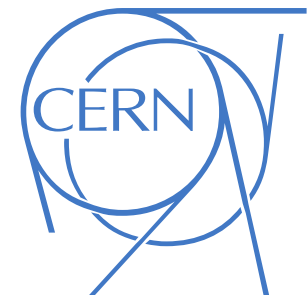


Electroweak symmetry breaking

Search for the missing piece of the Standard Model II

Pedro Ferreira da Silva – psilva@cern.ch

(CERN/LIP)



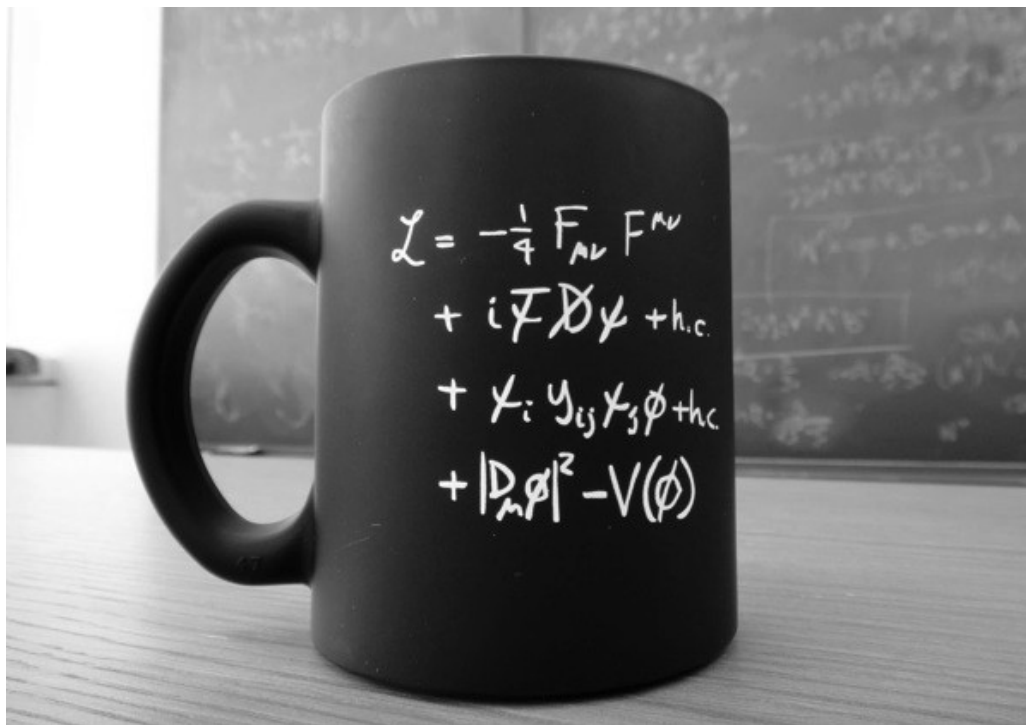
Plan for today

Summary from 1st lecture

ZZ production at the LHC

The golden channel: $H \rightarrow ZZ \rightarrow 4l$

High mass search: $H \rightarrow ZZ \rightarrow 2l2\nu$



Summary from 1st lecture

Gauge bosons have self interactions

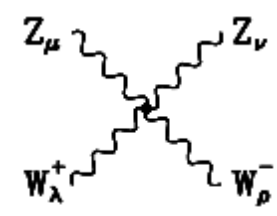
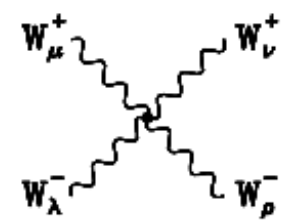
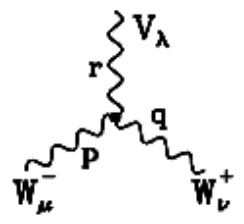
- The Standard Model allows for **pure gauge-bosons interactions**

$$\mathcal{L}_{\text{gauge-fixing}} = -\frac{1}{4}W_{\mu\nu}^i W^{\mu\nu i} - \frac{1}{4}B_{\mu\nu} B^{\mu\nu}$$

$F_{\mu\nu}$ is the field strength tensor which for the electroweak sector is given by:

$$W_{\mu\nu}^i = \partial_\mu W_\nu^i - \partial_\nu W_\mu^i - g_W \epsilon^{ijk} W_\mu^j W_\nu^k \quad B_{\mu\nu} = \partial_\mu B_\nu - \partial_\nu B_\mu$$

- This allows for **triple and quartic gauge boson interactions**



etc.

$$+ig_V [(p-q)_\lambda g_{\mu\nu} + (q-r)_\mu g_{\nu\lambda} + (r-p)_\nu g_{\lambda\mu}]$$

(all momenta incoming,
 $g_\lambda = e, g_z = g_W \cos\theta_W$)

$$+ig_W^2 [2g_{\mu\nu}g_{\lambda\rho} - g_{\mu\lambda}g_{\nu\rho} - g_{\mu\rho}g_{\nu\lambda}]$$

$$-ig_W^2 \cos^2\theta_W [2g_{\mu\nu}g_{\lambda\rho} - g_{\mu\lambda}g_{\nu\rho} - g_{\mu\rho}g_{\nu\lambda}]$$

Longitudinal vector boson scattering

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- Longitudinal polarization is possible for the massive vector bosons
- **Scattering of longitudinal polarized W bosons breaks unitarity at high $s^{1/2}$**

$$\sigma(W_L^+ W_L^- \rightarrow W_L^+ W_L^-) \sim s$$

- At $s^{1/2} \sim 1$ TeV interactions become strong unless unitarity is restored
- Scalar boson (H) interaction is a possible mechanism provided that:

$$g_{HWW} \sim M_W \quad g_f \sim M_f \quad M_H < 1 \text{ TeV}$$

- Then the cross section saturates (i.e. becomes constant) at high $s^{1/2}$

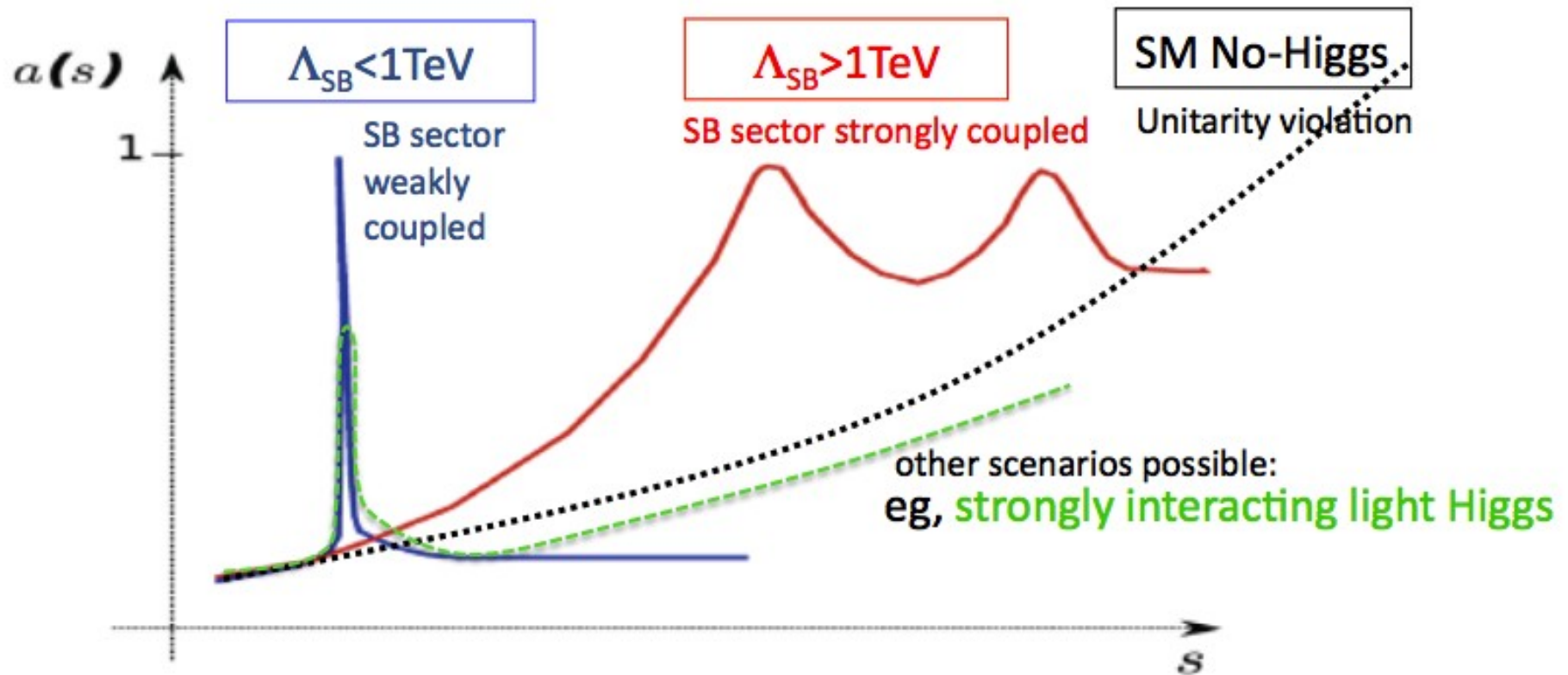
$$A(W^+ W^- \rightarrow W^+ W^-) \xrightarrow{s \gg M_W^2} \frac{1}{v^2} \left[s + t - \frac{s^2}{s - M_H^2} - \frac{t^2}{t - M_H^2} \right]$$

$$a_0 \xrightarrow{s \gg M_H^2} -\frac{M_H^2}{8\pi v^2} \Rightarrow M_H < 870 \text{ GeV}$$

Possible scenarios for VV scattering

6/46

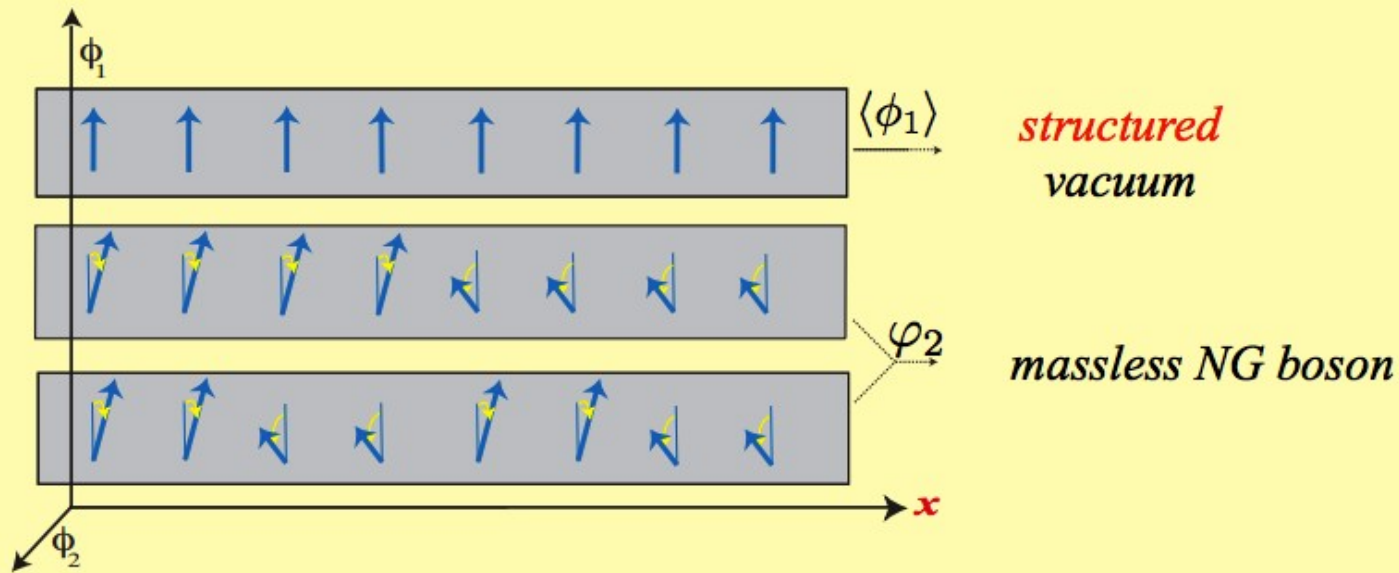
- If the scalar boson is strongly interacting / absent should observe distinct effects



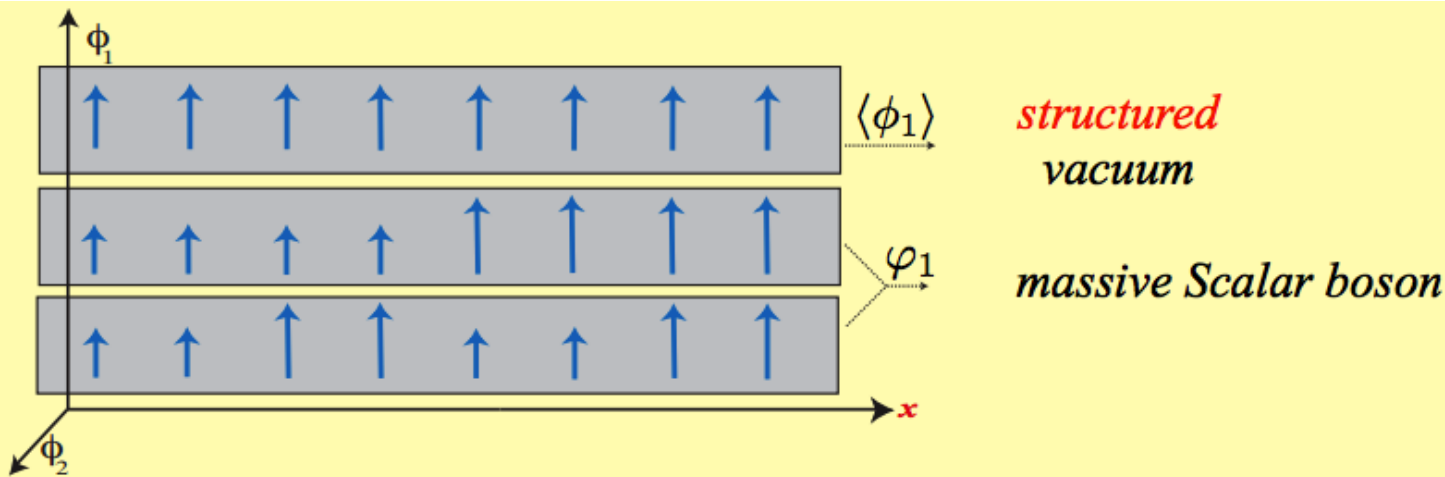
- VV scattering = fundamental probe of how the original EWK symmetry is broken

Original idea behind the Higgs mechanism

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► characterizes a continuous SSB



► measures the rigidity of the vacuum

The proponents

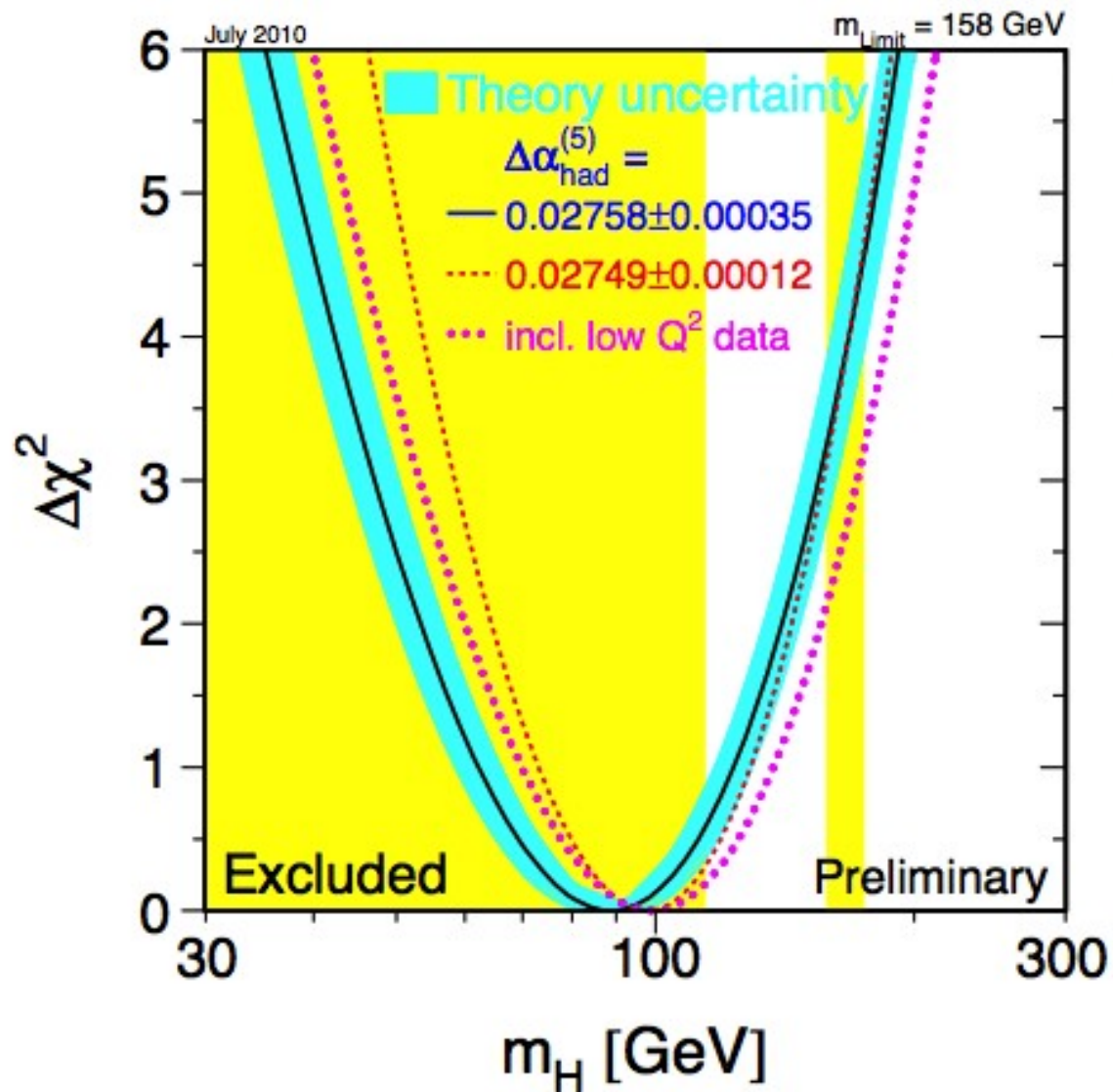
F. Englert and R. Brout PRL 13-[9] (1964) 321

P.W. Higgs PL 12 (1964) 132 and PRL 13-[16] (1964) 508

G.S. Guralnik, C.R. Hagen and T.W.B. Kibble PRL 13-[20] (1964) 585

What else do we know about the Higgs

8/46



EW-Fits:

$$M_H = 89^{+35}_{-26} \text{ GeV}$$

$$M_H < 158 \text{ GeV @ 95\% CL}$$

From direct
search at LEP:

$$M_H > 114 \text{ GeV}$$

@ 95% CL

From direct
search at Tevatron:

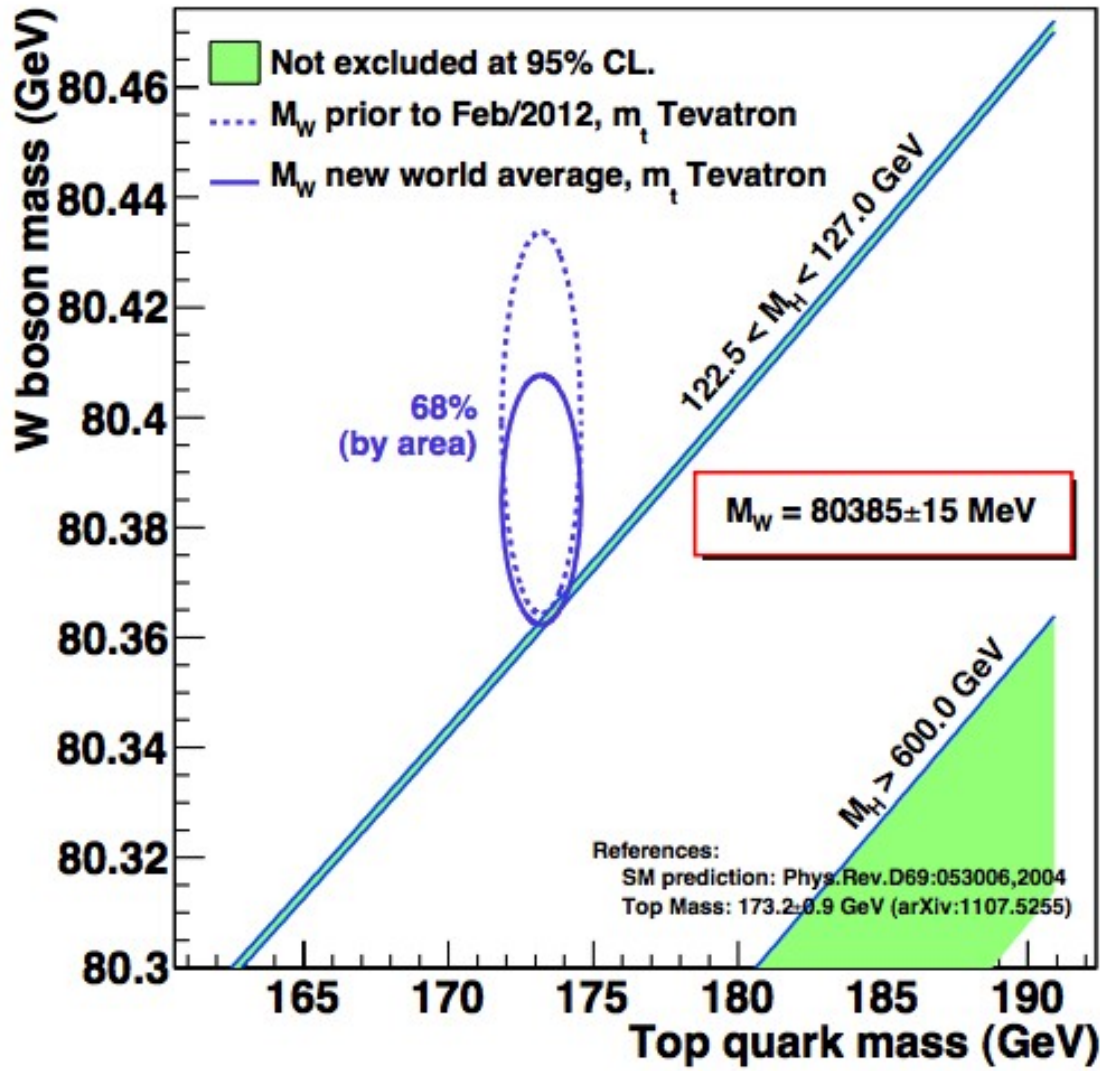
$$158 < M_H < 175 \text{ GeV}$$

@ 95% CL

[Updated: Summer 2010]

What else do we know about the Higgs

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- The mass of the heaviest particles is correlated from loop corrections including the Higgs boson
- The preferred region is still compatible at 68% CL with the not yet excluded SM Higgs mass at 95% CL

Higgs partial widths (tree level)

10/46

- **Fermions:** proportional to the mass and velocity dependent (1 factor from the matrix elem.+ 2 from phase space)

$$\Gamma_{f\bar{f}} = \frac{N_c G_F m_f^2 M_H}{4\sqrt{2}\pi} \beta^3$$

$$\text{where } \beta = \sqrt{1 - \frac{4m_f^2}{M_H^2}}$$

- **Vector bosons:** dominate due to the fact the longitudinal polarized bosons couple $\sim E \rightarrow$ coupling to Higgs as to rise as fast

$$\Gamma_{VV} = \frac{G_F M_H^3}{16\sqrt{2}\pi} \delta_V \beta \left(1 - x_V + \frac{3}{4} x_V^2\right)$$

$$\text{where } \begin{cases} \delta_{W,Z} = 2, 1 \\ \beta = \sqrt{1 - x_V} \\ x_V = \frac{4M_V^2}{M_H^2} \end{cases}$$

- **Gluons:** through top quark loops
- **Photons through top and W boson loops** (Z γ partial width is similar in structure)

$$\Gamma_{gg} = \frac{\alpha_s^2 G_F M_H^3}{16\sqrt{2}\pi^3} \left| \sum_i \tau_i [1 + (1 - \tau_i) f(\tau_i)] \right|^2$$

$$\text{with } \tau_i = \frac{4m_i^2}{M_H^2} \text{ and } f(\tau) = \begin{cases} [\sin^{-1} \sqrt{1/\tau}]^2 & \tau \geq 1 \\ -\frac{1}{4} [\ln \frac{1+\sqrt{1-\tau}}{1-\sqrt{1-\tau}} - i\pi]^2 & \tau < 1 \end{cases}$$

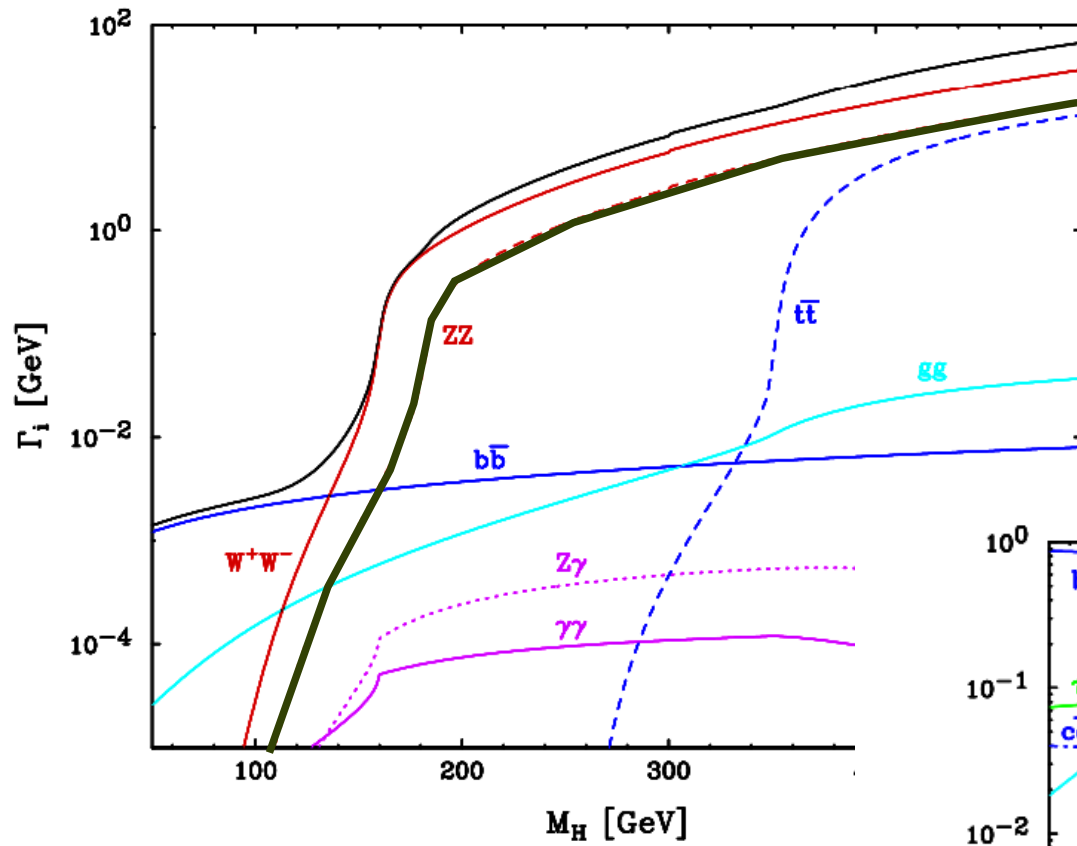
$$\Gamma_{\gamma\gamma} = \frac{\alpha^2 G_F M_H^3}{128\sqrt{2}\pi^3} \left| \sum_i N_{c,i} Q_i^2 F_i \right|^2$$

$$F_1 = 2 + 3\tau[1 + (2 - \tau)f(\tau)]$$

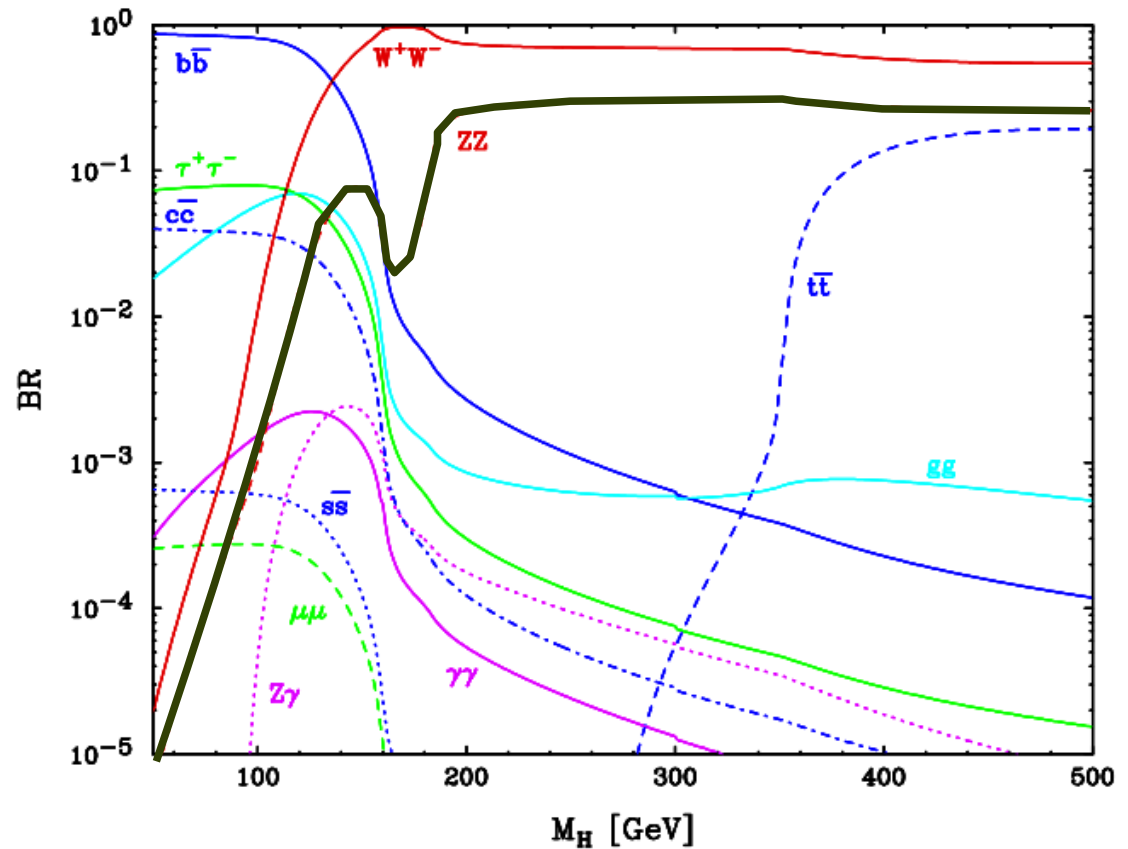
$$F_{1/2} = -2\tau[1 + (1 - \tau)f(\tau)]$$

$$F_0 = \tau[1 - \tau f(\tau)]$$

Higgs partial widths and branching ratios



Today's lecture



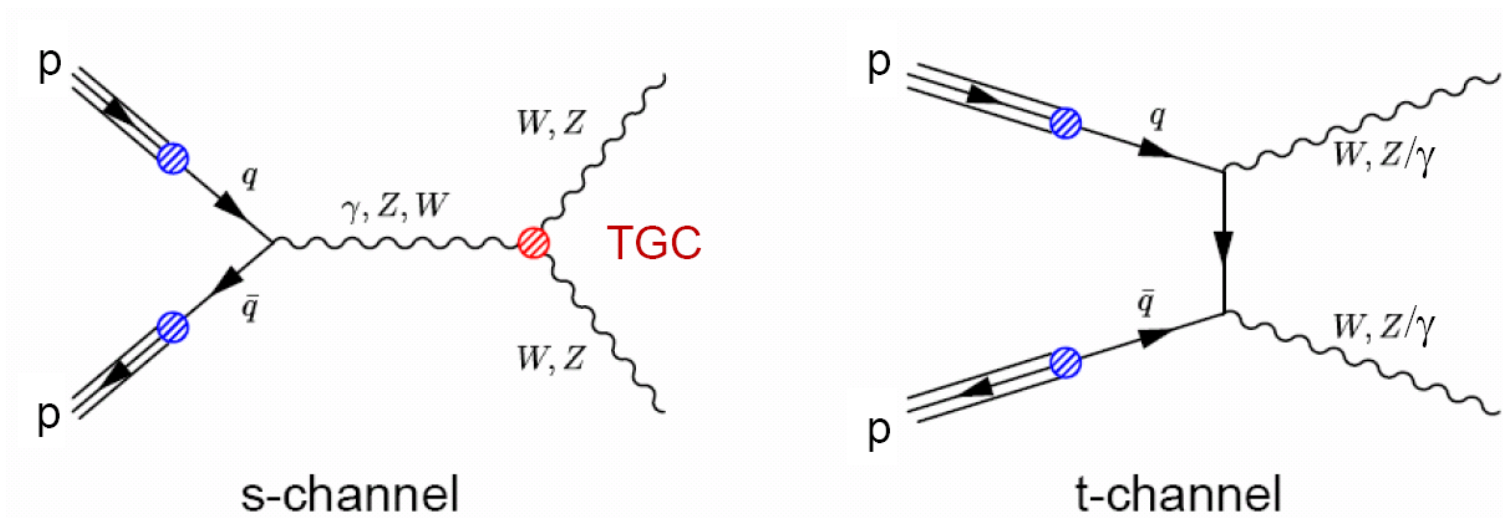
zzzzz **ZZ production at the LHC**



Diboson production

13/46

- Processes which produce **WW, WZ or ZZ** final states can help answering
 - why are EWK bosons massive?
 - how does the **EWK symmetry breaking** occur?
- New Physics expected to lead to EWKSB may be sought in di-boson production:
 - direct evidence of new particles
 - indirect evidence of observing anomalous TGCs



Anomalous Triple Gauge Couplings

$$L/g_{WWV} = ig_1^V (W_{\mu\nu}^* W^{\mu\nu} V^\nu - W_{\mu\nu} W^{*\mu\nu} V^\nu) + i\kappa^V W_\mu^* W_\nu^\mu V^{\nu\rho} + \frac{\lambda^V}{M_W^2} W_{\rho\mu}^* W_\nu^\mu V^{\nu\rho} \quad \text{charged}$$

$$L = -\frac{e}{M_Z^2} [f_4^V (\partial_\mu V^{\mu\beta}) Z_\alpha (\partial^\alpha Z_\beta) + f_5^V (\partial^\sigma V_{\sigma\mu}) \tilde{Z}^{\mu\beta} Z_\beta] \quad \text{neutral}$$

- All these anomalous terms are allowed in the SM lagrangian

→ Couplings are usually proportional to s or $s^{1/2}$ and lead to tree level unitarity

→ Apply effective cut-off scale

Final State	WZ	W γ	WW	ZZ	Z γ
SM					
an.TGC					

Coupling	Source	L (fb ⁻¹)	λ_Z	$\Delta\kappa_Z$	$\Delta\kappa_\gamma$	λ_γ
WW γ from W $^\pm\gamma$	D0 [27]	0.16			[-0.88, 0.96]	[-0.2, 0.2]
WWZ from W $^\pm Z$	D0 [24]	1.0	[-0.17, 0.21]	[-0.12, 0.29]		
WWZ from W $^\pm Z$	CDF	1.9	[-0.13, 0.14]	[-0.82, 1.27]		
WWZ = WW γ						
from W $^+W^-$	D0 [30]	0.25	[-0.31, 0.33]	[-0.36, 0.33]		
from W $^+W^-$, W $^\pm Z$	CDF [31]	0.35	[-0.18, 0.17]	[-0.46, 0.39]		

Note: more recent results are available, also from LHC experiments

ZZ production

- t-channel production dominates**

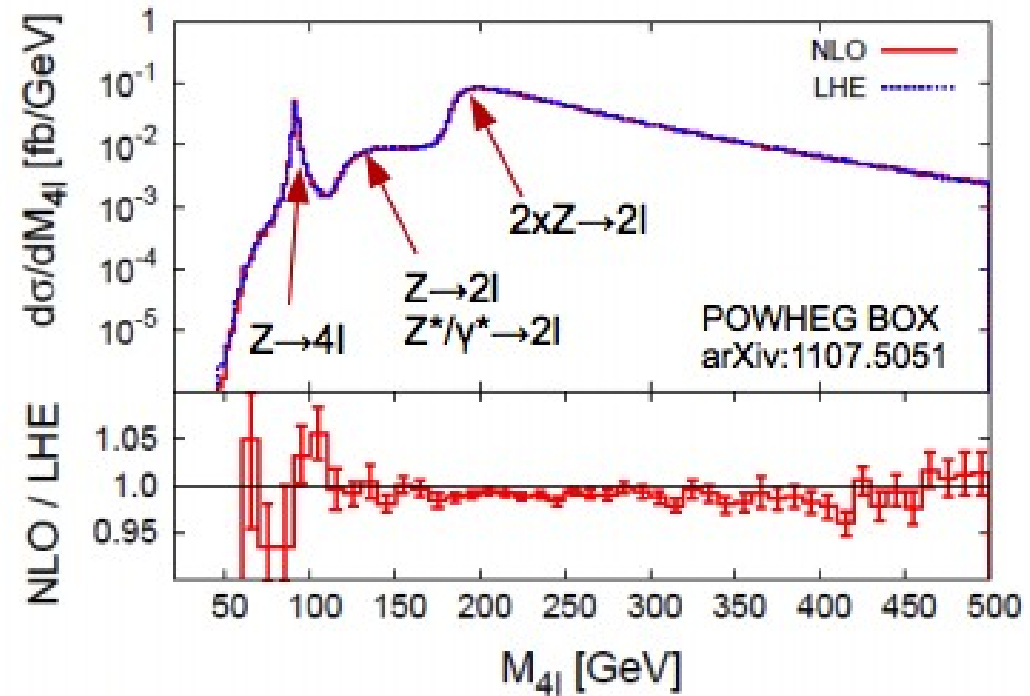
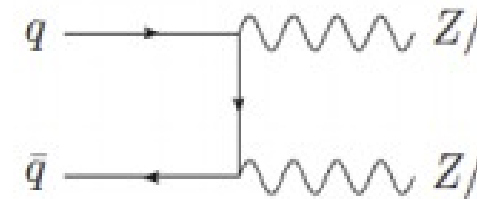
- unlike WZ and WW which have contribution from triple gauge coupling

- s-channel production is suppressed**

- $O(10^{-4})$
- if in excess: **anomalous gauge couplings**
- **resonant production**: Higgs, gravitons

- Cross section measurement**

- At the LHC only measured in 4l ▶
- important to compare decay channels (in particular when looking for deviations)



$$\sigma(pp \rightarrow ZZ) \times BR(Z \rightarrow ll)^2 = 6.4 \pm 0.3 \text{ (NLO)}$$

ATLAS: $ZZ \rightarrow ll$

CMS: $ZZ \rightarrow ll$

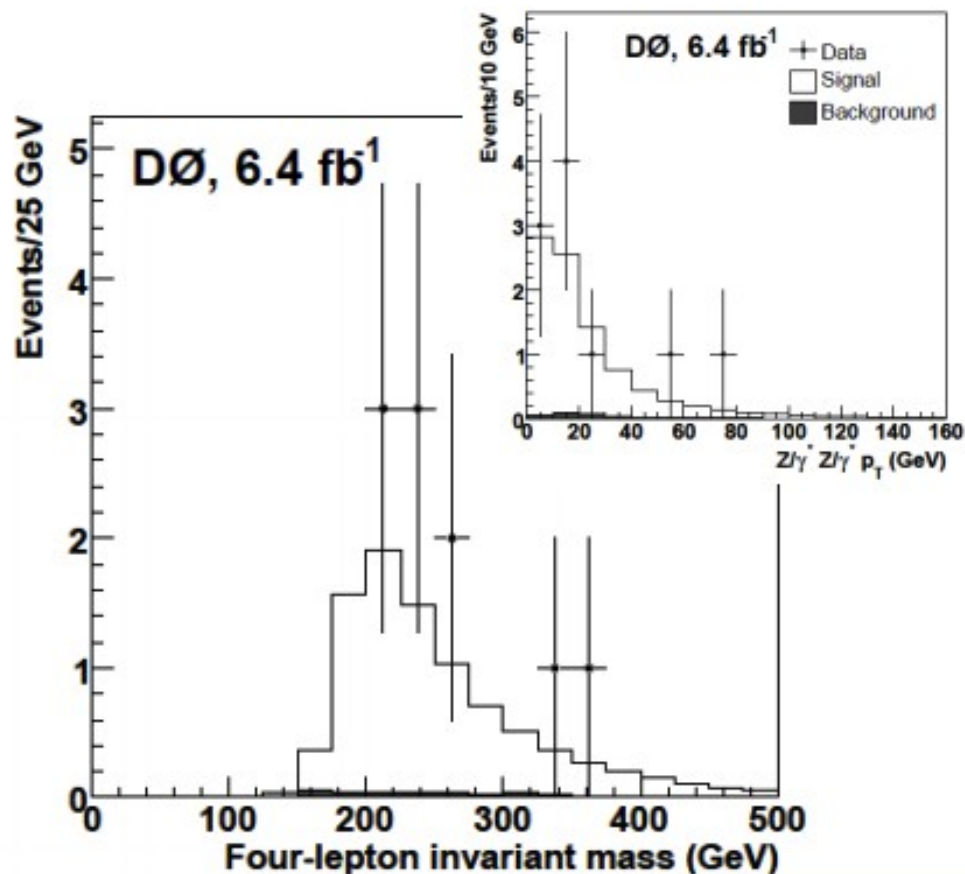
$$\left| \begin{array}{l} 8.4^{+2.7}_{-2.3}(\text{stat})^{+0.4}_{-0.7}(\text{syst}) \pm 0.3(\text{lumi}) \\ 3.8^{+1.5}_{-1.2}(\text{stat}) \pm 0.2(\text{syst}) \pm 0.2(\text{lumi}) \end{array} \right.$$

Search for resonant ZZ production at the Tevatron

16/46

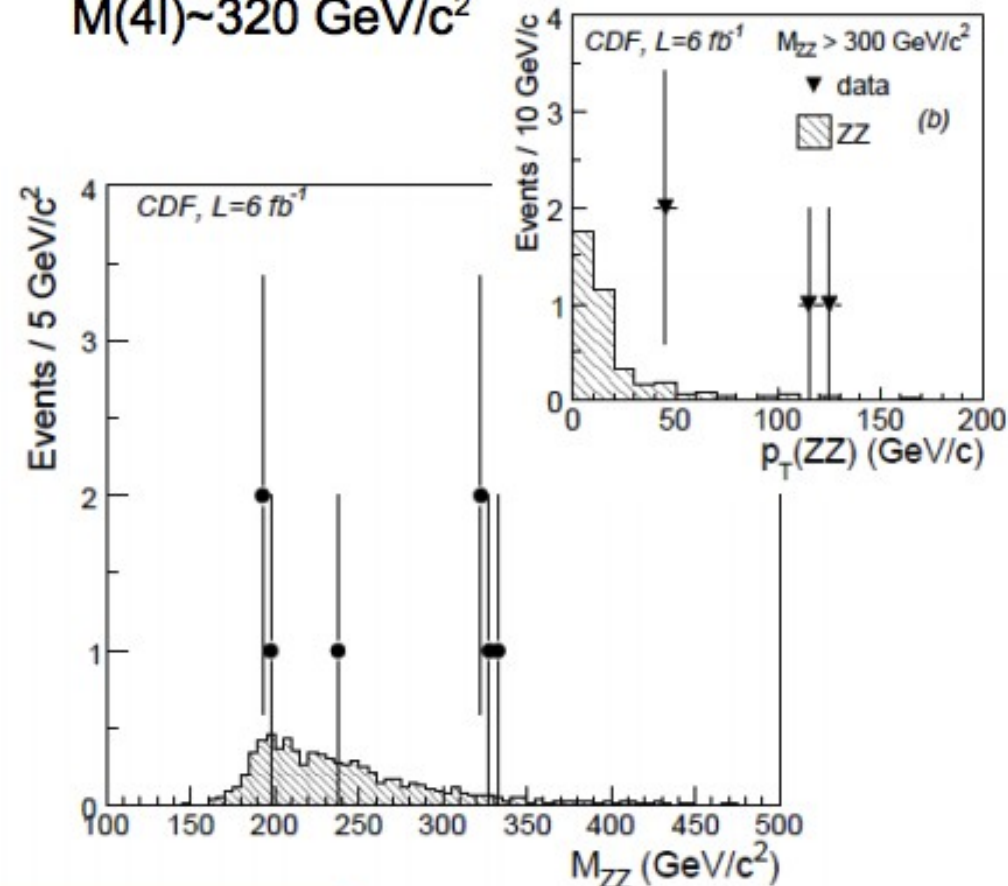
DØ – arXiv:1104.3078

- cross section in agreement with SM
- No particular excess in $M(4l)$



CDF – arXiv:1111.3432

- Cross section in agreement with SM
- Clustering of high p_T events at $M(4l) \sim 320 \text{ GeV}/c^2$



Crucial to compare decay channels

Search for resonant ZZ production at the Tevatron

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Privileged role in resonant production searches for $M_X > 200$ GeV/c²

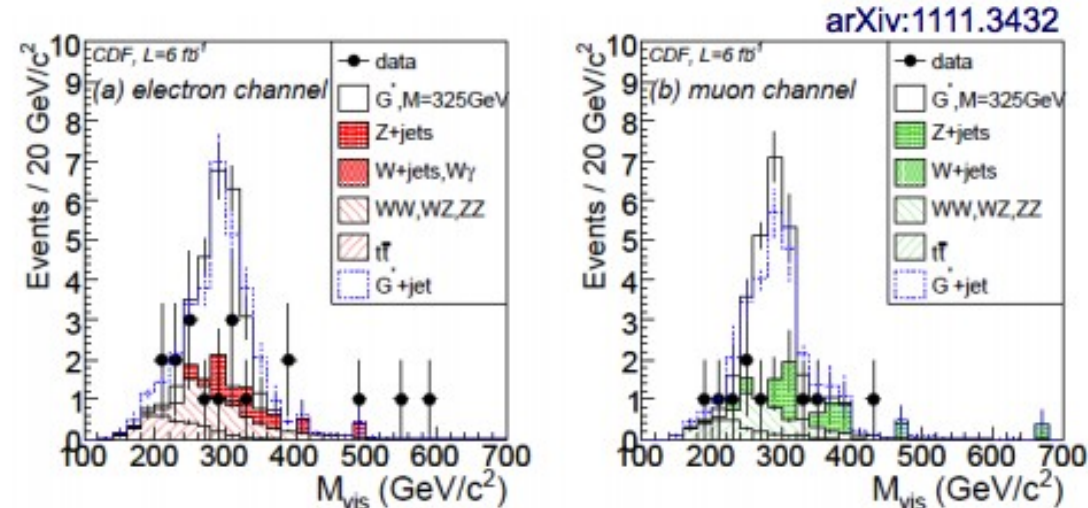
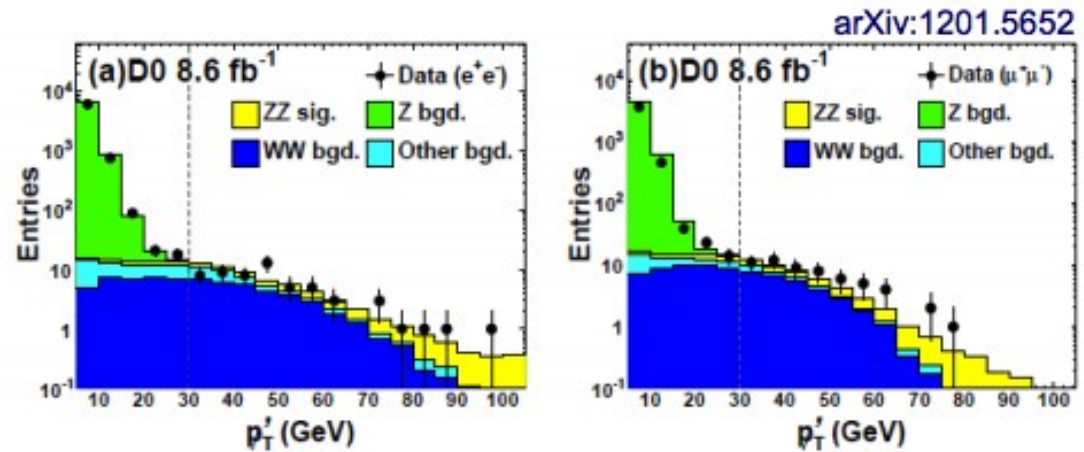
- $BR(ZZ \rightarrow 2l2\nu) \approx 6 \times BR(ZZ \rightarrow 4l)$
- Expected to lead exclusion limits

Challenges:

- partial reconstruction of the kinematics $E_T^{miss} = p_T(\nu_1) + p_T(\nu_2)$
- resonant signature through transverse or visible mass spectrum

$$M_T^2(ZZ) = [\sqrt{M_Z^2 + p_T^2(\ell\ell)} + \sqrt{M_Z^2 + p_{Tmiss}^2}]^2 - |\vec{p}_T(\ell\ell) + \vec{p}_{Tmiss}|^2$$

- backgrounds: Z+jets (instrumental) and di-bosons (mainly SM ZZ)



Neither CDF or D0 observe significant excesses in the 2l2ν final states

(similar results in the 2l2q channel)





The golden channel
H → ZZ → 4I

CMS Experiment at the LHC, CERN

Data recorded: 2011-May-25 08:00:19.229673 GMT (10:00:19 CEST)

Run / Event: 165633 / 394010457

2e2μ candidate

$m_{41} = 244.6 \text{ GeV}/c^2$

$m_{Z1} = 91.2 \text{ GeV}/c^2$

$m_{Z2} = 93.2 \text{ GeV}/c^2$

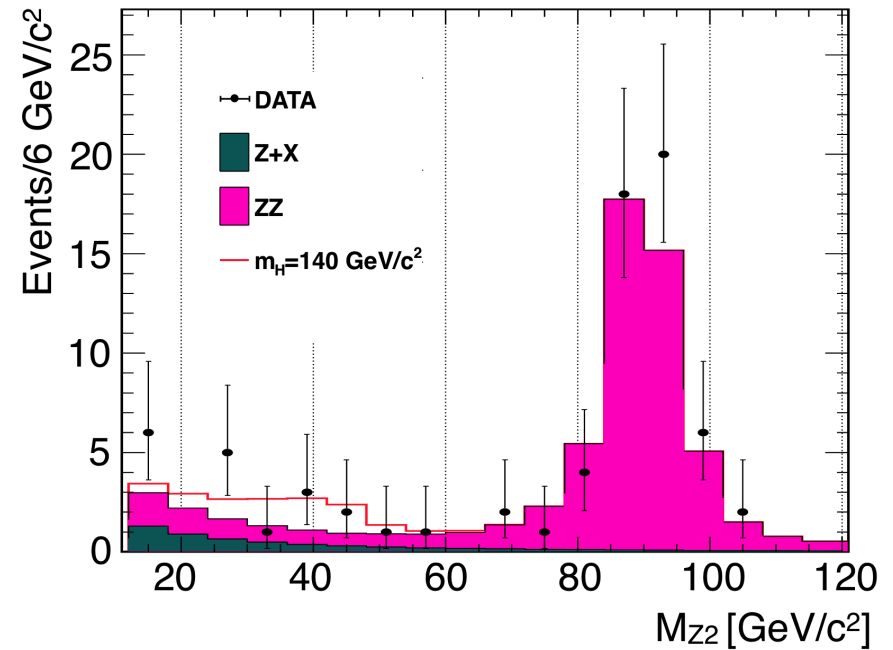
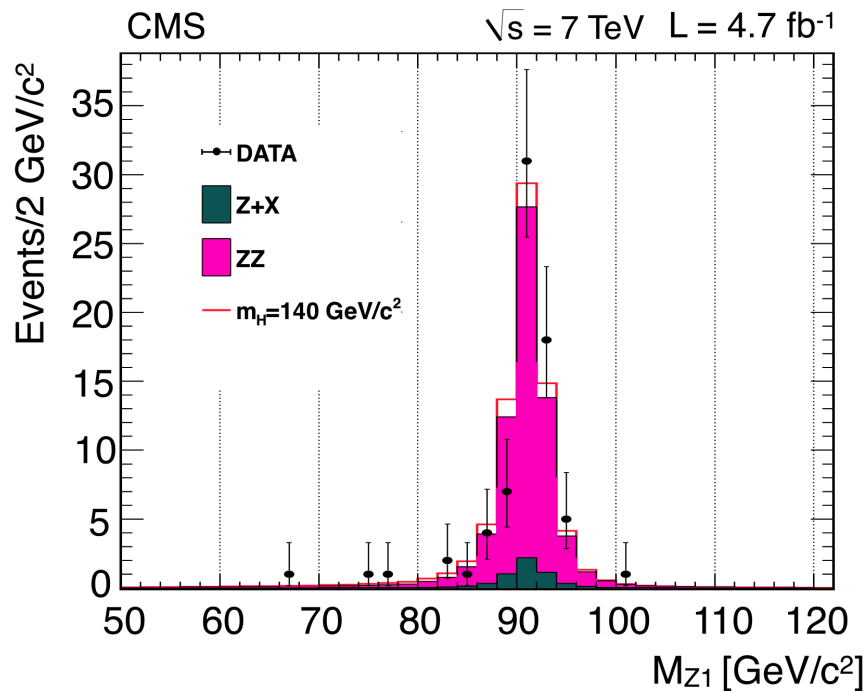
Channel signature - 1

20/46

- It's the **cleanest channel among all**

→ **2 high mass lepton pairs** : $50 < m_{Z_1} < 120 \text{ GeV}/c^2$ and $12 < m_{Z_2} < 120 \text{ GeV}/c^2$

(the second pair is allowed to be an off-shell Z)

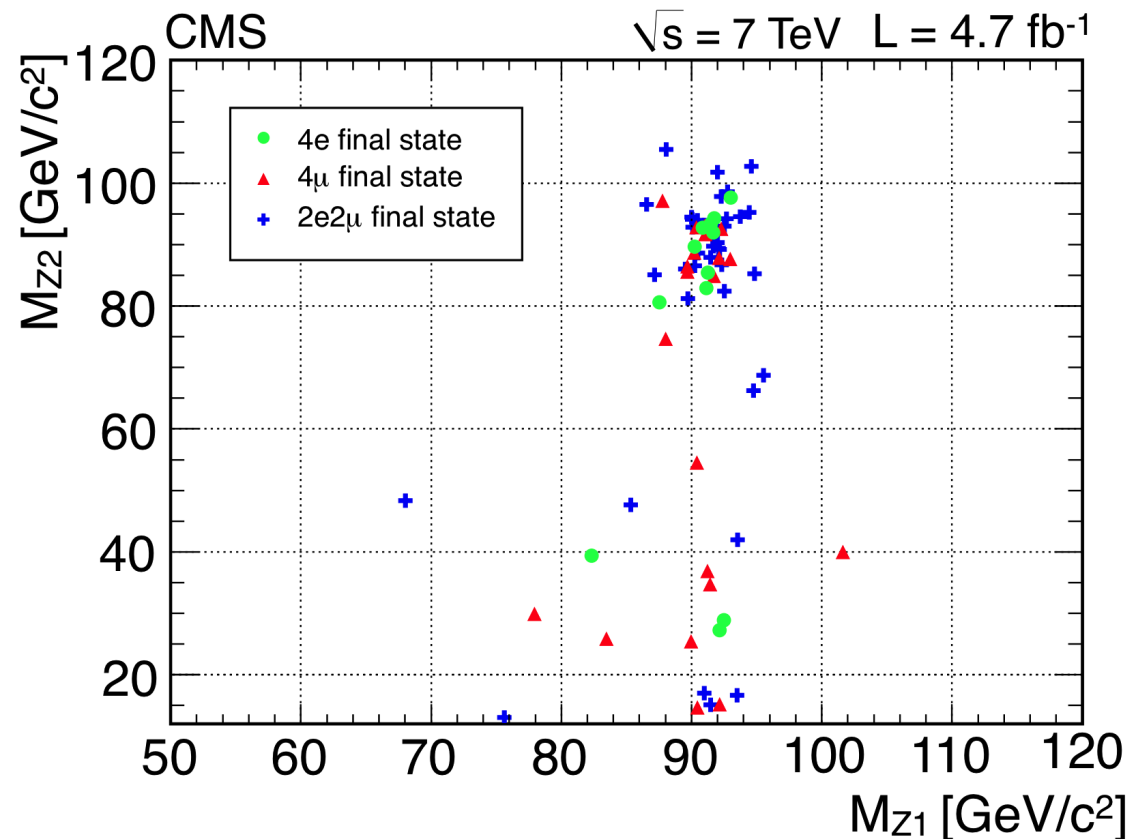


Channel signature - 2

- It's the **cleanest channel among all**

→ **2 high mass lepton pairs** : $50 < m_{Z_1} < 120 \text{ GeV}/c^2$ and $12 < m_{Z_2} < 120 \text{ GeV}/c^2$

(the second pair is allowed to be an off-shell Z)



Channel signature - 3

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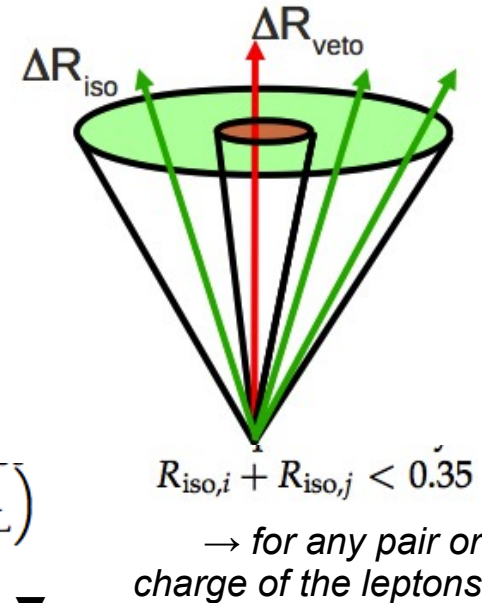
- 4 isolated leptons in the final state: 4e / 2e2μ / 4μ

→ Leptons are soft in p_T : $p_T^e > 7$ GeV/c $p_T^\mu > 5$ GeV/c

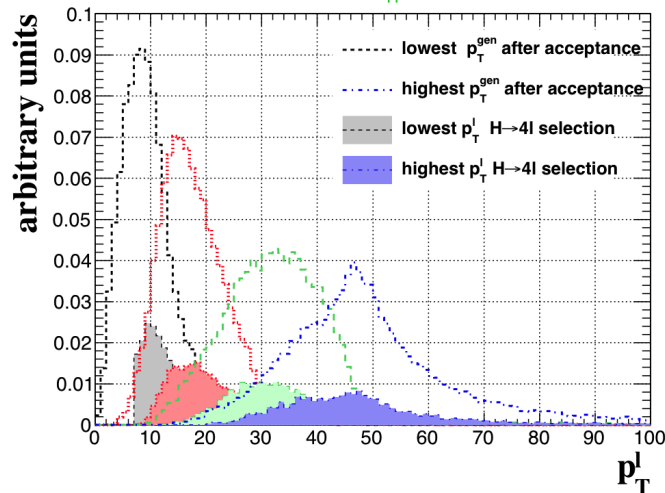
→ Relative isolation is defined from the sum of tracks / calorimeter deposits in a $R=0.3$ cone built around the lepton thrust

$$R_{\text{iso}} = (1/p_T^\ell) \times \left(\sum_i p_{T,\text{track}}^i + \sum_j E_{T,\text{ECAL}}^j + \sum_k E_{T,\text{HCAL}}^k \right)$$

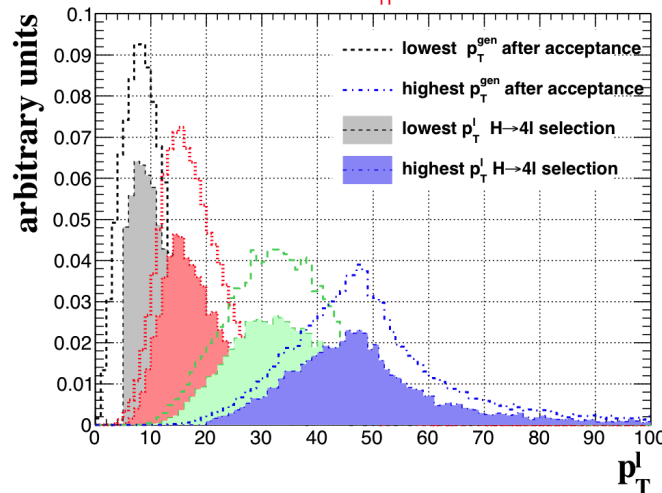
→ Selection efficiency is affected acceptance, p_T and mass requirements ▼



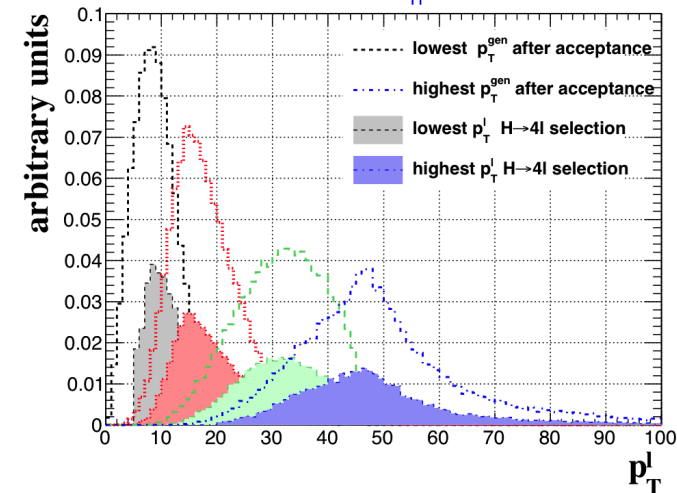
$H \rightarrow ZZ^* \rightarrow 4e$, $m_H = 120$ GeV



$H \rightarrow ZZ^* \rightarrow 4\mu$, $m_H = 120$ GeV



$H \rightarrow ZZ^* \rightarrow 2e2\mu$, $m_H = 120$ GeV



Lepton efficiencies

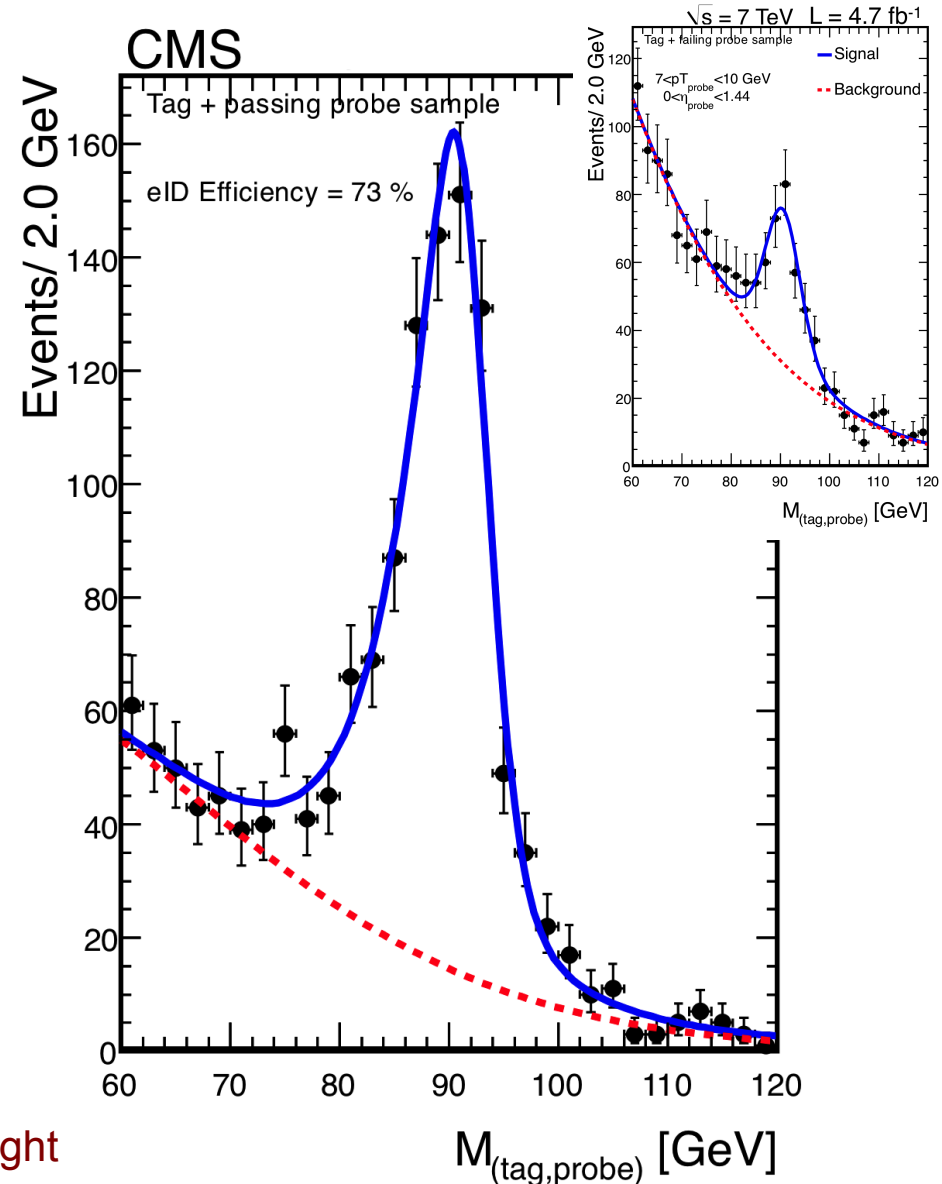
- Usually determined with a tag and probe method
- Choose a dilepton candle: $Z, J/\psi \rightarrow \ell\ell$
- Select tightly the first lepton (=tag) and loosely the second lepton (=probe) constrained by the resonance mass
- Efficiency is measured from:

$$\mathcal{E}^{\text{"tag and probe"}} = \frac{N_{1p} + 2N_{2p}}{N_{1f} + N_{1p} + 2N_{2p}} = \frac{N_p}{N_T}$$

1 fails tight requirement

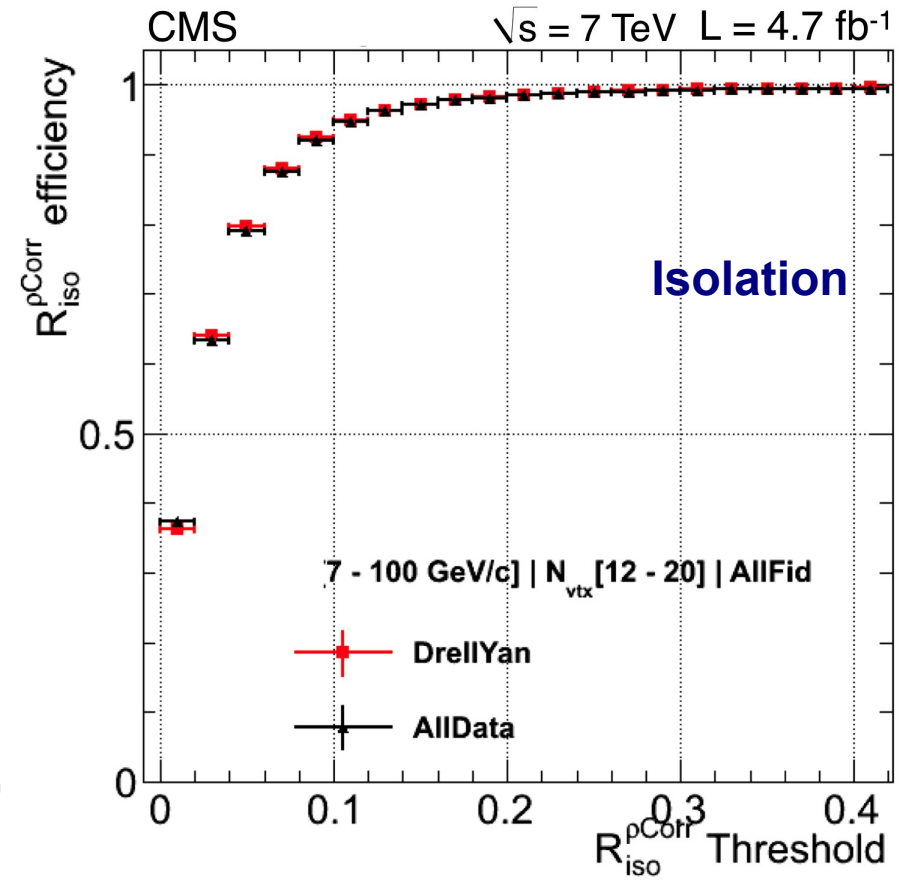
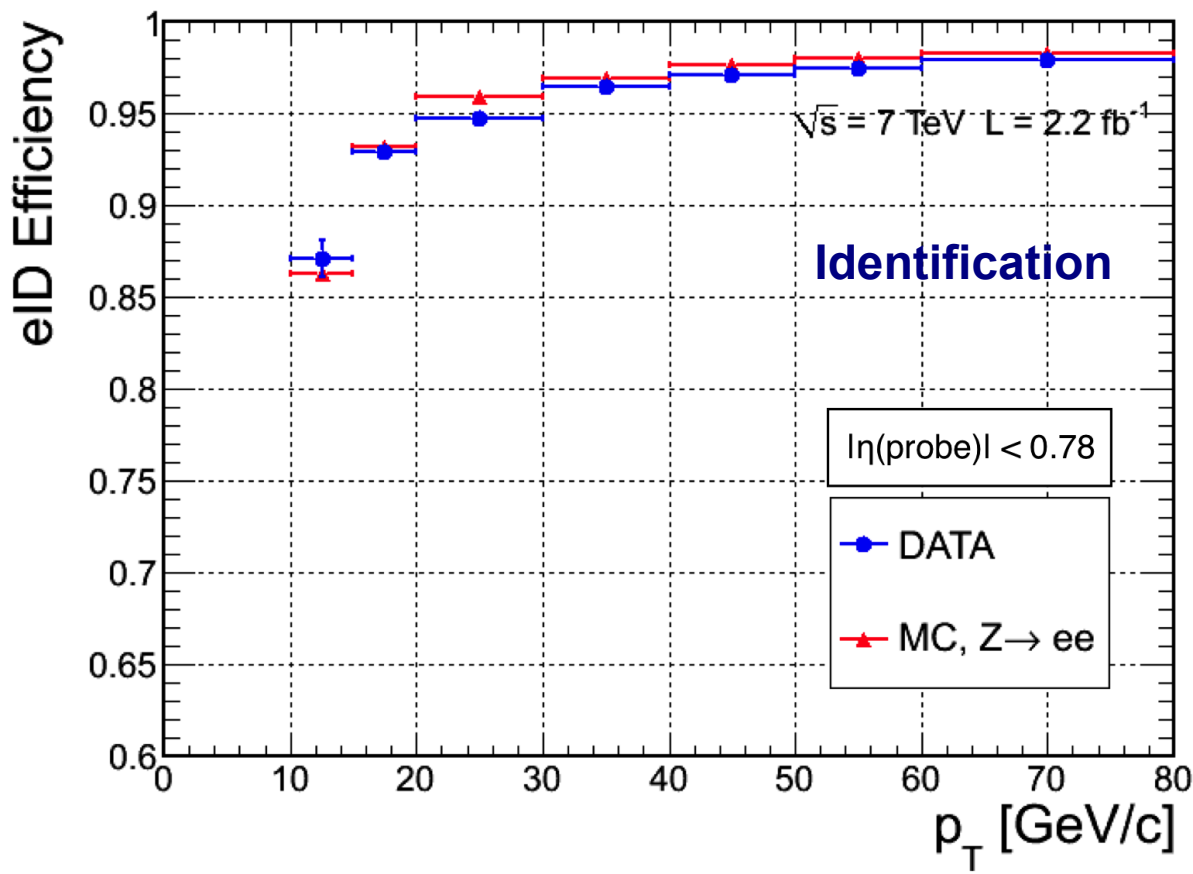
1 passes tight requirement

Both passes tight requirements



Lepton efficiencies

- Tag and probe used to derive separately trigger, reconstruction, identification and isolation efficiencies (efficiency of other cuts can be evaluated the same way)
- The Data/MC ratio is used to correct the simulation (re-weighting)
- Efficiencies drop fast at low p_T (<80%)



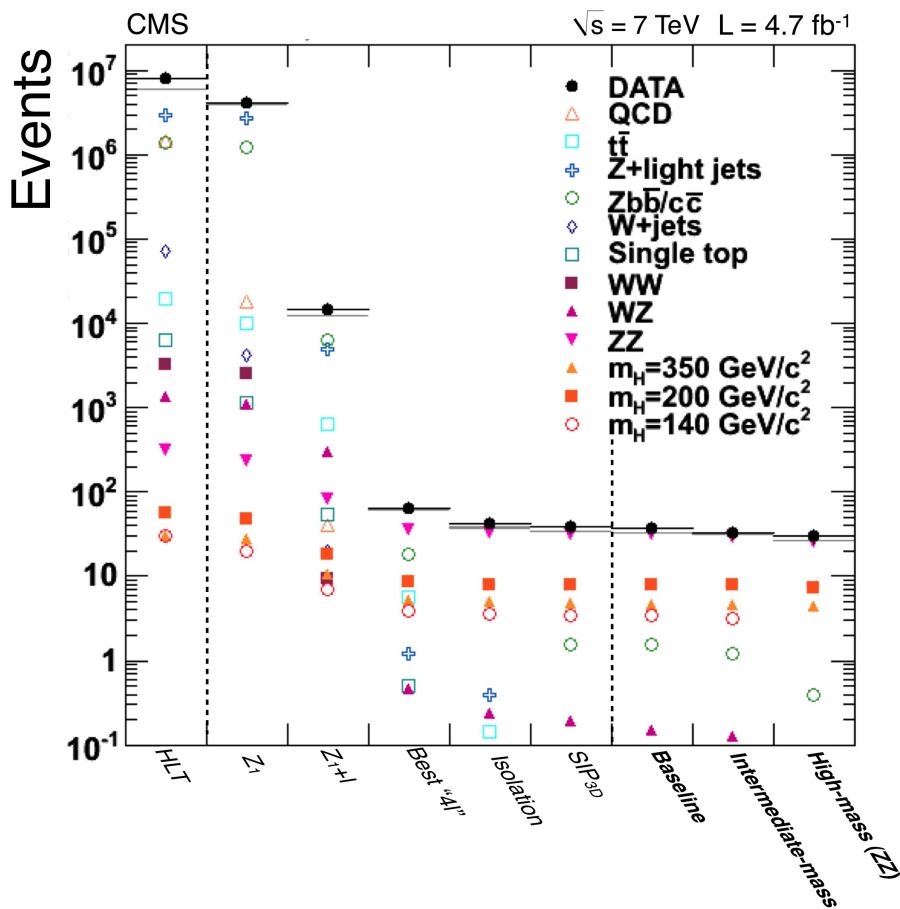
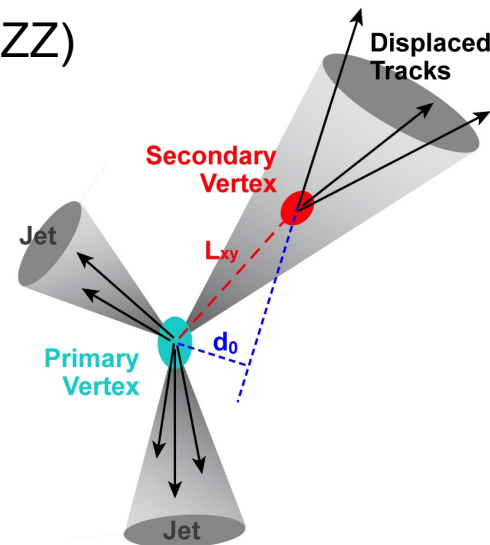
Event selection

- After selection **almost no background expected** (besides SM ZZ)

→ Residual backgrounds from Z + heavy flavor production removed

with impact parameter significance cut $|SIP_{3D}| < 4$

(reject displaced leptons from B hadron decays)



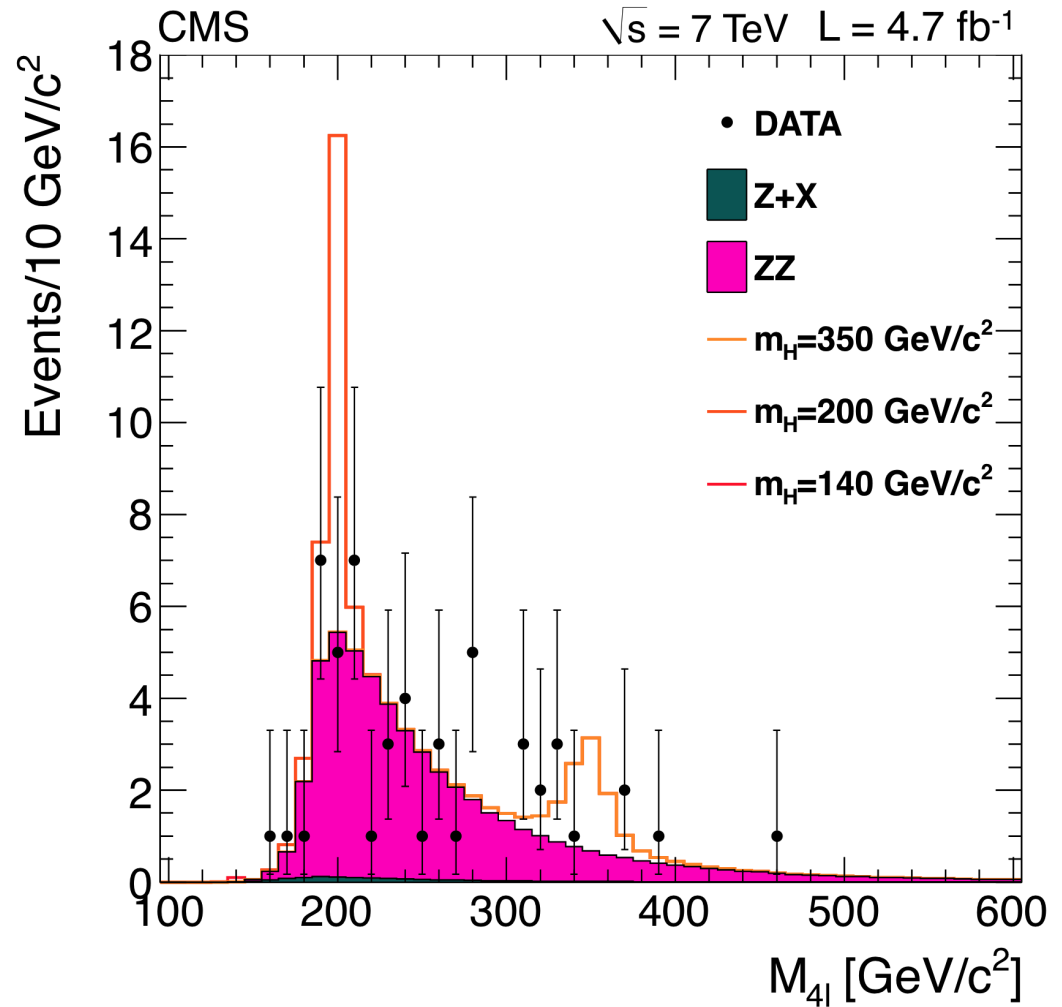
- **Number of selected events is consistent with expectations → look for resonance**

Baseline	4e	4μ	2e2μ
ZZ	12.27 ± 1.16	19.11 ± 1.75	30.25 ± 2.78
Z+X	1.67 ± 0.55	1.13 ± 0.55	2.71 ± 0.96
All background	13.94 ± 1.28	20.24 ± 1.83	32.96 ± 2.94
$m_H = 120 \text{ GeV}/c^2$	0.25	0.62	0.68
$m_H = 140 \text{ GeV}/c^2$	1.32	2.48	3.37
$m_H = 350 \text{ GeV}/c^2$	1.95	2.61	4.64
Observed	12	23	37

4 lepton invariant mass

High mass selection

- We select 72 events
- Expect 67 ± 6 events from background
- Slight excess around 350 GeV
(similar to CDF observation)

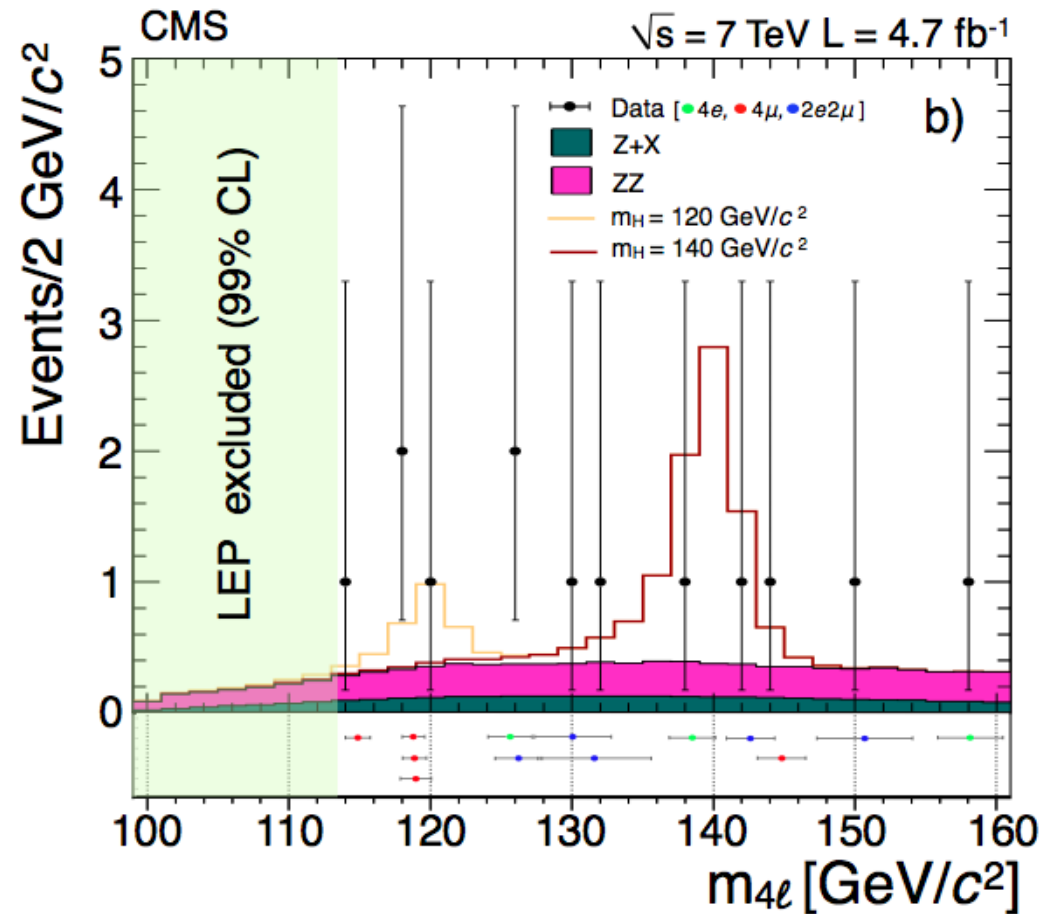
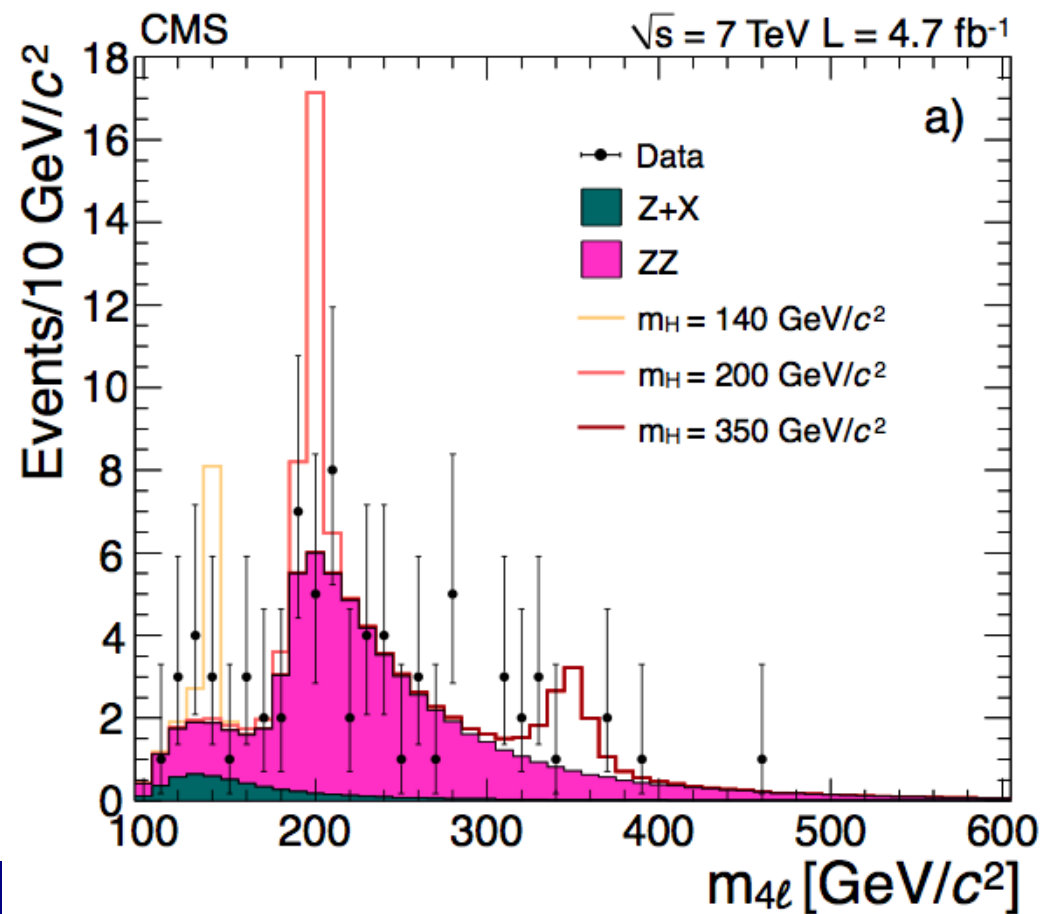


4 lepton invariant mass

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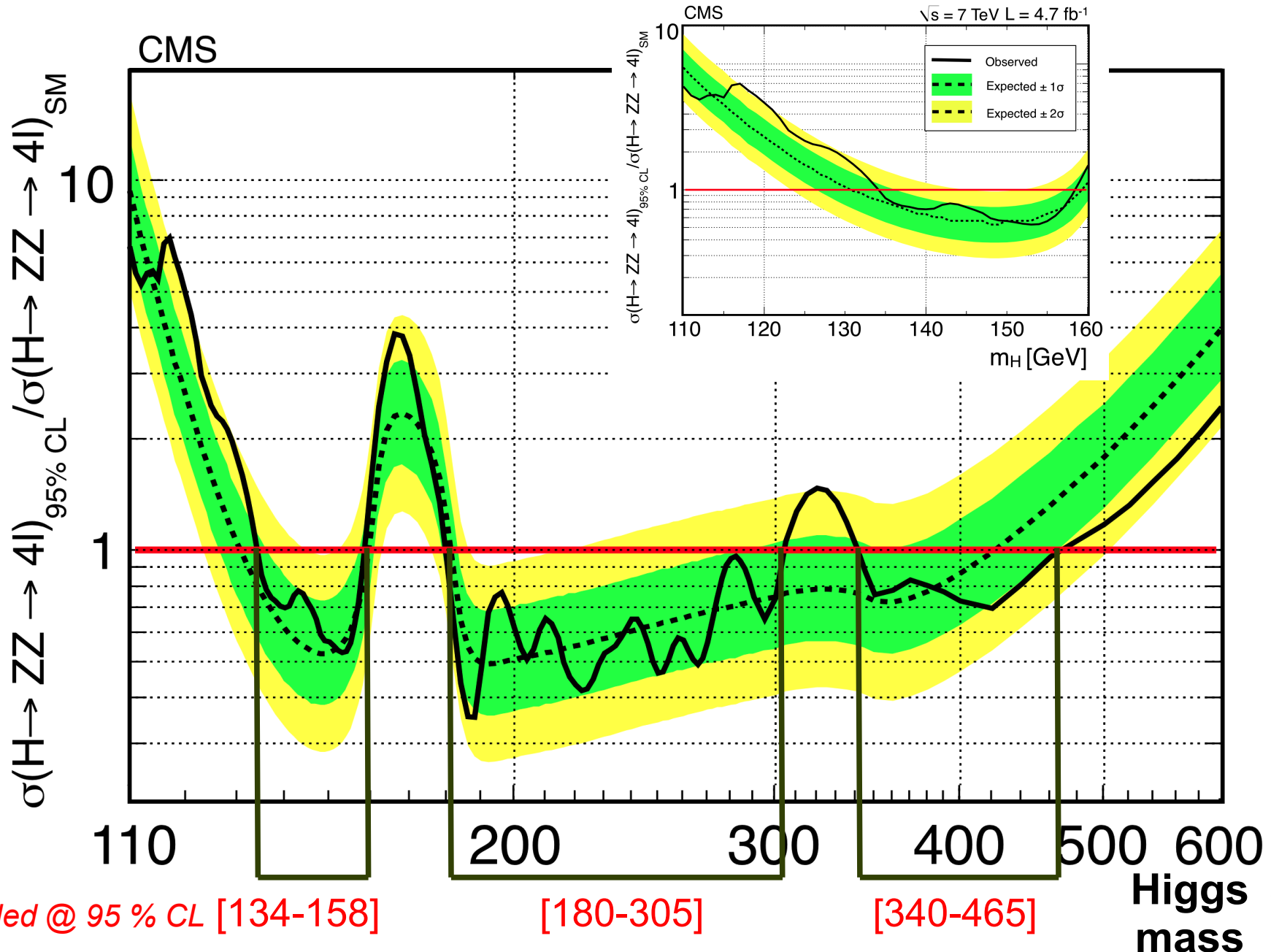
Baseline selection and zoom on low mass

- Observe 13 events between 100 and 160 GeV (expect 9.5 ± 1.3 events)
- Most significant clustering at $119.5 \text{ GeV}/c^2$



Statistical interpretation

- With the full 2011 data the statistics is still low \rightarrow set limits on Higgs production



Excluded @ 95 % CL [134-158]

[180-305]

[340-465]

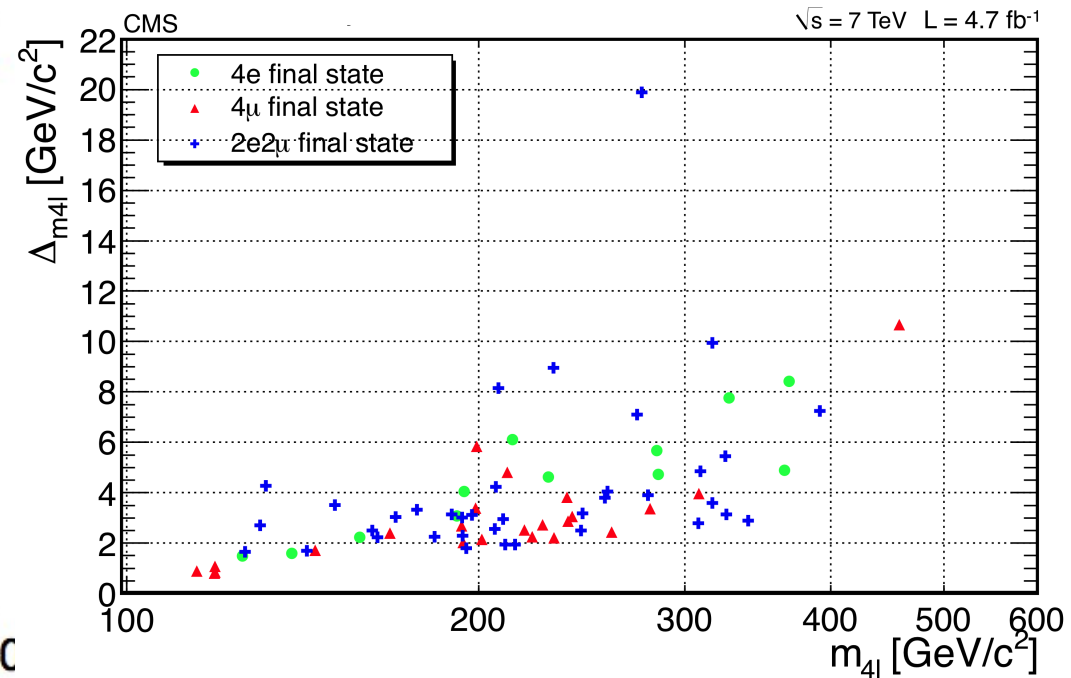
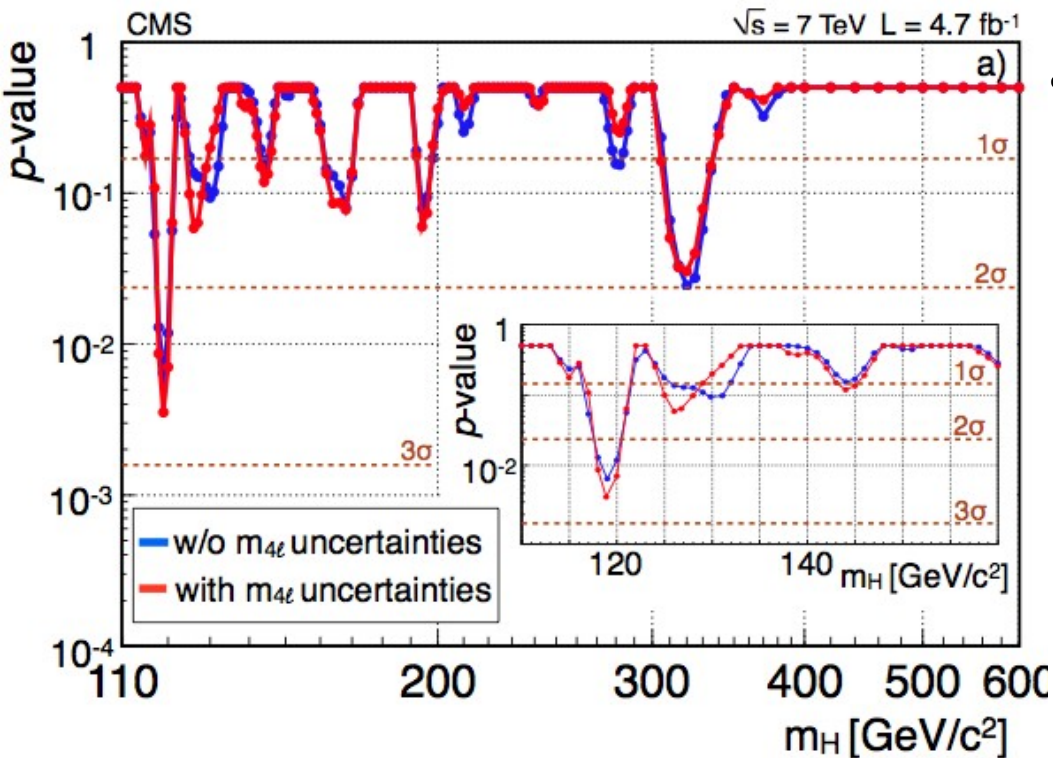
Higgs
mass

Statistical interpretation

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- Interpret excesses as p-values (probability that the background fluctuates upward)

→ Compare the result taking into account or not the uncertainty on $M(4l)$



- Largest excess observed at 119.5 GeV with local significance 2.5σ

→ global significance 1.0σ in the full mass range, 1.6σ in the mass range 100-160 GeV

Next steps: angular discriminant analysis

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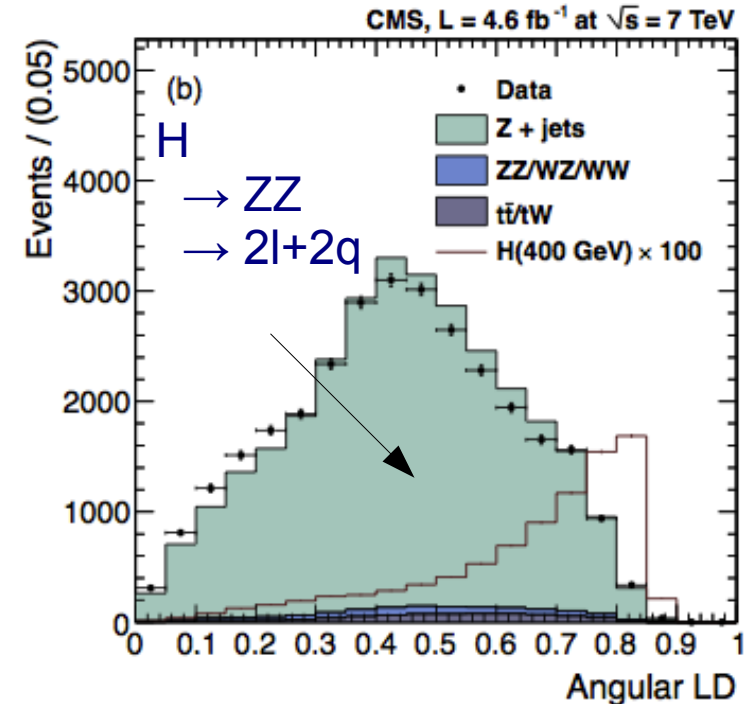
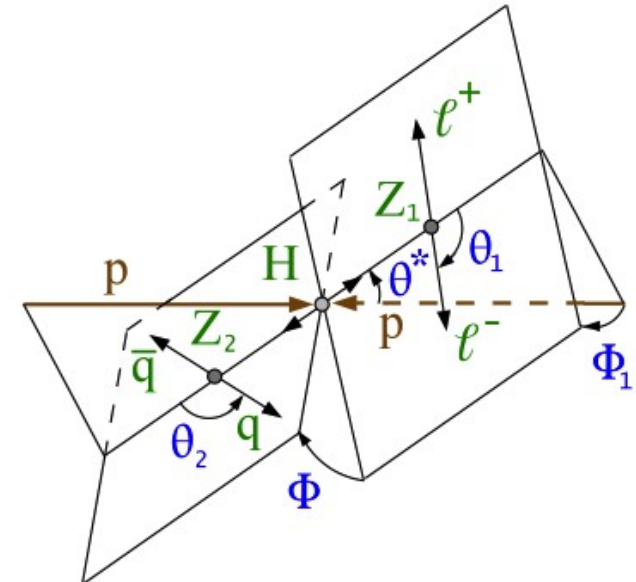
- We have more handles on the signal than just the invariant mass
- Expand the resonant decay in all possible angles

→ Build the expected distribution of the signal and the main backgrounds

→ Use the PDFs to **construct a LLR**

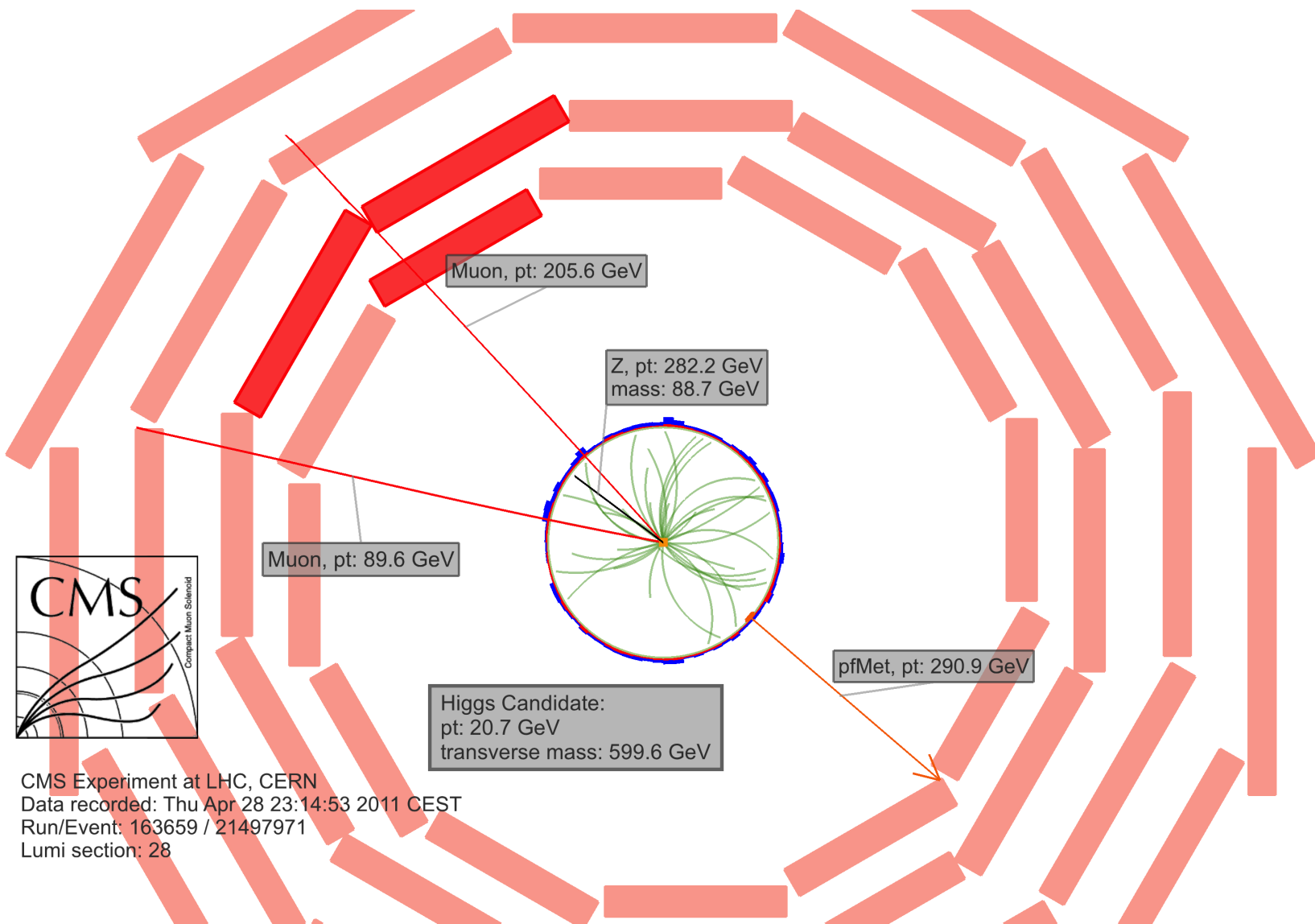
$$LR_S(x_{obs}) \equiv \frac{P_S(x_{obs})}{P_S(x_{obs}) + \sum_i k_i P_i(x_{obs})}$$

- Result can be used in two ways: **select events or find evidence for properties of the signal** (e.g. spin or parity)





High mass search:
 $H \rightarrow ZZ \rightarrow 2l2\nu$



CMS Experiment at LHC, CERN
Data recorded: Thu Apr 28 23:14:53 2011 CEST
Run/Event: 163659 / 21497971
Lumi section: 28

Channel signature

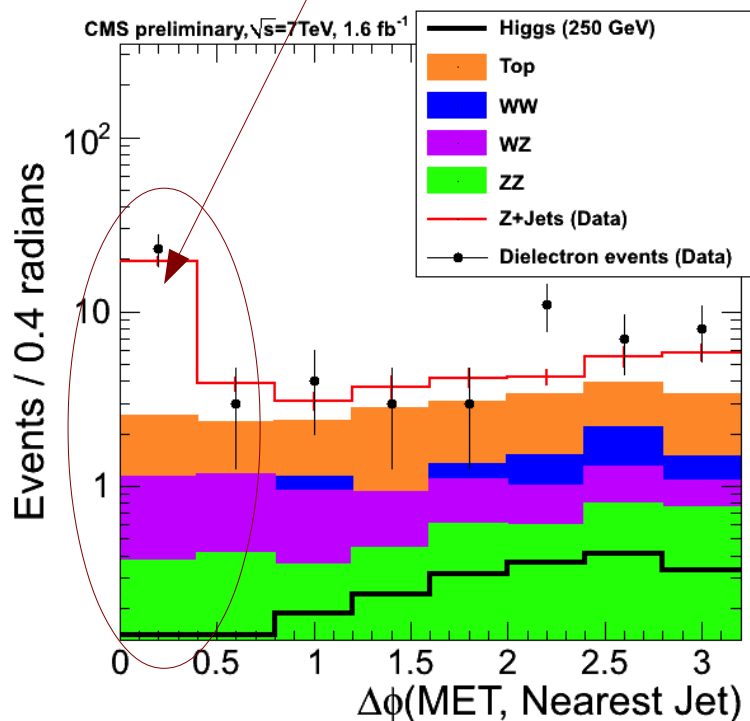
- A dilepton compatible with $Z \rightarrow \ell\ell$ decay, recoiling against nothing

→ $\text{BR}(ZZ \rightarrow 2\ell 2\nu) / \text{BR}(ZZ \rightarrow 4\ell) \approx 6$

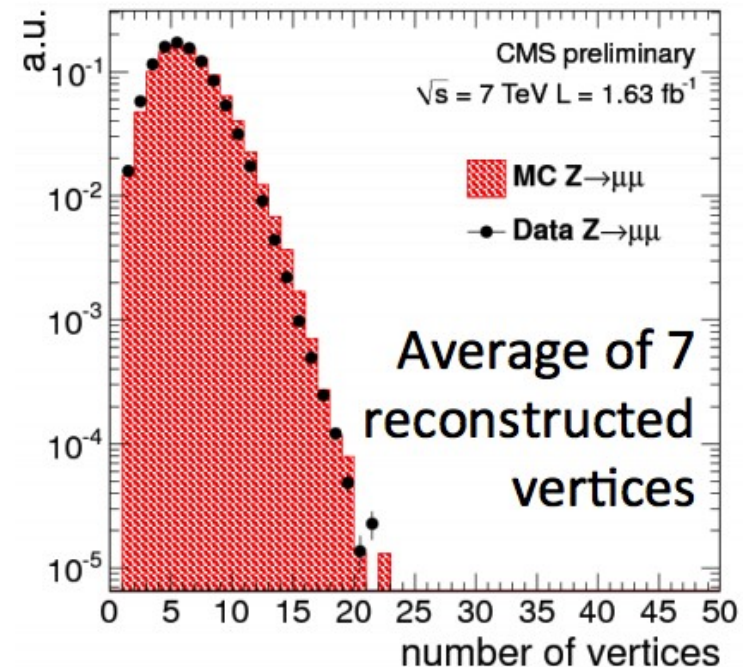
→ Large branching ratio, but also large background contamination

- Need a **robust handle** against two main contaminations

Mis-reconstruction of the Z recoil



Pileup contamination



Missing transverse energy

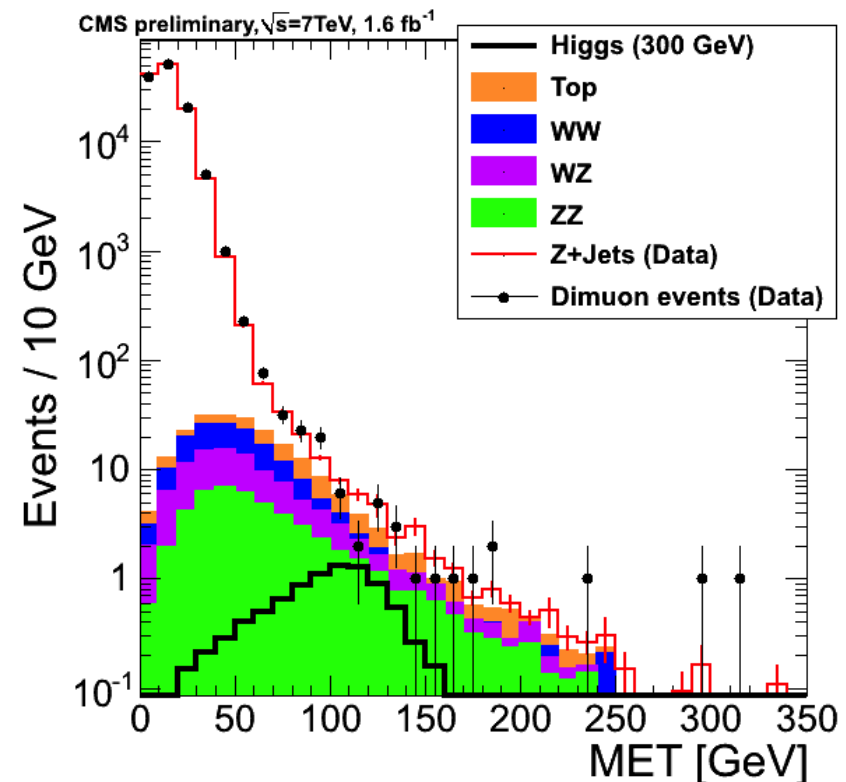
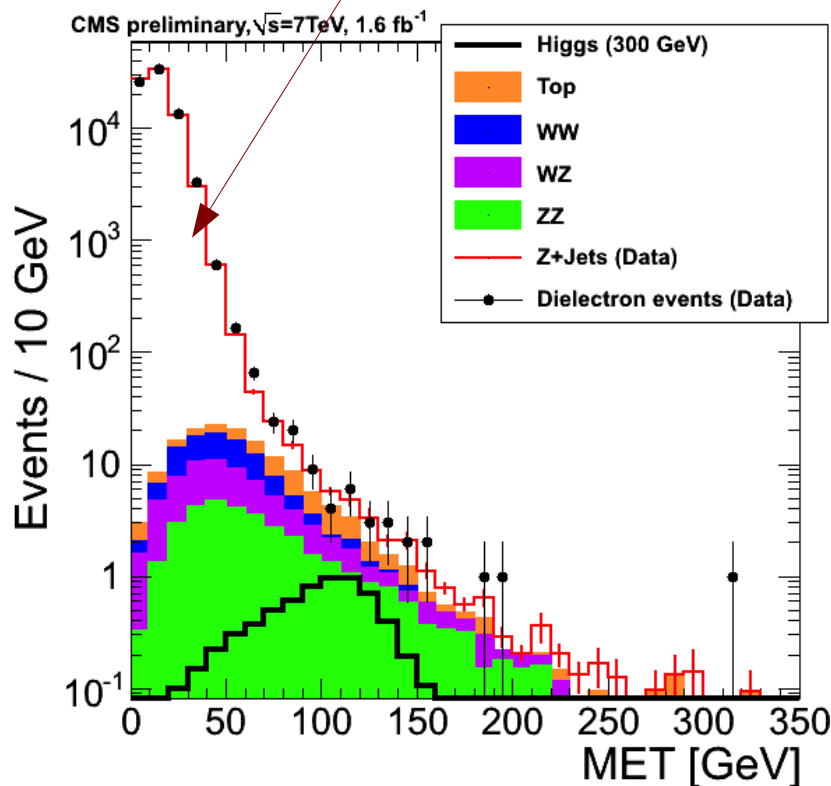
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- Built from the flux of the reconstructed particle momenta:

$$\vec{E}_T^{\text{miss}} = - \sum_{\text{particles}} \vec{p}_{T,i}$$

- The main background is $Z \rightarrow \ell\ell$ production, similar to γ +jets production:

→ Use photon sample and re-weight to match the $Z p_T$ spectrum to derive E_T^{miss} shape



Missing transverse energy

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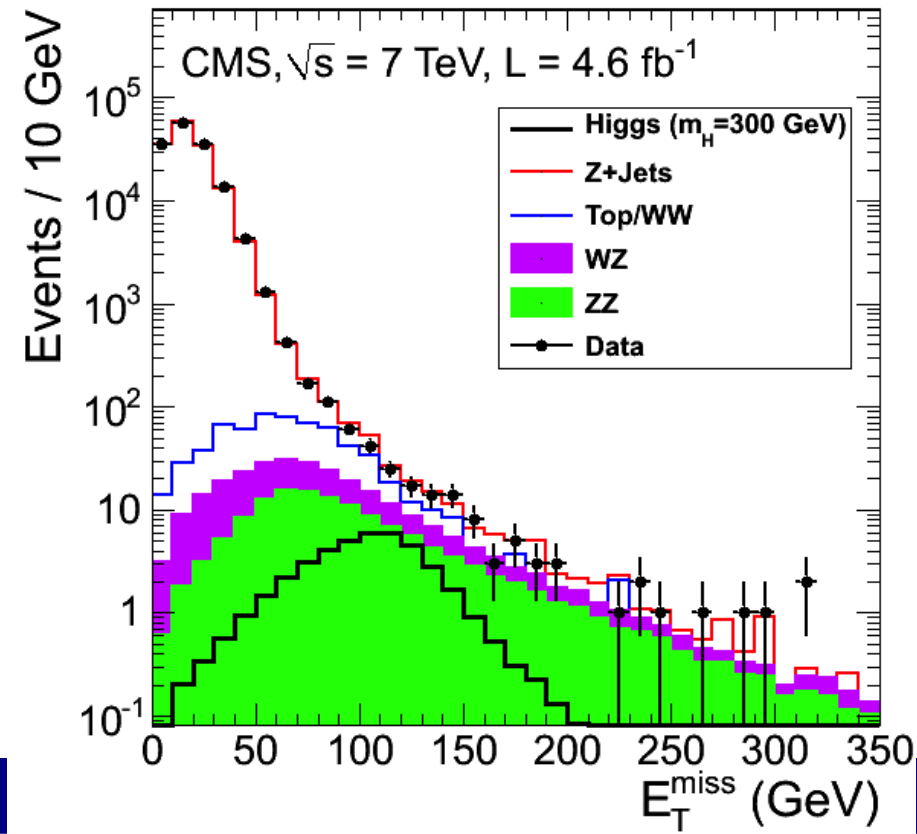
- Built from the flux of the reconstructed particle candidates momenta

$$\vec{E}_T^{\text{miss}} = - \sum_{\text{leptons}} \vec{p}_{T,i} - \sum_{\text{jets}} \vec{p}_{T,i} - \sum_{\text{unclustered}} \vec{p}_{T,i}$$

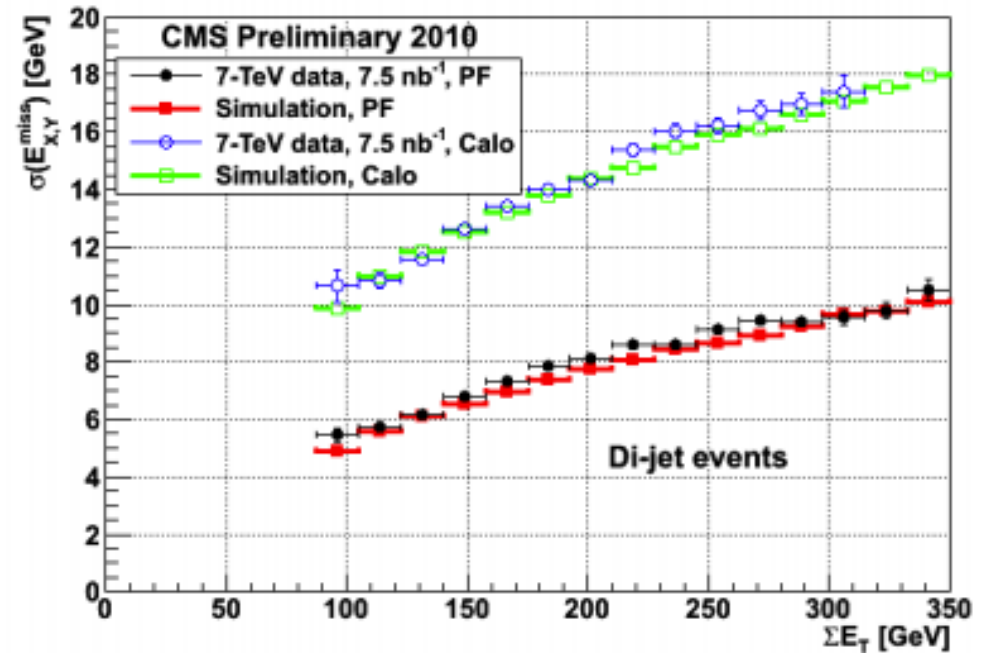
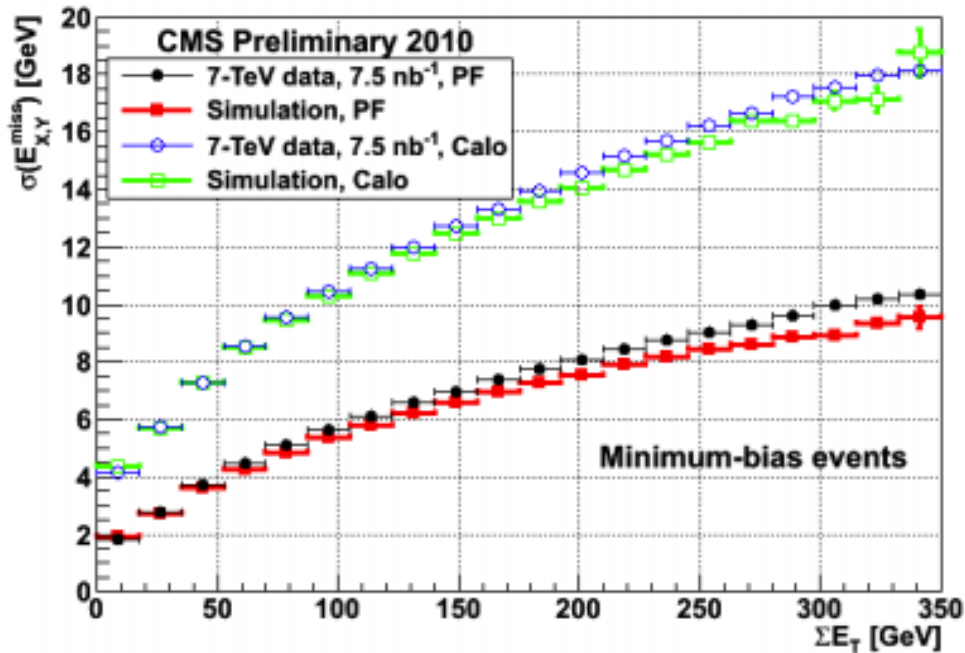
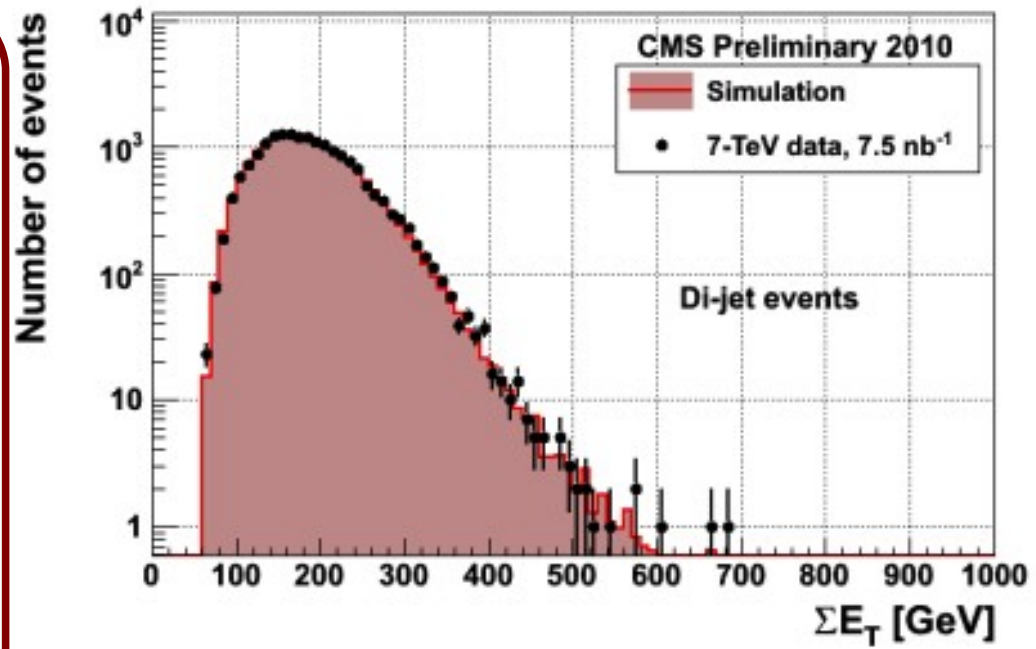
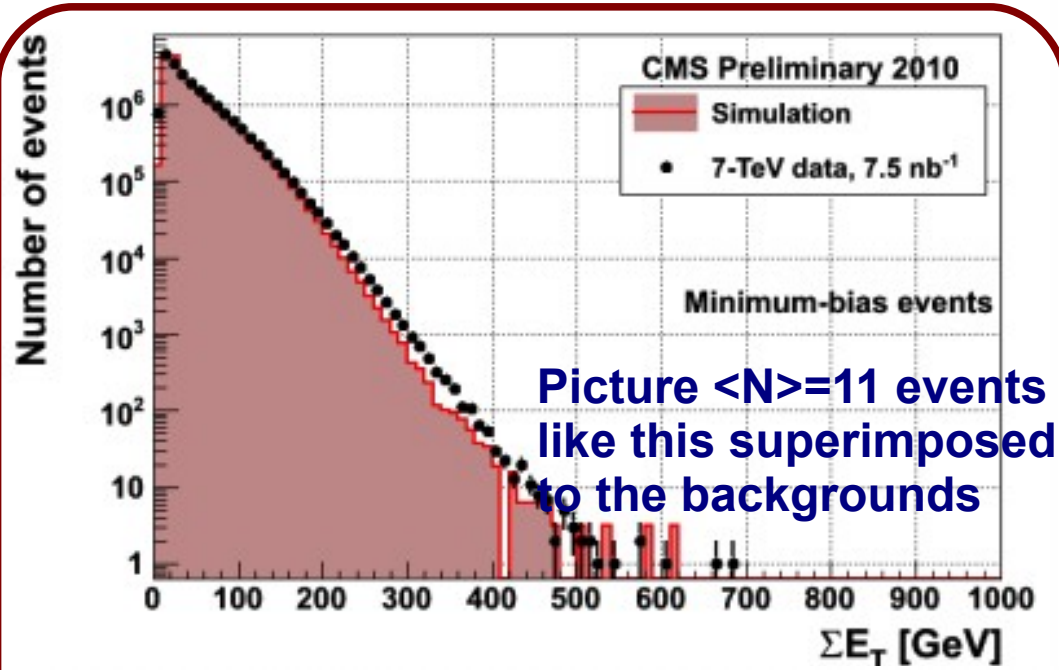
- Charged particles and jets with associated **tracks** can be constrained to the **primary vertex**
- **Neutrals can't be easily associated**
- All must be taken into account, otherwise additional imbalance is found in the event

- E_T^{miss} measurement is furthermore affected by:

- **jet energy scale/resolution** effects
- noise in the calorimeters, dead cells,



E_T^{miss} resolution in minimum bias events



Key distribution: transverse mass

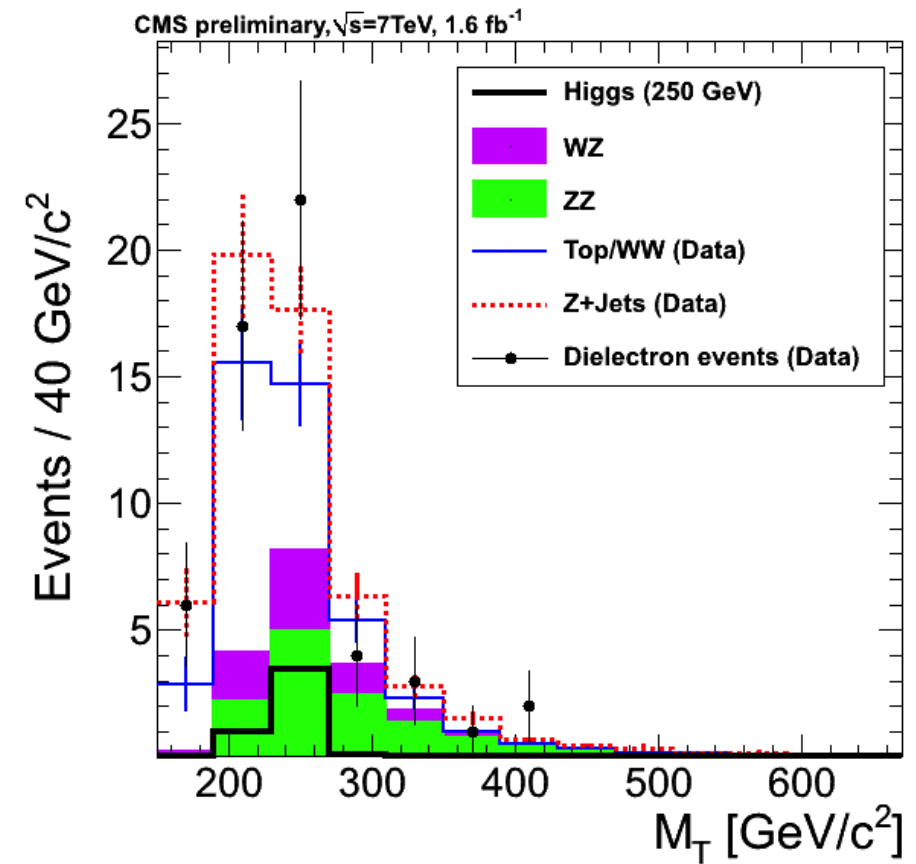
- Due to the presence of two neutrinos the mass the **full kinematics can't be reconstructed** (one degree of freedom left)

- Use the **transverse mass of the dilepton**

+ E_T^{miss} system:

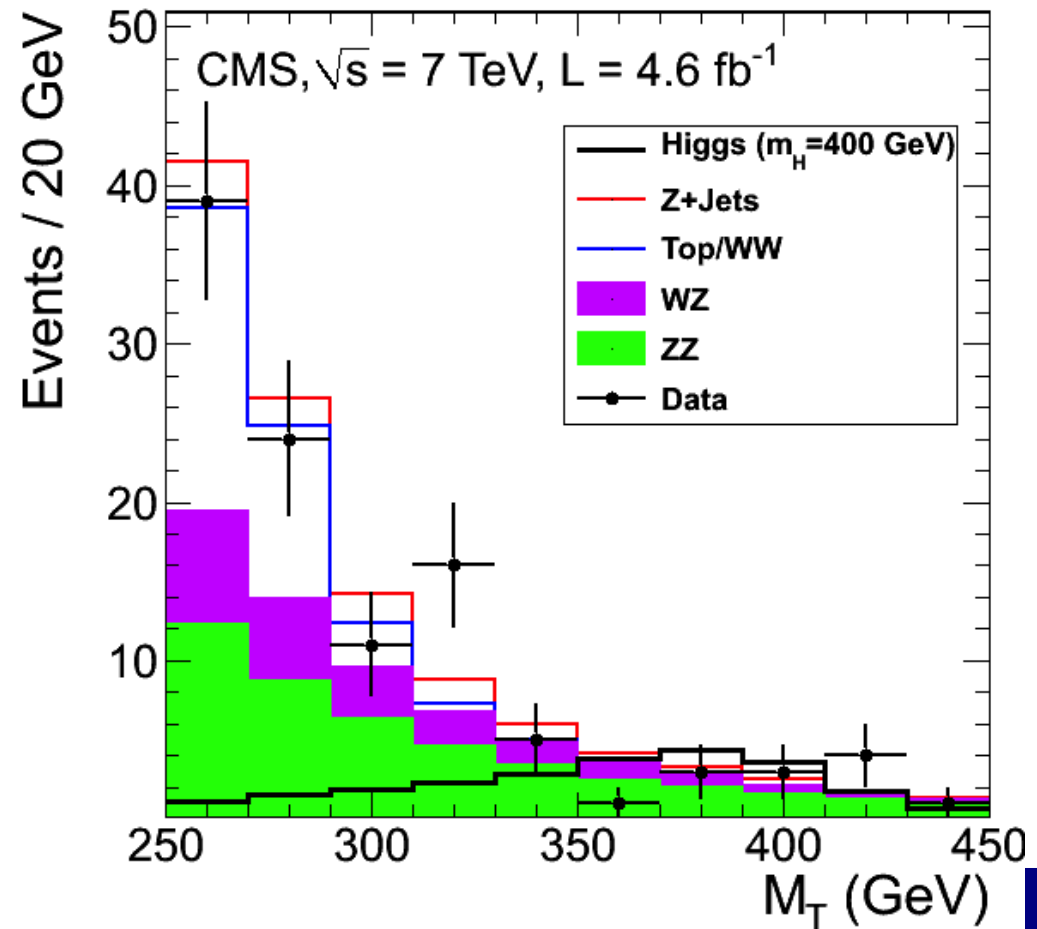
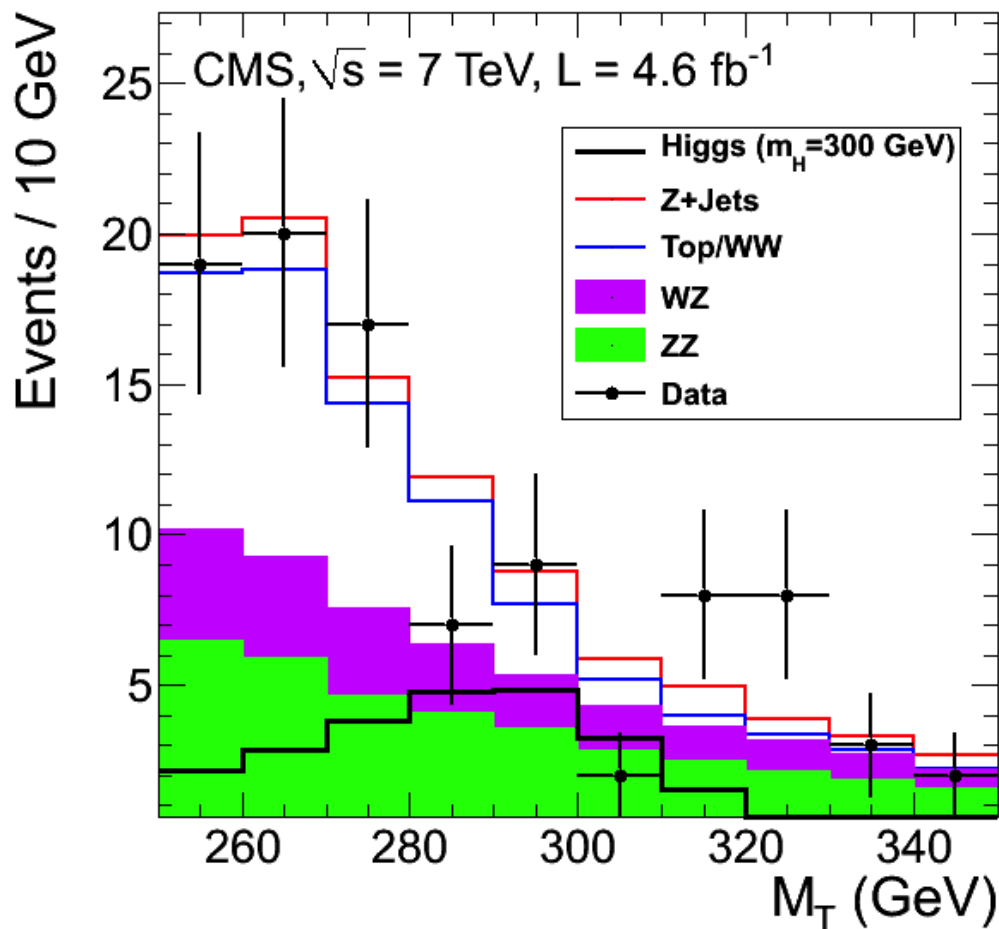
$$M_T^2 = \left(\sqrt{p_T(\ell\ell)^2 + M(\ell\ell)^2} + \sqrt{E_T^{\text{miss}^2} + M(\ell\ell)^2} \right)^2 - (\vec{p}_T(\ell\ell) + \vec{E}_T^{\text{miss}})^2$$

- assuming same mass for the two decay legs
- Lower bound for will be $M_T \sim 2 M_Z$



Transverse mass analysis

- The M_T distribution can be analyzed in two ways:
 - **Count** the number of events in a given region: simple, robust analysis
 - **Analyze the observed shape**: fit components, use sidebands to constrain backgrounds



Final event selection from cut and count

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- **Optimize E_T^{miss} and M_T cuts for best limits** → run several pseudo-experiments before looking at real data, then apply cuts to data
- **Number of selected events is compatible with background expectations**
 - Z+jets modeled directly from the γ +jets sample
 - Non-resonant background extracted from dilepton mass side-band in the $e\mu$ channel

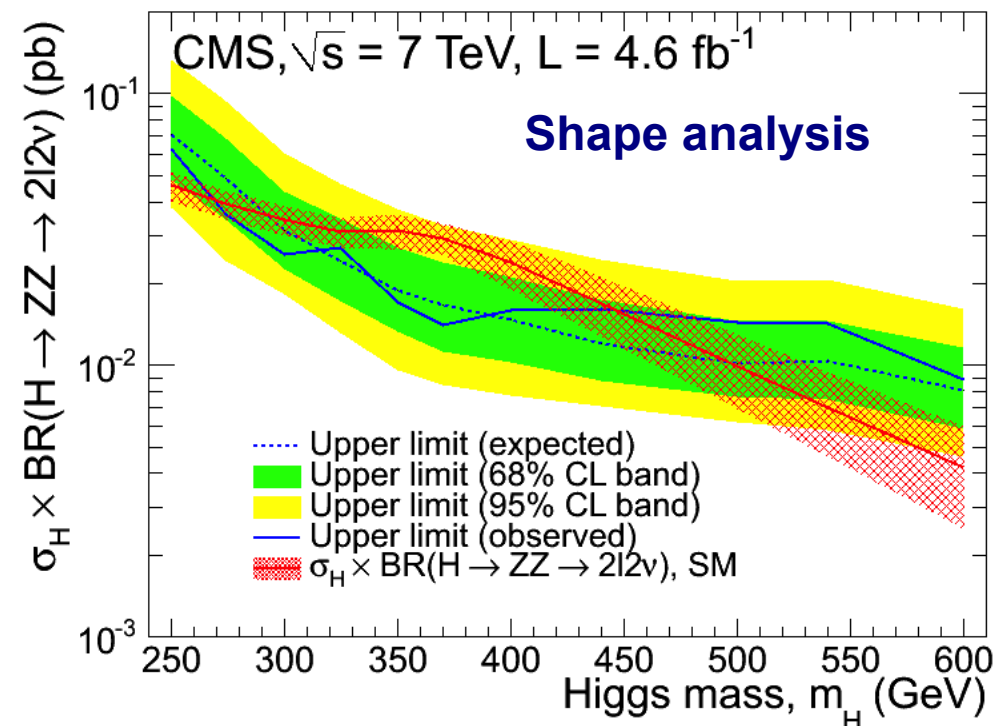
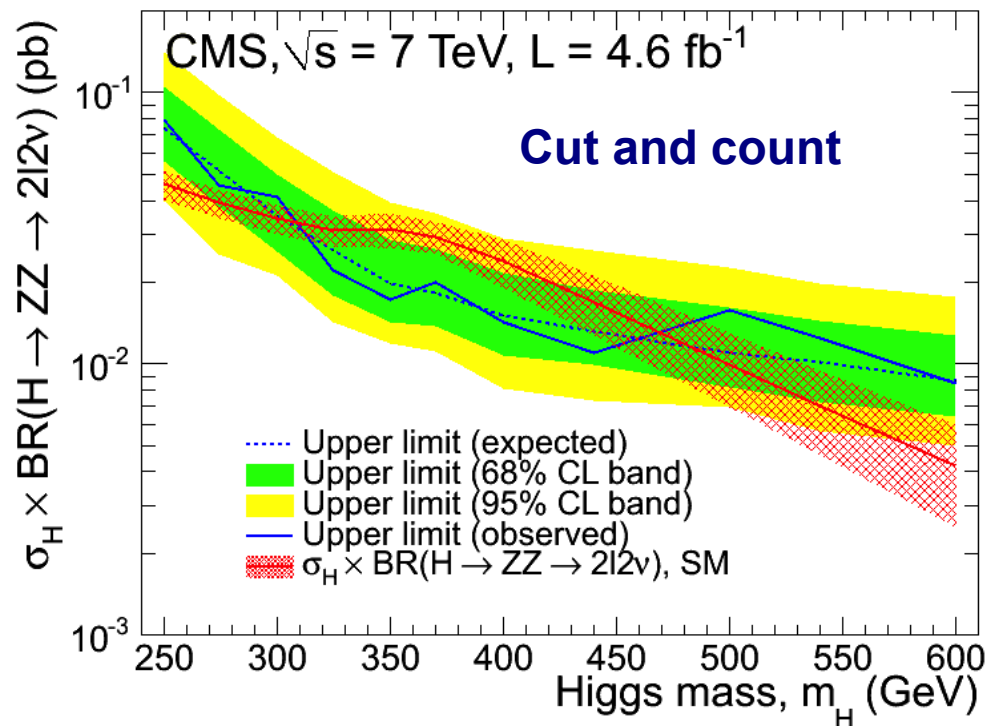
m_H (GeV)	ZZ	WZ	Top/WW/ W+jets/Z→ $\tau\tau$	Z+Jets	Total Background	Expected Signal	Data
250	$36.0 \pm 0.2 \pm 2.6$	$24.0 \pm 0.3 \pm 2.0$	$65.0 \pm 3.8 \pm 5.8$	15.0 ± 15.0	$140.0 \pm 3.8 \pm 16.0$	22.0 ± 2.2	142
300	$23.0 \pm 0.2 \pm 1.7$	$13.0 \pm 0.2 \pm 1.1$	$18.0 \pm 1.1 \pm 3.0$	6.3 ± 6.3	$60.0 \pm 1.1 \pm 7.3$	21.0 ± 2.1	64
350	$16.0 \pm 0.1 \pm 1.1$	$7.0 \pm 0.2 \pm 0.6$	$2.0 \pm 0.1 \pm 1.0$	4.1 ± 4.1	$29.0 \pm 0.3 \pm 4.4$	21.0 ± 2.5	26
400	$12.0 \pm 0.1 \pm 0.9$	$4.6 \pm 0.1 \pm 0.4$	< 1.1	2.7 ± 2.7	$19.0 \pm 0.2 \pm 2.9$	17.0 ± 2.0	18
500	$7.5 \pm 0.1 \pm 0.5$	$2.0 \pm 0.1 \pm 0.2$	< 1.1	1.4 ± 1.4	$11.0 \pm 0.1 \pm 1.5$	7.4 ± 1.3	14
600	$3.9 \pm 0.1 \pm 0.3$	$0.8 \pm 0.1 \pm 0.1$	< 1.1	0.6 ± 0.6	$5.3 \pm 0.1 \pm 0.7$	2.9 ± 0.7	5

Results

- Set limits on cross section for resonant $ZZ \rightarrow 2l2\nu$ production from Higgs

→ At low mass: Z+jets background overwhelms the signal, hard to probe

→ At high mass: dominated by theory uncertainties in particular in the Higgs mass shape

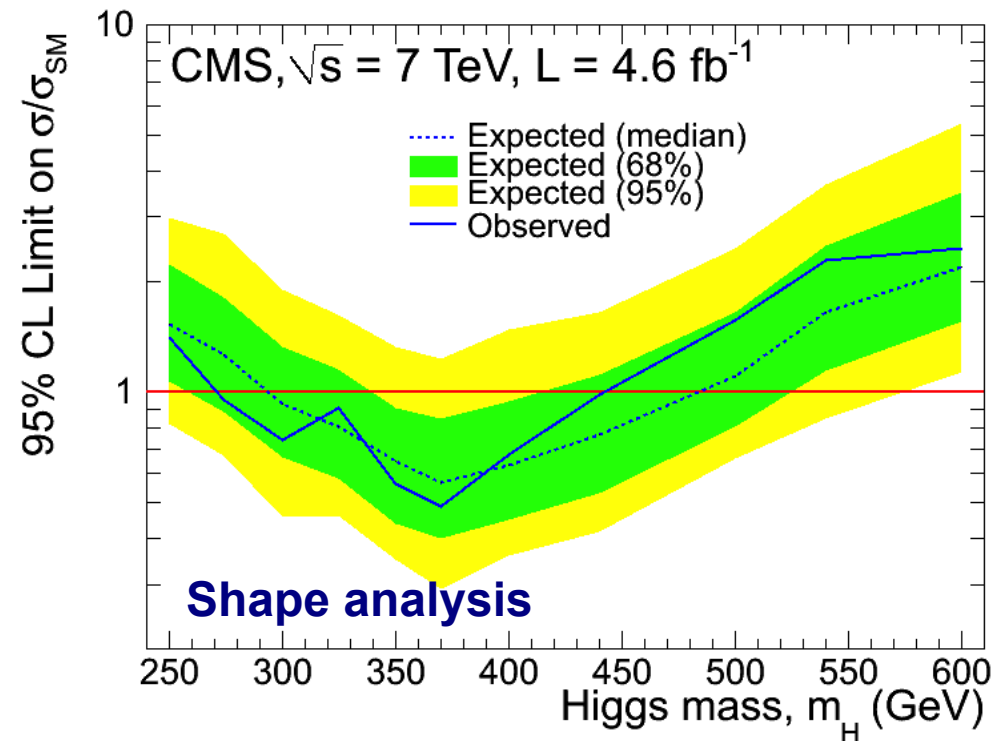
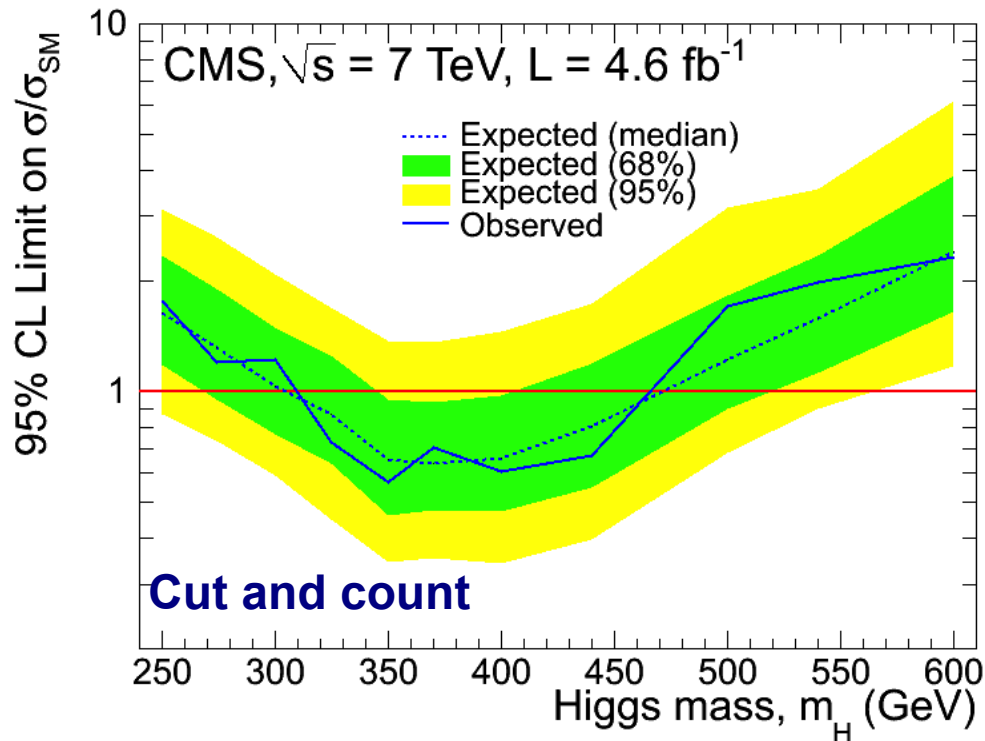


Results - 2

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- Limits on $R = \sigma/\sigma_{\text{SM}}$ at 95% CL

- no particular excess in the mass range analyzed: all is compatible with background only
- **exclude 270-440 GeV/c² mass range**



Conclusions

- Today I have focused on the ZZ process and the search for Higgs in its production
- **ZZ s-channel production is highly suppressed** – deviations can be interpreted as

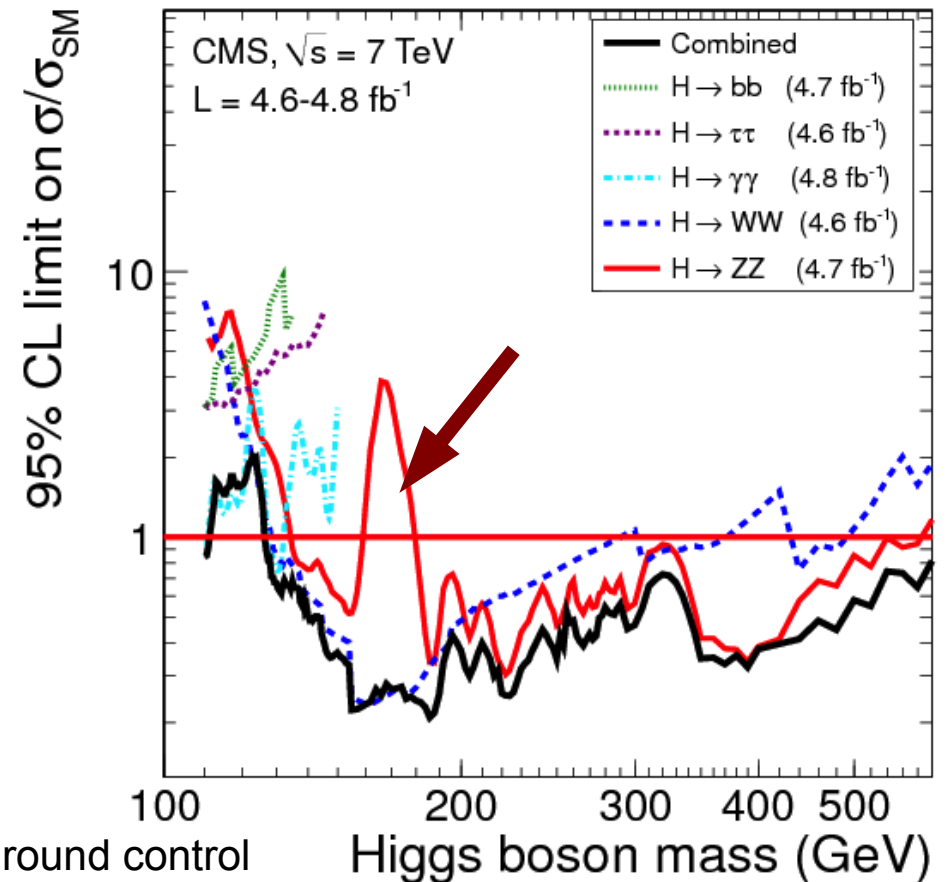
- anomalous triple gauge couplings
- resonant production of the ZZ channel

- **ZZ → 4l is the golden channel**

- Optimal resolution, but lowest branching ratio
- Slight excess $\sim 119.5 \text{ GeV}/c^2$

- **ZZ → 2l2v leads the search at high mass**

- Larger branching ratio but needs robust background control
- Observations are compatible with background only hypothesis



End of Lecture II on Higgs Physics

- ATLAS Collaboration, “*Expected Performance of the ATLAS Experiment: Detector, Trigger and Physics*”, CERN-OPEN-2008-020
- CMS Collaboration, “*Search for the standard model Higgs boson in the decay channel $H \rightarrow ZZ \rightarrow 4l$ in pp collisions at $s^{1/2} = 7$ TeV*”, *arXiv:1202.1997*
- CMS Collaboration, “*Search for the standard model Higgs boson in the $H \rightarrow ZZ \rightarrow 2l2\nu$ channel in pp collisions at $s^{1/2} = 7$ TeV*”, *arXiv:1202.3478*
- CMS Collaboration, “*Search for a Higgs boson in the decay channel $H \rightarrow ZZ^{(*)} \rightarrow qq ll$* ”, *arXiv:1202.1416*

Setting limits on Higgs production

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- **When no excess is observed the strategy is to set limits on $\sigma(H)$**

- Assess from data what is the allowed signal strength i.e. $\mu = \sigma/\sigma_{SM}$

- We measure the compatibility of the data with the signal hypothesis using a test statistics

- **Likelihood and test statistics definition**

- The data vs S+B hypothesis is tested with a likelihood

$$\mathcal{L}(\text{data} | \mu, \theta) = \text{Poisson}(\text{data} | \mu \cdot s(\theta) + b(\theta)) \cdot p(\tilde{\theta} | \theta)$$

Signal
expected

background
expected

nuisance paramters:
uncertainty on rates,
shapes, etc.

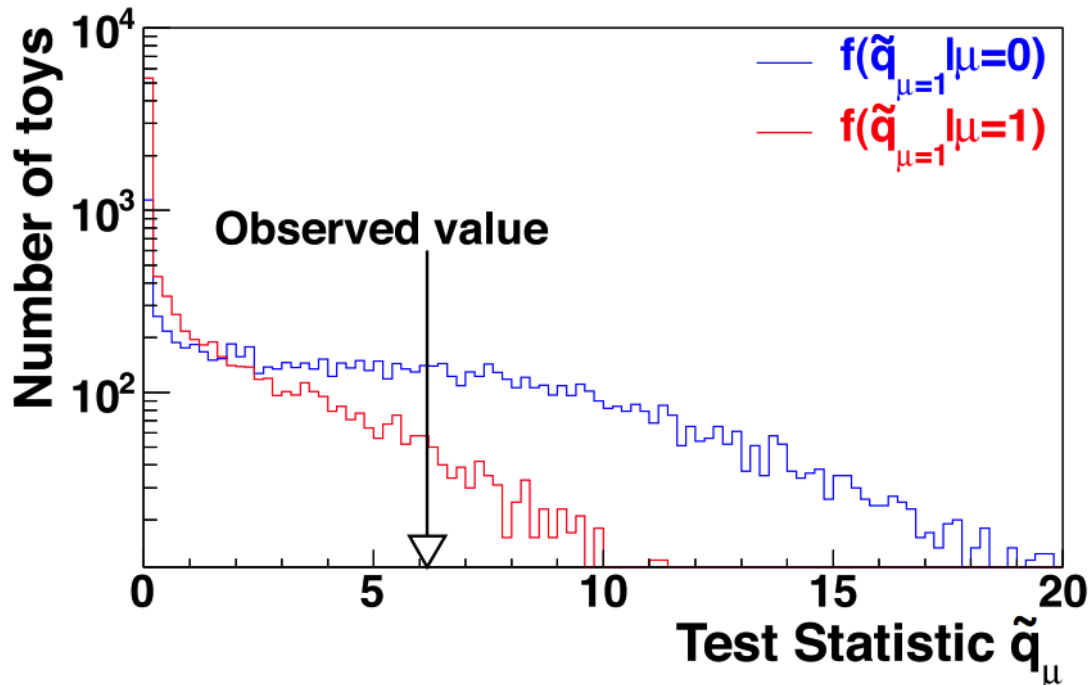
- Signal+background and background only hypothesis tested with

$$\tilde{q}_\mu = -2 \ln \frac{\mathcal{L}(\text{data} | \mu, \hat{\theta}_\mu)}{\mathcal{L}(\text{data} | \hat{\mu}, \hat{\theta})}$$

maximize likelihood

Setting limits – CL_s method

- In data we **compute** the **observed value of the test statistics** and find the best values of all nuisance parameters to fit background and background only hypothesis $\hat{\theta}_0^{obs}$ and $\hat{\theta}_\mu^{obs}$
- From MC/data-driven expectations we generate pseudo-experiments for each hypothesis



$$CL_s(\mu) = \frac{p_\mu}{1 - p_b}$$

Probability that
S+B test statistics
exceeds the
observed value in
S+B hypothesis

Similar, for
background
only
hypothesis

- If $CL_s < 5\%$ signal is excluded at 95%
- We call **upper endpoint**, i.e. μ_{up} , to the **signal strength for which $CL_s=5\%$**