

# Jets

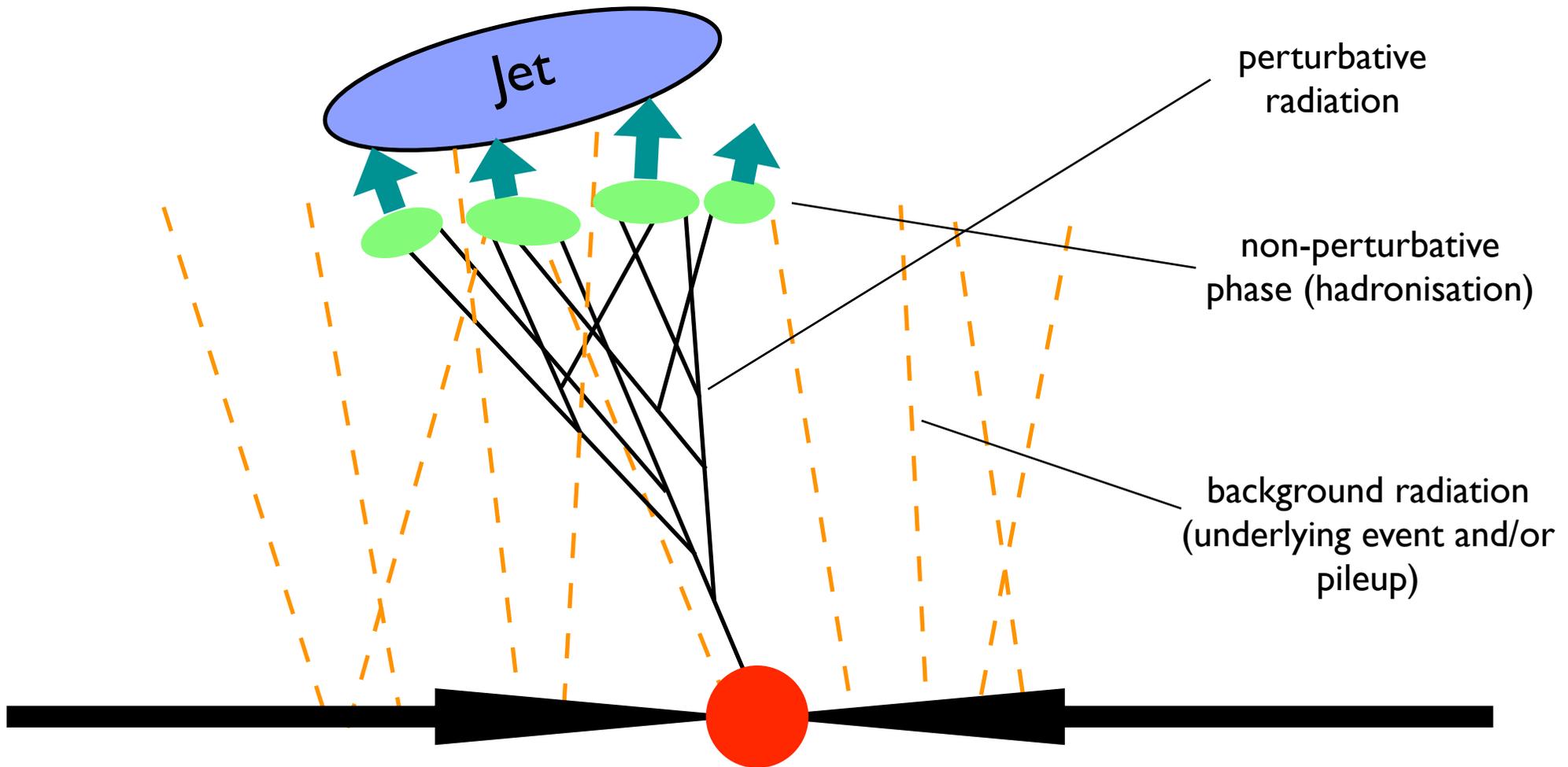
## Lecture 2

Matteo Cacciari

LPTHE Paris

- ▶ **Lecture 2: jet properties and substructure**
  - ▶ *Corrections to a jet transverse momentum*
    - ▶ *jet areas*
    - ▶ *background estimation and subtraction*
  - ▶ *Jet substructure techniques*
    - ▶ *taggers and groomers for boosted objects*

## What contributes to a jet's transverse momentum?

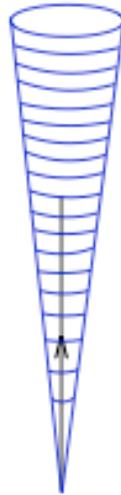


# Contributions to a jet $p_t$

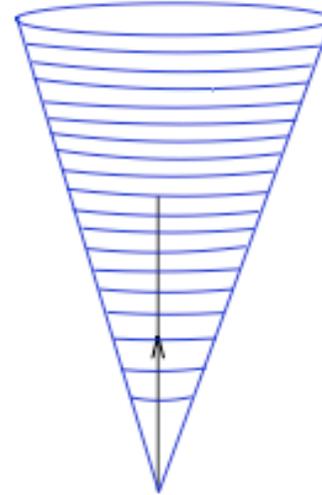
- ▶ Ideally, we'd like a jet to contain all radiation emitted by a parton, i.e. reproduce exactly the parton. This is impossible for a number of reasons:
  - ▶ First and foremost, *a parton is not a physical observable, while a jet is*. There cannot be exact equivalence between them
  - ▶ *A jet has finite extent*, and part of the radiation (perturbative or non-perturbative) emitted by the parton may end up beyond the jet's boundaries, leading to a decrease of the jet momentum (with respect to the parton's)
  - ▶ A jet does not fragment in a vacuum. *Background radiation* (underlying event and/or pileup) can affect its momentum in at least two ways:
    - ▶ Some background radiation will be clustered with the jet, increasing its momentum
    - ▶ Some hard particles (from the parton fragmentation process) may not cluster with the rest of the jet because of the disturbing presence of other nearby particles (they will form other jets with them). Alternatively, hard particles that in the vacuum would not have clustered with the jet will

# Effects of jet 'radius'

Small jet radius



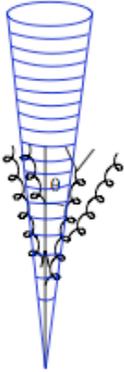
Large jet radius



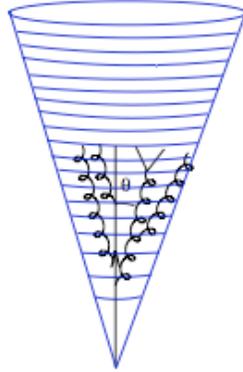
Irrelevant for a single-particle jet

# Effects of jet 'radius'

Small jet radius

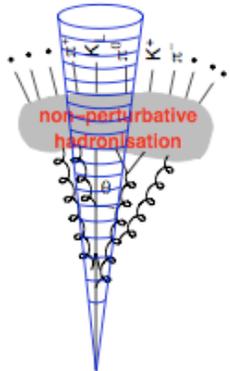


Large jet radius

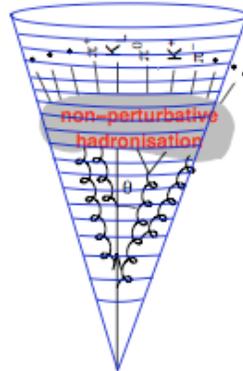


**perturbative** radiation:  
large radius **better** (lose less)

Small jet radius

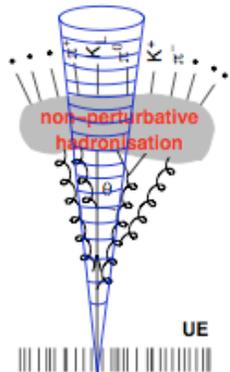


Large jet radius

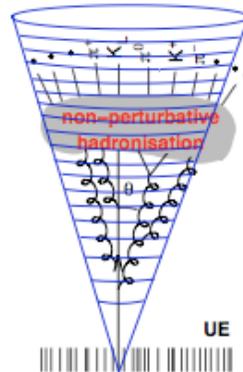


**non-perturbative** hadronisation:  
large radius **better** (lose less)

Small jet radius



Large jet radius



**underlying event:**  
large radius **worse** (capture more)

# R-dependent effects

**Perturbative radiation:**  $\Delta p_t \simeq \frac{\alpha_s (C_F, C_A)}{\pi} p_t \ln R$

**Hadronisation:**  $\Delta p_t \simeq -\frac{(C_F, C_A)}{R} \times 0.4 \text{ GeV}$

**Underlying Event:**  $\Delta p_t \simeq \frac{R^2}{2} \times \begin{matrix} (2.5 & \text{---} & 15 \text{ GeV}) \\ \text{Tevatron} & & \text{LHC} \end{matrix}$

(small-R limit results)

Analytical estimates: [Dasgupta, Magnea, Salam, arXiv:0712.3014](#)

# Jets 'reach'

Algorithmically, a jet is simply a collection of particles

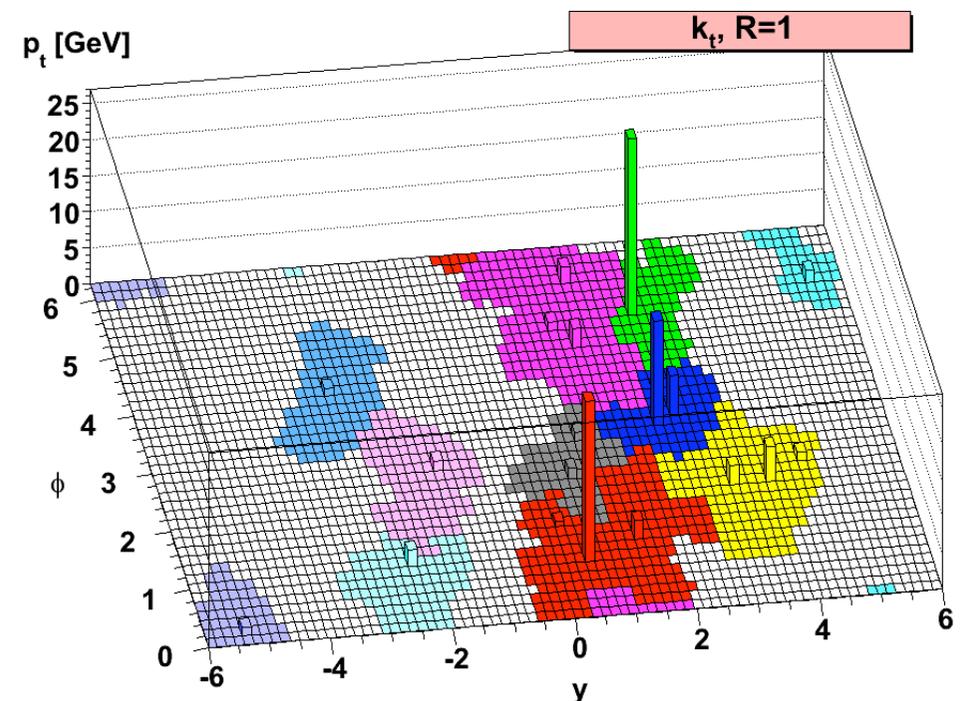
For a number of reasons, it is however useful to consider its **spatial extent**, i.e. given the position of its axis, up to where does it collect particles? What is its shape?

These details are important for a number of corrections of various origin: perturbative, non-perturbative (hadronisation), detector related, etc

Note that the intuitive picture of a jet being a cone (of radius  $R$ ) is **wrong**.

This is what  $k_t$  jets can look like:

(more later about what this plot really means)



# From jet 'reach' to jet areas

MC, Salam, Soyez, 0802.1188

Not one, but three **definitions** of a jet's size:

## ▶ **Passive area**

Place a single soft particle in the event, measure the extent of the region where it gets clustered within a given jet

Reach of jet for **pointlike** radiation

## ▶ **Active area**

Fill the events with many soft particles, cluster them together with the hard ones, see how many get clustered within a given jet

Reach of jet for **diffuse** radiation

## ▶ **Voronoi area**

Sum of areas of intersections of Voronoi cells of jet constituents with circle of radius  $R$  centred on each constituent

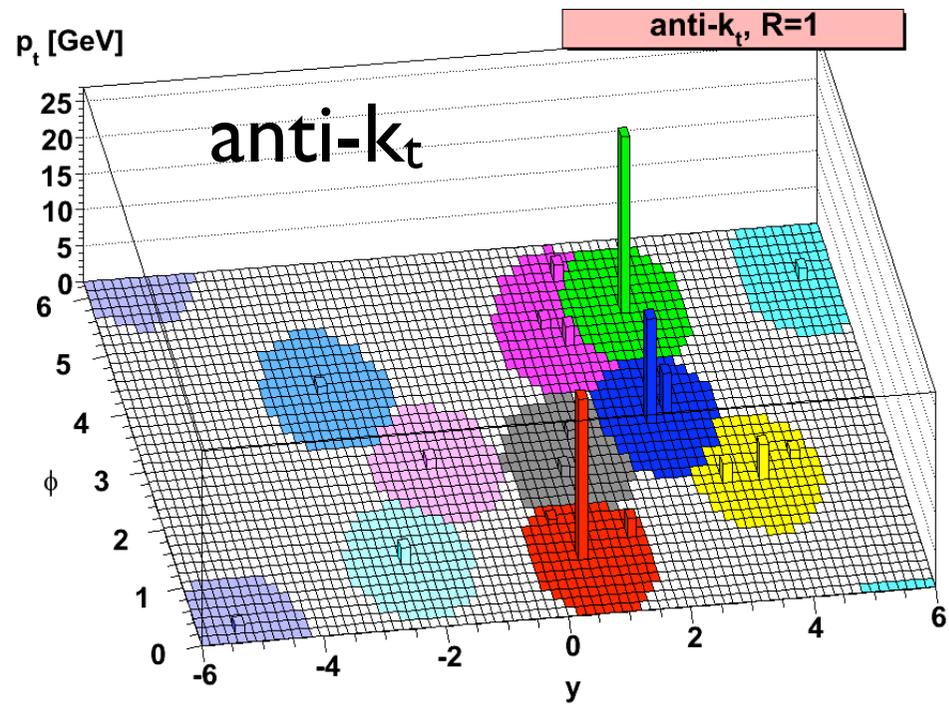
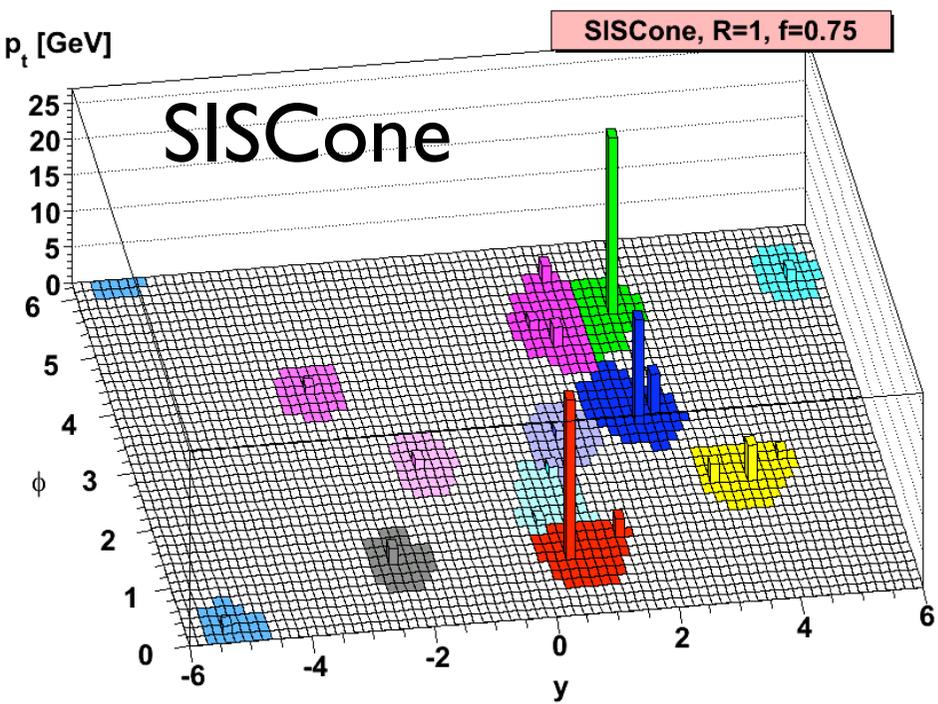
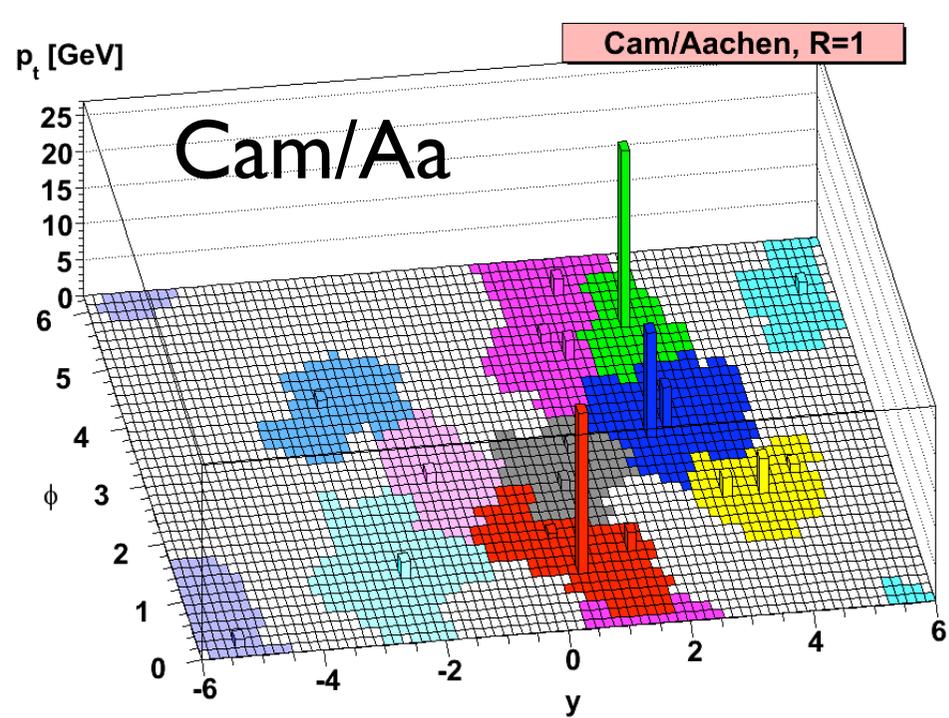
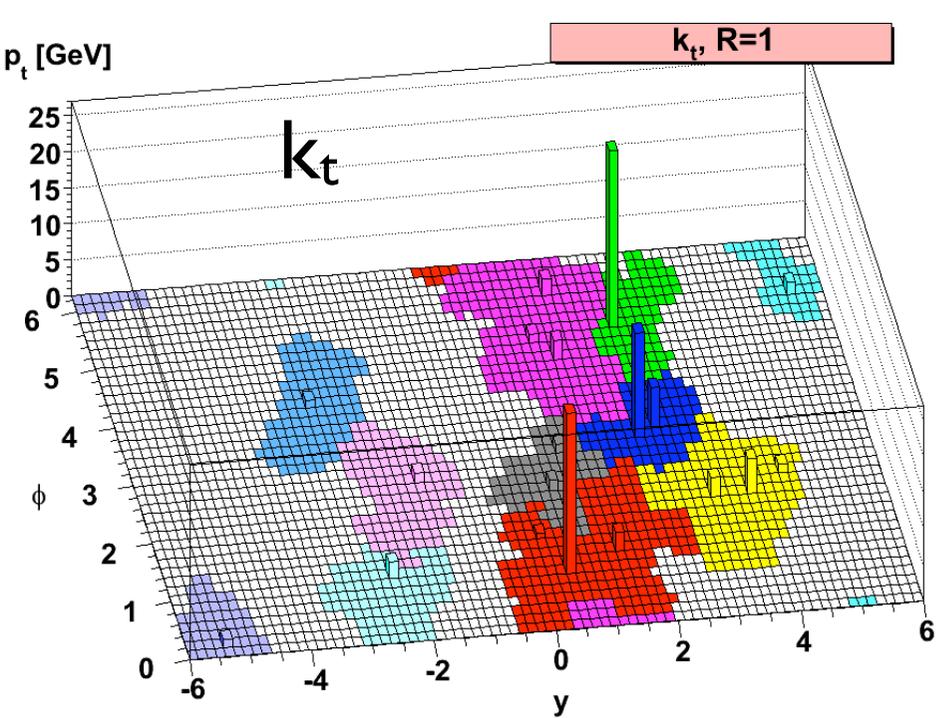
Coincides with passive area for  $k_t$  algorithm

(In the large number of particles limit all areas converge to the same value)

```
/// specify where to place the ghosts, their area, how many times
/// to repeat the clustering
double rapmax = 2.5;
int nrepeat = 1;
double ghost_area = 0.01;
GhostedAreaSpec gas(rapmax, nrepeat, ghost_area);

/// use this configuration to define the area
/// A sensible default for gas is provided
AreaDefinition area_def(active_area, gas);

/// construct cluster sequence with areas
ClusterSequenceArea(input_particles, jet_def, area_def);
....
jets[0].area()
```



# Jet areas: the single hard particle case

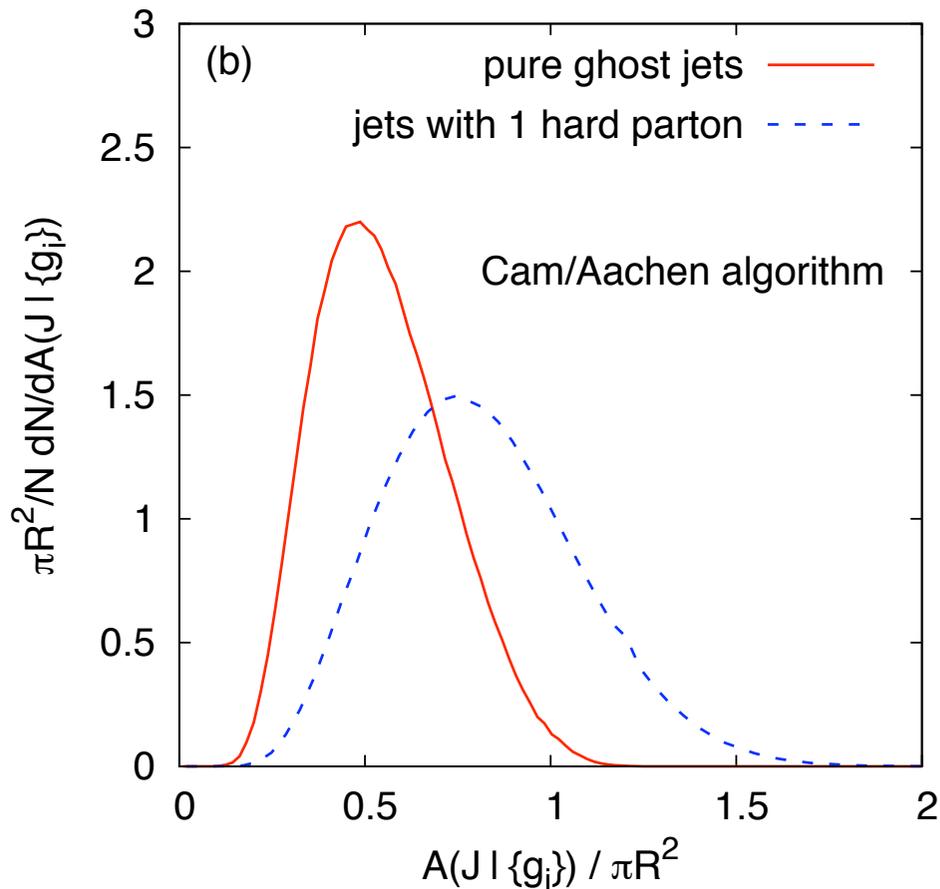
It is worth noting that, for a jet made of a **single hard particle**, while **passive** areas are indeed  $\pi R^2$ , **active** areas are **not**

Active areas	$k_t$	Cam/Aa	SISCone	anti- $k_t$
$\langle A \rangle / \pi R^2$	0.81	0.81	1/4	1

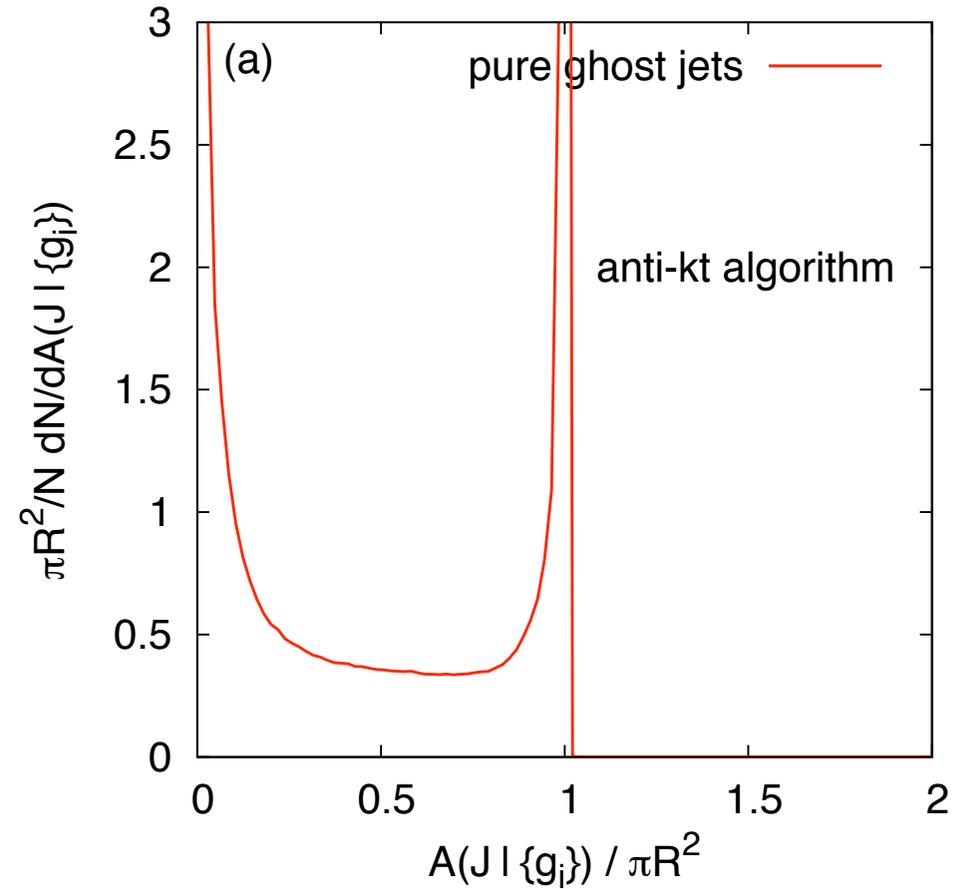
Only **anti- $k_t$**  has the behaviour one would naively expect, i.e. area =  $\pi R^2$

# Active area distributions

## $k_t$ and Cam/AA



## anti- $k_t$



For a roughly uniformly soft background, anti- $k_t$  gives many small jets and many large ones (you can't fill a plane with circles!)

# Jet area: summary

- ▶ Jets CAN have an area, but one must define it
- ▶ The jet (active) area expresses the susceptibility of a jet to contamination from a uniform background
- ▶ Different jet algorithms can have very different area properties:
  - ▶ Jet areas in many algorithms can fluctuate significantly from a jet to another. Isolated hard jets in anti- $k_t$  are one exception
  - ▶ Jet areas can depend on a jet's  $p_t$ , driven by a (calculable) anomalous dimension that is specific to each jet algorithm. Anti- $k_t$  jets are again an exception, in that the anomalous dimension is zero

# Effect of background

How are the hard jets modified by the background?

**Susceptibility**

(how much bkgd gets picked up)

**Jet areas**

**Resiliency**

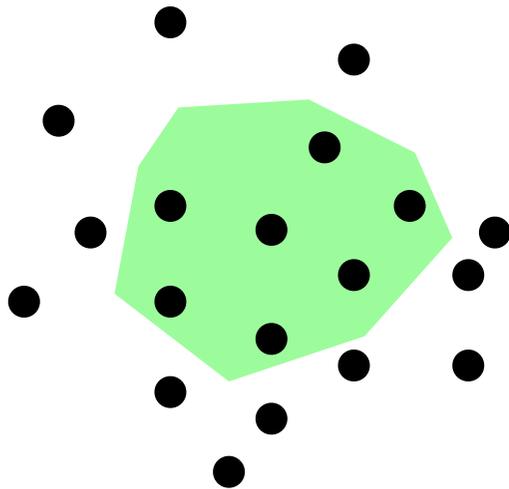
(how much the original jet changes)

**Backreaction**

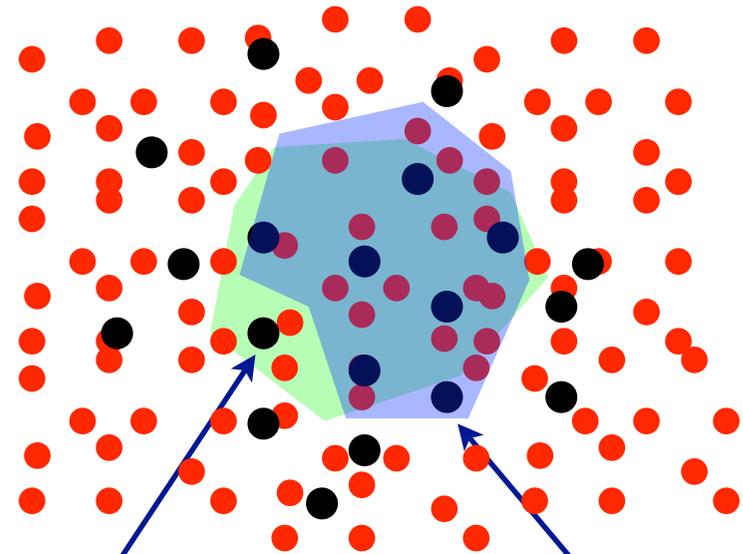
# Resiliency: backreaction

“How (much) a jet changes when immersed in a background”

Without  
background



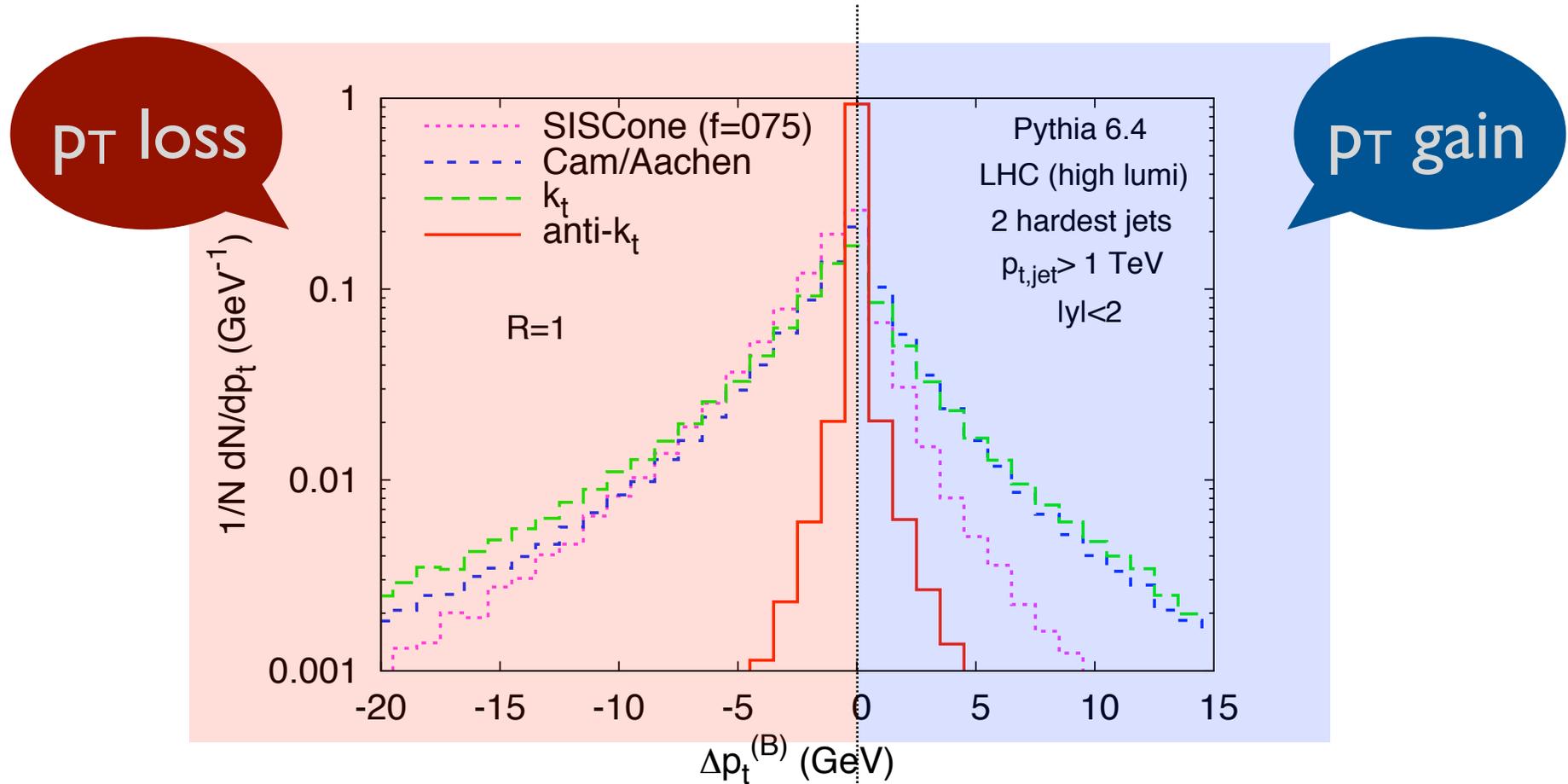
With  
background



Backreaction **loss**

Backreaction **gain**

# Resiliency: backreaction



Anti- $k_t$  jets are much more resilient to changes from background immersion

(NB. Backreaction is a minimal issue in pp background and at large  $p_t$ .  
Can be much more important in Heavy Ion collisions)

# Hard jets and background

## Modifications of the hard jet

$$\Delta p_t = \underbrace{\rho A \pm (\sigma \sqrt{A} + \sigma_\rho A + \rho \sqrt{\langle A^2 \rangle - \langle A \rangle^2})}_{\text{background}} + \underbrace{\Delta p_t^{BR}}_{\text{back-reaction}}$$

Background momentum density (per unit area)

**background**

**back-reaction**

**'susceptibility'**

**'resiliency'**

# Background subtraction

Once  $\rho$  has been measured, it can be used to **correct** the transverse momentum of the hard jets:

$$p_T^{\text{hard jet, corrected}} = p_T^{\text{hard jet, raw}} - \rho \times \text{Area}_{\text{hard jet}}$$

MC, Salam, 0707.1378

If  $\rho$  is measured on an **event-by-event basis**, and each jet subtracted **individually**, this procedure will remove many fluctuations and generally **improve the resolution of, say, a mass peak**

$$\Delta p_t = \rho A \pm (\underbrace{\sigma \sqrt{A}}_{\text{Irreducible fluctuations:}} + \underbrace{\sigma_\rho A}_{\text{uncertainty of the subtraction}} + \underbrace{\rho \sqrt{\langle A^2 \rangle - \langle A \rangle^2}}_{\text{uncertainty of the subtraction}}) + \Delta p_t^{BR}$$

Irreducible fluctuations:  
uncertainty of the subtraction

# Background estimation and subtraction

```
// constructor for a background estimator
JetMedianBackgroundEstimator bge(Selector sel,
                                   JetDefinition jet_def,
                                   AreaDefinition area_def);

// an alternative background estimator
//GridMedianBackgroundEstimator bge(Selector sel, grid_step);

bge.set_particles(input_particles);
....
double rho = bge.rho(jet);    // extract rho estimation
```

```
// define a subtractor
Subtractor sub(&bge);

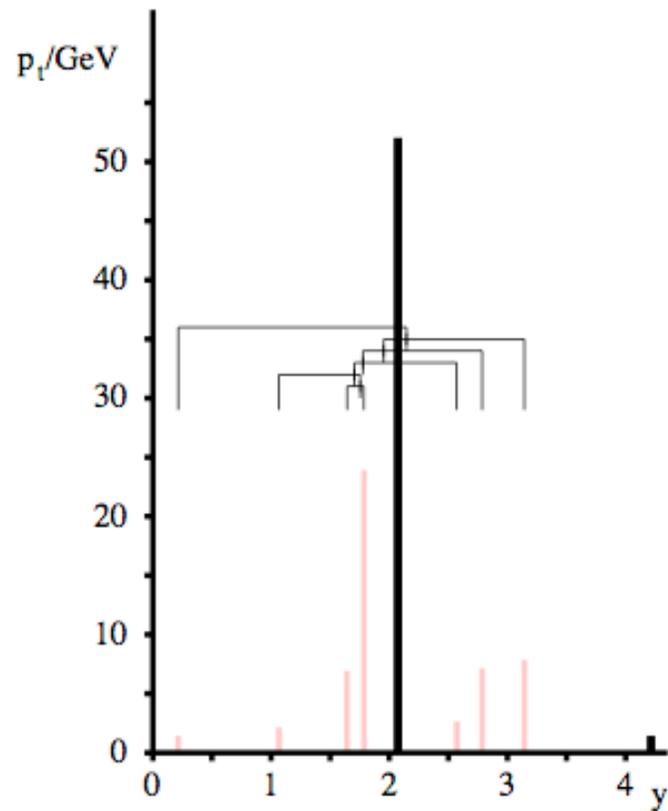
// apply it to a jet (or a vector of jets)
PseudoJet subtracted_jet = sub(jet);
```

# The IRC safe algorithms

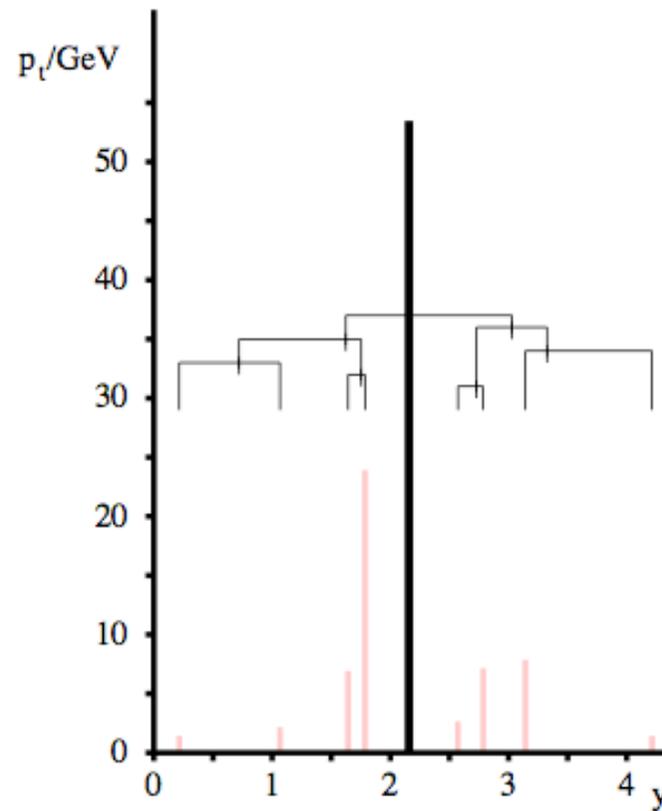
	Speed	Regularity	UE contamination	Backreaction	Hierarchical substructure
$k_t$	☺ ☺ ☺	☂	☂ ☂	☁ ☁	☺ ☺
Cambridge /Aachen	☺ ☺ ☺	☂	☂	☁ ☁	☺ ☺ ☺
anti- $k_t$	☺ ☺ ☺	☺ ☺	☁ / ☺	☺ ☺	✗
SISCone	☺	☁	☺ ☺	☁	✗

# Hierarchical substructure

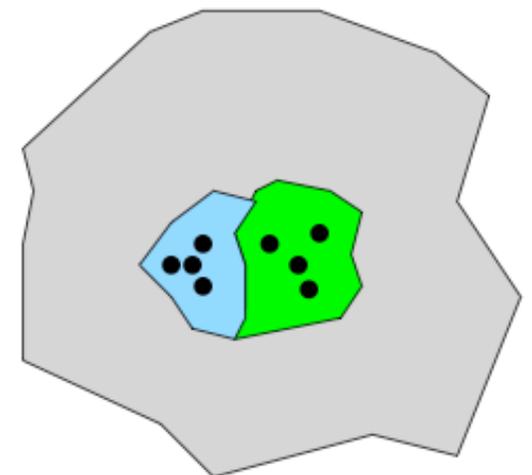
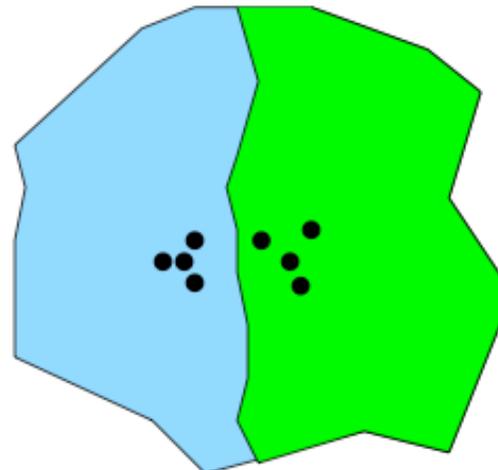
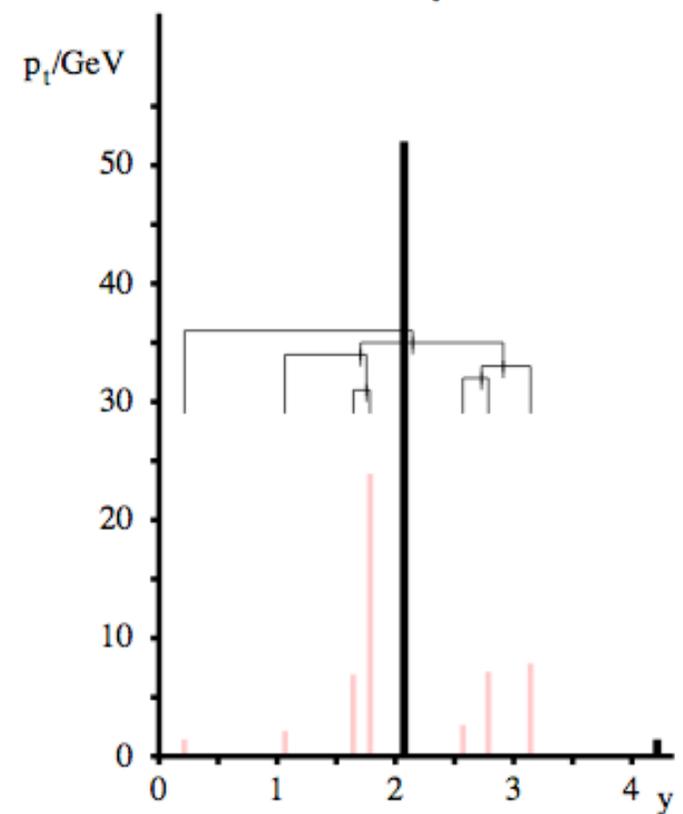
## anti- $k_t$ algorithm



## $k_t$ algorithm



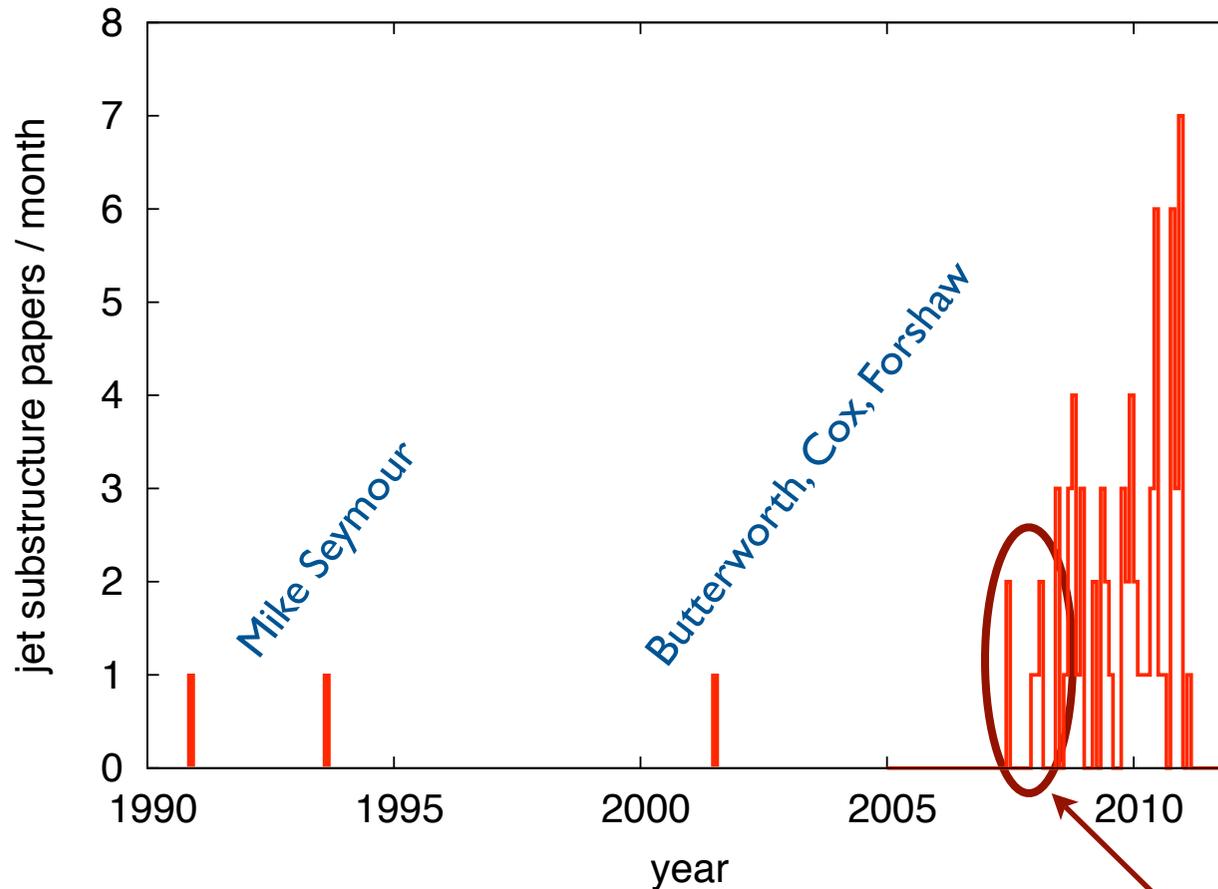
## Cambridge/Aachen



Slide by  
Gavin Salam

# 'Jet substructure' papers in SPIRES

Number of papers containing the words 'jet substructure' and 'LHC'



More than 70 papers after 2008.  
(the histogram dates from early 2011)

Pioneered by M. Seymour in the early '90s, rebooted by BDRS paper

## 15. Jet substructure as a new Higgs search channel at the LHC.

Jonathan M. Butterworth, Adam R. Davison (University Coll. London), Mathieu Rubin, Gavin P. Salam (Paris, LPTHE).

Published in *Phys.Rev.Lett.* 100 (2008) 242001

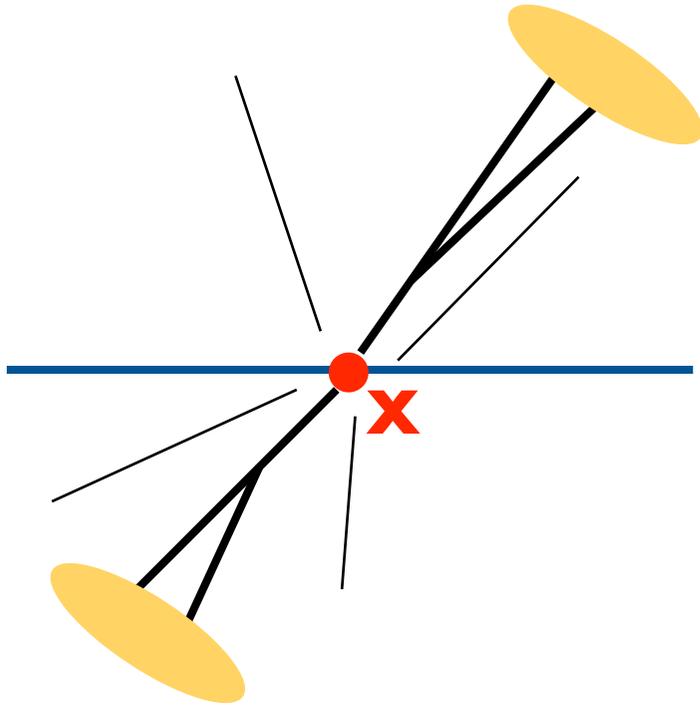
e-Print: [arXiv:0802.2470](https://arxiv.org/abs/0802.2470) [hep-ph]

# Jet substructure

- ▶ The substructure of a jet (i.e. the ability to further resolve smaller components) can be exploited to
  - ▶ *tag a particular structure inside the jet, i.e. a massive particle*
    - ▶ *Examples: Higgs (2-prongs decay), top (3-prongs decay)*
  - ▶ *remove background contamination from the jet or its components*
    - ▶ *Examples: filtering, trimming, pruning*

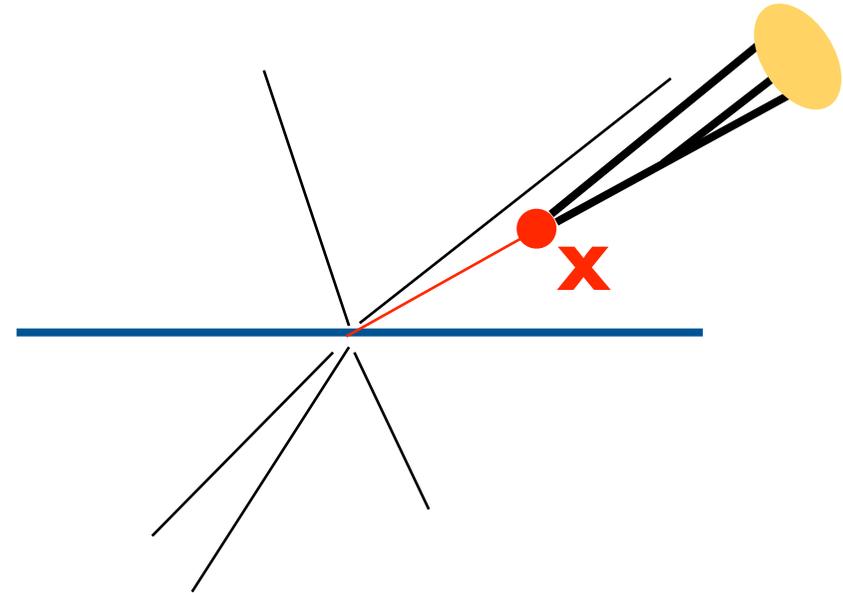
In the following I'll be mainly illustrating the BDRS tagger/filter as a pedagogical example, and also list other approaches

# Why boosted objects



Heavy particle X at **rest**

Easy to resolve jets and calculate invariant mass, but signal very likely swamped by background (eg  $H \rightarrow bb$  v.  $tt \rightarrow WbWb$ )

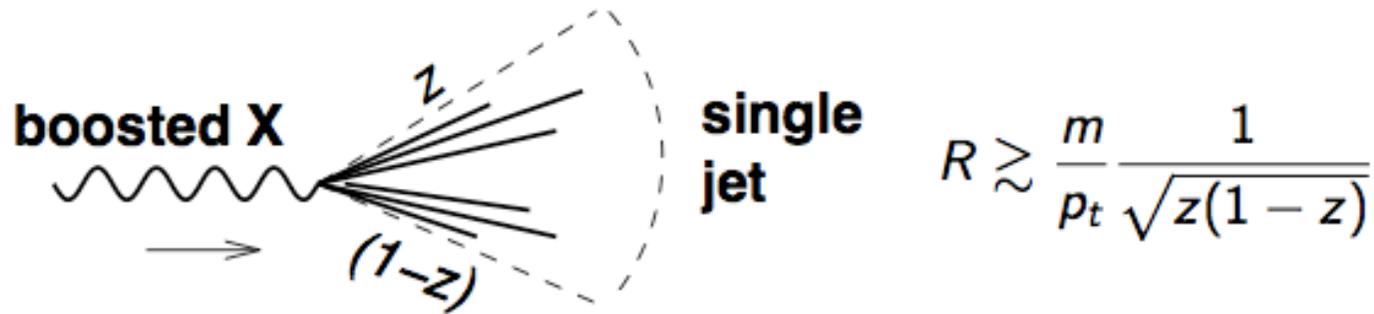


**Boosted** heavy particle X

Cross section very much reduced, but acceptance better and some backgrounds smaller/reducible

# Why substructure

Scales:  $m \sim 100 \text{ GeV}$ ,  $p_t \sim 500 \text{ GeV}$



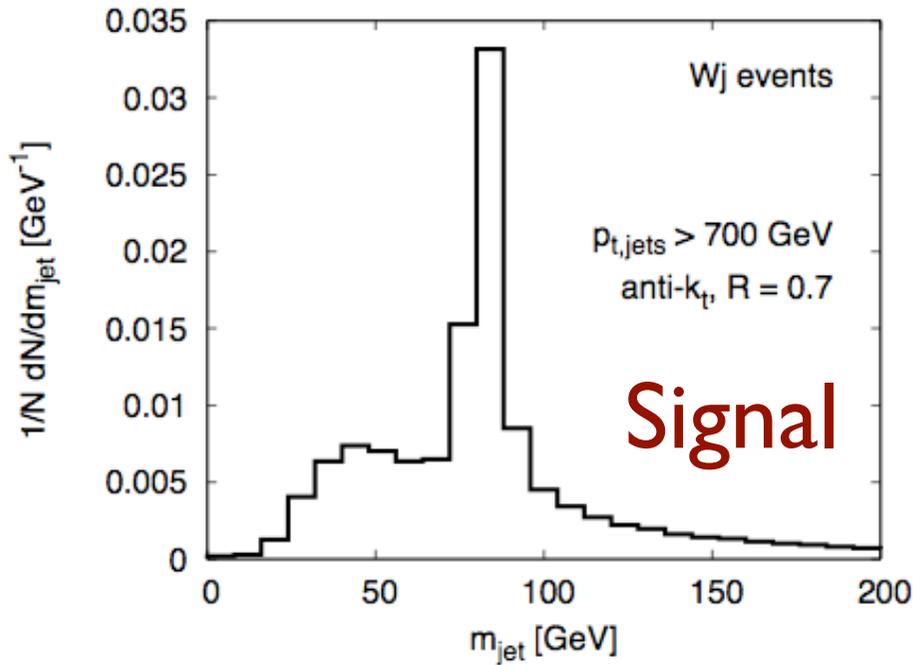
- ▶ need **small R** ( $< 2m/p_t \sim 0.4$ ) to resolve **two** prongs
- ▶ need **large R** ( $> \sim 3m/p_t \sim 0.6$ ) to cluster into a **single** jet

## Possible strategies

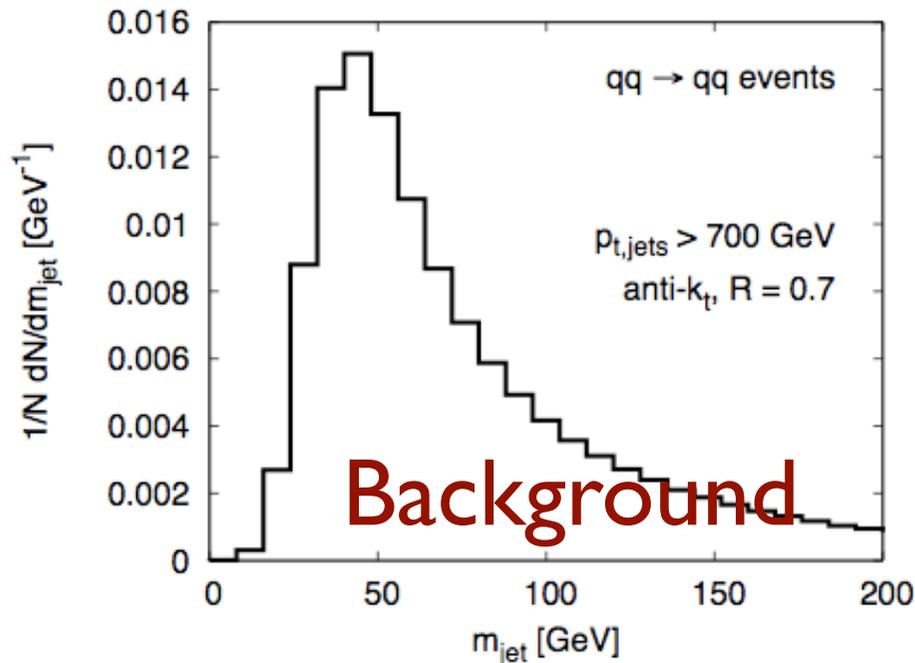
- ▶ Use large R, get a single jet : **background large**
- ▶ Use small R, resolve the jets : **what is the right scale?**
- ▶ **Let an algorithm find the 'right' substructure**

# Mass of a single jet

G. Salam



A heavy object decaying into a single jet naturally gives it a mass...

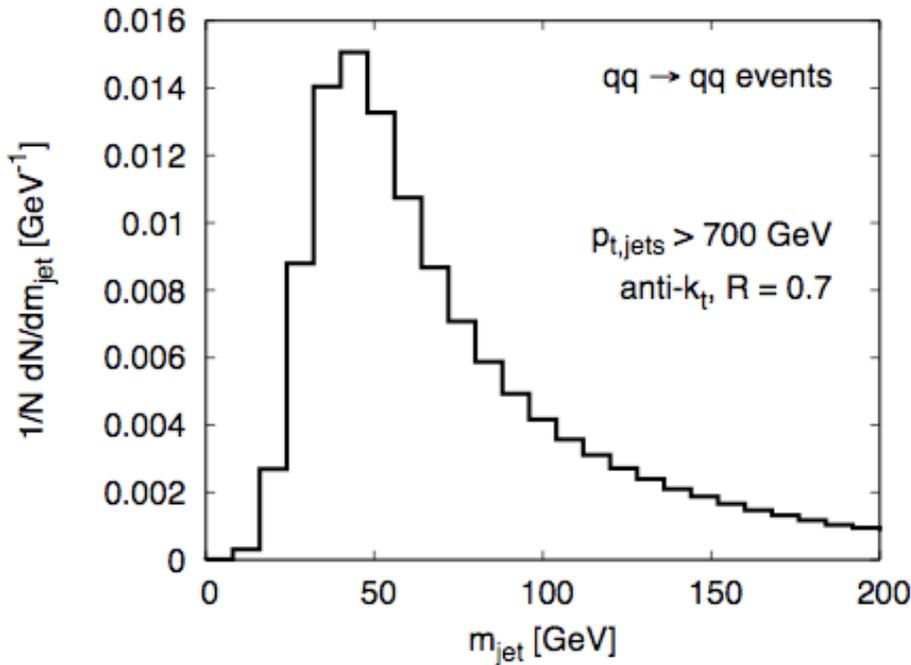


... but pure QCD jets can be massive too:

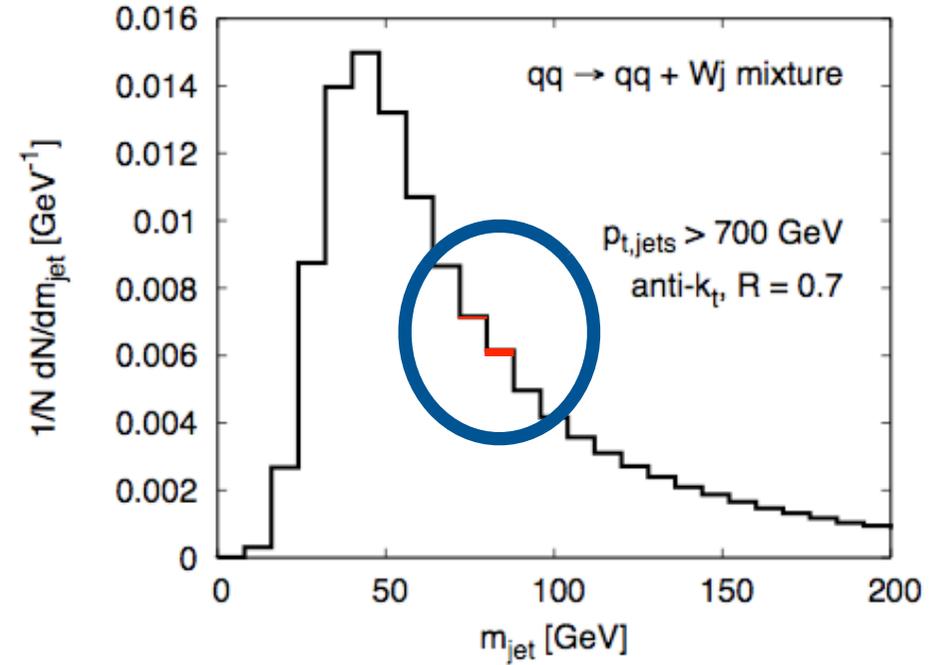
$$\frac{dN}{d \ln m} \sim \alpha_s \ln \frac{p_t R}{m} \times \text{Sudakov}$$

# Mass of a single jet

Summing 'signal' and 'background' (with appropriate cross sections) shows how much the background dominates



**Background only**



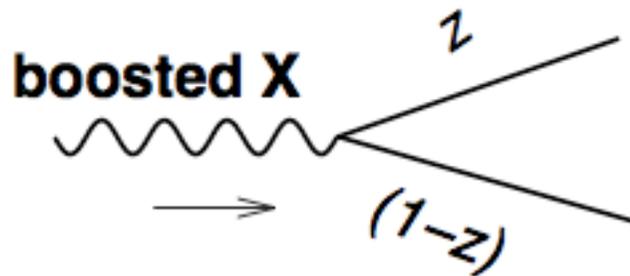
**Signal + background**

**Practically identical**

This means that one can't rely on the invariant mass only.  
An appropriate strategy must be found to reduce the background  
and enhance the signal

# QCD v. heavy decay

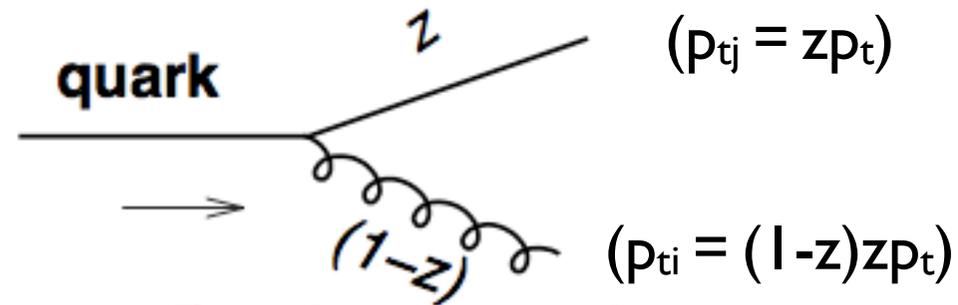
A possible approach for reducing the QCD background is to identify the two prongs of the heavy particle decay, and put a cut on their momentum fraction



**Signal:**

$$P(z) \sim \text{constant}$$

Will split mainly  
**symmetrically**



**Background:**

$$P(z) \sim \frac{1+z^2}{1-z} \quad P(z) \sim \frac{1+(1-z)^2}{z}$$

Will split mainly  
**asymmetrically**

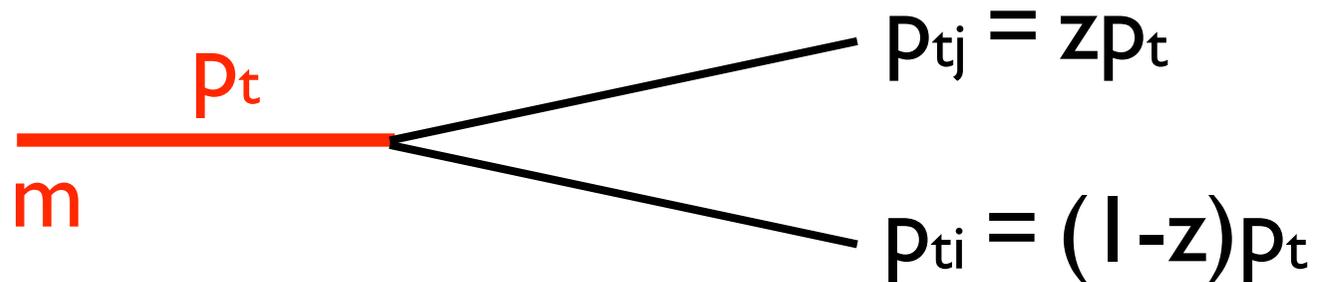
**Potential tagger: asymmetric splitting**

Possibly  
implemented  
via a cut on

$$y = \min(p_{ti}^2, p_{tj}^2) \frac{\Delta R_{ij}^2}{m^2} \simeq \frac{\min(p_{ti}, p_{tj})}{\max(p_{ti}, p_{tj})}$$

# Splittings and distances

Quasi-collinear  
splitting ( $p_{tj} < p_{ti}$ )



Invariant mass:  $m^2 \simeq p_{ti} p_{tj} \Delta R_{ij}^2 = (1-z) z p_t^2 \Delta R_{ij}^2$

$k_t$  distance:  $d_{ij}^{(p_{tj} < p_{ti})} \simeq z^2 p_t^2 \Delta R_{ij}^2 \simeq \frac{z}{1-z} m^2$

For a given mass, the background will have **smaller distance**  $d_{ij}$  than the signal, therefore it will tend to **cluster earlier** in the  $k_t$  algorithm

**Potential tagger: last clustering in  $k_t$  algorithm**

This is where the hierarchy of the  $k_t$  algorithm becomes relevant. QCD radiation is clustered first, and only at the end the symmetric, large-angle splittings due to decays are reclustered

# Alternative algorithms

- ▶ Suppose that for some reasons (which will become clearer later) one does not wish to use the  $k_t$  algorithm
  - ▶ *One must then find a way to determine what the **relevant splitting** (i.e. the one due to the decay, not to QCD radiation) is.*

A possible approach is to use a **Mass-Drop** requirement: the clustering is **progressively undone**, and a splitting is **the relevant one** if **both subjects are much less massive than their combination**

## A generic substructure approach to tagging will

- ▶ **Cluster initially with a large  $R$** , so as to collect all the decay products of a boosted heavy particle into a single jet
- ▶ **Decluster this jet into subjets**, using some conditions to decide when to stop the declustering (i.e. find the ‘relevant splitting’), possibly including kinematical cuts to reduce the QCD background.
  - ▶ *The stopping criterion automatically finds the ‘right size’ for the distance between the two prongs of the heavy particle decay*
- ▶ Optionally add a final ‘cleaning’ procedure to remove as much as possible spurious soft/background radiation

# Generic tagging/grooming

Fat-jet finding

Usually anti-kt,  $R \approx 1$



large  $p_t$ , large mass fat-jet,  
signal or background

Tagging step



signal jet candidate, still  
background-contaminated

grooming step



final candidate, potentially with little  
background contamination

Note that in some taggers  
(i.e. pruning) the tagging  
and grooming steps are  
not explicitly factorised

Also, some tools may  
actually not follow rigidly  
this scheme

# The BDRS tagger

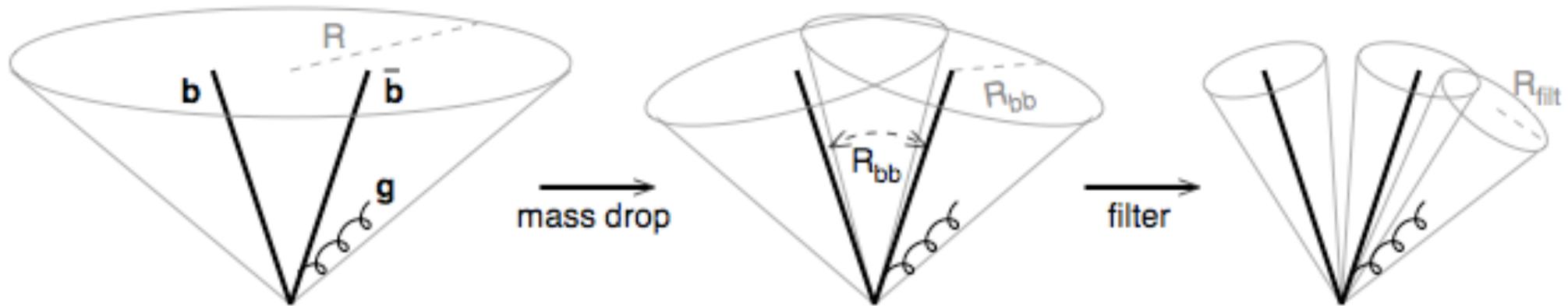
These ideas led to the first ‘modern’ implementation of a boosted tagger

## 15. Jet substructure as a new Higgs search channel at the LHC.

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e-Print: [arXiv:0802.2470](https://arxiv.org/abs/0802.2470) [hep-ph]



It's a two-prongs tagger for boosted Higgs, which

- ▶ Uses the **Cambridge/Aachen** algorithm (see why in the next slide)
- ▶ Employs a **Mass-Drop** condition (as well as an **asymmetry cut**) to find the relevant splitting (i.e. ‘tag’ the heavy particle)
- ▶ Includes a post-processing step, using ‘**filtering**’ (introduced in the same paper) to clean as much as possible the resulting jets of UE contamination

# Why C/A and not $k_t$

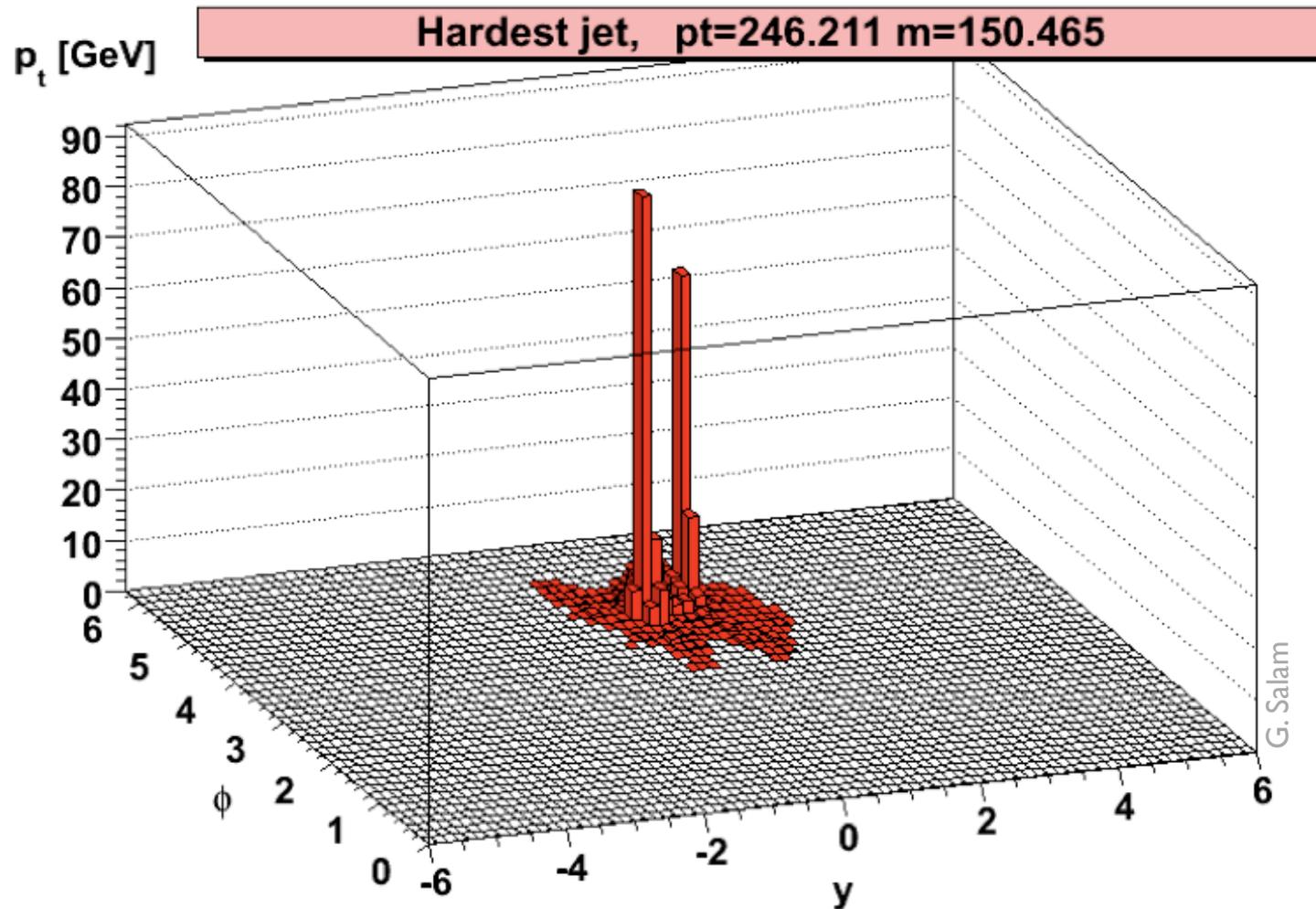
While  $k_t$  gives a ‘natural relevant splitting’ at its first declustering, there are a number of reasons why Cambridge/Aachen has been preferred

- ▶  $k_t$ 's ‘relevant subjects’ tend to include more soft radiation than needed, eventually leading to poor resolution (large areas and fluctuations)
- ▶ The angle-based clustering distance of Cambridge/Aachen ensures that at the relevant splitting the radii of the jets of the two prongs are similar to the distance between the two prongs themselves. This ensures that, because of angular ordering, these jets contain essentially all the radiations emitted by the decay products of the heavy particles (b quarks, in the case of BDRS)
- ▶ Cambridge/Aachen allows one to obtain naturally clustering sequences for any  $R$  with a single run, which is useful in the filtering step

# Boosted Higgs tagger

Butterworth, Davison, Rubin, Salam, 2008

$$pp \rightarrow ZH \rightarrow \nu\bar{\nu}b\bar{b}$$



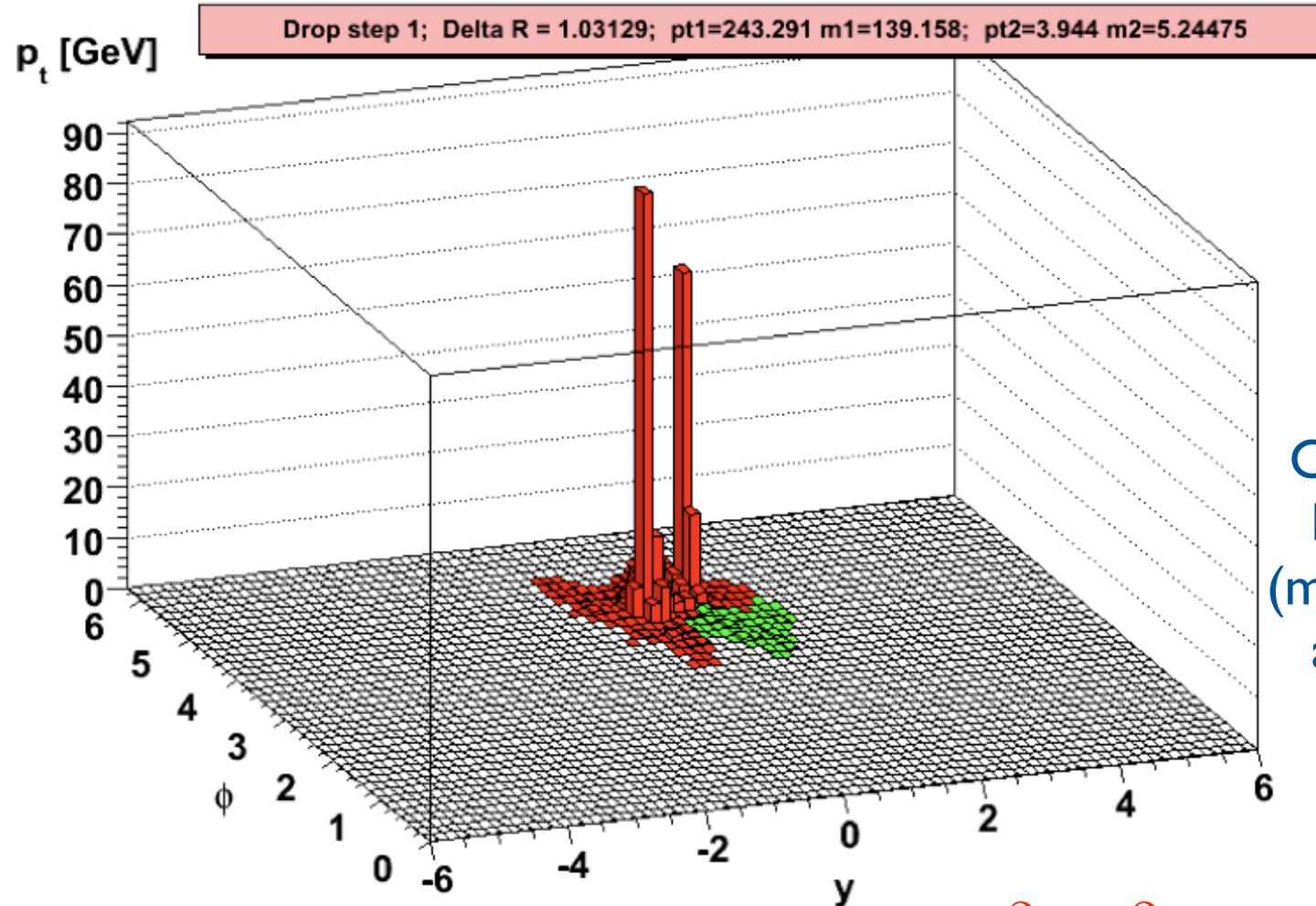
Start with the  
hardest jet

Use C/A with  
large  $R=1.2$

$m_j = 150$  GeV

# Boosted Higgs tagger

$pp \rightarrow ZH \rightarrow \nu\nu b\bar{b}$



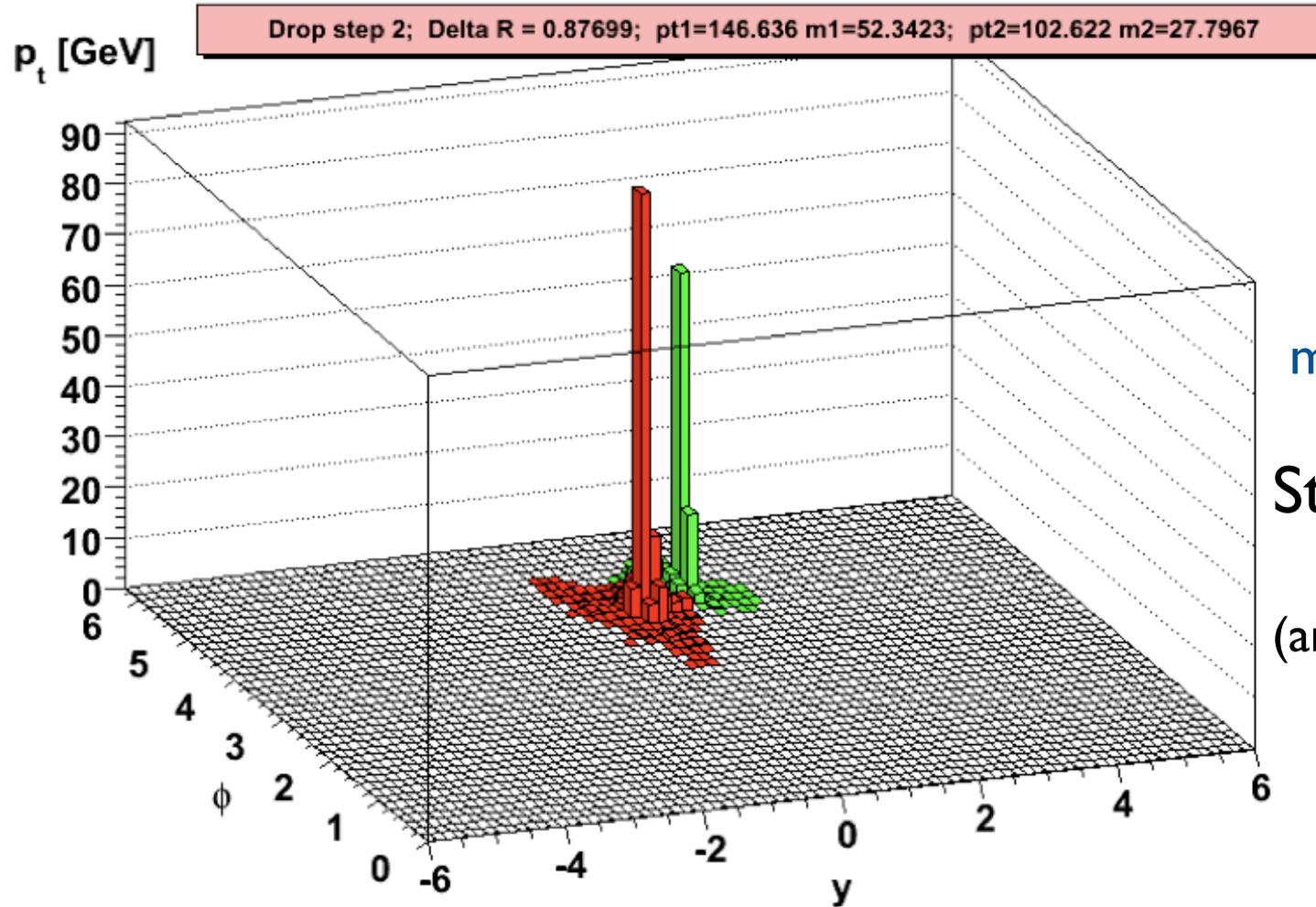
Undo last step of clustering

Check how the mass splits between the two subjects ( $m_1 = 139$  GeV,  $m_2 = 5$  GeV) and how asymmetric the splitting is

If  $\frac{\max(m_1, m_2)}{m_j} > \mu$  or  $\frac{\min(p_{t1}^2, p_{t2}^2)}{m_j^2} \Delta R_{12}^2 < y_{cut}$  repeat

# Boosted Higgs tagger

$pp \rightarrow ZH \rightarrow \nu\nu b\bar{b}$



$m_1 = 52 \text{ GeV}, m_2 = 28 \text{ GeV}$

Stop when a **large mass drop** is observed  
(and recombine these two jets)

[NB. Parameters used  $\mu = 0.67$  and  $y_{\text{cut}} = 0.09$ ]

# Jet substructure as filter

The **jet substructure** can be exploited to help **removing contamination** from a soft background

▶ Jet ‘filtering’

Butterworth, Davison, Rubin, Salam, 2008

Break jet into subjets at distance scale  $R_{\text{filt}}$ , retain  $n_{\text{filt}}$  hardest subjets

▶ Jet ‘trimming’

Krohn, Thaler, Wang, 2009

Break jet into subjets at distance scale  $R_{\text{trim}}$ , retain subjets with  $p_{t,\text{subjet}} > \epsilon_{\text{trim}} p_{t,\text{jet}}$

▶ Jet ‘pruning’

S. Ellis, Vermilion, Walsh, 2009

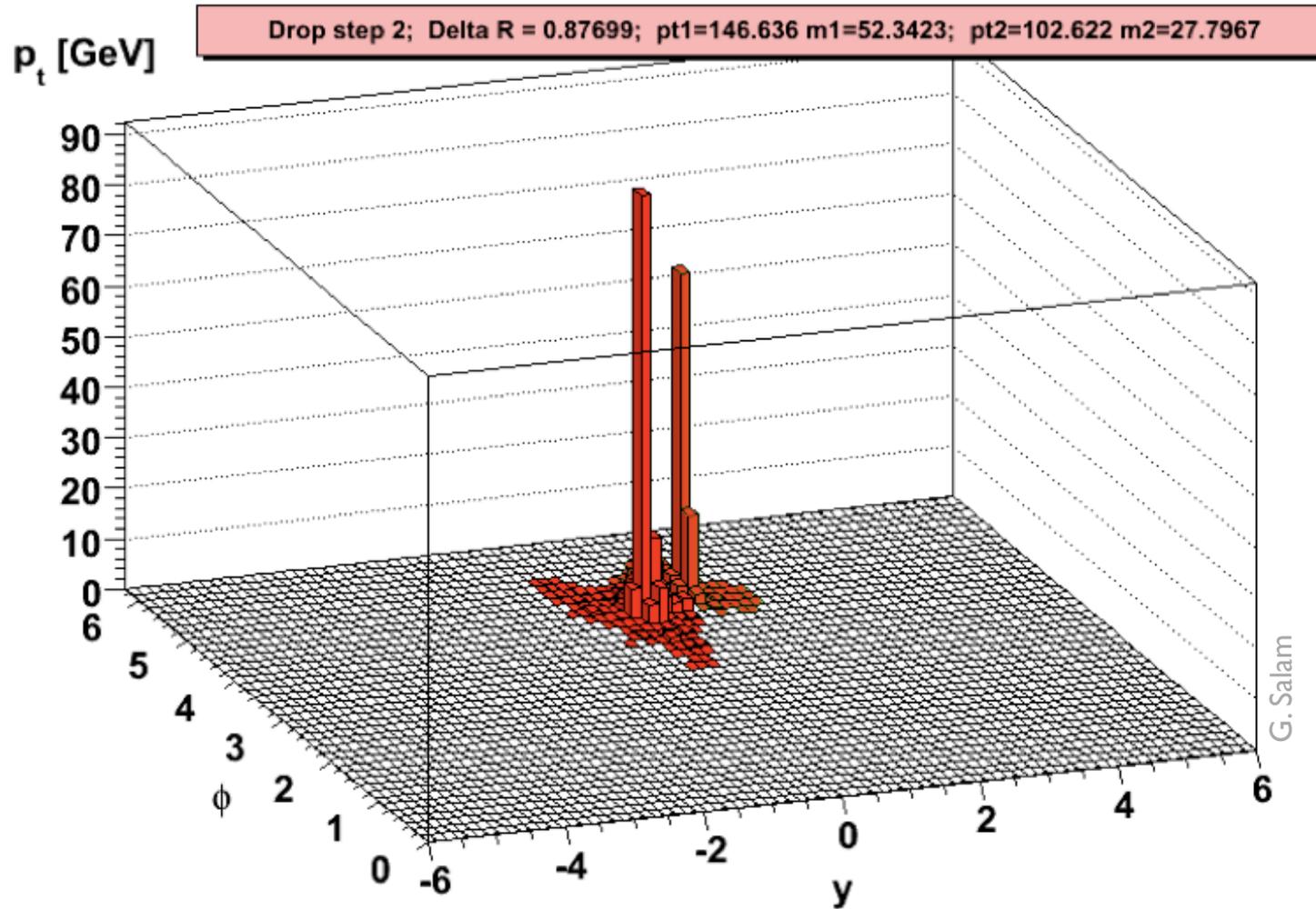
While building up the jet, discard softer subjets when  $\Delta R > R_{\text{prune}}$  and  $\min(p_{t1}, p_{t2}) < \epsilon_{\text{prune}} (p_{t1} + p_{t2})$

**Aim: limit sensitivity to background while retaining bulk of perturbative radiation**

Filtering, trimming and pruning are in the end effectively quite similar.  
These and similar tools are collectively called **groomers**

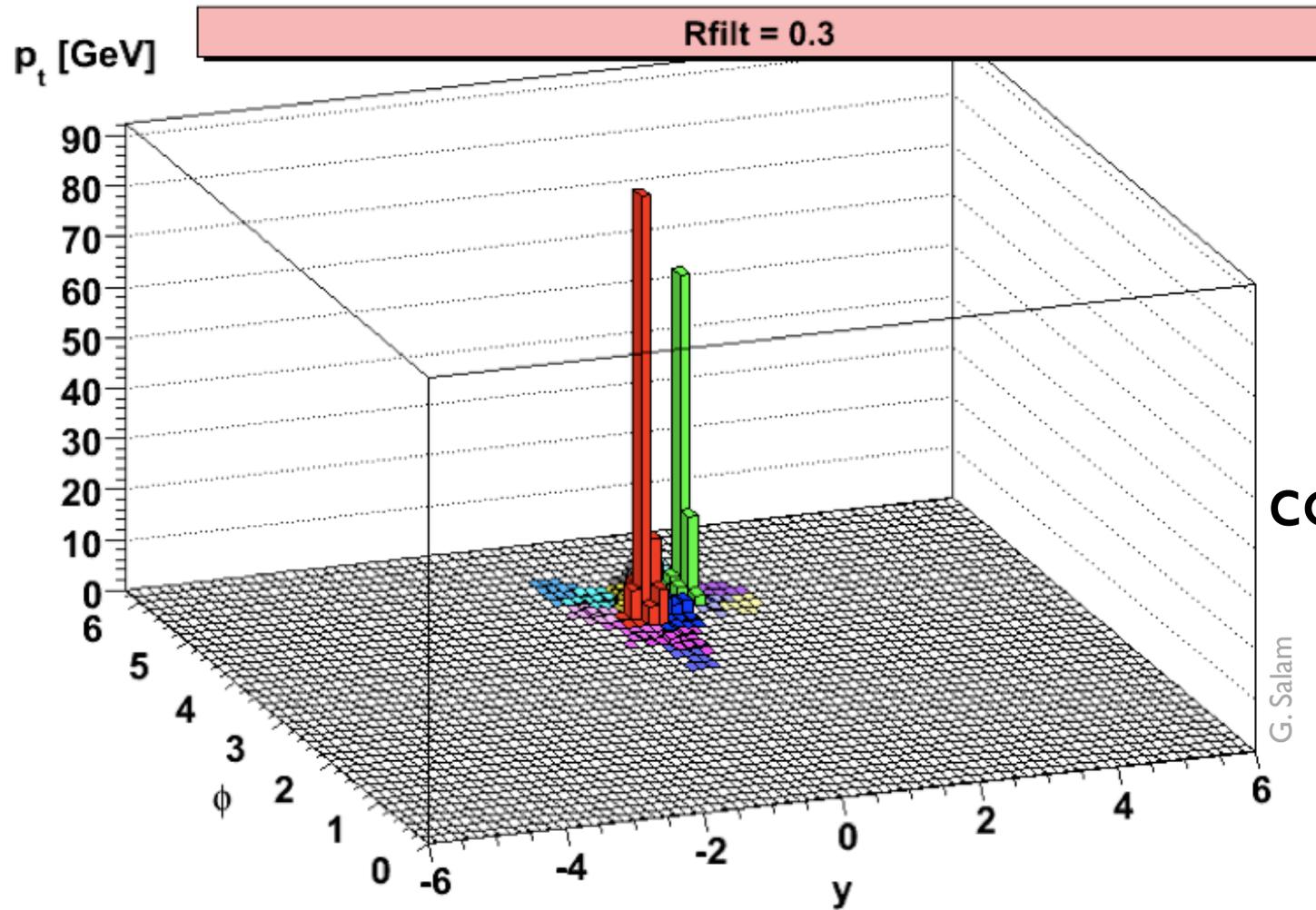
# Filtering in action

Butterworth, Davison, Rubin, Salam, arXiv:0802.2470



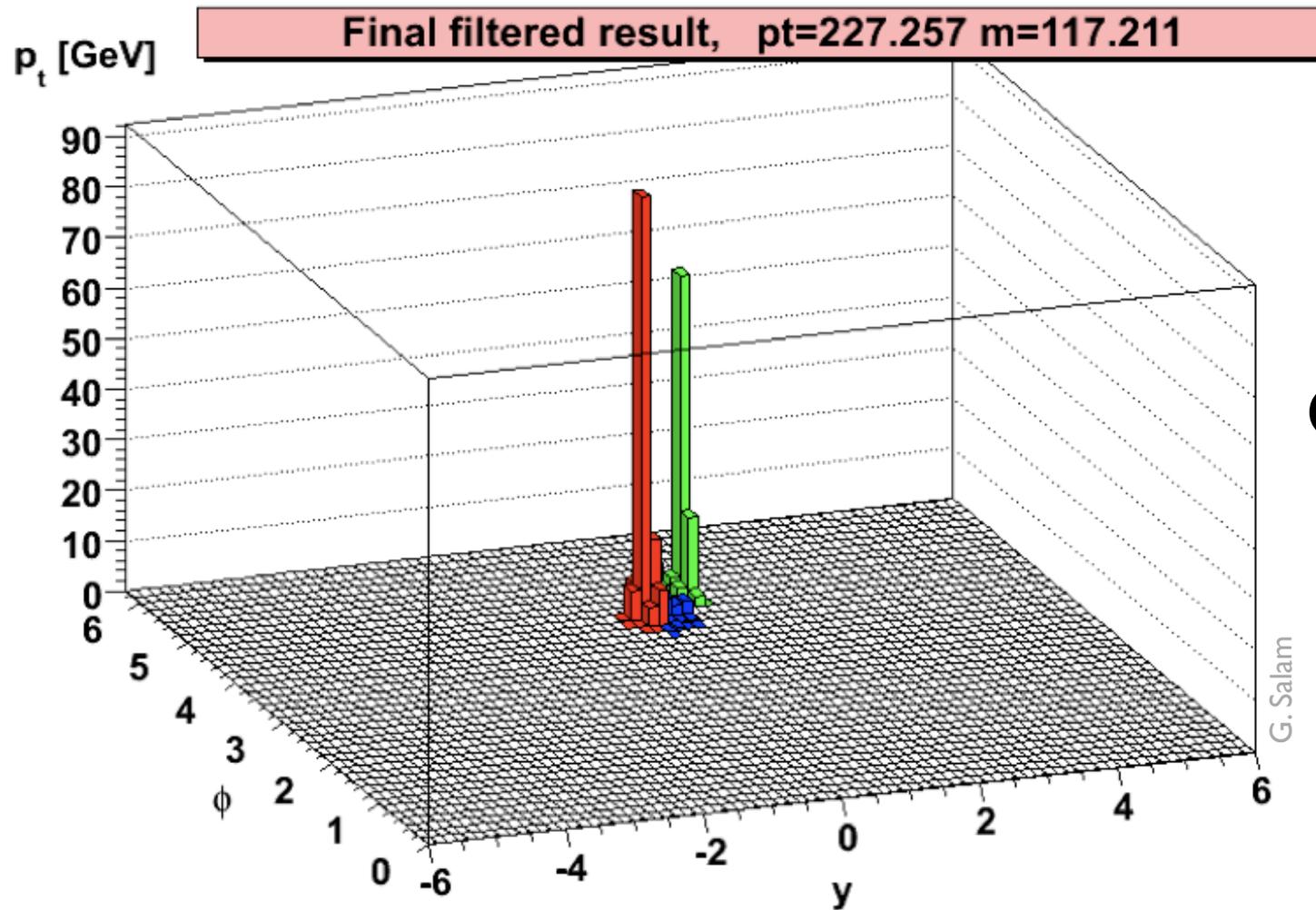
Start with a jet

# Filtering in action



Recluster the constituents with  $R_{\text{filt}}$

# Filtering in action



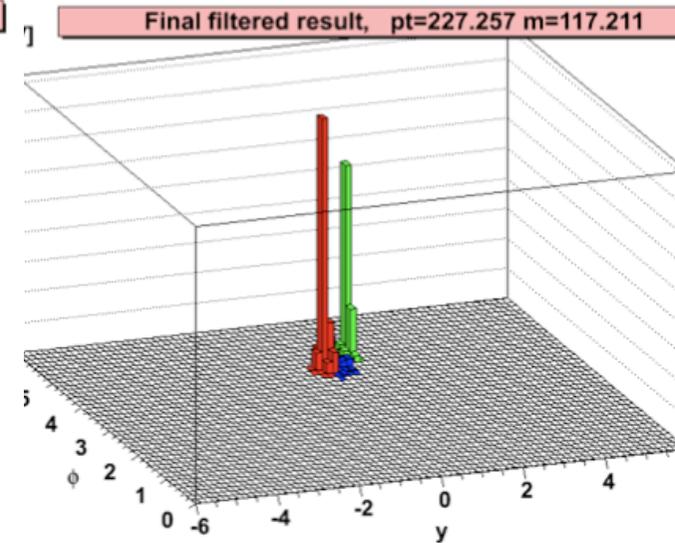
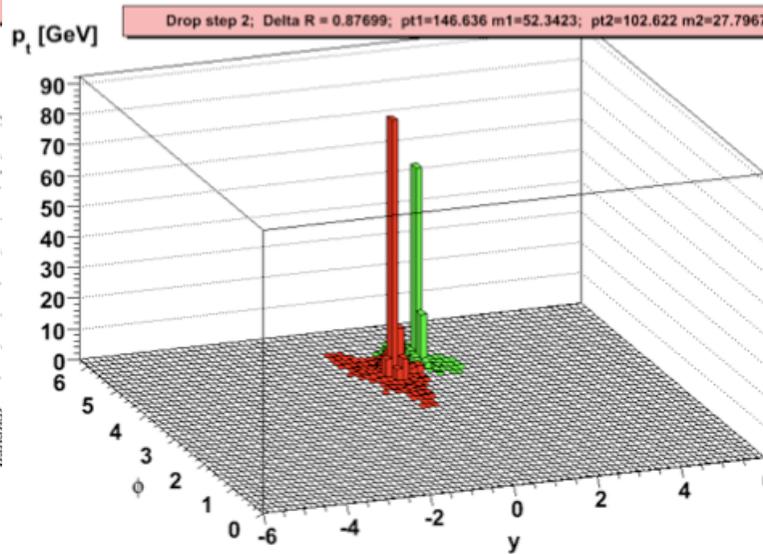
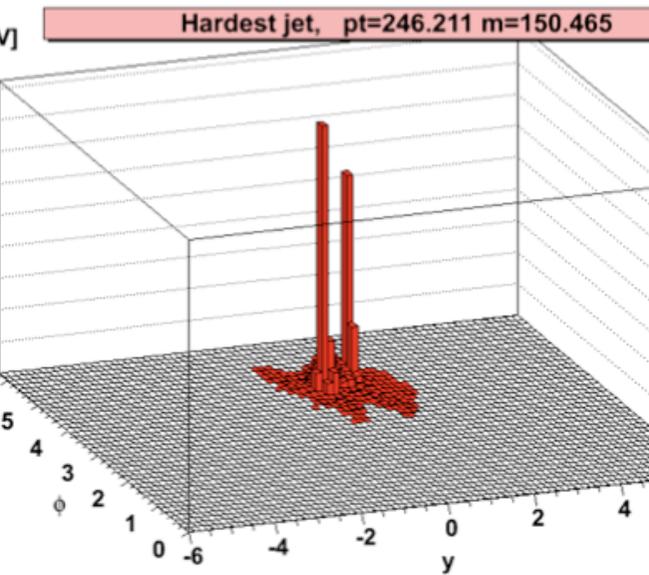
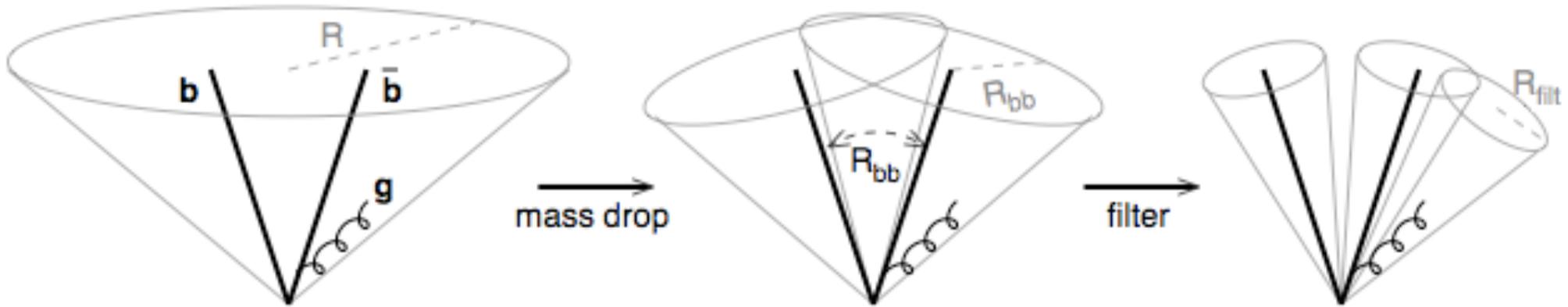
Only keep the  $n_{\text{filt}}$   
hardest jets

The low-momentum stuff surrounding the hard particles has been removed

# Boosted Higgs analysis

Butterworth, Davison, Rubin, Salam, 2008

$$pp \rightarrow ZH \rightarrow \nu\bar{\nu}b\bar{b}$$



Cluster with a large R

Undo the clustering into subjects,  
until a large mass drop  
is observed

Re-cluster with smaller R,  
and keep only 3 hardest  
jets

# BDRS in FastJet

```
#include "fastjet/tools/MassDropTagger.hh"
#include "fastjet/tools/Filter.hh"

JetDefinition jet_def(cambridge_algorithm, 1.2);
ClusterSequence cs(input_particles, jet_def);

// define the tagger and use it
MassDropTagger md_tagger(0.667, 0.09);
PseudoJet tagged = md_tagger(jets[0]);

// define the filter and use it
Filter filter(0.3, SelectorNHardest(3));
Pseudojet higgs = filter(tagged);           // this is the Higgs!!
```

The real analysis is slightly more refined (b-tagging, dynamical filter radius, etc)  
but the main features are already present here

# Pruning in FastJet

```
#include "fastjet/tools/Pruner.hh"

JetDefinition jet_def(cambridge_algorithm, 1.2);
ClusterSequence cs(input_particles, jet_def);

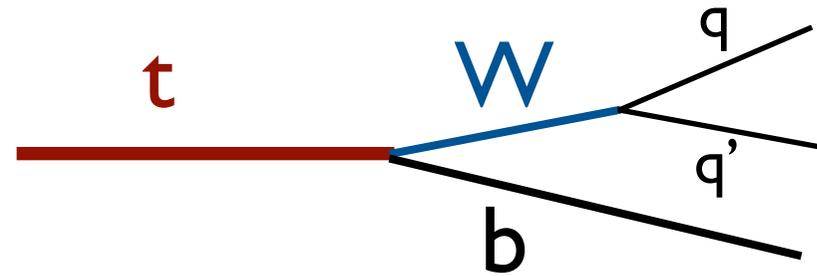
// define the pruner and use it
double zcut = 0.1;
double rcut_factor = 0.5;

Pruner pruner(cambridge_algorithm, zcut, rcut_factor);

PseudoJet tagged = pruner(jets[0]);
```

# Top tagging

In order to tag a (boosted) top one must now identify 3-prongs structures originating from the top decay



Simplest approach: iterate declustering of  $k_t$  jets, beyond a first ‘relevant splitting’

Early examples (2008):

- ▶ *ATLAS top tagger: put cuts on jet mass and  $d_{ij}$  scale*
- ▶ *Thaler-Wang: decluster to exactly 2 or 3 jets, put cuts on jet mass*

Many more top taggers after these

# Example: top taggers

[1 jet  $\gtrsim$  2 partons]

↳ [An example]

## Tagging boosted top-quarks

Many papers on top tagging in '08-'11: jet mass + something extra.

### Questions

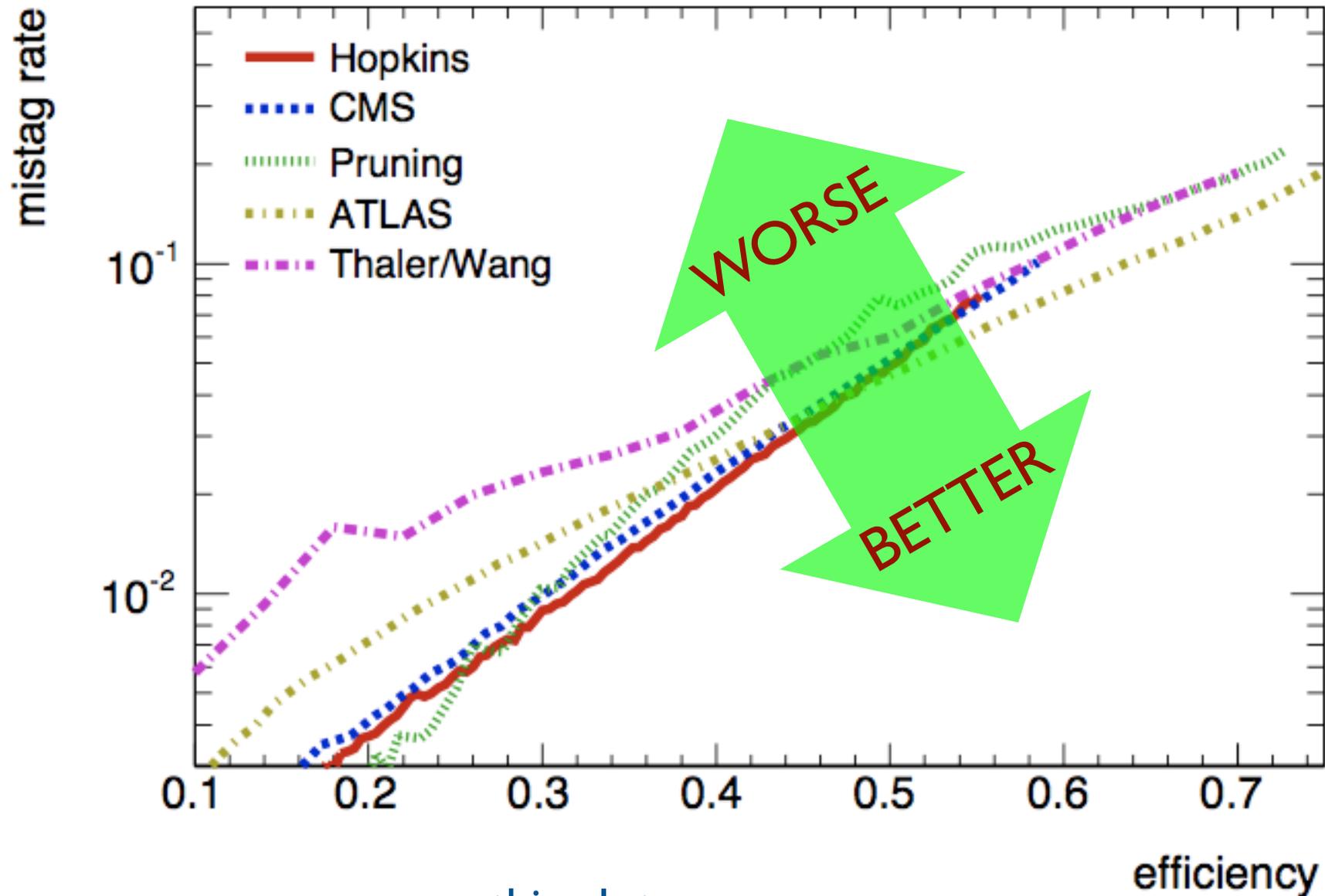
- ▶ What efficiency for tagging top?
- ▶ What rate of fake tags for normal jets?

### Rough results for top quark with $p_t \sim 1$ TeV

	"Extra"	eff.	fake
[from T&W]	just jet mass	50%	10%
Brooijmans '08	3,4 $k_t$ subjets, $d_{cut}$	45%	5%
Thaler & Wang '08	2,3 $k_t$ subjets, $z_{cut}$ + various	40%	5%
Kaplan et al. '08	3,4 C/A subjets, $z_{cut}$ + $\theta_h$	40%	1%
Ellis et al. '09	C/A pruning	10%	0.05%
ATLAS '09	3,4 $k_t$ subjets, $d_{cut}$ MC likelihood	90%	15%
Chekanov & P. '10	Jet shapes	60%	10%
Almeida et al. '08-'10	Template + shapes	13%	0.02%
Thaler & v Tilburg '10	Subjettiness	40%	2%
Plehn et al. '09-'10	C/A MD, $\theta_h$ /Dalitz [busy evs, $p_t \sim 300$ ]	35%	2%

# Comparison of top taggers

Boost 2010 proceedings, arXiv:1012.5412

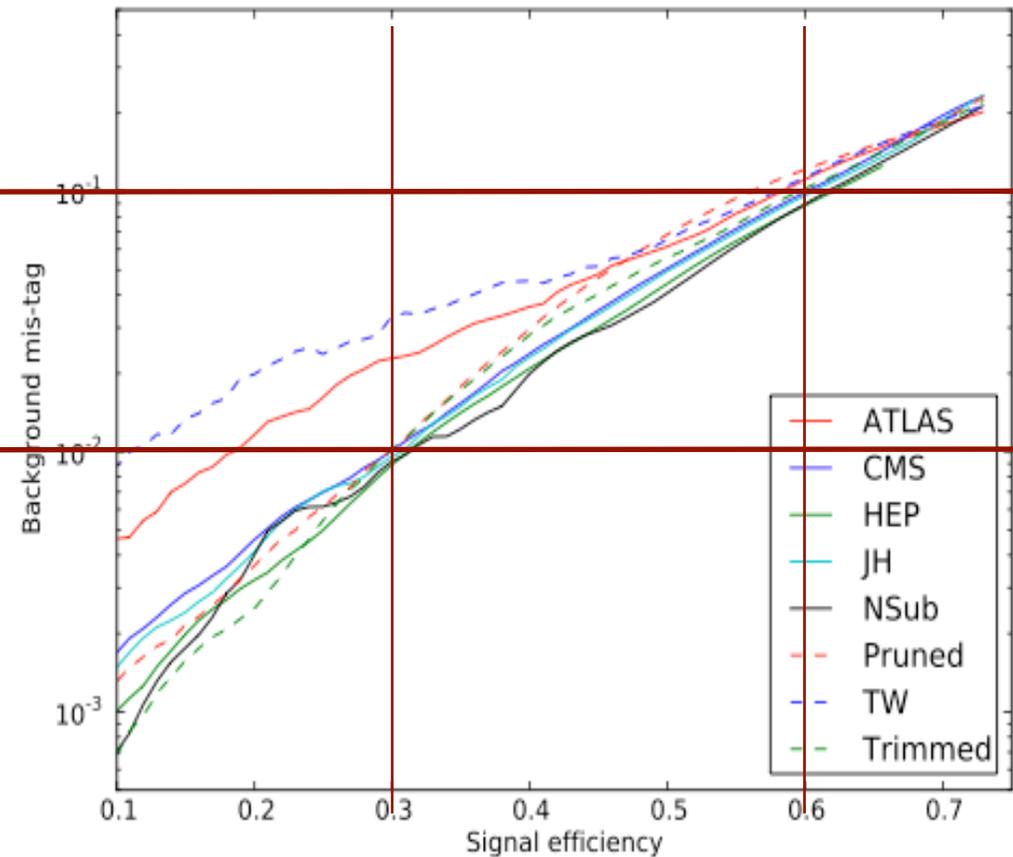
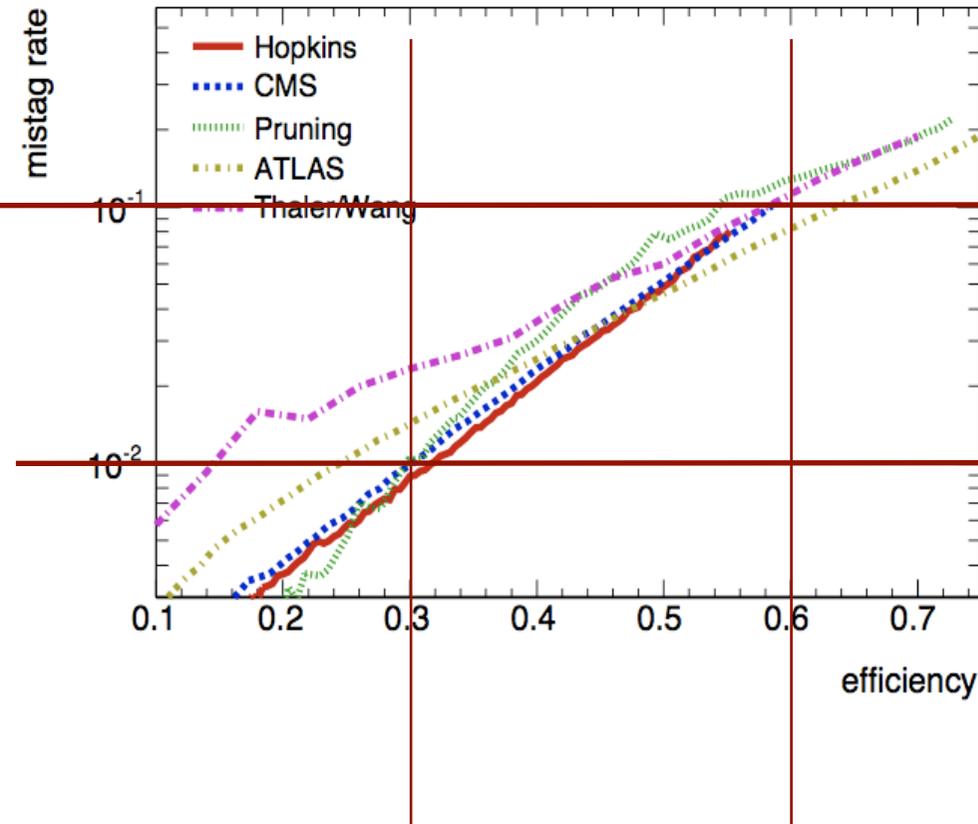


Even more curves now on this plot

# Comparison of top taggers

Boost **2010** proceedings, arXiv:1012.5412

Boost **2011** proceedings, arXiv:1201.0008



Law of diminishing returns: improvement has become very hard

Boost 2012 is under way right now in Valencia.

# Concluding remarks

- ▶ Proper (IRC-safe) definition of jet algorithms and efficient implementations have allowed for the study and exploitation of a number of jet properties
  - ▶ jet areas
    - ▶ *background determination and subtraction*
  - ▶ jet substructure
    - ▶ *taggers*
    - ▶ *groomers*
- ▶ Many new physics search strategies based on jet substructure are being explored and commissioned right now at the LHC
  - ▶ *As soon as more data (and more boosted particles) are available, we should see the first results from them*
- ▶ While these tools seem mature and are being refined, one can hope that new radical ideas will again start a 'revolution' and deeply change the field



# Areas as a dynamical jet property

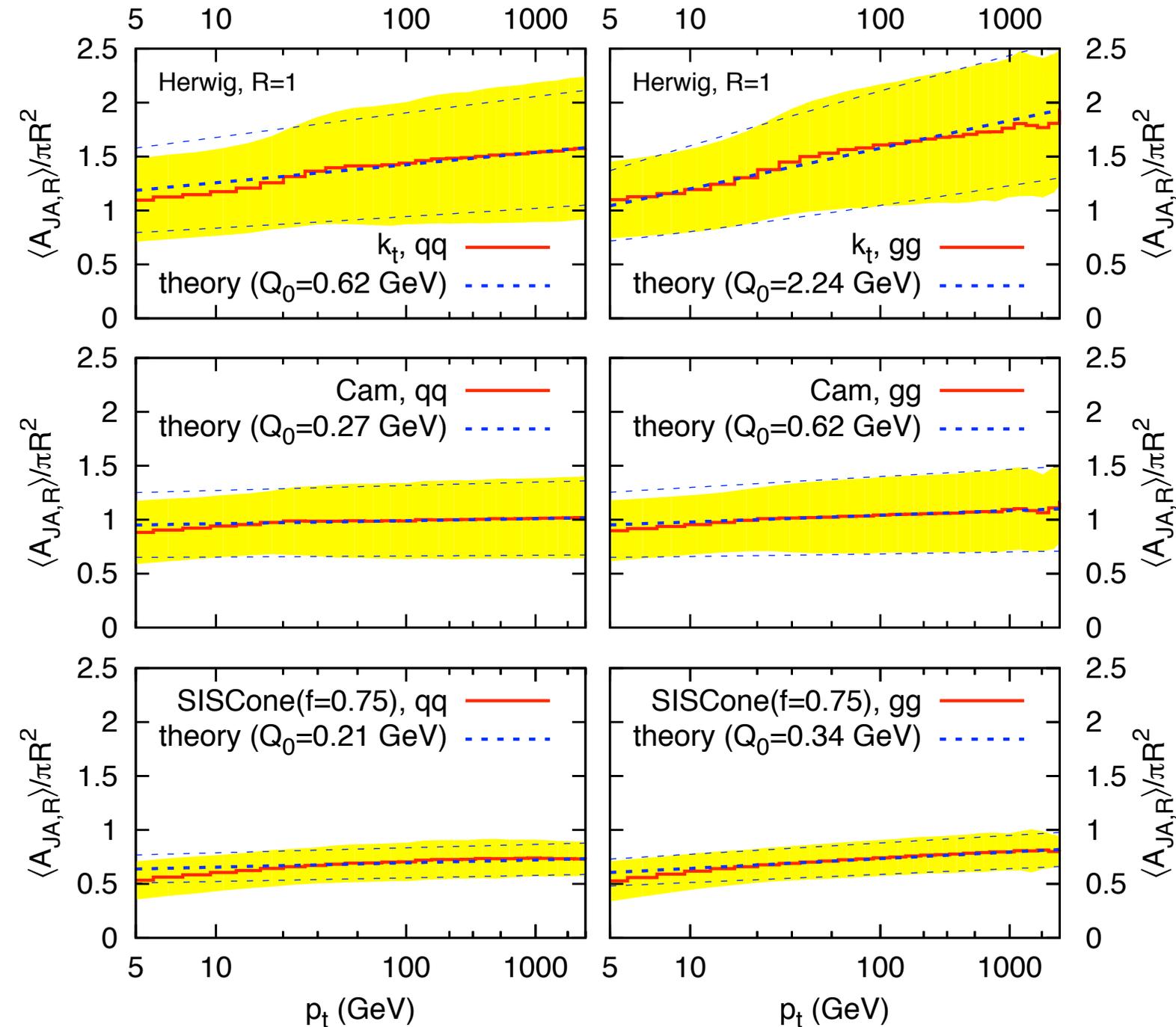
The average area of a jet can change with its  $p_t$ :

$$\langle \Delta A \rangle = \mathbf{D} \frac{C_1}{\pi b_0} \ln \frac{\alpha_s(Q_0)}{\alpha_s(Rp_{t1})}$$

	$k_t$	Cam/Aa	SISCone	anti- $k_t$
D	0.52	0.08	0.12	0

Again, only **anti- $k_t$**  has a typical area that does **not** increase with  $p_t$

# Jet areas scaling violations



Averages and dispersions evolution from Monte Carlo simulations (dijet events at LHC) in good agreement with simple LL calculations

**Area scaling violations are a legitimate observable.**

(Though they might not be the best place where to measure  $\alpha_s$  ...)

# Jet areas scaling violations

MC, Salam, Soyez, arXiv:0802.1189

Check anti- $k_t$  behaviour: **scaling violations** indeed **absent**, as predicted

