



FIAS Frankfurt Institute
for Advanced Studies



SMFNS2013

“Neutron stars (and more) within
the pseudo-complex General Relativity”

Neutron stars: Isaac Rodriguez, Peter Otto Hess,
Walter Greiner, S. Schramm

pc-GR and experiment: T. Schönenbach, G. Caspar, M. Schäfer
T. Boller, A. Müller,

CONTENT

- Introduction: Motivation, what is pseudo-complex (pc) and justification.
- pc-GR and pc-Einstein equations.
- Possible origin of the dark energy.
- Some experimental predictions of pc-GR
- Ansatz for the energy-momentum tensor.
- TOV equations.
- Results
- Conclusions

Motivation

FIRST ATTEMPTS.

○ A. Einstein, *Ann.Math.* **46** (1945), 518.

○ A. Einstein, *Rev. Mod. Phys.* **20** (1948), 35.

(Unification of gravitation and electrodynamics)

○ C. Mantz, T. Prokopec, (2008); arXiv:0804.0213

more recently: C. Mantz, T. Prokopec, *Found. Phys.* **41** (2011), 1597

(hermitian gravity and cosmology)

$$X^\mu = x^\mu + i \frac{l}{m} p^\mu, \quad i^2 = -1 \quad (\text{Introduction of the Planck length, } l)$$

Born's equivalence principle: $[x^k, p^j] = i\hbar \delta_k^j, [x^k, x^j] = 0, [p^k, p^j] = 0$

but $d^2 \Rightarrow g_\mu d_\nu^\mu dx^\nu$

(M. Born, *Proc. Roy. Soc. A* **165** (1938), 291 and
M. Born, *Rev. Mod. Phys.* **21** (1949), 463.)

PROPOSAL

(M. BORN)

$$d\Omega^2 = d_{\mu} d x^{\mu} + l^2 d_{\mu} d u^{\mu} = d_{\mu} d x^{\mu} \left(1 + l^2 \frac{d_{\mu} d u^{\mu}}{d\tau} \right)$$

$$\rightarrow d_{\mu} d x^{\mu} (1 + l^2 a^2)$$

$$a^2 = -a_{\mu} a^{\mu} = \frac{d_{\mu} d u^{\mu}}{d\tau} \quad \square \quad l^2 \leq \frac{1}{a^2}$$

l is a minimal length

a = maximal acceleration!!!

E.R. Caianiello (1981), H.E. Brandt, R.G. Beil (1980's)
S.G. Low (1990's and more recently: representation theory)

Pseudo-complex

$$x^\mu \rightarrow X^\mu = x^\mu + Iy^\mu \quad , \quad I^2 = 1$$
$$= X_+^\mu \sigma_+ + X_-^\mu \sigma_- \quad \sigma_\pm = \frac{1}{2}(1 \pm I)$$
$$\sigma_\pm^2 = \sigma_\pm \quad , \quad \sigma_+ \sigma_- = 0 \quad \rightarrow \text{Zero divisor}$$

PF Kelly and RB Mann, Class. Quant. Grav. 3 (1986), 705:

From all possible algebraic extensions only the pseudo-complex does not have ghost solutions in the limit of weak gravitational fields.

Final projection to the real 4-dimensional space: see talk of M. Schäfer at STARS2013.

Alternatively: modify variational principle such that the variation of the action is within the zero divisor.

The Theory

I. NEW VARIATIONAL PRINCIPAL:

$$S = \int L d\tau$$

$$\delta S = \delta S_+ \sigma_+ + \delta S_- \sigma_- \in z \quad e \quad d \quad r \quad i$$

$$\rightarrow \quad = \xi \quad - \quad \alpha c \quad o \quad) \quad n \quad v$$

If we use $\omega=0$, then: $\delta S_{\pm} = 0$ independently, thus two independent theories, i.e., no connection!

This results in the equations of motion

$$\frac{D}{D} \left(\frac{D}{D \dot{X}^{\mu}} \right) - \left(\frac{D}{D} \right)^{\mu} \in \frac{L}{X} z \quad e \quad d \quad r \quad i$$

(F. Schuller, PhD thesis, University of Cambridge (2003);

F. Schuller, *Ann. Phys. (N.Y.)* **299** (2002), 174,

F. S. has proposed this general variation principle in his thesis at Cambridge.)

Extension of the Theory of General Relativity:

- The metric is pseudo-complex, without torsion:

$$g_{\mu\nu} = g_{\mu\nu}^+ \sigma_+ + g_{\mu\nu}^- \sigma_- \quad , \quad g_{\mu\nu} = g_{\nu\mu}$$

* pseudo-complex length element

$$\begin{aligned} d\omega^2 &= g_{\mu\nu} DX^\mu DX^\nu \\ &= g_{\mu\nu}^+ DX_+^\mu DX_+^\nu \sigma_+ + g_{\mu\nu}^- DX_-^\mu DX_-^\nu \sigma_- \end{aligned}$$

$$\rightarrow g_{\mu\nu} \left(dx^\mu dx^\nu + \left[\frac{l}{c} \right]^2 du^\mu du^\nu \right) \quad , \quad \left(\frac{l}{c} u^\mu \ll 1 \right)$$

and $g_{\mu\nu} dx^\mu du^\nu = 0$

← dispersion relation!

Einstein equation

$$L = \sqrt{-g} R$$

$$G^{\mu\nu} = R^{\mu\nu} - \frac{1}{2} g^{\mu\nu} R = -\frac{8\pi\kappa}{c^2} T^{\mu\nu} \sigma_-$$

Origin of the dark energy?

M Visser, Phys. Rev. D **54** (1996) 5103 +...

C. Barceló et al., Phys. Rev. D **77** (2008), 044032

- 1) Semiclassical QM, with a Schwarzschild metric as background. Vacuum fluctuations are building up (dark energy!), which finally stop the collapse of a star (if the collapse is slow enough)
- 2) Advantage: They determine the density of the dark energy
Disadvantage: They do not have a re-coupling of the dark energy density to the metric and thus have to take care of the event horizon

Our theory:

- 1) We assume a dark energy density and determine the re-coupling of the dark energy density to the metric.
- 2) Advantage: We can determine the final metric
Disadvantage: We can not determine the form of the dark energy density → assumed

Both ways are complementary! Both explain the stop of a collapse due to the *distorsion of space due to the presence of the dark energy distribution such that space itself retains the collapse of a star!*

Schwarzschild and Kerr metrics

(G. Caspar et al., Int. J. Mod. Phys. E 21 (2012), 1250015)

Schwarzschild:

$$\begin{aligned}g_{00}^S &= \left(1 - \frac{2m}{r} + \frac{B}{2r^3}\right) \\g_{11}^S &= -\left(1 - \frac{2m}{r} + \frac{B}{2r^3}\right)^{-1} \\g_{22}^S &= -r^2 \\g_{33}^S &= -r^2 \sin^2 \vartheta.\end{aligned}$$

B is a measure of the dark energy Contribution:

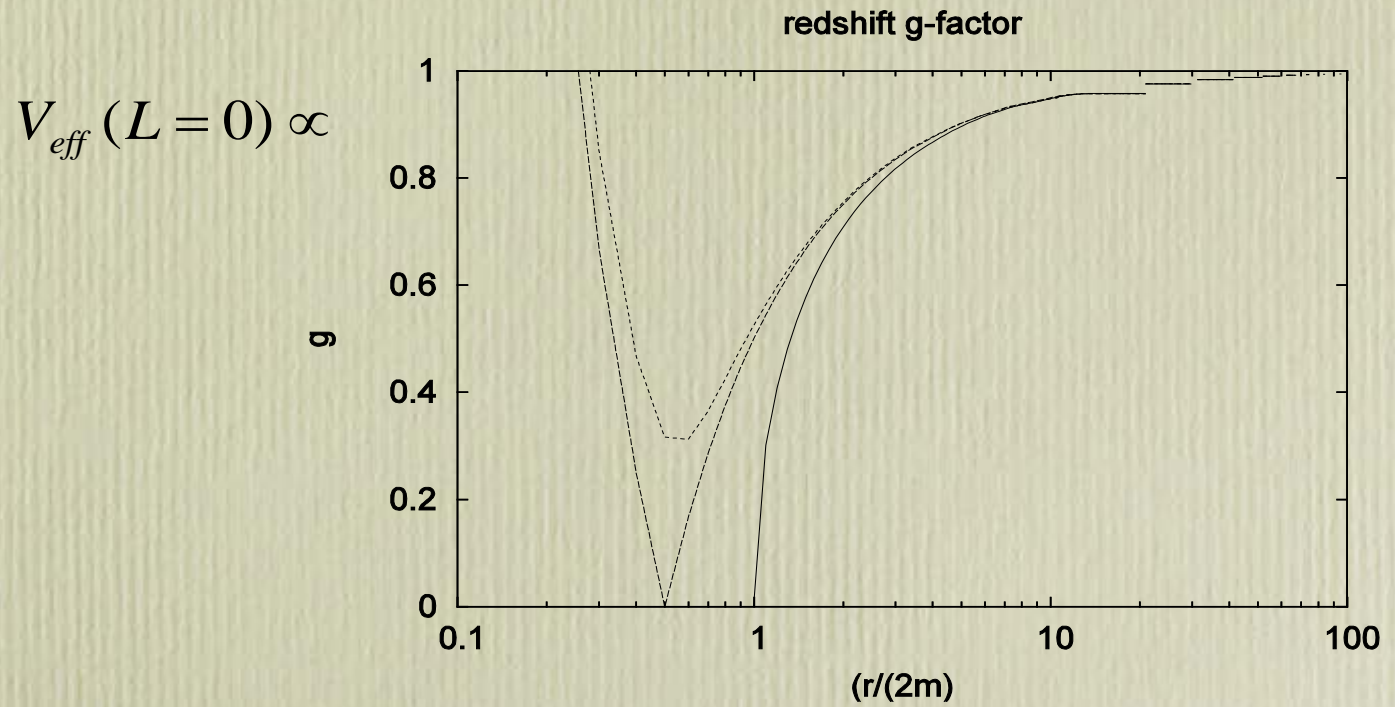
Kerr:

$$\begin{aligned}g_{00}^K &= \frac{r^2 - 2mr + a^2 \cos^2 \vartheta + \frac{B}{2r}}{r^2 + a^2 \cos^2 \vartheta} \\g_{11}^K &= -\frac{r^2 + a^2 \cos^2 \vartheta}{r^2 - 2mr + a^2 + \frac{B}{2r}} \\g_{22}^K &= -r^2 - a^2 \cos^2 \vartheta \\g_{33}^K &= -(r^2 + a^2) \sin^2 \vartheta - \frac{a^2 \sin^4 \vartheta (2mr - \frac{B}{2r})}{r^2 + a^2 \cos^2 \vartheta} \\g_{03}^K &= \frac{-a \sin^2 \vartheta 2mr + a \frac{B}{2r} \sin^2 \vartheta}{r^2 + a^2 \cos^2 \vartheta},\end{aligned}$$

$$B = bm^3$$

Schwarzschild:

$$g = \sqrt{g_{00}} \propto V_{eff}(L=0)$$



The radius of the dark disk is **SMALLER** in pc-GR!

Circular orbital of a particle around a large mass

(T. Schönembach et al., MNRAS 430 (2013), 2999)

Lagrangian: $L = g_{00}c^2\dot{t}^2 + g_{11}\dot{r}^2 + g_{22}\dot{\vartheta}^2 + g_{33}\dot{\varphi}^2 + 2g_{03}c\dot{t}\dot{\varphi} = \frac{ds^2}{ds^2} = 1$

→ geodesic

$$\frac{d}{ds} (2g_{11}\dot{r}) = g'_{00}c^2\dot{t}^2 + g'_{11}\dot{r}^2 + g'_{22}\dot{\vartheta}^2 + g'_{33}\dot{\varphi}^2 + 2g'_{03}c\dot{t}\dot{\varphi}$$

Circular motion: $\dot{r} = 0$, $\vartheta = \frac{\pi}{2}$ → $0 = g'_{00}(r_0)c^2\dot{t}^2 + g'_{33}(r_0)\omega^2\dot{t}^2 + 2g'_{03}(r_0)\omega c\dot{t}^2$

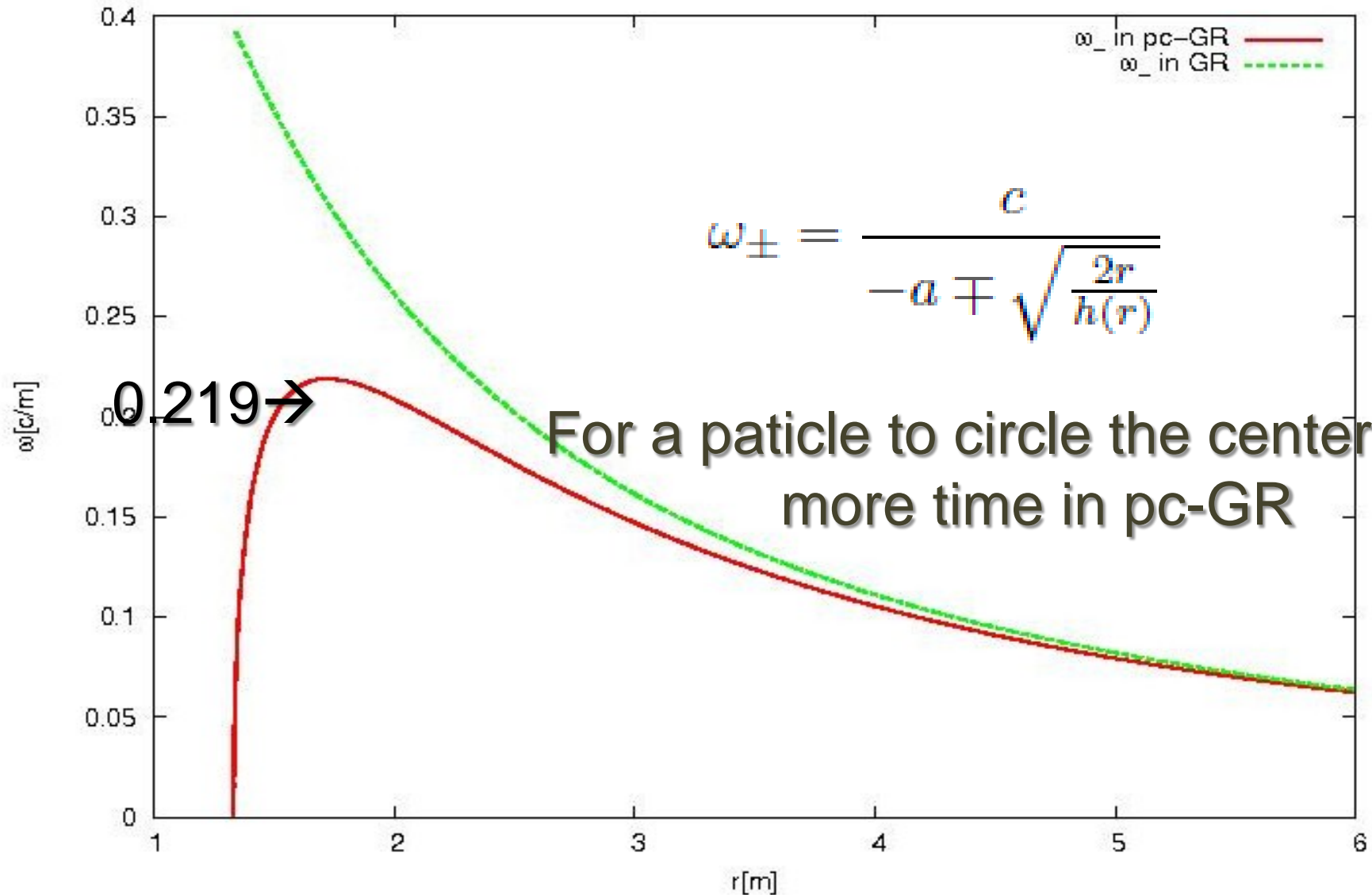
Resolve for the frequency:

$$\omega_{\pm} = c \frac{-g'_{03} \pm \sqrt{(g'_{03})^2 - g'_{00}g'_{33}}}{g'_{33}}$$

OR:

$$\omega_{\pm} = \frac{c}{-a \mp \sqrt{\frac{2r}{h(r)}}}$$

$a = -0.995m$



For a particle to circle the center takes more time in pc-GR

0.219 for 4 Million sun masses = 9.4 minutes

Redshift

$$d\tau^2 = g_{00}dt^2$$

,

$$\tau_0 = \sqrt{g_{00}}t_{obs}$$



$$\nu_{obs} = \sqrt{g_{00}}\nu_0$$

Redshift:

$$z := \frac{\nu_0 - \nu_{obs}}{\nu_{obs}} = \frac{1}{\sqrt{g_{00}}} - 1$$

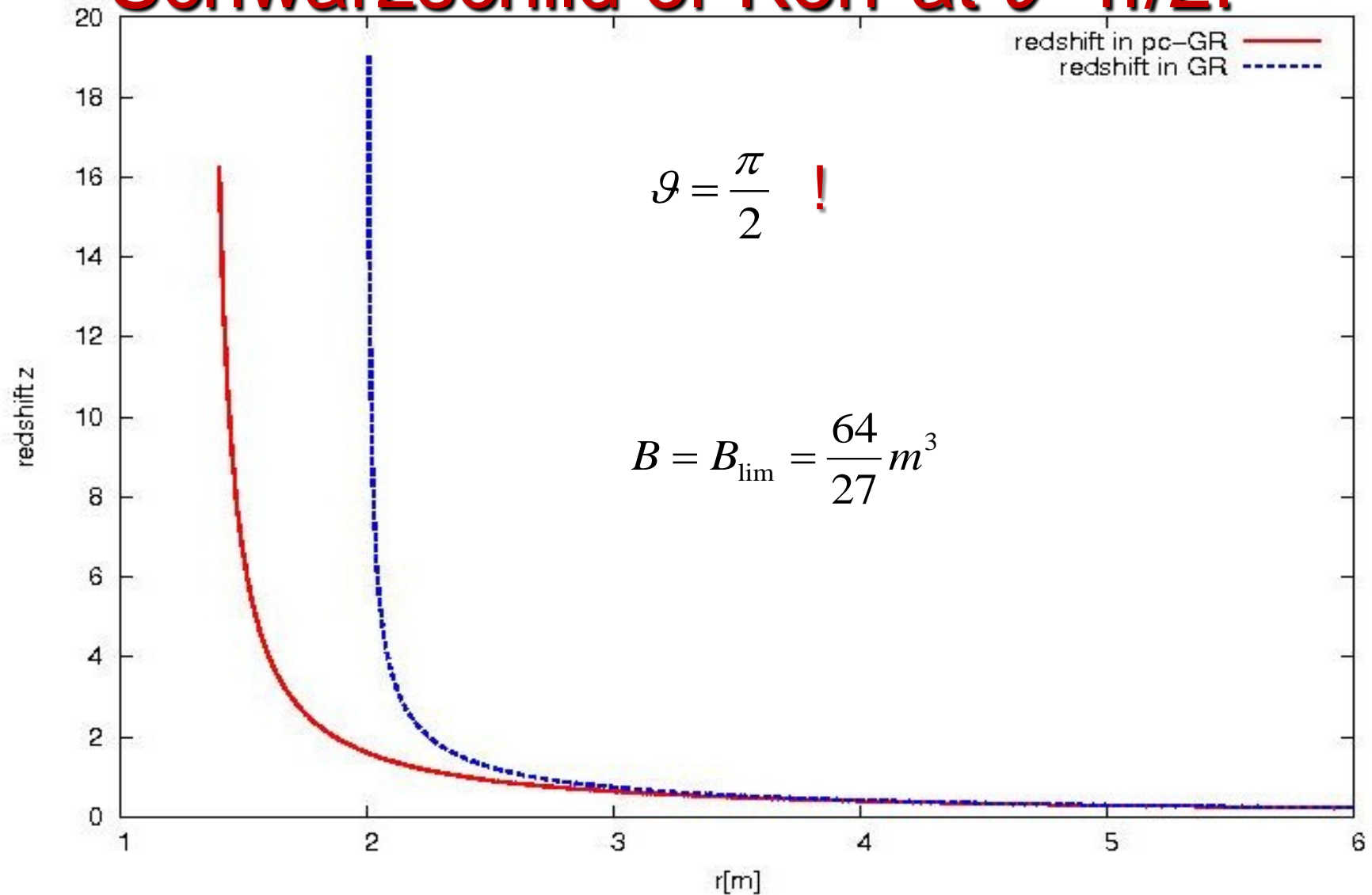
Kerr:

$$z = \frac{\sqrt{r^2 + a^2 \cos^2(\vartheta)}}{\sqrt{r^2 - 2mr + a^2 \cos^2(\vartheta) + \frac{B}{2r}}} - 1$$

Schwarzschild:

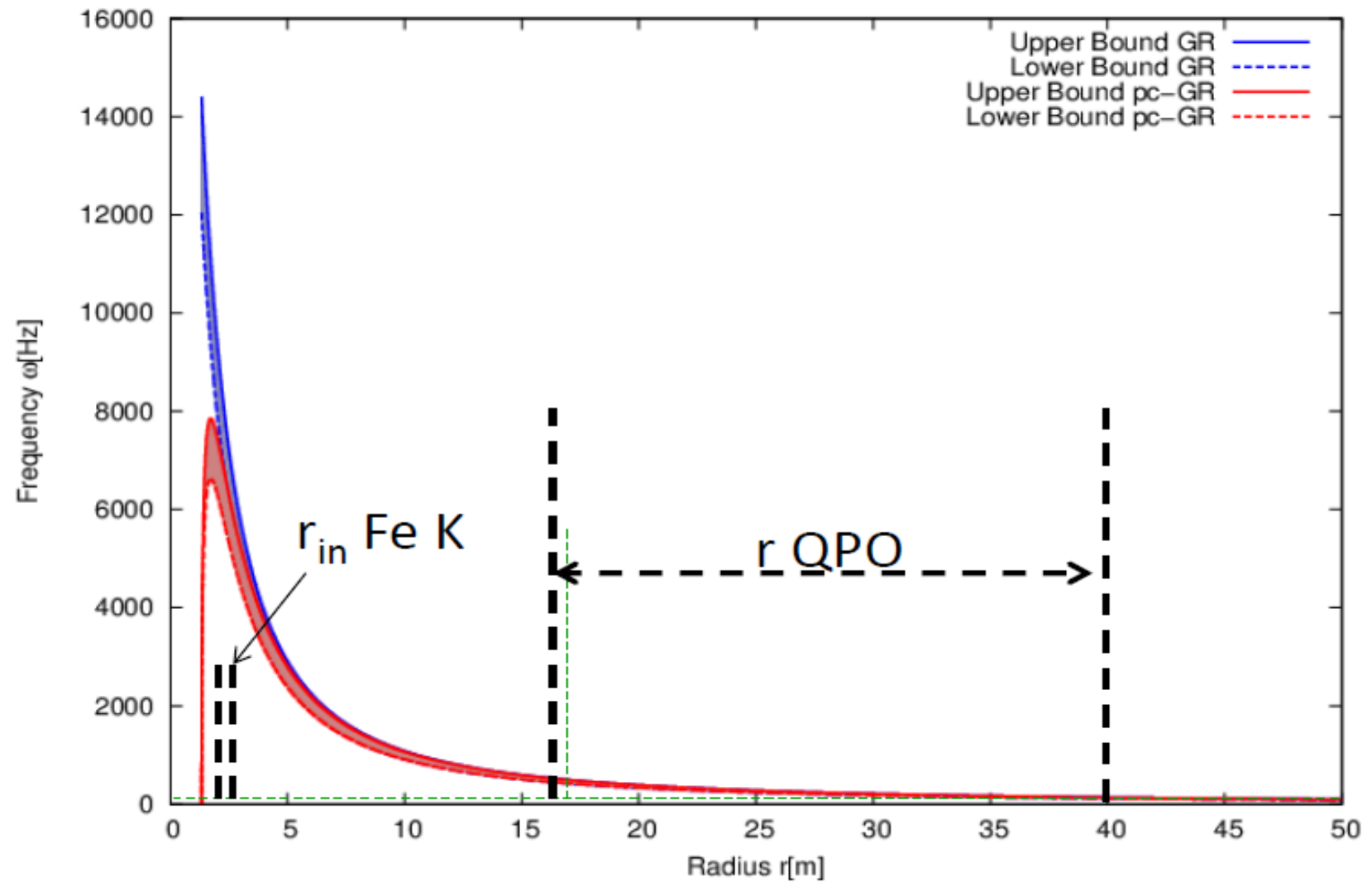
$$z = \frac{1}{\sqrt{1 - \frac{2m}{r} + \frac{B}{2r^3}}} - 1$$

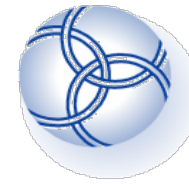
Schwarzschild or Kerr at $\vartheta = \pi/2$:



GRO J1655-40

$r_{in}(2.0-2.4)$ $a=(0.89..0.94)$, $M=6.3\pm 0.5$ 68 per cent confidence





Pseudo-Complex General Relativity (pc-GR):

$$\left(\mathcal{R}_{\nu i}^{\mu+} - \frac{1}{2}g_{\nu i}^{\mu+}\mathcal{R}_i^+\right)\sigma_+ + \left(\mathcal{R}_{\nu i}^{\mu-} - \frac{1}{2}g_{\nu i}^{\mu-}\mathcal{R}_i^-\right)\sigma_- = -\frac{8\pi k}{c^4}T_i^{de\mu}{}_{\nu}\sigma_- - \frac{8\pi k}{c^4}T^{m\mu}{}_{\nu}(\sigma_+ + \sigma_-)$$



$$\mathcal{R}_{\nu i}^{\mu+} - \frac{1}{2}g_{\nu i}^{\mu+}\mathcal{R}_i^+ = -\frac{8\pi k}{c^4}T^{m\mu}{}_{\nu}$$

$$\mathcal{R}_{\nu i}^{\mu-} - \frac{1}{2}g_{\nu i}^{\mu-}\mathcal{R}_i^- = -\frac{8\pi k}{c^4}T_i^{de\mu}{}_{\nu} - \frac{8\pi k}{c^4}T^{m\mu}{}_{\nu}$$



Pseudo-Complex General Relativity (pc-GR):

Perfect anisotropic fluid at rest

$$T_{\nu}^{\mu} = \begin{bmatrix} \varepsilon & 0 & 0 & 0 \\ 0 & -p_r & 0 & 0 \\ 0 & 0 & -p_{\theta} & 0 \\ 0 & 0 & 0 & -p_{\theta} \end{bmatrix}$$



Isotropic assumption $p_r = p_{\theta}$



$$T_{\nu}^{m\mu} = \begin{bmatrix} \varepsilon_m & 0 & 0 & 0 \\ 0 & -p_m & 0 & 0 \\ 0 & 0 & -p_m & 0 \\ 0 & 0 & 0 & -p_m \end{bmatrix}$$

$$T_i^{de\mu} = \begin{bmatrix} \varepsilon_{\Lambda i} & 0 & 0 & 0 \\ 0 & -p_{\Lambda i} & 0 & 0 \\ 0 & 0 & -p_{\Lambda i} & 0 \\ 0 & 0 & 0 & -p_{\Lambda i} \end{bmatrix}$$



Pseudo-Complex TOV system:

Anisotropic interior


$$\frac{dp_{rI}}{dr} = -\frac{(\varepsilon_I(r) + p_{rI}(r))}{r[r - 2m_m(r) + 2m_{\Lambda i}(r)]} \left[m_m(r) - m_{\Lambda i}(r) + \frac{4\pi k}{c^4} r^3 p_{rI}(r) \right]$$

$$\Delta p_I = p_{\theta I} - p_{rI}$$

$$p_{rI} = p_m + p_{\Lambda ri}$$

$$\varepsilon_I = \varepsilon_m + \varepsilon_{\Lambda i}$$

$$\frac{dp_m}{dr} = -\frac{(\varepsilon_m(r) + p_m(r))}{r[r - 2m_m(r) + 2m_{\Lambda i}(r)]} \left[m_m(r) - m_{\Lambda i}(r) + \frac{4\pi k}{c^4} r^3 (p_{\Lambda ri} + p_m)(r) \right]$$


$$\frac{dm_m}{dr} = 4\pi r^2 \varepsilon_m(r) \qquad p_m = p_m(\varepsilon_m)$$

$$\frac{dp_{\Lambda ri}}{dr} = -\frac{(\varepsilon_{\Lambda i}(r) + p_{\Lambda ri}(r))}{r[r - 2m_m(r) + 2m_{\Lambda i}(r)]} \left[m_m(r) - m_{\Lambda i}(r) + \frac{4\pi k}{c^4} r^3 (p_{\Lambda ri} + p_m)(r) \right]$$

$$\frac{dm_{\Lambda}}{dr} = -4\pi r^2 \varepsilon_{\Lambda}(r) \qquad p_{\Lambda} = p_{\Lambda}(\varepsilon_{\Lambda})$$



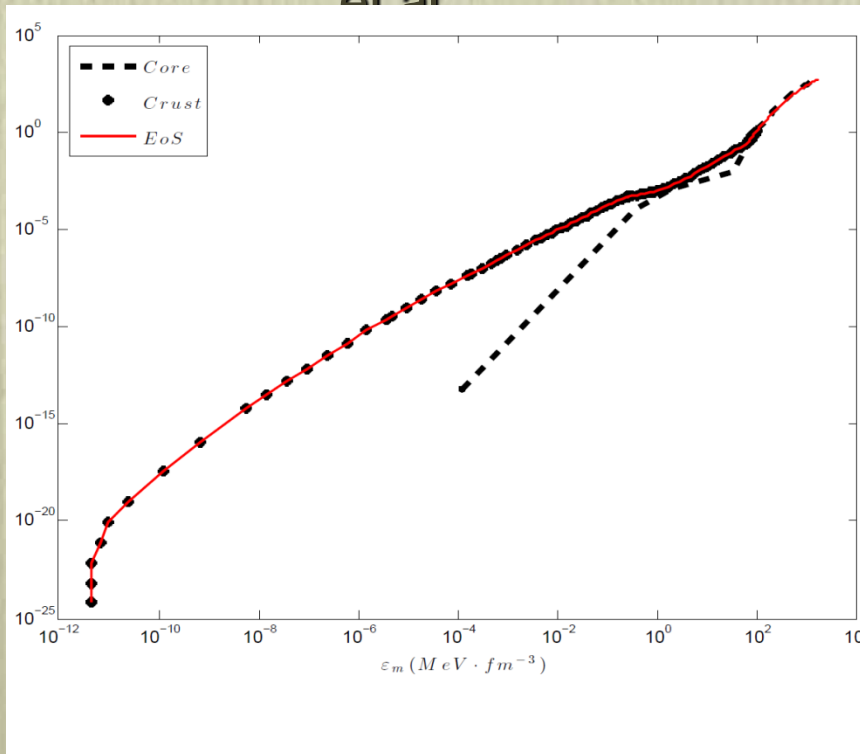
Isotropic assumption $p_r = p_{\theta}$



Equations of State:

Matter

chiral SU(3) Schramm
et al



pc-component

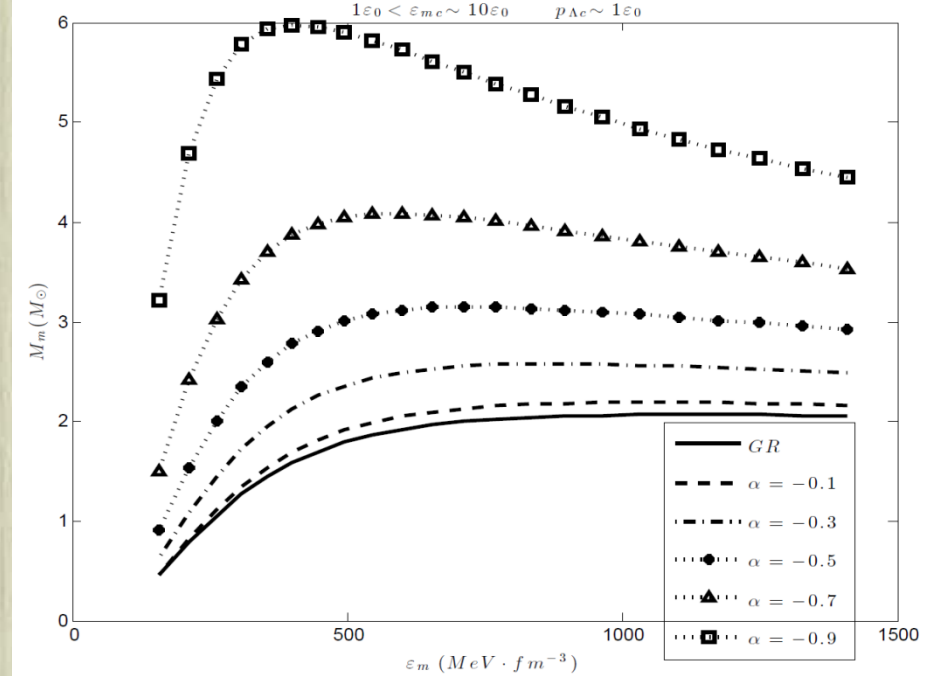
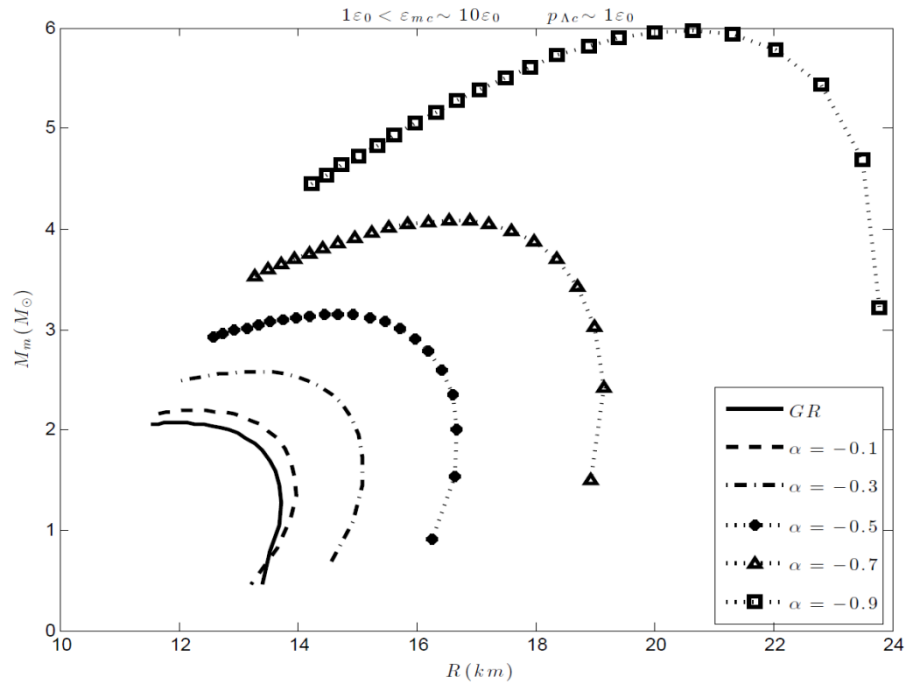
?

$$\epsilon_{\Lambda} = \alpha \epsilon_m \quad \alpha < 0$$

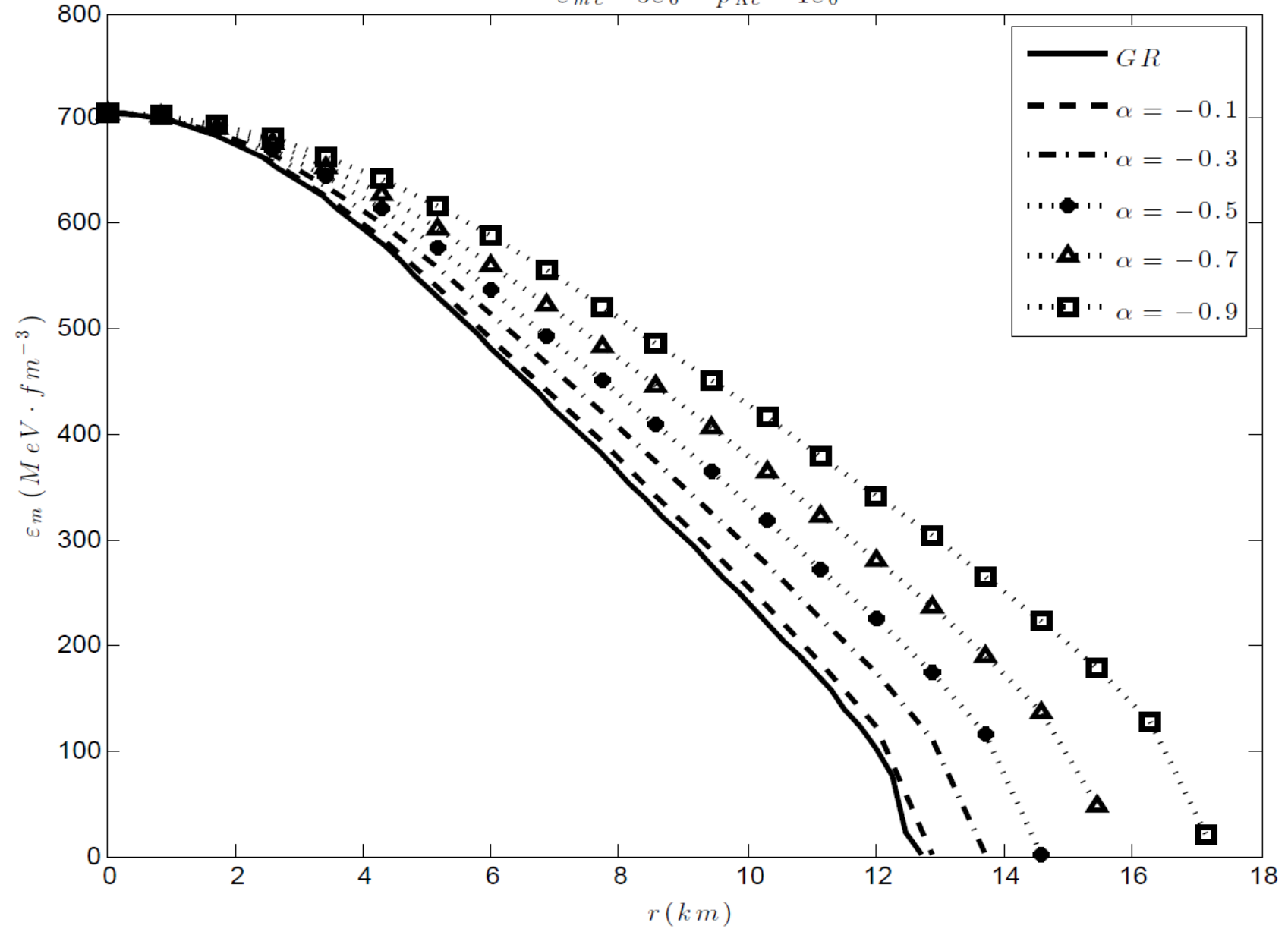
C- R- Ghezzi, „Anisotropic Dark Energy Stars“,
gr-qc:0908.0779



Results:



$\varepsilon_{me} \sim 5\varepsilon_0$ $p_{\Lambda e} \sim 1\varepsilon_0$

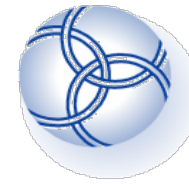




Conclusion & Future Work:

- It was shown that pc-GR can produce neutron stars with a significant larger mass as obtained in standard GR.
- Problems of the correct equation of state of the mass and the coupling of the dark energy density to the mass density remain.
- In future we try to solve this problem.

Thanks!



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