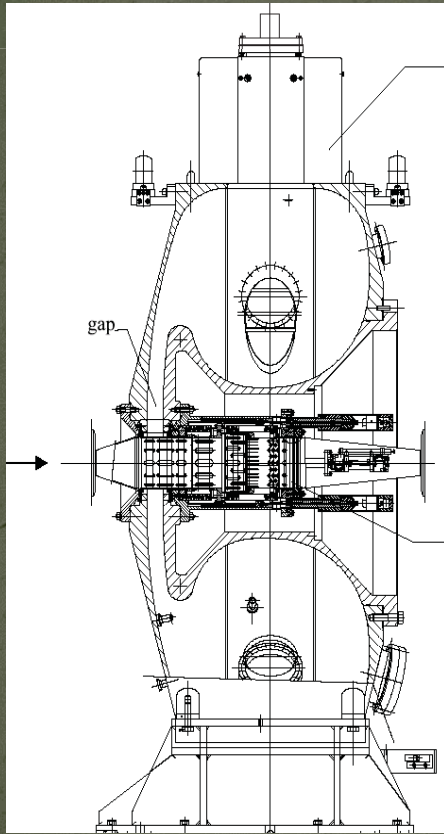


2nd Nov, 2012

CAS Granada - EJ: RF Systems II

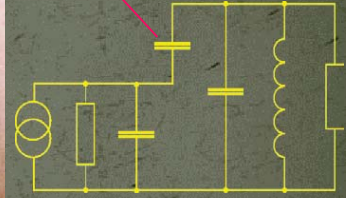
18

CERN/PS 40 MHz cavity (for LHC)

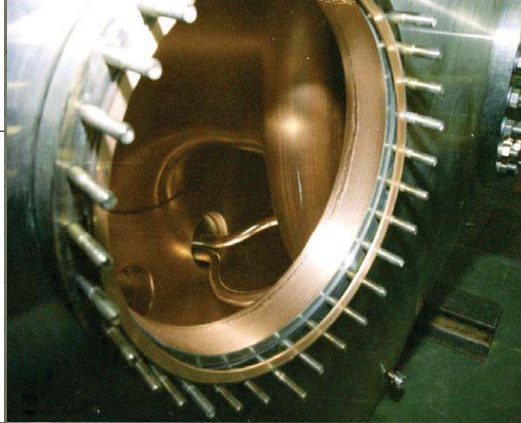


example for
capacitive coupling

coupling C



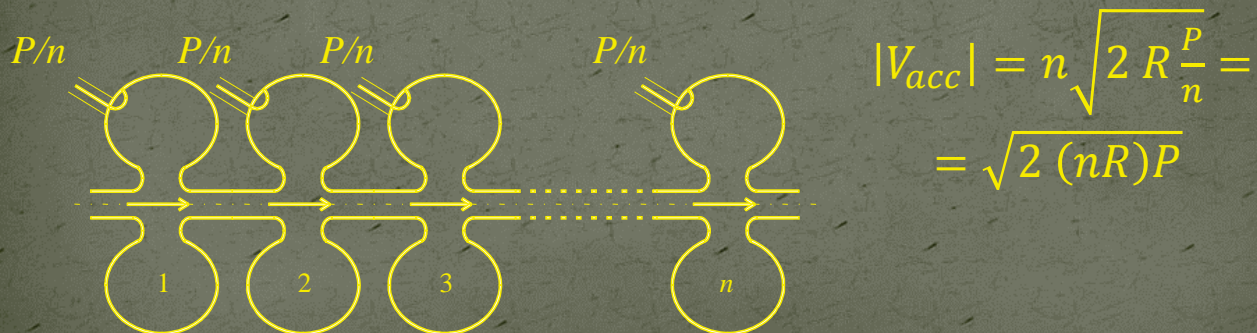
cavity



Many gaps

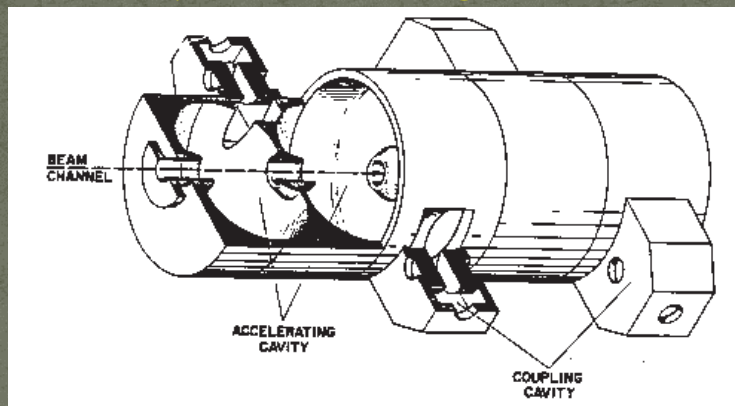
What do you gain with many gaps?

- The R/Q of a single gap cavity is limited to some 100 Ω . Now consider to distribute the available power to n identical cavities: each will receive P/n , thus produce an accelerating voltage of $\sqrt{2RP/n}$. The total accelerating voltage thus increased, equivalent to a total equivalent shunt impedance of nR .



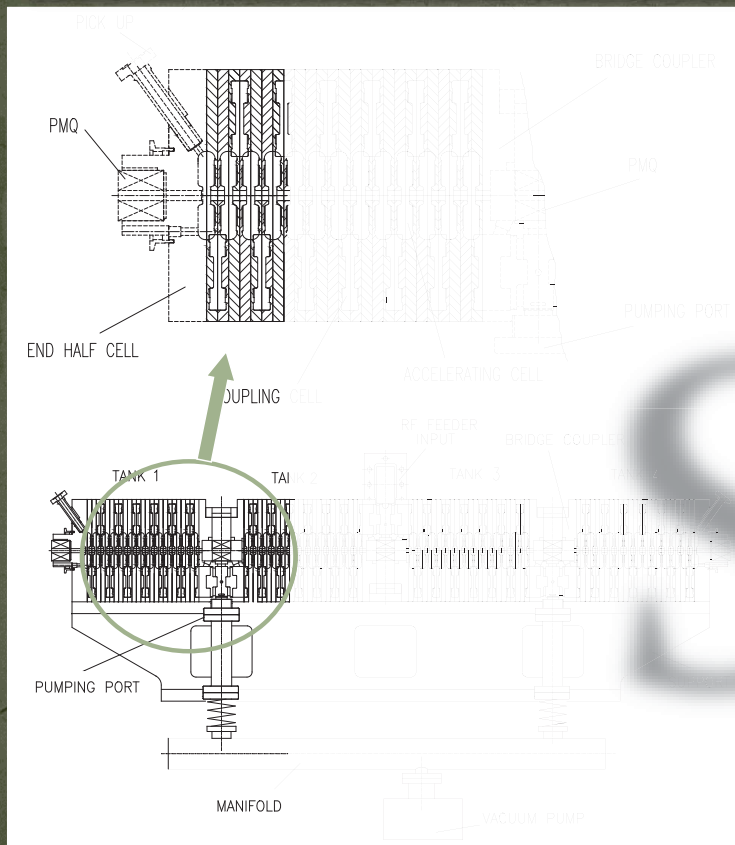
Standing wave multicell cavity

- Instead of distributing the power from the amplifier, one might as well couple the cavities, such that the power automatically distributes, or have a cavity with many gaps (e.g. drift tube linac).
- Coupled cavity accelerating structure (side coupled)



- The phase relation between gaps is important!

Side Coupled Structure : example LIBO



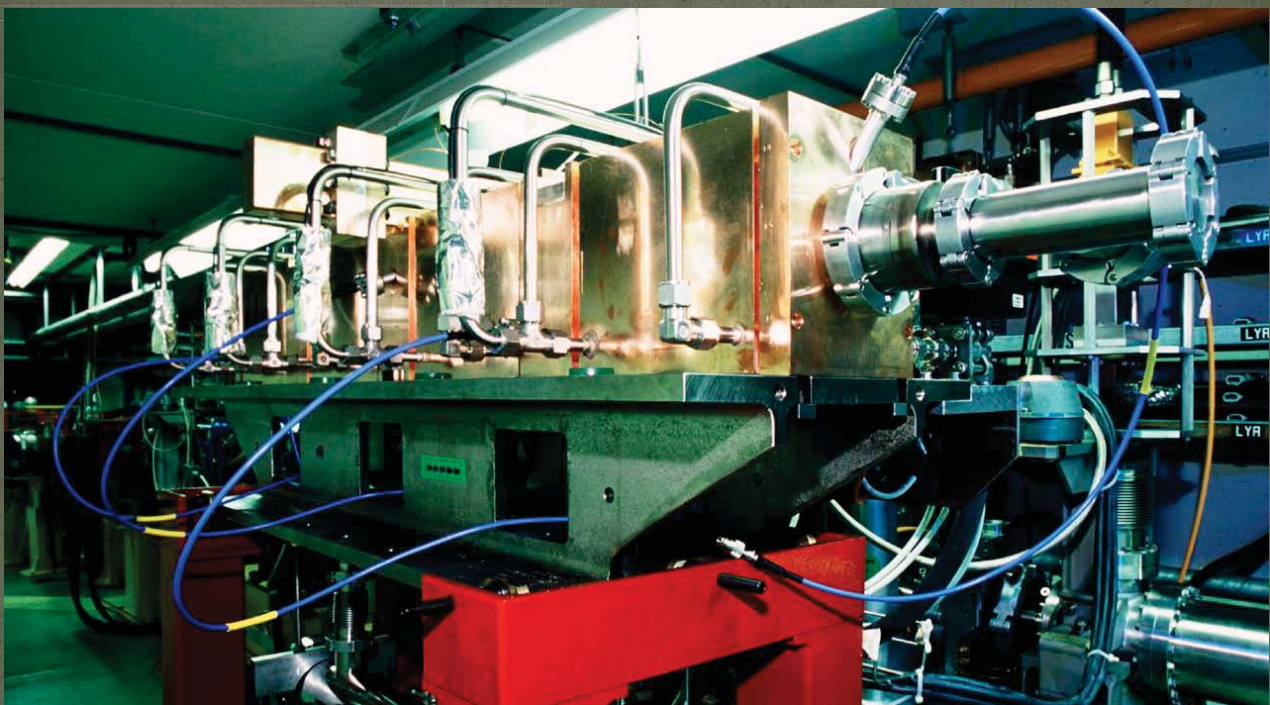
A 3 GHz Side Coupled Structure to accelerate protons out of cyclotrons from 62 MeV to 200 MeV

Medical application: treatment of tumours.

Prototype of Module 1 built at CERN (2000)

Collaboration CERN/INFN/Tera Foundation

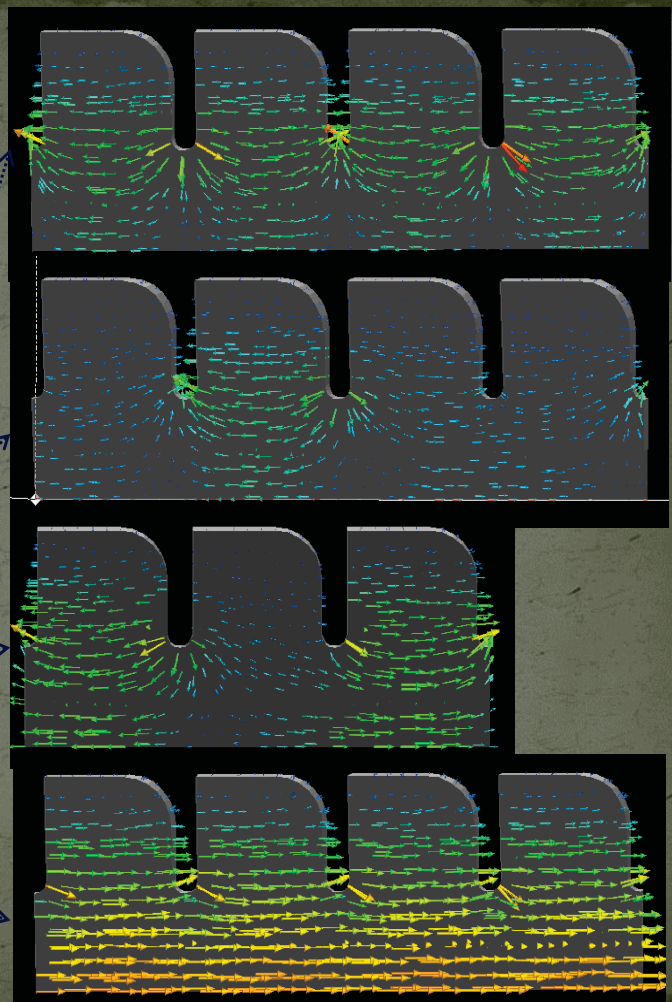
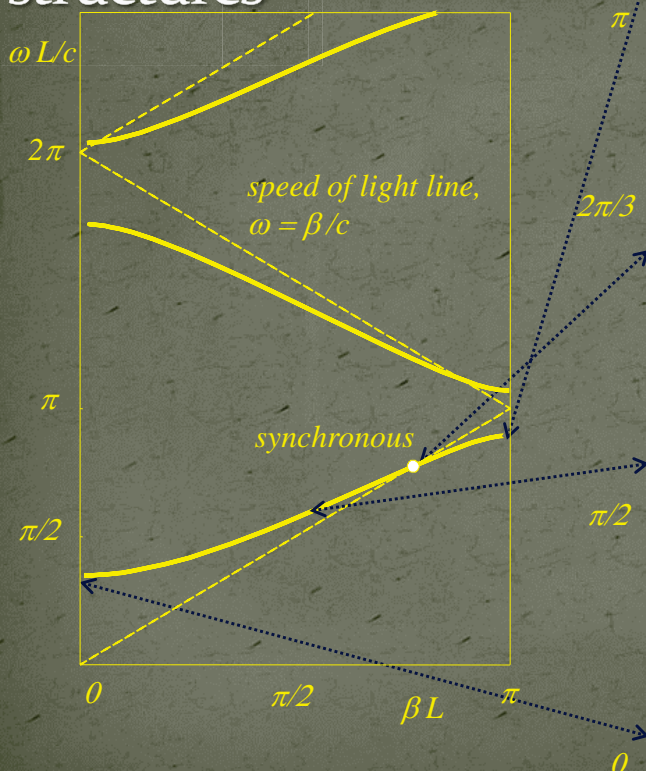
LIBO prototype



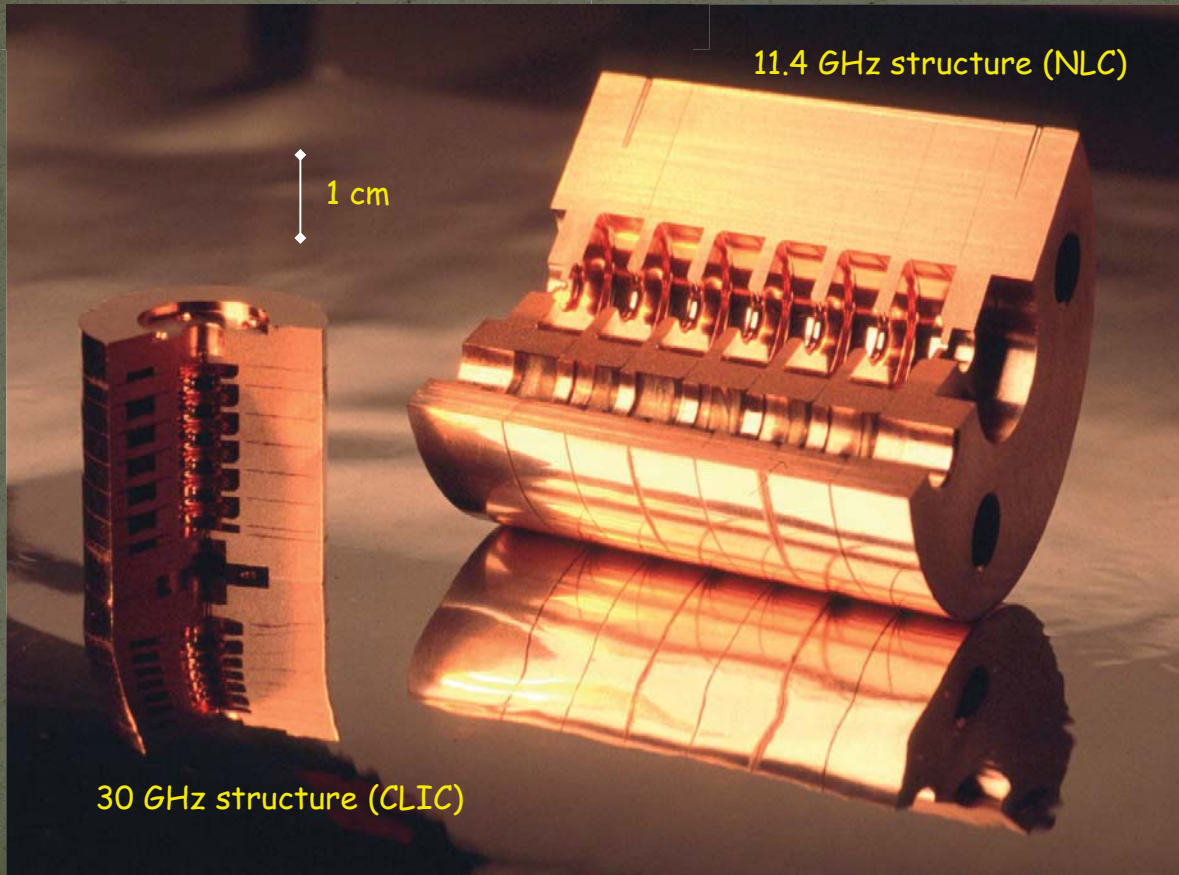
This Picture made it to the title page of CERN Courier vol. 41 No. 1 (Jan./Feb. 2001)

Travelling wave structures

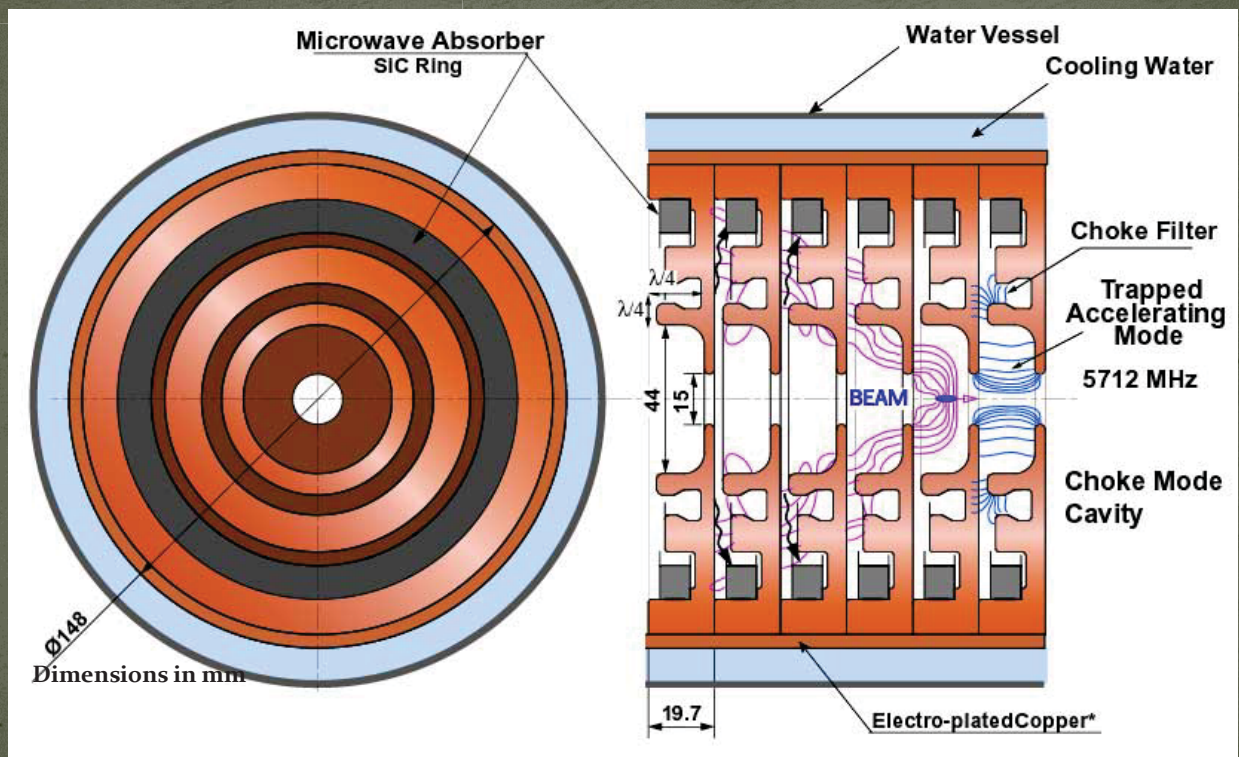
Brillouin diagram – Travelling wave structures



Iris loaded waveguide

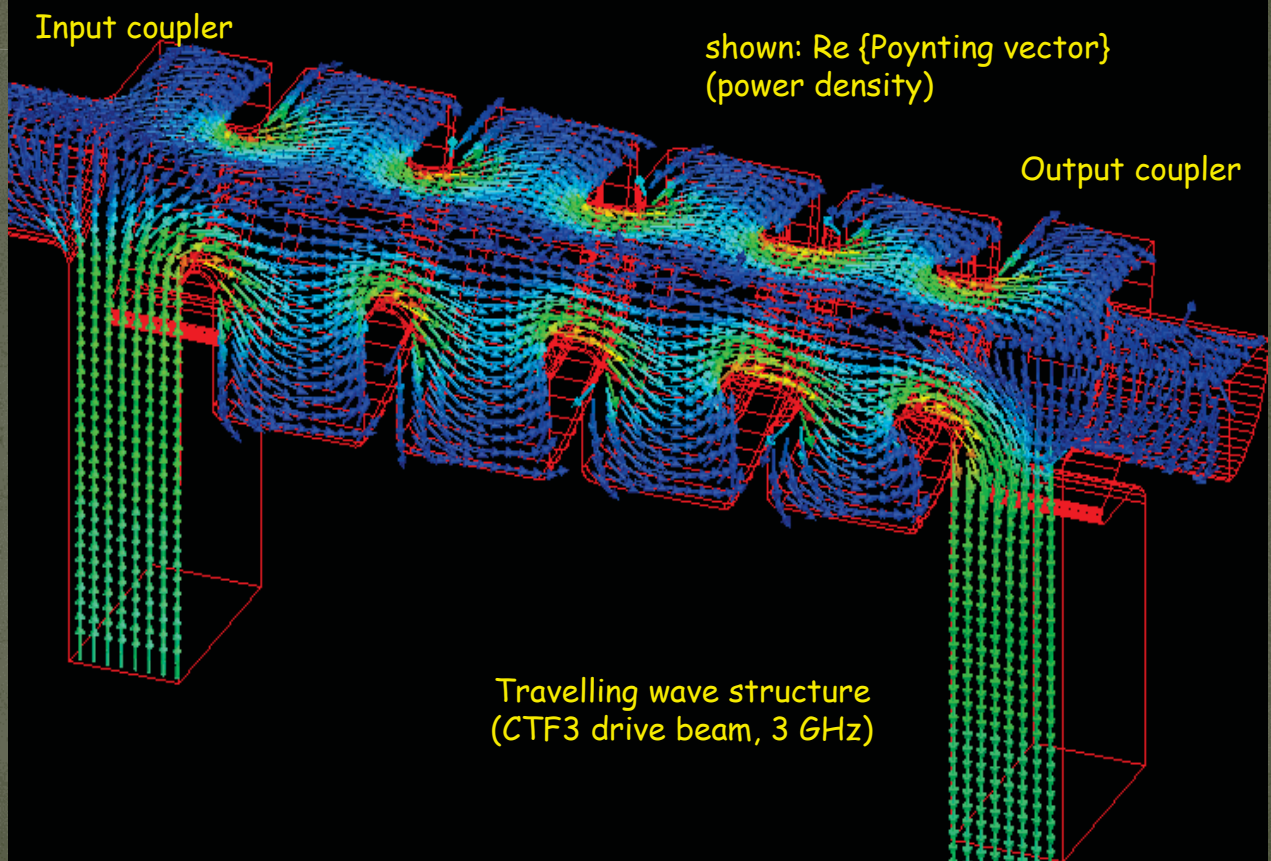


Disc loaded structure with strong HOM damping “choke mode cavity”

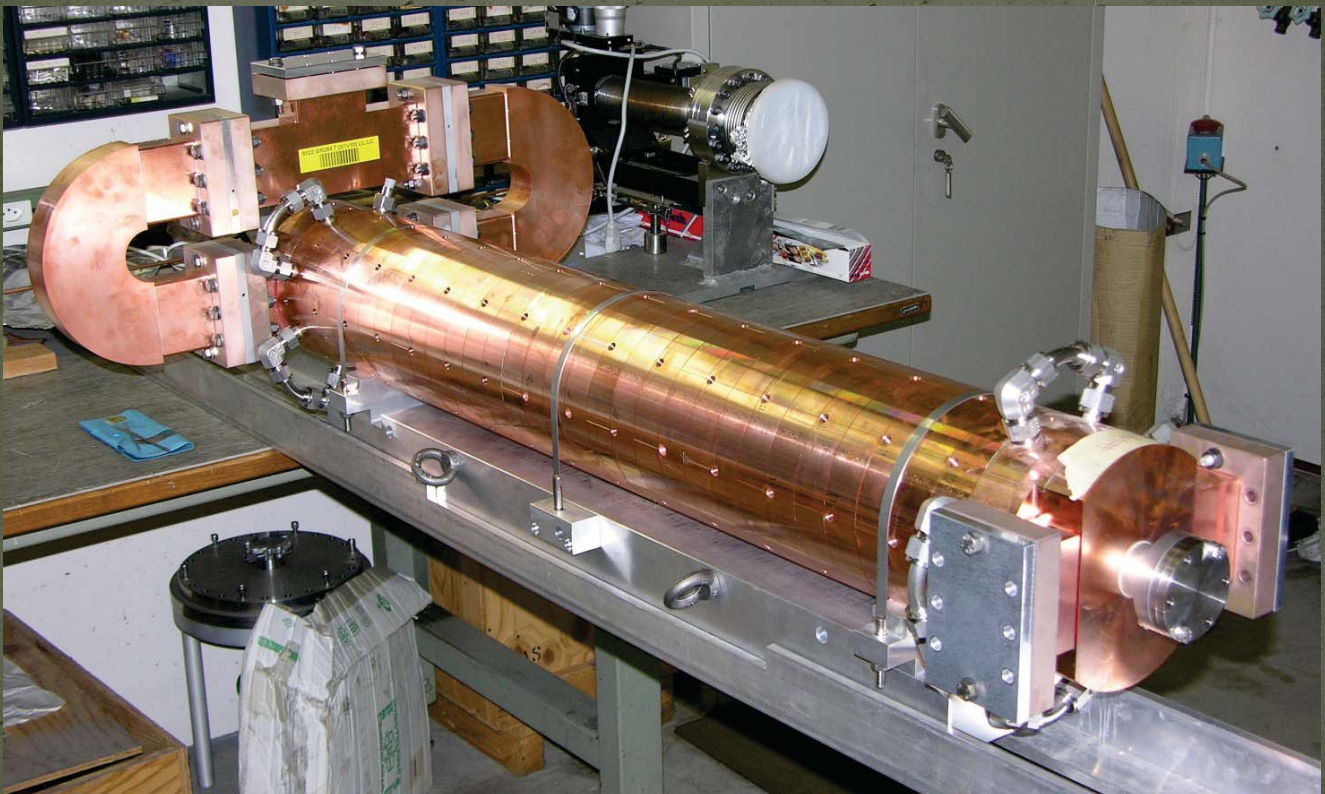


Waveguide coupling

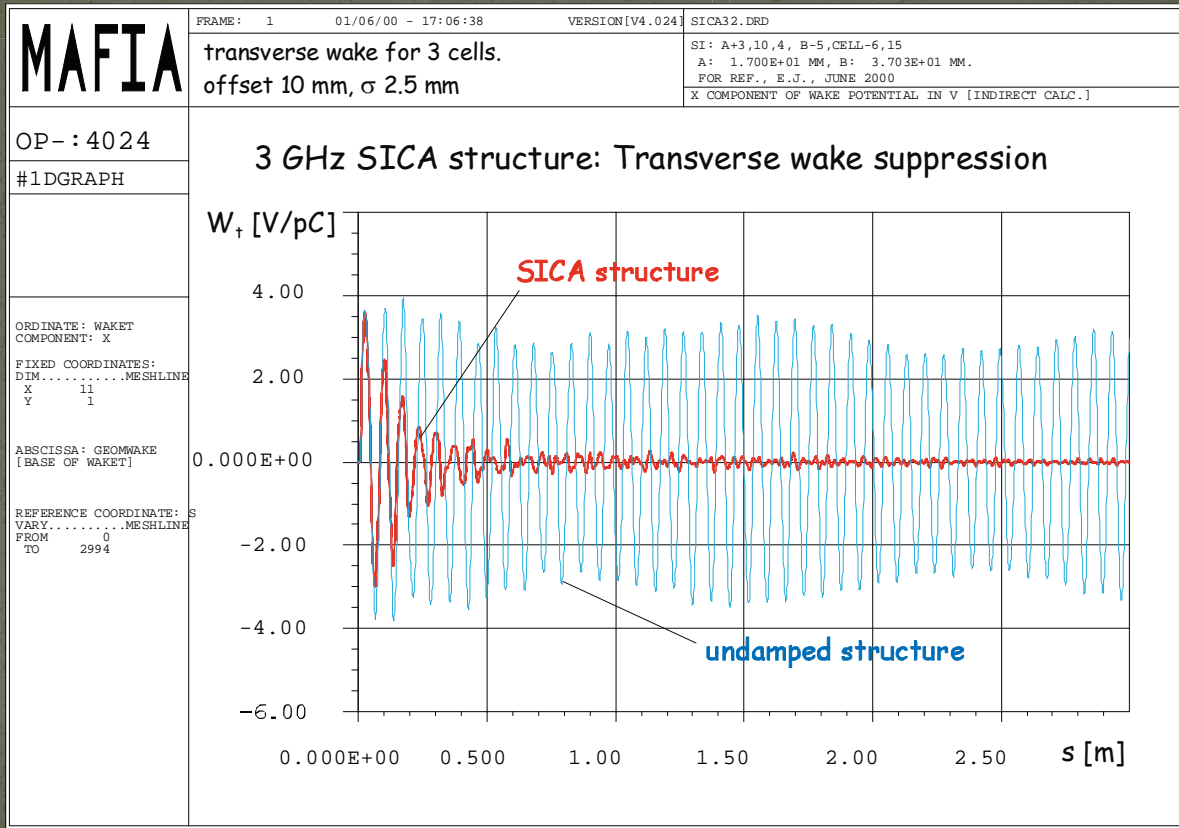
$\frac{1}{4}$ geometry shown



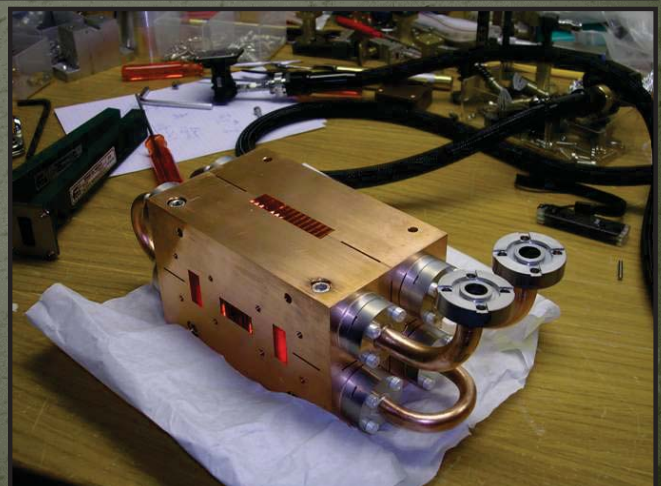
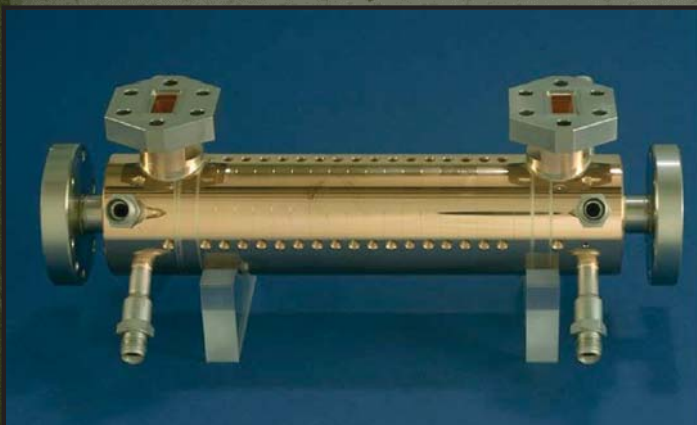
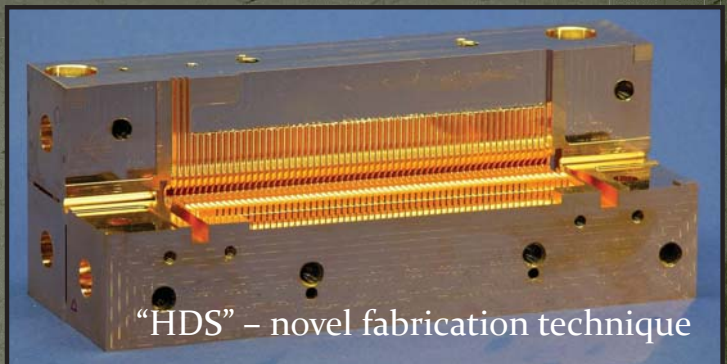
3 GHz Accelerating structure (CTF₃)



HOM damping at work



Recent CLIC structures (11.4, 12 and 30 GHz)



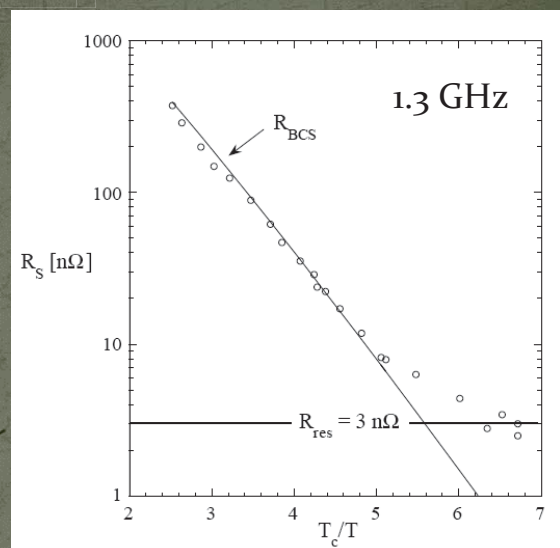
Superconducting Linacs

RF Superconductivity

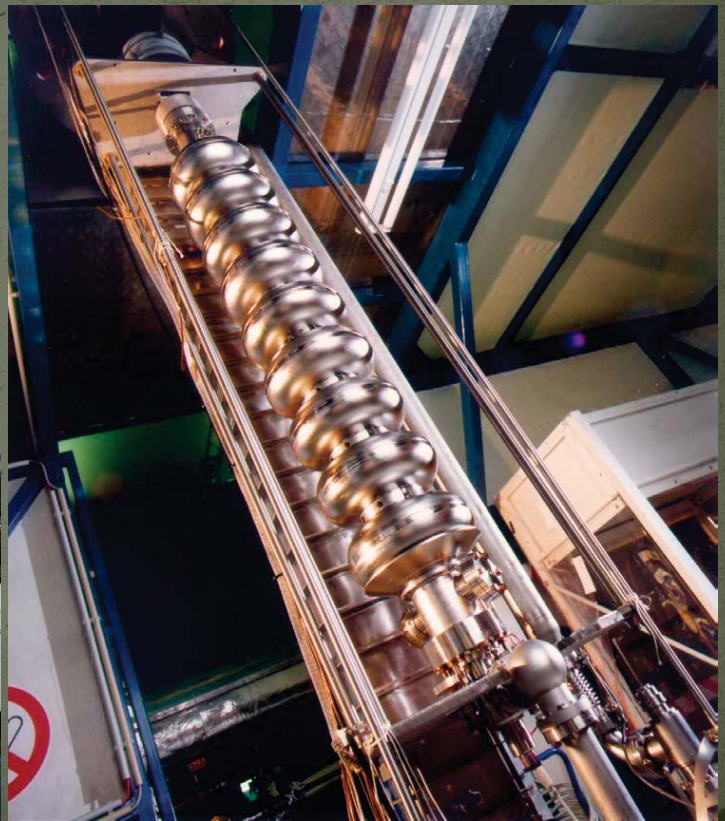
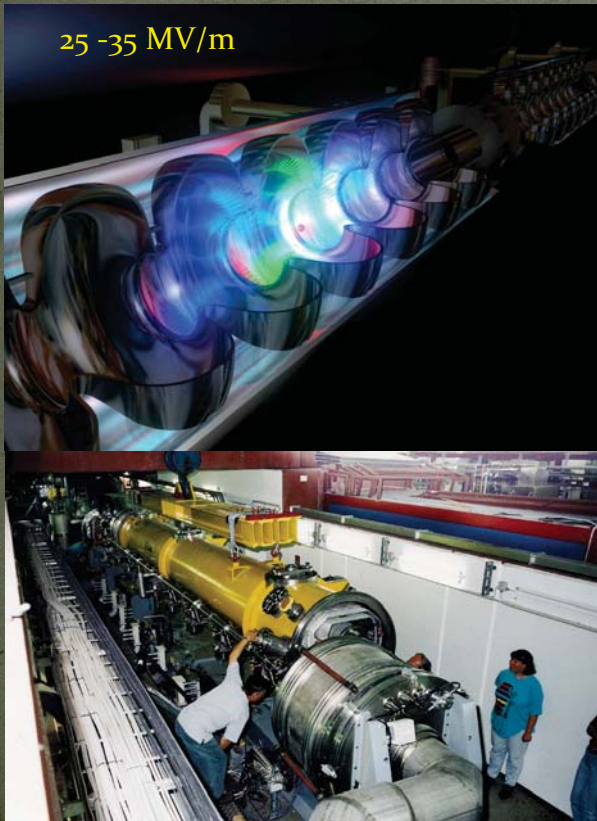
- Different from DC, at RF the resistance is not exactly zero, but just very small. It is

$$R_{surf} = R_{BCS} + R_{res} \quad R_{BCS} \propto \omega^2 e^{-1.76T_c/T}$$

- The maximum accelerating gradient is normally limited by the maximum possible surface magnetic field (the “superheating field”, 180 mT for Nb, 400 mT for Nb₃Sn).
- Maximum acc. gradients are however obtained for Nb (ILC, ≈ 40 MV/m).



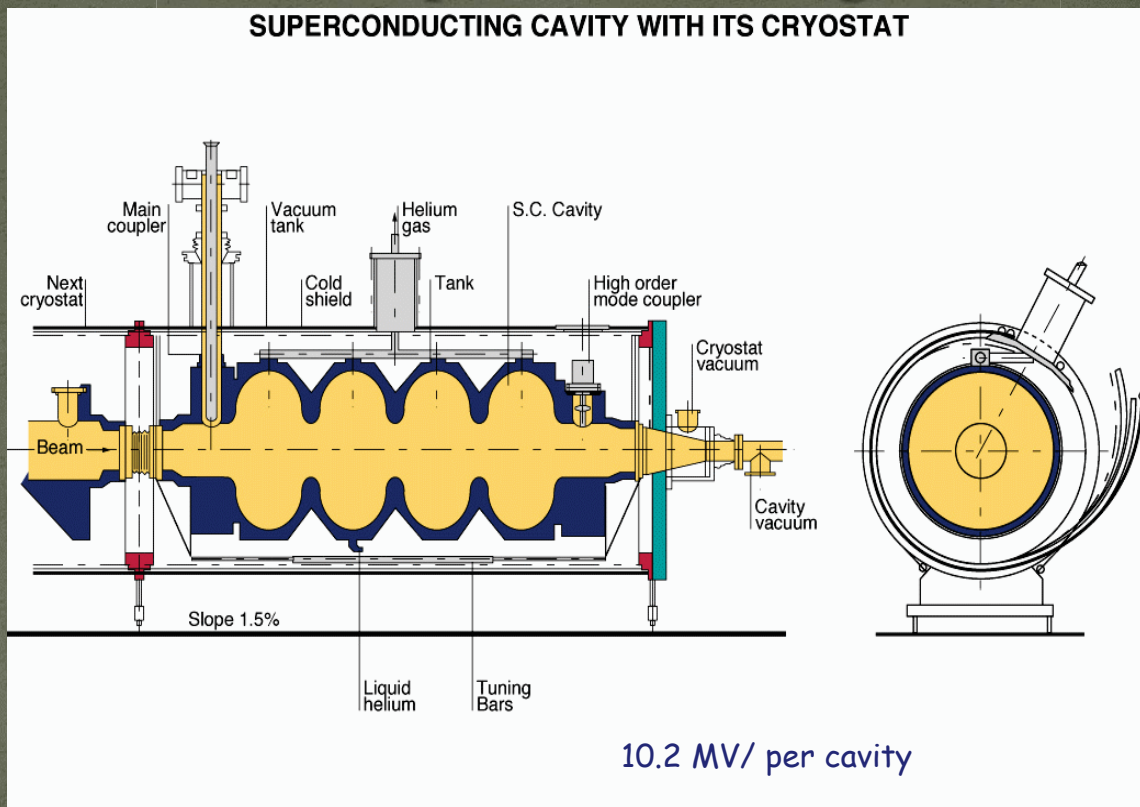
TESLA/ILC/X-FEL SC cavities, 1.3 GHz



CAS Granada - EJ: RF Systems II

LEP Superconducting cavities

SUPERCONDUCTING CAVITY WITH ITS CRYOSTAT



Nb coating techniques

- Sputtering Nb on Cu
 - Advantages:
 - Due to the high cost of Nb, this can reduce cost!
 - The Cu substrate increases the mechanical & thermal stability (quench resistance).
 - Technology initially developed at CERN (Benvenuti, LEP, 1980); experts today at JLAB, Legnaro, Saclay, Sheffield & CERN
 - Technique used today for ALPI (LNL), Soleil, LHC & HIE-Isolde
 - Today, the max. fields are still smaller than for bulk Nb – is this an intrinsic limitation? An interesting field of R&D!
 - Can this technique be extended to new materials? (NbTiN, V₃Si, Nb₃Sn, HTS?)
- “Energetic Condensation” - HiPIMS
 - Gas phase deposition of Nb with additional kinetic energy to slow ions.
- Cathodic Arc Deposition
- ...

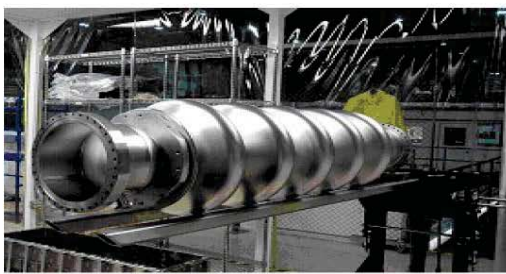
2nd Nov, 2012

CAS Granada - EJ: RF Systems II

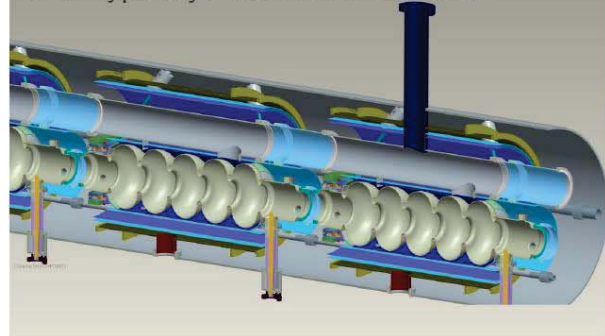
37

704 MHz cavities and cryomodule

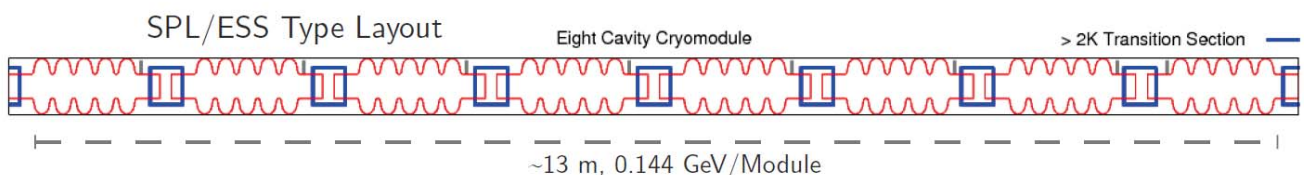
ESS, eRHIC, SPL



SNS type cryomodule for 700 MHz



5-cell cavities (1.6 m long), 8 per cryomodule



2nd Nov, 2012

CAS Granada - EJ: RF Systems II

38

LHC SC RF, 4 cavity module, 400 MHz



2nd Nov, 2012

CAS Granada - EJ: RF Systems II

39

Energy Recovery Linac

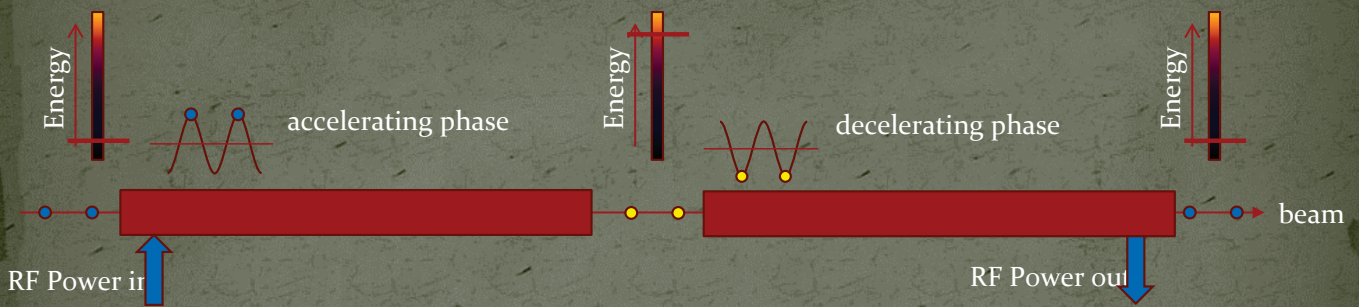
How to reach “power grid \rightarrow beam” efficiencies above 100%

2nd Nov, 2012

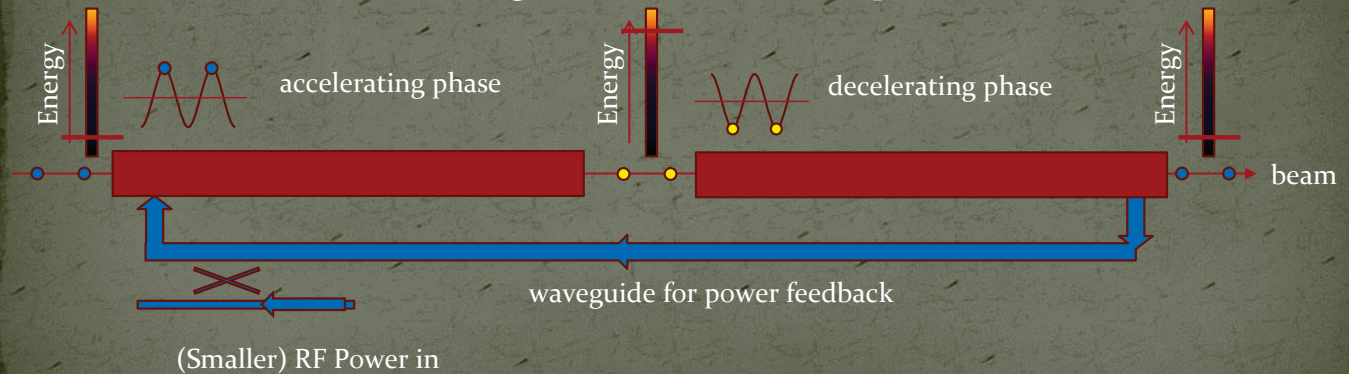
CAS Granada - EJ: RF Systems II

40

Recovering the energy from the beam – the concept



One could use a waveguide and reuse the RF power!

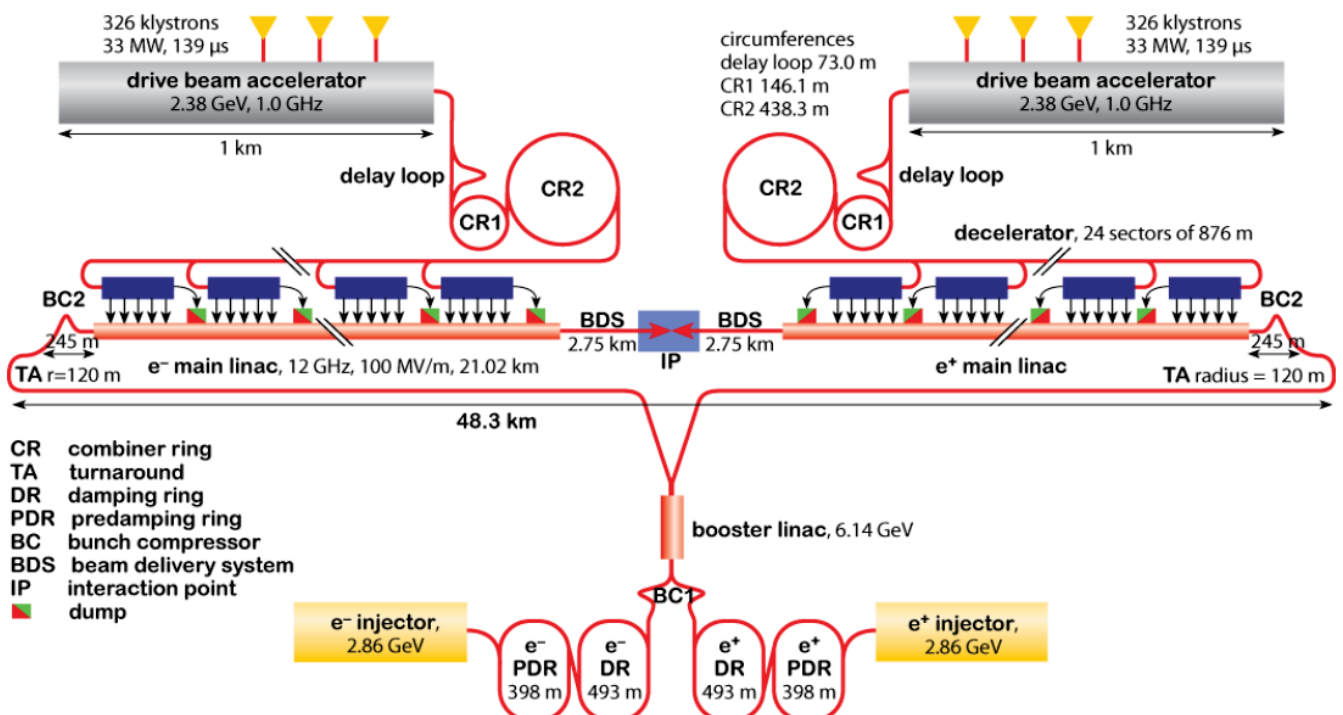


A word about CLIC

In the CLIC scheme, 90% of the drive beam power is recovered (to produce the RF power for the main beam)

CLIC general layout

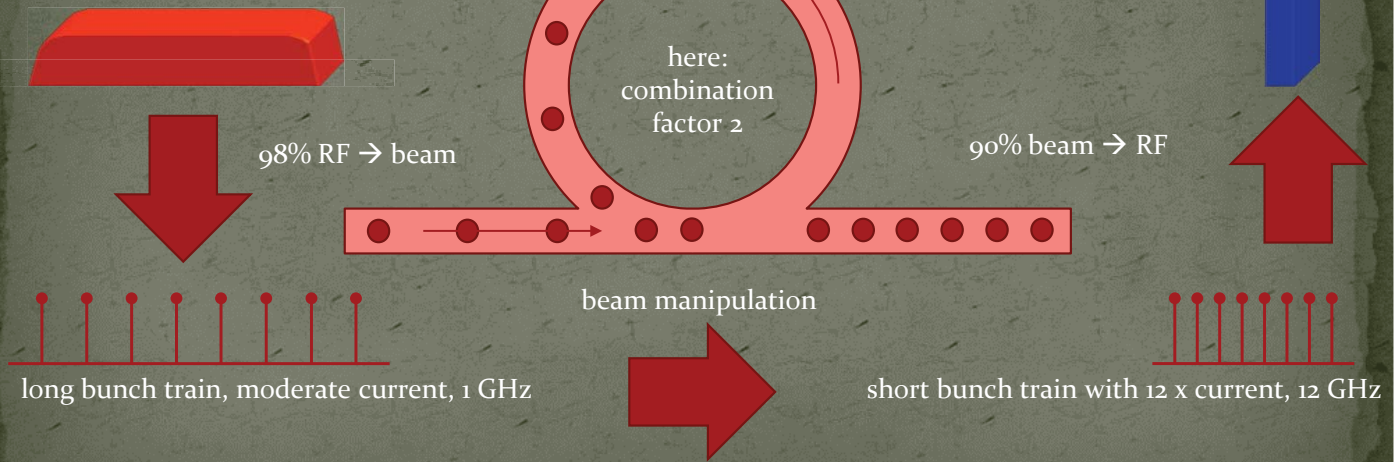
<http://clic-study.org/accelerator/CLIC-inaNutshell.php>



The CLIC power source idea

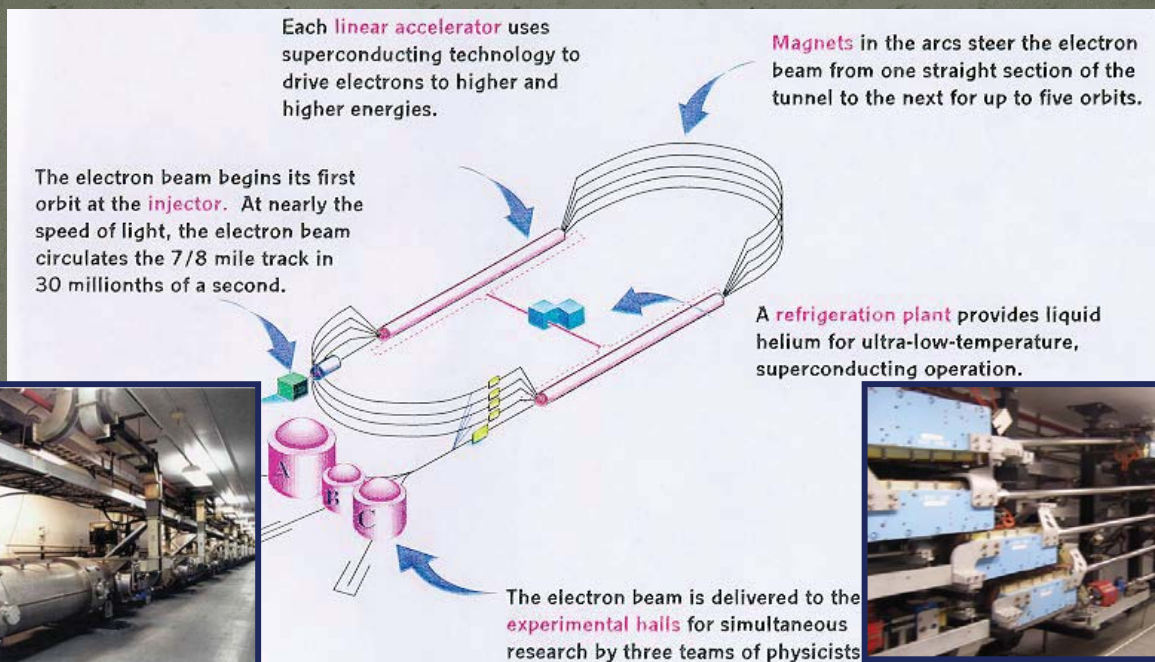
Long RF Pulses
 P_o, f_o, τ_o
 to accelerate the drive beam

Short RF Pulses
 $P_A = P_o \times N$
 $\tau_A = \tau_o / N$
 $f_A = f_o \times N$
 extracted from
 recombined drive beam



Recirculating Linac

- One could use the same accelerating structure more than once!
- CEBAF (Continuous Electron Beam Accelerator Facility) at JLAB, Newport News, VA, USA has been using this scheme successfully for many years.



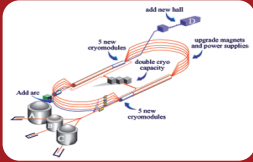
Recirculating Linac compared to linac and synchrotron

Linac



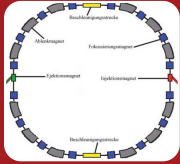
- Accelerating Structure used for 1 passage
- Less efficient
- Only single pass instabilities

Recirculating Linac



- Accelerating Structure used for some (2-10) passages
- Return arcs different for different energies
- Concerning instabilities, a good compromise

Synchrotron

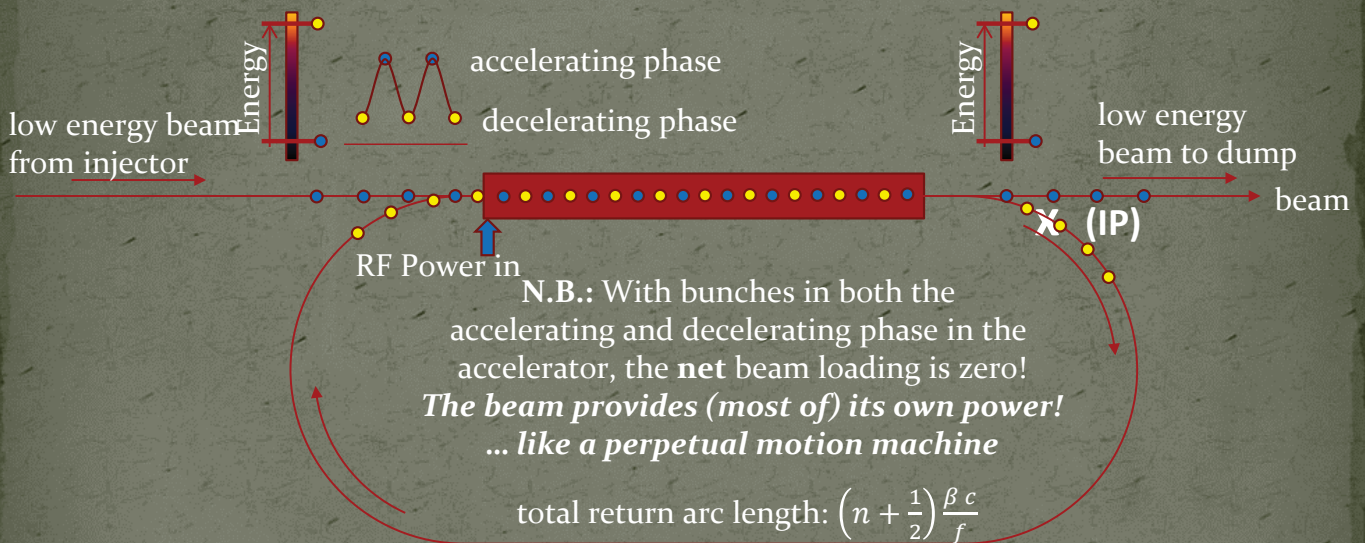


- Accelerator Structure used many times
- Periodic lattice
- Instabilities develop over many turns (coupled bunch, mode coupling)

L. Merminga '07: In a storage ring, electrons are stored for hours in an equilibrium state, whereas in an ERL it is the energy of the electrons that is stored. The electrons themselves spend little time in the accelerator (from ~1 to 10's of μ s) thus never reach equilibrium. As a result, in common with linacs, the 6-dimensional phase space in ERLs is largely determined by the electron source properties by design. On the other hand, in common with storage rings, ERLs have high current carrying capability enabled by the energy recovery process, thus promising high efficiencies. <http://accelconf.web.cern.ch/AccelConf/p07/PAPERS/MOYK103.PDF>

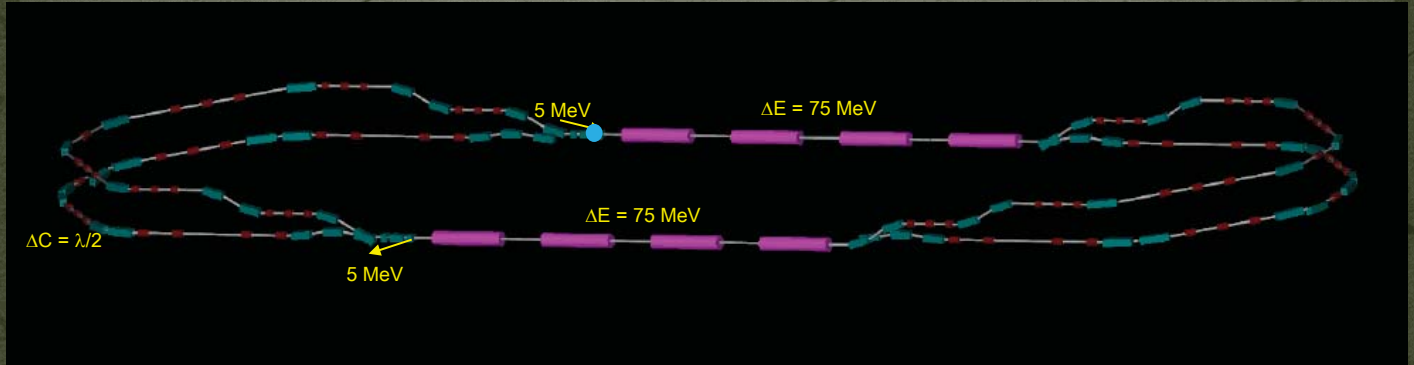
Energy Recovery Linac:

Combine "Energy recovery" and "recirculating"



LHeC ERL-TF (300 MeV) – Layout

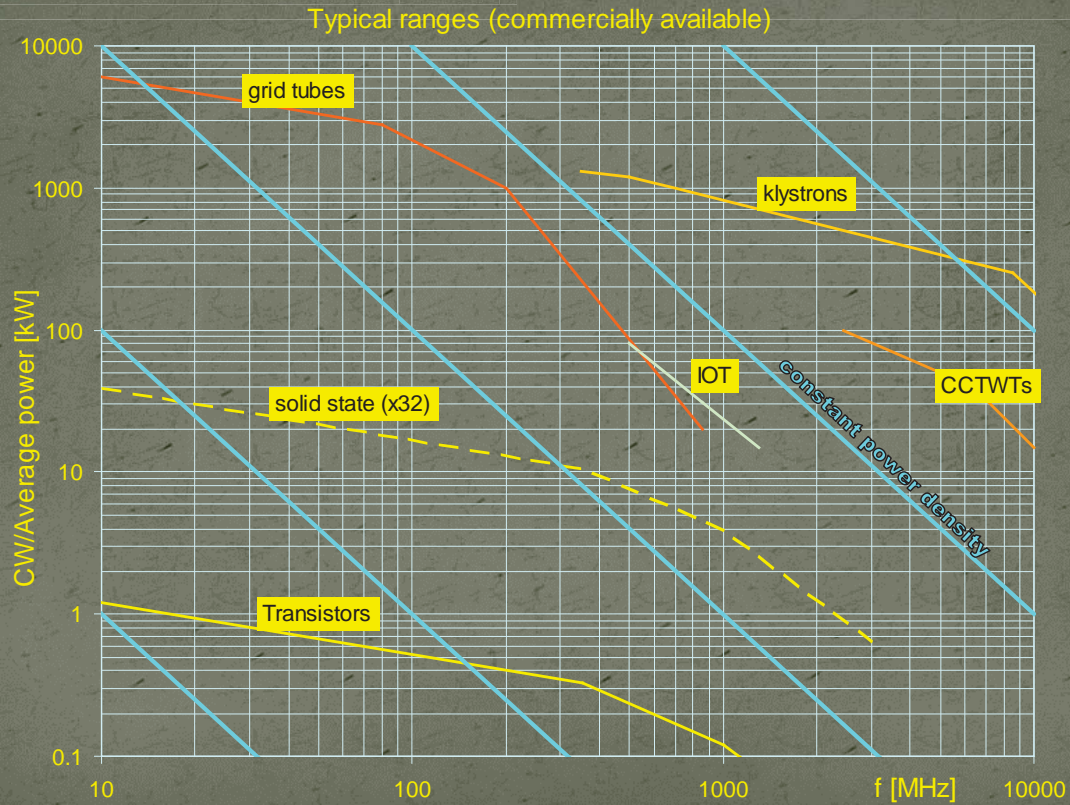
This model and animation by Alex Bogacz, Jefferson Lab



Two passes 'up' + Two passes 'down'

RF power sources

RF power sources



LEIR SSPA, 1 kW, 0.2 - 50 MHz

Taken from M. Paoluzzi, LEIR RF system



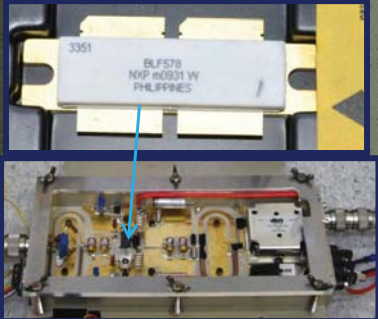
MRF151G

Soleil/ESRF Booster SSPA, 150 kW, 352 MHz

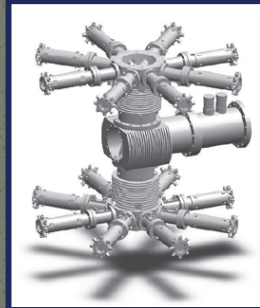
Taken from J. Jacob, CWRF 2012

- Initially developed by SOLEIL
- Transfer of technology to ELTA / AREVA

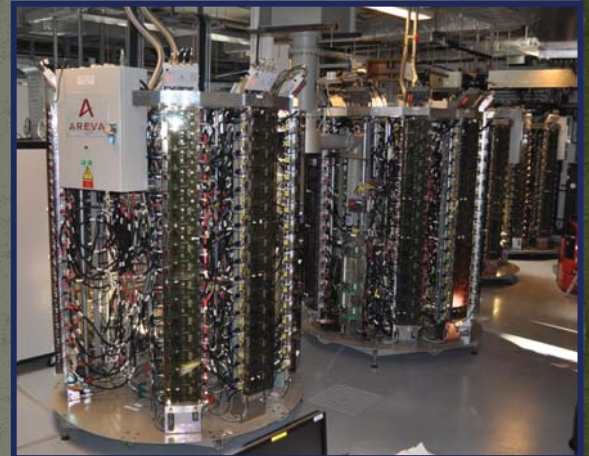
Pair of push-pull transistors



X 128



X 2



650 W RF module

- 6th generation LDMOSFET (BLF 578 / NXP), $V_{ds} = 50$ V
- Efficiency: 68 to 70 %

75 kW Coaxial combiner tree

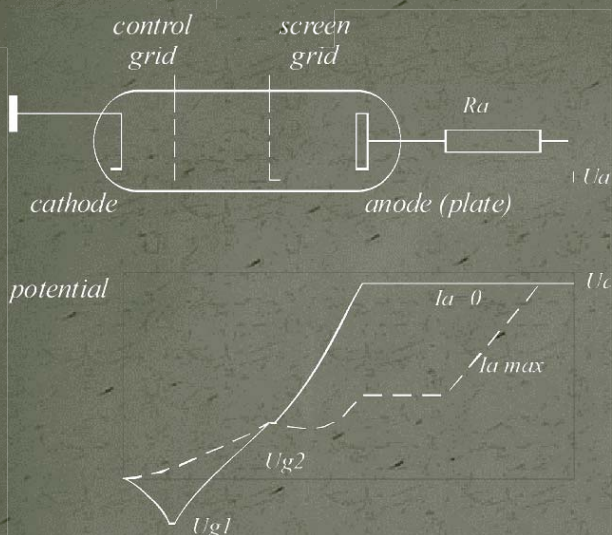
with $\lambda/4$ transformers

150 kW - 352.2 MHz Solid State Amplifiers for the ESRF booster

Efficiency: > 55 % at nominal power

- 1st batch of 4 x 150 kW SSAs from ELTA in operation on ESRF booster since March 2012
- 2nd batch of 3 x 150 kW SSAs in fabrication, will power 3 new cavities on ESRF storage ring

Tetrode



4CX250B (Eimac/CPI), < 500 MHz, 600 W (Anode removed)



RS 1084 CJ (ex Siemens, now Thales), < 30 MHz, 75 kW

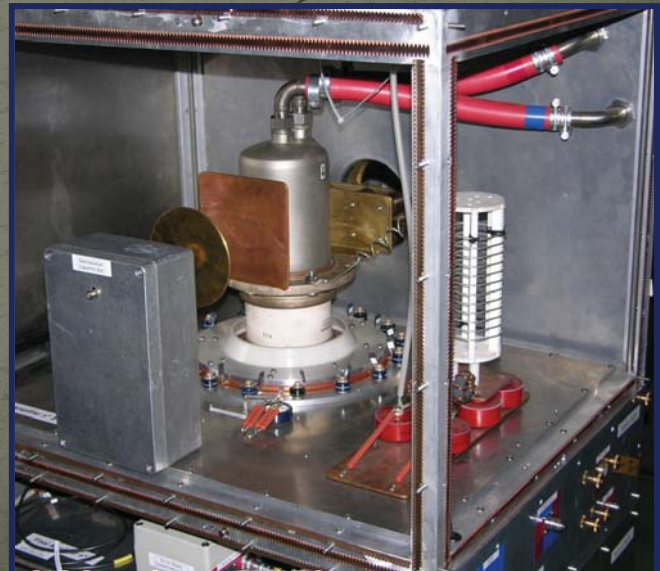
YL1520 (ex Philips, now Richardson), < 260 MHz, 25 kW

High power tetrode amplifier



CERN Linac3: 100 MHz, 350 kW
50 kW Driver: TH345, Final: RS 2054 SK

CERN PS: 13-20 MHz, 30 kW
Driver: solid state 400 W, Final: RS 1084 CJSC

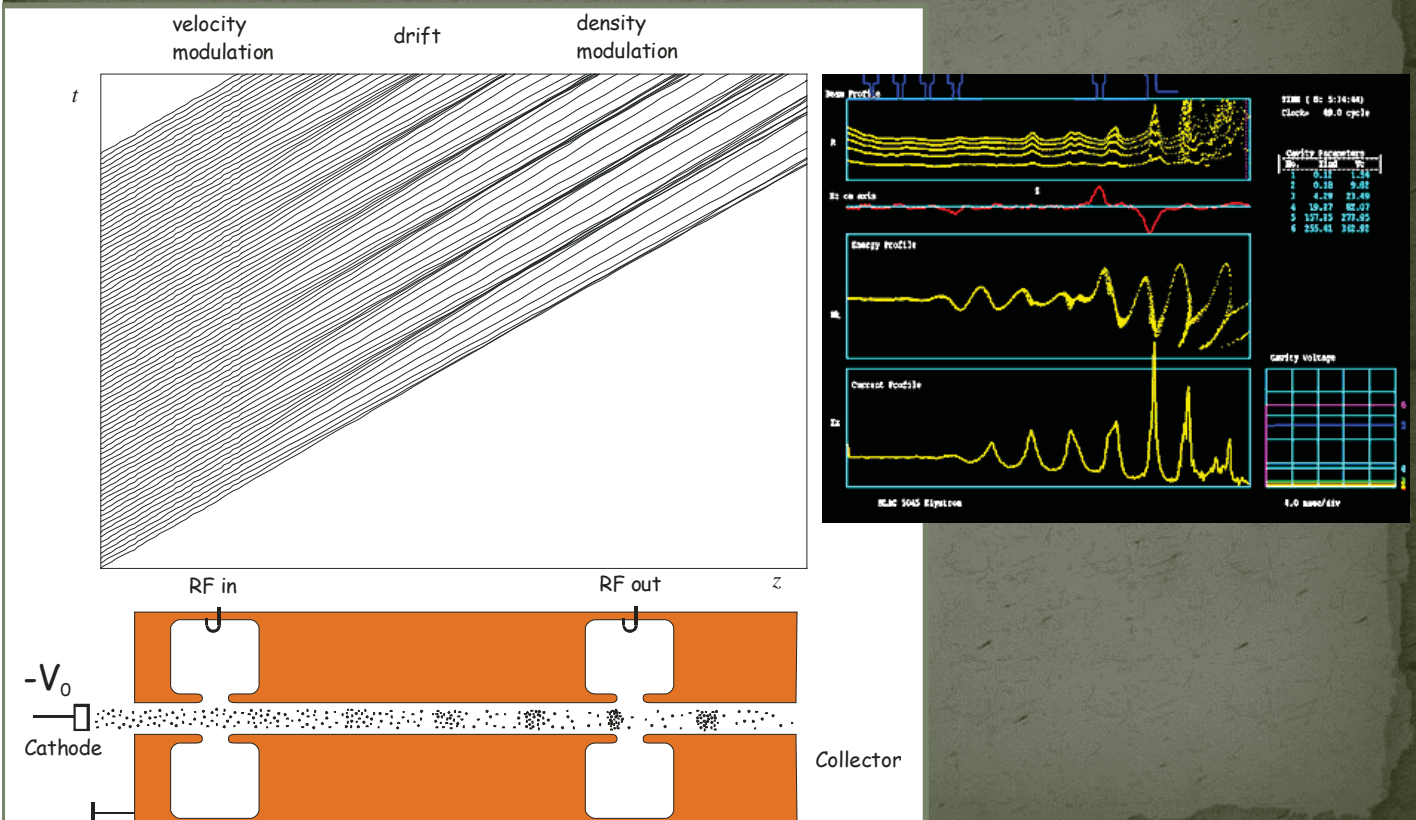


CAS Granada - EJ: RF Systems II

2nd Nov, 2012

53

Klystron principle



2nd Nov, 2012

CAS Granada - EJ: RF Systems II

54

Klystrons



CERN CTF3 (LIL):
3 GHz, 45 MW,
4.5 μ s, 50 Hz, η 45 %



CERN LHC:
400 MHz, 300 kW,
CW, η 62 %

2nd Nov, 2012

CAS Granada - EJ: RF Systems II

55

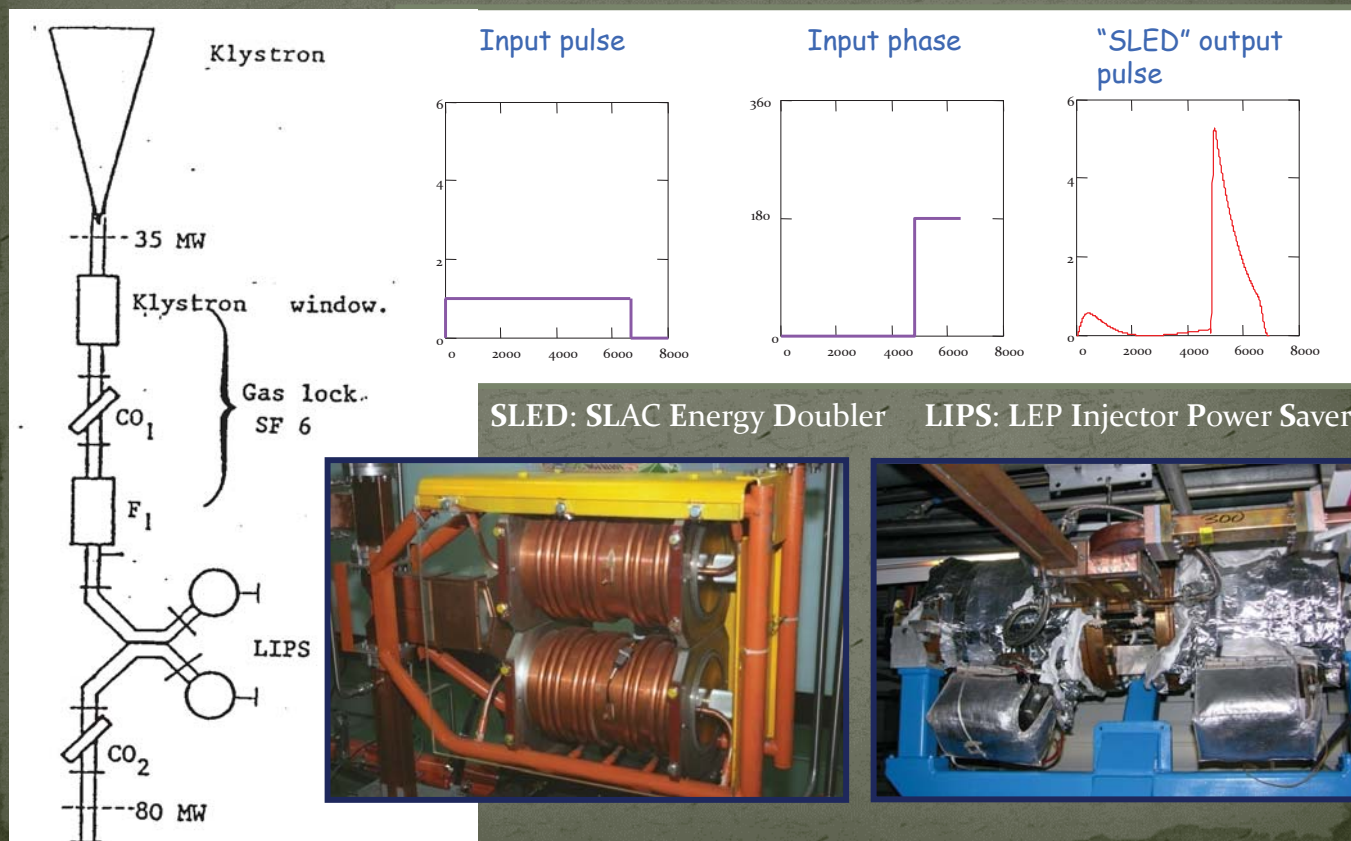
RF pulse compression

2nd Nov, 2012

CAS Granada - EJ: RF Systems II

56

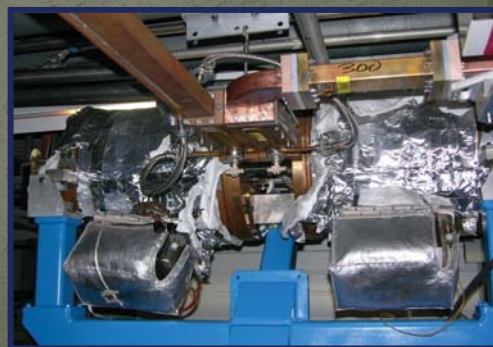
RF Pulse Compression



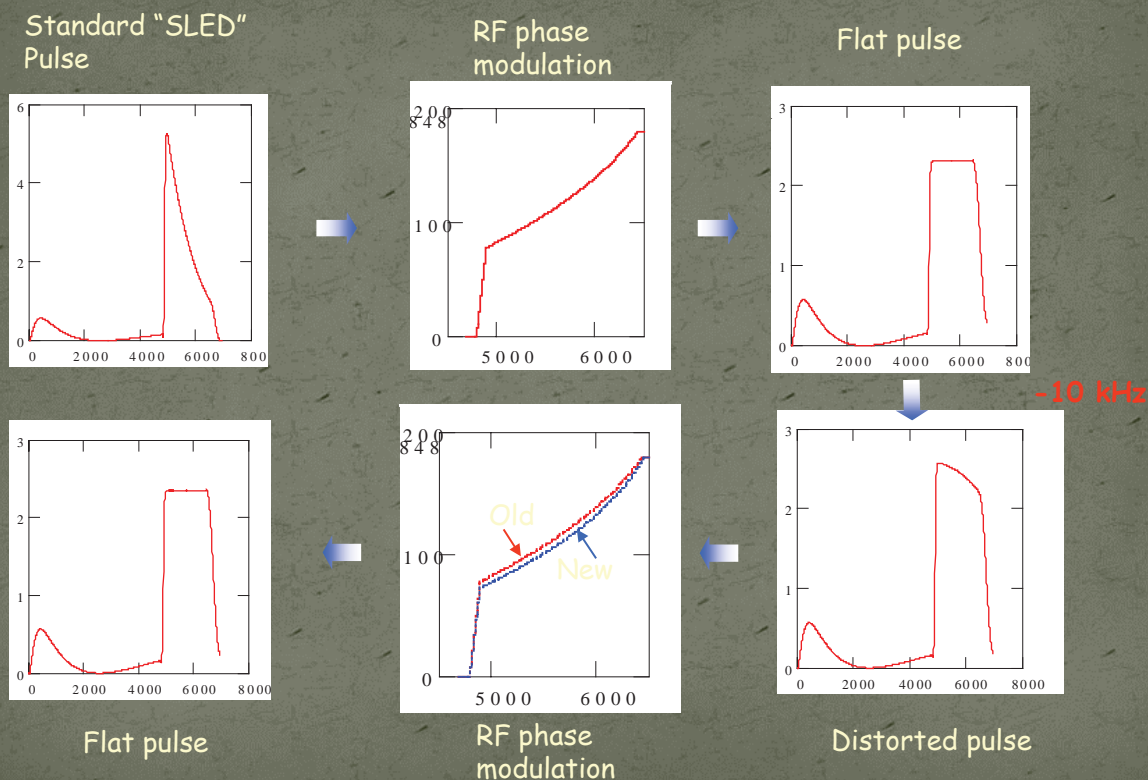
2nd Nov, 2012

CAS Granada - EJ: RF Systems II

57



Flat output pulses



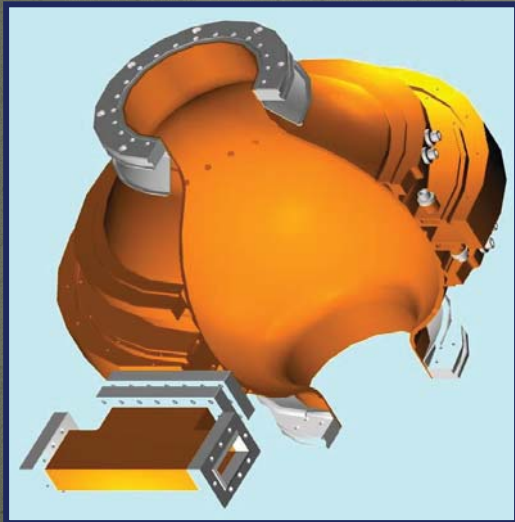
2nd Nov, 2012

CAS Granada - EJ: RF Systems II

58

Pulse compressor

BOC „Barrel Open Cavity“



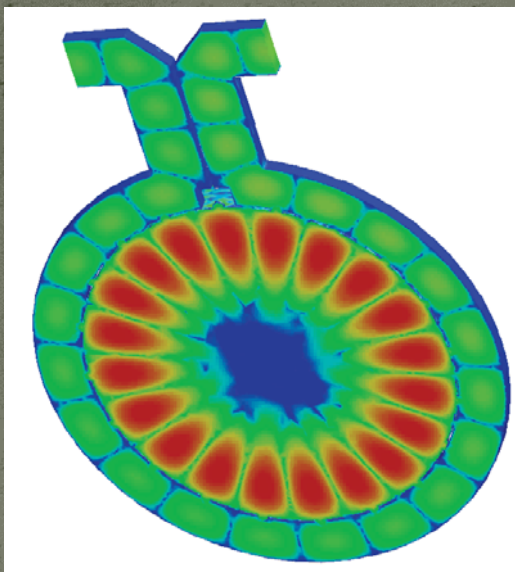
2nd Nov, 2012

CAS Granada - EJ: RF Systems II

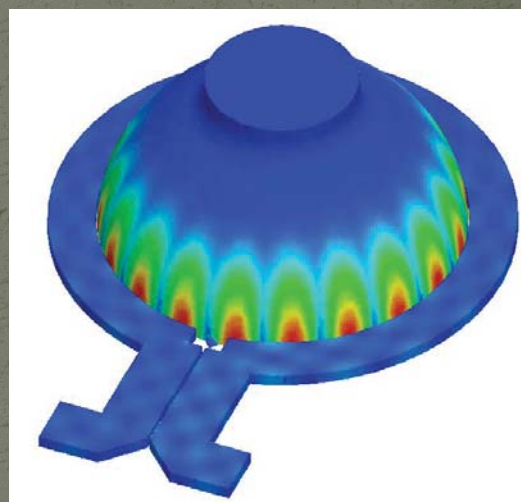
59

BOC

2.99848 GHz,
S11: -12.9 dB



Electric field, logarithmic scale



Magnetic field

2nd Nov, 2012

CAS Granada - EJ: RF Systems II

60

Some examples

Some Power RF systems at CERN

2nd Nov, 2012

CAS Granada - EJ: RF Systems II

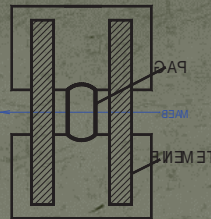
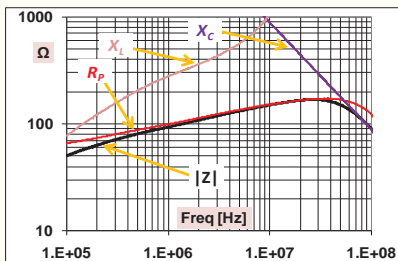
61

Finemet[®] based wide-band cavity

Finemet exhibits wideband response

C_p mostly depends on geometry and drives the high frequency response. The capacitive effect is enhanced by the final stage output capacitance.

R_p and L_p drive the low frequency response. They are mostly dependent on Finemet[®] Characteristics.



On-going project for CERN PSB upgrade
Instantaneous bandwidth:
0.6 MHz ... 4 MHz, 0.7 kV per gap

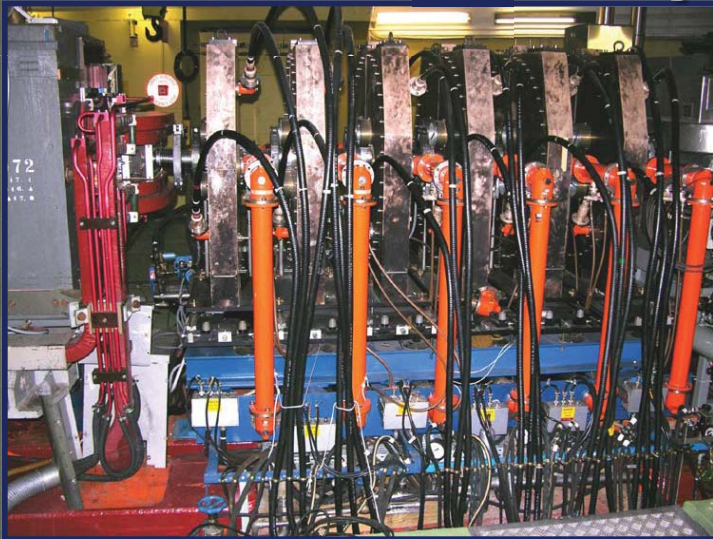


2nd Nov, 2012

CAS Granada - EJ: RF Systems II

62

CERN PS 200 MHz System



2nd Nov, 2012

CAS Granada - EJ: RF Systems II

63

CERN PS 10 MHz cavity (1 of 10)

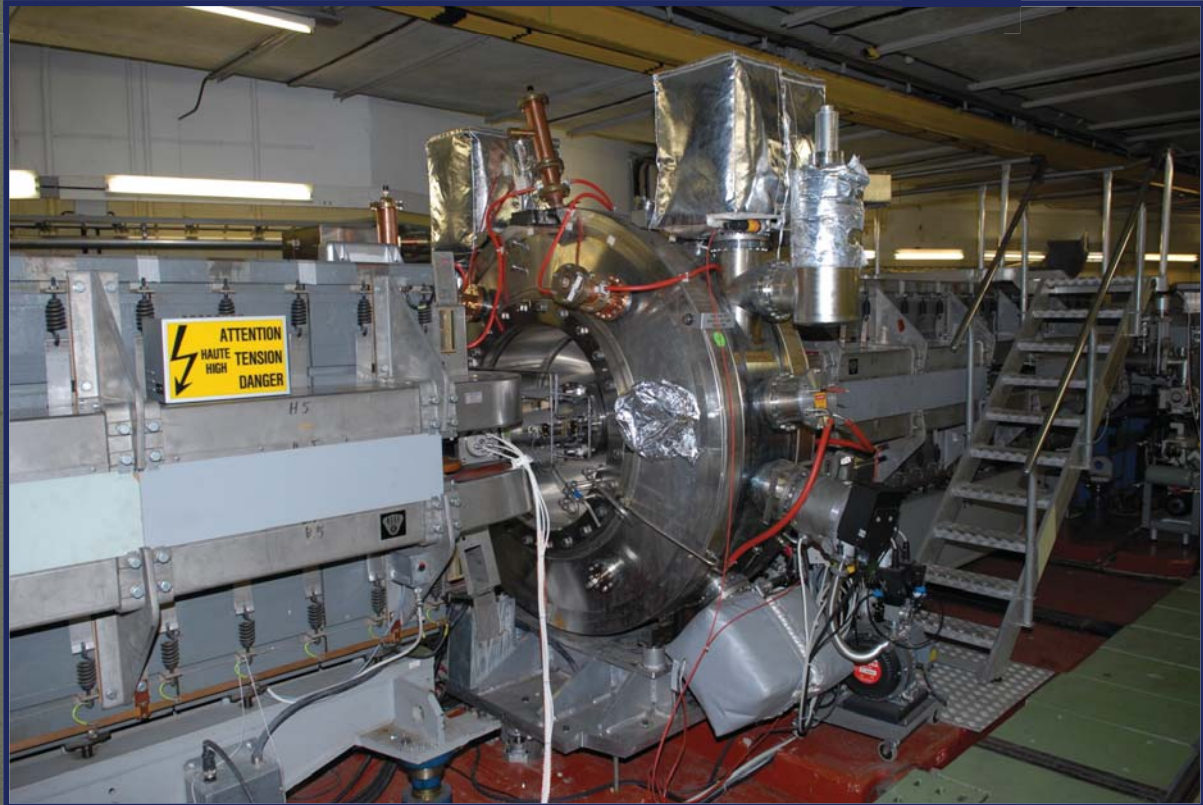


2nd Nov, 2012

CAS Granada - EJ: RF Systems II

64

CERN PS 80 MHz Cavity (1997)



2nd Nov, 2012

CAS Granada - EJ: RF Systems II

65

SPS 200 MHz RF system



4 TW cavities



Siemens:
4 x 550 kW (28 tetrode amplifiers)

Philips:
4 x 550 kW (72 tetrode amplifiers)

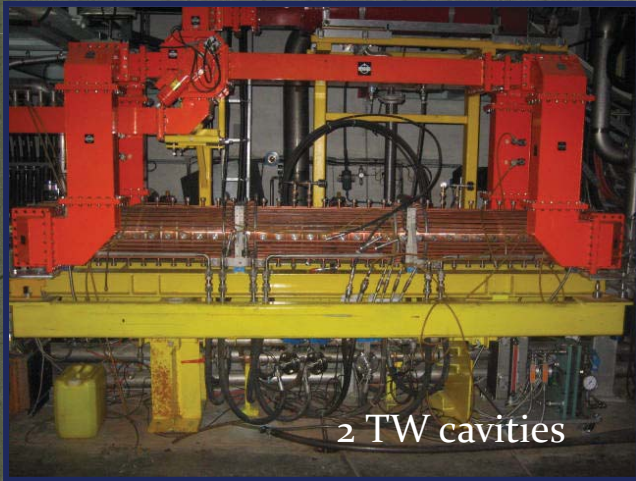


2nd Nov, 2012

CAS Granada - EJ: RF Systems II

66

SPS 800 MHz RF System



2 TW cavities

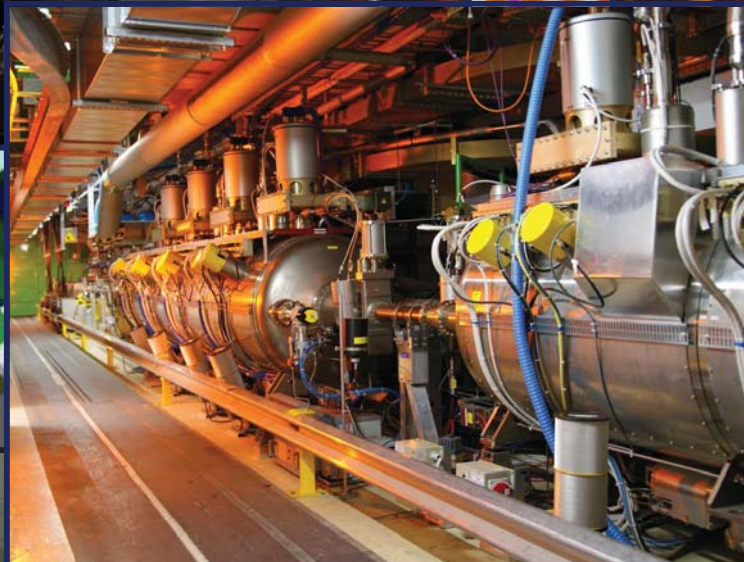


8 (old) klystrons

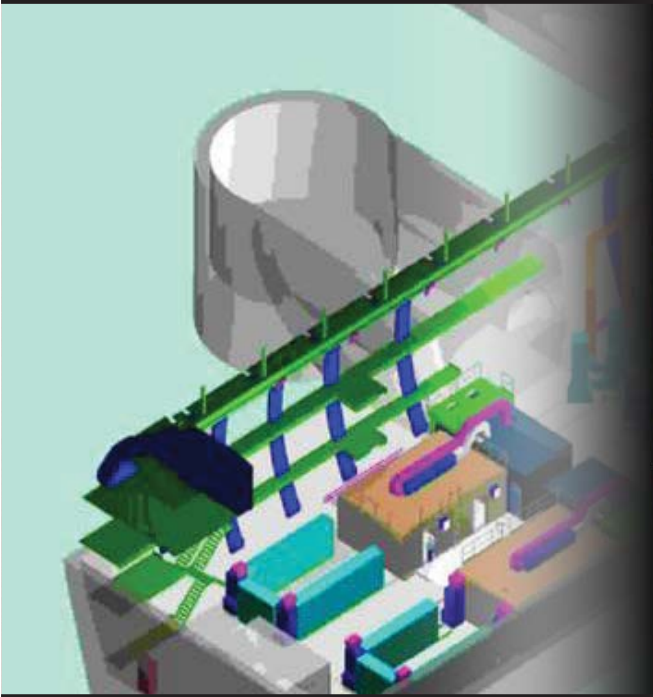
... soon to be replaced by an IOT based system



LHC High-Power RF System



LHC High-Power RF System



- 16 Klystrons
- 4 SC Cavity Modules
- 300 kW @ 400 MHz
- 1000 Interlocks
- ... of which each could dump the beam

End of RF Systems II

Thank you for your attention!