

# Signal of FCNC and CPV in 2HDM

**José Halim Montes de Oca**

Facultad de Estudios Superiores Cuautitlán  
UNAM

josehalim@gmail.com

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# Outline

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SM of the  
particle physics  
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$t \rightarrow d^{+}l^{-}$

Final remarks

- 1 CP violation in Standard Model
  - Standard Model of the particle physics
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- 2 FCNC and CPV in Two Higgs Doublet Model
  - Two Higgs Doublet Model
  - Flavor Changing Neutral Currents and CP-violation
  - Spontaneous CPV

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# Standard Model of the particle physics

- Gauge group  $SU(3)_C \otimes SU(2)_L \otimes U(1)_Y$
- Gauge invariance introduces interactions
- Higgs mechanism

Field Content:

$$q_L = \begin{pmatrix} u_{t,c,u} \\ d_{b,s,d} \end{pmatrix}_L, \quad l_L = \begin{pmatrix} \nu_{t,c,u} \\ e_{\tau,\mu,e} \end{pmatrix}_L$$

$$u_{R(t,c,u)}, \quad d_{R(b,s,d)} \quad \text{and} \quad e_{R(\tau,\mu,e)}$$

where  $\psi_{L,R} = \frac{1 \mp \gamma_5}{2} \psi$ .

# Higgs mechanism

The basic idea was introduce a complex scalar doublet field with  $Y = 1$

$$\phi = \begin{pmatrix} \varphi^+ \\ \varphi_0 \end{pmatrix}$$

and a potential,

$$V = -\mu^2 (\phi^\dagger \phi) + \frac{\lambda}{2} (\phi^\dagger \phi)^2$$

Then, SBB happens when vev is taken  $\langle \phi \rangle^T = \left( 0, \frac{1}{\sqrt{2}} v \right)$  and

$$M_{W^\pm} = \frac{ev}{2 \sin \theta_W},$$

$$M_Z = \frac{M_{W^\pm}}{\cos \theta_W},$$

$$M_H = \lambda v.$$

# Experimentally

- $\sin^2 \theta_W = 0.23116(12)$ ,
- $M_W = 80.385(15)$  GeV,
- Observation of new boson like SM Higgs boson is given by ATLAS<sup>1</sup> and CMS<sup>2</sup> with a mass of the order of 126.0(0.4)GeV and 125.3(0.4)GeV, respectively.

<sup>1</sup>ATLAS collaboration, G. Aad et al., Phys. Lett. **B 716**, 1, (2012).

<sup>2</sup>CMS collaboration, S. Chatrchyan et al., Phys. Lett. **B 716**, 30, (2012).



# Fermions-Higgs boson interaction

The interaction is introduced as

$$\mathcal{L}_{Yukawa} = \bar{q}_{Li}^0 Y_{ij}^u \tilde{\phi} u_{Rj}^0 + \bar{d}_{Li}^0 Y_{ij}^d \phi d_{Rj}^0 + \bar{l}_{Li}^0 Y_{ij}^l \phi e_{Rj}^0 + h.c.,$$

with  $\phi = \begin{pmatrix} \varphi^+ \\ \varphi^0 \end{pmatrix}$  and  $\tilde{\phi} = i\sigma_2 \phi$ .

After getting a correct SSB,  $\langle \phi \rangle^T = \left( 0, \frac{1}{\sqrt{2}} v \right)$ ,

$$\mathcal{L}_{Yukawa} = \bar{u}_i M_{ij}^u u_j + \bar{d}_i M_{ij}^d d_j + \bar{e}_i M_{ij}^e e_j,$$

where  $\psi_i^0 = U_{ij}^\psi \psi_j$  for  $\psi = u, d, e$ .

# CKM matrix

The quark and lepton kinetic terms are  $\bar{\psi}\gamma^\mu D_\mu\psi$ , with  $D_\mu = \partial_\mu - ig_1\frac{Y}{2}B_\mu - ig_2\frac{\tau_i}{2}W_\mu^i$ .  
Then, quarks are written in mass eigenstates

$$-\frac{e}{\sqrt{2}\sin\theta_W}\left[\bar{u}_{Ll}\gamma^\mu W_\mu^+(V_{CKM})_{ij}u_{Lj} + h.c.\right],$$

where  $V_{CKM} = U_L^u (U_L^d)^\dagger$ . The standard parametrization is

$$V_{CKM} = \begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta} & c_{23}c_{13} \end{pmatrix},$$

where  $c_{ij} = \cos\theta_{ij}$ ,  $s_{ij} = \sin\theta_{ij}$  and  $\delta$  is the CP-violating phase.

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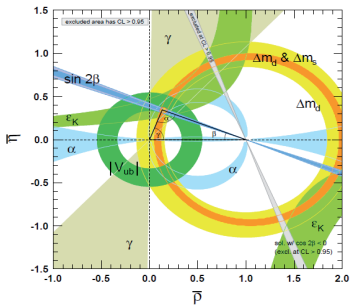
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The angles of unitary triangle are

$$\beta = \arg\left(-\frac{V_{cd}V_{cb}^*}{V_{td}V_{tb}^*}\right), \quad \sin 2\beta = 0.679(20)$$

$$\alpha = \arg\left(-\frac{V_{ud}V_{ub}^*}{V_{ud}V_{ub}^*}\right), \quad \alpha = (89.0^{+4.4}_{-4.2})^\circ$$

$$\gamma = \arg\left(-\frac{V_{ud}V_{ub}^*}{V_{cd}V_{cb}^*}\right), \quad \gamma = (68^{+10}_{-11})^\circ$$

# General 2HDM potential

- Two complex  $SU(2)_L$  doublets:  $\Phi_{1,2}$  ( $Y = 1$ ).

$$\begin{aligned} V = & m_{11}^2 \Phi_1^\dagger \Phi_1 + m_{22}^2 \Phi_2^\dagger \Phi_2 - \left[ m_{12}^2 \Phi_1^\dagger \Phi_2 + \text{h.c.} \right] \\ & + \frac{1}{2} \lambda_1 \left( \Phi_1^\dagger \Phi_1 \right)^2 + \frac{1}{2} \lambda_2 \left( \Phi_2^\dagger \Phi_2 \right)^2 \\ & + \lambda_3 \left( \Phi_1^\dagger \Phi_1 \right) \left( \Phi_2^\dagger \Phi_2 \right) + \lambda_4 \left( \Phi_1^\dagger \Phi_2 \right) \left( \Phi_2^\dagger \Phi_1 \right) \\ & + \left[ \lambda_5 \left( \Phi_1^\dagger \Phi_2 \right)^2 + \text{h.c.} \right] \\ & + \left[ \lambda_7 \left( \Phi_1^\dagger \Phi_1 \right) \left( \Phi_1^\dagger \Phi_2 \right) + \lambda_8 \left( \Phi_2^\dagger \Phi_2 \right) \left( \Phi_1^\dagger \Phi_2 \right) + \text{h.c.} \right] \end{aligned}$$

Explicit CP-Violation from  $m_{12}^2$ ,  $\lambda_{5,6,7}$ .

# Neutral Higgs mass-eigenstates

The most general  $U(1)_{EM}$ -conserving vacuum expectation value is:  $\langle \Phi_a \rangle = \frac{v}{\sqrt{2}} \begin{pmatrix} 0 \\ \hat{v}_a \end{pmatrix}$ , with  $\hat{v}_a = e^{i\eta} \begin{pmatrix} \cos \beta \\ \sin \beta e^{i\xi} \end{pmatrix}$  for  $a = 1, 2$ .

The neutral Higgs mass-eigenstate are

$$\begin{pmatrix} h_1 \\ h_2 \\ h_3 \end{pmatrix} = R \begin{pmatrix} \phi_1^0 \\ \phi_2^0 \\ a^0 \end{pmatrix}.$$

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# Flavor Changing Neutral Currents and CP-violation

The most general structure of the Yukawa interactions,

$$\mathcal{L}_{Yukawa} = \sum_{i,j=1}^3 \sum_{a=1}^2 \left[ \bar{q}_{Li}^0 [Y_a^u]_{ij}^0 \tilde{\Phi}_a u_{Rj}^0 + \bar{q}_{Li}^0 [Y_a^d]_{ij}^0 \Phi_a d_{Rj}^0 + \bar{l}_{Li}^0 [Y_a^l]_{ij}^0 \Phi_a e_{Rj}^0 + h.c. \right],$$

Then,

$$M^{u,l} = \frac{v_1}{\sqrt{2}} Y_1^{u,l} + e^{-i\xi} \frac{v_2}{\sqrt{2}} Y_2^{u,l}$$

and

$$M^d = \frac{v_1}{\sqrt{2}} Y_1^d + e^{i\xi} \frac{v_2}{\sqrt{2}} Y_2^d$$

where  $Y_a^f = V_L^f Y_a^{f,0} (V_R^f)^\dagger$  and  $V_{L,R}^f$  are dimensionless  $3 \times 3$  matrices such as  $f_{(L,R)i}^0 = \sum_j V_{(L,R)ij}^f f_{(L,R)j}^0$ .

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We can write the Yukawa interactions between neutral Higgs bosons and up-quarks as

$$\begin{aligned} \mathcal{L}_{Yukawa} = & \sum_{i,j,k=1}^3 \bar{u}_i \left[ \frac{1}{v} \left( \frac{1}{\cos \beta} R_{k1} + i R_{k3} \tan \beta \gamma_5 \right) M_{ij}^u \right. \\ & - R_{k1} \tan \beta J_{+ij} + R_{k2} K_{+ij} - i R_{k3} \tan \beta \sin \beta J_{-ij} \\ & \left. - i R_{k3} \cos \beta K_{-ij} \right] h_k u_j, \end{aligned}$$

where  $J_{\pm} = e^{-i\xi} Y_2^u P_R \pm e^{i\xi} (Y_2^u)^{\dagger} P_L$  and  
 $K_{\pm} = Y_2^u P_R \pm (Y_2^u)^{\dagger} P_L$

# Usual scalar mixing and Spontaneous CPV

Yukawa interactions for  $\alpha_2 = \alpha_3 = 0$ ,  $\alpha_1 = \alpha + \pi/2$ ,  $h_1 = h^0$  and  $h_2 = -H^0$ , become as

$$\begin{aligned} \mathcal{L}_{Yukawa}^{up} = & \sum_{i,j=1}^3 \bar{u}_i \left[ -\frac{1}{v \cos \beta} (h_0 \sin \alpha + H_0 \cos \alpha) M_{ij}^u \right. \\ & + \tan \beta (h_0 \sin \alpha - H_0 \cos \alpha) J_{+ij} \\ & + (h_0 \cos \alpha + H_0 \sin \alpha) K_{+ij} \\ & \left. + i \left( \frac{\tan \beta}{v} \gamma_5 M_{ij}^u - \tan \beta \sin \beta J_{-ij} - \cos \beta K_{-ij} \right) A_0 \right]. \end{aligned}$$

where  $J_{\pm} = e^{-i\xi} Y_2^u P_R \pm e^{i\xi} (Y_2^u)^\dagger P_L$  and  $K_{\pm} = Y_2^u P_R \pm (Y_2^u)^\dagger P_L$ .



# Rare top decay

The decay width for rare top decay is given by

$$\Gamma_{t \rightarrow d^+ l^-} = \frac{m_t^5 m_l^2 |Y_{23}^u|^2}{512 \pi^3 m_{h_1}^4 v^2} \left| \frac{\sqrt{2} v \cos(\alpha - \beta)}{2 m_l} Y_{33}^l - \frac{\sin \alpha}{\cos \beta} \right|^2$$
$$\left[ \frac{R_{13}^2}{\cos^2 \beta} (\cos^4 \beta + \sin^4 \beta + 2 \cos^2 \beta \sin^2 \beta \cos \xi) \right. \\ \left. - 2 R_{11} R_{12} \tan \beta \cos \xi - 2 R_{11} R_{13} \sin \beta \cos \xi \right. \\ \left. + 2 R_{12} R_{13} \frac{\sin \xi}{\cos \beta} (\cos^2 \beta \cos \xi + \sin^2 \beta) \right. \\ \left. + R_{12}^2 + R_{11}^2 \tan^2 \beta \right] F(\mu_1, \mu_2),$$

where

$$F(\mu_1, \mu_2) = 2(1 + \mu_1 - 4\mu_2) \sqrt{\mu_1} l_1 \\ + (1 + \mu_1 - 4\mu_2 - 2\sqrt{\mu_1}) l_2 - l_3$$

with  $\mu_1 = \frac{m_\xi^2}{s}$ ,  $\mu_2 = \frac{m_l^2}{s}$ . NN

# SCP-violation phase and usual cases

In particular,

$$\Gamma_{t \rightarrow cl^+l^-} = \frac{m_t^5 |Y_{23}^u|^2}{256\pi^3 m_{h^0}^4} \left| \frac{\cos(\alpha - \beta)}{\sqrt{2} \cos \beta} Y_{33}^l - \frac{m_l \sin \alpha}{v \cos \beta} \right|^2 (\tan^2 \beta \sin^2 \alpha + 2 \tan \beta \cos \alpha \sin \alpha \cos \xi + \cos^2 \alpha) F$$

Then, branching ratio is  $Br(t \rightarrow cl^+l^-) \approx \frac{\Gamma_{t \rightarrow cl^+l^-}}{\Gamma_{top}}$ , where  $\Gamma_{top} = 1.2 \text{ GeV}$ ,

$$Br(t \rightarrow cl^+l^-) \approx \frac{m_t^5 (I_2 - I_3)}{512\pi^3 m_{h^0}^4} |Y_{23}^u|^2 |Y_{33}^l|^2 \left( \frac{\cos(\alpha - \beta)}{\cos \beta} \right)^4, \quad (2)$$

The mixing angle  $\beta$  has been bounded as  $1 \leq \tan \beta \leq 50$  and  $-\pi/2 \leq \alpha \leq \pi/2$

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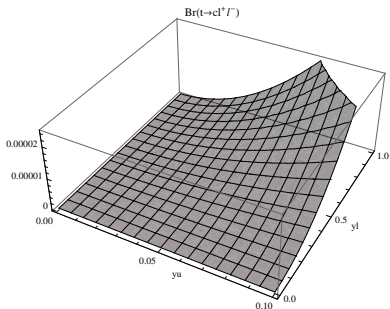


Figure:  $Br(t \rightarrow c \tau^+ \tau^-)$  for  
values  $\alpha = 0.1$ ,  $\tan \beta = 1$  and  
 $\xi = 0$ .

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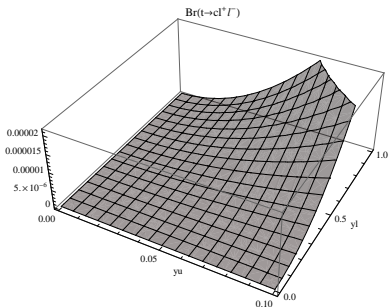


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- 2HDM is a source for FCNC and CPV
- Top factories
- SM has been probed successfully

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